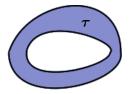


THE INITIAL STRESS SYMMETRY

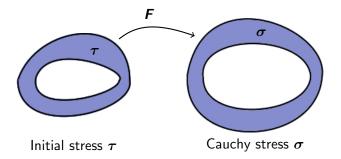
Author: Artur L. Gower Co-Authors:
Dr. Pasquale Ciarletta
Prof. Michel Destrade

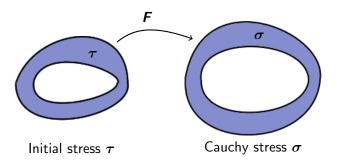




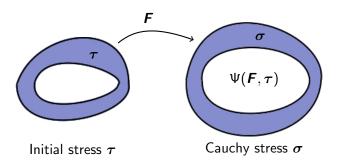


Initial stress au

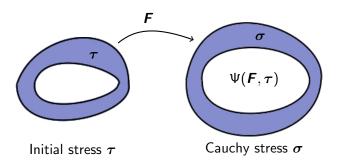




"To characterize the arterial wall or any other biological soft tissue, we need a stress-free state." – Chuong and Fung (1986)



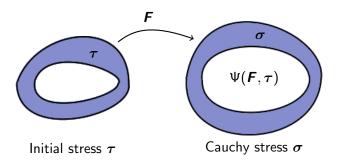
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Something like:
$$\Psi(\mathbf{F}, \boldsymbol{\tau}) = C_1 \operatorname{tr}(\mathbf{F}^T \mathbf{F}) + C_2 \operatorname{tr} \boldsymbol{\tau}$$
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Something like:
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. (Great for in-vivo characterization!)

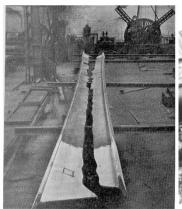
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Exploding residual stress

"Residual-stresses arising from metal forming and machining and the grounds for their induction have challenged the minds of engineers and scientists since the Industrial Revolution." – Upshaw et al. (2011)

Exploding residual stress

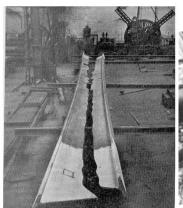
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Typically measuring residual stress in metal matrix composites involves cutting the component in half or drilling small holes.

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With no other intrinsic anisotropy other than the initial stress there are 10 invariants (Shams et al., 2011):

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$$\begin{split} I_1 &= \operatorname{tr} {\pmb C}, \quad I_2 = \frac{1}{2} [(I_1^2 - \operatorname{tr} ({\pmb C}^2)], \quad I_3 = \det {\pmb C}, \\ I_{\tau_1} &= \operatorname{tr} {\pmb \tau}, \quad I_{\tau_2} = \frac{1}{2} [(I_{\tau_1}^2 - \operatorname{tr} ({\pmb \tau}^2)], \quad I_{\tau_3} = \det {\pmb \tau}, \\ J_1 &= \operatorname{tr} ({\pmb \tau} {\pmb C}), \quad J_2 = \operatorname{tr} ({\pmb \tau} {\pmb C}^2), \quad J_3 = \operatorname{tr} ({\pmb \tau}^2 {\pmb C}), \quad J_4 = \operatorname{tr} ({\pmb \tau}^2 {\pmb C}^2). \end{split}$$

where $\mathbf{C} = \mathbf{F}^T \mathbf{F}$.

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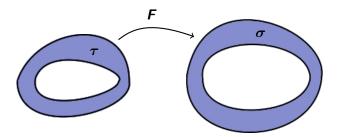
$$egin{aligned} I_1 &= \operatorname{tr} m{C}, \quad I_2 = rac{1}{2}[(I_1^2 - \operatorname{tr}(m{C}^2)], \quad I_3 = \det m{C}, \ I_{ au_1} &= \operatorname{tr} m{ au}, \quad I_{ au_2} = rac{1}{2}[(I_{ au_1}^2 - \operatorname{tr}(m{ au}^2)], \quad I_{ au_3} = \det m{ au}, \ J_1 &= \operatorname{tr}(m{ au}m{C}), \quad J_2 = \operatorname{tr}(m{ au}m{C}^2), \quad J_3 = \operatorname{tr}(m{ au}^2m{C}), \quad J_4 = \operatorname{tr}(m{ au}^2m{C}^2). \end{aligned}$$

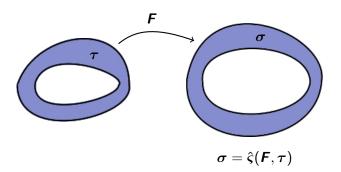
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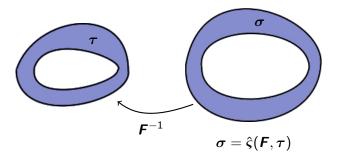
Shall we just drop a few invariants?

