



01OAIQD – Dynamic Design of Machines

Academic year 2018-2019

Discrete linear systems

TUTORIAL 2

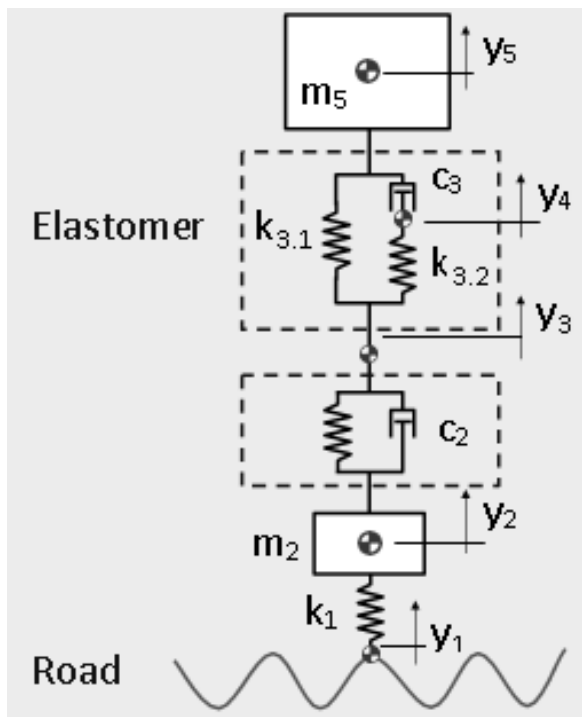
Emission date: November 22, 2019		GROUP:
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EXERCISE 1: Quarter car model

Figure 1 shows the so-called quarter car model, one of the simplest models used to study the dynamic behavior of motor vehicle suspensions. The upper mass m_5 simulates the part of the mass of the car body (the sprung mass) that can be considered supported by a given wheel, while the lower mass m_2 simulates the wheel and all the parts that can be considered as rigidly connected with the unsprung mass. The two masses are connected by a spring-damper system simulating the suspension (k_2, c_2) and the silent block (elastometer, $k_{3,1}, k_{3,2}, c_3$). The unsprung mass is connected to the ground with a second spring simulating the radial stiffness of the tire. The point at which the tire contacts the ground is assumed to move in a vertical direction with a given law $y_1(t)$, and it simulates the motion on uneven ground.

For the given quarter car model

- determine the equations of dynamic equilibrium using the Lagrangian approach,
- arrange the equations in matrix form,
- find the natural frequencies and the mode shapes of the system,
- Consider as input a harmonic excitation with amplitude y_1 of 3 mm @ 5 Hz. Compute the power dissipated in the elastomeric member.



sprung mass	$m_5 = 400 \text{ kg};$
unsprung mass	$m_2 = 30 \text{ kg};$
spring stiffness	$k_2 = 24 \text{ kN/m}$
spring damping	$c_2 = 1200 \text{ Ns/m};$
elastomer stiffness	$k_{3.1} = 150 \text{ kN/m};$
elastomer stiffness	$k_{3.2} = 90 \text{ kN/m};$
elastomer damping	$c_3 = 800 \text{ Ns/m};$
tire stiffness	$k_1 = 190 \text{ kN/m};$

Figure 1

SOLUTION

1 -

To compute the equation of motion using the Lagrangian approach it is necessary to determine:

- the kinetic energy $T = \frac{1}{2} m_2 \dot{y}_2^2 + \frac{1}{2} m_5 \dot{y}_5^2$
- the potential energy $U = \frac{1}{2} k_1 (y_2 - y_1)^2 + \frac{1}{2} k_2 (y_3 - y_2)^2 + \frac{1}{2} k_{3.1} (y_5 - y_3)^2 + \frac{1}{2} k_{3.2} (y_4 - y_3)^2$
- $\mathcal{F} = \frac{1}{2} C_2 (\dot{y}_3 - \dot{y}_2)^2 + \frac{1}{2} C_3 (\dot{y}_5 - \dot{y}_4)^2$

Thus, the Lagrange equation related to the 4 DoFs are:

$$\frac{d}{dt} \left(\frac{\partial T}{\partial \dot{q}_i} \right) + \frac{\partial U}{\partial q_i} + \frac{\partial \mathcal{F}}{\partial \dot{q}_i} = \frac{\partial (\delta L)}{\partial (\delta q_i)}$$

$$1) \frac{d}{dt} \left(\frac{\partial T}{\partial \dot{y}_2} \right) = m_2 \ddot{y}_2$$

$$\frac{\partial U}{\partial y_2} = k_1 (y_2 - y_1) - k_2 (y_3 - y_2)$$

$$\frac{\partial \mathcal{F}}{\partial \dot{y}_2} = -C_2 (\dot{y}_3 - \dot{y}_2)$$

