

# **Classical Black Holes**

09. Accretion in Binary Systems

**Edward Larrañaga** 

#### **Outline for Part 1**

- 1. Roche Lobe Overflow
  - 1.1 Binary Systems
  - 1.2 Roche Lobe Overflow

Formation of the Accretion Disk

Wind Accretion

# Accretion by Roche Lobe Overflow

# **Binary Systems**

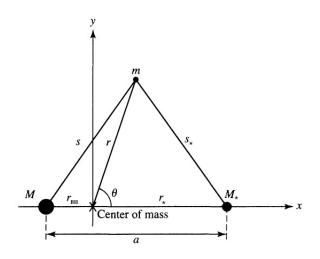
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- Thermodynamical processes in the star's interior may damp (or amplify) the pulsations
- This process may dissipate orbital and rotational energy until the system reaches the state of minimum energy for constant angular momentum: synchronous rotation in circular orbits



Center of mass:

$$r_{BH} + r_* = a$$

$$Mr_{BH} = M_*r_*$$

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Angular velocity of the system:

$$\omega = \frac{\mathsf{v}_*}{\mathsf{r}_*} = \frac{\mathsf{v}_{\mathsf{BH}}}{\mathsf{r}_{\mathsf{BH}}}$$

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$$s^2 = r^2 + r_{BH}^2 + 2rr_{BH}\cos\theta$$

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$$\omega^{2} = \left(\frac{2\pi}{\tau}\right)^{2} = \frac{G(M + M_{*})}{a^{3}} = \frac{G(M + M_{*})}{(r_{BH} + r_{*})^{3}}$$

# **Lagrange Points**

Lagrange Points:

$$\begin{cases} \frac{\partial \Phi}{\partial x} &= 0\\ \frac{\partial \Phi}{\partial y} &= 0 \end{cases}$$

#### **Lagrange Points**

The location of the point  $L_1$  from  $M_{BH}$  and from  $M_*$  are given by the approximations

$$b \approx a [0.500 - 0.227 \log_{10}(q)]$$

$$b_* \approx a [0.500 + 0.227 \log_{10}(q)]$$

- B. Peterson. An Introduction to Active galactic Nuclei. Cambrdige University Press. (1997)
- B. W. Carroll and D. A. Ostlie. An Introduction to Modern Astrophysics. Addison-Wesley (1996)

#### **Equipotential Surfaces**

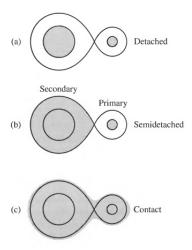
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For  $0.1 \le q \le 0.8$  there is the Paczynski approximation

$$\frac{R_{*R}}{a} \approx \frac{2}{3^{4/3} \left(\frac{q}{1+q}\right)^{1/3}}$$

$$\frac{R_{*R}}{a} \approx 0.462 \left(\frac{M_*}{M+M_*}\right)^{1/3}$$

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If Roche lobe shrinks → overflow continues!

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$$a = \frac{J^2}{G} \frac{M + M_*}{M^2 M_*^2}$$

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For conservative systems: BH grows putting more mas near the CM and the star moves in a wider orbit, increasing a, in order to conserve J.

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• Conservative systems  $(\dot{J} = 0)$  with  $q > \frac{5}{6}$ : Roche lobe of the star shrinks and the overflow continues.

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The Roche overflow is rapid and violent but stops when q is smaller than  $\frac{5}{4}$ .

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- The binary system loses angular momentum. There are many mechanisms for losing *J*:
  - Gravitational radiation
  - Tidal forces on the star
  - Wind (magnetically linked to the star)

#### Outline for Part 2

- 1. Roche Lobe Overflow
  - 1.1 Binary Systems
  - 1.2 Roche Lobe Overflow

2. Formation of the Accretion Disk

3. Wind Accretion

# Formation of the Accretion Disk

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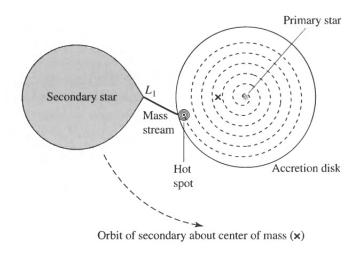
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 $v_{\parallel}$ : parallel component of the velocity with respect to the line of centers.

 $v_{\perp}$ : perpendicular component of the velocity with respect to the line of centers.





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Using

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we get

$$v_{\perp} \gtrsim 305.6 \left(\frac{M}{M_{\odot}}\right)^{1/3} (1+q)^{1/3} \tau_{hr}^{-1/3} \text{ [km/s]}$$

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- The particle will describe an elliptical trajectory around the BH but the presence of the star will make this ellipse to precess slowly.

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- Hence, particles will go into the trajectories with minimum energy for a given angular momentum, i.e. circular orbits!
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- The radius of the resulting circular orbit is called *circularization* radius,  $r_{circ}$ .

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- In the ring the particles have collisions, shocks, viscous dissipation and other processes that transform some of the potential energy into heat (producing radiation).
- However, this release of energy needs the loosing of angular momentum.
- In the absence of external torques, the only possible process is a *transfer of angular momentum* from inner regions outwards by internal torques.
- The redistribution of angular momentum makes particles in the outer parts move outwards (gaining angular momentum) and the particles in the inner particles spiral inwards.

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  - Turbulence in the fluid may be an origin of angular momentum transport. Turbulence may be produced by mechanisms as Themally driven convection, Pure hydrodynamic instabilities or Magnetohydrodynamic (MHD) turbulence.

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## **Outline for Part 3**

- 1. Roche Lobe Overflow
  - 1.1 Binary Systems
  - 1.2 Roche Lobe Overflow

2. Formation of the Accretion Disk

3. Wind Accretion

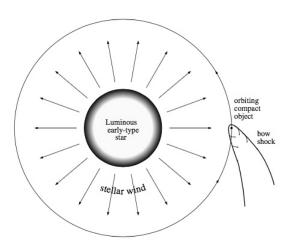
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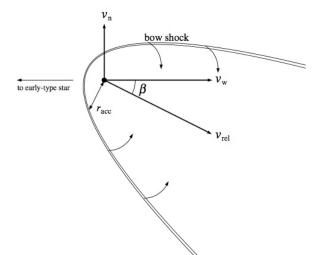
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Remember that for the typical conditions of a star,

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This implies that

- we may neglect gas pressure
- the flow may be considered as a collection of particles



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• β: Angle in which the winds moves

$$\beta = \tan^{-1}\left(\frac{v_n}{v_w}\right)$$

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The region of capture is a cylinder with axis along the relative wind direction,  $v_{rel}$ , and radius

$$r_{acc} \sim \frac{2GM}{v_{rel}^2}$$

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a: Separation of the elements of the binary system.

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$$\frac{\dot{M}}{-\dot{M}_{w}} \sim \frac{1}{4} \left(\frac{M}{M_{*}}\right)^{2} \left(\frac{R_{*}}{a}\right)^{2}$$

# Wind Accretion Rate. Example

Consider  $M_* \sim 5M$  and  $R_* \sim 0.5a$ .

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Wind accretion is very inefficient!

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Remember that in the Roche lobe accretion this is  $\sim b^2 \omega \sim 0.500 a^2 \omega$ . Therefore the angular momentum in the wind accretion is smaller and the chances of disk formation are smaller.

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- $\lambda(r)$  is unknown but may affect strongly the estimate of  $r_{circ}$
- The value of r<sub>acc</sub> used here is an approximation. The highly supersonic speed of the wind produces a strong bow shock that may change the description.

**Next Lecture** 

10. Viscous Torques