

Q. Derive the expression for energy stored in Magnetic field.

Ans: A simple expression for magnetic energy in the field of inductor:  $W_m = \frac{1}{2} LI^2$  — (1)

Consider a differential volume in a magnetic field. Let the volume be covered with conducting sheets at the top and bottom surfaces with current  $\Delta I$ .

→ Assuming the whole region is filled with such differential volume.

Each volume has an inductance

$$\Delta L = \frac{\Delta \Phi}{\Delta I} \quad \text{--- (2)}$$

$$\Delta \Phi = \vec{B} \cdot \Delta S \quad \text{--- (3)} \quad \Delta S = \Delta x \cdot \Delta z \quad \text{--- (4)}$$

③ and ④ in ②;

$$\Delta L = \frac{\vec{B} \cdot \Delta x \cdot \Delta z}{\Delta I} \quad \text{--- (5)}$$

⑤ and ⑤ in ①;

$$\Delta W_m = \frac{1}{2} \frac{\vec{B} \cdot \Delta x \cdot \Delta z \cdot (\Delta I)^2}{\Delta I}$$

$$= \frac{1}{2} \vec{B} \cdot \Delta x \cdot \Delta z \cdot \Delta I$$

$$= \frac{1}{2} \vec{B} \cdot \Delta x \cdot \Delta z \cdot \vec{H} \cdot \Delta y \quad \{ \vec{B} = \mu \vec{H} \}$$

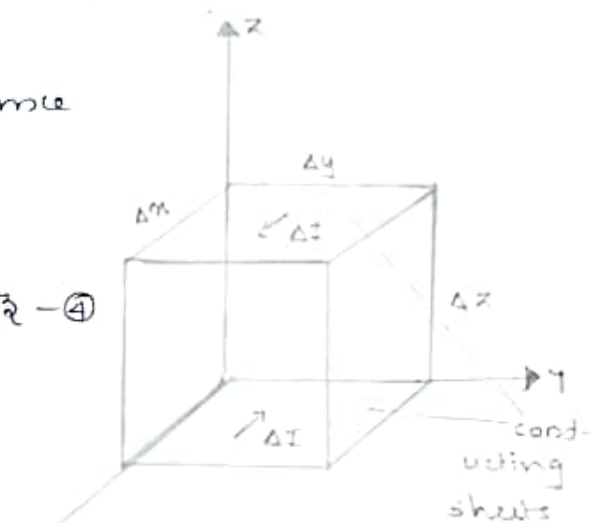
$$\Delta W_m = \frac{1}{2} \mu \vec{H} \cdot \Delta x \cdot \Delta z \cdot \vec{H} \cdot \Delta y = \frac{1}{2} \mu \vec{H}^2 \cdot \Delta x \cdot \Delta y \cdot \Delta z$$

$$\Delta W_m = \frac{1}{2} \mu \vec{H}^2 \Delta V$$

The magnetostatic energy density  $w_m$  (in J/m<sup>3</sup>) is;

$$w_m = \lim_{\Delta V \rightarrow 0} \frac{\Delta W_m}{\Delta V} = \frac{1}{2} \mu \vec{H}^2$$

$$w_m = \frac{1}{2} \mu \vec{H}^2 = \frac{1}{2} \vec{B} \cdot \vec{H} = \frac{\vec{B}^2}{2\mu}$$



$$\left\{ \begin{array}{l} \Delta I = H \cdot \Delta y \Rightarrow \text{By using} \\ \text{Ampere's law} \\ [\oint \vec{H} \cdot d\vec{l} = I_{enc}] \end{array} \right.$$

Thus the energy in magnetostatic field in linear medium;

$$W_m = \int w_m dv = \frac{1}{2} \int \vec{B} \cdot \vec{H} dv = \frac{1}{2} \int \mu \vec{H}^2 dv.$$

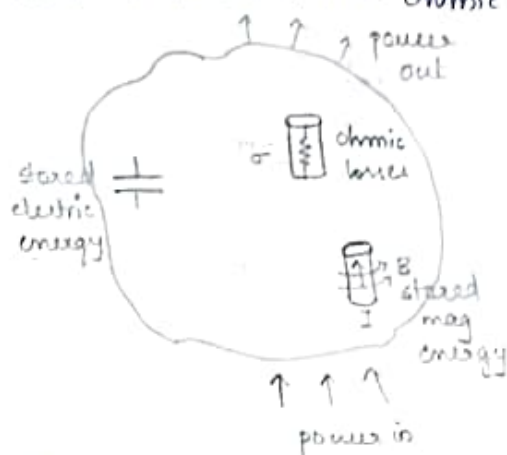
Q. State Poynting theorem. Derive Poynting theorem starting from Maxwell's equations.

