

First spectral moments of collision-induced rototranslational absorption band of CO₂-Ar pair: Classical theory with the use of ab initio calculated anisotropic interactions

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Introduction

Much attention is presently devoted to the study of dipole-forbidden molecular absorption caused by weak intermolecular interaction. This interest is significantly motivated by the need to refine the accuracy of climate models which are being developed for the planetary or exoplanetary atmospheres. Actual knowledge of binary absorption coefficients is still fragmentary (see e.g. [1]) and is thus fairly unsatisfactory in regard of great number of pairs potentially interesting for climate modeling. Experimental observation of the weak pressure-induced absorption is especially difficult in the far-infrared where the most important so-called rototranslational collision-induced absorption (RT CIA) bands are conventionally situated. Direct quantum calculations are still barely feasible for interacting polyatomics, although significant progress is worth mentioning that has been achieved recently in both quantum calculations of the spectral profiles [2,3] and quantum theory of spectral moments [4,5]. Molecular dynamics simulation and classical trajectory approaches are presently the most promising (see e.g. [6,7]) in terms of purely classical methods. The use of the formerly popular semi-empirical methods is getting more and more senseless nowadays as far as unprecedented computational facility and advancement in quantum chemical methods become available.

Present paper aims at theoretical examination of the first spectral moments in rototranslational CIA band of the prototype CO₂-Ar system. Our approach relies on the use of anisotropic potential energy and induced dipole surfaces obtained by virtue of sophisticated *ab initio* calculations. The knowledge of these surfaces permits direct classical calculation of the first spectral moments without address to any adjustable parameters. Our ultimate goal consists of development of such a rigorous classical formalism that could permit consideration of CIA in any arbitrary pair formed by typical atmospheric molecules.

