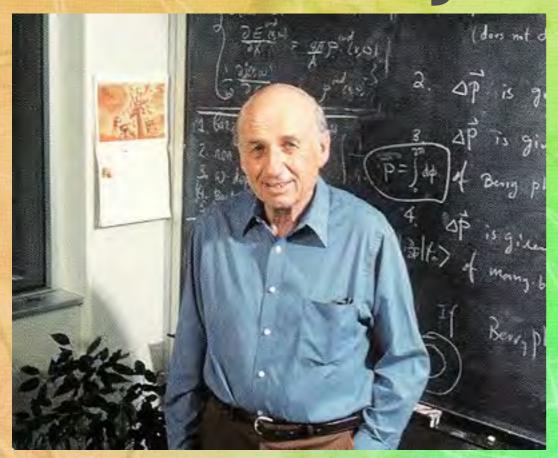




# Outlines

- Correlated materials
- DFT+U
- Haemoglobin
- Hemocyanin
- Conclusions

# Density functional theory



Simple idea: Transform a problem of interacting electrons (N-body problem) to a problem of individual electron interacting with a medium

Walter Kohn: Nobel prize 1998

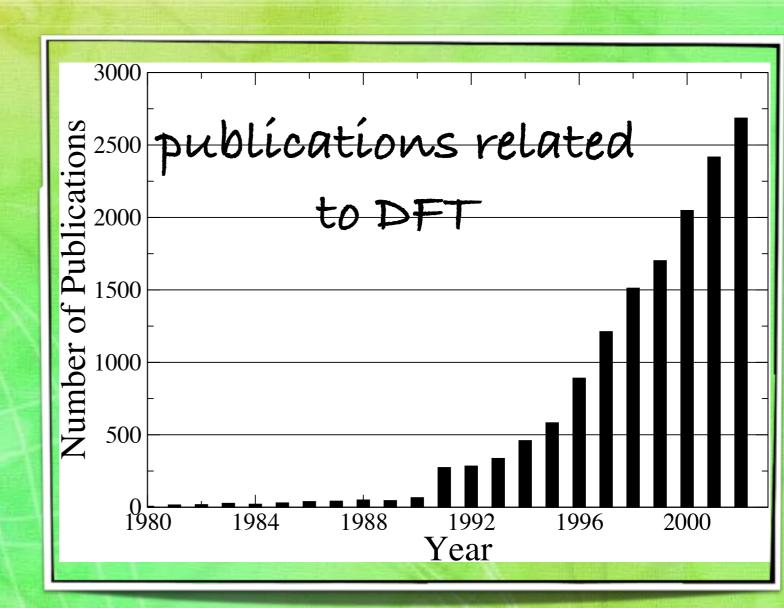
$$\Psi_0 pprox \Psi_{HF} = rac{1}{\sqrt{N!}}$$

$$\Psi_0 \approx \Psi_{HF} = \frac{1}{\sqrt{N!}} \left| \begin{array}{ccccc} \psi_1(\vec{x}_1) & \psi_2(\vec{x}_1) & ... & \psi_N(\vec{x}_1) \\ \psi_1(\vec{x}_2) & \psi_2(\vec{x}_2) & ... & \psi_N(\vec{x}_2) \\ \vdots & & \vdots & & \vdots \\ \psi_1(\vec{x}_N) & \psi_2(\vec{x}_N) & ... & \psi_N(\vec{x}_N) \end{array} \right|$$

$$\rho(\vec{r}) = N \int ... \int |\Psi(\vec{x}_1, \vec{x}_2, ..., \vec{x}_N)|^2 ds_1 d\vec{x}_2 ... d\vec{x}_N$$

#### W. Kohn (Nobel Lecture):

"In the intervening decades enormous progress has been made in finding approximate solutions of Shrodinger's wave equation for systems with several electrons [...]. DFT is an alternative approach to the theory of electronic structure, in which the electron density, rather than the many-body electron wave-function plays a central role."

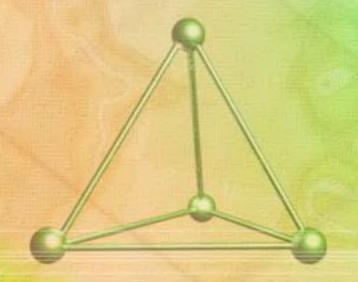


Any limiting cases where DFT needs improvement?