Ans: From Maxwell's third and fourth equations;

Poynting theorem states that the net power flowing out of a given volume  $V$  is equal to the time rate of decrease in the energy stored within  $V$  minus the ohmic losses.

→ Poynting vector: The cross product of  $E$  and  $H$  at any point gives the power per unit area.

$$\underline{P = E \times H} \text{ watts/m}^2$$



From Maxwell's 3rd and 4th equations:

$$\nabla \times \vec{E} = -\mu \frac{\partial \vec{H}}{\partial t} \quad \text{--- (1)}$$

$$\nabla \times \vec{H} = \sigma \vec{E} + \epsilon \frac{\partial \vec{E}}{\partial t} \quad \text{--- (2)}$$

Taking dot product with  $\vec{E}$  on both sides of (2)

$$\vec{E} \cdot (\nabla \times \vec{H}) = \sigma \vec{E}^2 + \epsilon \left[ \vec{E} \cdot \frac{\partial \vec{E}}{\partial t} \right] \quad \text{--- (3)}$$

Consider a vector identity;

$$\nabla \cdot (\vec{A} \times \vec{B}) = \vec{B} \cdot (\nabla \times \vec{A}) - \vec{A} \cdot (\nabla \times \vec{B}) \quad \text{--- (4)}$$

$$\text{Let } \vec{A} = \vec{H} \text{ and } \vec{B} = \vec{E}$$

$$\text{(4)} \Rightarrow \nabla \cdot (\vec{H} \times \vec{E}) = \vec{E} \cdot (\nabla \times \vec{H}) - \vec{H} \cdot (\nabla \times \vec{E}) \quad \text{--- (5)}$$

$$\vec{E} \cdot (\nabla \times \vec{H}) = \nabla \cdot (\vec{H} \times \vec{E}) + \vec{H} \cdot (\nabla \times \vec{E}) = \sigma \vec{E}^2 + \epsilon \left[ \vec{E} \cdot \frac{\partial \vec{E}}{\partial t} \right] \quad \text{--- (6)}$$

$$\vec{H} \cdot (\nabla \times \vec{E}) = \nabla \cdot (\vec{E} \times \vec{H}) - \sigma \vec{E}^2 + \epsilon \left( \vec{E} \cdot \frac{\partial \vec{E}}{\partial t} \right)$$

$$-\nabla \cdot (\vec{E} \times \vec{H}) = \sigma E^2 + \epsilon \left( \vec{E} \cdot \frac{\partial \vec{E}}{\partial t} \right) - \mu (\nabla \times \vec{E}) \quad \left\{ \nabla \times \vec{E} = -\mu \frac{\partial \vec{H}}{\partial t} \right\}$$

$$-\nabla \cdot (\vec{E} \times \vec{H}) = \sigma E^2 + \epsilon \left( \vec{E} \cdot \frac{\partial \vec{E}}{\partial t} \right) + \mu \mu \cdot \frac{\partial \vec{H}}{\partial t} \quad \text{--- (4)}$$

let's take;

$$\frac{\partial}{\partial t} (\vec{E} \cdot \vec{E}) = \vec{E} \cdot \frac{\partial \vec{E}}{\partial t} + \vec{E} \frac{\partial \vec{E}}{\partial t} = 2 \vec{E} \cdot \frac{\partial \vec{E}}{\partial t}$$

$$\therefore \underline{\underline{\vec{E} \cdot \frac{\partial \vec{E}}{\partial t} = \frac{1}{2} \frac{\partial \vec{E}^2}{\partial t}}}$$

Similarly;

$$\vec{H} \cdot \frac{\partial \vec{H}}{\partial t} = \frac{1}{2} \frac{\partial \vec{H}^2}{\partial t}$$

$$\therefore -\nabla \cdot (\vec{E} \times \vec{H}) = \sigma E^2 + \frac{\epsilon}{2} \frac{\partial \vec{E}^2}{\partial t} + \frac{\mu}{2} \frac{\partial \vec{H}^2}{\partial t}$$

Take volume integral;

$$-\int_V \nabla \cdot (\vec{E} \times \vec{H}) dV = \int_V \sigma E^2 dV + \int_V \left( \frac{\epsilon}{2} \frac{\partial \vec{E}^2}{\partial t} + \frac{\mu}{2} \frac{\partial \vec{H}^2}{\partial t} \right) dV$$

Applying divergence theorem to LHS;

$$-\oint_S (\vec{E} \times \vec{H}) \cdot d\vec{s} = + \frac{\partial}{\partial t} \int_V \sigma E^2 dV + \frac{\partial}{\partial t} \int_V \left( \frac{\epsilon E^2}{2} + \frac{\mu H^2}{2} \right) dV$$

multiplying by -ve sign;

$$\underline{\underline{\oint_S (\vec{E} \times \vec{H}) \cdot d\vec{s} = - \frac{\partial}{\partial t} \int_V \left[ \frac{1}{2} \epsilon \vec{E}^2 + \frac{1}{2} \mu \vec{H}^2 \right] dV - \int_V \sigma E^2 dV}} \quad \text{--- (5)}$$

i.e., Total power leaving the volume = Rate at which stored electromagnetic energy is decreasing in the volume + Power dissipated in the volume.

