

$\Rightarrow$  Two isotropic point source of same Amplitude

and Phase. that means two isotropic

Point sources ① and ② having Current Supplied.  
with equal amplitude and same phase.

$\Rightarrow$  due to current supplied to both is equal,  
electric field magnitude generated by both  
elements will be equal.

$$E_1 = E_2 = E_0 \quad \text{current is same.}$$

$\Rightarrow$  Path difference =  $d \cos \theta$

Total Phase difference =  $(2\pi/\lambda) d \cos \theta$

$$\beta = 2\pi/\lambda \Rightarrow \beta d \cos \theta$$

Generally  $\beta d \cos \theta + \alpha$

$\Downarrow$  associated same phase  
 $\therefore \alpha = 0$

Electric field by source 1 =  $E_1 = E_0 e^{-j\psi_{1/2}}$

Hence field of Point source 1 lags by  $\psi_{1/2}$  Phase.

Electric field by source 2 =  $E_2 = E_0 e^{j\psi_{1/2}}$

Hence E-field of Point source 2 lead by  $\psi_{1/2}$  Phase.

⇒ Total Electric field

$$E = E_1 + E_2$$

$$= E_0 e^{-j\psi_{1/2}} + E_0 e^{j\psi_{1/2}}$$

$$= E_0 \left( e^{-j\psi_{1/2}} + e^{j\psi_{1/2}} \right)$$

$$= 2E_0 \left( \frac{e^{-j\psi_{1/2}} + e^{j\psi_{1/2}}}{2} \right)$$

$$\boxed{\frac{e^{j\theta} + e^{-j\theta}}{2} = \cos\theta}$$

$$E = 2E_0 \cos(\psi_{1/2}) \quad \textcircled{1}$$

$E_n$  is normalized Electric field.

$$E_n = 2E_0$$

then  $\textcircled{1}$  become  $E = E_n \cos(\psi_{1/2}) \quad \textcircled{2}$

We know  $\psi = \beta d \cos\theta + \alpha$

$$E = E_n \cos\left(\frac{\beta d \cos \theta + \alpha}{2}\right)$$

This is equation of farfield Pattern  
of two Isotropic point sources of  
Same amplitude and Phase

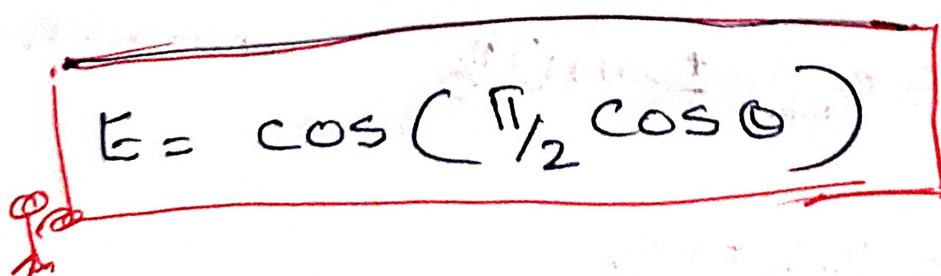
putting  $E_n = 1$  (Pattern is said to be normalized)

Here  $\alpha = 0^\circ$ ,

$$E = \cos\left(\frac{\beta d \cos \theta}{2}\right)$$

$$\text{if } d = \frac{\lambda}{2}, \beta = \frac{2\pi}{\lambda}$$

$$E = \cos\left(\frac{\frac{2\pi}{\lambda} \times \frac{\lambda}{2} \times \cos \theta}{2}\right)$$



$$E = \cos\left(\frac{\pi}{2} \cos \theta\right)$$

$$\boxed{E = \cos(\frac{\pi}{2} \cos\theta)}$$

$\cos(\frac{\pi}{2})$  (broadside) - two isotropic Point Source radiator.

### Maxima direction

$$\cos(\frac{\pi}{2} \cos\theta) = \pm 1$$

$$\frac{\pi}{2} \cos\theta_{\text{max}} = \pm n\pi, \quad \text{where } n = 0, 1, 2, \dots$$

$$\frac{\pi}{2} \cos\theta_{\text{max}} = 0$$

$$\cos\theta_{\text{max}} = 0$$

$$\boxed{\theta_{\text{max}} = 90^\circ \text{ and } 270^\circ}$$

### minima Direction

$$\cos(\frac{\pi}{2} \cos\theta) = 0$$

$$\frac{\pi}{2} \cos\theta_{\text{min}} = \pm(2n+1)\frac{\pi}{2}$$

where  $n = 0, 1, 2, \dots$

$$\frac{\pi}{2} \cos\theta_{\text{min}} = \pm\frac{\pi}{2}$$

$$\cos\theta_{\text{min}} = \pm 1$$

$$\boxed{\theta_{\text{min}} = 0^\circ \text{ and } 180^\circ}$$

## Half Power Point direction

[HPPD]

Power is  $\frac{1}{2}$  or Voltage or current is  $\frac{1}{\sqrt{2}}$  times of maximum value of Voltage or Current.

$\left[ \frac{P_{\text{max}}}{2} \text{ or } \frac{V_{\text{max}} \text{ or } I_{\text{max}}}{\sqrt{2}} \right] \text{ called Half Power Point}$

$$\cos(\frac{\pi}{2} \cos \theta) = \pm \frac{1}{\sqrt{2}}$$

$$\frac{\pi}{2} \cos \theta_{\text{HPPD}} = \pm (2n+1) \frac{\pi}{4}; \quad \text{where } n=0, 1, 2, \dots$$

$$\frac{\pi}{2} \cos \theta_{\text{HPPD}} = \pm \frac{\pi}{4}$$

$$\cos \theta_{\text{HPPD}} = \pm \frac{1}{2}$$

$$\textcircled{1} \quad \theta_{\text{HPPD}} = \pm 60^\circ, \pm 120^\circ$$

## Field Pattern [E vs θ]

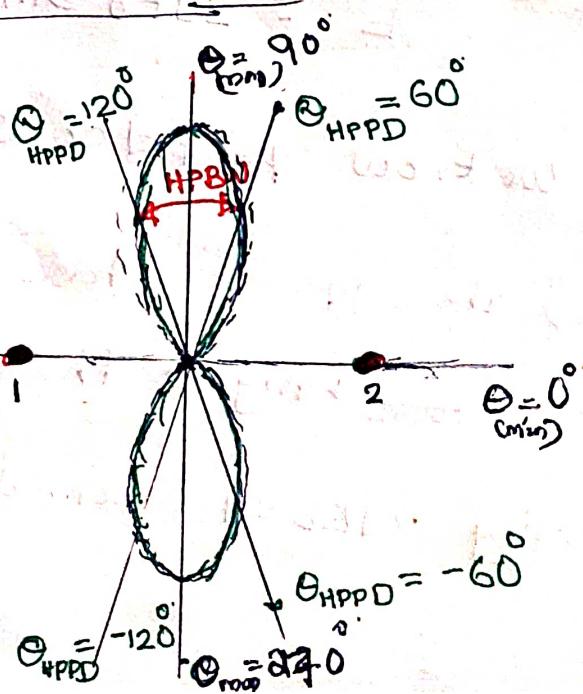
\* This is Simplest type of broad Side array

\* also known as

broadside couplet

as two isotropic antennas are in Phase.

\*  $\theta_{\text{HPPD}} = \pm 90^\circ$  &  $\alpha = 0^\circ$



case (ii) : Array of Point Source with  
Equal amplitude and opposite Phase (end fire)

\* (Sometime)

$$\text{Source 1} \Rightarrow E_1 = -E_0 e^{-j\psi/2}$$

$$\text{Source 2} \Rightarrow E_2 = E_0 e^{j\psi/2}$$

Total Electric field (farfield E)

$$E = E_1 + E_2$$

$$= -E_0 e^{-j\psi/2} + E_0 e^{j\psi/2}$$

$$E = E_0 (e^{j\psi/2} - e^{-j\psi/2})$$

$$E = 2jE_0 \sin(\psi/2)$$

$$\begin{aligned} & \frac{e^{j\theta} - e^{-j\theta}}{2j} = \cos\theta \\ & e^{j\theta} - e^{-j\theta} = 2j \sin\theta \end{aligned}$$

We know  $\psi = Bd \cos\theta$

term  $j \Rightarrow$  the opposite Phase fed at source 1 and 2,  
 cause bring a Phase shift of  $90^\circ$  in the total field

putting  $2jE_0 = 1$  (normalized).

Let  $d = \frac{\lambda}{2}$  &  $2jE_0 = 1$ ,  $\beta = \frac{2\pi}{\lambda}$

$$E = \sin\left(\frac{2\pi}{\lambda} \times \frac{\lambda}{2} \frac{\cos\theta}{2}\right)$$

$$E = \sin\left(\frac{\pi}{2} \cos\theta\right)$$

Cylinder (end fire) two isotropic point source radiators.

