

The impact of cosmic variance on simulating weak lensing surveys

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25 August 2014

ABSTRACT

According to our current understanding, galaxy shapes and morphologies should depend on various factors such as the local environment. Realistic image simulations for calibration of weak lensing analysis methods that use training samples from the Hubble Space Telescope can therefore be affected by these trends, due to the limited volume of the universe that has been surveyed by Hubble. We will show how redshift slices in a volume-limited subsample of COSMOS can be classified as overdense or underdense (or neither), and how the statistical properties of various morphological parameters such as ellipticity, Sérsic n , Bulge-to-Total ratio and color differ in these bins. This study requires a careful distinction between environment effects from large-scale structure, which we do not wish to include in simulations, and general trends in the galaxy population with redshift. We conclude with some guidance for how upcoming surveys can use COSMOS data as the basis for weak lensing simulations without having their conclusions overly affected by cosmic variance. **I have some comments on this abstract, but prefer to do all revision of the abstract at the end once the paper is finalized.**

Key words: Gravitational lensing: weak, Cosmology: Large-scale structure of Universe, Galaxies: evolution ? **Are these from the MNRAS list?**

1 INTRODUCTION

DES is far from the only survey; since I'm in HSC and the work was supported by Euclid project, it seems strange to only mention DES. I recommend taking one of my papers (e.g., the one on image simulations from 2012) and lifting the entire list of references to surveys directly from there. Major surveys such as the Dark Energy Survey (DES) are planned to be carried out over the next decade or two to make measurements that constrain the nature of dark energy. Weak gravitational lensing probes the nature of dark energy through both the growth of structure in the universe with redshift and the dependence of distances on the expansion rate or the Hubble parameter. However, most of the matter in the universe is made up of dark matter which cannot be detected by most of the standard astronomical/cosmological probes. Weak lensing happens to be the cleanest probe for this purpose. As light from the background sources passes through the foreground cluster of galaxies and dark matter profiles, the shapes get distorted. Weak lensing also causes magnification in addition to introducing a shear.

For the upcoming surveys to achieve their promise, the systematic error budget must be below the statistical error budget. This would mean that the systematic errors need to be reduced from their current typical level. One method that is commonly used to test for the presence of systematics is image simulation. Realistic simulations use samples based on images taken from the Hubble Space Telescope (HST) such as COSMOS or UDF (Ultra Deep Field), which are deep (high redshift) surveys taken over small areas of the sky. The goal of this work is to understand the systematic uncertainties that might arise due to finite size of these areas when dividing the sample into redshift slices to test for shear systematics as a function of redshift.

Previous studies have shown that any shape measurement method that uses second moments cannot be independent of the morphology and substructure of the galaxies involved. (**rephrased from GREAT3**). Hence, it is important to check any method that is proposed to measure the shape of the galaxies to check against a simulation that generates images as realistic as possible. With this in mind, the GREAT3 challenge (Mandelbaum et al. 2014) was held with the aim to test and facilitate methods for measuring weak gravitational lensing by analyzing astronomical im-

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ages. Software packages like GalSim¹ can generate galaxy images fairly accurately, given realistic input parameters. The training dataset used in the challenge comes from the COSMOS survey, which we describe in sufficient detail in §2. While providing a good quality (high S/N) images to be trained over, there could be some bias introduced due to the survey being a narrow one. It is known **cite refs.** that the shape and morphology of galaxies depend very much on their local environment. Hence, local overdensities or underdensities observed in the COSMOS' field of view may cause the properties of the galaxies to be associated with the redshift, thus unable to distinguish between redshift evolution and large scale structure.

In this paper, we show that these large scale structure effects mentioned are indeed present in the data. The rest of the paper is structured as follows. In §3.1, we explain how we identify overdensities and underdensities, volume-limit our data in §3.2 and in §3.3, we explain how we obtain axis-ratios from the images of the galaxies. Finally, we present our results in §4 and end with conclusions in §5.

2 DATA

COSMOS (Scoville et al. 2007; Koekemoer et al. 2007; Leauthaud et al. 2007) is a flux-limited, narrow deep field survey covering a contiguous area of 1.64 deg^2 of sky, with images taken using the Advanced Camera for Surveys (ACS) Wide Field Channel (WFC) in the Hubble Space Telescope (HST). Precise shape measurements, when compared to ground-based surveys, can be made since the full width half-maximum (FWHM) of the point-spread function (PSF) is $0.12''$. High resolution images taken through the wide F814W filter (broad I), after correcting for the PSF, have led to collection of postage stamp images² for the GREAT3 challenge (Mandelbaum et al. 2014), using the procedure described in Mandelbaum et al. (2011). Parametric models to most of these galaxies including Sérsic n profile fits, 2 component bulge+disk fits, axis ratios etc were done for GREAT3 challenge (Mandelbaum et al. 2014) using the method described in Lackner and Gunn (2012). This is described briefly in Sec. 3.3. In addition to the ACS/WFC (F814W) imaging, the COSMOS field has also been imaged by Subaru Suprime-Cam ($B_j, V_j, g^+, r^+, i^+, z^+$, NB816), the Canada-French Hawaii Telescope (CFHT; u^*, i^*) and the KPNO/CTIO (K_s).

Photometric redshifts were determined by Ilbert et al. (2009). The accuracy of photometric redshifts for F814W ≤ 22.5 is $\sigma_{\Delta z} = 0.007(1 + z_s)$ and at fainter magnitudes, i.e. F814W ≤ 24 , $\sigma_{\Delta z} = 0.012(1 + z_s)$. The photometric redshift values become noisy beyond z of 1 for our purposes and the various fits to the galaxies are also not very reliable beyond the apparent magnitude (m) value of 23.5. **Make a clear and explicit statement that we always require $z < 1$, we require the MAG_AUTO in F814W to be < 23.5 , and we use the parent sample that goes to 25.2 for stellar masses and luminosities only.**

Stellar mass estimates were obtained (Leauthaud et al.

