

Initial overview of the L-unification

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July 19, 2024

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1 Disadvantages of the standard model

Non-sufficient self-consistency of the underlying version of quantum field theory.
Incapability of describing the gravitational sector at quantum level.
The standard model is too complicated to be the dreamlike distinguished fundamental theory.

2 Two basic ways to create a new theory

One can think that the particles of the standard model are not really elementary, but are made from a smaller number of simpler components – this leads to development of composite, e.g. technicolor models.

One can think that the particles of the standard model are only fragments of some bigger, more symmetric and universal entities – this is the idea behind unification models.

3 Unification models

The basic idea of all unification models is to enhance the symmetry, embedding the non-simple gauge group $G_{st} = SU(3) \times SU(2) \times U(1)$ of the standard model into a larger simple group G , which inevitably has larger dimension and thus introduces extra particles.

4 Main idea

The central idea of Lisi's approach is to consider a minimal modification – ideally just a reformulation – of the standard model, with as few additional fields as only possible.

No extra essences are introduced, like superpartners, extra dimensions, branes, strings etc.

Since the number of fields in the standard model is of the order of 200, the number of different generators of "unification algebra" should exceed this number. If Lie algebra is classical, like $SU(N)$, $SO(N)$ or $Sp(N)$, this requires it to have a large rank $= \sqrt{200} > 14$. If instead the group is exceptional, then the smallest group with dimension > 200 is actually the largest: E_8 with $\dim(E_8) = 248$ and its rank is as low as $\text{rank}(E_8) = 8$ (dimension of the next smaller exceptional group $\dim(E_7) = 133 < 200$).

5 L and G - unification

Furthermore, while in G-unification models, in order to minimize the number of 3 additional gauge bosons, one attempts to find a group of the lowest dimension with a given rank, in L-unification for the same purpose one tries to minimize the rank for a given dimension, which is now defined by the number of standard-model particles, predominantly fermions.

6 The number of known elementary particles

Standard model fermions numbers $n_F = 96$ (degrees of freedom) and $N_F = 192$ (fields).

Standard model bosons numbers $n_B = 30$ and $N_B = 92$.

7 Without extra particles and fields?

Extra particles, are, however, necessary in any case: because of the rank mismatch: there is simply no Lie group which could both match the strict lower limit $\dim(G) \geq n_B + n_F > n_F = 90$ and saturate the rank condition $\text{rank}(G) = \text{rank}(U(1)) + \text{rank}(SU(2)) + \text{rank}(SU(3)) + \text{rank}(SO(3, 1)) = 1 + 1 + 2 + 2 = 6$.

8 The main claim

The action of the standard model coupled to Einstein gravity,

$$\begin{aligned} & \int \sqrt{\det g} d^4x \{ M_{\text{Pl}}^2 R(g) - \Lambda + \frac{1}{4g_3^2} \text{Tr}_{SU(3)} G_{\mu\nu} G^{\mu\nu} + \frac{1}{4g_2^2} \text{Tr}_{SU(2)} W_{\mu\nu} W^{\mu\nu} + \frac{1}{4g_1^2} V_{\mu\nu} V^{\mu\nu} + \\ & + \sum_{p=1}^{n_g=3} \left(\bar{l}_L^{(p)} \hat{D} l_L^{(p)} + \bar{l}_R^{(p)} \hat{D} l_R^{(p)} + \bar{q}_L^{(p)} \hat{D} q_L^{(p)} + \bar{q}_R^{(p)} \hat{D} q_R^{(p)} \right) + \frac{1}{2} D_\mu \phi^\dagger D^\mu \phi + V(\phi) + \\ & + \sum_{P,Q=1}^{2n_g=6} \left(M_l^{(PQ)} \bar{l}_R^{(P)} \phi_L^{(Q)} + M_q^{(PQ)} \bar{q}_R^{(P)} \phi_L^{(Q)} + \text{c.c.} \right) \} \quad (1) \end{aligned}$$

can be rewritten in a compact form as

$$\int \text{Tr}_{E_8} B \wedge F + \int Q(B) \quad (2)$$

where

$$F = dA + A \wedge A \quad (3)$$

is a curvature of the "generalized-connection" field $A(x)$, which is a linear combination of the standard-model fields, distributed over elements of the E_8 -algebra

matrix (adjoint representation of E_8), and $Q(B)$ is a quadratic function of the auxiliary field B , which is also an adjoint E_8 matrix.