Maxima direction

maximum value of Sine function is  $\pm 1$

$$\sin\left(\frac{\pi}{2} \cos\theta\right) = \pm 1$$

$$\frac{\pi}{2} \cos\theta_{\text{max}} = \pm (2n+1) \frac{\pi}{2}$$

$$\frac{\pi}{2} \cos\theta_{\text{max}} = \pm \frac{\pi}{2}$$

$$\cos\theta_{\text{max}} = \pm 1$$

$$\theta_{\text{max}} = 0^\circ \text{ and } 180^\circ$$

## minima direction

$$\sin(\frac{\pi}{2} \cos\theta) = 0$$

$$\frac{\pi}{2} \cos\theta_{\min} = \pm n\pi \quad \text{where } n=0, 1, 2, \dots$$

$$\cos\theta_{\min} = 0$$

$$\theta_{\min} = 90^\circ \text{ and } 270^\circ$$

## Half Power point directions

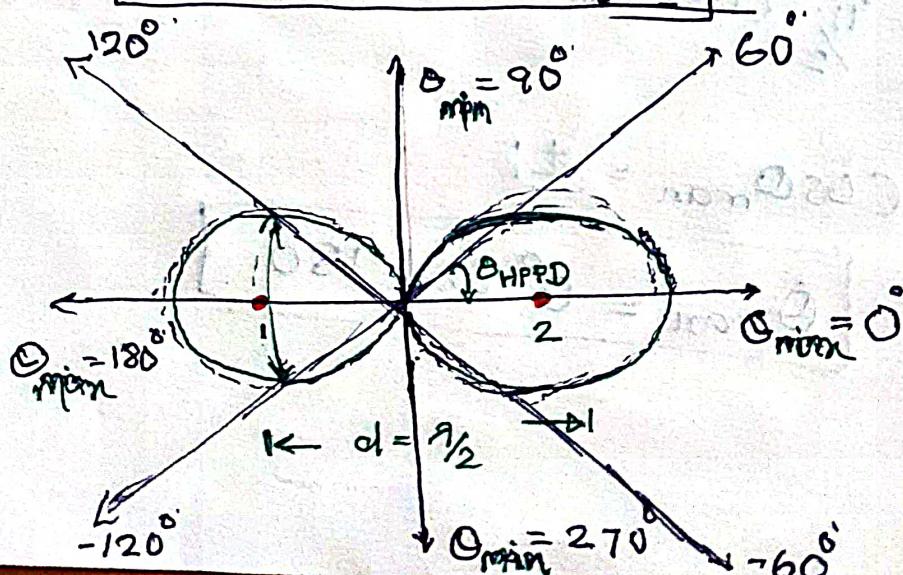
$$\sin(\frac{\pi}{2} \cos\theta) = \pm \frac{1}{\sqrt{2}}$$

$$\frac{\pi}{2} \cos\theta_{HPPD} = \pm (2n+1) \frac{\pi}{4}$$

$$\frac{\pi}{2} \cos\theta_{HPPD} = \pm \frac{\pi}{4}$$

$$\cos\theta_{HPPD} = \pm \frac{1}{2}$$

$$\theta_{HPPD} = \pm 60^\circ, \pm 120^\circ$$



Case I

## Radiation Pattern of two isotropic Point

Source Radiator Separated by Wavelength ( $\lambda$ ) $[d = \lambda]$ 

Electric field by two elements

$$E = E_n \cos(\frac{\psi}{2})$$

$$= E_n \cos\left(\frac{\beta d \cos \theta + \alpha}{2}\right)$$

$$\beta = \frac{2\pi}{\lambda}, \alpha = \lambda, \alpha = 0$$

$$E = E_n \cos\left(\frac{\frac{2\pi}{\lambda} \times \lambda \cos \theta + 0}{2}\right)$$

$$E = E_n \cos(\pi \cos \theta)$$

For Maxima

$$\cos(\pi \cos \theta) = \pm 1$$

$$\pi \cos \theta_{\text{max}} = \pm n\pi$$

$$\cos \theta_{\text{max}} = \pm n, n = 0, 1, 2, \dots$$

 $\Rightarrow$  for  $n=0$ 

$$\cos \theta_{\text{max}} = 0$$

$$\theta_{\text{max}} = 90^\circ, 270^\circ$$

 $\Rightarrow$  for  $n=1$ 

$$\cos \theta_{\text{max}} = \pm 1$$

$$\theta_{\text{max}} = 0^\circ, 180^\circ$$

$$\boxed{\theta_{\text{max}} = 0^\circ, 90^\circ, 180^\circ, 270^\circ}$$

→ for minima (Nulls)

$$\cos(\pi \cos \theta) = 0$$

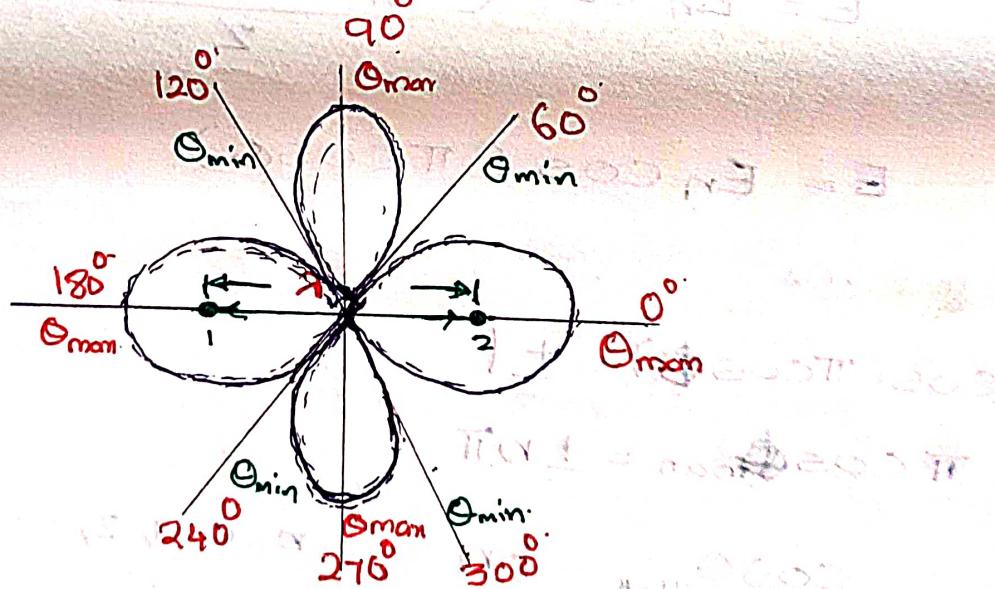
$$\pi \cos \theta_{\min} = \pm (2n+1)\frac{\pi}{2}, \quad n=0, 1, 2, \dots$$

chances exist at first position

⇒ for  $n=0$ :

$$\cos \theta_{\min} = \pm \frac{1}{2}$$

$$\theta_{\min} = 60^\circ, 120^\circ, 240^\circ, 300^\circ$$



This is farfield - E - pattern at  $d=2$

case II  $d = 2\lambda, \alpha = 0$

$$E = E_0 \cos(2\pi \cos \theta)$$

$\Rightarrow \theta_{\text{min}}$  or Nulls at:  $48^\circ, 76^\circ$

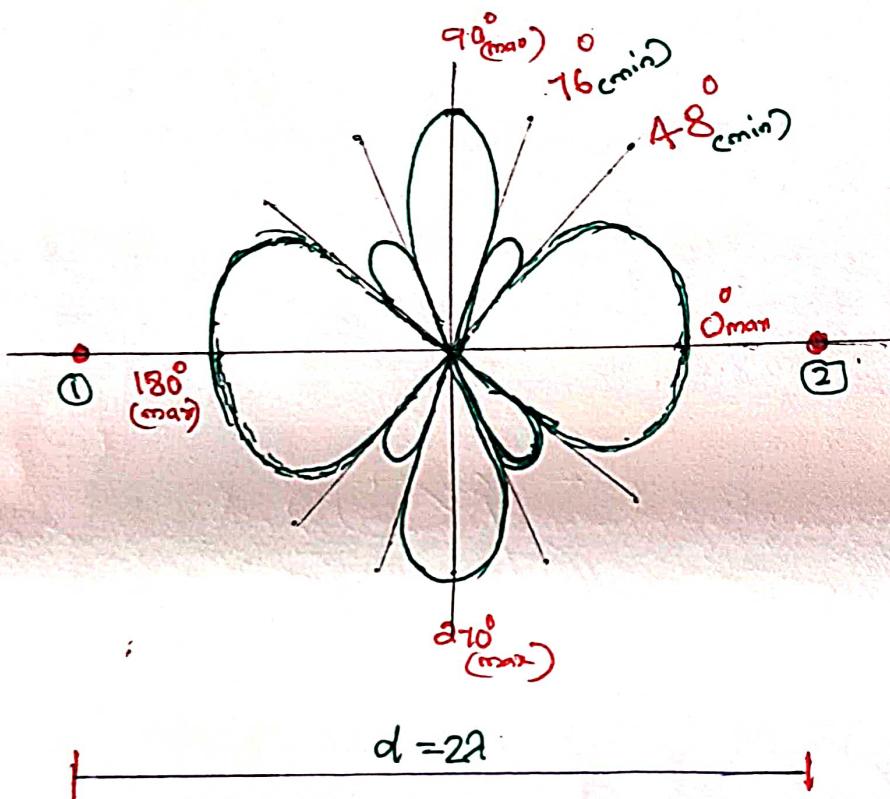


Fig. Radiation pattern of isotropic radiator spaced  $2\lambda$ .

## Pattern multiplication

Principle of Pattern Multiplication States that the radiation Pattern of an array is the Product of

Pattern of the Individual Pattern with array Pattern.

The array pattern is the function of the location of the antennas in the array and their relative complex enciation of amplitude.

### Advantage

- \* It helps to sketch the radiation pattern of array antennas rapidly from the simple product of element pattern and array pattern.

Generalised Expression of Principle of Pattern multiplication

$$E = \{ E_i(\theta, \phi) \times E_a(\theta, \phi) \} \times \{ E_{pi}(\theta, \phi) + E_{pa}^{(0, \phi)} \}$$

E = (multiplication of field pattern)  $\times$  (addition of Phase Pattern)

$E \rightarrow$  Total field

$E_i(\theta, \phi) \Rightarrow$  field Pattern of individual source.