2010) using the Bayesian code described in Bundy et al. (2006). Basically, a grid of models that vary in age, star formation history, dust content and metallicity are constructed to which the observed galaxy's spectral energy distribution (SED) and the photometric redshift are referenced to. At each grid point, the probability that the SED fits the model are calculated and by marginalizing over all the parameters in the grid, the stellar mass probability distribution is obtained. The median of this distribution is taken as the stellar mass estimate. A Chabrier IMF is assumed.

We apply the following set of initial cuts to the data, all of which is defined in Leauthaud et al. (2007):

(i) **MU_CLASS=1**: Unlike SExtractor's stellar index **CLASS_STAR** which is continuous, the **MU_MAX** method that compares the peak surface brightness to the background level, provides a clear differentiation of galaxies (**MU_CLASS=1**) from stars (**MU_CLASS=2**) and other spurious objects like cosmic rays (**MU_CLASS=3**)

(ii) **CLEAN=1**: Objects near bright stars or those containing saturated pixels were removed and the rest classify as *clean*.

(iii) **GOOD_ZPHOT_SOURCE=1**: This cut requires that there be a good photometric redshift, which typically is equivalent to requiring that the galaxy not be located within the masked regions of the Subaru $BVIz$ imaging used for photometric redshifts, and that it have a successful match against an object in the Subaru imaging **copied from Mandelbaum 2011**.

The cuts you just gave are really basic or fundamental ones. Those should go before any discussion of more complicated derived quantities like redshifts, stellar mass, Sérsic fits, etc.

We will analyze how the galaxy morphology and intrinsic shape depends on the environment in which they reside.

3 METHODS

In order to study the variation in the intrinsic ellipticity distribution and other morphological indicators with environment, there are three main steps to be carried out.

(i) To identify overdense and underdense environments in our survey from the redshift distribution of galaxies (§3.1),

(ii) To estimate the galaxy axis ratios and other morphological indicators such as Sérsic index and Bulge-to-Total ratios (§3.3),

(iii) To limit the sample such that Malmquist bias is minimized before comparing galaxies in different redshift slices (§3.2)

3.1 Finding overdensities

Weak lensing simulations will use galaxies selected in 2D redshift slices and we split our sample in a similar way to assign overdensities.

For our sample of galaxies, upto $z = 1.2$, we fit parametric models to the histogram of photometric redshift in order to assign values of overdensity. Figures show Chi-Squared and the functional form given in Baugh and Efstathiou (1993) fits to the histograms. For low z , the volume is too small to rely the overdensity values from our model fits and hence the fit is made for $z \geq 0.3$. Our binning width, Δz

¹ <https://github.com/GalSim-developers/GalSim>

² <http://irsa.ipac.caltech.edu/data/COSMOS/images/galaxy-postage-stamps/>

must be small enough so as to be able to identify localized overdensities/underdensities but large enough to not let our conclusions be affected by the errors in photometric redshift. We choose our bins to be 0.05 wide starting from $z = 0.3$. **Were you going to use other fitting functions as a test of the method? If so, you should introduce them here.**

The curves are normalized such that the area under the curves is equal to the number of galaxies considered. **Fig 3 - Histogram with fits** .

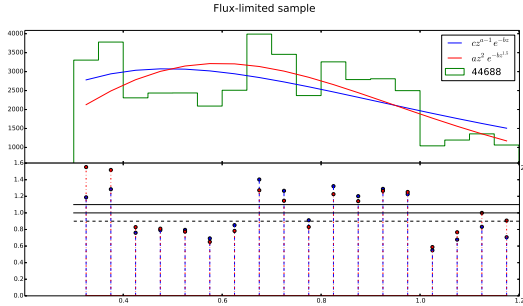


Figure 1.

Overdensity in a redshift bin is defined as the ratio of the difference between the value given by the histogram and the value predicted by the model to the latter: $\delta = (N - N_{\text{mod}})/N_{\text{mod}}$. We leave a 10% margin of safety i.e. iff δ from both the fits is greater than 0.1, then we call that redshift bin as an overdense region whereas if δ from both the fits is less than -0.1 , then we call that as an underdense region. Redshift bins with $-0.1 < \delta < 0.1$ are neither too overdense nor too underdense and we call them as ‘unclassified regions’ and discard them from our analysis. **Fig 4 - overdensities vs redshift bins. Let’s not discard them. We should just treat them as “environmentally neutral.” You do after all show them on your plots.**

We thus identify the regions $z = 0.30 - 0.40, 0.65 - 0.75, 0.80 - 0.85$ as overdense, $z = 0.55 - 0.65, 0.75 - 0.80$ as underdense and $z = 0.40 - 0.55$ as unclassified. The classification of overdensities and underdensities agrees with the work done by Kovač et al. (2010) on the zCOSMOS field using spectra and 3D density field reconstruction, except for the last two high redshift bins. We believe that this disagreement is due to the errors in our photometric redshifts and the overdensity reported in $z = 0.875 - 1$ slice is observed by us in $z = 0.80 - 0.85$ slice.

We see that the region between $0.85 < z < 1.0$ is neither overdense nor underdense according to our model and hence is not going to be considered for further analysis.

We use therefore relax our luminosity cut so that the sample is volume-limited *not* until $z = 1$ but until $z = 0.85$. We choose to impose the cut at $M_I = -20.8$, with 95.3% completeness. This increases the sample size significantly, from 7,418 galaxies to 11,169.

