Question 1(a)

The Hamiltonian for the hydrogen atom, one proton and one electron with masses m_p and m_e ,

$$H = \frac{p_p^2}{2m_p} + \frac{p_e^2}{2m_e} + V(\vec{x}_p - \vec{x}_e)$$

We transform them into 2 equivalent particles M and m:



We let

$$M=m_p+m_epprox m_p, \qquad m=rac{m_pm_e}{m_p+m_e}pprox m_e$$

Then we get

$$M\vec{X} = m_p \vec{x}_p + m_e \vec{x}_e, \qquad \vec{x} = \vec{x}_p - \vec{x}_e$$

The momentum,

$$\vec{P} = M\dot{\vec{X}}, \qquad \vec{p} = m\dot{\vec{x}}$$

Rearranging, we get

$$\vec{p}_p = \frac{m}{m_e} \vec{P} + \vec{p} \approx \vec{P} + \vec{p}, \qquad \vec{p}_e = \frac{m}{m_p} \vec{P} - \vec{p} \approx -\vec{p}$$

Using the above relations, the Hamiltonian now becomes

$$\begin{split} H &= \frac{1}{2m_{p}} \left(\frac{m}{m_{e}} \vec{P} + \vec{p} \right)^{2} + \frac{1}{2m_{e}} \left(\frac{m}{m_{p}} \vec{P} - \vec{p} \right)^{2} + V(\vec{x}) \\ &= \frac{1}{2m_{p}} \left[\frac{m^{2}}{m_{e}^{2}} P^{2} + \frac{m}{m_{e}} (\vec{P} \cdot \vec{p} + \vec{p} \cdot \vec{P}) + p^{2} \right] + \frac{1}{2m_{e}} \left[\frac{m^{2}}{m_{p}^{2}} P^{2} + \frac{m}{m_{e}} (\vec{P} \cdot \vec{p} - \vec{p} \cdot \vec{P}) + p^{2} \right] + V(\vec{x}) \\ &= \frac{P^{2}}{2(m_{p} + m_{e})} + \frac{p^{2}}{2} \left(\frac{m_{p} + m_{e}}{m_{p} m_{e}} \right) + V(\vec{x}) \\ &= \frac{P^{2}}{2M} + \frac{p^{2}}{2m} + V(\vec{x}) \end{split}$$

Where M, P describes the motion of the center of mass, and m, p describes the relative motion of the electron and the proton. [shown]

Question 1(b)

We know that

$$\frac{\partial r}{\partial x_i} = \frac{x_i}{r}, \qquad \frac{\partial \theta}{\partial x_i} = \frac{1}{r \sin \theta} \left(\frac{x_i}{r} \cos \theta - \delta_{i3} \right), \qquad \frac{\partial \phi}{\partial x_i} = \frac{1}{r \sin \theta} \left(\delta_{i2} \cos \phi - \delta_{i1} \sin \phi \right)$$

Therefore we have

$$\begin{split} l_3 &= x_1 p_2 - x_2 p_1 \\ &= \frac{\hbar}{i} \left(x_1 \frac{\partial}{\partial x_2} - x_2 \frac{\partial}{\partial x_1} \right) \\ &= \frac{\hbar}{i} \left[x_1 \left(\frac{\partial r}{\partial x_2} \frac{\partial}{\partial r} + \frac{\partial \theta}{\partial x_2} \frac{\partial}{\partial \theta} + \frac{\partial \phi}{\partial x_2} \frac{\partial}{\partial \phi} \right) - x_2 \left(\frac{\partial r}{\partial x_1} \frac{\partial}{\partial r} + \frac{\partial \theta}{\partial x_1} \frac{\partial}{\partial \theta} + \frac{\partial \phi}{\partial x_1} \frac{\partial}{\partial \phi} \right) \right] \end{split}$$

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$$\begin{split} &=\frac{\hbar}{i}\bigg[x_1\bigg(\frac{x_2}{r}\frac{\partial}{\partial r}+\frac{x_2\cos\theta}{r^2\sin\theta}\frac{\partial}{\partial\theta}+\frac{\cos\phi}{r\sin\theta}\frac{\partial}{\partial\phi}\bigg)-x_2\bigg(\frac{x_1}{r}\frac{\partial}{\partial r}+\frac{x_1\cos\theta}{r^2\sin\theta}\frac{\partial}{\partial\theta}-\frac{\sin\phi}{r\sin\theta}\frac{\partial}{\partial\phi}\bigg)\bigg]\\ &=\frac{\hbar}{i}\bigg(\frac{x_1\cos\phi}{r\sin\theta}\frac{\partial}{\partial\phi}+\frac{x_2\cos\phi}{r\sin\theta}\frac{\partial}{\partial\phi}\bigg)\\ &=\frac{\hbar}{i}\bigg(\cos^2\phi\frac{\partial}{\partial\phi}+\sin^2\phi\frac{\partial}{\partial\phi}\bigg)\\ &=\frac{\hbar}{i}\frac{\partial}{\partial\phi}\end{split}$$