$E_a(\theta, \phi) \Rightarrow$  field Pattern of array of isotropic point source [array factor of Isotropic element]

$E_{pi}(\theta, \phi) \Rightarrow$  Phase pattern of individual source

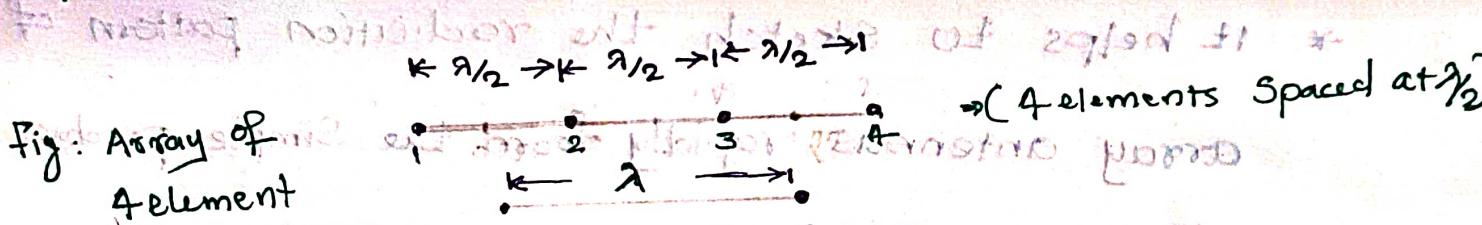
$E_{pa}(\theta, \phi) \Rightarrow$  Phase pattern of array of isotropic Point source.

### Disadvantage

- \* The Principle is applicable only for arrays containing identical elements

### Example

- ① Radiation Pattern of 4-isotropic element fed in Phase, spaced  $\lambda/2$  apart.



\* element 1 and 2 Considered as one unit ' $A$ '

\* Element 3 & 4 considered as one unit ' $B$ '

\* let  $A$  and  $B$  (are two isotropic antennas spaced at  $\lambda$ )

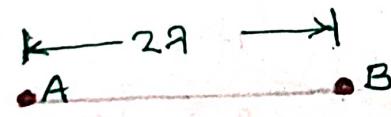
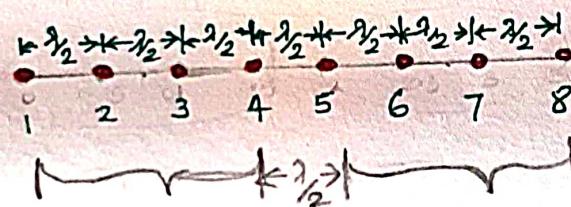
⇒ According to Pattern of 4 elements is obtained by multiplying the radiation pattern of individual element and array of two units spaced  $\lambda$ , cw shown in below-

Unit pattern  
( $d = \frac{\lambda}{2}$ )

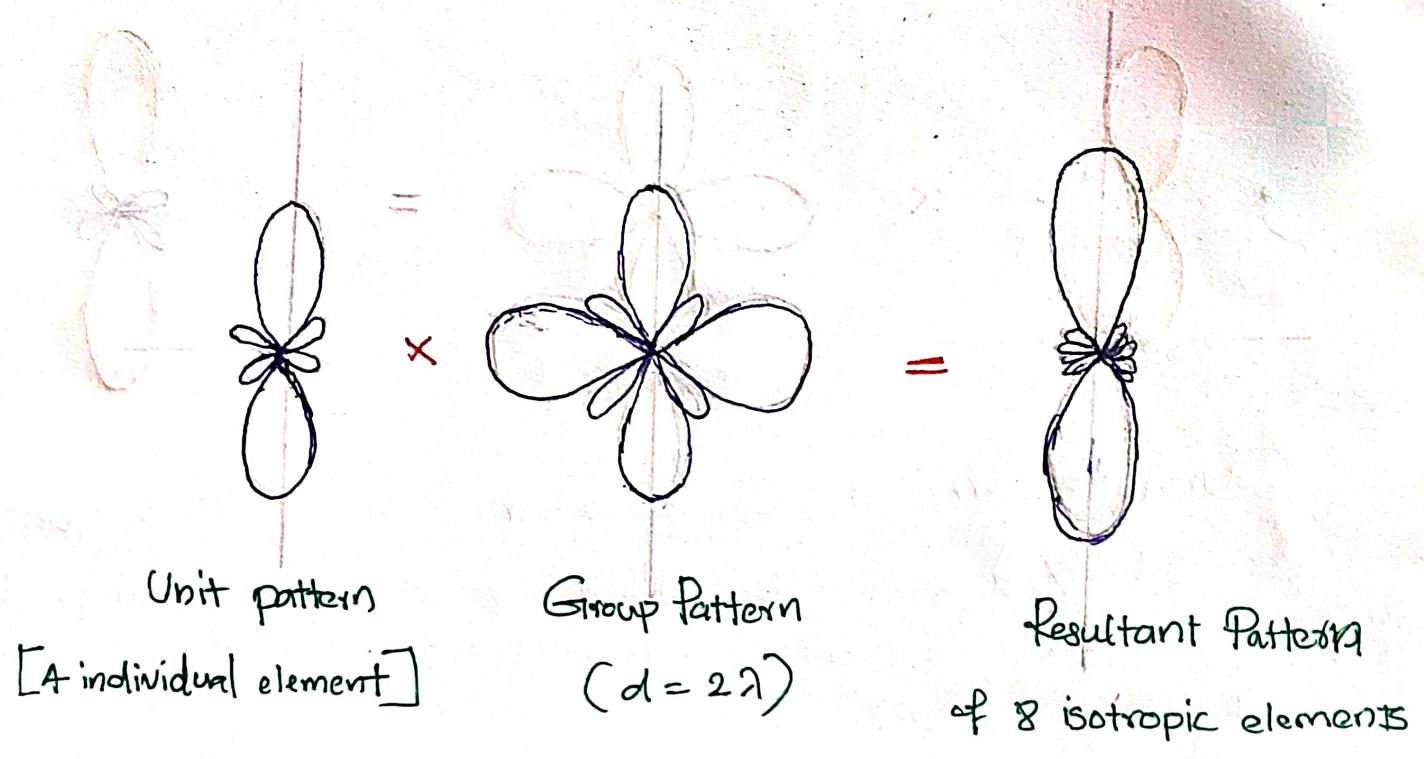
Group Pattern  
( $d = \lambda$ )

Resultant Pattern  
of 4 Isotropic elements

Example 2: Array of 8 elements spaced at  $\frac{\lambda}{2}$



\* The Radiation pattern of 8 isotropic elements is obtained by multiplying the Unit pattern of 4 individual elements and group pattern of two isotropic radiators spaced  $2\lambda$ , as shown below.



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## Uniform Linear array

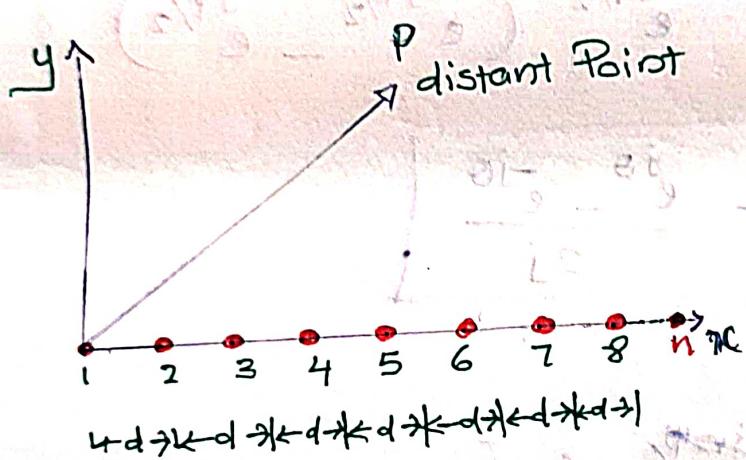
It is an array where the elements are spaced and excited equally along a straight line.

Linear: Equal Spacing  
 Uniform: Fed with current of equal amplitude and uniform progressive phase shift  
 Linear array with  $n$  isotropic point source of equal amplitude and spacing

⇒ These array are suitable for production of narrow radiation beam

⇒ These are required for point to point communication

⇒ They are also used in high angular resolution radars



Linear array with  $n$  isotropic point sources

with equal amplitude and spacing

→ Total far-field pattern at distant point  $P$ , given by

$$E_t = E_0 e^{0j\psi} + E_0 e^{1j\psi} + E_0 e^{2j\psi} + \dots + E_0 e^{j(n-1)\psi}$$

$$E_t = E_0 \left( e^{0j\psi} + e^{1j\psi} + e^{2j\psi} + \dots + e^{j(n-1)\psi} \right)$$

This is the form of G.P (Sum of the first  $n$  terms of G.P)  $\Rightarrow S_n = \frac{a(1-r^n)}{1-r}$   $r < 1$

$$\text{Here } \begin{cases} a = 1 \\ r = e^{-j\psi} \end{cases}$$

$$E_t = E_0 \frac{e^{jn\psi}}{(1 - e^{j\psi})}$$

$$E_t = E_0 \frac{(1 - e^{jn\psi/2} e^{jn\psi/2})}{(1 - e^{j\psi/2} e^{j\psi/2})}$$

$$E_t = E_0 \frac{e^{jn\psi/2} (e^{-jn\psi/2} - e^{+jn\psi/2})}{e^{j\psi/2} (e^{-j\psi/2} - e^{j\psi/2})} \quad \text{--- (1)}$$

We know

$$\sin \theta = \frac{e^{j\theta} - e^{-j\theta}}{2j}$$

~~$$E_t = E_0 e^{\frac{j(n-1)\psi}{2}} \left( \frac{e^{jn\psi/2} - e^{-jn\psi/2}}{2j} \right)$$~~

(1) becomes,

$$E_t = E_0 e^{\frac{j(n-1)\psi}{2}} \left\{ \frac{\left( \frac{e^{jn\psi/2} - e^{-jn\psi/2}}{2j} \right)}{\left( \frac{e^{j\psi/2} - e^{-j\psi/2}}{2j} \right)} \right\}$$

$$E_t = E_0 e^{\frac{j(n-1)\psi}{2}} \frac{\sin n \psi/2}{\sin \psi/2}$$

$$E_T = E_0 \frac{\sin(n\psi_{1/2})}{\sin(\psi_{1/2})} e^{j\phi}$$

curve  $\phi = \frac{(n-1)\psi}{2}$

$$E_T = E_0 \frac{\sin(n\psi_{1/2})}{\sin(\psi_{1/2})} e^{j\phi}$$

$\rightarrow$  ②

This is the equation of total far field pattern of linear array of n-isotropic point source with source 1 as reference point for source.