$$\frac{\partial (\delta L)}{\partial (\delta y_2)} = 0$$

The first equation will be :

$$m_2 \ddot{y}_2 + C_2 \dot{y}_2 - C_2 \dot{y}_3 + y_2 (k_1 + k_2) - k_2 y_3 = k_1 y_1$$

$$2) \frac{d}{dt} \left(\frac{\partial T}{\partial \dot{y}_3} \right) = 0$$

$$\frac{\partial U}{\partial y_3} = k_2 (y_3 - y_2) - k_{3.1} (y_5 - y_3) - k_{3.2} (y_4 - y_3)$$

$$\frac{\partial \mathcal{F}}{\partial \dot{y}_3} = C_2 (\dot{y}_3 - \dot{y}_2)$$

$$\frac{\partial (\delta L)}{\partial (\delta y_3)} = 0$$

The second equation will be:

$$-C_2\dot{y}_2 + C_2\dot{y}_3 - k_2y_2 + y_3(k_{3.1} + k_{3.2} + k_2) - k_{3.2}y_4 - k_{3.1}y_5 = 0$$

$$3) \frac{d}{dt} \left(\frac{\partial T}{\partial \dot{y}_4} \right) = 0$$

$$\frac{\partial U}{\partial y_4} = k_{3.2}(y_4 - y_3)$$

$$\frac{\partial \mathcal{F}}{\partial \dot{y}_4} = -C_3(\dot{y}_5 - \dot{y}_4)$$

$$\frac{\partial(\delta L)}{\partial(\delta y_4)} = 0$$

The third equation will be :

$$C_3\dot{y}_4 - C_3\dot{y}_5 - k_{3.2}y_3 + k_{3.2}y_4 = 0$$

$$4) \frac{d}{dt} \left(\frac{\partial T}{\partial \dot{y}_5} \right) = m_5\ddot{y}_5^2$$

$$\frac{\partial U}{\partial y_5} = k_{3.1}(y_5 - y_3)$$

$$\frac{\partial \mathcal{F}}{\partial \dot{y}_5} = C_3(\dot{y}_5 - \dot{y}_4)$$

$$\frac{\partial(\delta L)}{\partial(\delta y_5)} = 0$$

The forth equation will be :

$$m_5\ddot{y}_5^2 - C_3\dot{y}_4 + C_3\dot{y}_5 - k_{3.1}y_3 + k_{3.1}y_5 = 0$$

2 -

In matrix form:

$$\begin{bmatrix} m_2 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & m_5 \end{bmatrix} \begin{Bmatrix} \ddot{y}_2 \\ \ddot{y}_3 \\ \ddot{y}_4 \\ \ddot{y}_5 \end{Bmatrix} + \begin{bmatrix} C_2 & -C_2 & 0 & 0 \\ -C_2 & C_2 & 0 & 0 \\ 0 & 0 & C_3 & -C_3 \\ 0 & 0 & -C_3 & C_3 \end{bmatrix} \begin{Bmatrix} \dot{y}_2 \\ \dot{y}_3 \\ \dot{y}_4 \\ \dot{y}_5 \end{Bmatrix} + \begin{bmatrix} k_1 + k_2 & -k_2 & 0 & 0 \\ -k_2 & k_2 + k_{3.1} + k_{3.2} & -k_{3.2} & -k_{3.1} \\ 0 & -k_{3.2} & k_{3.2} & 0 \\ 0 & -k_{3.1} & 0 & k_{3.1} \end{bmatrix} \begin{Bmatrix} y_2 \\ y_3 \\ y_4 \\ y_5 \end{Bmatrix} = \begin{Bmatrix} k_1 y_1 \\ 0 \\ 0 \\ 0 \end{Bmatrix}$$

3 -

For computing the natural frequency

$$\{y_i\} = \{y_0\}e^{i\omega t} \quad \{\dot{y}_i\} = i\omega\{y_0\}e^{i\omega t} \quad \{\ddot{y}_i\} = -\omega^2\{y_0\}e^{i\omega t}$$

$$([k] - [m]\omega_{ni}^2)\{y_0\} = \{0\}$$

$$\det([k] - [m]\omega_{ni}^2)\{y_0\} = \{0\}$$

$$\omega_{n1} = 0, \quad \omega_{n2} = 0.2159, \quad \omega_{n3} = 2.5611, \quad \omega_{n4} = inf$$

The mode shapes:

First mode:

$$\begin{Bmatrix} y_{02} \\ y_{03} \\ y_{04} \\ y_{05} \end{Bmatrix} = \begin{Bmatrix} 0 \\ 1 \\ 1 \\ 0 \end{Bmatrix}$$

Second mode:

$$\begin{Bmatrix} y_{02} \\ y_{03} \\ y_{04} \\ y_{05} \end{Bmatrix} = \begin{Bmatrix} -0.989 \\ 0.8757 \\ 0.8757 \\ -1 \end{Bmatrix}$$