Let ψ be an arbitrary function, hence we have

$$\begin{split} \left[\phi, -i\hbar\frac{\partial}{\partial\phi}\right]\psi &= -i\hbar\phi\frac{\partial\psi}{\partial\phi} + i\hbar\frac{\partial}{\partial\phi}(\phi\psi) = -i\hbar\phi\frac{\partial\psi}{\partial\phi} + i\hbar\psi + i\hbar\phi\frac{\partial\psi}{\partial\phi} = i\hbar\psi \\ & \div\left[\phi, -i\hbar\frac{\partial}{\partial\phi}\right] = i\hbar \quad [\text{shown}] \end{split}$$

Suppose $[\phi, l_3] = i\hbar$, then $\langle l, m | [\phi, l_3] | l, m \rangle = i\hbar \langle l, m | l, m \rangle$ $LHS = \langle l, m | [\phi, l_3] | l, m \rangle = \langle l, m | (\phi l_3 - l_3 \phi) | l, m \rangle = \langle l, m | (\phi m\hbar - m\hbar \phi) | l, m \rangle = 0$ $RHS = i\hbar \langle l, m | l, m \rangle \neq 0$

$$\therefore [\phi, l_3] \neq i\hbar$$

In coordinate representation,

$$l_{+} = \hbar e^{i\phi} \left(\frac{\partial}{\partial \theta} + i \cot \theta \frac{\partial}{\partial \phi} \right), \qquad l_{3} = \frac{\hbar}{i} \frac{\partial}{\partial \phi}$$

With
$$l_3 Y_l^l(\theta, \phi) = l\hbar Y_l^l(\theta, \phi)$$
, we have
$$\frac{\hbar}{i} \frac{\partial}{\partial \phi} Y_l^l(\theta, \phi) = l\hbar Y_l^l(\theta, \phi), \qquad \Rightarrow \quad Y_l^l(\theta, \phi) = f(\theta) e^{il\phi}$$

Where $f(\theta)$ is an arbitrary function of θ . Substitute into the representation of l_+ ,

$$l_{+}Y_{l}^{l}(\theta,\phi) = 0$$

$$\hbar e^{i\phi} \left(\frac{\partial}{\partial \theta} + i\cot\theta \frac{\partial}{\partial \phi}\right) f(\theta) e^{il\phi} = 0$$

$$\left(\frac{\partial}{\partial \theta} - l\cot\theta\right) f(\theta) = 0$$

$$\frac{\partial f}{\partial \theta} = lf\cot\theta$$

$$\int \frac{1}{f} df = \int l\cot\theta d\theta$$

$$\ln f = l\ln(\sin\theta) + c$$

$$f = N\sin^{l}\theta$$

$$\therefore Y_l^l(\theta, \phi) = N \sin^l \theta \, e^{il\phi} \quad [\text{shown}]$$

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Question 2(a)

$$H = \frac{mv^2}{2} + V = \frac{(\vec{p} - q\vec{A})^2}{2m} + qA_0$$

Using the formula given, we have

Formula given, we have
$$\frac{d}{dt}\langle x_i \rangle = \frac{i}{\hbar} \langle [H, x_i] \rangle + \langle \frac{\partial x_i}{\partial t} \rangle$$

$$= \frac{i}{\hbar} \langle \left[\frac{(\vec{p} - q\vec{A})^2}{2m} + qA_0, x_i \right] \rangle$$

$$= \frac{i}{2m\hbar} \langle ([p - qA)_j (p - qA)_j, x_i] \rangle$$

$$= \frac{i}{2m\hbar} \langle (p - qA)_j [(p - qA)_j, x_i] + [(p - qA)_j, x_i] (p - qA)_j \rangle$$

$$= \frac{i}{2m\hbar} \langle (p - qA)_j (-i\hbar\delta_{ij}) + (-i\hbar\delta_{ij}) (p - qA)_j \rangle$$

$$= \frac{i}{2m\hbar} \langle -2i\hbar (p - qA)_j \rangle$$

$$= \frac{1}{m} \langle (p - qA)_j \rangle$$

$$\therefore \frac{d}{dt} \langle \vec{x} \rangle = \frac{1}{m} \langle (\vec{p} - q\vec{A}) \rangle$$

Question 2(b)(i)

We know that $\vec{B} = \vec{\nabla} \times \vec{A}$. Then also,

$$\vec{A} = \frac{B_0}{2}(x_1\hat{j} - x_2\hat{i}) = \frac{1}{2}\vec{B} \times \vec{x}, \qquad \vec{x} = x_1\hat{i} + x_2\hat{j} + x_3\hat{k}, \qquad \vec{B} = B_0\hat{k}$$