\* If the reference point (source 1) is shifted to the centre of the array from the source or origin of the coordinate, then phase angle  $\phi$

is automatically eliminated,

$$E_T = E_0 \frac{\sin(n\psi_{1/2})}{\sin(\psi_{1/2})}$$

$\rightarrow$  ③

(Reference source located at the centre of co-ordinates)

\* If individual elements are non-isotropic but similar sources then,  
1)  $E_0$  will represent individual source pattern or primary pattern.

2)  $\frac{\sin(n\psi_{1/2})}{\sin(\psi_{1/2})}$  the array factor or secondary pattern.

for  $\psi = 0$ , then equation ③ becomes indeterminate.  
and hence L' hospital rule must be applied to evaluate the factor.

$$\begin{aligned} \lim_{\psi \rightarrow 0} (E_t) &= E_0 \cdot \lim_{\psi \rightarrow 0} \frac{\frac{d}{d\psi} (\sin n \psi/2)}{\frac{d}{d\psi} (\sin \psi/2)} \\ &= E_0 \lim_{\psi \rightarrow 0} \frac{\cos n \psi/2 \cdot \frac{1}{2}}{\cos \psi/2 \cdot \frac{1}{2}} \end{aligned}$$

$E_{t\text{man}} = E_0 \cdot n$

— ④

The normalised field pattern is obtained as,

$$\textcircled{3} \div \textcircled{4} \Rightarrow E_{\text{norm}} = \frac{E_t}{E_{t\text{man}}} = \frac{E_0 \frac{\sin n \psi/2}{\sin \psi/2}}{E_0 \cdot n}$$

$E_{\text{norm}} = \frac{\sin n \psi/2}{n \sin \psi/2} = (AF)_n$

$AF \Rightarrow$  Array Factor.

## ARRAY OF $n$ ISOTROPIC SOURCES OF EQUAL AMPLITUDE AND SPACING [BROADSIDE CASE]

(1) Maximum Radiation Intensity to array axis.

The equation of field pattern is given by

$$E_t = \frac{E_0 \sin n \Psi / 2}{\sin \Psi / 2}$$

$$\Psi = \beta d \cos \theta + \alpha$$

For broadside array individual elements are fed at same phase (ie  $\alpha = 0$ , and  $\beta d \cos \theta = 0$ )

$$\alpha = 0 \quad \beta d \cos \theta_{\text{max}} = 0$$

$$\Psi = \beta d \cos \theta \quad \theta_{\text{max}} = 90^\circ \text{ and } 270^\circ$$

Direction of major lobe (Pattern maxima)

$$E_t = \frac{E_0 \sin n \Psi / 2}{\sin \Psi / 2} \quad (1)$$

broadside  $\alpha = 0, \Psi = 0$

$$\Psi = \beta d \cos \theta = 0$$

$\cos \theta = 0$  major lobe  
 $\theta = 90^\circ \text{ or } 270^\circ$  [Principal maxima occurs in these directions]

Now with  $\Psi = 0$ ; the above eqn (1) becomes indeterminate

and hence D' Hospital rule must be applied to

evaluate the fun.

$$E_{t \text{ max}} = \lim_{\Psi \rightarrow 0} \frac{E_0 \sin n \frac{\Psi}{2}}{\sin \frac{\Psi}{2}} = \lim_{\Psi \rightarrow 0} \frac{E_0 \frac{\cos n \frac{\Psi}{2}}{-\frac{1}{2}} \cdot \frac{n}{2}}{\cos \frac{\Psi}{2} \cdot \frac{1}{2}}$$

$E_{t \text{ max}} = E_0 n$

maximum value of  $E_t$  is  $nE_0$  and it occurs at

$$\psi = 0 \text{ ie } \beta d \cos \theta = 0$$

or

$$\cos \theta = 0$$

$$\theta_{\text{max}} = 90^\circ \text{ and } 270^\circ$$

These are major lobe maxima.

minor lobe maxima

$$E_t = \frac{E_0 \sin n\psi/2}{\sin \psi/2}$$

This is maximum when numerator is maximum

$$\sin n\psi/2 = 1$$

$$(\text{Secondary max}) \text{ or } n\psi/2 = (2N+1)\frac{\pi}{2}$$

$$N = 1, 2, 3, 4, \dots$$

$N=0$  corresponds to major lobe ma-

$$\psi/2 = \pm (2N+1)\frac{\pi}{2} \cdot \frac{1}{n}$$

$$\boxed{\psi = \pm (2N+1)\frac{\pi}{n}}$$

$$\text{we know } \psi = \beta d \cos \theta$$

$$\beta d \cos(\theta_{\text{min}}) = \pm (2N+1)\frac{\pi}{n}$$

② we know  $B = \frac{2\pi}{\lambda}$

(3)

$$\frac{2\pi}{\lambda} d \cos(\theta_{\text{main}})_{\text{minor}} = \pm (2N+1) \frac{\pi}{n}$$

$$(\theta_{\text{main}})_{\text{minor}} = \cos^{-1} \left\{ \pm \frac{(2N+1) \frac{\pi}{n}}{2d} \right\}$$

$$\Rightarrow N = 1, 2, 3, \dots$$

$\Rightarrow n = \text{number of array elements}$

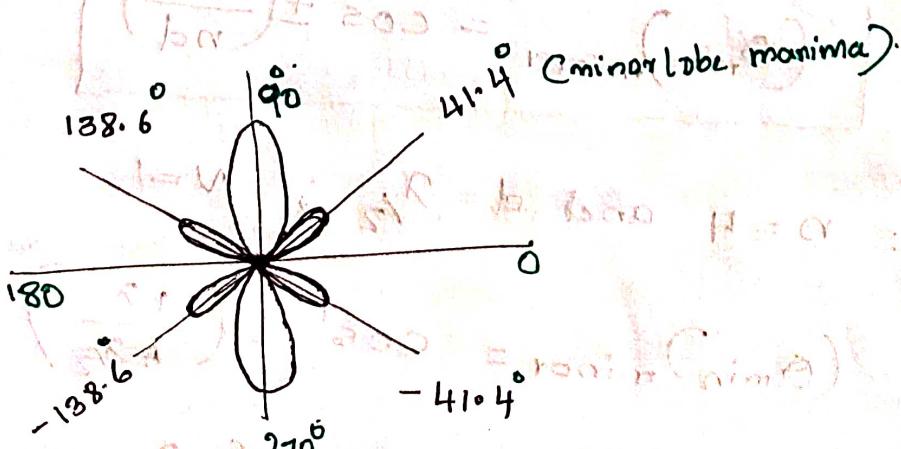
for eg.  $n=4$  and  $d = \frac{\lambda}{2}, N=1$

$$(\theta_{\text{main}})_{\text{minor}} = \cos^{-1} \left\{ \pm \frac{3\pi}{2 \times 4 \times \frac{\lambda}{2}} \right\}$$

$$= \cos^{-1} (\pm \frac{3\pi}{8}) \quad \cos^{-1} (\pm \frac{5\pi}{4})$$

$$(\theta_{\text{main}})_{\text{min}} = \pm 41.4^\circ \quad (\text{or } \pm 138.6^\circ)$$

are the four minor lobe maxima



## Direction of Pattern minima

$$E_T = E_0 \frac{\sin n\psi/2}{\sin \psi/2}$$

For  $E_T = 0$ ;  $\sin n\psi/2 = 0$  Provided  $\sin \psi/2 \neq 0$   
where  $N = 1, 2, 3, \dots$

$$\sin \frac{n\psi}{2} = 0$$

$$n\psi/2 = \pm N\pi$$

$$\boxed{\psi = \pm \frac{2N\pi}{n}}$$

$$\beta d(\cos \theta_{\min})_{\text{minor}} = \pm \frac{2N\pi}{n}$$

$$(\cos \theta_{\min})_{\text{minor}} = \pm \frac{2N\pi}{n} \times \frac{2}{2\pi} \times \frac{1}{d}$$