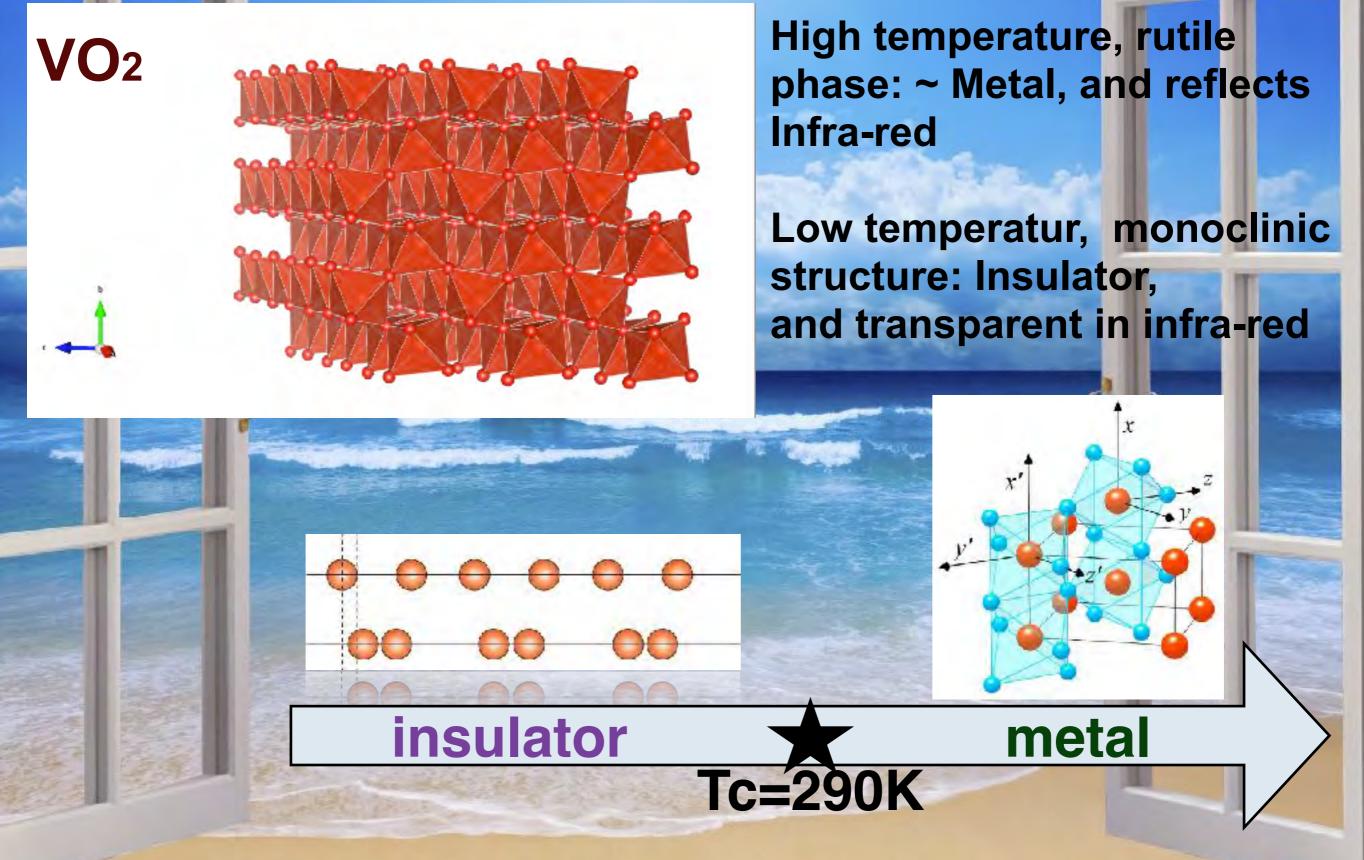
Yes! for localised d- of f- atoms where Coulomb repulsion between electrons is large! Particularly 3d and 4f are close to the nucleus for orthogonality reasons.

# Why study d- or f- materials?

- 2) They raise fundamental questions
- 2) Applications
- 3) They require new formalisms, new ideas