We substitute \vec{A} into the formula,

$$\vec{\nabla} \times \vec{A} = \vec{\nabla} \times \left(\frac{1}{2}\vec{B} \times \vec{x}\right)$$

$$= \frac{1}{2} \left[(\vec{x} \cdot \vec{\nabla})\vec{B} - (\vec{B} \cdot \vec{\nabla})\vec{x} + \vec{B}(\vec{\nabla} \cdot \vec{x}) - \vec{x}(\vec{\nabla} \cdot \vec{B}) \right]$$

$$= \frac{1}{2} \left[0 - B_0 \hat{k} + 3B_0 \hat{k} - 0 \right]$$

$$= B_0 \hat{k}$$

$$= \vec{B}$$

Next, we know that $\vec{E} = -2Cx_3\hat{k}$. Substituting A_0 into the formula

$$-\vec{\nabla}A_0 + \frac{\partial \vec{A}}{\underbrace{\partial t}} = -\vec{\nabla}(Cx_3^2) = -2Cx_3\hat{k}$$

$$\vec{A} = \frac{B_0}{2}(x_1\hat{j} - x_2\hat{i}), \quad A_0 = Cx_3^2 \text{ [shown]}$$

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Question 2(b)(ii)

The Hamiltonian now becomes

$$\begin{split} H &= \frac{\left(\vec{p} - q\vec{A}\right)^{2}}{2m} + qA_{0} \\ &= \frac{\left[\vec{p} - q\frac{B_{0}}{2}(x_{1}\hat{\jmath} - x_{2}\hat{\imath})\right]^{2}}{2m} + qCx_{3}^{2} \\ &= \frac{1}{2m} \left[p^{2} + \frac{q^{2}B_{0}^{2}}{4}(x_{1}^{2} + x_{2}^{2}) - qB_{0}l_{3}\right] + qCx_{3}^{2} \\ &= \left[\frac{1}{2m}(p_{1}^{2} + p_{2}^{2}) + \frac{q^{2}B_{0}^{2}}{8m}(x_{1}^{2} + x_{2}^{2})\right] + \left[\frac{p_{3}^{2}}{2m} + qCx_{3}^{2}\right] - \frac{qB_{0}l_{3}}{2m} \\ &= \left[\frac{1}{2m}(p_{1}^{2} + p_{2}^{2}) + \frac{1}{2}m\underbrace{\left(\frac{q^{2}B_{0}^{2}}{4m^{2}}\right)}_{=\omega_{1}^{2}}(x_{1}^{2} + x_{2}^{2})\right] + \left[\frac{p_{3}^{2}}{2m} + \frac{1}{2}m\underbrace{\left(\frac{2qC}{m}\right)}_{=\omega_{2}^{2}}x_{3}^{2}\right] - \frac{qB_{0}l_{3}}{2m} \end{split}$$

Which is a combination of 2 harmonic oscillators. We then find the energies to be

$$E(n_1, n_2) = \left(n_1 + \frac{1}{2}\right)\hbar\omega_1 + \left(n_2 + \frac{1}{2}\right)\hbar\omega_2 \quad [\text{shown}]$$

When $\vec{E}=0$, we get the energy $E=\left(n_1+\frac{1}{2}\right)\hbar\omega_1$, which is the Landau level of the particle.

Question 3(a)

