

PROBLEM 5

$$a) \delta \Gamma_{\nu\lambda}^{\mu} = \Gamma_{\nu\lambda}^{\mu} - \bar{\Gamma}_{\nu\lambda}^{\mu}$$

$$= \frac{1}{2} g^{\mu\sigma} [\partial_{\lambda} g_{\sigma\nu} + \partial_{\nu} g_{\sigma\lambda} - \partial_{\sigma} g_{\lambda\nu}]$$

$$- \frac{1}{2} \bar{g}^{\mu\sigma} [\partial_{\lambda} \bar{g}_{\sigma\nu} + \partial_{\nu} \bar{g}_{\sigma\lambda} - \partial_{\sigma} \bar{g}_{\lambda\nu}]$$

$$= \frac{1}{2} \bar{g}^{\mu\sigma} [\partial_{\lambda} \bar{g}_{\sigma\nu} + \partial_{\nu} \bar{g}_{\sigma\lambda} - \partial_{\sigma} \bar{g}_{\lambda\nu}]$$

$$+ \frac{1}{2} \bar{g}^{\mu\sigma} [\partial_{\lambda} h_{\sigma\nu} + \partial_{\nu} h_{\sigma\lambda} - \partial_{\sigma} h_{\lambda\nu}]$$

$$+ \frac{1}{2} h^{\mu\sigma} [\partial_{\lambda} h_{\sigma\nu} + \partial_{\nu} h_{\sigma\lambda} - \partial_{\sigma} h_{\lambda\nu}]$$

SECOND ORDER.

$$+ \frac{1}{2} h^{\mu\sigma} [\partial_{\lambda} \bar{g}_{\sigma\nu} + \partial_{\nu} \bar{g}_{\sigma\lambda} - \partial_{\sigma} \bar{g}_{\lambda\nu}]$$

$$- \frac{1}{2} \bar{g}^{\mu\nu} [\partial_{\lambda} \bar{g}_{\sigma\nu} + \partial_{\nu} \bar{g}_{\sigma\lambda} - \partial_{\sigma} \bar{g}_{\lambda\nu}]$$

$$= \frac{1}{2} \bar{g}^{\mu\sigma} [\partial_{\lambda} h_{\sigma\nu} + \partial_{\nu} h_{\sigma\lambda} - \partial_{\sigma} h_{\lambda\nu}]$$

$$+ \frac{1}{2} h^{\mu\sigma} [\partial_{\lambda} \bar{g}_{\sigma\nu} + \partial_{\nu} \bar{g}_{\sigma\lambda} - \partial_{\sigma} \bar{g}_{\lambda\nu}]$$

$$= \left(\frac{1}{2} \bar{g}^{\mu\sigma} h_{\sigma\alpha} \bar{g}^{\alpha\sigma} \right) [\dots] = -\frac{1}{2} \bar{g}^{\mu\sigma} h_{\sigma\alpha} \cdot 2 \bar{\Gamma}_{\nu\lambda}^{\alpha}$$

$$= \frac{1}{2} \bar{g}^{\mu\sigma} [\partial_{\lambda} h_{\sigma\lambda} + \partial_{\nu} h_{\sigma\lambda} + \partial_{\sigma} h_{\lambda\nu} - 2 h_{\sigma\alpha} \bar{\Gamma}_{\nu\lambda}^{\alpha}]$$

$$5b) \delta \Gamma_{00}^i =$$

$$= \frac{1}{2} \bar{g}^{i\beta} \left[\partial_0 h_{\beta 0} + \partial_0 h_{\beta 0} - \partial_\beta h_{00} - 2h_{\beta 0} \bar{\Gamma}_{00}^\beta \right]$$

$$\left[\begin{array}{l} \bar{g}^{\mu\nu} = \text{diag} \\ \Downarrow \\ g = i \end{array} \right] = \frac{1}{2} \delta^{ij} a^{-2} \left[\partial_0 h_{j0} + \partial_0 h_{j0} - \partial_j h_{00} - 2h_{j0} \bar{\Gamma}_{00}^j \right]$$