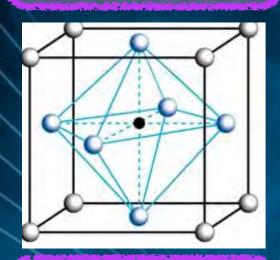


# Big thing example: The intelligent window



## DFT extension required for transition metals

transition metal ions



ion+oxygen cage =transition metal oxide

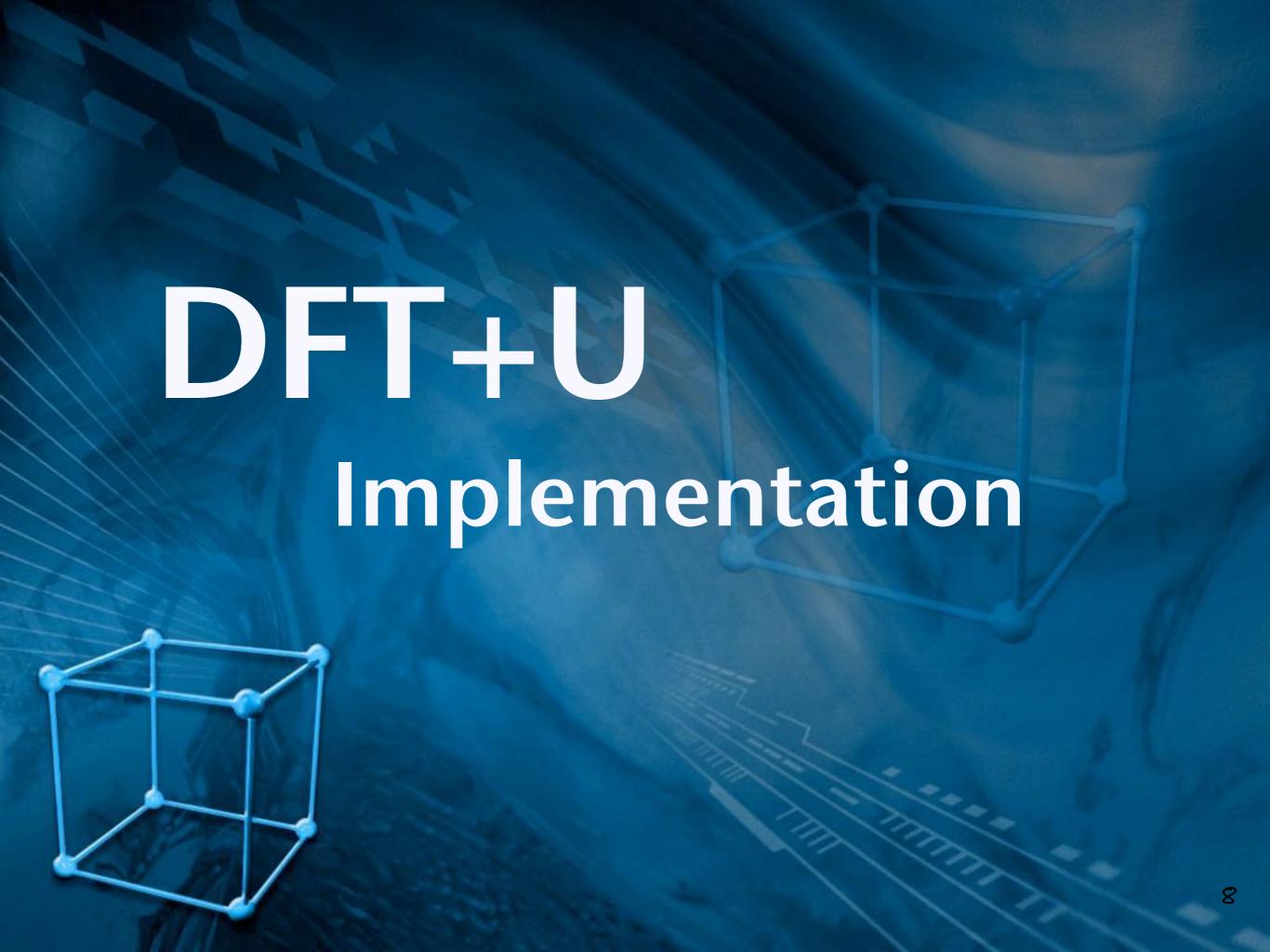


#### V Cr Mn Fe Co Ni Cu

VO<sub>2</sub> Room temperature MIT La<sub>1-x</sub>Sr<sub>x</sub>MnO<sub>3</sub> Colossal <u>Magnetoresistance</u>

Na<sub>x</sub>CoO<sub>2</sub> Thermoelectrics La<sub>1-x</sub>Sr<sub>x</sub>CuO<sub>4</sub> High temperature superconductor

**Extensions of DFT are required** 



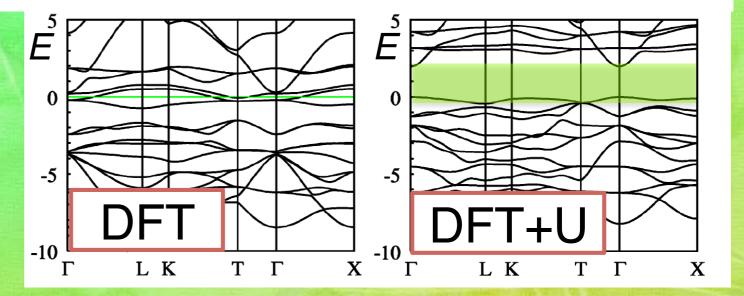
## Basic Hubbard model Hamiltonian

$$H_{\mathrm{Hub}} = t \sum_{\langle i,j \rangle, \sigma} (c_{i,\sigma}^{\dagger} c_{j,\sigma} + \mathrm{h.c.}) + U \sum_{i} n_{i,\uparrow} n_{i,\downarrow}$$

Conductor to insulator transition

$$E_{\text{DFT}+U} = E_{\text{DFT}} + E_{U} = E_{\text{DFT}} + E_{\text{Hub}} - E_{\text{dc}}$$

DFT conductors to DFT+U insulators



conductors ↔ insulators

### How good is DFT?

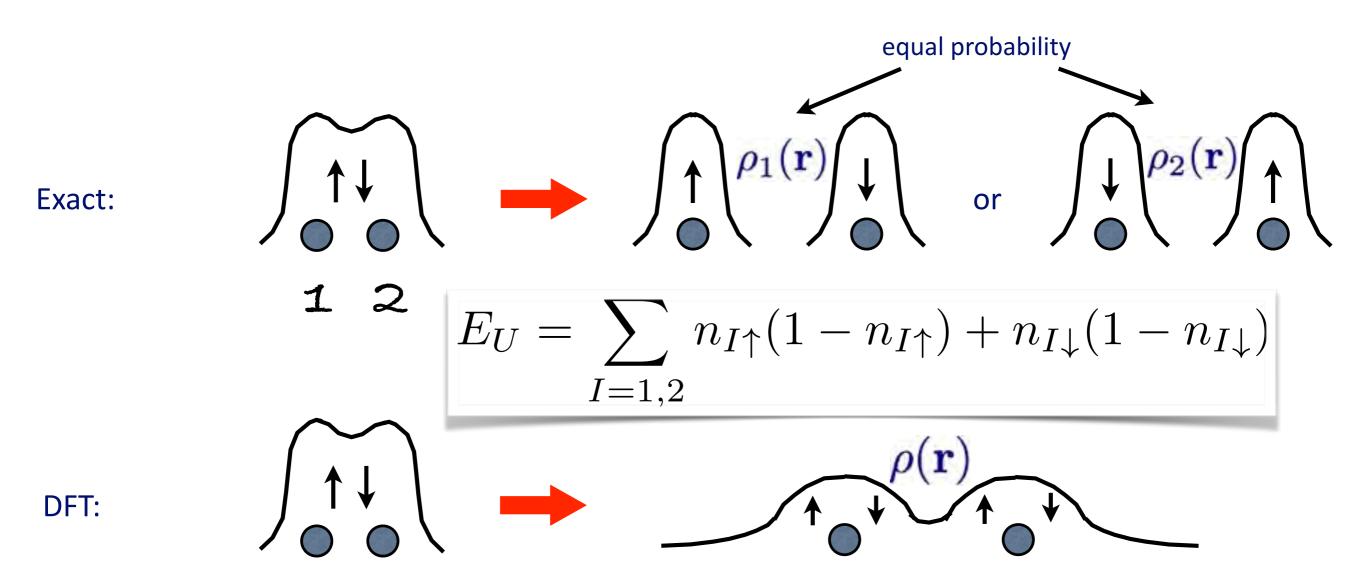
A "simple" case: the dissociation of H<sub>2</sub>

