## Tutorial - 3: Quantum Mechanics

## To be discussed on 19 November 2018

1. Consider a particle in a one-dimensional box of length a, defined by the potential,

$$V(x) = 0$$
, for  $(0 \le x \le a)$   
=  $\infty$ , for  $(0 > x, x > a)$ 

The energy eigenvalues of the particle are given by  $\psi_n(x) = \sqrt{\frac{2}{a}} \sin\left(\frac{n\pi}{a}x\right)$ , where  $n = 1, 2, 3, \cdots$ 

- (a) At time t=0, the state of the particle is  $\psi(x,t=0)=\sqrt{\frac{1}{3}}\ \psi_1(x)+\sqrt{\frac{2}{3}}\ \psi_2(x)$  at time t=0.
  - i. What is the average energy,  $\langle E \rangle$  of the particle at t=0 ?
  - ii. In a measurement of energy, what is the probability to get the value  $E_1 = \frac{\pi^2 \hbar^2}{2ma^2}$ ?
  - iii. What is its state  $\psi(x,t)$  at a later time t?
  - iv. What is the value of  $\langle E \rangle$  of the particle at time t?
- (b) Consider a particle with energy  $E_1$  in the box.
  - i. What is the expectation value of the position of the particle?
  - ii. What is the probability to find it in the region  $0 \le x \le \frac{a}{2}$ ?
- (c) Repeat the above for the particle of energy  $E_2$ .

## Solution:

(a) i. We have  $H\psi_n = E_n\psi_n$  with  $E_n = \frac{n^2\phi^2\hbar^2}{2ma^2}$ . The average value of energy at t=0,

$$\langle E \rangle = \int_0^a \psi(x, t = 0)^* H \psi(x, t = 0) \, dx$$

$$= \int_0^a \left( \sqrt{\frac{1}{3}} \, \psi_1(x) + \sqrt{\frac{2}{3}} \, \psi_2(x) \right) \, H \left( \sqrt{\frac{1}{3}} \, \psi_1(x) + \sqrt{\frac{2}{3}} \, \psi_2(x) \right) \, dx$$

$$= \int_0^a \left( \sqrt{\frac{1}{3}} \, \psi_1(x) + \sqrt{\frac{2}{3}} \, \psi_2(x) \right) \, \left( \sqrt{\frac{1}{3}} \, E_1 \, \psi_1(x) + \sqrt{\frac{2}{3}} \, E_2 \, \psi_2(x) \right) \, dx$$

$$= \frac{E_1}{3} + \frac{2E_2}{3}$$

ii. The probability to yield  $E_1$  in a measurement is given by

$$\mathcal{P}_{E1} = \left| \int_0^a \psi_1^* \ \psi(x, t = 0) \ dx \right|^2 = \frac{1}{3}$$
 (1)

iii. Time evolution of the energy eigenstates are  $\psi_n(x,t) = e^{\frac{-iE_nt}{\hbar}} \psi_n(x)$ . That gives

$$\psi(x,t) = \sqrt{\frac{1}{3}} \psi_1(x) e^{\frac{-iE_1t}{\hbar}} + \sqrt{\frac{2}{3}} \psi_2(x) e^{\frac{-iE_2t}{\hbar}}$$
 (2)

iv.  $\langle E \rangle_t = \langle E \rangle_{t=0}$ .

(b) i.

$$\langle x \rangle = \frac{2}{a} \int_0^a x \sin^2\left(\frac{\pi x}{a}\right) dx = \frac{a}{2}$$

ii.

$$\mathcal{P}_{0 \le x \le \frac{a}{2}} = \frac{2}{a} \int_0^a \sin^2\left(\frac{\pi x}{a}\right) dx = \frac{1}{2}$$

(c) i.

$$\langle x \rangle = \frac{2}{a} \int_0^a x \sin^2 \left( \frac{2\pi x}{a} \right) dx = \frac{a}{2}$$

ii.

$$\mathcal{P}_{0 \le x \le \frac{a}{2}} = \frac{2}{a} \int_0^a \sin^2\left(2\frac{\pi x}{a}\right) dx = \frac{1}{2}$$

