

**MEASUREMENT OF MULTIJET CROSS-SECTION RATIOS  
IN PROTON-PROTON COLLISIONS WITH THE  
CMS DETECTOR AT THE LHC**

A THESIS

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*Dedicated to  
my Parents*









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# Chapter 1

## Measurement of the Differential Inclusive Multijet Cross Sections and their Ratio

The inclusive  $n$ -jet event samples include the events with number of jets  $\geq n$ , where  $n = 2$  and 3 in the current study. The inclusive differential jet event cross sections are studied as a function of the average transverse momentum,  $H_{T,2}/2 = \frac{1}{2}(p_{T,1} + p_{T,2})$ , where  $p_{T,1}$  and  $p_{T,2}$  denote the transverse momenta of the two leading jets, and are defined by :

$$\frac{d\sigma}{d(H_{T,2}/2)} = \frac{1}{\epsilon \mathcal{L}_{\text{int,eff}}} \frac{N_{\text{event}}}{\Delta(H_{T,2}/2)} \quad (1.1)$$

where  $N_{\text{event}}$  is the number of inclusive 2- or 3-jet events counted in an  $H_{T,2}/2$  bin,  $\epsilon$  is the product of the trigger and jet selection efficiencies, which are greater than 99%,  $\mathcal{L}_{\text{int,eff}}$  is the effective integrated luminosity, and  $\Delta(H_{T,2}/2)$  are the bin widths. The measurements are reported in units of (pb/GeV).

The cross section ratio  $R_{32}$ , defined in Eq. 1.2 is obtained by dividing the differential cross sections of inclusive 3-jet events to that of inclusive 2-jet one, for

each bin in  $H_{T,2}/2$ .

$$R_{32} = \frac{\frac{d\sigma_{3-jet}}{d(H_{T,2}/2)}}{\frac{d\sigma_{2-jet}}{d(H_{T,2}/2)}} \quad (1.2)$$

For inclusive 2-jet events ( $n_j \geq 2$ ) sufficient data are available up to  $H_{T,2}/2 = 2000$  GeV, while for inclusive 3-jet events ( $n_j \geq 3$ ) and the ratio  $R_{32}$ , the accessible range in  $H_{T,2}/2$  is limited to  $H_{T,2}/2 < 1680$  GeV.

## 1.1 Data Samples

This measurement uses the data which was collected at the center of mass energy of 8 TeV by CMS experiment in the 2012 run period of the LHC. The 2012 data is taken in four periods A, B, C, D and the data sets are divided into samples according to the run period. Further each sample is grouped into subsets based on the trigger decision. For run B-D, the **JetMon** stream datasets contain prescaled low trigger threshold paths (HLTPFJet40, 80, 140, 200 and 260) while the **JetHT** stream datasets contain unprescaled high threshold trigger paths (HLT PFJet320 and 400). For run A, the **Jet** stream contains all the above mentioned trigger paths. The data to be used in physics analysis must satisfy a certain criteria according to which it should fulfill the validation requirements of data quality monitoring procedure. CMS uses JSON (Java Script Object Notation) format files to store the range of good lumi sections within a run. In the current analysis, the applied certification file<sup>1</sup> is based on the final event reconstruction of the 2012 CMS data sets. The datasets used in the current study are mentioned in the Table 1.1 along with the luminosity of each dataset.

The data sets have the LHC luminosity increasing with period, full data sample of 2012 corresponds to an integrated luminosity of  $19.71 \text{ fb}^{-1}$ .

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<sup>1</sup>Cert\_190456-208686\_8TeV\_22Jan2013ReReco\_Collisions12.JSON

Table 1.1: Run range and luminosity of the proton-proton collisions data collected at the center of mass energy of 8 TeV by CMS experiment in the year of 2012 in four different run periods A, B, C and D.

Run	Run range	Data set	Luminosity $\text{fb}^{-1}$
A	190456-193621	/Jet/Run2012A-22Jan2013-v1/AOD	0.88
B	193834-196531	/Jet[Mon,HT]/Run2012B-22Jan2013-v1/AOD	4.41
C	198022-203742	/Jet[Mon,HT]/Run2012C-22Jan2013-v1/AOD	7.06
D	203777-208686	/Jet[Mon,HT]/Run2012D-22Jan2013-v1/AOD	7.37

### 1.1.1 Monte Carlo Samples

To have a comparison of data results with the simulated events, the MADGRAPH5 [1] Monte-Carlo (MC) event generator has been used. The MADGRAPH5 generates matrix elements for High Energy Physics processes, such as decays and  $2 \rightarrow n$  scatterings. It has been interfaced to PYTHIA6 [2] by the LHE event record [3], to generate the rest of the higher-order effects using the Parton Showering (PS) model, with tune Z2\* to model the underlying event. Matching algorithms make sure that there is no double-counting between the tree-level and the PS-model-generated partons. The MC samples are processed through the complete CMS detector simulation to allow studies of the detector response and compare to measured data on detector level.