~~$$(\cos \theta_{\min})_{\text{minor}} = \pm \frac{N\lambda}{nd}$$~~

$$\boxed{(\theta_{\min})_{\text{minor}} = \cos^{-1} \pm \left( \frac{N\lambda}{nd} \right)}$$

eg:  $n=4$  and  $d = \lambda/2$ ;  $N=1$

$$(\theta_{\min})_{\text{minor}} = \cos^{-1} \pm \left( \frac{1\lambda}{4\lambda/2} \right)$$

$$(\theta_{\min})_{\text{minor}} = \cos^{-1} \pm \left( \frac{1}{2} \right)$$

$$(\theta_{\min})_{\text{minor}} = 60^\circ, 120^\circ, -60^\circ, -120^\circ$$

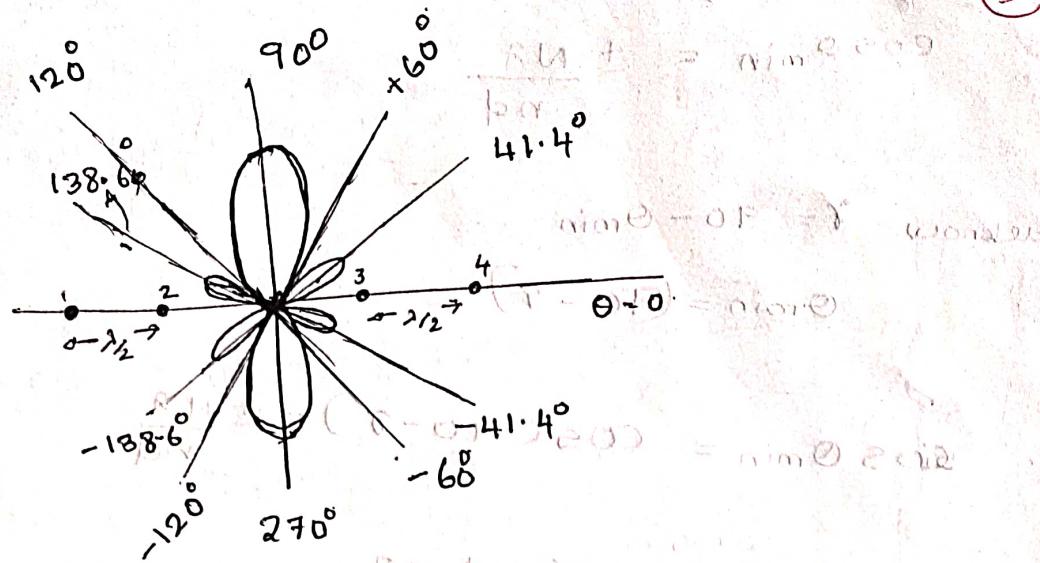
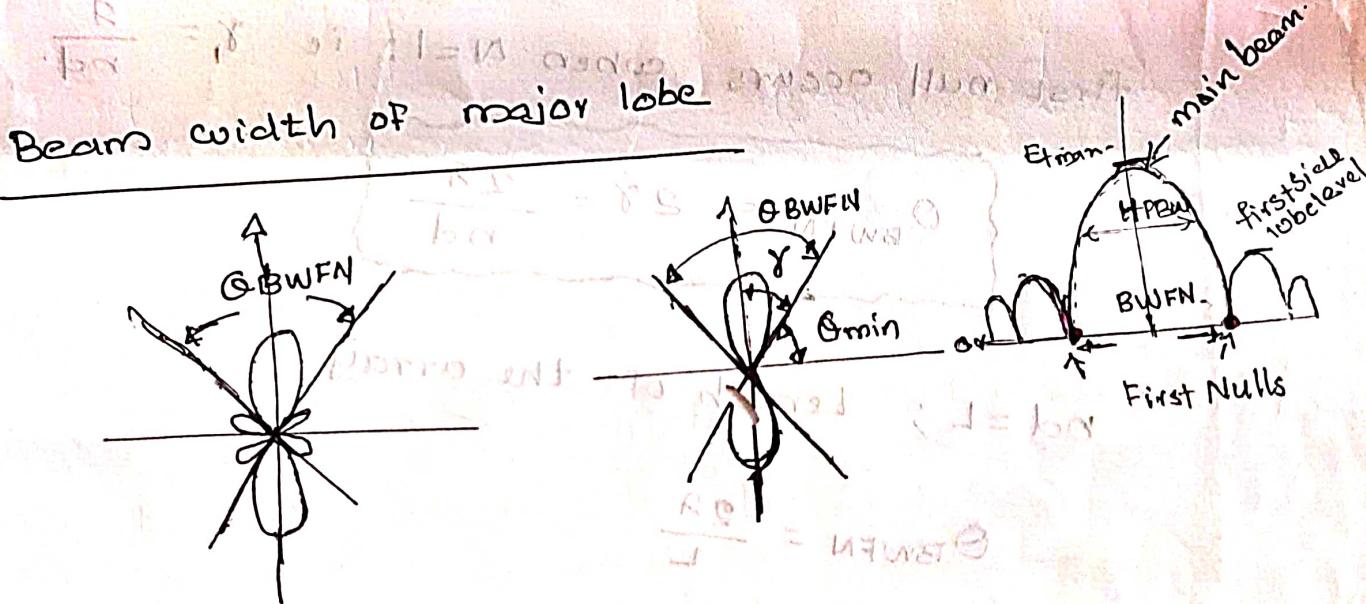


Fig : four isotropic sources of equal amplitude and phase, field pattern in broadside case.



it is the angle b/w first nulls

or  
double angle between first null and major lobe maximum

$$\Theta_{BWFN} = 2\gamma$$

$$= 2(90 - \Theta_{min})$$

$$\gamma = 90 - \Theta_{min}$$

complementary angle.

$$\text{we know } \Theta_{min} = \cos^{-1} \left( \pm \frac{N \lambda}{4 \pi d} \right)$$

$$\cos \theta_{\min} = \pm \frac{N\lambda}{nd}$$

we know  $\gamma = 90 - \theta_{\min}$

$$\theta_{\min} = (90 - \gamma)$$

$$\sin \theta_{\min} = \cos(90 - \gamma) = \pm \frac{N\lambda}{nd}$$

$$= \sin \gamma = \pm \frac{N\lambda}{nd}$$

$\because \gamma$  is very small;  $\sin \gamma = \gamma$

$$\gamma = \pm \frac{N\lambda}{nd}$$

first null occurs when  $N=1$ ; ie  $\gamma_1 = \frac{\lambda}{nd}$ .

$$\theta_{\text{BWFN}} = 2\gamma = \frac{2\lambda}{nd}$$

$nd = L$ ; Length of the array

$$\theta_{\text{BWFN}} = \frac{2\lambda}{L}$$

$$= \frac{2}{(L/\lambda)} \text{ radian}$$

$$\text{radian} = \left(\frac{180}{\pi}\right) \text{ degree}$$

$$\text{radian} = 57.3^\circ \text{ degree}$$

$$= \frac{2 \times 57.3^\circ}{(L/\lambda)} \text{ degree}$$

Beam width b/w

First Null

$$\text{BWFN} = \frac{114.6^\circ}{L/\lambda} = \frac{114.6 \times \lambda}{nd}$$

$$\text{Half Power Beam width HPBW} = \frac{\text{BWFN}}{2}$$

4

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$$\textcircled{5} \text{ HPBW} = \frac{57.3^\circ}{(L/\lambda)}$$

\* The Null-to-Null beamwidth of broadside array [BWN] =

$$\text{array [BWN]} = \frac{2\lambda}{nd}$$

$$\text{Directivity } D = 2\left(\frac{L}{\lambda}\right) = 2\left(\frac{nd}{\lambda}\right)$$

Array of n sources of equal amplitude and spacing

[END-FIRE CASE]

We know the General expression for  $E_T$ , is

$$E_T = E_0 \sin \frac{n\psi}{2}$$

$$\sin \frac{\psi}{2}$$

- \* for end-fire array maximum radiation occurs when  $\psi=0$  and  $\theta=0^\circ$  or  $180^\circ$ . (i.e. The direction of maximum radiation coincides with the array axis).

when  $\psi=0$ ;

$$\lim_{\psi \rightarrow 0} E_T = \lim_{\psi \rightarrow 0} E_0 \frac{\cos n \frac{\psi}{2} \cdot \frac{n}{2}}{\cos \frac{\psi}{2} \cdot \frac{1}{2}}$$

$$E_{T_{\text{max}}} = n E_0$$

$$\theta_{\text{max}} = 0^\circ \text{ and } 180^\circ$$

Direction of main lobe  
( $\psi=0, \alpha=-\beta d$ )

We know  $\psi = \beta d \cos \theta + \alpha$

$$\alpha = \beta d \cos \theta + \alpha$$

$$\alpha = -\beta d$$

→ This is the Condition for end-fire array

- \* If spacing b/w two sources  $\lambda/2$  or  $\lambda/4$ , then Phase angle by which source 2 lags behind source 1 is  $2\pi \cdot \frac{\lambda}{2}$  or  $\frac{2\pi}{\lambda} \cdot \frac{\lambda}{4} = \pi$  or  $\pi/2$ .