I think “arbitrary” is not the right word to use in this paragraph. It’s obviously well-motivated in a general sense, and the question is whether it’s robust or not. Comparison

with a more robust 3d method is what we use as validation to say that it’s robust enough for our purposes. Done.

3.2 Volume limiting

COSMOS is a flux-limited survey and is therefore affected by Malmquist bias, i.e. at higher redshifts, the brighter galaxies are preferentially observed. Our analysis involves comparing galaxies in different redshift slices to find significant differences in morphology, if any, so with a flux-limited sample, we would be comparing only the bright galaxies at high redshifts with bright and faint galaxies at low redshifts. For a fair comparison, we must restrict ourselves to bright galaxies at all redshifts and this is achieved by volume-limiting the sample. We generate, using the method below, a volume-limited sample that is complete upto $z = 1$ by applying a cut on luminosity such that only galaxies intrinsically brighter than a certain threshold is considered. This threshold is set on K -corrected i -band magnitudes, denoted as M_I . Since the parent sample contains fainter galaxies, upto $F814W=25.2$, we compare the distribution of the $F814W=23.5$ sample with the samples containing fainter galaxies for high redshift bins, to see where the sample is no longer complete.

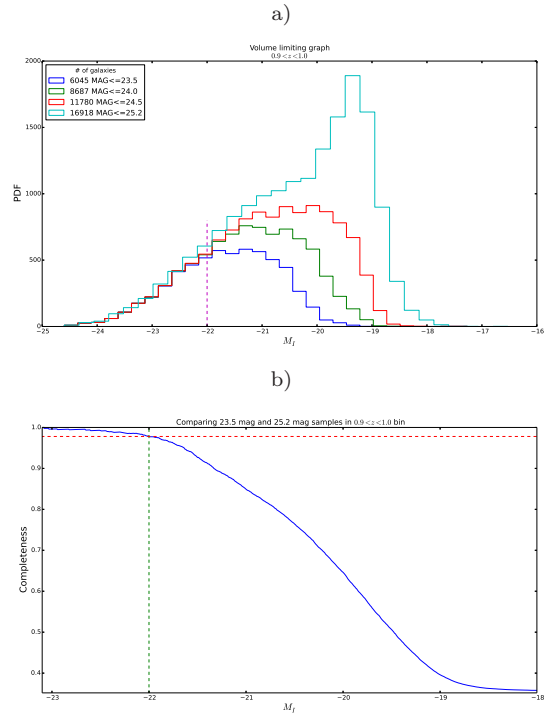


Figure 2. ...

At $M_I \sim -22.0$, we see the sample is beginning to be biased in the $0.9 < z < 1.0$ bin due to the flux limit. We obtain 97.84% completeness in this bin for M_I , where completeness is defined as the ratio between the number of galaxies in $F814W \leq 23.5$ and in $F814W \leq 25.2$ samples. Thus the sample $z \leq 1$ and $M_I < -22$ is at least 97.84% complete. Fig. 5 shows that at $M_I < -22$, we are not affected by flux limit yet.

However, previous studies [cite](#) have shown that absolute luminosities evolve with redshift. Thus, we must also let the luminosity cut evolve with the redshift. There has been no published literature on the LF for the *I*-band, particularly for the COSMOS survey. We used the published result (Faber et al. 2007) for the ‘rate’ of evolution of B-band magnitudes for DEEP2 and COMBO-17 surveys, which is $\Delta M_B^* \sim -1.23$ mag per unit redshift, for both blue and red galaxies combined together. Typically, estimates of evolution in K-band are lesser than the estimates of evolution in B and V bands. Assuming that the evolution is a smooth function of the wavelength, the evolution in I-band is expected to be lesser than that of the B-band. Thus, by considering no evolution and an upper bound on the evolution, we can interpolate what the results would be like for the true *I*-band evolution.

Alternatively, one could get around the problem of considering redshift evolution by imposing cuts on stellar mass instead of absolute luminosity in a particular band. In Fig. 3, we show the stellar mass function (SMF) of our sample for various F814W flux limit. Tomczak et al. (2014) report the SMFs for the ZFOURGE survey, which includes COSMOS. They considered for this work, a single stellar population following a Chabrier [cite](#) IMF. We plot their SMF for *all* in 3 for reference. Their SMF is higher than ours since they reach K_s -band 5σ depth of 24.9. **What exactly does this mean?** As in the case of M_I s, we compare the stellar mass in the $F814W \leq 23.5$ sample with that of the $F814W \leq 25.2$ sample. The sample $\log(M/M_\odot) > 10.15$ is about 95% complete in the redshift bin $[0.75 - 0.85]$ and has 10,341 galaxies in total.

There is a minor pitfall with this method. We analyse only those galaxies for which postage stamp images exist. Postage stamps may not exist either because the image occurs near the edge of the CCD chips or because the image is too large to fit in a bounding box. [Rachel, please correct my terminology here.](#) While the former is purely random and doesn’t introduce bias, the latter is not. Typically galaxies that are nearby and intrinsically very bright do not have postage stamps associated with them and this effect is dominant at lower redshifts. Thus, we use the full $F814W \leq 23.5$ sample, irrespective of the existence of postage stamps, and believe that our conclusions are not affected by this bias.

Not surprisingly, we see the overdense and underdense regions obtained by the above method have local average number density higher than the global average number density after volume limiting. And the environmentally neutral regions fall below or at least lie close to the 10% margin.