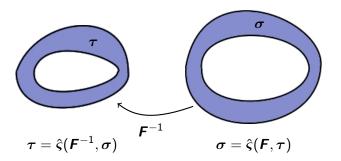


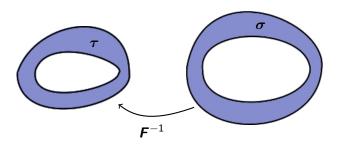
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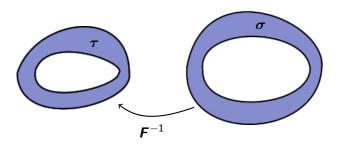




Initial Stress Symmetry (ISS):

$$\sigma = \hat{\varsigma}(F, \tau)$$
 and $\tau = \hat{\varsigma}(F^{-1}, \sigma)$ for every F and τ .

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Turns out that ISS is automatically satisfied if τ is due to the deformation of a virtual stress-free configuration!

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Elasticity

For a free energy density $\Psi(\mathbf{F}, \tau)$, ISS can be written as 9 scalar equations involving the invariants of \mathbf{F} and τ .

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For a free energy density $\Psi(\mathbf{F}, \tau)$, ISS can be written as 9 scalar equations involving the invariants of \mathbf{F} and τ .

For example, if $\Psi({\pmb F}, {\pmb au})$ is independent of ${\rm tr}({\pmb au}^2 {\pmb C})$ and ${\rm tr}({\pmb au}^2 {\pmb C}^2)$ then

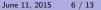
$$4\Psi_{J_1}\Psi^{\sigma}_{J_1}=1, \ \frac{\Psi^{\sigma}_{I_2}}{\sqrt{\Psi^{\sigma}_{J_1}}}+\frac{\Psi_{I_2}}{\sqrt{\Psi_{J_1}}}=0, \ \ p\Psi^{\sigma}_{J_1}=\Psi^{\sigma}_{I_1}, \ \ p_{\tau}\Psi_{J_1}=\Psi_{I_1},$$

where

$$\Psi_{I_k} = \Psi_{I_k}(\mathbf{F}, \tau), \quad \Psi_{J_1} = \Psi_{J_1}(\mathbf{F}, \tau), \tag{1}$$

$$\Psi_{I_k}^{\sigma} := \Psi_{I_k}(\mathbf{F}^{-1}, \boldsymbol{\sigma}), \ \Psi_{J_m}^{\sigma} := \Psi_{J_m}(\mathbf{F}^{-1}, \boldsymbol{\sigma}), \tag{2}$$

p and p_{τ} are Lagrange multipliers due to incompressibility.



Examples of $\Psi(\boldsymbol{F}, \boldsymbol{\tau})$

Examples that do not satisfy ISS:

$$\Psi = \frac{1}{2}\mu(I_1 - 3) + \frac{1}{2}(J_1 - I_{\tau 1}),$$

$$\Psi = \frac{1}{2}\mu(I_1 - 3) + \frac{1}{4}(J_2 - I_{\tau 1}),$$

$$\Psi = \frac{1}{2}\mu(I_1 - 3) + \frac{1}{2}\overline{\mu}(J_1 - I_{\tau 1})^2 + \frac{1}{2}(J_1 - I_{\tau 1}).$$

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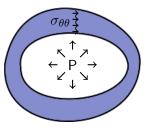
Examples that **do** satisfy ISS:

$$\Psi=rac{1}{2}\left(p_{ au}(au,1)I_1+J_1-3\mu
ight) ext{ (Incompressible neo-Hookean)},
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ight) ext{ (Compressible)}.$$

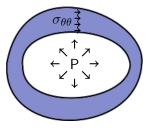


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Suppose you want a homogeanious circumferential Cauchy stress:

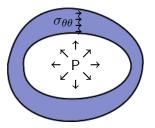


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ISS in the form $\tau = \hat{\varsigma}(\mathbf{F}^{-1}, \sigma)$ suggests that for any choice of σ , there will exist an initial stress τ that supports σ .