$$\begin{split} \mathcal{R}_{\vec{u}}(\varepsilon) &= 1 + \varepsilon \vec{u} \times + \frac{1}{2} \varepsilon^2 \vec{u} \times \vec{u} \times + O(\varepsilon^3) \\ \mathcal{R}_{\vec{v}} \mathcal{R}_{\vec{u}} &= \left(1 + \varepsilon \vec{v} \times + \frac{1}{2} \varepsilon^2 \vec{v} \times \vec{v} \times + O(\varepsilon^3) \right) \left(1 + \varepsilon \vec{u} \times + \frac{1}{2} \varepsilon^2 \vec{u} \times \vec{u} \times + O(\varepsilon^3) \right) \\ &= 1 + \varepsilon (\vec{u} + \vec{v}) \times + \varepsilon^2 \vec{v} \times \vec{u} \times + \frac{1}{2} \varepsilon^2 \vec{u} \times \vec{u} \times + \frac{1}{2} \varepsilon^2 \vec{v} \times \vec{v} \times + O(\varepsilon^3) \\ \mathcal{R}_{\vec{u}}^{-1} \mathcal{R}_{\vec{v}} \mathcal{R}_{\vec{u}} &= \left(1 - \varepsilon \vec{u} \times + \frac{1}{2} \varepsilon^2 \vec{u} \times \vec{u} \times + O(\varepsilon^3) \right) \left[1 + \varepsilon (\vec{u} + \vec{v}) \times + \varepsilon^2 \vec{v} \times \vec{u} \times + \frac{1}{2} \varepsilon^2 \vec{u} \times \vec{v} \times + O(\varepsilon^3) \right] \\ &= 1 + \varepsilon \vec{v} \times + \varepsilon^2 \vec{v} \times \vec{u} \times - \varepsilon^2 \vec{u} \times \vec{v} \times + \frac{1}{2} \varepsilon^2 \vec{v} \times \vec{v} \times + O(\varepsilon^3) \\ \mathcal{R}_{\vec{v}}^{-1} \mathcal{R}_{\vec{u}}^{-1} \mathcal{R}_{\vec{v}} \mathcal{R}_{\vec{u}} &= \left(1 - \varepsilon \vec{v} \times + \frac{1}{2} \varepsilon^2 \vec{v} \times \vec{v} \times + O(\varepsilon^3) \right) \left(1 + \varepsilon \vec{v} \times + \varepsilon^2 \vec{v} \times \vec{u} \times - \varepsilon^2 \vec{u} \times \vec{v} \times + \frac{1}{2} \varepsilon^2 \vec{v} \times \vec{v} \times + O(\varepsilon^3) \right) \\ &= 1 + \varepsilon^2 \vec{v} \times \vec{u} \times - \varepsilon^2 \vec{u} \times \vec{v} \times + O(\varepsilon^3) \\ \mathcal{R}_{\vec{v}}^{-1} \mathcal{R}_{\vec{u}}^{-1} \mathcal{R}_{\vec{v}} \mathcal{R}_{\vec{u}} \vec{x} &= \vec{x} + \varepsilon^2 \vec{v} \times (\vec{u} \times \vec{x}) - \varepsilon^2 \vec{u} \times (\vec{v} \times \vec{x}) + O(\varepsilon^3) \\ &= \vec{x} + \varepsilon^2 [\vec{u} (\vec{v} \cdot \vec{x}) - \vec{v} (\vec{u} \cdot \vec{x})] + O(\varepsilon^3) \\ &= \vec{x} + \varepsilon^2 \vec{x} \times (\vec{u} \times \vec{v}) + O(\varepsilon^3) \\ &= \vec{x} - \varepsilon^2 \vec{w} \times \vec{x} + O(\varepsilon^3) \end{split}$$

But we know that $\mathcal{R}_{\overrightarrow{w}}(-\varepsilon^2) = \overrightarrow{x} - \varepsilon^2 \overrightarrow{w} \times \overrightarrow{x} + O(\varepsilon^4)$. $\therefore \mathcal{R}_{\overrightarrow{v}}^{-1} \mathcal{R}_{\overrightarrow{u}}^{-1} \mathcal{R}_{\overrightarrow{v}} \mathcal{R}_{\overrightarrow{u}}$ differs from $\mathcal{R}_{\overrightarrow{w}}(-\varepsilon^2)$ only by terms of higher order than ε^2 . [shown]

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Question 3(b)

$$J_1^2|j_1, m_1\rangle = \frac{3}{4}\hbar^2|j_1, m_1\rangle = j_1(j_1 + 1)\hbar^2|j_1, m_1\rangle$$
$$j_1^2 + j_1 - \frac{3}{4} = 0 \quad \Rightarrow \quad j_1 = -\frac{3}{2}, \frac{1}{2}$$

$$J_2^2|j_2, m_2\rangle = 6\hbar^2|j_2, m_2\rangle = j_2(j_2 + 1)\hbar^2|j_2, m_2\rangle$$

 $j_2^2 + j_2 - 6 = 0 \implies j_2 = -3,2$

So we have $j_1 = \frac{1}{2}$, $j_2 = 2$. Using the Clebsch-Gordan Coefficient table, and looking for $|j,m\rangle \Rightarrow |j_2,j_1,m_2,m_1\rangle$, we get a set of simultaneous normalized eigenvectors,

$$|j,m\rangle \Rightarrow |j_{2},j_{1},m_{2},m_{1}\rangle, \text{ we get a set of simult}$$

$$|\frac{5}{2},\frac{5}{2}\rangle = |2,\frac{1}{2},2,\frac{1}{2}\rangle$$

$$|\frac{5}{2},\frac{3}{2}\rangle = \frac{1}{\sqrt{5}}|2,\frac{1}{2},2,-\frac{1}{2}\rangle + \frac{2}{\sqrt{5}}|2,\frac{1}{2},1,\frac{1}{2}\rangle$$

$$|\frac{3}{2},\frac{3}{2}\rangle = \frac{2}{\sqrt{5}}|2,\frac{1}{2},2,-\frac{1}{2}\rangle - \frac{1}{\sqrt{5}}|2,\frac{1}{2},1,\frac{1}{2}\rangle$$

$$|\frac{5}{2},\frac{1}{2}\rangle = \sqrt{\frac{2}{5}}|2,\frac{1}{2},1,-\frac{1}{2}\rangle + \sqrt{\frac{3}{5}}|2,\frac{1}{2},0,\frac{1}{2}\rangle$$

$$|\frac{3}{2},\frac{1}{2}\rangle = \sqrt{\frac{3}{5}}|2,\frac{1}{2},1,-\frac{1}{2}\rangle - \sqrt{\frac{2}{5}}|2,\frac{1}{2},0,\frac{1}{2}\rangle$$