(6)

## Direction of minor lobe Maximum

$$E_t = E_0 \frac{\sin n\psi/2}{\sin \psi/2}$$

which is maximum when

$$\sin n\psi/2 = 1 \quad \text{and} \quad \sin \psi/2 \neq 0$$

$$n\psi/2 = \pm (2N+1)\frac{\pi}{2}$$

$$\psi = \pm \frac{(2N+1)\pi}{n}$$

$$\beta d \cos(\theta_{\text{max}})_{\text{minor}} + \alpha = \pm \frac{(2N+1)\pi}{n}$$

we know  $\alpha = -\beta\phi$

$$\beta d \cos(\theta_{\text{max}})_{\text{minor}} - \beta d = \pm \frac{(2N+1)\pi}{n}$$

$$\beta d [\cos(\theta_{\text{max}})_{\text{minor}} - 1] = \pm \frac{(2N+1)\pi}{n}$$

$$\cos(\theta_{\text{max}})_{\text{minor}} - 1 = \pm \frac{(2N+1)\pi}{n\beta d}$$

$$\cos(\theta_{\text{max}})_{\text{minor}} = \pm \frac{(2N+1)\pi}{n\beta d} + 1$$

$$(\theta_{\text{man}})_{\text{minor}} = \cos^{-1} \left\{ \pm \frac{(2N+1)\pi}{nBd} + 1 \right\}$$

example.

$$\text{If } n=4, d=\frac{\lambda}{2}, \alpha=-Bd=-\pi, N=1$$

$$(\theta_{\text{man}})_{\text{minor}} = \cos^{-1} \left\{ \pm \frac{(2 \times 1 + 1)\pi}{4 \times 2\pi \times \frac{\lambda}{2}} + 1 \right\}$$

$$= \cos^{-1} \left\{ \pm \frac{3}{4} + 1 \right\}$$

$$= \left( \frac{2 \times N + 1}{4} + 1 \right)$$

$$= \cos^{-1} \left\{ \frac{7}{4}, \frac{1}{4} \right\}$$

$$= \cos^{-1} \left\{ \frac{7}{4} \right\}, \therefore \cos^{-1} \left( \frac{7}{4} \right) \text{ does not exist.}$$

$$(\theta_{\text{man}})_{\text{minor}_1} = 75 \cdot 5^\circ$$

→ If  $N=2$

$$(\theta_{\text{man}})_{\text{minor}_2} = \cos^{-1} \left\{ \pm \frac{5}{4} + 1 \right\}$$

$$= \cos^{-1} \left\{ \frac{9}{4}, -\frac{1}{4} \right\}$$

$$= \cos^{-1} \left( -\frac{1}{4} \right) \because \cos^{-1} \left( \frac{9}{4} \right) \text{ does not exist}$$

$$(\theta_{\text{man}})_{\text{minor}_2} = 104 \cdot 5^\circ$$

## Direction of pattern minima

$$E_t = E_0 \frac{\sin n\psi_{1/2}}{\sin \psi_{1/2}}$$

$E_t = 0$ , when  $\sin n\psi_{1/2} = 0$  and  $\sin \psi_{1/2} \neq 0$

$$\sin n\psi_{1/2} = 0$$

$$n\psi_{1/2} = \pm N\pi$$

$$\psi = \pm \frac{2N\pi}{n}$$

We know  $\psi = \beta d \cos \theta + \alpha$ ,  ~~$\alpha = \beta d$~~

$$\beta d \cos(\theta_{\min})_{\text{min.}} + \alpha = \pm \frac{2N\pi}{n}$$

$$\text{hence, } \alpha = -\beta d$$

$$\beta d \cos(\theta_{\min})_{\text{min.}} - \beta d = \pm \frac{2N\pi}{n}$$

$$\beta d [\cos(\theta_{\min})_{\text{min.}} - 1] = \pm \frac{2N\pi}{n}$$

$$[\cos(\theta_{\min})_{\text{min.}} - 1] = \pm \frac{2N\pi}{n\beta d}$$

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$$\beta = 2\pi/2$$

$$\cos \theta_{\min} - 1 = \pm \frac{2N\pi}{2\pi/2 nd} = \pm \frac{N\lambda}{nd}$$

**We know**  $\cos \theta = 1 - 2 \sin^2 \theta/2$

$$1 - 2 \sin^2 \theta_{\min} - 1 = \pm \frac{N\lambda}{nd}$$

$$2 \sin^2 \frac{\theta_{\min}}{2} = \pm \frac{N\lambda}{nd}$$

$$\sin^2 \frac{\theta_{\min}}{2} = \pm \frac{N\lambda}{2nd}$$

$$\sin \frac{\theta_{\min}}{2} = \pm \sqrt{\frac{N\lambda}{2nd}}$$

$$\frac{\theta_{\min}}{2} = \sin^{-1} \left( \pm \sqrt{\frac{N\lambda}{2nd}} \right)$$

$$\theta_{\min} = 2 \sin^{-1} \left( \pm \sqrt{\frac{N\lambda}{2nd}} \right)$$

Example  $n = 4, d = \lambda/2, N = 1$

(10) (8)

$$(\theta_{\min})_0 = 2 \sin^{-1} \left( \pm \sqrt{\frac{N^2}{2nd}} \right)$$

$$(\theta_{\min})_1 = 2 \sin^{-1} \left( \pm \sqrt{\frac{1 \times 2}{2 \times 4 \times 2/2}} \right)$$

$$= 2 \sin^{-1} \left( \pm \sqrt{\frac{1}{4}} \right) = 2 \sin^{-1} \left( \pm \frac{1}{2} \right)$$

$$= 2 \times (\pm 30^\circ)$$

$$(\theta_{\min})_1 = \pm 60^\circ$$

If  $N=2$ .

$$(\theta_{\min})_2 = 2 \sin^{-1} \left( \pm \sqrt{\frac{2}{4}} \right) = 2 \sin^{-1} \left( \pm \sqrt{\frac{1}{2}} \right)$$

$$(\theta_{\min})_2 = 2 \times (\pm 45^\circ)$$

$$(\theta_{\min})_2 = \pm 90^\circ$$

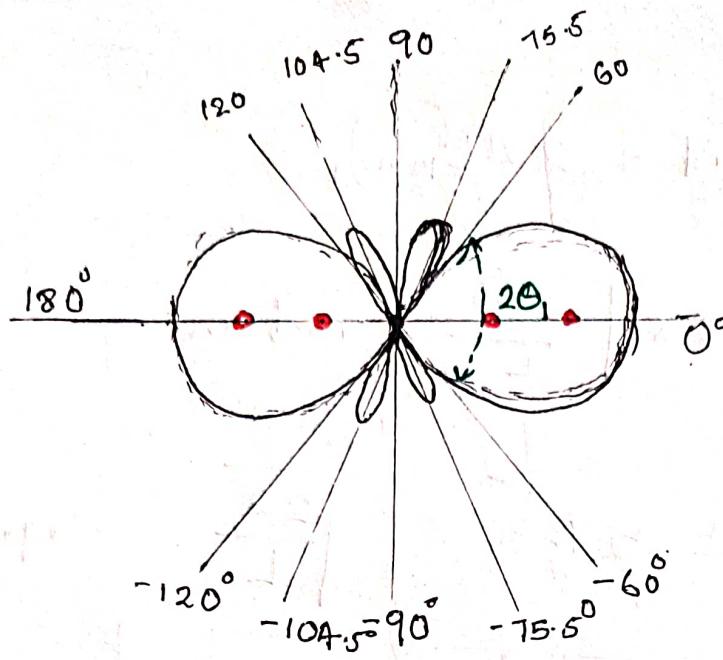
If  $N=3$ 

$$(\theta_{\min})_3 = 2 \sin^{-1} \left( \pm \sqrt{\frac{3}{4}} \right) = 2 \sin^{-1} \left( \pm \frac{\sqrt{3}}{2} \right)$$

$$= 2 \times (\pm 60^\circ)$$

$$(\theta_{\min})_3 = \pm 120^\circ$$

$$\boxed{\theta_{\min} = 60^\circ, 90^\circ, 120^\circ, -60^\circ, -90^\circ, -120^\circ}$$



### Beam Width of Major lobe

Beam Width =  $2 \times$  angle b/w first null and maximum of major lob

$$BWFN = 2\theta_1$$

$$\theta_{min} = 2 \sin^{-1} \left( \pm \sqrt{\frac{Nr}{2nd}} \right)$$

$$\sin \theta_{min} = 2 \left( \pm \sqrt{\frac{Nr}{2nd}} \right)$$

If  $\theta_{min}$  very small,  $\therefore \sin \theta_{min} \approx \theta_{min}$

$$\theta_{min} = \pm \sqrt{\frac{2Nr}{nd}}$$

If the array is long of length L then

$$L = (n-1)d \approx nd.$$

$$(\theta_{min}) = \pm \sqrt{\frac{2Nr}{L}}$$

→ IF N=1

$$(\theta_{min})_1 = \pm \sqrt{\frac{2R}{L}}$$

$$BWNFN = 2 \times (\theta_{\min}),$$

$$= \pm 2 \times \sqrt{\frac{2\lambda}{L}} \text{ radian}$$

$$BWNFN = \pm 57.3^\circ \times 2 \sqrt{\frac{2\lambda}{L}} \text{ degree}$$

$$\boxed{BWNFN = \pm 114.6 \sqrt{\frac{2}{L/\lambda}}}$$

\* Null-to-Null beam width of end-fire array =  $2 \sqrt{\frac{2\lambda}{nd}}$

\* The directivity of end-fire array,  $D = 4 \left( \frac{nd}{\lambda} \right)$

$$\boxed{D = 4 \left( \frac{nd}{\lambda} \right)}$$

## Binomial arrays

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→ Binomial array is an array whose elements are excited according to the current levels determined by the binomial coefficient,  $n_r$ ,  $n$  being the number of element in the array.

→ It is a non-uniform array.

→ Advantage: ~~no~~ no side lobes in the resultant

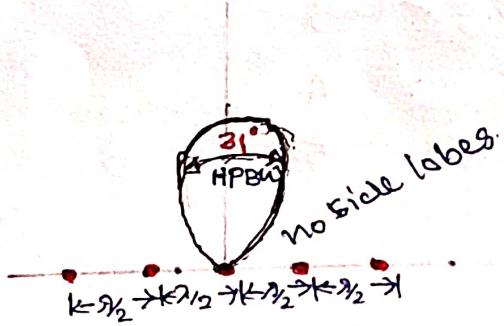
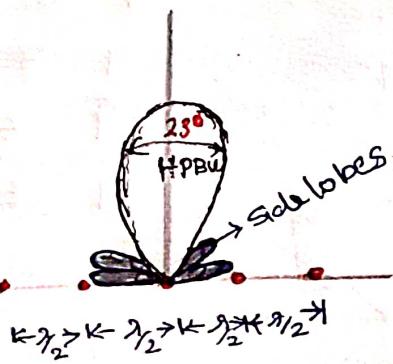
### Pattern of array

(Note: A uniform distribution yields the maximum directivity

but the minor lobes are relatively large. So is the

main objective of  
constructed binomial array to eliminate this minor or

side lobes in the resultant pattern of array.



eg:  
 $n = 5$ ,  
 $d = \lambda/2$

### Uniform array

Binomial array  
amplitude ratio  $1:4:6:4:1$

→ Disadvantage: 1) HPBW increases and hence the directivity decreases  
2) For design of a large array, larger amplitude ratio of sources is required.