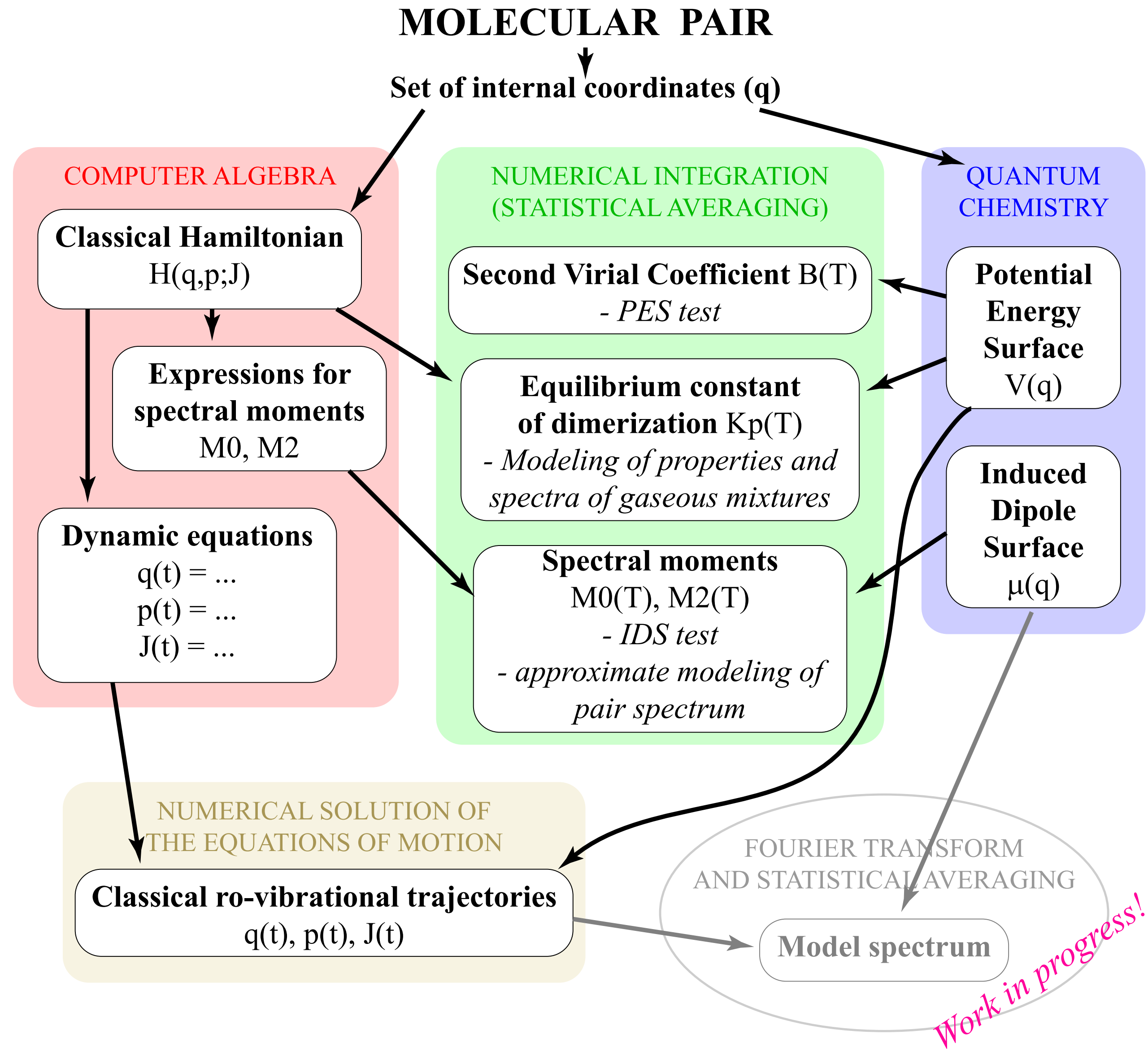


Figure 1: A scheme illustrating main areas of work and their interrelations

Calculation of spectral moments

Spectral moments are integral values, which are widely used to characterise CIA spectra. The usage of these values is attributed to the possibility to represent them either in terms of integrals over experimentally measurable spectral profiles or in terms of Boltzmann weighted functions of induced dipole. Thus the knowledge of complete potential energy (PES) and induced dipole (IDS) surfaces can help both verification of *ab initio* obtained data and characterization of interaction-induced absorption.

Two systems of coordinates are conventionally employed in theoretical consideration of collisional dynamics among polyatomic molecules. These are so-called laboratory- and body-fixed frames. Salient feature of our approach consists of development and subsequent use of rigorous classical Hamiltonian in the body-fixed frame. All kinetic energy terms that are responsible for Coriolis interaction are kept in derived Hamiltonian. In its general form the latter can be represented as

$$H = \frac{1}{2} \mathbf{p}^\top \mathbb{G}_{11} \mathbf{p} + \mathbf{p}^\top \mathbb{G}_{12} \mathbf{J} + \frac{1}{2} \mathbf{J}^\top \mathbb{G}_{22} \mathbf{J} + U(\mathbf{q}), \quad (1)$$

where \mathbf{q} , \mathbf{p} denote the generalized coordinates and conjugated momenta, \mathbf{J} denotes the vector of total angular momentum (matrices \mathbb{G}). For the analysis of collision-induced spectra and the comparison with theory, certain integrals of the spectra, the spectral moments, are used. The n -th spectral moment of the spectral function, $G(\nu)$, is defined by [8]:

$$M_n = \int_{-\infty}^{\infty} \nu^n V G(\nu) d\nu \quad (2)$$

In the classical limit the following expressions are used to theoretically calculate spectral moments [8]:

$$M_{2n} = V (2\pi c)^{-n} \frac{1}{4\pi \varepsilon_0} \left\langle \left| \frac{d^n}{dt^n} \boldsymbol{\mu}(t) \right|^2 \right\rangle \Big|_{t=0} \quad (3)$$