Third mode:

$$\begin{pmatrix} y_{02} \\ y_{03} \\ y_{04} \\ y_{05} \end{pmatrix} = \begin{pmatrix} 1 \\ 0.1315 \\ 0.1315 \\ -0.0074 \end{pmatrix}$$

Fourth mode:

$$\begin{pmatrix} y_{02} \\ y_{03} \\ y_{04} \\ y_{05} \end{pmatrix} = \begin{pmatrix} 0 \\ -1 \\ 1 \\ 0 \end{pmatrix}$$

4 –

Starting directly from the equation of motion in matrix form:

$$M\ddot{Y} + C\dot{Y} + KY = F$$

Harmonic excitation

$$F = F_0 e^{j\omega t}$$

Where F_0 is something like $F_0 = [0 \dots k_1 y_1 \dots 0]^T$

Hence

$$\begin{aligned} Y &= Y_0 e^{j\omega t} \\ \dot{Y} &= j\omega Y_0 e^{j\omega t} \\ \ddot{Y} &= -\omega^2 Y_0 e^{j\omega t} \\ \Rightarrow (-\omega^2 M + j\omega C + K)Y_0 &= F_0 \end{aligned}$$

Let's define $(-\omega^2 M + j\omega C + K) = K_{dyn}(\omega)$

So,

$$Y_0 = \text{inv}(K_{dyn}(\omega)) F_0$$

The response to the excitation at $\omega_0 = 2\pi 5$ rad/s is

$$Y_0 = \text{inv}(K_{dyn}(\omega_0)) F_0$$

Y_0 is the amplitude of the response. It is something like $Y_0 = [y_{02} \ y_{03} \ y_{04} \ y_{05}]^T$, where the y_{0i} component could be complex quantities.

The power dissipated in the elastomeric member is due to the lumped damping c_3 since it is the dissipative member (the spring does not dissipate energy). The power is force times velocity. Hence, the power dissipated by the elastomeric member is

$$P = F_3 \dot{y}_4 = (c_3 \dot{y}_4) \dot{y}_4 = c_3 \dot{y}_4^2$$

Where $y_4 = \text{Re}(y_{04} e^{j\omega_0 t})$ (note that in the frequency domain we have complex quantities but in the time domain we are interested in the real part only). Consequently, we have

$$\dot{y}_4 = \text{Re}(j\omega_0 y_{04} e^{j\omega_0 t}) = \text{Re}\left(\omega_0 y_{04} e^{j\omega_0 t + \frac{\pi}{2}}\right)$$

y_{04} could be a complex quantity. Hence, it can be written as

$$y_{04} = |y_{04}| e^{j\phi}$$

By substituting, we have

$$\dot{y}_4 = \text{Re}\left(\omega_0 |y_{04}| e^{j(\omega_0 t + \frac{\pi}{2} + \phi)}\right)$$

Where, physically speaking, ϕ is the phase delay between the harmonic input y_1 and the output displacement y_4 .

By picking the real part only, it follows

$$\dot{y}_4 = |\dot{y}_4| \cos\left(\omega_0 t + \frac{\pi}{2} + \phi\right) = \omega_0 |y_{04}| \sin(\omega_0 t + \phi)$$

Therefore, the instant power is

$$P = c_3 \omega_0^2 |y_{04}|^2 (\sin(\omega_0 t + \phi))^2$$

The peak power dissipated by the elastomeric member is

$$P_{peak} = c_3 \omega_0^2 |y_{04}|^2$$

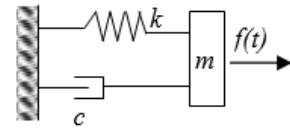
And the RMS value is

$$P_{RMS} = P_{peak} / \sqrt{2}$$

EXERCISE 2

Determine the unit impulse response of the mass-damper-spring system in figure. Plot the response and

- evaluate the first peaks of the response;
- calculate the damping ratio of the system



$$k = 100 \text{ N/m}, \quad m = 2 \text{ kg}, \quad c = 2 \text{ N/m s}.$$

SOLUTION

Determine the unit impulse response of the system

Since the impulse determines a variation of the momentum,

$$m(v_2 - v_1) = \int_v f dt$$

And if $v_1 = 0$, the system is at rest in the origin

$$m v_2 = f \rightarrow v_2 = \frac{f}{m}$$

Then, the initial conditions are $\begin{cases} \dot{x}(0) = f/m \\ x(0) = 0 \end{cases}$

Since impulse is not a proper external force acting on the system, I need to study the associated homogeneous differential equation:

$$m \ddot{x} + c \dot{x} + kx = 0$$

Whose solution is:

$$x(t) = e^{-\xi \omega_n t} [x_{01} \cos(\omega_p t) + x_{02} \sin(\omega_p t)]$$

