# Vectorial Calculus A1 Recap

# 10/04/2025

# Spherical and Cylindrical coordinates

Everything described in this section has an easier version in  $\mathbb{R}^2$ 

### Cylindrical coordinates

The transformation from the ordinary space  $\mathbb{R}^3$  to  $C(\mathbb{R}^3)$  is nearly analogous to polar coordinates in  $\mathbb{R}^2$  described below:

$$T(x, y, z) = (r \cos \theta, r \sin \theta, z)$$

, and if we were integrating some function  $f: \mathbb{R}^3 \to \mathbb{R}$  under the region  $A \subset \mathbb{R}^3$  in the space (x, y, z), the equivalent integral in the new system  $C(\mathbb{R}^3)$  is:

$$\iiint_{C(A)} f(r\cos\theta, r\sin\theta, z) r dz dr d\theta$$

### **Spherical coordinates**

From the reasoning used in cylindrical coordinates, we have:

$$\iiint_A f(x,y,z) dx dy dz = \iiint_{S(A)} f(\rho \sin \varphi \cos \theta, \rho \sin \varphi \sin \theta, \rho \cos \theta) \rho^2 \sin \varphi d\rho d\theta d\varphi$$

# Jacobian Determinant For Multiple Integrals

More general change of variables require some technicalities:

If  $T: \mathbb{R}^3 \to \mathbb{R}^3$  is a bijective function that morphes a region in the space xyz into the space uvw throught the equations that follow:

$$x = g(u, v, w)$$
$$y = h(u, w)$$
$$z = k(u, v, w)$$

Then the **jacobian determinant of T** is:

$$\frac{\partial(x,y,z)}{\partial(u,v,w)} = \det \begin{pmatrix} \frac{\partial x}{\partial u} & \frac{\partial x}{\partial v} & \frac{\partial x}{\partial w} \\ \frac{\partial y}{\partial u} & \frac{\partial y}{\partial v} & \frac{\partial y}{\partial w} \\ \frac{\partial z}{\partial u} & \frac{\partial z}{\partial v} & \frac{\partial z}{\partial w} \end{pmatrix}$$

And the equivalent integral onto the new system is:

$$\iiint_A f(x,y,z) dx dy dz = \iiint_{T(A)} f(x(u,v,w),y(u,v,w),z(u,v,w)) \frac{\partial (x,y,z)}{\partial (u,v,w)} du dv dw$$

# **Physics**

## **Mass Center and Centroid**

Given a plane region  $S^2 \subset \mathbb{R}^2$ , its **centroid** is the point  $(x_c, y_c)$ , where:

$$x_c = \frac{1}{\operatorname{area}(S)} \iint x dx dy$$

$$y_c = \frac{1}{\operatorname{area}(S)} \iint_S y dx dy$$

If the region/body/whatever has **constant** density  $\mu(x,y)$ ,  $\forall x,y \in \mathbb{R}$ , then its centroid is the same as the mass center, which is the point  $(\hat{x_c}, \hat{y_c})$  where:

$$\hat{x_c} = \frac{\iint_S \mu(x, y) x dx dy}{\iint_S \mu(x, y) dx dy}$$

$$\hat{y_c} = \frac{\iint_S \mu(x, y) y dx dy}{\iint_S \mu(x, y) dx dy}$$

Notice that the **mass** of S is given by:

$$\iint_{S} \mu(x,y) dx dy$$

## **Moment of Inertia**

Let  $X \subset \mathbb{R}^3$  be a body, rotating over a given axis, let  $\mu(x)$  be the mass density in  $x, \forall x \in X$ , if r(X) is the distance to axis it is rotating over, and v(x) is the speed,  $\forall x \in X$ , then  $|v(x)| = \omega r(x)$ , where  $\omega$  is the angular speed

It follows that the **kinetic energy** of the body is  $\frac{mv^2}{2}$ :

$$\frac{1}{2}\cdot\int_X \mu(x)\ |v(x)|^2\ dX = \frac{1}{2}\omega^2\cdot\int_X \mu(x)r(x)^2dX$$

We define now the **Moment of Inertia** as:

$$I = \int_X \mu(x) r(x)^2 dX.$$

 $L = \omega I$  is called the **angular momentum**, it is conserved if there are no external rotational forces.

# **Vector Calculus**

#### Curves

A Curve is a continuous function  $\gamma:[a,b]\to\mathbb{R}^n$ , it is of class  $C^1$  if  $\gamma'$  exists and is continuous in [a,b], if  $\gamma(a)=\gamma(b)$ , the curve is **closed**.

A curve is said to be  $C^1$  by parts if there is a partition of [a, b] in a finite number of subintervals such that the curve is  $C^1$  in each of tese subintervals.