2. The wave function of a particle in a one-dimensional box (described in Question 1) has the wave function  $\psi(x) = A x^2(a-x)$ . What is the probability to find the particle in energy state  $E_1$  in a measurement? Solution:

$$\psi(x) = \sum_{n=0}^{\infty} c_n \sin\left(\frac{n\pi}{a}x\right)$$

$$c_n = A \int_0^a x^2 (a-x) \sin\left(\frac{n\pi}{a}x\right)$$

$$= \begin{cases} -\frac{2Aa^4}{n^3\pi^3}, & \text{for } n=1, 3, 5, \cdots \\ \frac{6Aa^4}{n^3\pi^3}, & \text{for } n=2, 4, 6, \cdots \end{cases}$$

Probability to find the particle in state of energy  $E_1$ ,  $\mathcal{P}_{E_1} = |c_1|^2 = \left(\frac{2Aa^4}{\pi^3}\right)^2$ .

## 3. Consider the one-dimensional potential,

$$V(x) = \infty,$$
 for  $(x < 0)$   
=  $-V_0$ , for  $(0 \le x \le a)$   
=  $0$ , for  $(x > a)$ .

Find the wave function of a particle of energy  $-V_0 < E < 0$  in all the three regions. Solution:

In region 1: x < 0,  $\psi_1(x) = 0$ .

In region 2:  $0 \le x \le a$ , we have the Schrödinger equation,

$$-\frac{\hbar^2}{2m}\frac{\partial^2}{\partial x^2}\psi - V_0\psi = E\psi.$$

$$\frac{\partial^2}{\partial x^2}\psi = -\frac{2m(E+V_0)}{\hbar^2}\psi$$

Note that  $E + V_0 > 0$ . This give the solution:  $\psi_2(x) = A\cos(kx) + B\sin(kx)$ , with  $k = \frac{\sqrt{2m(E+V_0)}}{\hbar}$ 

In region 3: x > a. We have V(x) = 0. That leads to  $\frac{\partial^2}{\partial x^2} \psi = -\frac{2mE}{\hbar^2} \psi = +\kappa^2 \psi$ , where  $\kappa = \frac{\sqrt{-2mE}}{\hbar} > 0$ .  $\implies \psi_3(x) = C \ e^{-\kappa x} + D \ e^{\kappa x}$ , but D = 0, lest  $\psi \to \infty$  as  $x \to \infty$ .

That is,

$$\begin{array}{lll} \psi_1(x) & = & 0, & \text{for } (x < 0) \\ \psi_2(x) & = & A\cos(kx) + B\sin(kx), & \text{for } (0 \le x \le a) \\ \psi_3(x) & = & C \; e^{-\kappa x}, & \text{for } (x > a). \end{array}$$

Boundary conditions:  $\psi_1(0) = \psi_2(0)$ ,  $\psi_2(a) = \psi_3(a)$ ,  $\frac{\partial \psi_2}{\partial x}\Big|_{x=a} = \frac{\partial \psi_3}{\partial x}\Big|_{x=a}$ 

$$A = 0$$

$$B\sin(ka) = Ce^{-\kappa a}$$

$$k B \cos(ka) = -\kappa Ce^{-\kappa a}$$

$$\Rightarrow \tan(ka) = -\frac{k}{\kappa}$$
(3)

In k and  $\kappa$ , E is the parameter. Energy cannot be any value, but those that satisfy the above equation. Redefine the parameter:

z = ka.

Notice that  $k^2 + \kappa^2 = \frac{2mV_0}{\hbar^2}$ .  $\Longrightarrow \kappa a = \sqrt{z_0^2 - z^2}$ .

Eqn.3 is now:  $\tan(z) = -\frac{z}{\sqrt{z_0^2 - z^2}}$ . Plot below shows  $\tan z$  (blue) and  $-\frac{z}{\sqrt{z_0^2 - z^2}}$  (orange) with z along the horizontal axis. The intersection points are the allowed values of z, and therefore (the corresponding) E.