To calculate x_{01} and x_{02} , the initial conditions are considered

$$\begin{cases} x(0) = x_{01} = 0 \\ \dot{x}(0) = -\xi \omega_n e^{-\xi \omega_n t} [x_{01}] + e^{-\xi \omega_n t} [-x_{01} \omega_p \sin(\omega_p t) + \omega_p x_{02} \sin(\omega_p t)] = -\xi \omega_n x_{01} + \omega_p x_{02} = \frac{f}{m} \end{cases}$$

So, $\begin{cases} x_{01} = 0 \\ x_{02} = \frac{f}{m \omega_p} \end{cases}$

$$\text{Then, } x(t) = e^{-\xi \omega_n t} \left[\frac{f}{m \omega_p} \sin(\omega_p t) \right].$$

The first peak response is when $\sin(\omega_p t)$ is maximum, so:

$$\sin(\omega_p t) = 1, \text{ then } \omega_p t = \frac{\pi}{2} \rightarrow t = \omega_p^{-1} \frac{\pi}{2}$$

And substituting in $x(t)$:

$$x(t) = e^{-\frac{\xi \omega_n \pi}{\omega_p^2}} \left[\frac{f}{m \omega_p} \sin\left(\frac{\pi}{2}\right) \right] \quad \text{where } \omega_n = 5\sqrt{2} \text{ rad/s}$$

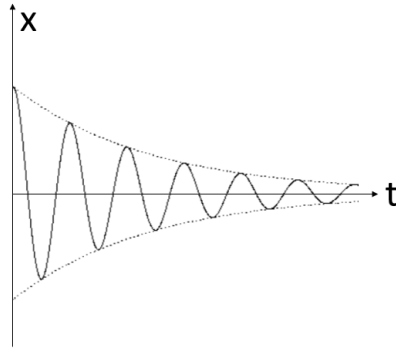
$$\omega_p = \omega_n \sqrt{1 - \xi^2} = 6.816$$

$$\xi = \frac{c}{2\sqrt{km}} = \sqrt{2}/20$$

Then, $x_{max} = 0.065 \text{ m}$

2)

The damping ratio of the system is given by the ratio between the amplitudes of two subsequent peaks:

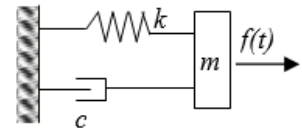


$$\delta = \ln \frac{x_i}{x_{i+1}} = \ln \frac{e^{-\xi \omega_n t_i}}{e^{-\xi \omega_n t_{i+1}}} = \ln e^{-\xi \omega_n (t_i - t_{i+1})} = \ln e^{\xi \omega_n \frac{2\pi}{\omega_n \sqrt{1-\xi^2}}} = \frac{2\pi\xi}{\sqrt{1-\xi^2}}$$

EXERCISE 3

Calculate the unit step response $s(t)$ of the mass-damper-spring system of exercise 1 integrating the impulse response. Plot $s(t)$ versus t .

- Determine the response for the undamped system ($c=0$) and,
- for the damped system as given in the picture.
- Which is the response at infinite?



SOLUTION

Calculate the unit step response $s(t)$ of the system

The step response is the integral of the impulse response

$$\begin{aligned}
 x(t) &= \int_{-\infty}^t e^{-\xi \omega_n t} \left[\frac{f}{m \omega_p} \sin(\omega_p t) \right] dt = \frac{f}{m \omega_p} \int_{-\infty}^t \underbrace{e^{-\xi \omega_n t}}_{b'} \underbrace{\sin(\omega_p t)}_a dt \\
 &\quad \int a b' = ab - \int a' b \\
 x(t) &= \frac{f}{m \omega_p} [-\sin(\omega_p t)] \frac{e^{-\xi \omega_n t}}{\xi \omega_n} + \frac{\omega_p}{\xi \omega_n} \int_{-\infty}^t \cos(\omega_p t) e^{-\xi \omega_n t} dt = \\
 &= \frac{f}{m \omega_p} \left[-\sin(\omega_p t) \frac{e^{-\xi \omega_n t}}{\xi \omega_n} + \frac{\omega_p}{\xi \omega_n} \left[-\frac{e^{-\xi \omega_n t}}{\xi \omega_n} \cos(\omega_p t) + \int_{-\infty}^t -\omega_p \frac{\sin(\omega_p t)}{\xi \omega_n} e^{-\xi \omega_n t} dt \right] \right] = \\
 &= \frac{f}{m \omega_p} \left[\int_{-\infty}^t e^{-\xi \omega_n t} \sin(\omega_p t) dt + \frac{\omega_p^2}{\xi^2 \omega_n^2} \int_{-\infty}^t \sin(\omega_p t) e^{-\xi \omega_n t} dt \right] = \\
 &= \frac{f}{m \omega_p} \left[\frac{e^{-\xi \omega_n t}}{\xi \omega_n} \sin(\omega_p t) - \frac{\omega_p}{\xi^2 \omega_n^2} e^{-\xi \omega_n t} \cos(\omega_p t) \right] = \\
 &= \frac{f}{m \omega_p} \left(-\frac{e^{-\xi \omega_n t}}{\xi \omega_n} \right) \left[\sin(\omega_p t) + \frac{\omega_p}{\xi \omega_n} \cos(\omega_p t) \right] = \\
 \int_0^t e^{-\xi \omega_n t} \sin(\omega_p t) dt &= -\left(\frac{\xi^2 \omega_n^2}{\xi^2 \omega_n^2 + \omega_p^2} \right) \frac{f}{m \omega_p} \frac{e^{\xi \omega_n t}}{\xi^2 \omega_n^2} (\xi \omega_n \sin(\omega_p t) + \omega_p \cos(\omega_p t)) = \\
 &= -\left(\frac{1}{\xi^2 \omega_n^2 + \omega_p^2} \right) \frac{f}{m \omega_p} e^{\xi \omega_n t} (\xi \omega_n \sin(\omega_p t) + \omega_p \cos(\omega_p t)) \Big|_0^t = \\
 &= -\left(\frac{1}{\omega_n (\xi^2 \omega_n - \sqrt{1 - \xi^2})} \right) \frac{f}{m} e^{\xi \omega_n t} \left(\frac{\xi}{\sqrt{1 - \xi^2}} \sin(\omega_p t) + \cos(\omega_p t) \right)
 \end{aligned}$$