(5) is referred to as Poynting theorem.

Q. Derive the Conductor-dielectric and Conductor-free space boundary conditions for electric field.

Ans: - DIELECTRIC-CONDUCTOR BOUNDARY CONDITIONS:

Under static conditions, the following assumptions can be made about perfect conductor.

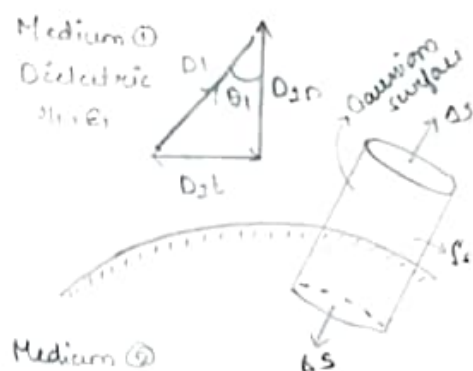
- No electric field exists within a perfect conductor ( $E=0$ )
- Conductor is considered as equipotential surface.

Applying Gauss's law to the Gaussian surface;

$$\oint_S \vec{D} \cdot d\vec{s} = Q_{enc}$$

$$D_{sn} \Delta S - 0(\Delta S) = \rho_s \Delta S$$

$$\boxed{D_{sn} = \rho_s} \quad \text{--- (1)}$$



Medium ②  
Perfect conductor  
 $\vec{D} \cdot \vec{E} = 0$   
 $\rho_v = 0$

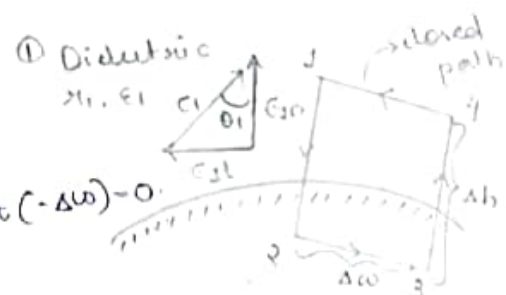
Applying  $\oint_C \vec{E} \cdot d\vec{l} = 0$  to the path 1-2-3-4-1.

$$\int_1^2 \vec{E} \cdot d\vec{l} + \int_2^3 \vec{E} \cdot d\vec{l} + \int_3^4 \vec{E} \cdot d\vec{l} + \int_4^1 \vec{E} \cdot d\vec{l} = 0$$

$$= E_1 n \left( -\frac{\Delta h}{2} \right) + 0 + E_1 n \left( \frac{\Delta h}{2} \right) + E_1 t (-\Delta w) = 0$$

$$-\cancel{E_1 n \frac{\Delta h}{2}} + \cancel{E_1 n \frac{\Delta h}{2}} + E_1 t - \Delta w = 0$$

$$\boxed{E_1 t = 0}$$



② Perfect conductor  
 $\vec{D} \cdot \vec{E} = 0$   
 $\rho_v = 0$

No electric field exist within a conductor, under static conditions.

# CONDUCTOR - FREE SPACE BOUNDARY CONDITIONS.

$$\oint_C \vec{E} \cdot d\vec{l} = 0.$$

$$\int_a^b + \int_b^c + \int_c^d + \int_d^a = 0.$$

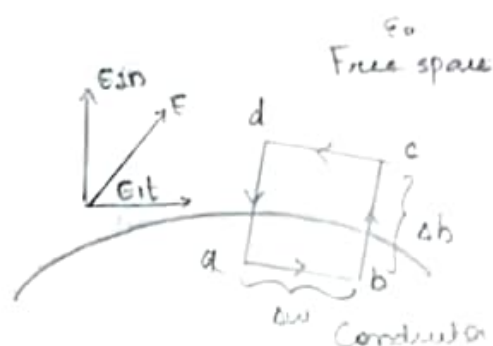
$$E_{1t} \Delta w + (E_{1n} \frac{\Delta b}{2} + E_{2n} \frac{\Delta b}{2}) +$$

$$E_{1t} \Delta w + (E_{1n} \frac{\Delta b}{2} + E_{2n} \frac{\Delta b}{2}) = 0.$$