$$|\frac{5}{2},-\frac{1}{2}\rangle = \sqrt{\frac{3}{5}}|2,\frac{1}{2},0,-\frac{1}{2}\rangle + \sqrt{\frac{2}{5}}|2,\frac{1}{2},-1,\frac{1}{2}\rangle$$

$$|\frac{3}{2},-\frac{1}{2}\rangle = \sqrt{\frac{2}{5}}|2,\frac{1}{2},0,-\frac{1}{2}\rangle - \sqrt{\frac{3}{5}}|2,\frac{1}{2},-1,\frac{1}{2}\rangle$$

$$|\frac{5}{2},-\frac{3}{2}\rangle = \frac{2}{\sqrt{5}}|2,\frac{1}{2},-1,-\frac{1}{2}\rangle + \frac{1}{\sqrt{5}}|2,\frac{1}{2},-2,\frac{1}{2}\rangle$$

$$|\frac{3}{2},-\frac{3}{2}\rangle = \frac{1}{\sqrt{5}}|2,\frac{1}{2},-1,-\frac{1}{2}\rangle - \frac{2}{\sqrt{5}}|2,\frac{1}{2},-2,\frac{1}{2}\rangle$$

$$|\frac{5}{2},-\frac{5}{2}\rangle = |2,\frac{1}{2},-2,-\frac{1}{2}\rangle$$

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Question 4(a)

Let c_m be the number of states $|j,m\rangle$ for a fixed $m=m_1+m_2$, and d_m be the number of states $|j,m\rangle$ for a fixed m and j. With $|m| \le j$, and assume $m \ge 0$,

$$c_m = \sum_{m \geq 0} d_m = d_m + d_{m+1} + \cdots$$

$$c_{m+1} = d_{m+1} + d_{m+2} + \cdots$$

$$c_m - c_{m+1} = d_m$$

m	m_1	m_2	c_m
$j_1 + j_2$	j_1	j_2	1
$j_1 + j_2 - 1$	$j_1 - 1$ j_1	$j_2 \\ j_2 - 1$	2
$j_1 + j_2 - 2$	$j_1 - 2$ $j_1 - 1$ j_1	j_2 $j_2 - 1$ $j_2 - 2$	3
•••			
$j_1 + j_2 - n$			n+1

So if you know c_m , we can get d_m . As in the case of a simple harmonic oscillator, we know that $d_j = 1$ if j is allowed, and $d_j = 0$ if j is not allowed. The value of n to terminate the counting is $-2j_2$ if $j_2 < j_1$, or $-2j_1$ if $j_1 < j_2$.

$$: j = j_1 + j_2, j_1 + j_2 - 1, j_1 + j_2 - 2, ..., |j_1 - j_2|$$
 [shown]

Question 4(b)

$$E = \frac{\hbar^2 \pi^2}{2mb^2} (k^2 + l^2 + m^2 + n^2)$$

(i) For distinguishable identical particles,

Ground state,
$$\psi = \frac{4}{b^2} \sin\left(\frac{k\pi x_1}{b}\right) \sin\left(\frac{l\pi x_2}{b}\right) \sin\left(\frac{m\pi x_3}{b}\right) \sin\left(\frac{n\pi x_4}{b}\right)$$

$$\psi = \frac{2\hbar^2 \pi^2}{mb^2}$$

$$\psi = \frac{4}{b^2} \sin\left(\frac{\pi x_1}{b}\right) \sin\left(\frac{\pi x_2}{b}\right) \sin\left(\frac{\pi x_3}{b}\right) \sin\left(\frac{\pi x_4}{b}\right)$$
1st excited state,
$$E = \frac{7\hbar^2 \pi^2}{2mb^2}$$

$$\psi = \frac{4}{b^2} \sin\left(\frac{\pi x_1}{b}\right) \sin\left(\frac{\pi x_2}{b}\right) \sin\left(\frac{\pi x_3}{b}\right) \sin\left(\frac{2\pi x_4}{b}\right)$$

$$\psi = \frac{4}{b^2} \sin\left(\frac{\pi x_1}{b}\right) \sin\left(\frac{\pi x_2}{b}\right) \sin\left(\frac{2\pi x_3}{b}\right) \sin\left(\frac{\pi x_4}{b}\right)$$

$$\psi = \frac{4}{b^2} \sin\left(\frac{\pi x_1}{b}\right) \sin\left(\frac{2\pi x_2}{b}\right) \sin\left(\frac{\pi x_3}{b}\right) \sin\left(\frac{\pi x_4}{b}\right)$$

$$\psi = \frac{4}{b^2} \sin\left(\frac{2\pi x_1}{b}\right) \sin\left(\frac{2\pi x_2}{b}\right) \sin\left(\frac{\pi x_3}{b}\right) \sin\left(\frac{\pi x_4}{b}\right)$$

$$\psi = \frac{4}{b^2} \sin\left(\frac{2\pi x_1}{b}\right) \sin\left(\frac{\pi x_2}{b}\right) \sin\left(\frac{\pi x_3}{b}\right) \sin\left(\frac{\pi x_4}{b}\right)$$

