

# Subatomic Physics II: Problem set 7

Arthur Adriaens

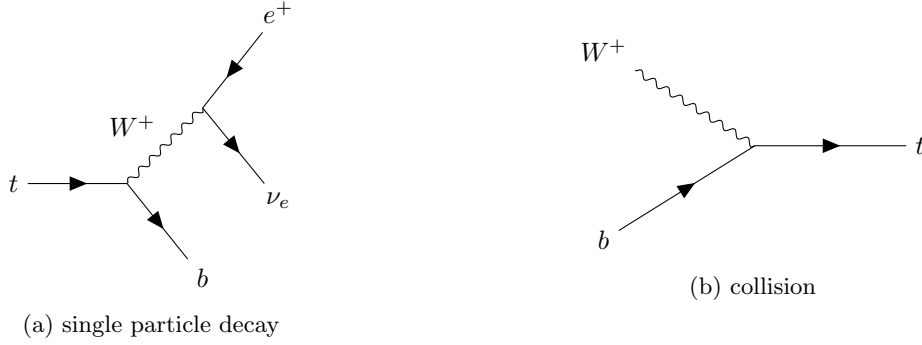
## 7.1 Feynman rules

### 7.1.1 $Ze\mu$

We can say that this is not possible as lepton flavour violation would occur.

### 7.1.2 $Wtb$

This is a possible vertex with the following examples for feynmann diagrams:



#### Single particle decay

Here we have (COM frame):

$$s = (P_{W^+}^\mu + P_b^\mu)^2 g_{\mu\nu} \quad (1)$$

$$s = m_{W^+}^2 + m_b^2 + 2(E_{W^+}E_b - p_{W^+} \cos(\theta)p_b \cos(\theta') + p_{W^+} \sin(\theta)p_b \sin(\theta')) \quad (2)$$

$$s \geq m_{W^+}^2 + m_b^2 + 2(E_{W^+}E_b)|_{\vec{p}_{w,b}=0} = m_{W^+}^2 + m_b^2 + 2(m_{W^+}m_b) \quad (3)$$

$$\sqrt{s} \geq 84.56 \text{ GeV} \quad (4)$$

$$(5)$$

Where  $m_b = 4.18 \text{ GeV}$  and  $m_{W^+} = 80.379 \text{ GeV}$ <sup>1</sup>, now as  $m_t = 172.76 \text{ GeV}$  we see that we definitely have the required center of mass energy, even if the  $W^+$  boson were on-shell (which it doesn't have to be).

#### Collision

Here we have the following four-vectors (COM frame):

$$P_{W^+}^\mu = (E_{W^+}, 0, -p, 0) \quad (6)$$

$$P_b^\mu = (E_b, 0, p, 0) \quad (7)$$

Doing analogues calculations as the previous one:

$$s = (P_{W^+}^\mu + P_b^\mu)^2 g_{\mu\nu} \quad (8)$$

$$s = m_{W^+}^2 + m_b^2 + 2(E_{W^+}E_b + p^2) \quad (9)$$

$$s = m_{W^+}^2 + m_b^2 + 2 \left( \sqrt{m_{W^+}^2 + p^2} \sqrt{m_b^2 + p^2} + p^2 \right) \quad (10)$$

$$(11)$$

Now as  $m_\tau^2 \leq s$  we have that  $p \geq 67.6 \text{ GeV}$  for the t-quark to be on shell, in which  $W^+$  has to be on-shell.