Then,

$$x(t) = \frac{f}{k} \left(1 - e^{-\xi \omega_n t} \left(\frac{\xi}{\sqrt{1 - \xi^2}} \sin(\omega_p t) + \cos(\omega_p t) \right) \right)$$

Determine the response for $c=0$

$$\begin{aligned}
 x(t) &= \int_0^t \frac{f}{m \omega_p} \sin(\omega_p t) dt = -\frac{f}{m \omega_p} \frac{\omega_p \cos(\omega_p t)}{\omega_p^2} + \frac{f}{m \omega_p^2} = \frac{f}{m \omega_p^2} (1 - \cos(\omega_p t)) = \\
 &= \frac{f}{k} (1 - \cos(\omega_n t))
 \end{aligned}$$

$$\omega_p = \omega_n \quad \text{if } c = 0$$

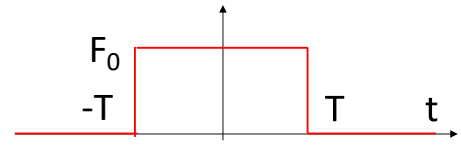
Which is the response at infinite?

$$X_{steadyState} = f/k$$

EXERCISE 4

Use the concept of unit step function and calculate the response of system in exercise 2 to the rectangular pulse shown in figure. Plot the first 8 peaks around $-T$ and T .

$$F_0 = 50 \text{ N}, \quad T = 5 \text{ s}$$



SOLUTION

The initial conditions of the Step input are:

$$x(0) = 0, \quad \dot{x}(0) = 0, \quad \omega_n = 7.0711 \text{ (rad/s)}$$

$$f_0 = 1$$

The response of the system to the excitation can be computed by adding the solution obtained for free oscillations to the steady-state response to the constant force f_0 .

$$x(t) = e^{-\sigma t} [X_{01} \cos(\omega_p t) + X_{02} \sin(\omega_p t)] + \frac{f_0}{k}$$

Imposing the initial conditions, it is possible to compute X_{01}, X_{02} from the previous equation:

$$X_{01} = -\frac{f_0}{k}$$

$$X_{02} = \frac{f_0}{k} \cdot \frac{\xi}{\sqrt{1-\xi^2}}$$

Thus,

$$x(t) = \frac{f_0}{k} g(t)$$

$$g(t) = 1 - e^{-\xi \omega_n t} \left[\cos\left((\omega_n \sqrt{1-\xi^2}) t\right) + \frac{\xi}{\sqrt{1-\xi^2}} \sin\left((\omega_n \sqrt{1-\xi^2}) t\right) \right]$$