The cross section measured as a function of the transverse momentum  $p_T$  or the scalar sum of the transverse momentum of all jets  $H_T$  falls steeply with the increasing  $p_T$ . So in the reasonable time, it is not possible to generate a large number of high  $p_T$  events. Hence, the events are generated in the different phase-space region binned in  $H_T$  or the leading jet  $p_T$ . Later on, the different phase-space regions are added together in the data analyses by taking into account the cross section of the different phase-space regions. The official CMS MADGRAPH5+PYTHIA6 (MG5+P6) MC samples used in this analysis are generated as slices in the  $H_T$  phase-space are tabulated in Table 1.2 along with their cross sections and number of events generated.

Table 1.2: The official Monte Carlo samples are produced in phase space slices in  $H_T$  with the generator MADGRAPH5 and interfaced to PYTHIA6 for the parton shower and motorization of the events. The cross section and number of events generated are mentioned for each sample.

Generator	Sample	Events	Cross Section pb
MADGRAPH5 + PYTHIA 6	/QCD_HT-100To250_TuneZ2star_8TeV-madgraph-pythia6/ Summer12_DR53X-PU_S10_START53_V7A-v1/AODSIM	50129518	$1.036 \times 10^7$
	/QCD_HT-250To500_TuneZ2star_8TeV-madgraph-pythia6/ Summer12_DR53X-PU_S10_START53_V7A-v1/AODSIM	27062078	$2.760 \times 10^5$
	/QCD_HT-500To1000_TuneZ2star_8TeV-madgraph-pythia6/ Summer12_DR53X-PU_S10_START53_V7A-v1/AODSIM	30599292	$8.426 \times 10^3$
	/QCD_HT-1000ToInf_TuneZ2star_8TeV-madgraph-pythia6/ Summer12_DR53X-PU_S10_START53_V7A-v1/AODSIM	13843863	$2.040 \times 10^2$

## 1.2 Event Selection

The events are selected according to several quality criteria which ensures the high purity and high selection efficiency of the sample to be studied. This event selection also reduces beam induced background, detector-level noise and jets arising from fake calorimeter energy deposits.

### 1.2.1 Trigger Selection

CMS implements a two-level trigger system to reduce the amount of recorded events to a sustainable rate. In this analysis the jets are the final objects to study. So single jet trigger paths with varying thresholds are used to select events in data. It consists of one L1 trigger seed and multiple HLT filters. The L1 jet trigger uses transverse energy sums computed using both HCAL and ECAL in the central region ( $|\eta| < 3.0$ ) or HF in the forward region ( $|\eta| > 3.0$ ). The single jet triggers (HLT\_PFJetX) used for the current study are tabulated in Table 1.3. A single jet trigger selects an event in which at least one jet has the transverse momentum above the threshold. HLT\_PFJetX implies that there is at-least one jet in the event, whose  $p_T > X$  (GeV). The L1 trigger has a lower threshold to ensure full efficiency versus  $p_T$  of the HLT trigger. The  $p_T$  spectrum is steeply falling and hence the rates for low- $p_T$  jets are very high. So it is not feasible to use a single unprescaled trigger for the selection of



all required events. To collect sufficient data in the lower part of the  $p_T$  spectrum, different five prescaled low- $p_T$  trigger paths, each with different prescale value, are used. Also, one unprescaled trigger i.e. HLT\_Jet320 is used in the high  $p_T$  region, in which the rate is sufficiently small to collect and store all events. During the reconstruction of the spectrum, the prescales have been taken into the account.

Table 1.3: The single jet trigger paths used in the analysis are listed here. The column  $H_{T,2}/2$ , 99% indicates the value of  $H_{T,2}/2$  at which each trigger exhibits an efficiency larger than 99%. The last column gives the effective luminosity seen by each trigger which divided by the total integrated luminosity of  $19.71 \text{ fb}^{-1}$ , gives the effective prescale applied on a trigger over the whole run period.