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(ii) For identical bosons,

$$\begin{split} \overline{\psi} &= \frac{1}{b^2 \sqrt{6}} \left[\sin \left(\frac{k\pi x_1}{b} \right) \sin \left(\frac{l\pi x_2}{b} \right) \sin \left(\frac{m\pi x_3}{b} \right) \sin \left(\frac{n\pi x_4}{b} \right) + \sin \left(\frac{k\pi x_2}{b} \right) \sin \left(\frac{l\pi x_3}{b} \right) \sin \left(\frac{m\pi x_4}{b} \right) \sin \left(\frac{n\pi x_1}{b} \right) \\ &+ \sin \left(\frac{k\pi x_3}{b} \right) \sin \left(\frac{l\pi x_4}{b} \right) \sin \left(\frac{m\pi x_1}{b} \right) \sin \left(\frac{n\pi x_2}{b} \right) + \sin \left(\frac{k\pi x_4}{b} \right) \sin \left(\frac{l\pi x_1}{b} \right) \sin \left(\frac{m\pi x_3}{b} \right) \\ &+ \sin \left(\frac{k\pi x_2}{b} \right) \sin \left(\frac{l\pi x_1}{b} \right) \sin \left(\frac{m\pi x_2}{b} \right) \sin \left(\frac{n\pi x_2}{b} \right) + \sin \left(\frac{k\pi x_4}{b} \right) \sin \left(\frac{l\pi x_3}{b} \right) \sin \left(\frac{n\pi x_4}{b} \right) \sin \left(\frac{n\pi x_3}{b} \right) \\ &+ \sin \left(\frac{k\pi x_3}{b} \right) \sin \left(\frac{l\pi x_4}{b} \right) \sin \left(\frac{m\pi x_2}{b} \right) \sin \left(\frac{n\pi x_1}{b} \right) + \sin \left(\frac{k\pi x_4}{b} \right) \sin \left(\frac{l\pi x_2}{b} \right) \sin \left(\frac{m\pi x_4}{b} \right) \sin \left(\frac{n\pi x_3}{b} \right) \\ &+ \sin \left(\frac{k\pi x_1}{b} \right) \sin \left(\frac{l\pi x_3}{b} \right) \sin \left(\frac{m\pi x_2}{b} \right) \sin \left(\frac{n\pi x_4}{b} \right) + \sin \left(\frac{k\pi x_4}{b} \right) \sin \left(\frac{l\pi x_2}{b} \right) \sin \left(\frac{m\pi x_4}{b} \right) \sin \left(\frac{n\pi x_2}{b} \right) \\ &+ \sin \left(\frac{k\pi x_2}{b} \right) \sin \left(\frac{l\pi x_4}{b} \right) \sin \left(\frac{m\pi x_1}{b} \right) \sin \left(\frac{n\pi x_3}{b} \right) + \sin \left(\frac{k\pi x_4}{b} \right) \sin \left(\frac{l\pi x_4}{b} \right) \sin \left(\frac{m\pi x_3}{b} \right) \sin \left(\frac{n\pi x_2}{b} \right) \\ &+ \sin \left(\frac{k\pi x_4}{b} \right) \sin \left(\frac{l\pi x_3}{b} \right) \sin \left(\frac{m\pi x_4}{b} \right) \sin \left(\frac{n\pi x_2}{b} \right) + \sin \left(\frac{k\pi x_4}{b} \right) \sin \left(\frac{l\pi x_4}{b} \right) \sin \left(\frac{n\pi x_4}{b} \right) \sin \left(\frac{n\pi x_4}{b} \right) \\ &+ \sin \left(\frac{k\pi x_4}{b} \right) \sin \left(\frac{l\pi x_3}{b} \right) \sin \left(\frac{m\pi x_4}{b} \right) \sin \left(\frac{n\pi x_2}{b} \right) + \sin \left(\frac{k\pi x_4}{b} \right) \sin \left(\frac{l\pi x_4}{b} \right) \sin \left(\frac{n\pi x_4}{b} \right) \sin \left(\frac{n\pi x_4}{b} \right) \\ &+ \sin \left(\frac{k\pi x_4}{b} \right) \sin \left(\frac{l\pi x_3}{b} \right) \sin \left(\frac{m\pi x_4}{b} \right) \sin \left(\frac{n\pi x_4}{b} \right) \sin \left(\frac{l\pi x_4}{b} \right) \sin \left(\frac{n\pi x_4}{b} \right) + \sin \left(\frac{l\pi x_4}{b} \right) \sin \left(\frac{l\pi x_4}{b} \right) \sin \left(\frac{n\pi x_4}{b} \right) \sin \left(\frac{n\pi x_4}{b} \right) \\ &+ \sin \left(\frac{k\pi x_4}{b} \right) \sin \left(\frac{l\pi x_4}{b} \right) \sin \left(\frac{m\pi x_4}{b} \right) \sin \left(\frac{n\pi x_4}{b} \right) \sin \left(\frac{l\pi x_4}{b} \right) \sin \left(\frac{n\pi x_4}{b} \right) \sin \left(\frac{n\pi x_4}{b} \right) \\ &+ \sin \left(\frac{k\pi x_4}{b} \right) \sin \left(\frac{l\pi x_4}{b} \right) \sin \left(\frac{m\pi x_4}{b} \right) \sin \left(\frac{l\pi x_4}{b$$