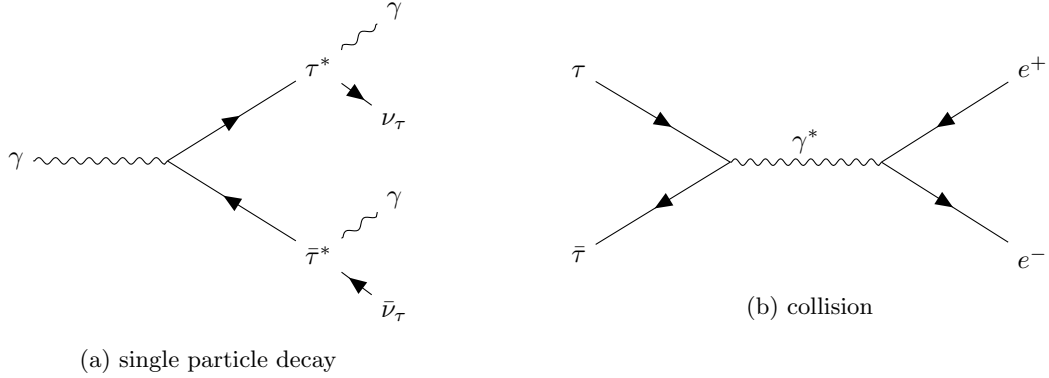
<sup>1</sup>All the masses used in this work are obtained from the PDG [2]

### 7.1.3 $H\gamma\gamma$

Although this is in principle 'possible' as "the Higgs boson can decay to all Standard Model particles" [1], the coupling strength of the Higgs boson is proportional to the mass of the particles involved and thus this will tend to zero for the photon.

### 7.1.4 $\gamma\tau\tau$

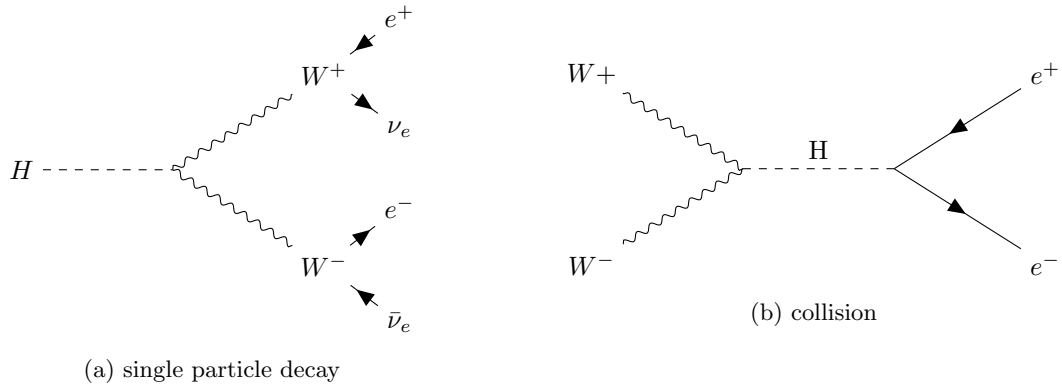
The decay isn't possible for on-shell particles as we could always Lorentzboost to a frame of reference where 3-momentum  $\mathbf{p}$  isn't conserved (e.g the  $\tau$  particles' COM frame). It is however possible for the photon to decay into 2 off-shell  $\tau$  leptons (denoted  $\tau^*$  and  $\bar{\tau}^*$ ) as then we can momentarily violate 3-momentum conservation:



And, as seen from the figure, it's also possible for 2 on-shell tau leptons to decay into a virtual photon which decays into an electron-positron pair, as  $m_\tau \gg m_e$  there's no need to concern ourselves with the COM calculations.

### 7.1.5 $HWW$

This is a possible vertex with the following examples:



#### Single particle decay

This is possible but as  $m_H < 2m_W$  (125.25GeV < 160.758GeV) One of the  $W$  bosons produced is off-shell.

#### Collision

As  $m_H < 2m_W$  at least one of the  $W$  bosons has to be virtual, we can understand this by looking at the COM frame. before the interaction we have:

$$s = (P_{W^+}^\mu + P_{W^-}^\mu)^2 g_{\mu\nu} \quad (12)$$

$$s = m_{W^+}^2 + m_{W^-}^2 + 2(E_{W^+}E_{W^-} + pp) \quad (13)$$

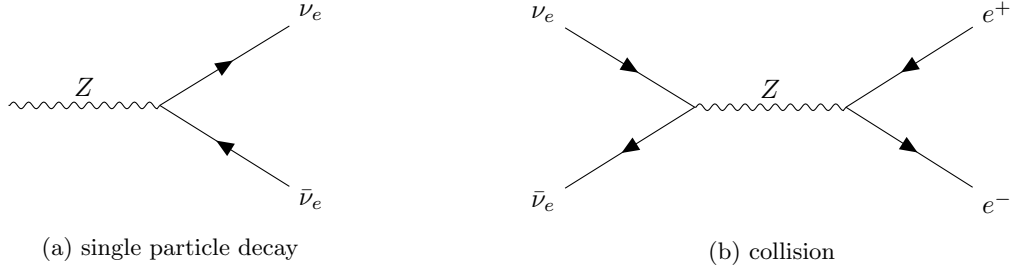
$$s_{\min} \stackrel{p \rightarrow 0}{=} m_{W^+}^2 + m_{W^-}^2 + 2m_{W^+}m_{W^-} \quad (14)$$

$$(15)$$

Which implies a minimal COM energy way bigger than 125.25GeV if both bosons are on-shell.

### 7.1.6 $Z\nu_e\nu_e$

This is a possible vertex with example diagrams:



#### Single particle decay

As we're working with the tree-level Feynman rules of the Standard Model, we assume the neutrino to be massless (not that this matters much as  $m_{\nu_e} < 1.1\text{eV}$ ). In the COM frame we thus have:

$$m_Z^2 = s = (P_{\nu_e}^\mu + P_{\bar{\nu}_e}^\mu)^2 g_{\mu\nu} \quad (16)$$

$$= 2E^2 \quad (17)$$

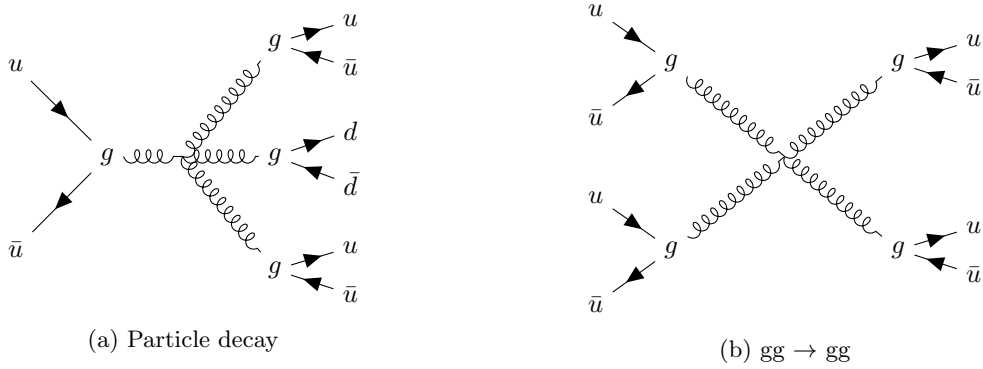
I.e the neutrinos have energies of  $\frac{1}{\sqrt{2}}m_Z = \frac{1}{\sqrt{2}}91.1876\text{GeV} = 64.48\text{GeV}$ .

#### Collision

If the neutrinos in this interaction have energies of 64.48GeV, then this is a possible on-shell Z-boson production mechanism. It's also possible to produce off-shell Z-bosons with less energetic neutrinos.

### 7.1.7 $gggg$

This is a possible vertex with the following feynmann diagrams:



As all the gluons are off-shell (due to the specific nature of gluons and color confinement), the kinematic condition is reduced to if there's enough COM energy to produce the observed pions:

#### Decay

Here we have that the end COM energy  $\sqrt{s} \geq 3 \times m_{\pi^0} = 404.91 \text{ MeV}$ , the com energy on the left is:

$$\sqrt{s} = \sqrt{2m_u^2 + 2P_u P_{u'}} = \sqrt{2m_u^2 + 2m_u^2 + |\mathbf{p}|^2} \approx |\mathbf{p}| \quad (18)$$

And thus  $|\mathbf{p}| \geq 405 \text{ MeV}$ .

#### Collision

In the collision, the incoming and outgoing particles are the same and so because of conservation of energy there are no kinematic terms that need to be computed.

## References

- [1] Mark Thomson. *Modern Particle Physics*. 2019.
- [2] P.A. Zyla et al. Review of Particle Physics. *PTEP*, 2020(8):083C01, 2020.