Trigger Path	L1 threshold GeV	HLT threshold GeV	$H_{T,2}/2$ , 99% GeV	Eff. Lumi $\text{fb}^{-1}$
HLT_PFJet80	36	80	120.0	$0.21 \times 10^{-2}$
HLT_PFJet140	68	140	187.5	$0.56 \times 10^{-1}$
HLT_PFJet200	92	200	262.5	0.26
HLT_PFJet260	128	260	345.0	1.06
HLT_PFJet320	128	320	405.0	19.71

The efficiency of each trigger path as a function of  $H_{T,2}/2$  is described by the turn-on curves with a rising part where the trigger is partly inefficient, until a plateau region where the trigger is fully efficient. Hence it is important to determine the threshold above which a trigger becomes fully efficient. The threshold is the value at which the trigger efficiency exceeds 99%. The trigger efficiency for HLT\_PFJetY is given by Eq. 1.3 where HLT\_PFJetX is the reference trigger and is assumed to be fully efficient in the considered phase space region. The value of X is chosen previous to that of Y in  $p_T$  ordering from the trigger list so that the higher trigger condition can be emulated from the lower trigger path.

$$\epsilon_{\text{HLT\_PFJetY}} = \frac{H_{T,2}/2 \left( \text{HLT\_PFJetX} + (\text{L1Object\_p}_T > Z) + (\text{HLTObject\_p}_T > Y) \right)}{H_{T,2}/2(\text{HLT\_PFJetX})} \quad (1.3)$$

where Y is the  $p_T$  threshold of HLT\_PFJetY and Z is the L1 seed value corresponding

to the trigger path HLT\_PFJetY. The denominator represents the number of events for which the reference trigger path HLT\_PFJetX has been fired. The numerator is the number of events for which HLT\_PFJetX has been fired along the  $p_T$  of L1Object  $\geq Z$  and the  $p_T$  of HLTOBJECT  $\geq Y$ . For example, to obtain turn-on curve for HLT\_PFJet260, HLT\_PFJet200 is the reference HLT path. The  $p_T$  cut on L1Object is 128 GeV and  $p_T$  cut on HLTOBJECT is 260 GeV. The threshold point at which the trigger efficiency is larger than 99% is determined by fitting the turn-on distribution with a sigmoid function described in Eq. 1.4. The trigger turn-on curves as a function of  $H_{T,2}/2$  can be seen in Fig. 1.1 which are described by a sigmoid function (blue line). The error bars give the uncertainty on the efficiency which is calculated using Clopper-Pearson confidence intervals [4].

$$f_{fit}(x) = \frac{1}{2} \left( 1 + \operatorname{erf} \left( \frac{x - \mu}{\sqrt{2}\sigma} \right) \right) \quad (1.4)$$

### 1.2.2 Primary Vertex Selection

The reconstructed tracks, selected based on the transverse impact parameter significance with respect to the beam line, number of strip and pixel hits and the normalized track  $\chi^2$ , identify the primary vertex (PV). The tracks are clustered according to the z-coordinate of their point of closest approach to the beam axis. A selection criteria for primary vertex should be followed which helps to identify and reject the beam background events. At-least one good primary vertex reconstructed from at least four tracks within a distance of  $|z(PV)| < 24$  cm to the nominal interaction point in a collision, is required in each event. The radial distance in x-y plane,  $\rho(PV)$  should be not be greater than 2 cm. The number of degrees of freedom in fitting for the position of each vertex using its associated tracks should be at-least four in number.