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ISS in the form $\tau = \hat{\varsigma}(\mathbf{F}^{-1}, \sigma)$ suggests that for any choice of σ , there will exist an initial stress τ that supports σ .

So first solve for your ideal σ then find τ !



So we can drop hyperelasticity and minimize

$$\frac{1}{b-a}\int_a^b \|\nabla \boldsymbol{\sigma}\|^2 dr,$$

with inner and outer radius a and b.

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For radial symmetry the solution is

$$\sigma_{rr} = -P(1-\rho)(1-\rho\beta+\rho^{2}\beta^{2}) + O(\beta^{3}) P,$$

$$\sigma_{\theta\theta} = P\beta^{-1}(1+\beta^{3}\rho^{2}(4\rho-3)) + O(\beta^{3}) P,$$

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 $\sigma_{zz} = 0$, where

$$\rho = \frac{r-a}{b-a}$$
 and $\beta = \frac{b-a}{a}$ (≈ 0.07).

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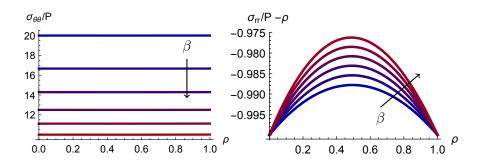


Figure: the ideal adimensional Cauchy stress against the adimensional radius ρ , where β goes from 0.05 to 0.1.

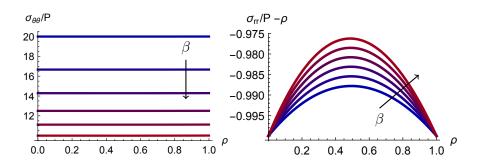


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To make predictions we need to choose $\Psi(\mathbf{F}, \tau)$,

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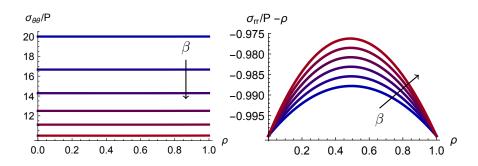


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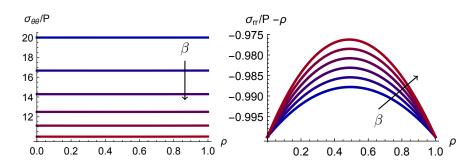


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and unloaded boundary conditions. Now on to predictions!

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Predictions for ideal σ

For the neo-Hookean $\Psi(\mathbf{F}, \boldsymbol{\tau})$.

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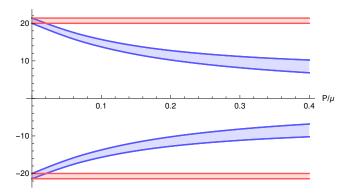


Figure: The loaded aorta wall is red and the unloaded aorta wall is blue. The parameters are for the descending thoracic aorta: loaded inner radius = 20mm, outer radius = 21.4mm and $\beta = 0.07$.

Predictions for ideal σ

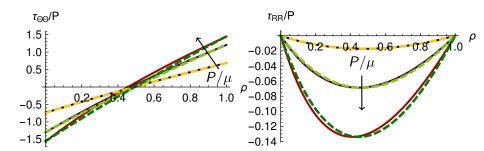


Figure: the adimensional residual stress against the adimensional radius ρ . The dashed curves are from the opening angle method while the solid curves are from the ideal stress σ , both with $P/\mu=0.05,\,0.2$ or 0.35.

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■ Initial Stress Symmetry:

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See more details in the arXiv entitled:

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Thanks for listening! Any questions?



- Chuong, C. and Y. Fung (1986). "Residual stress in arteries". In: *Frontiers in Biomechanics*. Springer, pp. 117–129.
- Shams, M., M. Destrade, and R. W. Ogden (2011). "Initial stresses in elastic solids: Constitutive laws and acoustoelasticity". In: *Wave Motion* 48, pp. 552–567.
- Upshaw, D., M. Steinzig, and J. Rasty (2011). "Influence of drilling parameters on the accuracy of hole-drilling residual stress measurements". In: *Engineering Applications of Residual Stress, Volume 8.* Springer, pp. 95–109.