PART I

Question 1

Adiabatic,
$$dQ = 0$$
, $\Rightarrow dW = P \ dV$
 $PV^{\gamma} = k$, $P_1V_1^{\gamma} = P_2V_2^{\gamma}$

$$\begin{split} \therefore W &= \int_{V_1}^{V_2} \frac{k}{V^{\gamma}} dV \\ &= \left[\frac{k}{\gamma - 1} \frac{1}{V^{\gamma - 1}} \right]_{V_1}^{V_2} \\ &= \frac{k}{\gamma - 1} \left(\frac{1}{V_2^{\gamma - 1}} - \frac{1}{V_1^{\gamma - 1}} \right) \\ &= \frac{P_1}{\gamma - 1} \left(\frac{V_1^{\gamma}}{V_2^{\gamma - 1}} - \frac{V_1^{\gamma}}{V_1^{\gamma - 1}} \right) \\ &= \frac{P_1}{\gamma - 1} \left[\left(\frac{V_1}{V_2} \right)^{\gamma - 1} - 1 \right] \end{split}$$

Question 2

$$F = E - TS = n\varepsilon - Tk \ln \left[\frac{N!}{n! (N-n)!} \right]^{2}$$

$$\frac{1}{T} = \frac{\partial S}{\partial E} = \frac{\partial S}{\partial n} \frac{\partial n}{\partial E} = \varepsilon \frac{\partial S}{\partial n}$$

$$S = 2k [\ln N! - \ln n! - \ln(N-n)!]$$

$$\frac{\partial S}{\partial n} = 2k \left[-\frac{\partial}{\partial n} \ln n! - \frac{\partial}{\partial n} \ln (N - n)! \right]$$

$$= -2k \left[\frac{\partial}{\partial n} (n \ln n - n) + \frac{\partial}{\partial n} ((N - n) \ln (N - n) - N + n) \right]$$

$$= -2k [\ln n + 1 - 1 - 1 - \ln (N - n) + 1]$$

$$= -2k \ln \left(\frac{n}{N - n} \right)$$

$$\frac{\varepsilon}{2kT} = \ln \left(\frac{N - n}{n} \right)$$

$$e^{\frac{\varepsilon}{2kT}} = \frac{N}{n} - 1$$

$$\frac{n}{N} = \frac{1}{e^{\frac{\varepsilon}{2kT}} + 1}$$

Question 3

$$\begin{split} Z &= \sum_{r} (2r+1) e^{-\beta \left[\varepsilon_{r}^{tr} + \frac{\hbar}{2l} r(r+1) \right]} = \sum_{r} e^{-\beta \varepsilon_{r}^{tr}} \sum_{r} (2r+1) e^{\beta \varepsilon_{r}^{rot}} = Z_{trans} Z_{rot} \\ Z_{1}^{trans} &= \int_{0}^{\infty} f(p) e^{-\beta \left(\frac{p^{2}}{2m} \right)} dp = \int_{0}^{\infty} \frac{V 4\pi p^{2}}{h^{3}} e^{-\beta \left(\frac{p^{2}}{2m} \right)} dp = V \left(\frac{2\pi m k T}{h^{2}} \right)^{\frac{3}{2}} \\ Z_{1}^{rot} &= 1 + 3 e^{-\beta \frac{\hbar^{2}}{l}} + 5 e^{-\beta \frac{3\hbar^{2}}{l}} + \cdots \\ \text{Let } x &= r(r+1), \qquad dx = (2r+1) \, dr \\ Z_{1}^{rot} &\approx \int_{0}^{\infty} (2r+1) \, e^{-\beta \frac{\hbar^{2}}{2l} r(r+1)} dr = \int_{0}^{\infty} e^{-\beta \frac{\hbar^{2}}{2l} x} \, dx \Rightarrow \frac{2l}{\beta \hbar^{2}} \\ \text{for } \frac{\varepsilon_{r}}{kT} \ll 1. \end{split}$$