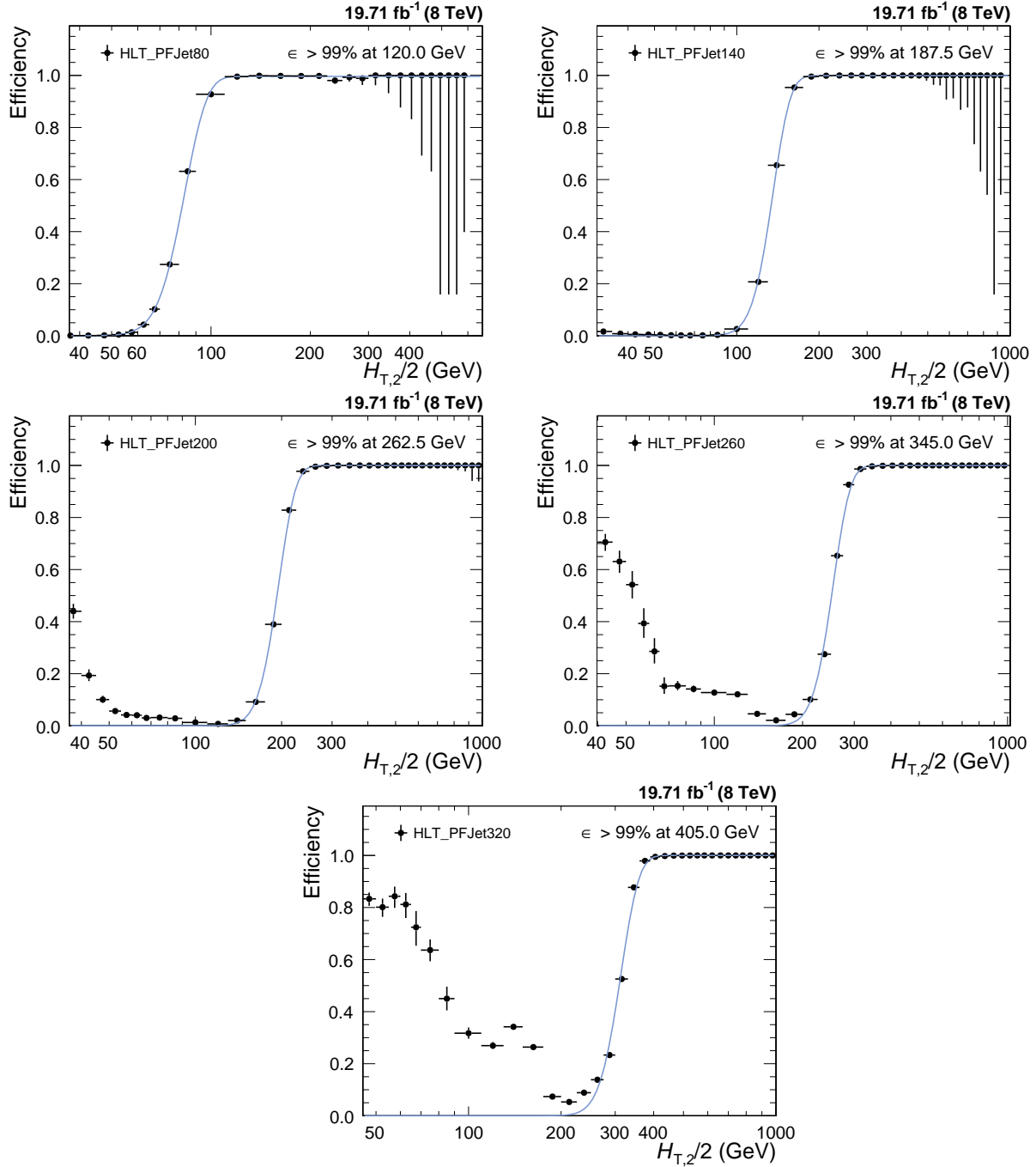


Figure 1.1: Trigger efficiencies turn-on curves for the single jet trigger paths are fitted with a sigmoid function (blue line) to obtain the 99% efficiency threshold. The error bars give the uncertainty on the efficiency which is calculated using Clopper-Pearson confidence intervals [4].

### 1.2.3 Missing Transverse Energy

In an ideal detector where all particles could be identified and perfectly measured, the transverse momentum of all particles would sum up to zero. But the neutral weakly interacting particles, such as neutrinos, escape from typical collider detectors and do not produce any direct response in the detector elements. The imbalance of total momentum of all visible particles can give the hints of the presence of such particles. The vector momentum imbalance in the plane perpendicular to the beam direction is known as missing transverse momentum or energy ( $E_T^{\text{miss}}$ ). It is one of the most important observables for discriminating leptonic decays of W bosons and top quarks from background events which do not contain neutrinos, such as multijet and Drell–Yan events or searches for physics beyond the Standard Model.

The ratio of missing transverse energy to the total transverse energy  $E_T^{\text{miss}}/\sum E_T$ , shown in Fig. 1.2 for  $n_j \geq 2$  (left) and  $n_j \geq 3$  events (right), shows a discrepancy between data (black solid circles) and simulated MC (blue histogram), at the tail part of the distribution. This is because of a finite contribution from  $Z(\rightarrow \nu\bar{\nu}) + \text{jet}$  events which gives rise to non-zero  $E_T$  in the events in data. Such events are absent in QCD simulated events in MC. Hence  $E_T^{\text{miss}}/\sum E_T$  is required to be less than 0.3 to reject events with high  $E_T^{\text{miss}}$ .

### 1.2.4 Jet Identification

In order to suppress fake jets, arising from detector noise or misreconstructed particles, jet identification criteria (ID) has been applied. Instead of applying it event-wise, it is applied it on each jet. The algorithm works on reconstructed jets using information of the clustered particle candidates. The official tight jet ID [5], recommended by JETMET group [6] is used. Due to pileup and electronic noise the jet constituent fractions may vary from event to event. In order to reject the noisy jets, some jet selection criteria are optimized to select only good quality jets. The selec-

