# MATH 316 Lecture 12

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# 1 Wave Equations: Neumann Boundary Conditions

Solve IBVP:  $y_{tt} = a^2 y_{xx}$ 

Boundary conditions:  $y_x(0,t) = y_x(L,t) = 0 \Rightarrow$  Homogeneous Neumann boundary conditions.

Initial conditions: y(x,0) = f(x),  $y_t(x,0) = 0 = g(x)$  (Therefore zero initial velocity, with specified initial displacement).

Split into two problems (w and z). Because g(x) = 0, we only have the z equation. For the solution, we use separation of variables.

$$X'' + \lambda X = 0$$

$$\Rightarrow X_n(x) = \cos\left(\frac{n\pi x}{L}\right)$$

$$\lambda_n = \left(\frac{n\pi}{L}\right)^2$$

$$X_0(x) = 1 \Rightarrow \lambda_0 = 0$$

$$\ddot{T} + a^2 \lambda T = 0$$

$$\dot{T}_n(0) = 0$$

$$T_n(t) = A_n \cos\left(\frac{n\pi a}{L}t\right)$$

$$T_0(t) = 1$$

The solution would be:

$$y(x,t) = A_0 + \sum_{n=1}^{\infty} A_n \cos\left(\frac{n\pi x}{L}\right) \cos\left(\frac{n\pi a}{L}t\right)$$

Note that  $A_0$  is the multiplication of  $X_0T_0$  terms.

At 
$$t = 0 \longrightarrow y(x, 0) = A_0 + \sum_{n=1}^{\infty} A_n \cos\left(\frac{n\pi x}{L}\right) = f(x)$$

We construct Fourier cosine series for f(x).

$$f(x) = \frac{a_0}{2} + \sum_{n=1}^{\infty} n = 1^{\infty} a_n \cos\left(\frac{n\pi x}{L}\right)$$
$$\Rightarrow A_0 = \frac{a_0}{2}, A_n = a_n$$

Note: We can do a similar solution if initial velocity is given.

#### 1.1 Recap

- Introduced wave equation
- Developed separation of variables method to find its solution
  - Dirichlet and Neumann boundary conditions
  - Examples and normal modes

Now: New method.

• New look at the wave equation and we solve the wave equation using **D'Alembert's solution**.

<sup>&</sup>lt;sup>1</sup>Note that left and right side of table are entirely separate.

### 2 D'Alembert's Solution

$$y_{tt} = a^2 y_{xx}$$

Let's see if we can guess a solution of exponential format.

$$y(x,t) = e^{ikx + \sigma t}$$

where k and  $\sigma$  are constants. <sup>2</sup>

Substitute the guessed solution into the PDE.

$$y_t t = \sigma^2 e^{ikx + \sigma t}$$

$$y_{xx} = -k^2 e^{ikx + \sigma t}$$

Now, substitute this into the PDE:

$$(\sigma^2 + a^2 k^2) e^{ikt + \sigma t} = 0$$

$$\Rightarrow \sigma = \pm ika$$

$$y_1(x, t) = e^{ik(x - at)}$$

$$y_2(x,t) = e^{ik(x+at)}$$

 $x \pm at$  are known as characteristics, these are lines in x and t along which the initial conditions (and general information) travels.

The question here is this: Can this form of solution be more general such that we can apply it to any wave equation?

$$y_1(x,t) = F(x-at), y_2(x,t) = G(x+at)$$

Can we find a general equation that satisfies the wave equation?

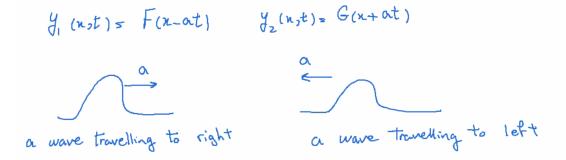


Figure 1: F(x - at) is a wave travelling to the right with a speed of a. G(x + at) is a wave travelling to the left with a speed of a.

Hence, a general solution:

$$y(x,t) = F(x - at) + G(x + at)$$

Does it satisfy the PDE?