Question 4(i)

$$\begin{split} Z(V,T) &= \sum_{n_1} \sum_{n_2} \dots \left(e^{-\beta n_1 \varepsilon_1} e^{-\beta n_2 \varepsilon_2} e^{-\beta n_3 \varepsilon_3} \dots \right) \\ &= \sum_{n_1} e^{-\beta n_1 \varepsilon_1} \sum_{n_2} e^{-\beta n_2 \varepsilon_2} \sum_{n_3} e^{-\beta n_3 \varepsilon_3} \dots \\ &= \frac{1}{1 - e^{-\beta \varepsilon_1}} \frac{1}{1 - e^{-\beta \varepsilon_2}} \frac{1}{1 - e^{-\beta \varepsilon_3}} \dots \\ &= \prod_{n=1}^{\infty} \frac{1}{1 - e^{-\beta \varepsilon_r}} \end{split}$$

Question 4(ii)

$$\begin{split} \bar{n}_r &= -\frac{1}{\beta} \frac{\partial (\ln Z)}{\partial \varepsilon_r} \\ &= -\frac{1}{\beta} \frac{\partial}{\partial \varepsilon_r} \left[\ln \left(\frac{1}{1 - e^{-\beta \varepsilon_1}} \frac{1}{1 - e^{-\beta \varepsilon_2}} \frac{1}{1 - e^{-\beta \varepsilon_3}} \dots \right) \right] \\ &= \frac{1}{\beta} \frac{\partial}{\partial \varepsilon_r} \sum_i \ln (1 - e^{-\beta \varepsilon_i}) \\ &= \frac{1}{\beta} \frac{\beta e^{-\beta \varepsilon_r}}{1 - e^{-\beta \varepsilon_r}} \\ &= \frac{1}{e^{\beta \varepsilon_r} - 1} \end{split}$$

Question 5(i)

$$\bar{n}_i = \frac{1}{e^{\beta(\varepsilon_i - \mu)} \pm 1}$$

where + is for Fermions, - is for Bosons.

$$\frac{1}{\beta} \left(\frac{\partial \bar{n}_i}{\partial \mu} \right) = \frac{e^{\beta(\varepsilon_i - \mu)}}{(e^{\beta(\varepsilon_i - \mu)} \pm 1)^2} = (\Delta n_i)^2$$

Question 5(ii)

Fluctuation is $\frac{\Delta n_i}{\bar{n}_i}$, actually almost the same for both statistic, since $e^{\frac{\beta(\varepsilon_i-\mu)}{2}}$ is quite large. Except the case of an extremely degenerate FD gas for which the energetically lowest lying single particle states are occupied, $\bar{n}_i \approx 1$.

Question 6

PART II

Question 1(i)

$$S = k \ln \Omega$$
= $k \ln \left(\frac{N!}{n! (N-n)!} \right)$
= $k [\ln N! - \ln n! - \ln(N-n)!]$
= $k [N \ln N - n \ln n - (N-n) \ln(N-n)]$

$$E = -n\mu B + (N - n)\mu B = (N - 2n)\mu B$$

$$\frac{1}{T} = \frac{\partial S}{\partial E} = \frac{\partial S}{\partial n} \frac{\partial n}{\partial E} = -\frac{k}{2\mu B} \ln\left(\frac{N - n}{n}\right)$$

Question 1(ii)

$$\mu B = \frac{E}{N - 2n}$$

$$\therefore \frac{1}{T} = \frac{k}{2E}(N - 2n) \ln \left(\frac{n}{N - n}\right)$$

$$E = kT \left(\frac{N}{2} - n\right) \ln \left(\frac{n}{N - n}\right)$$

Question 1(iii)

$$Z_{1} = e^{-\beta\mu B} + e^{\beta\mu B} = 2\cosh\beta\mu B = 2\cosh\frac{\mu B}{kT}$$
$$E = N\frac{\partial(\ln Z_{1})}{\partial\beta} - \frac{N\sinh\frac{\mu B}{kT}}{\cosh\frac{\mu B}{kT}} = N\mu B\tanh\frac{\mu B}{kT}$$

It is very different. For energy of 1 particle, it depends on μB . When there are many particles, it depends on kT.