Redshift	# of galaxies	Environment
0.30-0.35	726	Overdense
0.35-0.40	1000	Overdense
0.55-0.60	727	Underdense
0.60-0.65	1070	Underdense
0.65-0.70	2089	Overdense
0.70-0.75	1970	Overdense
0.75-0.80	1159	Underdense
0.80-0.85	2428	Overdense

In the following section, we will compare and analyze the distribution of properties of the galaxies residing in the overdense regions.

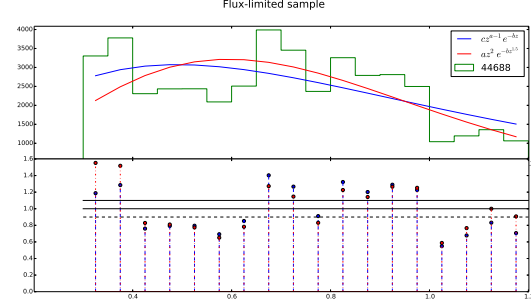


Figure 4.

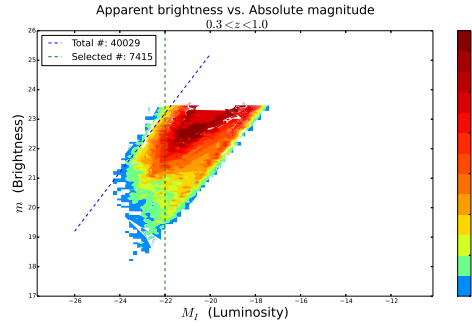


Figure 5. 2-D histogram of galaxies in apparent magnitude (m) and absolute magnitude (M_I) space.

Talk more about environments - Nature and nurture?

3.3 Finding fit parameters

Should explain in brief Claire’s work.

If galaxies have elliptical isophotes, its shape and size could be defined by the axis ratio and the area enclosed by a boundary isophote. However, in real galaxy images, the boundary may not be well defined and the shape may not be well approximated by an ellipse. Given the surface brightness (intensity) $I(\theta)$ of a galaxy image at every angular position θ , we can define the tensor of second brightness moments as

$$Q_{ij} = \frac{\int d^2\theta q_I(\theta_i - \bar{\theta}_i)(\theta_j - \bar{\theta}_j)}{\int d^2\theta q_I}, \quad (1)$$

where q_I is a suitable chosen kernel and $\bar{\theta}$ is the average angular coordinate, weighted by the same kernel function and $i, j \in \{1, 2\}$ are the components of $\vec{\theta}$. It is common to find in literature two kinds of *complex* ellipticities that quantify the shape of the galaxies - χ and ϵ defined as

$$\chi = \frac{Q_{11} - Q_{22} + 2iQ_{12}}{Q_{11} + Q_{22}} \quad (2)$$

$$\epsilon = \frac{Q_{11} - Q_{22} + 2iQ_{12}}{Q_{11} + Q_{22} + 2(Q_{11}Q_{22} - Q_{12}^2)^{1/2}} \quad (3)$$

If the image has elliptical isophotes with axis ratio q , then

$$|\chi| = \frac{1 - q^2}{1 + q^2}, \quad |\epsilon| = \frac{1 - q}{1 + q}. \quad (4)$$

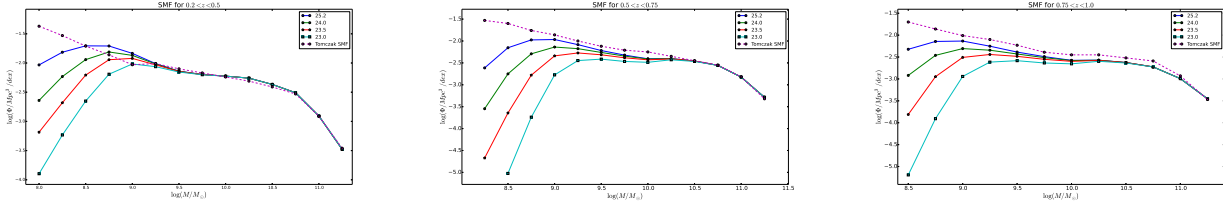


Figure 3. SMF plots

Conversely, one could compute χ or ϵ for a galaxy, and obtain an effective, azimuthally averaged (?) axis-ratio q . No, that isn't possible because of the PSF. One would have to use a PSF correction scheme. Recommend removing much of the above discussion. Should I remove these even if we plan to use the ellipticities from the second moments? However, our method of finding the value of q is little more complicated. Briefly, we fit two models to each image:

- (i) a Sérsic profile given by the expression

$$S = \Sigma_{1/2} \exp \left(-k(R/R_{\text{eff}})^{1/n} - 1 \right), \quad (5)$$

- (ii) two Sérsic component fits: de Vaucouleurs bulge ($n = 4$) + exponential disc profile ($n = 1$),

where

$$R^2 = ((x - x_0) \cos \Phi + (y - y_0) \sin \Phi)^2 + ((y - y_0) \cos \Phi - (x - x_0) \sin \Phi)^2 / q^2,$$

R_{eff} is the half-light radius of the profile, $\Sigma_{1/2}$ is the surface brightness at $R = R_{\text{eff}}$, (x_0, y_0) is the centroid of the image, Φ is the profile rotation angle, n is the Sérsic index, k is a n -dependent normalization factor and q is the axis ratio of the elliptical isophotes. Thus, the Sérsic profile has 7 free parameters. The bulge+disk model has 10 free parameters since the Sérsic indices are fixed as 1 and 4 and both the profiles are required to have the same centroid (x_0, y_0) . Best-fit parameters were found, by minimizing the weighted sum of the difference between the image and the PSF-convolved mode using Levenberg-Marquardt minimization, `mpfit2fun` in IDL [cite](#). More details about the fit can be found in Lackner and Gunn (2012).