$\rightarrow 0$
 $(0=0 \Rightarrow h_{00}=0)$

$$= \frac{\partial h_{i0}}{a^2}$$

$$\delta \Gamma_{0i}^0 = \delta \Gamma_{0i}^0 = \frac{1}{2} g^{0\beta} \left[\partial_0 h_{\beta 0} + \partial_0 h_{\beta 0} - \partial_\beta h_{00} - 2h_{\beta 0} \bar{\Gamma}_{0i}^\beta \right]$$

$$[g=0] = \frac{1}{2} g^{00} \left[\cancel{\partial_0 h_{00}} + \cancel{\partial_0 h_{00}} - \cancel{\partial_0 h_{00}} - 2h_{00} \bar{\Gamma}_{0i}^0 \right]$$

$=0$ $=0$

$$= \frac{1}{2} \cdot 2 h_{00} \bar{\Gamma}_{0i}^0 = \underline{\underline{H h_{0i}}}$$

$$\bullet \delta \Gamma_{j0}^i = \frac{1}{2} \bar{g}^{i\beta} \left[\partial_0 h_{\beta j} + \partial_j h_{\beta 0} - \partial_\beta h_{0j} - 2h_{\beta 0} \bar{\Gamma}_{j0}^\beta \right]$$

$$[g=k] = \frac{1}{2} \bar{g}^{ik} \delta_{ik} \left[\cancel{\partial_0 h_{k0}} + \partial_j h_{k0} - \cancel{\partial_k h_{0j}} - 2h_{k0} \bar{\Gamma}_{j0}^k \right]$$

$\rightarrow 0$ $\rightarrow 0$

$$= \frac{1}{2} a^{-2} \delta_{ik} \left[\partial_j h_{k0} - \partial_k h_{0j} \right]$$

$$= \underline{\underline{\frac{1}{2} a^{-2} \left[\partial_j h_{i0} - \partial_i h_{0j} \right]}}$$

$$\delta \Gamma_{ij}^0 = \frac{1}{2} \bar{g}^{0s} \left[\partial_j h_{si} + \partial_i h_{sj} - \partial_s h_{ji} - 2 h_{s0} \Gamma_{ij}^0 \right]$$

$[s=0]$

$$= -\frac{1}{2} \left[\partial_j h_{0i} + \partial_i h_{0j} - \cancel{\partial_s h_{ji}} - 2 \cancel{h_{s0}} \Gamma_{ij}^0 \right]$$