$$E_{1t} \Delta w + 0 + 0 + 0 = 0.$$

$$E_{1t} \Delta w = 0$$

$$\boxed{E_{1t} = 0}$$



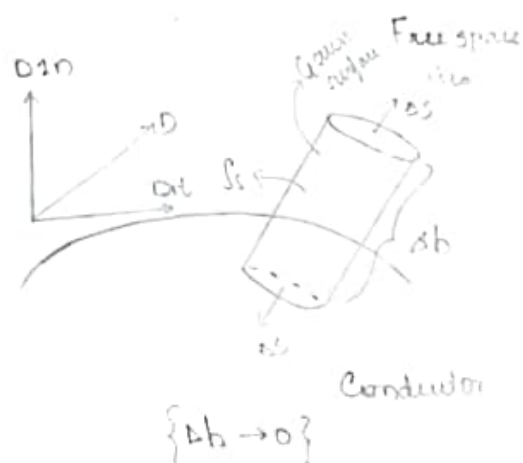
$$\oint_S \vec{D} \cdot d\vec{s} = Q_{enc}$$

$$\int_{top} + \int_{bottom} + \int_{side} = Q_{enc}$$

$$D_{1n} \Delta S + 0 + 0 = \rho_s \Delta S$$

$$D_{1n} \Delta S = \rho_s \Delta S$$

$$\boxed{D_{1n} = \rho_s}$$





Q. Find  $\nabla \cdot \vec{A}$  at  $(2, 30^\circ, 90^\circ)$  for the field:

$$\vec{A} = 0.2r^3 \sin^2 \alpha \hat{a}_r + 0.2r^3 \sin^2 \theta \hat{a}_\theta + 0.2r^3 \sin^2 \phi \hat{a}_\phi$$

Ans:  $\nabla \cdot \vec{A} = \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 A_r) + \frac{1}{r \sin \theta} \frac{\partial}{\partial \theta} (A_\theta \sin \theta) +$

$$\frac{1}{r \sin \theta} \frac{\partial A_\phi}{\partial \phi}$$

$$A_r = 0.2r^3 \sin^2 \theta$$

$$A_\theta = 0.2r^3 \sin^2 \theta$$

$$A_\phi = 0.2r^3 \sin^2 \theta$$

$$\nabla \cdot \vec{A} = \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 \times 0.2r^3 \sin^2 \theta) + \frac{1}{r \sin \theta} \frac{\partial}{\partial \theta} (0.2r^3 \sin^2 \theta) +$$

$$\frac{1}{r \sin \theta} \frac{\partial}{\partial \phi} (0.2r^3 \sin^2 \theta)$$

$$= \frac{1}{r^2} [0.2r^4 \times 3 \times \sin^2 \theta] + \frac{1}{r \sin \theta} [0.2r^3 \sin^2 \theta \times \cos \theta] +$$

$$\frac{1}{r \sin \theta} [0.2r^3 \sin^2 \theta]$$

Here,  $r = 2$ ,  $\theta = \frac{\pi}{6}$ ,  $\phi = \frac{\pi}{2}$ .

$$\nabla \cdot \vec{A} = 0.2r^2 \times 3 \times \sin^2 \theta + 0.2r^2 \times 3 \times \sin \theta \times \cos \theta + 0.2r^2 \sin \theta$$

$$= 0.2 \times 4 \times 3 \times \frac{\pi}{2} \times \sin^2 \frac{\pi}{6} + 0.2 \times 4 \times \frac{\pi}{2} \times 3 \times \sin \frac{\pi}{6} \times \cos \frac{\pi}{6} +$$

$$0.2 \times 4 \times \sin \frac{\pi}{6}$$

$$= 1.5704 + 1.6324 + 0.4$$

$$= \underline{\underline{3.6031}}$$

Q. A uniform line charge, infinite in extent, with  $\rho_L = 20 \text{ nC/m}$  lies along  $z$ -axis. Find  $\vec{E}$  at  $(6, 8, 3)$ ?

Ans:  $\vec{D} = \frac{\rho_L a_z}{2\pi r}$

$$\{\vec{D} = \epsilon \vec{E}\}$$

$$\vec{E} = \frac{\rho_L a_z}{2\pi \epsilon_0 r}$$

$$r = \sqrt{6^2 + 8^2}$$

$$= 10$$

$$= \frac{20 \times 10^{-9}}{6.28 \times 8.85 \times 10^{-12} \times 10} a_z$$

$$= \frac{20 \times 10^{-9}}{5.5578 \times 10^{-10}} a_z$$

$$= \underline{\underline{35.985 a_z \text{ V/m}}}$$