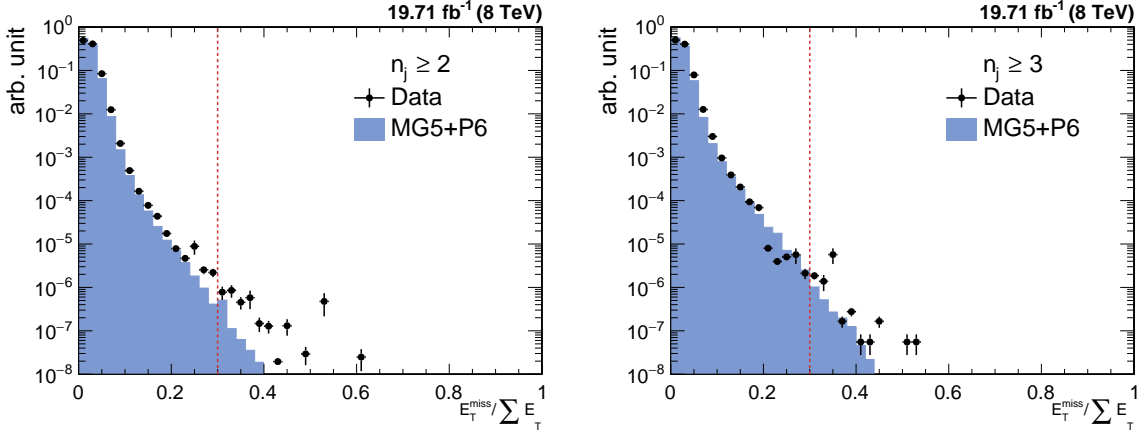


Figure 1.2: Missing transverse energy fraction of the total transverse energy per event in data (black solid circles) and simulated Monte Carlo events (blue histogram) in inclusive 2-jet (left) and 3-jet events (right). To remove background and noise, events with a fraction exceeding a certain threshold, here indicated with the red dashed line, are rejected.

tion criteria are implemented as selection cut on jet fractions. Table 1.4 summarizes the properties of the reconstructed jets and their respective cuts. Each jet should contain at least two particles, one of which should be a charged hadron. The cut on the fraction of neutral hadrons and photons removes HCAL noise and ECAL noise, respectively. Muons that are falsely identified and clustered as jets are removed by the muon fraction criterion. Based on information of the tracker, additional selection cuts are enforced in the region  $|\eta| < 2.4$ . The charged electromagnetic fraction cut removes the jets clustered from misidentified electrons. Furthermore, the fraction of charged hadrons in the jet must be larger than zero and jets without any charged hadrons are very likely to be pileup jets. The Figs. 1.3 and 1.4 show the distributions of the jet constituents observed in data (black solid circles) and simulated MC events (blue histogram) for  $n_j \geq 2$  and  $n_j \geq 3$ , respectively.

#### 1.2.4.1 Jet ID Efficiency

The efficiency of the jet ID as a function of  $H_{T,2}/2$  is studied using a tag-and-probe technique with dijet events. The two leading jets are required to be back-to-back in the azimuthal plane such that  $|\Delta\phi - \pi| < 0.3$ . One of the dijets is selected randomly