$\swarrow \searrow$
 $0 \quad 0$

$$= -\frac{1}{2} \left[\partial_j h_{0i} + \partial_i h_{0j} \right]$$

$$[\mu = i, \nu = 0]$$

$$\begin{aligned}
 5c) \delta R_{i0} &= \overbrace{\partial_\alpha \delta \Gamma_{i0}^\alpha}^{(1.)} - \overbrace{\partial_0 \delta \Gamma_{i\alpha}^\alpha}^{(2.)} + \overbrace{\delta \Gamma_{i0}^\beta \bar{\Gamma}_{\alpha\beta}^\alpha}^{(3.)} \\
 &+ \overbrace{\delta \Gamma_{\alpha\beta}^\alpha \bar{\Gamma}_{i0}^\beta}^{(4.)} - \overbrace{\delta \Gamma_{i\alpha}^\beta \bar{\Gamma}_{0\beta}^\alpha}^{(5.)} - \overbrace{\delta \Gamma_{0\beta}^\alpha \bar{\Gamma}_{i\alpha}^\beta}^{(6.)} \\
 &= \overbrace{(\partial_0 \delta \Gamma_{i0}^\alpha + \partial_j \delta \Gamma_{i0}^j)}^{(1.)} - \partial_0 (\delta \Gamma_{i0}^\alpha + \delta \Gamma_{i0}^j) \\
 &+ \underbrace{[\delta \Gamma_{i0}^\alpha \bar{\Gamma}_{00}^\alpha + \delta \Gamma_{i0}^\alpha \bar{\Gamma}_{j0}^j + \delta \Gamma_{i0}^j \bar{\Gamma}_{0j}^\alpha + \delta \Gamma_{i0}^j \bar{\Gamma}_{k0}^k]}_{(3.)} \\
 &+ \underbrace{[\delta \Gamma_{00}^\alpha \bar{\Gamma}_{i0}^\alpha + \delta \Gamma_{j0}^j \bar{\Gamma}_{i0}^\alpha + \delta \Gamma_{0j}^\alpha \bar{\Gamma}_{i0}^j + \delta \Gamma_{jk}^j \bar{\Gamma}_{i0}^k]}_{(4.)} \\
 &- \underbrace{[\delta \Gamma_{i0}^\alpha \bar{\Gamma}_{00}^\alpha + \delta \Gamma_{ij}^\alpha \bar{\Gamma}_{00}^\alpha + \delta \Gamma_{i0}^j \bar{\Gamma}_{0j}^\alpha + \delta \Gamma_{ij}^k \bar{\Gamma}_{0k}^j]}_{(5.)} \\
 &- \underbrace{[\delta \Gamma_{00}^\alpha \bar{\Gamma}_{i0}^\alpha + \delta \Gamma_{00}^j \bar{\Gamma}_{ij}^\alpha + \delta \Gamma_{0j}^\alpha \bar{\Gamma}_{i0}^j + \delta \Gamma_{0k}^j \bar{\Gamma}_{ij}^k]}_{(6.)} \\
 &\stackrel{(1.)}{=} \left[\cancel{\partial_0 (H h_{0i})} + \partial_j \left(\frac{\partial_i h_{j0} - \partial_j h_{0i}}{2a^2} \right) \right] - \stackrel{(2.)}{\left[\cancel{\partial_0 (H h_{0i})} + \partial_0 (\delta_{ij} h_{j0} + H) \right]} \\
 &\stackrel{(3.)}{+} \left[H h_{0i} \delta_j^j H \right] + \left[H h_{0j} \delta_i^j H - \delta_{ik} h_{j0} H \delta_k^i \right] \\
 &+ \left[\delta_{ij} h_{k0} H \delta_{ik} H \right] - \stackrel{(5.)}{\left[\frac{\partial_0 h_{j0}}{a^2} \delta_{ij} a^2 H + H h_{0j} \delta_j^i H \right]} \quad (6.)
 \end{aligned}$$

$$\begin{aligned}
 &= \left[\frac{\partial_j \partial_i h_{jo}}{2a^2} - \frac{\partial_j \partial_j h_{io}}{2a^2} \right]^{(1.)} + \left[\partial_0(h_{io} H) \right]^{(2.)} \\
 &+ \left[+3H^2 h_{io} \right]^{(3.)} + \left[\cancel{H^2 h_{io}} - \cancel{H^2 h_{io}} \right]^{(4.)} \\
 &+ \left[H^2 h_{io} \right]^{(5.)} - \left[H \partial_0 h_{io} + h_{io} H^2 \right]^{(6.)}
 \end{aligned}$$

$$= \frac{\partial_j \partial_i h_{jo}}{2a^2} - \frac{\partial_j \partial_j h_{io}}{2a^2} + 3H^2 h_{io} + \partial_0(h_{io} H)$$

$$\cancel{H \partial_0 h_{io}}$$

$$\begin{aligned}
 &h_{io} \partial_0(H) + H \partial_0(h_{io}) \\
 &= h_{io} \left(\frac{\ddot{a}}{a} - \frac{\dot{a}^2}{a^2} \right) + H \partial_0(h_{io}) \\
 &= h_{io} \frac{\ddot{a}}{a} - h_{io} H^2 + H \partial_0 h_{io}
 \end{aligned}$$

$$= h_{io} \left[\frac{\ddot{a}}{a} + 2H^2 \right] + \underbrace{\frac{\partial_j \partial_i h_{jo}}{2a^2}}_{= \partial_j \partial_i [a g_i] = 0} - \underbrace{\frac{\partial_j \partial_j h_{io}}{2a^2}}_{= \frac{\delta^{ij} \partial_i \partial_j h_{io}}{2a^2}}$$

$$= h_{oi} \left[\frac{\ddot{a}}{a} + 2H^2 \right] - \frac{\nabla^2 h_{io}}{2a^2}$$