If $C=0$ then $\xi = 0, X_{01} = -0.01, X_{02} = 0$

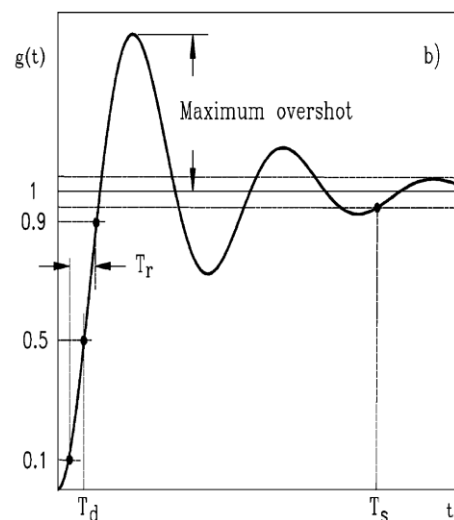
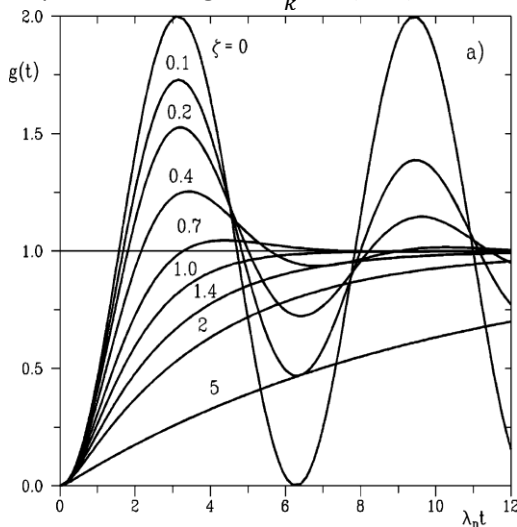
$$g(t) = 1 - [\cos(\omega_n t)]$$

If $C=2$ ($N/m.s$), $\xi = 0.0707, \omega_p = 7.0534$ (rad/s)

$$g(t) = 1 - e^{-\xi \omega_n t} \left[\cos(\omega_p t) + \frac{\xi}{\sqrt{1-\xi^2}} \sin(\omega_p t) \right]$$

$$\begin{cases} x(t) = 0 & , \quad t \leq 0 \\ x(t) = \frac{f_0}{k} g(t), & t > 0 \end{cases}$$

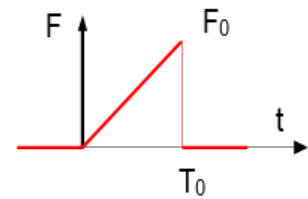
The infinite response converges to $\frac{f_0}{k} = 1$ (mm)



EXERCISE 5

Derive an expression for the response of the system in exercise 1 to a ramp force $F(t) = F_0 \frac{t}{T_0}$ in terms of the convolution integral. Consider both undamped and damped systems.

$$F_0 = 10 \text{ N/s}, T_0 = 5 \text{ ms}$$

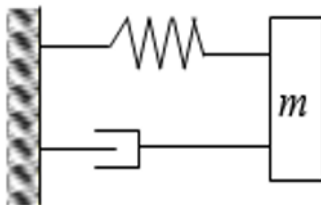


SOLUTION (complete solution in the dedicated pdf)

The force is given by:

$$F(t) = \begin{cases} \frac{F_0}{T_0} t & 0 \leq t \leq T_0 \\ F_0 & t > T_0 \end{cases}$$

- Damped system: $x(t) = A(t) \sin(\omega_p t) + B(t) \cos(\omega_p t)$



$$\begin{aligned} \omega_n &= \sqrt{\frac{K}{m}} = 5\sqrt{2} \text{ rad/s} \\ \omega_p &= 6.816 \text{ rad/s} \\ \xi &= \frac{\sqrt{2}}{20} \end{aligned}$$

$$A(t) = \frac{1}{m\omega_p} e^{\theta t} \int_0^t F(\tau) e^{-\theta \tau} \cos(\omega_p \tau) d\tau = \frac{1}{m\omega_p} e^{\theta t} \frac{F_0}{\Delta t} \int_0^t \tau e^{-\theta \tau} \cos(\omega_p \tau) d\tau$$

$$B(t) = \frac{1}{m\omega_p} e^{\theta t} \frac{F_0}{\Delta t} \int_0^t \tau e^{-\theta \tau} \sin(\omega_p \tau) d\tau$$

Undamped system

By introducing the impulse response of an undamped system, the particular integral of the equation of motion can be expressed as

$$x(t) = A(t) \sin(\omega_n t) - B(t) \cos(\omega_n t)$$

Starting from: $A(t) = \frac{1}{m\omega_n} F_0 \int_0^t \tau \cos(\omega_n \tau) d\tau$

Integrating by parts,

$$\int a b' = ab - \int a' b \rightarrow a = t, b' = \cos(\omega_n t):$$

$$\begin{aligned} \int_0^t \tau \cos(\omega_n \tau) d\tau &= \frac{\tau}{\omega_n} \sin(\omega_n \tau) - \frac{1}{\omega_n} \int_0^t \sin(\omega_n \tau) d\tau = t \frac{1}{\omega_n} \sin(\omega_n t) + \frac{1}{\omega_n^2} \cos(\omega_n t) \Big|_0^t = \\ &= \frac{1}{\omega_n^2} (\omega_n \sin(\omega_n t) t + \cos(\omega_n t) - 1) \end{aligned}$$