Ground state,
$$E = \frac{2\hbar^2 \pi^2}{mb^2}$$

$$\psi = \frac{24}{b^2 \sqrt{6}} \sin\left(\frac{\pi x_1}{b}\right) \sin\left(\frac{\pi x_2}{b}\right) \sin\left(\frac{\pi x_3}{b}\right) \sin\left(\frac{\pi x_4}{b}\right)$$
1st excited state,
$$E = \frac{7\hbar^2 \pi^2}{2mb^2}$$

$$\psi = \frac{\sqrt{6}}{b^2} \left[\sin\left(\frac{\pi x_1}{b}\right) \sin\left(\frac{\pi x_2}{b}\right) \sin\left(\frac{\pi x_3}{b}\right) \sin\left(\frac{2\pi x_4}{b}\right) + \sin\left(\frac{\pi x_2}{b}\right) \sin\left(\frac{\pi x_3}{b}\right) \sin\left(\frac{\pi x_1}{b}\right) \sin\left(\frac{\pi x_2}{b}\right) \sin\left(\frac{\pi x_3}{b}\right) \sin\left(\frac{\pi x_3}{b$$

$$\psi = \frac{\sqrt{6}}{b^2} \left[\sin\left(\frac{\pi x_1}{b}\right) \sin\left(\frac{\pi x_2}{b}\right) \sin\left(\frac{\pi x_3}{b}\right) \sin\left(\frac{2\pi x_4}{b}\right) + \sin\left(\frac{\pi x_2}{b}\right) \sin\left(\frac{\pi x_3}{b}\right) \sin\left(\frac{2\pi x_4}{b}\right) \sin\left(\frac{\pi x_1}{b}\right) + \sin\left(\frac{\pi x_3}{b}\right) \sin\left(\frac{\pi x_2}{b}\right) \sin\left(\frac{\pi x_3}{b}\right) \sin\left(\frac{\pi x_3}{$$

