

# 1 Proof of $|\hat{\sigma} \cdot \hat{r}; \pm\rangle \equiv |\psi(\theta, \varphi)\rangle = \cos(\theta/2) |0\rangle + \sin(\theta/2)e^{i\varphi} |1\rangle$ .

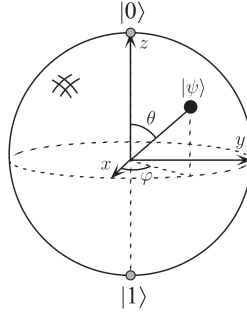


Figure 1.1: Bloch sphere representation of a qubit  $|\psi\rangle$ , as a function of the angles  $\theta$  and  $\varphi$ .

In the Bloch sphere representation of a qubit, Fig. 1.1,

$$\hat{r} = \frac{\vec{r}}{|\vec{r}|} = \frac{\vec{r}}{r} = \sin \theta \cos \varphi \hat{x} + \sin \theta \sin \varphi \hat{y} + \cos \theta \hat{z} = (\sin \theta \cos \varphi \hat{x}, \sin \theta \sin \varphi \hat{y}, \cos \theta \hat{z}). \quad (1.1)$$

Hence,<sup>1</sup>

$$\hat{\vec{\sigma}} \cdot \hat{r} = (\hat{\sigma}_1 \hat{x}, \hat{\sigma}_2 \hat{y}, \hat{\sigma}_3 \hat{z}) \cdot (\sin \theta \cos \varphi \hat{x}, \sin \theta \sin \varphi \hat{y}, \cos \theta \hat{z}) = \hat{\sigma}_1 \sin \theta \cos \varphi + \hat{\sigma}_2 \sin \theta \sin \varphi + \hat{\sigma}_3 \cos \theta. \quad (1.2)$$

From the definition of the Pauli matrices [10, p. 169, Eq. (3.2.32)], in the  $\{|\pm\rangle_z\} \equiv \{|0\rangle, |1\rangle\}$  computational basis,

$$\begin{aligned} \hat{\vec{\sigma}} \cdot \hat{r} &= \hat{\sigma}_1 \sin \theta \cos \varphi + \hat{\sigma}_2 \sin \theta \sin \varphi + \hat{\sigma}_3 \cos \theta \\ &\doteq \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \sin \theta \cos \varphi + \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \sin \theta \sin \varphi + \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \cos \theta \\ &= \begin{pmatrix} \cos \theta & \sin \theta \cos \varphi - i \sin \theta \sin \varphi \\ \sin \theta \cos \varphi + i \sin \theta \sin \varphi & -\cos \theta \end{pmatrix} = \begin{pmatrix} \cos \theta & e^{-i\varphi} \sin \theta \\ e^{i\varphi} \sin \theta & -\cos \theta \end{pmatrix}. \end{aligned} \quad (1.3)$$

In order to find the eigenvalues ( $\lambda$ ) and the pertaining eigenvectors corresponding to the operator  $\hat{\vec{\sigma}} \cdot \hat{r}$  ( $|\hat{\vec{\sigma}} \cdot \hat{r}; \pm\rangle$ ),

$$\begin{aligned} |\hat{\vec{\sigma}} \cdot \hat{r} - \lambda \hat{1}| &\doteq \begin{vmatrix} \cos \theta - \lambda & e^{-i\varphi} \sin \theta \\ e^{i\varphi} \sin \theta & -\cos \theta - \lambda \end{vmatrix} = (\cos \theta - \lambda)(-\cos \theta - \lambda) - \sin^2 \theta = -\cos^2 \theta - \lambda \cos \theta \\ &\quad + \lambda \cos \theta + \lambda^2 - \sin^2 \theta = -\cos^2 \theta - \sin^2 \theta + \lambda^2 = \lambda^2 - 1 = 0 \leftrightarrow \lambda = \pm 1, \end{aligned} \quad (1.4)$$

as we could have inferred directly from  $\hat{\vec{\sigma}} \cdot \hat{r} |\hat{\vec{\sigma}} \cdot \hat{r}; \pm\rangle = \pm |\hat{\vec{\sigma}} \cdot \hat{r}; \pm\rangle = \lambda |\hat{\vec{\sigma}} \cdot \hat{r}; \pm\rangle \leftrightarrow \lambda = \pm 1$ .

Let us find the eigenvectors of  $\hat{\vec{\sigma}} \cdot \hat{r}$ ,  $|\hat{\vec{\sigma}} \cdot \hat{r}\rangle$ . If we substitute  $\lambda = \pm 1$  in Eq. (1.4),

$$\begin{pmatrix} \cos \theta \mp 1 & e^{-i\varphi} \sin \theta \\ e^{i\varphi} \sin \theta & -\cos \theta \mp 1 \end{pmatrix} \begin{pmatrix} c_1 \\ c_2 \end{pmatrix} = 0 = \begin{pmatrix} (\cos \theta \mp 1)c_1 + (e^{-i\varphi} \sin \theta)c_2 \\ (e^{i\varphi} \sin \theta)c_1 + (-\cos \theta \mp 1)c_2 \end{pmatrix}. \quad (1.5)$$

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<sup>1</sup> $\hat{\vec{\sigma}} \cdot \hat{r}$  is an operator, the *hat* in  $\hat{\vec{\sigma}}$  means that it is an operator, and the hat in  $\hat{r}$ , that it is a unitary vector.