So, $A(t) = \frac{1}{\sqrt{2} 50} (5\sqrt{2} \sin(5\sqrt{2} t) t + \cos(5\sqrt{2} t) - 1)$

The term B(t) is computed:

$$\begin{aligned}
 B(t) &= \frac{1}{m\omega_p} F_0 \int_0^t \tau \sin(\omega_n \tau) d\tau = \int_0^t \tau \sin(\omega_n \tau) d\tau = -\frac{\tau}{\omega_n} \cos(\omega_n \tau) \Big|_0^t + \frac{1}{\omega_n} \int_0^t \cos(\omega_n \tau) d\tau \\
 &= -t \frac{1}{\omega_n} \sin(\omega_n t) + \frac{1}{\omega_n} + \frac{1}{\omega_n^2} \sin(\omega_n t) = \frac{1}{\omega_n^2} (1 + \sin(\omega_n t) - \omega_n t \cos(\omega_n t)) \\
 B(t) &= \frac{1}{\sqrt{2} 50} (1 + \sin(5\sqrt{2} t) - 5\sqrt{2} t \cos(5\sqrt{2} t))
 \end{aligned}$$

So,

$$\begin{aligned}
 x(t) &= \frac{1}{\sqrt{2} 50} (5\sqrt{2} \sin^2(5\sqrt{2} t) t + \sin(5\sqrt{2} t) \cos(5\sqrt{2} t) - \sin(5\sqrt{2} t)) \\
 &\quad - \frac{1}{\sqrt{2} 50} (\cos(5\sqrt{2} t) + \cos(5\sqrt{2} t) \sin(5\sqrt{2} t) - 5\sqrt{2} t \cos^2(5\sqrt{2} t)) \\
 &= \frac{1}{\sqrt{2} 50} (5\sqrt{2} t - \sin(5\sqrt{2} t))
 \end{aligned}$$

Directly applying the Duhamel integral:

$$\begin{aligned}
 x(t) &= \int_0^t F_0(t-\tau) \frac{1}{\omega_n m} \sin(\omega_n \tau) d\tau = \frac{1}{\omega_n m} F_0 t \int_0^t \sin(\omega_n \tau) d\tau - \frac{1}{\omega_n m} F_0 \int_0^t \tau \sin(\omega_n \tau) d\tau = \\
 &= \frac{1}{\omega_n m} F_0 t \left(-\frac{1}{\omega_n} \cos(\omega_n t) + \frac{1}{\omega_n} \right) - \frac{1}{\omega_n m} F_0 \int_0^t \tau \sin(\omega_n \tau) d\tau \quad (1)
 \end{aligned}$$

Solving the integration by parts of the term $\int_0^t \tau \sin(\omega_n \tau) d\tau$

$$\int a b' = ab - \int a' b \rightarrow a = \tau \quad b' = \sin(\omega_n \tau)$$

$$\rightarrow -\frac{t}{\omega_n} (\cos(\omega_n t)) - \frac{1}{\omega_n} \int_0^t \cos(\omega_n \tau) d\tau = -\frac{t}{\omega_n} (\cos(\omega_n t)) - \frac{1}{\omega_n^2} \sin(\omega_n t) \quad (2)$$

Thus, substituting (2) in equation (1)

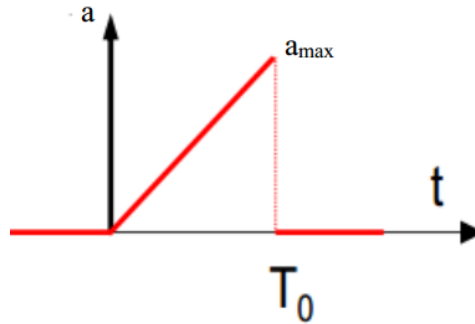
$$\begin{aligned}
 x(t) &= \frac{1}{\omega_n m} F_0 \left[t \left(\frac{1}{\omega_n} - \frac{1}{\omega_n} \cos(\omega_n t) \right) + \frac{t}{\omega_n} \cos(\omega_n t) - \frac{1}{\omega_n^2} \sin(\omega_n t) \right] \\
 &= \frac{1}{\omega_n m} F_0 \left[\frac{t}{\omega_n} - \frac{1}{\omega_n^2} \sin(\omega_n t) \right] = \frac{1}{\omega_n^3 m} F_0 [\omega_n t - \sin(\omega_n t)]
 \end{aligned}$$

EXERCISE 6

A single degree of freedom dynamic system ($k = 2000 \text{ N/m}$, $m = 0.3 \text{ kg}$, neglect the damping) is shocked by an acceleration having peak 4000 m/s^2 and duration 0.5 ms . Determine the equation of motion by using both the impulse response method and the Duhamel integral (assume a linear acceleration profile). Plot the displacements versus time and compare the results (at least in term of maximum amplitude).