$$\frac{\nabla^2 h_{io}}{2a^2}$$

$$5d) \quad \delta S_{io} = T_{io} - \cancel{\bar{T}_{io}} - \frac{1}{2} g_{io} T^{\lambda}_{\lambda} + \frac{1}{2} \cancel{\bar{g}_{io}} \bar{T}^{\lambda}_{\lambda}$$

$$\left[\bar{T}_{io} = \bar{p} \cancel{\bar{g}_{io}} + (\bar{s} + \bar{p}) \cancel{\bar{u}_i \bar{u}_o} \right]$$

$$T_{io} = \bar{p} (\cancel{\bar{g}_{io}} + h_{io}) + (\bar{s} + \bar{p}) u_i u_o = \bar{p} a G_i - (\bar{s} + \bar{p}) u_i$$

$$T^{\lambda}_{\lambda} = g^{\lambda\sigma} T_{\sigma\lambda} = g^{\lambda\sigma} [\bar{p} g_{\sigma\lambda} + (\bar{s} + \bar{p}) u_{\sigma} u_{\lambda}]$$

$$= \delta^{\lambda}_{\lambda} \bar{p} + (\bar{s} + \bar{p}) \cancel{u^{\lambda} u_{\lambda}} \rightarrow 0 \text{ (second order)}$$

$$= 4 \bar{p}$$

$$g_{io} = \cancel{\bar{g}_{io}} + h_{io} = a G_i$$

$$\delta S_{io} = \bar{p} a G_i - (\bar{s} + \bar{p}) u_i - 2 a G_i \bar{p}$$

$$= -(\bar{s} + \bar{p}) u_i - a G_i \bar{p}$$

$$R_{io} = 8\pi G S_{io}$$

$$\bar{R}_{io} + \delta R_{io} = 8\pi G (\bar{S}_{io} + \delta S_{io})$$

$$h_{oi} \left(\frac{\ddot{a}}{a} + 2H^2 \right) - \frac{\nabla^2 h_{io}}{2a^2} = -8\pi G \left([\bar{S} + \bar{P}] u_i + a G_i \bar{P} \right)$$

$$\frac{\nabla^2 (a G_i)}{2a^2} = 8\pi G (\bar{S} + \bar{P}) u_i + a G_i \left(\frac{\ddot{a}}{a} + 2H^2 \right) + 8\pi G a G_i \bar{S}$$

$$\frac{\nabla^2 G_i}{2a} = 8\pi G (\bar{S} + \bar{P}) a u_i + a^2 G_i \left(\frac{\ddot{a}}{a} + 2H^2 \right) + 8\pi G a^2 G_i \bar{S}$$

show that this is zero:

$$\frac{\ddot{a}}{a} + 2H^2 = -8\pi G \bar{S}$$

1st Friedmann eqn.

$$\frac{\ddot{a}}{a} + \frac{4\pi G \bar{S}}{3} = -8\pi G \bar{S}$$

second Friedmann eqn.

$$\underline{\underline{0 = 0}}$$

$$5.e) \quad \partial_0 \left[(\bar{g} + \bar{p}) u_i \right] + \frac{3\dot{a}}{a} \left[(\bar{g} + \bar{p}) u_i \right] = 0$$

$$\frac{\partial}{\partial t} V_i + \frac{3}{a} \frac{da}{dt} V_i = 0$$

$$\frac{1}{V_i} dV_i = -\frac{3}{a} da$$

$$\log V_i + V_{i0} = -3 \log a$$

$$V_i(x, a) = V_{i0}(x) a^{-3} \leftarrow \text{Runs as } a^{-3}.$$

Inserting into Einstein equation:

$$\frac{1}{2} \nabla^2 G_i = 8\pi G \frac{V_{i0}(x)}{a^2}$$

Fourier transforming both sides:

$$G_i k^2 \propto \frac{\tilde{V}_i(x)}{a^2} \Rightarrow \underline{G_i \propto \frac{1}{a^2 k^2} \tilde{V}_i(x)}$$

As we see, vector perturbations decay irregardless of scale, while tensor perturbations have exponential decay for small k , and harmonic oscillator evolution for large k .