4 RESULTS

The distribution of the axis ratios that are obtained by the above manner are compared between redshift bins. We use two statistical tests namely the Kolmogorov-Smirnov test and Anderson-Darling test to compare distributions.

We first compare the distribution of the axis ratio in *all* overdense bins and *all* underdense bins. The p-value from both the KS and AD tests are below 0.05, so we reject the ‘null hypothesis’ that the overdense and underdense regions have same axis ratio distributions at 95% significance level. Next, we will compare distributions between two overdense / underdense regions, where we expect to find similarity, and between an overdense and underdense regions, where we expect the distributions to differ. Figures 9 and 10 show that the distributions are indeed similar when the environments

are similar and different when the environments are different, confirming our expectation. The cumulative distinction functions are also plotted alongside in order to be able to visualize the ‘distance’ between the distributions.

It is useful to consider the root mean square ellipticity since it can characterize the shape of the population/sample of galaxies.

The RMS ellipticities of galaxies in each redshift bin are shown in Figure 12. As one can see, the underdense regions have significantly higher values for RMS ellipticities when compared to the overdense regions. Note in particular that we’ve been able to capture the narrow overdense bin $0. \leq z < 0.80$. There is (almost) no redshift dependence in the figure. The dependence on the local environment is consistent with **give physical argument..** From Figs. ?? and ??, the region $0.4 \leq z < 0.55$ show signs of being marginally underdense but have low RMS ellipticities too that agree with the rest of the overdense regions.

When the evolution is taken into account, there is a systematic increase in the ellipticity at lower redshifts.

Figures ?? show the median of the Sérsic index with and without taking into account of the luminosity evolution. Median values obtained using stellar mass selected samples are plotted alongside for reference. We chose median over mean since Sérsic n greater than 6 are truncated to 6. This affects the mean statistic more than the median.

5 CONCLUSION

We have shown that

ACKNOWLEDGMENTS

AK and RM acknowledge the support of NASA ROSES 12-EUCLID12-0004, and program HST-AR-12857.01-A, provided by NASA through a grant from the Space Telescope Science Institute, which is operated by the Association of Universities for Research in Astronomy, Incorporated, under NASA contract NAS5-26555. AK and RM thank Alexie Leauthaud for many useful discussions.

REFERENCES

- C. M. Baugh and G. Efstathiou. The Three-Dimensional Power Spectrum Measured from the APM Galaxy Survey - Part One - Use of the Angular Correlation Function. *MNRAS*, 265:145, November 1993.

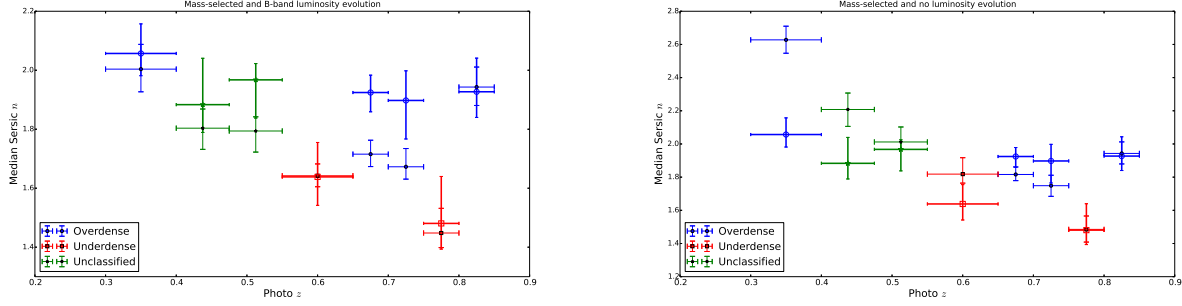


Figure 6. something

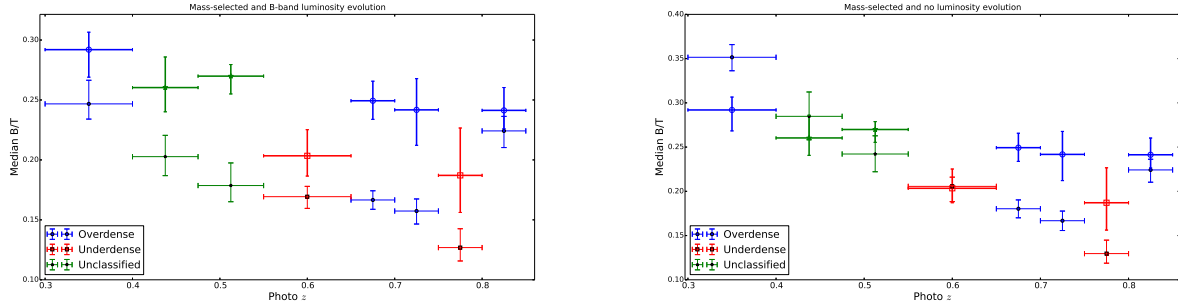


Figure 7. something else

K. Bundy, R. S. Ellis, C. J. Conselice, J. E. Taylor, M. C. Cooper, C. N. A. Willmer, B. J. Weiner, A. L. Coil, K. G. Noeske, and P. R. M. Eisenhardt. The Mass Assembly History of Field Galaxies: Detection of an Evolving Mass Limit for Star-Forming Galaxies. *ApJ*, 651:120–141, November 2006. doi: 10.1086/507456.

