

# Subatomic Physics II

## Problem Set 10

Due on December 16, 2021, 11:59 PM

### Problem 10: Geoneutrinos

Decay of radioactive elements such as  $^{238}\text{U}$  and  $^{232}\text{Th}$  inside the Earth is a source of terrestrial heat (radiogenic heat) and also of neutrinos, which are called *geoneutrinos*. Although the total heat production inside Earth is known ( $47 \pm 2$  TW), the radiogenic heat is very poorly constrained. According to different models, the radiogenic power can be anything from 3 to 30 TW. Measuring the geoneutrino flux will hugely improve this estimate.

Let us estimate whether such a measurement is feasible.

- Assume that the antineutrino is detected by the inverse  $\beta$ -decay  $\bar{\nu}_e p \rightarrow e^+ n$  it induces in the detector. Calculate the neutrino energy threshold  $E_0$  in MeV. (1pt)
- The uranium decay chain  $^{238}\text{U} \rightarrow ^{206}\text{Pb}$  is an important source of radiogenic heat and geoneutrinos. This decay releases 52 MeV of energy per  $^{238}\text{U}$  nucleus and produces a number of antineutrinos (how many?). On average one of these  $\bar{\nu}_e$  has an energy above the threshold  $E_0$ . Assuming that uranium decays are responsible for 8 TW of radiogenic heat, calculate the geoneutrino flux  $f$  (measured in  $\text{cm}^{-2}\text{s}^{-1}$ ) through the Earth's surface. Assume for simplicity that uranium is homogeneously distributed inside the Earth. (2pt)
- Neutrinos oscillate on their way from the Earth's interior to the detector. The probability that, after travelling a distance  $L$ , a  $\bar{\nu}_e$  remains  $\bar{\nu}_e$ , is given by

$$P_{ee} = \sin^4 \theta_{13} + \cos^4 \theta_{13} \left[ 1 - \sin^2 2\theta_{12} \sin^2 \left( \frac{\Delta m_{21}^2 L}{4E_\nu} \right) \right],$$

where  $\theta_{12}$ ,  $\theta_{13}$  and  $\Delta m_{21}^2$  are the neutrino oscillation parameters. Calculate the oscillation length for a 2.2 MeV neutrino. Comparing it with the Earth's radius, deduce the *average* survival probability  $\langle P_{ee} \rangle$  for geoneutrinos to reach the detector on the surface in the  $\bar{\nu}_e$  state. (2pt)

- The cross section of the process  $\bar{\nu}_e p \rightarrow e^+ n$  is given by the formula

$$\sigma = \frac{G_F^2}{\pi} (g_V^2 + 3g_A^2) (E_\nu - \Delta) \sqrt{(E_\nu - \Delta)^2 - m_e^2},$$

where  $G_F$  is the Fermi constant,  $g_V = 1$ ,  $g_A \approx 1.23$  are the vector and axial coupling constants, and  $\Delta = m_n - m_p$ . For an average geoneutrino energy  $E_\nu = 2.2$  MeV, calculate  $\sigma$  in  $\text{cm}^2$ . (1pt)

- Consider a neutrino detector filled with 1000 tons of water. Using the above results and assuming 100% detection efficiency, estimate the number of detected geoneutrino events per year. (2pt)
- It is known that the Sun produces a huge neutrino flux, of the order of  $10^{10} \text{ cm}^{-2}\text{s}^{-1}$  at the surface of the Earth, in the same energy range. Why is it nevertheless possible to measure geoneutrinos despite such a huge solar background? What do you think is another important source of background, and how can one control this? (2pt)