<sup>&</sup>lt;sup>2</sup>Try this guess solution with heat solution! You will find that it does work for heat equations.

$$y(x,0) = f(x) \Rightarrow F(x) + G(x) = f(x)$$
 (1)  
 $y_t(x,0) = g(x) \Rightarrow -aF'(x) + aG'(x) = g(x)$  (2)

We get (2) from:

$$-aF(x) + aG(x) = \int_0^x g(s)ds + A$$

$$(1)xa + 2 \Rightarrow 2aG(x) + f(x) = \int_0^x g(s)ds + A$$

$$\Rightarrow G(x) = \frac{1}{2}f(x) + \frac{1}{2a}\int_0^x g(s)ds + \frac{A}{2a}$$

To find F(x):

$$(1)xa - (2) \Rightarrow 2aF(x) = af(x) - \int_0^x g(s)ds - A$$
$$\Rightarrow F(x) = \frac{1}{2}f(x) - \frac{1}{2a}\int_0^x g(s)ds - \frac{A}{2a}$$

Now, substitute these into the general solution: (plug into y(x,t) = F(x-at) + G(x+at)) This gives us:

$$\frac{1}{2}f(x-at) - \frac{1}{2a} \int_0^{x-at} g(s)ds + \frac{1}{2}f(x+at) + \frac{1}{2a} \int_0^{x+at} g(s)ds$$

Note that  $-\frac{A}{2a}$  and  $\frac{A}{2a}$  cancel.

$$y(x,t) = \frac{1}{2} \left[ \underbrace{f(x-at)}_{\text{half of init cond travels right}} + \underbrace{f(x+at)}_{\text{half of init cond travels left}} \right] + \frac{1}{2a} \int_{x-at}^{x+at} g(s) ds$$
 (1)

N.B. Above analysis has no boundary conditions:  $-\infty < x < \infty$ 

What if the problem has boundary condition?

Let  $F^{o}(x)$  and  $G^{o}(x)$  be the odd<sup>3</sup> 2L-periodic extension of f(x) and g(x) respectively:

$$y(x,t) = \frac{1}{2} \left[ F^{o}(x-at) + F^{o}(x+at) \right] + \frac{1}{2a} \int_{x-at}^{x+at} G^{o}(s) ds$$

Boundary conditions: <sup>4</sup>

$$y(0,t) = \underbrace{\frac{1}{2} \left[ F^{o}(-at) + F^{o}(at) \right]}_{=0} + \underbrace{\frac{1}{2a} \int_{-at}^{at} G^{o}(s) ds}_{=0} = 0$$

What's the relationship between D'Alembert's formula and the separation of variables?

$$y(x,t) = \sum_{n=1}^{\infty} \sin\left(\frac{n\pi x}{L}\right) \left[b_n \frac{L}{n\pi a} \sin\left(\frac{n\pi a}{L}t\right) + b'_n \cos\left(\frac{n\pi a}{L}t\right)\right]$$

Recall trig formulae:

$$\sin(A)\sin(B) = \frac{1}{2}\left[\cos(A-B) - \cos(A+B)\right]$$

<sup>&</sup>lt;sup>3</sup>(Assumes Dirichlet boundary conditions)

<sup>&</sup>lt;sup>4</sup>Note that we are using the properties of odd functions to cancel out both F and G.

$$\sin(A)\cos(B) = \frac{1}{2}\left[\sin(A-B) + \sin(A+B)\right]$$

Let's apply these:

$$y(x,t) = \frac{1}{2} \sum_{n=1}^{\infty} \frac{b_n L}{n\pi a} \underbrace{\left[\cos\left(\frac{n\pi}{L}(x-at)\right) - \cos\left(\frac{n\pi}{L}(x+at)\right)\right]}_{\text{Let's write this in integral format}} \rightarrow \\ \leftrightarrow +b'_n \left[\sin\left(\frac{n\pi}{L}(x-at)\right) - \sin\left(\frac{n\pi}{L}(x+at)\right)\right] \\ y(x,t) = \frac{1}{2} \sum_{n=1}^{\infty} b'_n \left[\sin\left(\frac{n\pi}{L}(x-at)\right) - \sin\left(\frac{n\pi}{L}(x+at)\right)\right] \rightarrow \\ \leftrightarrow +\frac{1}{2a} \sum_{n=1}^{\infty} b_n \int_{x-at}^{x+at} \sin\left(\frac{n\pi s}{L}\right) ds$$

Recall:

$$\sum_{n=1}^{\infty} b'_n \sin(\frac{n\pi x}{L}) = f(x)$$

and

$$\sum_{n=1}^{\infty} b_n \sin\left(\frac{n\pi x}{L}\right) = g(x)$$

 $\Rightarrow$  d'Alembert's solution:

$$y(x,t) = \frac{1}{2} \left[ f(x-at) + f(x+at) \right] + \frac{1}{2a} \int_{x-at}^{x+at} g(s) ds$$

Both methods give similar solution.