S. M. Faber, C. N. A. Willmer, C. Wolf, D. C. Koo, B. J. Weiner, J. A. Newman, M. Im, A. L. Coil, C. Conroy, M. C. Cooper, M. Davis, D. P. Finkbeiner, B. F. Gerke, K. Gebhardt, E. J. Groth, P. Guhathakurta, J. Harker, N. Kaiser, S. Kassin, M. Kleinheinrich, N. P. Konidakis, R. G. Kron, L. Lin, G. Luppino, D. S. Madgwick, K. Meisenheimer, K. G. Noeske, A. C. Phillips, V. L. Sarajedini, R. P. Schiavon, L. Simard, A. S. Szalay, N. P. Vogt, and R. Yan. Galaxy Luminosity Functions to $z \sim 1$ from DEEP2 and COMBO-17: Implications for Red Galaxy Formation. *ApJ*, 665:265–294, August 2007. doi: 10.1086/519294.

O. Ilbert, P. Capak, M. Salvato, H. Aussel, H. J. McCracken, D. B. Sanders, N. Scoville, J. Kartaltepe, S. Arnouts, E. Le Floc'h, B. Mobasher, Y. Taniguchi, F. Lamareille, A. Leauthaud, S. Sasaki, D. Thompson, M. Zamojski, G. Zamorani, S. Bardelli, M. Bolzonella, A. Bongiorno, M. Brusa, K. I. Caputi, C. M. Carollo, T. Contini, R. Cook, G. Coppia, O. Cucciati, S. de la Torre, L. de Ravel, P. Franzetti, B. Garilli, G. Hasinger, A. Iovino, P. Kampeczyk, J.-P. Kneib, C. Knobel, K. Kovac, J. F. Le Borgne, V. Le Brun, O. L. Fèvre, S. Lilly, D. Looper, C. Maier, V. Mainieri, Y. Mellier, M. Mignoli, T. Murayama, R. Pellò, Y. Peng, E. Pérez-Montero, A. Renzini, E. Ricciardelli, D. Schiminovich, M. Scodeggio, Y. Shioya, J. Silverman, J. Surace, M. Tanaka,

L. Tasca, L. Tresse, D. Vergani, and E. Zucca. Cosmos Photometric Redshifts with 30-Bands for 2-deg^2 . *ApJ*, 690:1236–1249, January 2009. doi: 10.1088/0004-637X/690/2/1236.

A. M. Koekemoer, H. Aussel, D. Calzetti, P. Capak, M. Giavalisco, J.-P. Kneib, A. Leauthaud, O. Le Fèvre, H. J. McCracken, R. Massey, B. Mobasher, J. Rhodes, N. Scoville, and P. L. Shopbell. The COSMOS Survey: Hubble Space Telescope Advanced Camera for Surveys Observations and Data Processing. *ApJS*, 172:196–202, September 2007. doi: 10.1086/520086.

K. Kovač, S. J. Lilly, O. Cucciati, C. Porciani, A. Iovino, G. Zamorani, P. Oesch, M. Bolzonella, C. Knobel, A. Finoguenov, Y. Peng, C. M. Carollo, L. Pozzetti, K. Caputi, J. D. Silverman, L. A. M. Tasca, M. Scodeggio, D. Vergani, N. Z. Scoville, P. Capak, T. Contini, J.-P. Kneib, O. Le Fèvre, V. Mainieri, A. Renzini, S. Bardelli, A. Bongiorno, G. Coppia, S. de la Torre, L. de Ravel, P. Franzetti, B. Garilli, L. Guzzo, P. Kampeczyk, F. Lamareille, J.-F. Le Borgne, V. Le Brun, C. Maier, M. Mignoli, R. Pello, E. Perez Montero, E. Ricciardelli, M. Tanaka, L. Tresse, E. Zucca, U. Abbas, D. Bottini, A. Cappi, P. Cassata, A. Cimatti, M. Fumana, A. M. Koekemoer, D. Maccagni, C. Marinoni, H. J. McCracken, P. Memeo, B. Meneux, and R. Scaramella. The Density Field of the 10k zCOSMOS Galaxies. *ApJ*, 708:505–533, January 2010. doi: 10.1088/0004-637X/708/1/505.

C. N. Lackner and J. E. Gunn. Astrophysically motivated bulge-disc decompositions of Sloan Digital Sky Survey galaxies. *MNRAS*, 421:2277–2302, April 2012. doi: 10.1111/j.1365-2966.2012.20450.x.