DFT:

$$\rho(\mathbf{r}) = \frac{1}{2}\rho_1(\mathbf{r}) + \frac{1}{2}\rho_2(\mathbf{r})$$

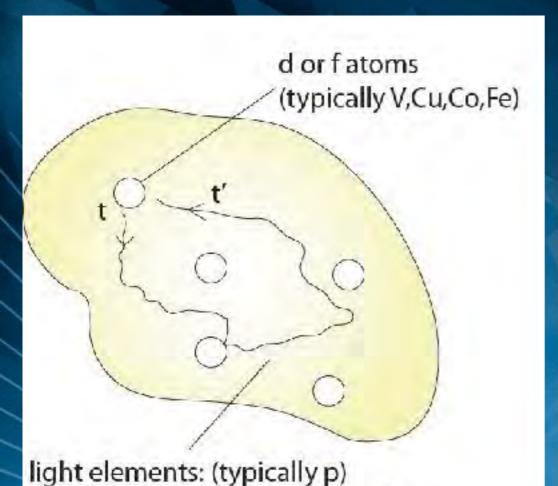
### How good is DFT?

A "simple" case: the dissociation of H<sub>2</sub>

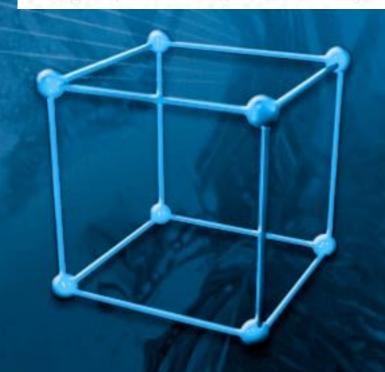


$$\rho(\mathbf{r}) = \frac{1}{2}\rho_1(\mathbf{r}) + \frac{1}{2}\rho_2(\mathbf{r})$$

# DFT+U



oxygen(if oxides), Carbon, Hydrogen



- 1) we define 'hubbard' atoms (the d or f elements)
- 2) we extract the occupation of the hubbard atoms, by projecting the total density n(r) on the d atomic orbitals, we obtain (α is the d orbital index):

$$n_{\alpha}^{I}$$

3) we add a correction to the DFT energy:

$$E_U = \frac{U}{2} \sum_{I,\alpha\sigma} n_{\alpha\sigma}^I (1 - n_{\alpha\sigma}^I)$$

# DFT potentials

derivative of the total energy with respect to the  $\epsilon_i = \frac{\delta E^{\rm LDA+U}}{\delta n_i^{\sigma}}$  occupation of the orbital: We obtain the orbital eigenvalue by taking the

$$\epsilon_i = \frac{\delta E^{\text{LDA+U}}}{\delta n_i^{\sigma}}$$

$$\epsilon_{i} = \frac{\delta E^{\text{LDA}+\text{U}}}{\delta n_{i}^{\sigma}} = \epsilon_{i}^{\text{LDA}} + \frac{\delta}{\delta n_{i}^{\sigma}} \left( \frac{U}{2} \sum_{j\sigma'} n_{j}^{\sigma'} (1 - n_{j}^{\sigma'}) \right)$$

$$\epsilon_i = \frac{\delta E^{\text{LDA+U}}}{\delta n_i^{\sigma}} = \epsilon_i^{\text{LDA}} + \frac{U}{2} \left( -n_i^{\sigma} + (1 - n_i^{\sigma}) \right)$$

$$\epsilon_i = \frac{\delta E^{\mathrm{LDA} + \mathrm{U}}}{\delta n_i^{\sigma}} = \epsilon_i^{\mathrm{LDA}} + U\left(\frac{1}{2} - n_i^{\sigma}\right) \tag{occupied: } -\frac{U}{2}$$
 empty:  $+\frac{U}{2}$ 