The 2 rows we have for each eigenvalue  $\lambda$  in the column matrix in the right-hand side of Eq. (1.5) should be equivalent. From the first row,  $c_2 = -\frac{(\cos\theta+1)c_1}{e^{-i\varphi}\sin\theta} = \frac{\pm 1 - \cos\theta}{\sin\theta} e^{i\varphi} c_1$ . So, if we choose  $c_1 = \sin\theta$ ,

$$|\hat{\sigma} \cdot \hat{r}; \pm\rangle \doteq N \begin{pmatrix} c_1 \\ c_2 \end{pmatrix} = N \begin{pmatrix} \sin\theta \\ (\pm 1 - \cos\theta)e^{i\varphi} \end{pmatrix}, \quad (1.6)$$

subject to the normalization condition

$$\begin{aligned} 1 &= \langle \hat{\sigma} \cdot \hat{r}; \pm | \hat{\sigma} \cdot \hat{r}; \pm \rangle \doteq N^* \begin{pmatrix} \sin\theta & (\pm 1 - \cos\theta)e^{-i\varphi} \end{pmatrix} N \begin{pmatrix} \sin\theta \\ (\pm 1 - \cos\theta)e^{i\varphi} \end{pmatrix} \\ &= N^* N (\sin^2\theta + (\pm 1 - \cos\theta)^2) = |N|^2 (\sin^2\theta + (\pm 1 - \cos\theta)^2) \\ \Leftrightarrow |N| &= \frac{1}{\sqrt{(\sin^2\theta + (\pm 1 - \cos\theta)^2)}} \Leftrightarrow N = \frac{1}{\sqrt{(\sin^2\theta + (\pm 1 - \cos\theta)^2)}} e^{i\delta} : \delta \in \mathbb{R}. \end{aligned} \quad (1.7)$$

We will choose the phase  $\delta$  so that  $N$  is real and positive. For example, if we choose  $\delta = 0$ , Eq. (1.6) becomes

$$|\hat{\sigma} \cdot \hat{r}; \pm\rangle \doteq \frac{1}{\sqrt{(\sin^2\theta + (\pm 1 - \cos\theta)^2)}} \begin{pmatrix} \sin\theta \\ (\pm 1 - \cos\theta)e^{i\varphi} \end{pmatrix}. \quad (1.8)$$

**Careful!** In Fig. 1.1, we use  $|+\rangle_z \equiv |0\rangle \equiv |\hat{\sigma}_z; +\rangle$  as a *reference axis*. Then, since we use the coordinate system chosen in Fig. 1.1, we should take the upper sign. Thus, Eq. (1.8) becomes

$$|\hat{\sigma} \cdot \hat{r}; +\rangle \doteq \frac{1}{\sqrt{(\sin^2\theta + (+1 - \cos\theta)^2)}} \begin{pmatrix} \sin\theta \\ (+1 - \cos\theta)e^{i\varphi} \end{pmatrix}, \quad (1.9)$$

where

- For the first row of Eq. (1.8),

$$\begin{aligned} N c_1 &= \frac{\sin\theta}{\sqrt{\sin^2\theta + (1 - \cos\theta)^2}} = \frac{2 \sin(\theta/2) \cos(\theta/2)}{\sqrt{(2 \sin(\theta/2) \cos(\theta/2))^2 + (2 \sin^2(\theta/2))^2}} = \frac{2 \sin(\theta/2) \cos(\theta/2)}{\sqrt{4 \sin^2(\theta/2) (\cos^2(\theta/2) + \sin^2(\theta/2))}} \\ &= \frac{2 \sin(\theta/2) \cos(\theta/2)}{2 \sin(\theta/2)} = \cos(\theta/2). \end{aligned} \quad (1.10)$$

- For the second row of Eq. (1.8),

$$N c_2 = \frac{1 - \cos\theta}{\sin\theta} e^{i\varphi} c_1 = \frac{1 - \cos\theta}{\sin\theta} e^{i\varphi} \cos(\theta/2) = \frac{2 \sin^2(\theta/2)}{2 \sin(\theta/2) \cos(\theta/2)} e^{i\varphi} \cos(\theta/2) = \sin(\theta/2) e^{i\varphi}. \quad (1.11)$$

Thus, we obtain the desired normalized eigenket that represents a qubit in the Bloch sphere in Fig. 1.1,

$$|\hat{\sigma} \cdot \hat{r}; \pm\rangle = \cos(\theta/2) |0\rangle + \sin(\theta/2) e^{i\varphi} |1\rangle \quad (1.12)$$

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