A. Leauthaud, R. Massey, J.-P. Kneib, J. Rhodes, D. E.

- Johnston, P. Capak, C. Heymans, R. S. Ellis, A. M. Koekemoer, O. Le Fèvre, Y. Mellier, A. Réfrégier, A. C. Robin, N. Scoville, L. Tasca, J. E. Taylor, and L. Van Waerbeke. Weak Gravitational Lensing with COSMOS: Galaxy Selection and Shape Measurements. *ApJS*, 172: 219–238, September 2007. doi: 10.1086/516598.
- A. Leauthaud, A. Finoguenov, J.-P. Kneib, J. E. Taylor, R. Massey, J. Rhodes, O. Ilbert, K. Bundy, J. Tinker, M. R. George, P. Capak, A. M. Koekemoer, D. E. Johnston, Y.-Y. Zhang, N. Cappelluti, R. S. Ellis, M. Elvis, S. Giodini, C. Heymans, O. Le Fèvre, S. Lilly, H. J. McCracken, Y. Mellier, A. Réfrégier, M. Salvato, N. Scoville, G. Smoot, M. Tanaka, L. Van Waerbeke, and M. Wolk. A Weak Lensing Study of X-ray Groups in the Cosmos Survey: Form and Evolution of the Mass-Luminosity Relation. *ApJ*, 709:97–114, January 2010. doi: 10.1088/0004-637X/709/1/97.
- R. Mandelbaum, C. M. Hirata, A. Leauthaud, R. J. Massey, and J. Rhodes. SHERA: SHEar Reconvolution Analysis, August 2011. Astrophysics Source Code Library.
- R. Mandelbaum, B. Rowe, J. Bosch, C. Chang, F. Courbin, M. Gill, M. Jarvis, A. Kannawadi, T. Kacprzak, C. Lackner, A. Leauthaud, H. Miyatake, R. Nakajima, J. Rhodes, M. Simet, J. Zuntz, B. Armstrong, S. Bridle, J. Coupon, J. P. Dietrich, M. Gentile, C. Heymans, A. S. Jurling, S. M. Kent, D. Kirkby, D. Margala, R. Massey, P. Melchior, J. Peterson, A. Roodman, and T. Schrabback. The Third Gravitational Lensing Accuracy Testing (GREAT3) Challenge Handbook. *ApJS*, 212:5, May 2014. doi: 10.1088/0067-0049/212/1/5.
- N. Scoville, H. Aussel, M. Brusa, P. Capak, C. M. Carollo, M. Elvis, M. Giavalisco, L. Guzzo, G. Hasinger, C. Impey, J.-P. Kneib, O. LeFevre, S. J. Lilly, B. Mobasher, A. Renzini, R. M. Rich, D. B. Sanders, E. Schinnerer, D. Schminovich, P. Shopbell, Y. Taniguchi, and N. D. Tyson. The Cosmic Evolution Survey (COSMOS): Overview. *ApJS*, 172:1–8, September 2007. doi: 10.1086/516585.
- A. R. Tomczak, R. F. Quadri, K.-V. H. Tran, I. Labbé, C. M. S. Straatman, C. Papovich, K. Glazebrook, R. Allen, G. B. Brammer, G. G. Kacprzak, L. Kawinwanichakij, D. D. Kelson, P. J. McCarthy, N. Mehrrens, A. J. Monson, S. E. Persson, L. R. Spitler, V. Tilvi, and P. van Dokkum. Galaxy Stellar Mass Functions from ZFOURGE/CANDELS: An Excess of Low-mass Galaxies since $z = 2$ and the Rapid Buildup of Quiescent Galaxies. *ApJ*, 783:85, March 2014. doi: 10.1088/0004-637X/783/2/85.

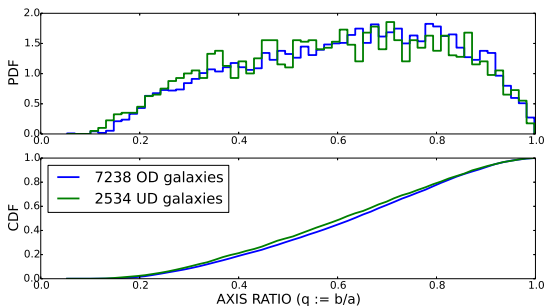


Figure 8. KS p-value = 0.004 & AD p-value = 0.001

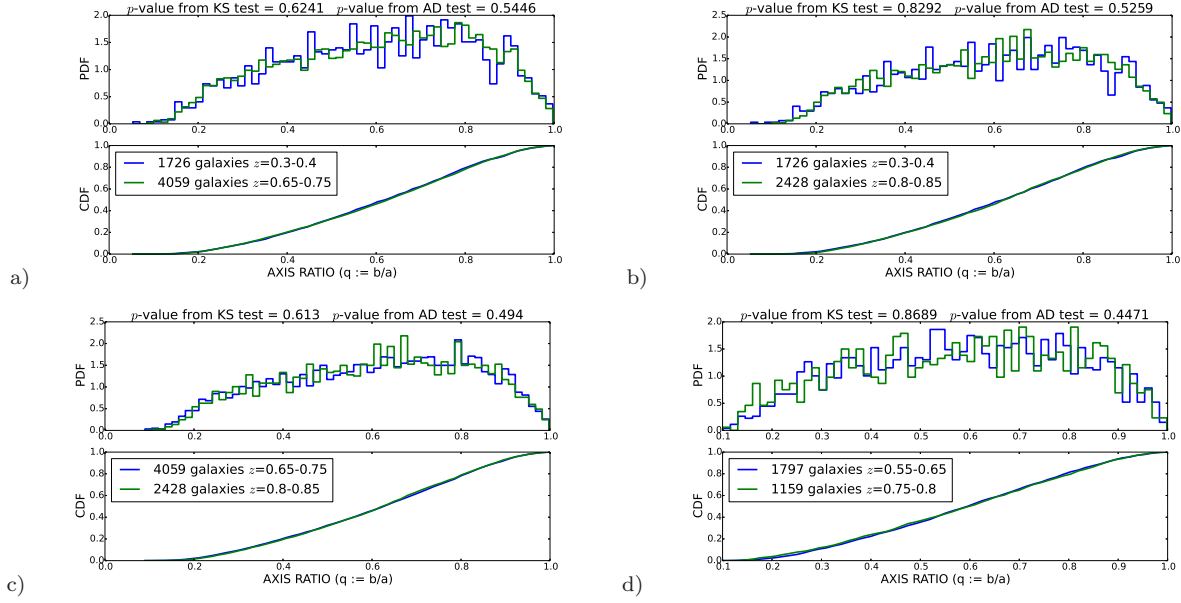


Figure 9. Comparison of axis ratios of galaxies in similar environments. p -values from the KS and AD test are (will be) given in the plot.