$$M_{2n+1} = 0$$

where $\boldsymbol{\mu}$ denotes the vector of dipole moment and angular brackets denote averaging over the complete phase space. Specifically, classical expressions for lowest-order zeroth and second spectral moments can be written as

$$M_0 = \frac{\int \boldsymbol{\mu}^2 \exp(-H(\mathbf{q}, \mathbf{p}, \mathbf{J})/kT) d\mathbf{q} d\mathbf{p}}{\int \exp(-H(\mathbf{q}, \mathbf{p}, \mathbf{J})/kT) d\mathbf{q} d\mathbf{p}}, \quad M_2 = \frac{\int \dot{\boldsymbol{\mu}}^2 \exp(-H(\mathbf{q}, \mathbf{p}, \mathbf{J})/kT) d\mathbf{q} d\mathbf{p}}{\int \exp(-H(\mathbf{q}, \mathbf{p}, \mathbf{J})/kT) d\mathbf{q} d\mathbf{p}}. \quad (4)$$

The expression for the second moment includes time derivative of the dipole which can be calculated from the Poisson bracket using classical Hamiltonian

$$\frac{d\boldsymbol{\mu}}{dt} = [\boldsymbol{\mu}, H] = \sum_i \left\{ \frac{\partial \boldsymbol{\mu}}{\partial q_j} \frac{\partial H}{\partial p_j} - \frac{\partial p_j}{\partial p_j} \frac{\partial H}{\partial q_j} \right\}. \quad (5)$$

Obviously, the IDS is determined by spatial position of interacting moieties and does not depend on their velocities (conjugated momenta). That is why we are interested only in derivatives of the Hamiltonian over momenta which can be calculated numerically at any point of the phase space using matrix formalism.

Assuming the Hamiltonian is written in the body-fixed frame (BF), the squared time derivative of the dipole in laboratory frame (LF) can be written as:

$$(\dot{\boldsymbol{\mu}}^{LF})^2 = (\dot{\boldsymbol{\mu}}^{BF})^2 + 2\boldsymbol{\mu}^{BF} \left(\frac{\partial H}{\partial \mathbf{J}} \times \boldsymbol{\mu}^{BF} \right) + (\boldsymbol{\mu}^{BF})^\top \mathbb{I}^J \boldsymbol{\mu}^{BF}, \quad (6)$$

$$\mathbb{I}_{ij}^J = \sum_k \left(\frac{\partial H}{\partial J_k} \right)^2 \delta_{ij} - \left(\frac{\partial H}{\partial J_i} \right) \left(\frac{\partial H}{\partial J_j} \right) \quad (7)$$