Question 2(i)

Consider a system immersed in a heat bath. We assume particle number N and volume V are fixed. The system microstates in equilibrium labeled 1,2,...,r and energies $E_1,E_2,...,E_r$. We let E_0 be the total energy of the system. The probability,

$$p_{r} = \frac{\Omega_{2}(E_{0} - E_{r})}{\sum_{r} \Omega_{2}(E_{0} - E_{r})}$$

$$= \text{const. } e^{\frac{S_{2}(E_{0} - E_{r})}{k}}$$

$$= \text{const. } e^{\frac{1}{k} \left[S_{2}(E_{0}) - \frac{\partial S_{2}(E_{0})}{\partial E_{0}} E_{r}\right]}$$

$$= \text{const. } e^{\frac{1}{k}S_{2}(E_{0}) - \beta E_{r}}$$

$$= \frac{e^{-\beta E_{r}}}{Z}$$

$$= \frac{e^{-\beta E_{r}}}{\sum_{r} e^{-\beta E_{r}}}$$

$$\begin{split} & \therefore Z = \sum_r e^{-\beta E_r} \\ & \bar{E} = \sum_r p_r E_r = \sum_r \frac{e^{-\beta E_r}}{Z} E_r = -\sum_r p_r \left[\frac{\partial (\ln Z_r)}{\partial \beta} \right] = -\frac{\partial (\ln Z)}{\partial \beta} \\ & (\Delta E)^2 = \overline{E^2} - \bar{E}^2 = \frac{\partial^2 \ln Z}{\partial \beta^2} - \left(\frac{\partial \ln Z}{\partial \beta} \right)^2 + \left(\frac{\partial \ln Z}{\partial \beta} \right)^2 = \frac{\partial^2 \ln Z}{\partial \beta^2} \end{split}$$

Question 2(ii)

$$\bar{E} = -\frac{\partial \ln Z}{\partial \beta} = -\frac{\partial}{\partial \beta} \left[-\ln N! + N \ln V - \frac{3}{2} N \ln \left(\frac{\hbar \beta}{2\pi m} \right) \right] = \frac{3}{2} N k T$$

$$(\Delta E)^2 = \frac{\partial}{\partial \beta} \left(\frac{\partial \ln Z}{\partial \beta} \right) = \frac{\partial}{\partial \beta} \left(-\frac{3}{2} \frac{N}{\beta} \right) = \frac{3}{2} N k^2 T^2$$

$$\frac{\Delta E}{\bar{E}} = \frac{\sqrt{\frac{3N}{2}} k T}{\frac{3}{2} k T} = \sqrt{\frac{2}{3} \frac{1}{\sqrt{N}}} \propto \frac{1}{\sqrt{N}}$$

Relative fluctuation is inversely proportional to the root of its size. $N \sim 10^{23}$, fluctuation is extremely small, so the energy of macroscopic body in heat bath, for practical purposes, is completely determined.

Question 3(i)

At equilibrium,

$$2\mu_{CO_2} = 2\mu_{CO} + \mu_{O_2}$$

Question 3(ii)

 $dN_i \propto v_i$

Where v_i is the stoichiometric coefficient.

$$dN(CO_2): dN(CO): dN(O_2) = 2: -2: -1$$

Question 3(iii)

Law of mass action,

$$\prod_{i} (c_i)^{v_i} = K_c(T)$$

$$c_i = \frac{N_i}{V}$$

$$K_c(T) = \frac{c_{CO}^2 c_{O_2}}{c_{CO_2}^2}$$

$$2(1 - f)CO_2 \rightleftharpoons f(2CO + O_2)$$

$$\rightleftharpoons 2fCO + fO_2$$

At a fixed, temperature, the total pressure P,

$$P = P_{CO_2} + P_{CO} + P_{O_2}$$

$$P_{CO_2} = \frac{2(1-f)}{2(1-f) + 2f + f} P = \frac{2(1-f)}{2+f} P$$

$$P_{CO} = \frac{2f}{2+f} P$$

$$P_{O_2} = \frac{f}{2+f} P$$

$$K_{P}(T) = \prod_{i} (P_{i})^{v_{i}} = \frac{\left(\frac{2f}{2+f}\right)^{2} \left(\frac{f}{2+f}\right) P^{3}}{\left[\frac{2(1-f)}{2+f}\right] P^{2}} = \frac{f^{3}}{2(1-f)^{2}} P$$

$$P = 2K_{P} \frac{(1-f)^{2}}{f^{3}} = 2K_{P} \left(\frac{1}{f^{3}} - \frac{2}{f^{2}} + \frac{1}{f}\right)$$

When P increases, f decreases.

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