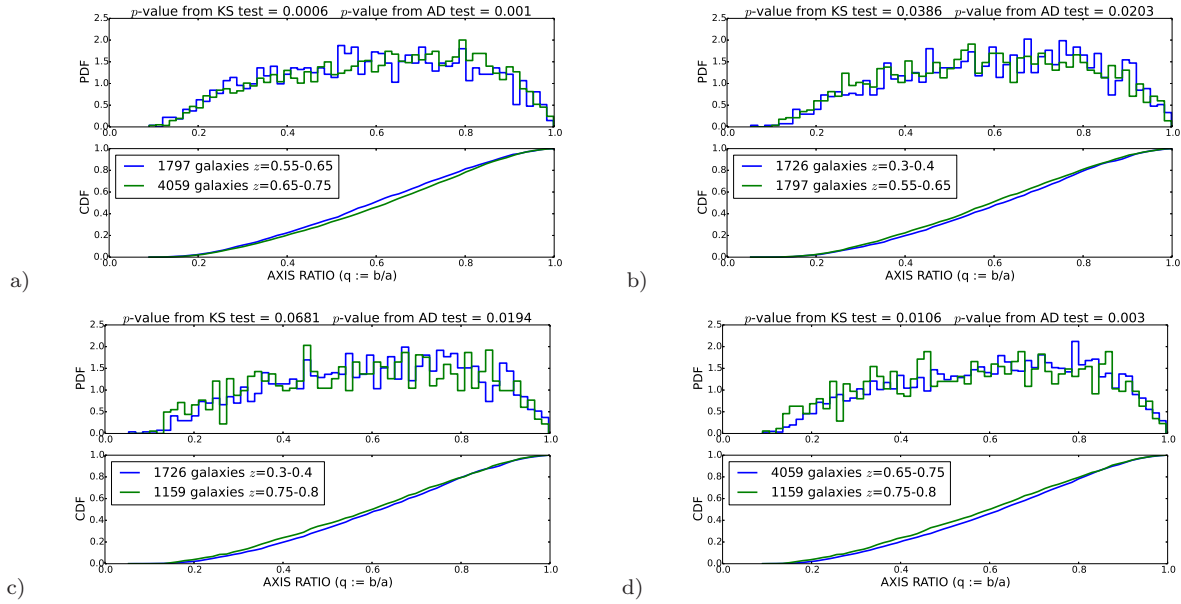


Figure 10. Comparison of axis ratios of galaxies in contrasting environments. p -values from the KS and AD test are (will be) given in the plot.

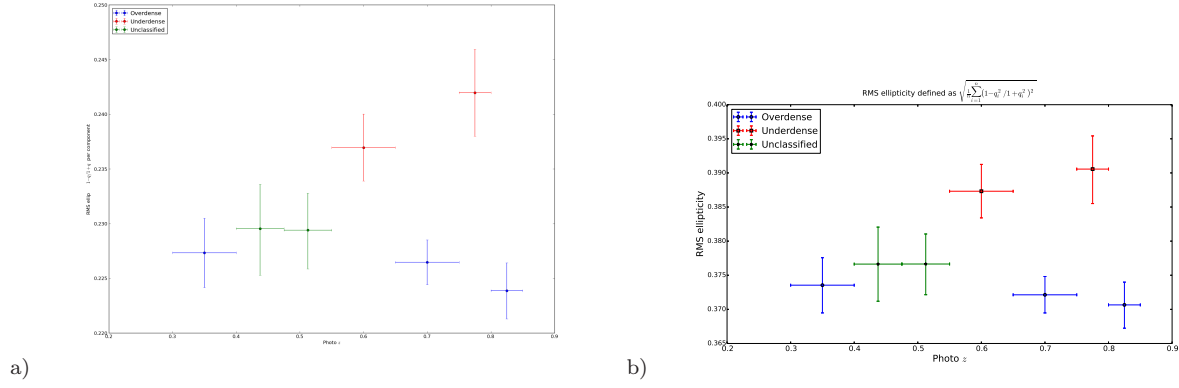


Figure 11. Left: RMS ellipticities with ellipticity defined as $\frac{1-q}{1+q}$. Right: RMS ellipticities with ellipticity defined as $\frac{1-q^2}{1+q^2}$. The horizontal errorbars simply correspond to the binwidth while the vertical ones are 1σ errorbars obtained by bootstrapping.

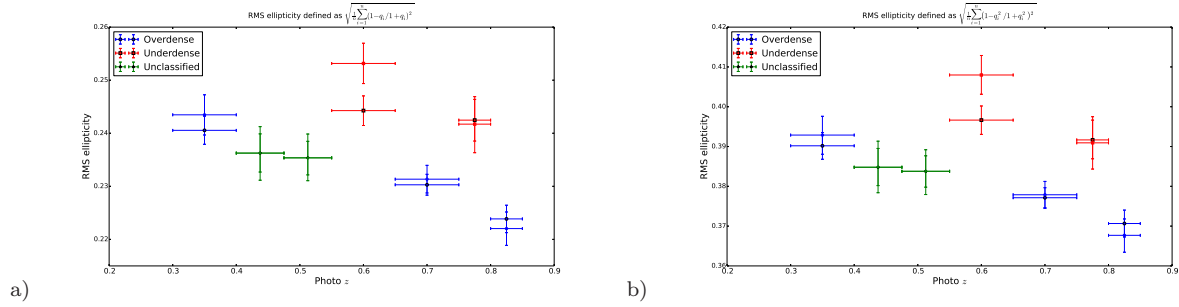


Figure 12. Left: RMS ellipticities with ellipticity defined as $\frac{1-q}{1+q}$. Right: RMS ellipticities with ellipticity defined as $\frac{1-q^2}{1+q^2}$. The horizontal errorbars simply correspond to the binwidth while the vertical ones are 1σ errorbars obtained by bootstrapping. The solid points correspond to the sample where the luminosity cut evolves by -1.23 per unit redshift and the unfilled points correspond to the sample obtained from stellar mass cuts.