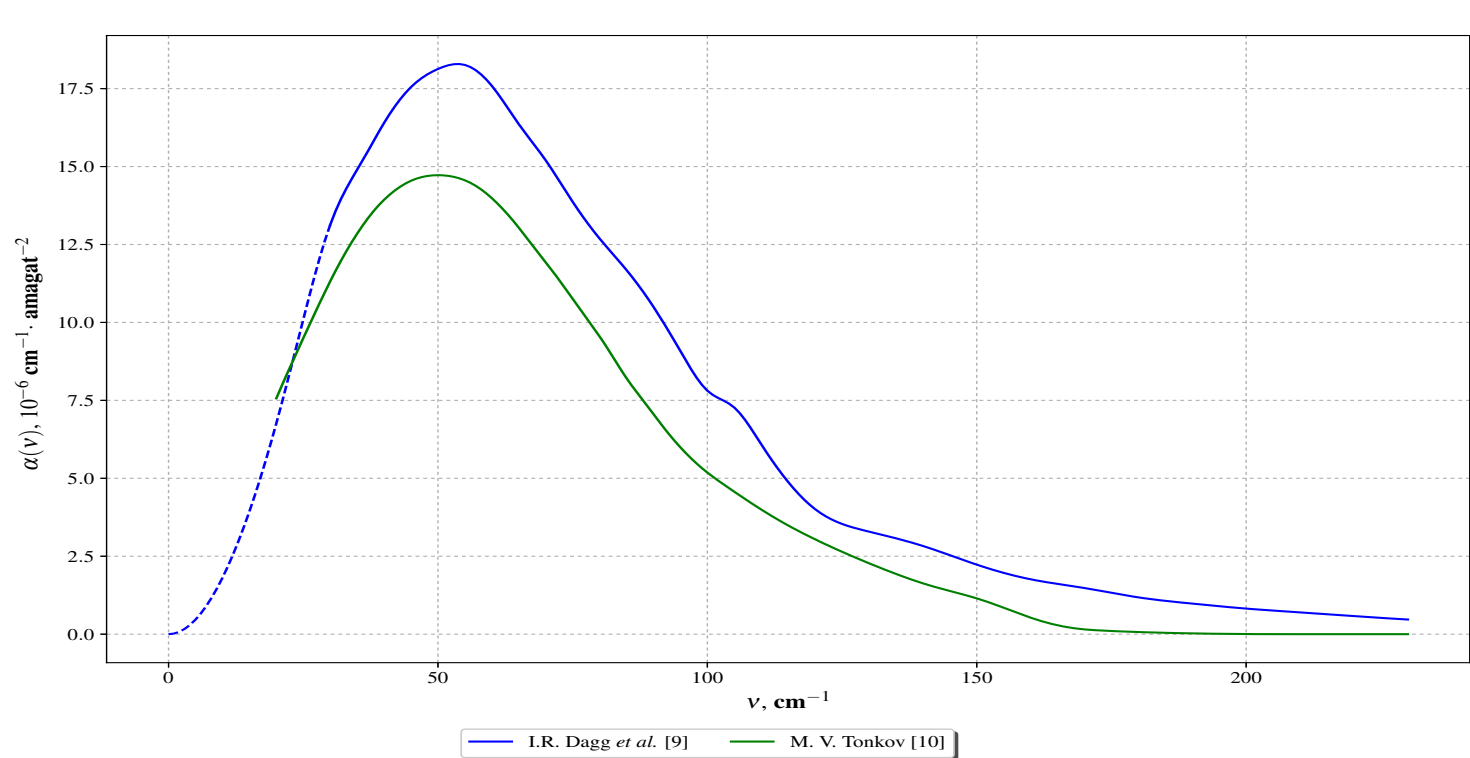


Figure 2. Experimental absorption spectra of CO₂-Ar

Spectral moments of CO₂-Ar rototranslational band

In the present paper we have chosen CO₂-Ar as an example of relatively simple atom-linear molecule system possessing an appreciable anisotropy of interaction (Fig. 1).