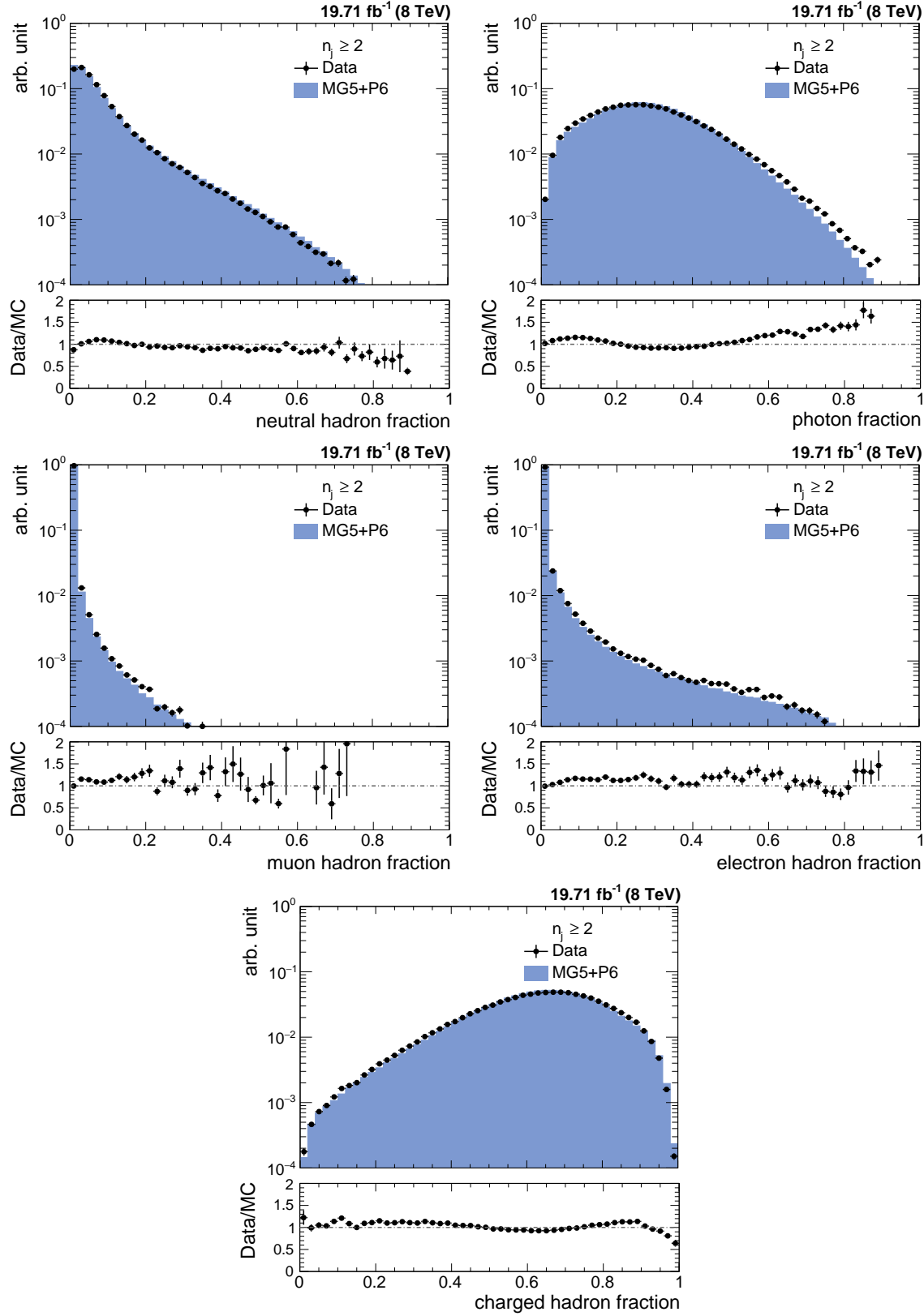


Figure 1.3: The fractions of jet constituents as observed in data (black solid circles) and simulated Monte Carlo events (blue histogram) for different types of PF candidates for inclusive 2-jet events. Data and simulations are normalized to the same number of events. The distributions are shown after the application of the jet ID.