(iii) For identical fermions,

$$\begin{split} \psi &= \frac{1}{b^2 \sqrt{6}} \left[\sin\left(\frac{k\pi x_1}{k}\right) \sin\left(\frac{l\pi x_2}{b}\right) \sin\left(\frac{m\pi x_3}{b}\right) \sin\left(\frac{n\pi x_4}{b}\right) + \sin\left(\frac{k\pi x_2}{b}\right) \sin\left(\frac{l\pi x_3}{b}\right) \sin\left(\frac{m\pi x_4}{b}\right) \sin\left(\frac{n\pi x_1}{b}\right) \right] \\ &+ \sin\left(\frac{k\pi x_3}{b}\right) \sin\left(\frac{l\pi x_4}{b}\right) \sin\left(\frac{m\pi x_1}{b}\right) \sin\left(\frac{n\pi x_2}{b}\right) + \sin\left(\frac{k\pi x_4}{b}\right) \sin\left(\frac{l\pi x_1}{b}\right) \sin\left(\frac{m\pi x_4}{b}\right) \sin\left(\frac{n\pi x_2}{b}\right) \\ &- \sin\left(\frac{k\pi x_2}{b}\right) \sin\left(\frac{l\pi x_4}{b}\right) \sin\left(\frac{m\pi x_2}{b}\right) \sin\left(\frac{n\pi x_4}{b}\right) - \sin\left(\frac{k\pi x_4}{b}\right) \sin\left(\frac{l\pi x_2}{b}\right) \sin\left(\frac{m\pi x_4}{b}\right) \sin\left(\frac{n\pi x_2}{b}\right) \\ &- \sin\left(\frac{k\pi x_3}{b}\right) \sin\left(\frac{l\pi x_4}{b}\right) \sin\left(\frac{m\pi x_2}{b}\right) \sin\left(\frac{n\pi x_1}{b}\right) - \sin\left(\frac{k\pi x_4}{b}\right) \sin\left(\frac{l\pi x_2}{b}\right) \sin\left(\frac{m\pi x_4}{b}\right) \sin\left(\frac{n\pi x_2}{b}\right) \\ &+ \sin\left(\frac{k\pi x_1}{b}\right) \sin\left(\frac{l\pi x_3}{b}\right) \sin\left(\frac{m\pi x_2}{b}\right) \sin\left(\frac{n\pi x_4}{b}\right) + \sin\left(\frac{k\pi x_4}{b}\right) \sin\left(\frac{l\pi x_2}{b}\right) \sin\left(\frac{m\pi x_4}{b}\right) \sin\left(\frac{n\pi x_4}{b}\right) \\ &+ \sin\left(\frac{k\pi x_2}{b}\right) \sin\left(\frac{l\pi x_3}{b}\right) \sin\left(\frac{m\pi x_4}{b}\right) \sin\left(\frac{n\pi x_3}{b}\right) + \sin\left(\frac{k\pi x_4}{b}\right) \sin\left(\frac{l\pi x_4}{b}\right) \sin\left(\frac{m\pi x_2}{b}\right) \\ &- \sin\left(\frac{k\pi x_4}{b}\right) \sin\left(\frac{l\pi x_3}{b}\right) \sin\left(\frac{m\pi x_4}{b}\right) \sin\left(\frac{n\pi x_2}{b}\right) - \sin\left(\frac{k\pi x_4}{b}\right) \sin\left(\frac{l\pi x_4}{b}\right) \sin\left(\frac{m\pi x_2}{b}\right) \\ &- \sin\left(\frac{k\pi x_4}{b}\right) \sin\left(\frac{l\pi x_4}{b}\right) \sin\left(\frac{m\pi x_4}{b}\right) \sin\left(\frac{n\pi x_2}{b}\right) - \sin\left(\frac{k\pi x_4}{b}\right) \sin\left(\frac{l\pi x_4}{b}\right) \sin\left(\frac{m\pi x_4}{b}\right) \sin\left(\frac{n\pi x_4}{b}\right) \\ &+ \sin\left(\frac{k\pi x_4}{b}\right) \sin\left(\frac{l\pi x_4}{b}\right) \sin\left(\frac{m\pi x_4}{b}\right) \sin\left(\frac{n\pi x_2}{b}\right) + \sin\left(\frac{k\pi x_4}{b}\right) \sin\left(\frac{l\pi x_4}{b}\right) \sin\left(\frac{m\pi x_4}{b}\right) \sin\left(\frac{n\pi x_4}{b}\right) \\ &+ \sin\left(\frac{k\pi x_4}{b}\right) \sin\left(\frac{l\pi x_4}{b}\right) \sin\left(\frac{m\pi x_4}{b}\right) \sin\left(\frac{n\pi x_4}{b}\right) + \sin\left(\frac{k\pi x_4}{b}\right) \sin\left(\frac{l\pi x_4}{b}\right) \sin\left(\frac{m\pi x_4}{b}\right) \sin\left(\frac{n\pi x_4}{b}\right) \\ &+ \sin\left(\frac{k\pi x_4}{b}\right) \sin\left(\frac{l\pi x_4}{b}\right) \sin\left(\frac{m\pi x_4}{b}\right) \sin\left(\frac{n\pi x_4}{b}\right) \sin\left(\frac{l\pi x_4}{b}\right) \sin\left(\frac{n\pi x_4}{b}\right) \sin\left(\frac{n\pi x_4}{b}\right) \\ &+ \sin\left(\frac{k\pi x_4}{b}\right) \sin\left(\frac{l\pi x_4}{b}\right) \sin\left(\frac{m\pi x_4}{b}\right) \sin\left(\frac{l\pi x_4}{b}$$

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Ground state,
$$E = \frac{1}{mb^2} \frac{\pi}{mb} = \frac{1}{b^2 \sqrt{6}} \left[\sin\left(\frac{\pi h}{b}\right) \sin\left(\frac{2\pi x_2}{b}\right) \sin\left(\frac{3\pi x_3}{b}\right) \sin\left(\frac{4\pi x_4}{b}\right) + \sin\left(\frac{\pi x_2}{b}\right) \sin\left(\frac{2\pi x_3}{b}\right) \sin\left(\frac{3\pi x_2}{b}\right) \sin\left(\frac{4\pi x_3}{b}\right) + \sin\left(\frac{\pi x_2}{b}\right) \sin\left(\frac{2\pi x_3}{b}\right) \sin\left(\frac{3\pi x_2}{b}\right) \sin\left(\frac{4\pi x_3}{b}\right) + \sin\left(\frac{\pi x_2}{b}\right) \sin\left(\frac{2\pi x_3}{b}\right) \sin\left(\frac{3\pi x_2}{b}\right) \sin\left(\frac{4\pi x_2}{b}\right) + \sin\left(\frac{\pi x_2}{b}\right) \sin\left(\frac{2\pi x_3}{b}\right) \sin\left(\frac{4\pi x_3}{b}\right) + \sin\left(\frac{\pi x_2}{b}\right) \sin\left(\frac{2\pi x_3}{b}\right) \sin\left(\frac{2\pi x_3}{b}\right) \sin\left(\frac{2\pi x_2}{b}\right) \sin$$

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