# **Scalar Line Integrals**

Given  $f: \mathbb{R}^n \to \mathbb{R}$  a function and  $\gamma: [a,b] \to \mathbb{R}^n$  a  $C^1$  class curve in  $\mathbb{R}^n$ , the **scalar line integral** of f along  $\gamma$  is:

$$\int_{\gamma} f ds = \int_{a}^{b} f(\gamma(t)) \, \left\| \gamma'(y) \right\| \, dt$$

If  $\gamma$  is  $C^1$  by parts, we integrate on the  $C^1$  partition-subintervals and sum each odf the smaller integrals.

## Centroid and Mass Center of a Curve

## **Mass Cernter**

Given  $\gamma \subset \mathbb{R}^3$  a curve, and let f(x, y, z) be the mass density per unit of length of  $\gamma$ , we know that the  $\gamma$ 's mass is given by:

$$M = \int_{\gamma} f(x, y, z) ds.$$

So the mass center of  $\gamma$  is the point  $c_m = (x_c, y_c, z_c)$  s.t:

$$\begin{split} x_c &= \frac{\int_{\gamma} x f(x,y,z) ds}{M} \\ y_c &= \frac{\int_{\gamma} y f(x,y,z) ds}{M} \\ z_c &= \frac{\int_{\gamma} z f(x,y,z) ds}{M} \end{split}$$

If f is constant (homogeneous curve), then the mass center is the centroid as well

### Centroid

Let  $\gamma \subset \mathbb{R}^3$  be a smooth or piecewise smooth curve, parameterized by  $\gamma(t):[a,b] \to \mathbb{R}^3$ . If the curve is **homogeneous**, meaning the mass density per unit length is constant, then its **centroid** is the point  $(x_c, y_c, z_c)$  given by:

$$x_c = \frac{1}{L} \int_{\gamma} x ds$$
 
$$y_c = \frac{1}{L} \int_{\gamma} y ds$$
 
$$z_c = \frac{1}{L} \int_{\gamma} z ds$$

Where L is the total arc length of the curve:

$$L = \int_{\gamma} ds = \int_a^b \|\gamma'(t)\| \ dt$$

# **Vectorial Line Integrals**

Consider now  $F: \mathbb{R}^n \to \mathbb{R}^n$ , usually called a *vector field*, and a class  $C^1$  curve  $\gamma: [a,b] \to \mathbb{R}^n$  in this field.

The integral of F along  $\gamma$  is:

$$\int_{\gamma} F = \int_{a}^{b} F(\gamma(t))\gamma'(t)dt$$

This line integral is linear:  $\int_{\gamma} (aF+bG) = a \int_{\gamma} F + b \int_{\gamma} G$ 

# **Conservative Vector Fields and Angle Variation**

#### **Foreword**

We have been trained in the mysterious and dark arts of newtonian and lagrangian mechanics by Master Paulo Verdasca Amorim himself, this is nothing to us.

#### **Conservative Vector Fields**

A field  $F: \Omega \subset \mathbb{R}^n \to \mathbb{R}^n$ ,  $\Omega$  an open and connected set, is said to be conservative if  $\exists U: \Omega \to \mathbb{R}$  s.t  $F = \Delta U$ , U is called F's **potential** 

**Theorem**: Let  $F: \Omega \to \mathbb{R}^n$  be a vector field and  $f: \Omega \to \mathbb{R}$  be its potential, i.e  $F = \Delta f$ . Let  $A, B \in \Omega$  and  $c: [a,b] \to \mathbb{R}^n$  be a  $C^1$  by parts curve s.t c(a) = A and c(b) = B, then:

$$\int_{C} F = f(B) - f(A)$$

This looks like Calculus' Fundamental Theorem

### **Angle Variation**

We conclude this section on line integrals with a counterexample: a vector field that satisfies  $\partial \frac{F_1}{\partial} y = \partial \frac{F_2}{\partial} x$  but is not conservative.

Let  $F(x,y) = \frac{1}{x^2+y^2}(-y,x)$ , defined over  $\mathbb{R}^2 \smallsetminus 0$ . Note that:

$$\partial \frac{F_1}{\partial} y = \partial \frac{F_2}{\partial} x = \frac{y^2 - x^2}{\left(x^2 + y^2\right)^2}$$

So the field meets the symmetry of mixed partials, but F is still not conservative — the domain is not simply connected.

Take the closed curve  $c:[0,2\pi]\to\mathbb{R}^2$  given by:

$$c(t) = (\cos t, \sin t)$$

Then:

$$\int_{c} F = \int_{0}^{\{2\pi\}} (-\sin t, \cos t) \cdot (-\sin t, \cos t) dt = \int_{0}^{\{2\pi\}} 1 dt = 2\pi \neq 0$$

Since the line integral over a closed curve is nonzero, F is not conservative.