→ To reduce the sidelobe level, John Stone

Proposed that sources have amplitudes proportional  
to the coefficients of a binomial series of the form.

$$(a+b)^{n-1} = a^{n-1} + (n-1) a^{n-2} b + \frac{(n-1)(n-2)(n-3)}{2!} a^2 b^2 + \dots$$

where  $n$  is the number of sources.

For array of 2 to 6 <sup>(elements)</sup> sources, the relative amplitude levels are given by

Number of elements

$n = 1$

2

3

4

5

6

Relative amplitude

1						
1	1					
1	2	1				
1	3	3	1			
1	4	6	4	1		
1	5	10	10	5	1	

→ These coefficients for any number of radiating sources can also be obtained from Pascal's triangle

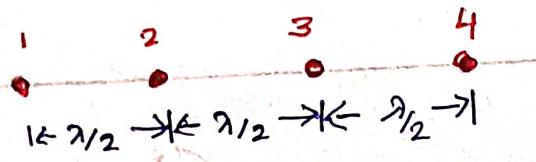
1	1	1			
	1	2	1		
		3	3	1	
			4	6	4
				10	10
					5

$$\Rightarrow n = 3, \quad L = 2d. \\ = (3-1)d$$

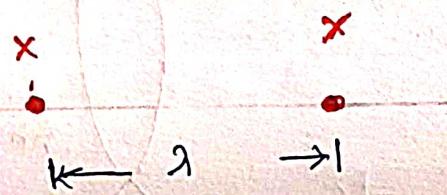
$$\text{length, } L = (n-1)d$$

eg: case I : Uniform array

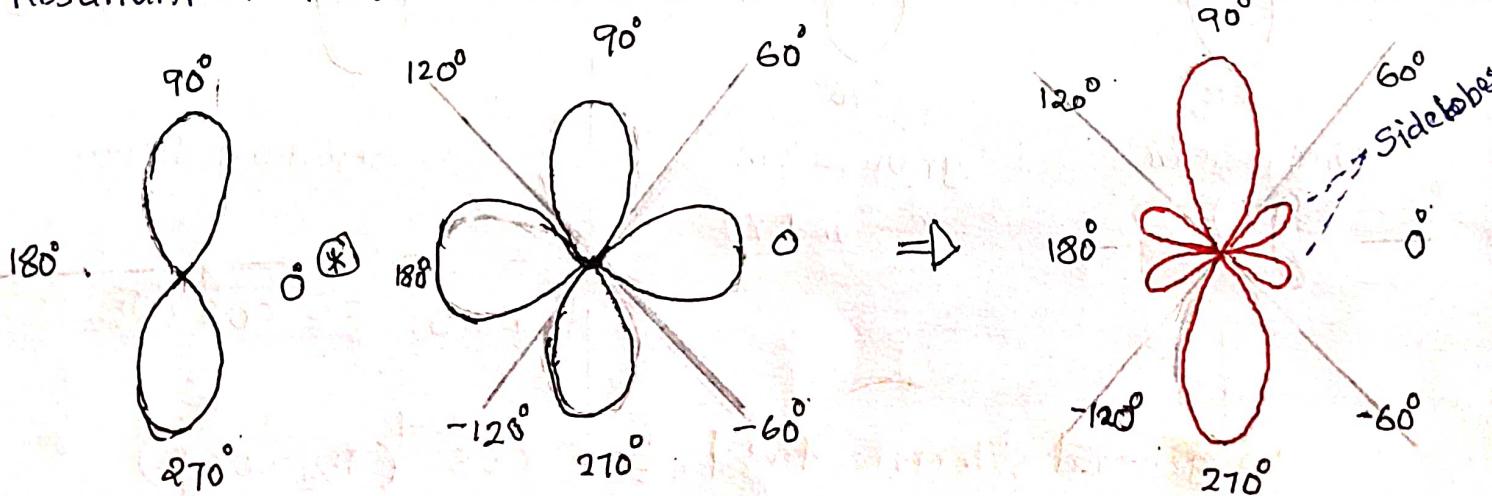
$$\text{If } n=4, d = \frac{1}{2}$$



$$L = (n-1)d$$



Resultant Pattern = Unit pattern of X or Y  $\oplus$  group pattern X and Y



Unit pattern  
(X or Y)

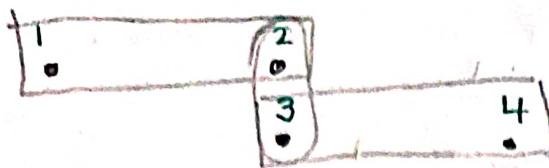
group pattern  
(X and Y)

$\Rightarrow$  Resultant Pattern

$$E = \cos(\pi/2 \cos \theta)$$

Case II:

For Binomial array



$$\leftarrow \pi_{1/2} \rightarrow 1 \leftarrow \pi_{1/2} \rightarrow 1$$

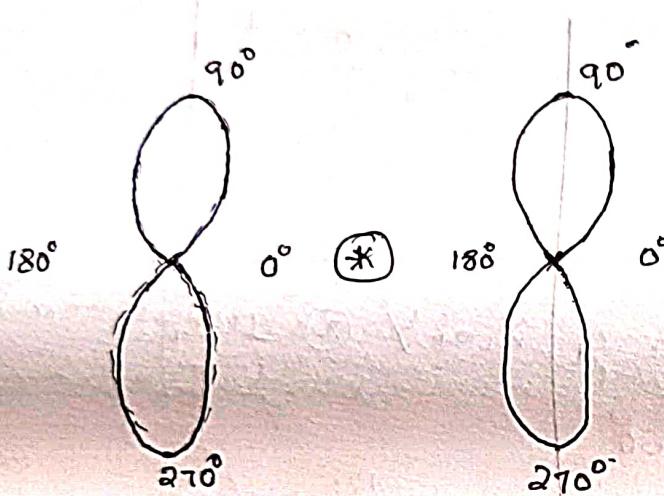
$$I \quad 2I \quad I$$

$$\leftarrow \pi_{1/2} \rightarrow 1 \leftarrow \pi_{1/2} \rightarrow 1$$

$$\text{Here } n=3$$

$$X \quad Y$$

$$\leftarrow \pi_{1/2} \rightarrow 1$$



Unit pattern  
(X or Y)

group pattern  
(X and Y)

Resultant Pattern

$$\text{if } n=3, E = \cos^2(\pi_{1/2} \cos \theta)$$

$$\text{In general Electric field, } E = \cos^{n-1}(\pi_{1/2} \cos \theta)$$

$$\text{Length of array, } L = (n-1) \lambda_{1/2}$$

$$\boxed{\text{HPBW} = \frac{1.06}{\sqrt{n-1}}}$$

$$\boxed{D = 1.77 \sqrt{n}}$$

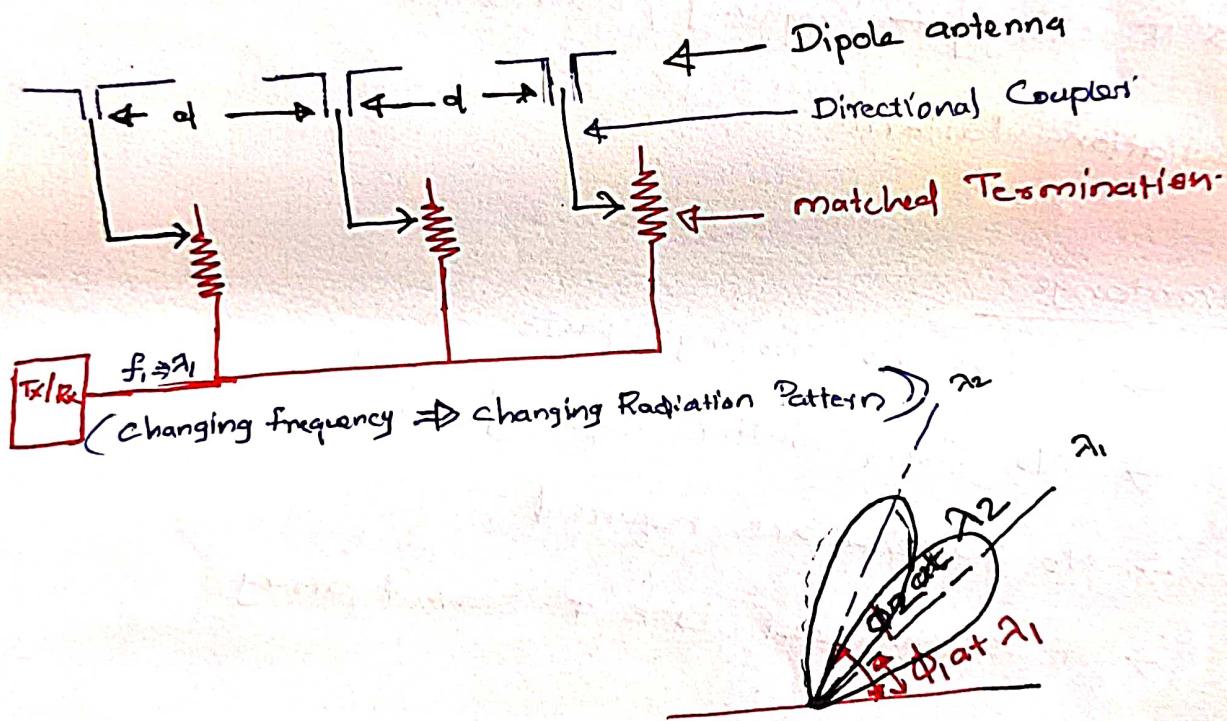
$$\text{HPBW} = \frac{1.06}{\sqrt{2L/\lambda}} = \frac{0.75}{\sqrt{L/\lambda}}$$

$$D = 1.77 \sqrt{1 + 2(L/\lambda)}$$

## Dolph - Tchebychev Array

→ In this antenna array, direction of the beam of radiation is swept or scanned by changing the frequency.