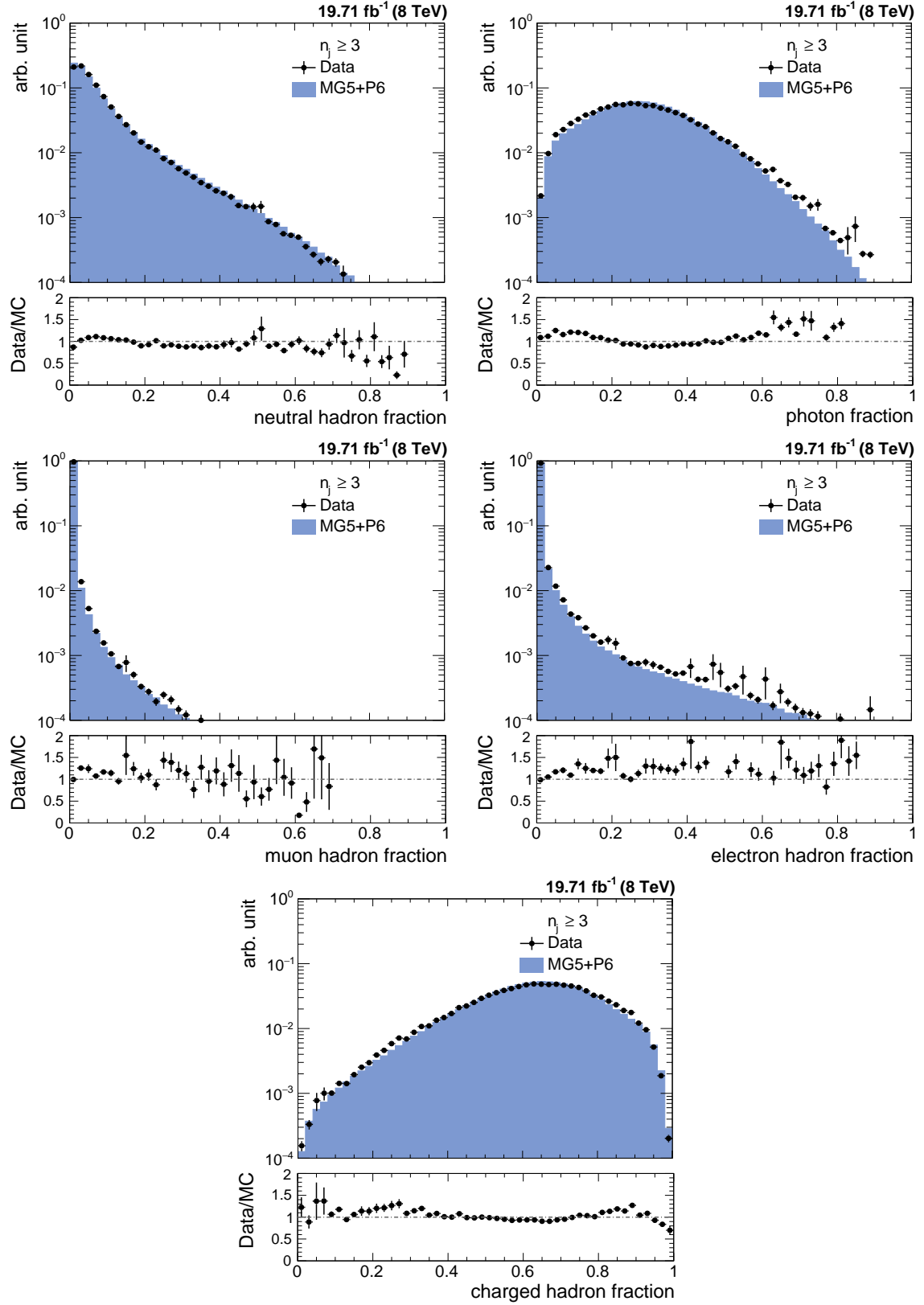


Figure 1.4: The fractions of jet constituents as observed in data (black solid circles) and simulated Monte Carlo events (blue histogram) for different types of PF candidates for inclusive 3-jet events. Data and simulations are normalized to the same number of events. The distributions are shown after the application of the jet ID.

Table 1.4: The jet ID removes noise and fake jets based on the properties of the reconstructed jets and the clustered particle candidates. All the selection cuts which are recommended by the JETMET group are applied [6].

	Property	Loose ID	Tight ID
Whole $\eta$ region	neutral hadron fraction	$< 0.99$	$< 0.90$
	neutral EM fraction	$< 0.99$	$< 0.90$
	number of constituents	$> 1$	$> 1$
	muon fraction	$< 0.80$	$< 0.80$
only $ \eta  < 2.4$	charged hadron fraction	$> 0$	$> 0$
	charged multiplicity	$> 0$	$> 0$
	charged EM fraction	$< 0.99$	$< 0.90$

as a “tag” jet which is required to fulfill the tight jet ID criteria. The other jet is called “probe” jet for which it is examined, whether it also passes the tight jet ID. The ID efficiency is defined as the ratio of events where the probe jet passes the ID requirements, over the total number of dijet events. It is shown as function of  $H_{T,2}/2$  in Fig. 1.5 and as expected, it is always greater than 99%. The QCD cross section decreases as a function of  $H_{T,2}/2$  and hence the number of events decrease on moving to higher  $H_{T,2}/2$ . Consequently the statistical fluctuations for ID efficiency are larger at higher  $H_{T,2}/2$ .

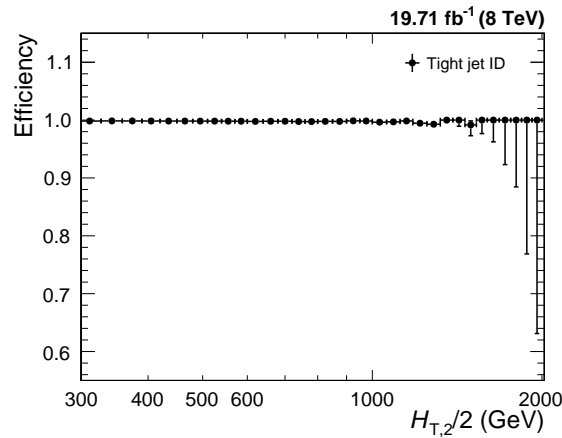


Figure 1.5: The jet ID efficiency is studied as a function of  $H_{T,2}/2$  with tag-and-probe technique using dijet event topologies and it always exceeds 99%.



### 1.2.5 Jet Selection

The measurement of differential cross sections and their ratio uses jets clustered from particle flow candidates using the anti- $k_t$  jet algorithm with a size parameter,  $R = 0.7$ . The energy scale of the jets is corrected with the CMS recommended jet energy corrections, described in Sec. [?](#). These corrections are applied to jets in both data<sup>[2](#)</sup> as well those in simulated events<sup>[3](#)</sup>. The jet selection, based on phase space cuts on transverse momentum and rapidity of jets in an event, is as follows :

- All jets having  $p_T > 150$  GeV and  $|y| < 5.0$  are selected.
- Events with at least two jets are selected.
- The two leading jets should have  $|y| < 2.5$  and further jets are counted only, if they lie within the same central rapidity range of  $|y| < 2.5$ .