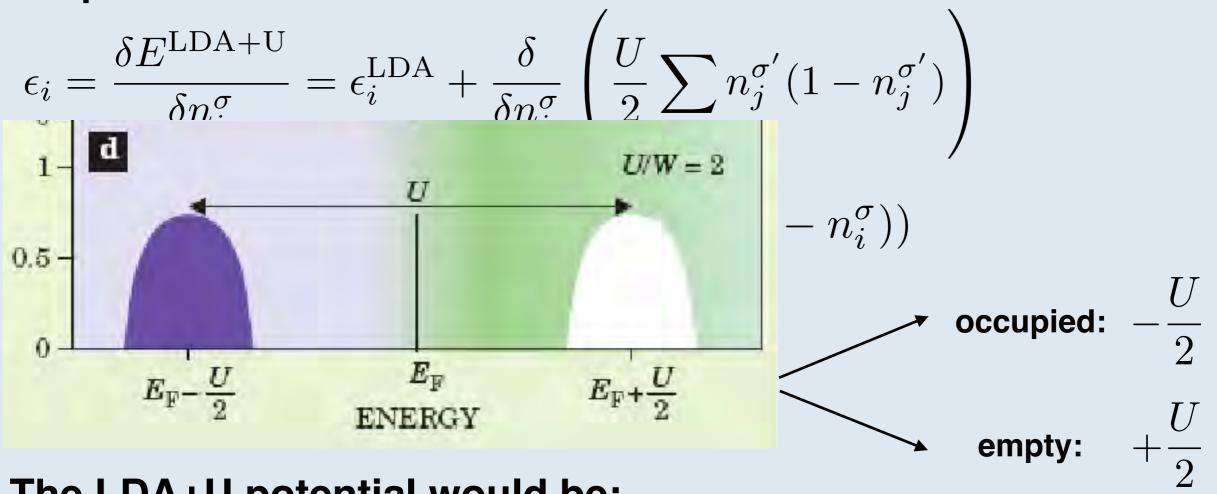
#### The LDA+U potential would be:

$$V^{\text{LDA+U}}(r) = V_i^{\text{LDA}} + U\left(\frac{1}{2} - n_i^{\sigma}\right) P_i$$

# DFT potentials

We obtain the orbital eigenvalue by taking the We obtain the orbital eigenvalue by taking the derivative of the total energy with respect to the  $\epsilon_i = \frac{\delta E^{\mathrm{LDA} + \mathrm{U}}}{\delta n_i^{\sigma}}$ occupation of the orbital:

$$\epsilon_i = \frac{\delta E^{\text{LDA+U}}}{\delta n_i^{\sigma}}$$



#### The LDA+U potential would be:

$$V^{\text{LDA+U}}(r) = V_i^{\text{LDA}} + U\left(\frac{1}{2} - n_i^{\sigma}\right) P_i$$

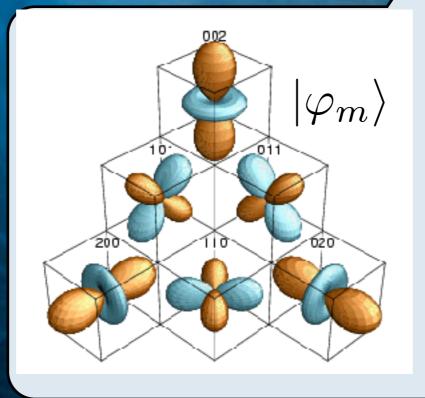
# Projectors

# eigenstate occupation

# Kohn-Sham eigenstate

$$n_{mm'}^{I\sigma} = \sum_{\mathbf{k},v} f_{\mathbf{k}v}^{\sigma} \langle \psi_{\mathbf{k}v}^{\sigma} | P_{mm'}^{I} | \psi_{\mathbf{k}v}^{\sigma} \rangle$$

Hubbard atom density matrix



### projector:

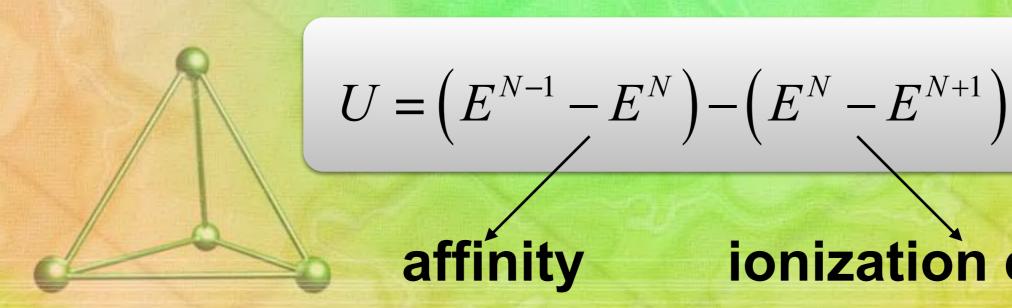
$$P_{mm'}^{I} = |\varphi_{m}^{I}\rangle\langle\varphi_{m'}^{I}|$$

# What about "U"?

What is U? U is the Coulomb repulsion between electron, strong in d orbitals (these orbitals are localized, so the Coulomb repulsion is larger than in s or p orbitals).