→ Consider a line-feed array of uniformly spaced elements with a Tx/Rx connected to the right end of the line as shown in below figure.



→ Here position of antenna and coupler is fixed and direction of radiation can be changed by changing frequency.

→ Angle of radiation  $\phi$  is expressed as

$\phi \Rightarrow$  Angle of Radiation  
 $c \Rightarrow$  Velocity of light

$v \Rightarrow$  Phase velocity

$m \Rightarrow$  mode number

$d \Rightarrow$  spacing b/w elements

$\lambda \Rightarrow$  wavelength.

$$\cos \phi = \frac{c}{v} + \frac{m}{d/\lambda}$$

for example

$$\Rightarrow \text{If } c = v, \lambda = d, m = -1, \phi = ? \quad \Rightarrow \text{If } c = v, \lambda = 0.9d, m = -1$$
$$\cos\phi = 1 + \frac{(-1)}{1} = 1 - 1 = 0 \quad \cos\phi = 1 + \frac{(-1)}{1/0.9}$$
$$\cos\phi = 0 \quad \cos\phi = 0.1$$
$$\phi = 90^\circ \quad \phi = 84.26^\circ$$

→ If changing in frequency / wavelength, we can steer the  
Radiation Pattern.

### Advantage

- It has no moving parts }  $\Rightarrow$  Low Cost
- No Phase shifter required
- No switches are required
- It is simplest type of scanned array

### Tchebyshov Polynomial

→ on the name of Tchebyshov, the first letter T is used  
as symbol.

$$\Rightarrow T_m(x) = \cos(m \cos^{-1} x), \text{ for } |x| \leq \pm 1$$

$$T_m(x) = \cosh(m \cosh^{-1} x), \text{ for } |x| \geq \pm 1$$

Consider

$$\delta = \cos^{-1} x \rightarrow x = \cos \delta$$

$$T_m(x) = \cos(m\delta)$$

for  $m=0$

$$T_0(x) = \cos 0 = 1$$

for  $m=1$

$$T_1(x) = \cos \delta = x$$

for  $m=2$

$$T_2(x) = \cos(2\delta)$$

$$= 2\cos^2 \delta - 1$$

$$= 2x^2 - 1$$

for  $m=3$

$$T_3(x) = \cos(3\delta)$$

$$= 4\cos^3 \delta - 3\cos \delta$$

$$= 4x^3 - 3x$$

Further higher term  
Relation b/w the Poly

$\Rightarrow$  From above equation we can find the

$$T_{m+1}(x) = 2x T_m(x) - T_{m-1}(x) \quad \text{--- (3)}$$

$$T_3(x) = 2x T_2(x) - T_1(x)$$

$$= 2x(2x^2 - 1) - x$$

$$= 4x^3 - 3x$$

→ For  $m=4$ , we must put  $m=3$  in above equation we

$$T_4(x) = 2 \times T_3(x) - T_2(x)$$

$$= 2 \times (4x^3 - 3x) - 2x^2 + 1$$

$$T_4(x) = 8x^4 - 8x^2 + 1$$

Similarly, the value of  $T_5$ ;  $T_6$ ,  $T_7$ ,  $T_8$ ,  ~~$T_9$~~ , and  $T_{10}$  also found and summarized as follows.

$$T_0(n) = 1$$

$$T_1(n) = n$$

$$T_2(n) = 2n^2 - 1$$

$$T_3(n) = 4n^3 - 3n$$

$$T_4(n) = 8n^4 - 8n^2 + 1$$

$$T_5(n) = 16n^5 - 20n^3 + 5n$$

$$T_6(n) = 32n^6 - 48n^4 + 18n^2 - 1$$

$$T_7(n) = 64n^7 - 112n^5 + 56n^3 - 7n$$

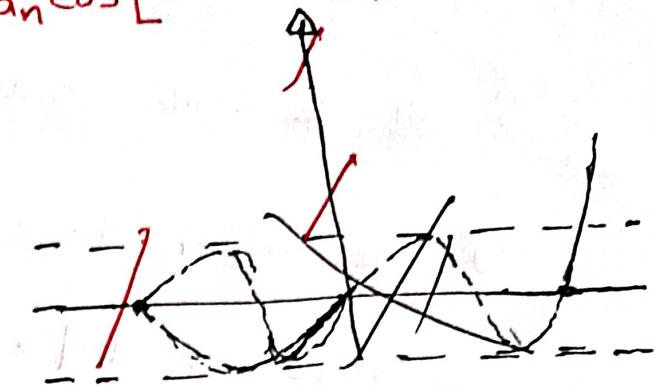
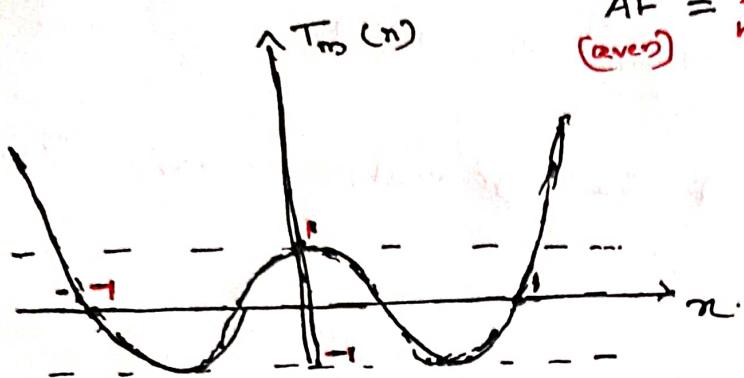
$$T_8(n) = 128n^8 - 256n^6 + 160n^4 - 32n^2 + 1$$

$$T_9(n) = 256n^9 - 576n^7 + 432n^5 - 120n^3 + 9n$$

$$T_{10}(n) = 512n^{10} - 1280n^8 - 608n^6 + 40n^4 - 14n$$

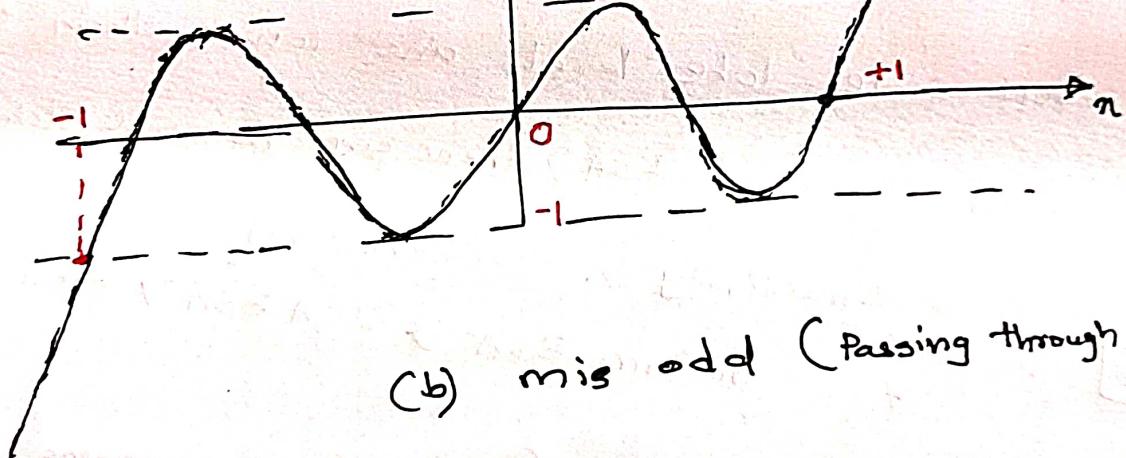
→ General Characteristics of DCA

$$AF = \sum_{n=1}^M a_n \cos[(2n-1)u], N=2M \text{ even}$$



(a)  $m$  is even.

$$AF = \sum_{n=1}^{M+1} a_n \cos[2(n-1)u], N=2M+1$$



(b)  $m$  is odd (Passing through origin)

→ Side lobes arise in the region  $x < 1$

→ The main lobe extends into the range  $n > 1$ .

→ The value of  $m$  and the degree of Tchebyshoff

Polynomial are the same.

→ Properties of chebyshev Polynomials:

- 1] All polynomials of any order 'm' pass through the point (1, 1)
- 2] All the nulls occur within  $-1 \leq n \leq 1$
- 3] The higher the order of the Polynomial (mainlobe), the steeper the slope for  $|n| > 1$

### Advantage

- 1] For a given number of elements directivity ~~rent~~ after that of uniform BSA, ~~is greater than Binomial BSA but side lobe levels~~ (for Binomial BSA, there are no side lobe at all).

Directivity & Side lobe level  
 $\text{uniform BSA} > \text{DCA BSA} > \text{Binomial BSA}$

2]

Lower side lobe (i.e. DCA provides minimum

~~Binomial BSA~~ optimum beam width for a specified degree of side-lobe level reduction)

2] It results in side lobes are lobes are all of the same amplitudes, unlike in Uniform BSA -

3) Ratio of current between centre element and end element is small, which provides ease in feeding design.