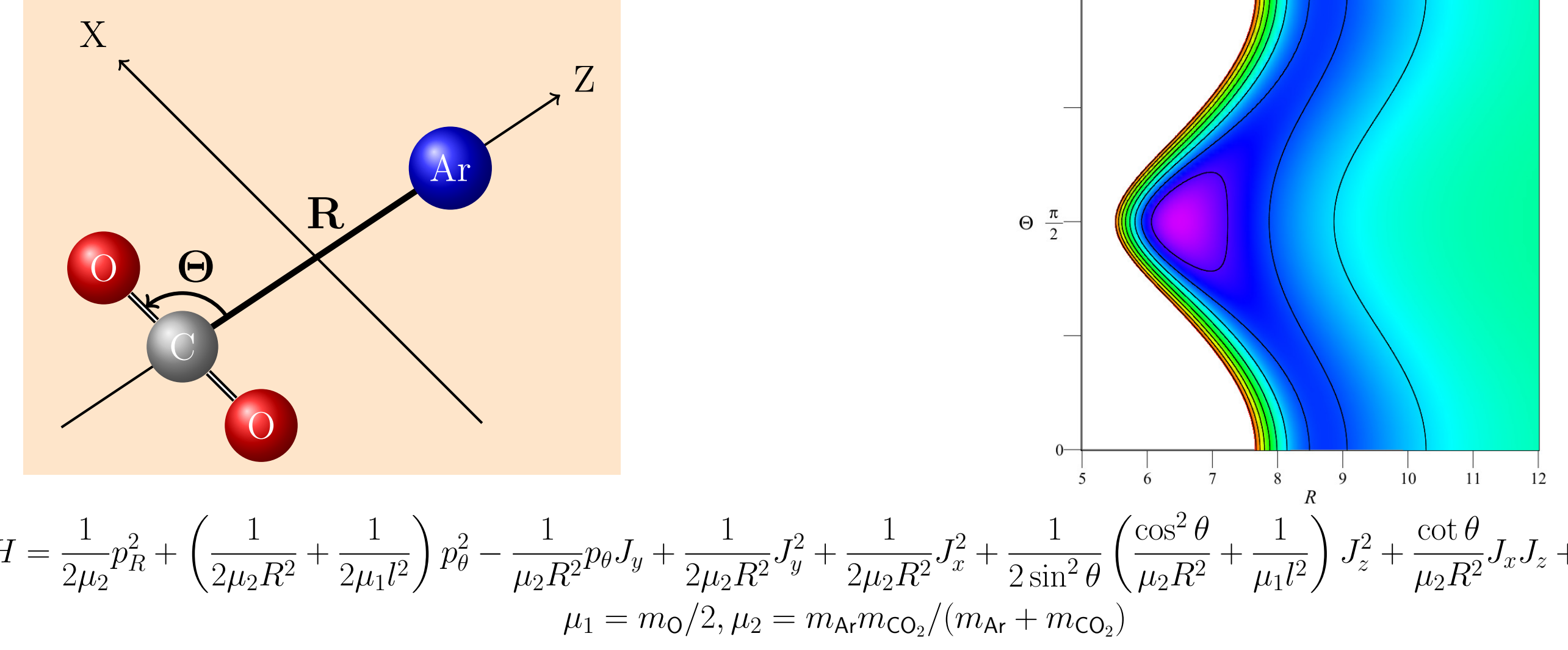


Figure 3: Coordinate system, potential energy surface (PES) and classical Hamiltonian for CO₂-Ar