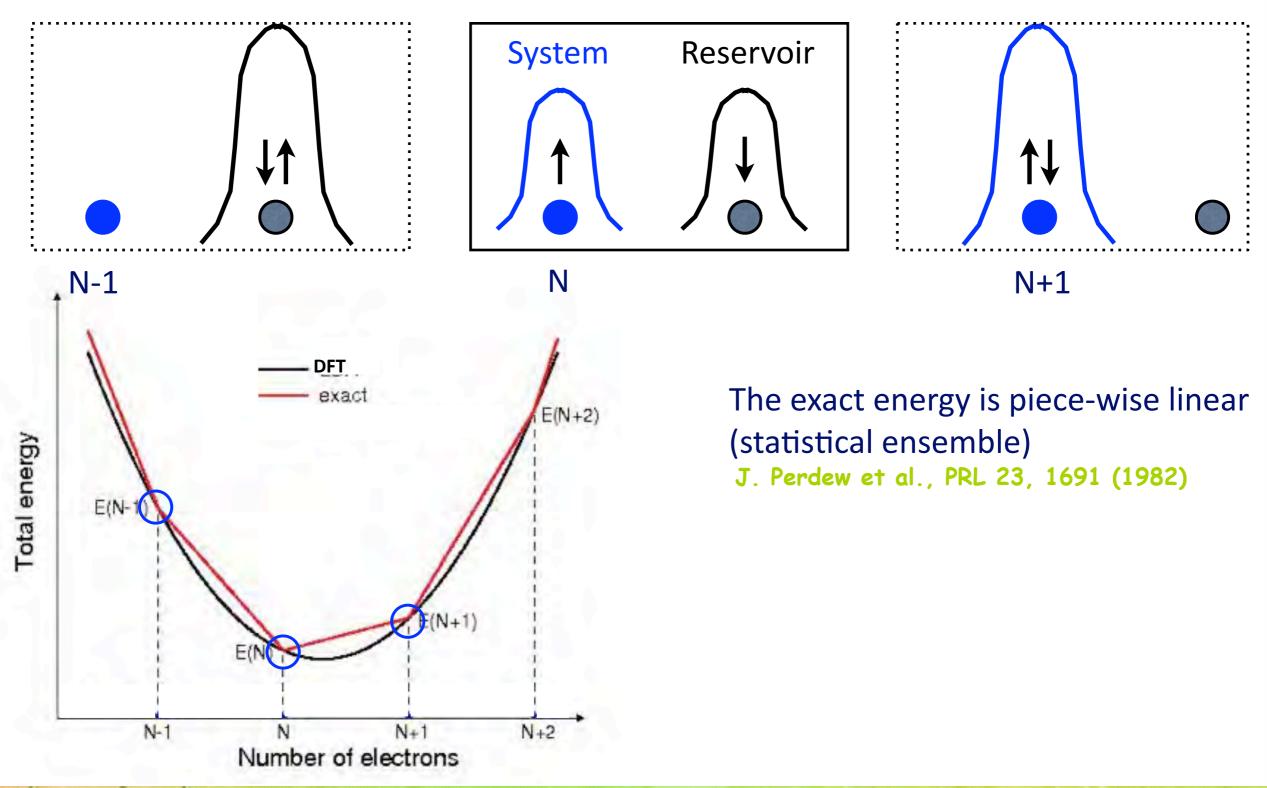
Considering the atoms are embedded in a polarizable surrounding, U is the energy required to move an electron from one atom to another.

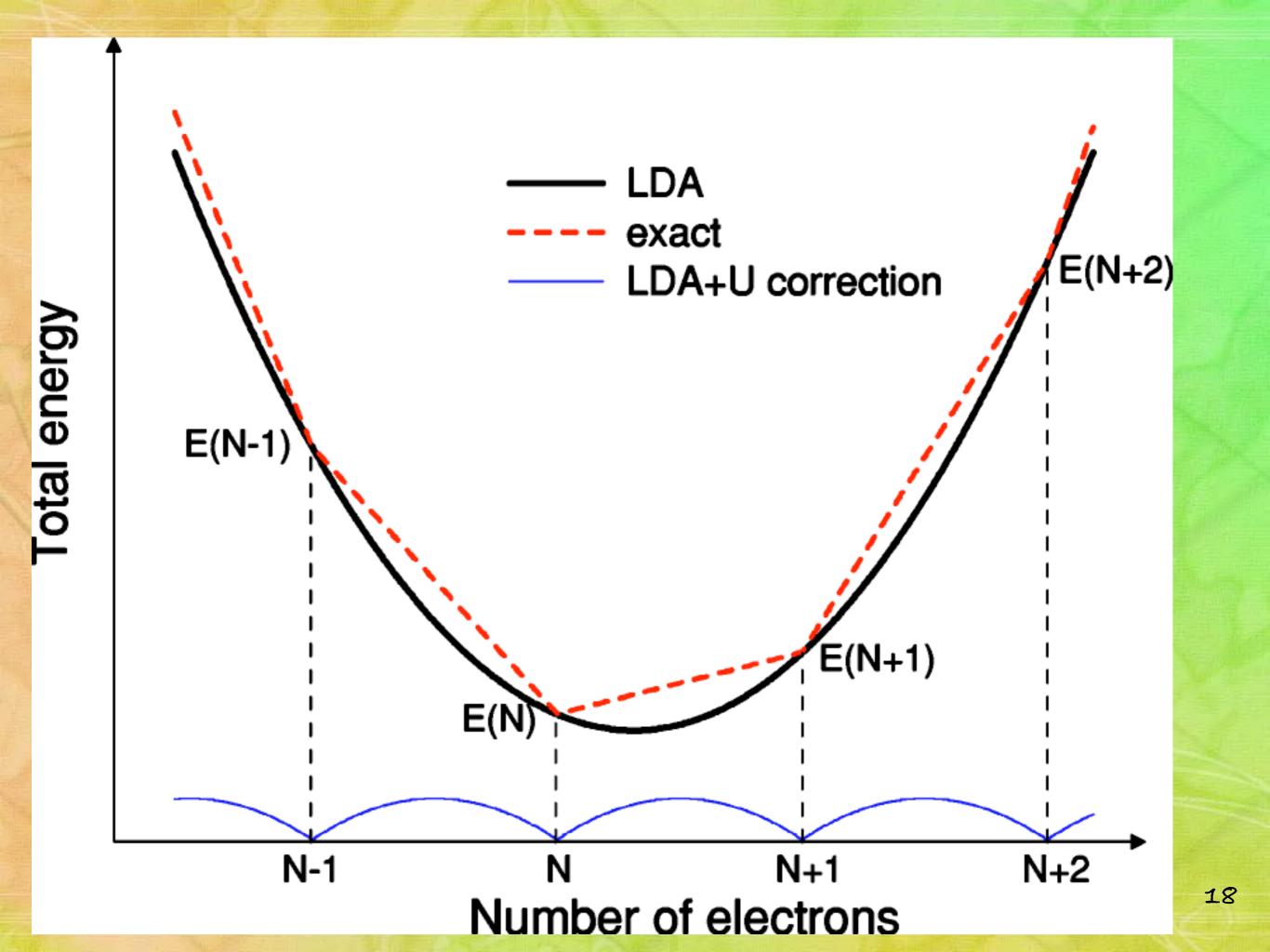
U is equal to the difference of ionization potential and electron affinity of the solid:



ionization energy

## **Energy linearity**

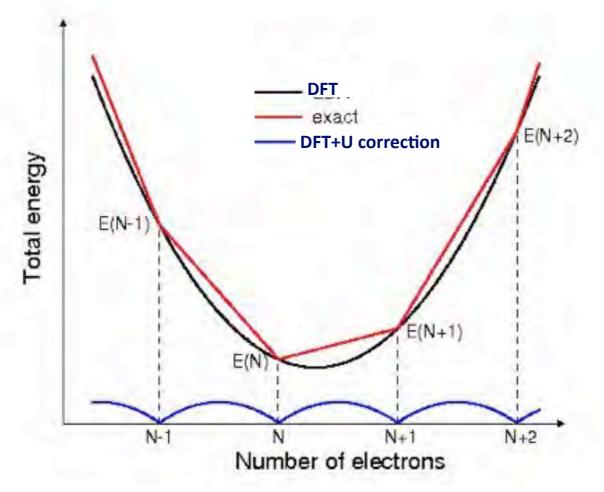




### Open system in contact with a reservoir

### The meaning of U in DFT+U

$$E_{exact} \approx E_{DFT} + \sum_{I} \frac{U^{I}}{2} \sum_{mm'\sigma} \left[ n_{mm'}^{I\sigma} (\delta_{mm'} - n_{mm'}^{I\sigma}) \right] = E_{DFT+U}$$



The (approximate) DFT energy has an unphysical curvature

The exact solution is piecewise linear

+U correction reproduces the exact solution

U and rotationally-invariant U: V.I. Anisimov and coworkers PRB (1991), PRB (1995); Dudarev, and coworkers PRB (1995) LRT U: M. Cococcioni PhD (2002), and M. Cococcioni and S. de

LRT U: M. Cococcioni PhD (2002), and M. Cococcioni and S. de Gironcoli. PRB (2005)

$$U = \frac{d^2 E_{DFT}}{dn^2}$$

### Linear response in practice

Apply a perturbation to the potential acting on the localized states of each Hubbard atom and compute the response of the occupations

$$V_{tot}|\psi_{kv}^{\sigma}\rangle = V_{KS}|\psi_{kv}^{\sigma}\rangle + \alpha^{I}\sum_{m}|\phi_{m}^{I}\rangle\langle\phi_{m}^{I}|\psi_{kv}^{\sigma}\rangle \implies \Delta n^{I}$$

Response matrices:

$$\chi^{IJ} = rac{dn^I}{dlpha^J} \qquad \chi^{IJ}_0 = rac{dn^I_0}{dlpha^J} \qquad \chi = egin{pmatrix} \chi^{11} & \ldots & \chi^{18} \ \ldots & \ldots \ \chi^{81} & \ldots & \chi^{88} \ \end{pmatrix}$$

Effective interactions:

$$U^{I} = \left(\chi_{0}^{-1} - \chi^{-1}\right)_{II}$$

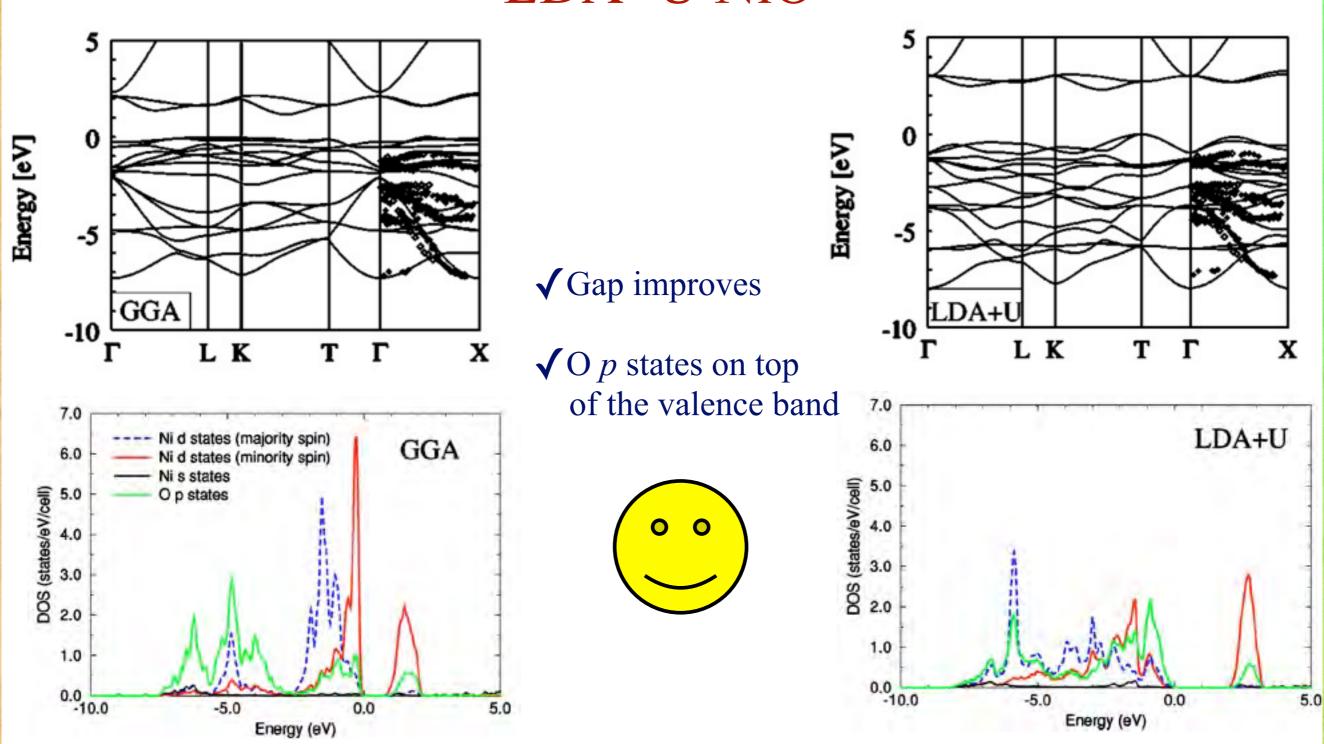
M. Cococcioni and S. de Gironcoli, *PRB* 71, 035105 (2005)

### U is a system-dependent property

Compounds	U
CaCuO <sub>2</sub>	7.5 eV (0.550 Ry)
NiO	8.0 eV (0.590 Ry)
CoO	7.8 eV (0.574 Ry)
FeO	6.8 eV (0.500 Ry)
MnO	6.9 eV (0.507 Ry)
VO	6.7 eV (0.493 Ry)
TiO	6.6 eV (0.485 Ry)

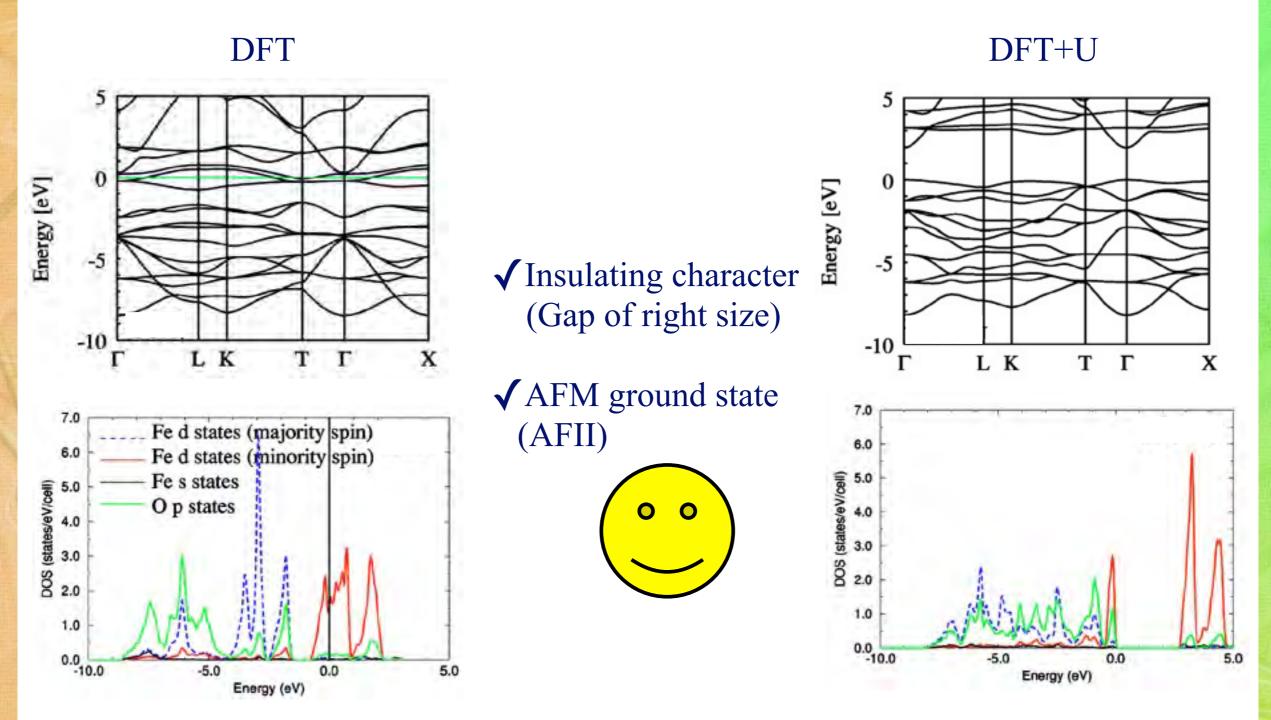
Electron configuration
Covalency/ionicity
Spin states/charge states
Element identity
Coordination numbers

#### LDA+U NiO



M. Cococcioni and S. de Gironcoli, PRB 71, 035105 (2005)

### FeO: DFT and DFT+U



M. Cococcioni and S. de Gironcoli, PRB 71, 035105 (2005)