In a little more detail:

$$A = \sum_{\alpha \in G}^{56} A^\alpha(x) T_\alpha \oplus \sum_{a \in E/G}^{192} \psi^a(x) T_a, \quad (4)$$

$$B = \sum_{\alpha \in G}^{56} B^\alpha(x) T_\alpha \oplus \sum_{a \in E/G}^{192} \chi^a(x) T_a \quad (5)$$

and

$$Q(B) = \sum_{\alpha, \alpha' \in G}^{56} \left(Q_{\text{grav}}^{\alpha, \alpha'} B^\alpha \wedge B^{\alpha'} + Q_{\text{YM}}^{\alpha, \alpha'} B^\alpha \wedge *B^{\alpha'} \right) \quad (6)$$

Here the sums are over 248 generators of E_8 , which are divided into two different groups of 56 and 192, labeled by Greek and small Latin letters respectively, which correspond to decomposition of adjoint (the minimal possible) representation of E_8 into a subalgebra $G = SO(7, 1) \times SO(8)$ and its representation $R = E/G$.

Bosonic 1-forms $A^\alpha = A_\mu^\alpha dx^\mu$ include the gauge fields of the standard model, the gravitational spin-connection and a mixture of Higgs scalars and gravitational tetrads. One of the two tetrad indices is included into the set of E_8 -indices $\{\alpha\}$, the remaining vector index, as well as the one of the gauge vectors is the 1-form index of A .

Fermionic (Grassmann-valued) 0-forms ψ^a and χ^a are made from the fermions of the standard model, with spinor indices included into the set of $\{a\}$. Auxiliary fields $B^\alpha = B_{\mu\nu}^\alpha dx^\mu \wedge dx^\nu$ and $\chi^a = \chi_{\mu\nu\lambda}^a dx^\mu \wedge dx^\nu \wedge dx^\lambda$ are bosonic 2-forms and fermionic 3-forms respectively. No 4-form component is included into B , therefore the ψ^2 term in $F = F^\alpha T_\alpha + (D\psi)^a T_a + \psi^a \psi^b [T_a, T_b]$ does not contribute to the action.

The main disadvantage is, however, more serious: E_8 unification-model not fully justified and, most probably, is simply wrong. Desired distribution of all the fields of the standard model among the generators of E_8 is not actually found.

9 Main drawbacks

Introduction of new fields beyond the standard model is still unavoidable.
Hierarchy and of quantum gravity are not resolved.
Cosmological constant is non-vanishing in BF -version of the Palatini formalism.
G-projector and the Hodge star, appear in the action – not fully topological.
Higher generations are not adequately described.

10 Features of L-unification

Unification algebra $G = SO(7,1) \times SO(8)$ contains gravitational $SO(3,1) = SL(2, C)$ and Yang-Mills $G_{YM} = SU(3)(SU(2) \times U(1)) \times 2$.
This is a vector-like theory, with right-hand copies of W , Z bosons and second photon – which should somehow decouple at low energies: along with other gauge bosons from G/G_{st} should acquire large masses.
Unification group $G = SO(7,1) \times SO(8)$ is not simple, with strong and electroweak groups belonging to different factors – $SO(8)$ and $SO(7,1)$ respectively, – there is no danger of proton decay.
The same non-simplicity implies that unification of coupling constants g_{str} and g_{weak} does not need to take place in this model.
Flavor-changing neutral currents.

11 From Palatini action to Einstein equations

In the case of arbitrary space-time dimension d Palatini action involves the curvature of $SO(d)$ connection

$$S_p = \epsilon_{a_1 \dots a_d} \int \exp a_1 \wedge \dots \wedge \exp a_{d-2} \wedge R^{a_{d-1} a_d} \quad (7)$$

$$\dots \quad (8)$$

$$R_\mu^a + \frac{1}{2} e_\mu^a R \sim \det(e) T_\mu^a - e_\mu^a T \quad (9)$$

It is easy to deduce that $F = -T$ and therefore previous equation is actually equivalent to the proper Einstein equation

12 Decompositions of E_8 algebra

$$E_8 = SO(16) + R_{128} \quad (10)$$

where R_{128} is spinor representation of $SO(16)$.

$$E_8 = D_8 + R_{128} = D_4 + D_4 + (8 \times 8 + R_{128}) \quad (11)$$

Gauge fields are distributed among the two D_4 factors, with gravity and two weak groups (left and right) belonging to one of them and with strong and two abelian groups belonging to another.

Fermionic fields are distributed among the spinor representation R_{128} of maximal subgroup $SO(16) \subset E_8$ and the 8×8 generators of the factor $D_8/D_4 \times D_4$. The three 64-plet constituents of the fermionic sector are related by discrete triality symmetry.