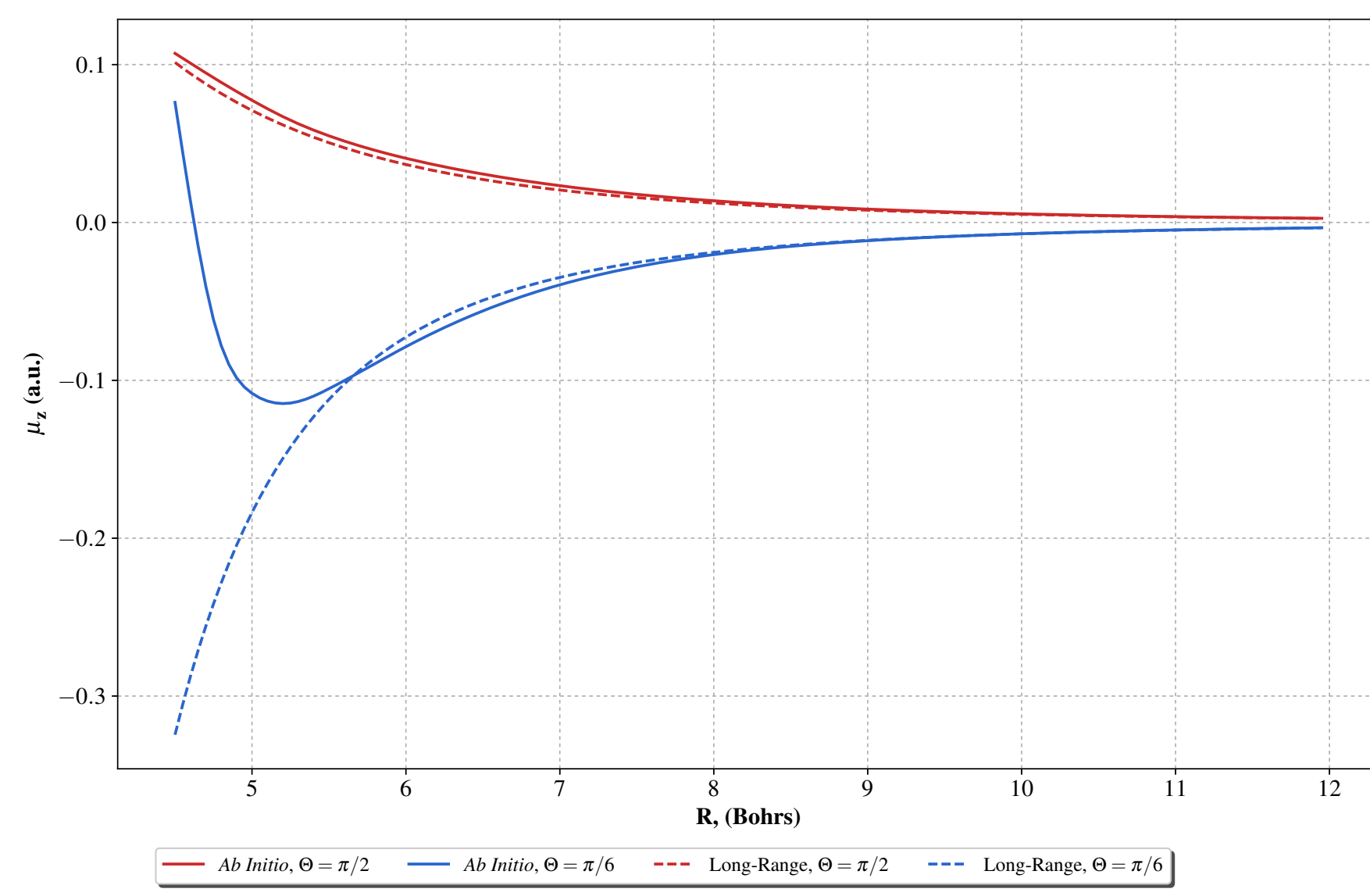


Figure 4: Z-components of the dipole moment

Our calculated T -dependences for the zeroth and second spectral moments are shown in Figs. 3, 4, respectively. The symbols in these figures relate to the data obtained from available [8, 9] experimental binary absorption coefficients $\alpha(\nu)$ using the integrals:

$$M_0 = \frac{1}{\rho_1 \rho_2} \int_0^\infty \alpha(\nu) \coth \left(-\frac{h c \nu}{2 k T} \right) \frac{d\nu}{\nu}, \quad M_2 = \frac{1}{\rho_1 \rho_2} \int_0^\infty \alpha(\nu) \nu d\nu. \quad (8)$$