Repeat the calculation when the same shock lasts for 20 ms . Explain the difference with the previous case.

SOLUTION



1) Equation of motion using the impulse response method

The initial conditions are $x(0) = 0$, $\dot{x}(0) = \frac{f_0}{m}$

$$x(t) = e^{-\xi\omega_n t} \frac{1}{\omega_p} \cdot \frac{f_0}{m} \sin(\omega_p t) \quad , C = 0 \text{ so } \omega_p = \omega_n = 25.82 \text{ (rad/s)}$$

$$f_0 = \frac{1}{2} a_m * m * \Delta t$$

$$\Delta t = 0.5 \text{ (ms)} \text{ so } f_0 = 0.3 \text{ (N.s)}$$

$$x(t) = 0.387 \sin(\omega_n t) \text{ and } x_{max}(t) = 0.387 \text{ (m)}$$

$$\Delta t = 20 \text{ (ms)} \text{ so } f_0 = 12 \text{ (N.s)}$$

$$x(t) = 0.0387 \sin(\omega_n t) \text{ and } x_{max}(t) = 0.0387 \text{ (m)}$$

2) Equation of motion using the Duhamel integral

The particular integral of the equation of motion can be expressed as

$$x(t) = A_1(t) \sin(\omega_p t) - B_1(t) \cos(\omega_p t)$$

where

$$A_1(t) = \frac{1}{m\omega_p e^{-\xi\omega_n t}} \int_0^{T_0} F(\tau) e^{\xi\omega_n \tau} \cos(\omega_p \tau) d\tau$$

$$B_1(t) = \frac{1}{m\omega_p e^{-\xi\omega_n t}} \int_0^{T_0} F(\tau) e^{\xi\omega_n \tau} \sin(\omega_p \tau) d\tau$$

$$F(\tau) = \frac{f_0}{T_0} \tau$$

$$\text{When } t > T_0 \text{ so } A_1(t) = A_1(T_0) \text{ and } B_1(t) = B_1(T_0)$$

Consider that there is no damping and so:

$$A(t) = \frac{1}{m\omega_p} F_0 \int_0^t \tau \cos(\omega_n \tau) d\tau$$

Integrating by parts,

$$\int a b' = ab - \int a' b \rightarrow a = t, b' = \cos(\omega_n t):$$

$$\int_0^t \tau \cos(\omega_n \tau) d\tau = \frac{\tau}{\omega_n} \sin(\omega_n \tau) - \frac{1}{\omega_n} \int_0^t \sin(\omega_n \tau) d\tau = t \frac{1}{\omega_n} \sin(\omega_n t) + \frac{1}{\omega_n^2} \cos(\omega_n t) \Big|_0^t =$$

$$A = \frac{1}{\omega_n^2} (\omega_n \sin(\omega_n t) t + \cos(\omega_n t) - 1)$$

The term B(t) is computed as follows:

$$\begin{aligned} B(t) &= \frac{1}{m\omega_p} F_0 \int_0^t \tau \sin(\omega_n \tau) d\tau = \int_0^t \tau \sin(\omega_n \tau) d\tau = -\frac{\tau}{\omega_n} \cos(\omega_n \tau) \Big|_0^t + \frac{1}{\omega_n} \int_0^t \cos(\omega_n \tau) d\tau \\ &= -t \frac{1}{\omega_n} \sin(\omega_n t) + \frac{1}{\omega_n} + \frac{1}{\omega_n^2} \sin(\omega_n t) = \frac{1}{\omega_n^2} (1 + \sin(\omega_n t) - \omega_n t \cos(\omega_n t)) \end{aligned}$$

MULTIPLE CHOICE PROBLEMS: select the right answer. Explain in few sentences (eventually with the support of graphs, schemes and formulas) the reason of the correct answer.

- 1) The hysteretic damping model is applicable to
[A] linear models subject to any type of excitation
[B] only linear models subject to harmonic excitation
[C] linear and non-linear models subject to harmonic or multi harmonic excitation

SOLUTION [C]

- 2) The Duhamel integral can be applied only
[A] For linear systems,
[B] with shock inputs
[C] if the system is undamped.

SOLUTION [A]

- 3) At resonance
[A] elastic forces balance exactly inertia forces;
[B] damping forces balance exactly inertia forces;
[C] elastic forces balance exactly damping forces.

SOLUTION [A]

- 4) The dynamic compliance of a system with viscous damping
[A] is expressed by a complex number;
[B] tends to zero when the forcing function tends to zero;
[C] is always expressed by a real number.

SOLUTION [A]

- 5) The response of an undamped linear system at its resonant frequency
[A] is infinitely large;
[B] grows linearly in time to infinity;
[C] grows exponentially in time to infinity.

SOLUTION [B]