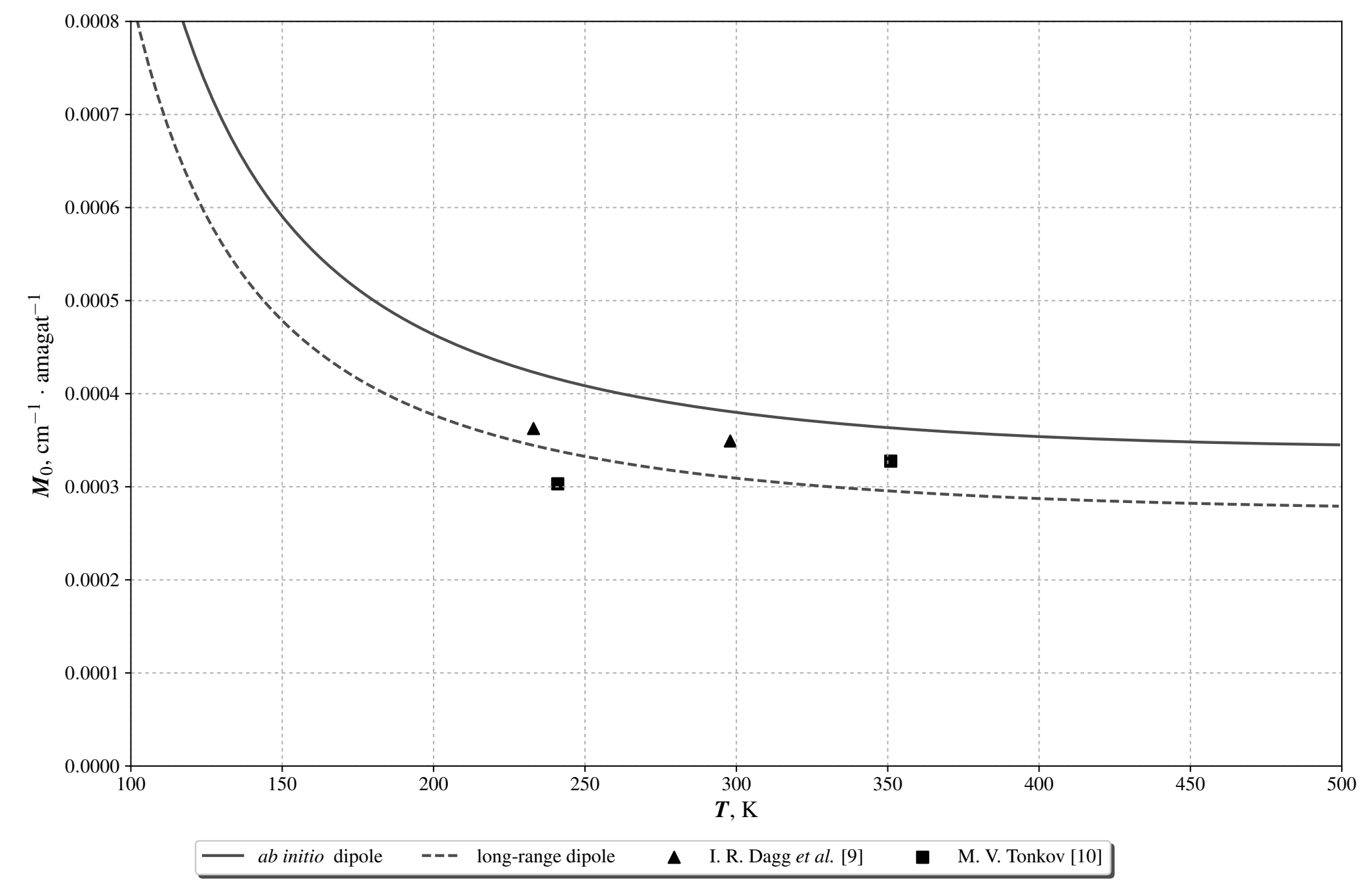


Figure 5: Temperature dependence of zeroth spectral moment

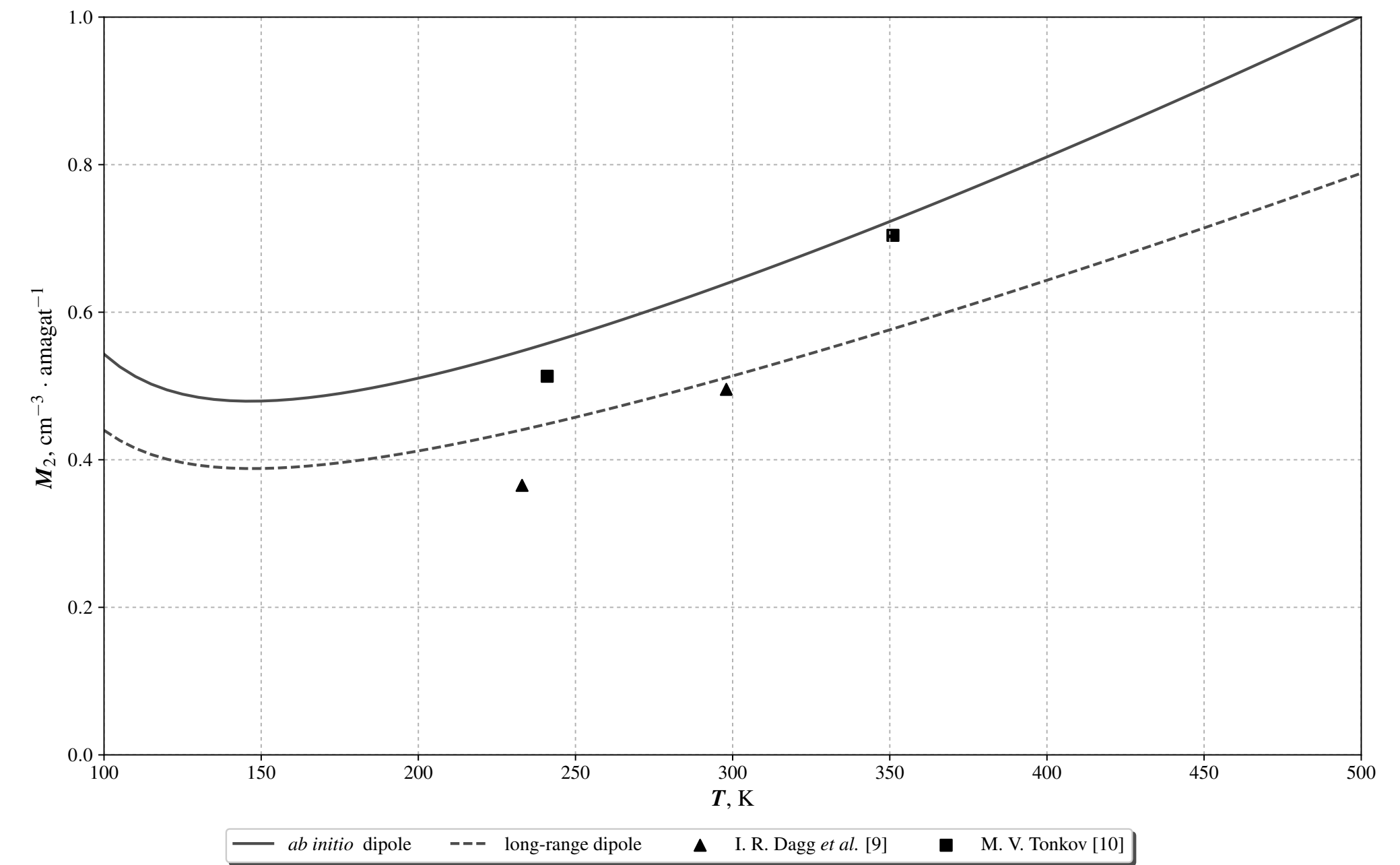


Figure 6: Temperature dependence of second spectral moment

Conclusion

The results reported in the present paper demonstrate capability of classical methods to characterize collision-induced absorption in the far-infrared spectral range with an accuracy comparable with that experimental. No use of adjustable parameters is required provided potential energy and induced dipole surfaces are obtained as a result of the high-level *ab initio* calculations and no omission is made in kinetic energy terms of the classical Hamiltonian. The results obtained so far encourage us to extend our approach to other molecular pairs the knowledge of the dipole-forbidden absorption in which is in high demand by the planetary atmospheres investigators.

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