These cuts assure high detector acceptance and exactly same selection is applied in the measurement, simulated events as well in theoretical calculations for a consistent comparison.

## 1.3 Comparison with Simulation

### 1.3.1 Pile-up Reweighting

While generating the official Monte-Carlo samples, the number of pileup interactions describing the conditions expected for each data-taking period are taken care of. But the number of pile-up events implemented in the simulation  $N_{MC}(N_{PU,truth})$ , does not match exactly with the one measured actually in data  $N_{data}(N_{PU,est.})$ . To match the pile-up distributions in data, a reweighting factor  $w_{PU}$ , as given by Eq. [1.5](#) is applied

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<sup>2</sup>Winter14\_V8 jet energy corrections

<sup>3</sup>START53\_V27 jet energy corrections

to the simulated events. In Fig. 1.6 the number of reconstructed vertices are shown before (left) and after pile-up reweighting (right). It is observed that before pile-up reweighting there was a significant mismatch of the pile-up distributions in data (black solid circles) and simulated MC events (blue histogram), which completely vanishes after reweighting.

$$w_{\text{PU}} = \frac{N_{\text{data}}(N_{\text{PU,est.}}) / \sum N_{\text{data}}}{N_{\text{MC}}(N_{\text{PU,truth}}) / \sum N_{\text{MC}}} \quad (1.5)$$

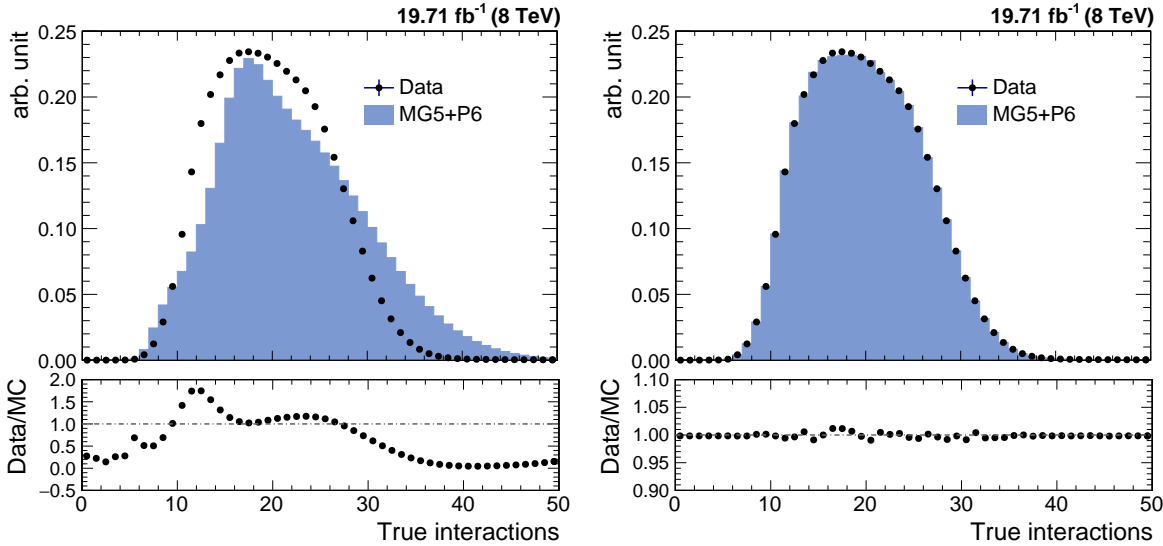


Figure 1.6: Number of reconstructed vertices in data (black solid circles) and simulated Monte Carlo events (blue histogram) before (left) and after (right) the pile-up reweighting.

### 1.3.2 Comparison of Cross Sections and their Ratio

The measured data distribution of differential cross section at detector level is compared to the predictions of Monte Carlo simulation using MADGRAPH5 generator interfaced with PYTHIA6 (MG5+P6) including the detector simulation as well as to a fixed-order theory prediction obtained using CT10-NLO PDF set. Figure 1.7 shows the comparison of differential cross section as a function of  $H_{\text{T},2}/2$  for  $n_j \geq 2$  (left) and  $n_j \geq 3$  events (right), for data (black solid circles), MG5+P6 MC (red empty

circles) and CT10-NLO (blue histogram). The bottom panel in each plot shows the ratio of data to the MC predictions (red line) as well as to the CT10-NLO theory predictions (blue line). The NLO predictions on parton level are not corrected for non-perturbative effects. Still the NLO predictions describe the data better as compared to the LO MC simulations which roughly describes the spectrum on detector level. The sufficient data for  $n_j \geq 2$  and  $n_j \geq 3$  events are available up to  $H_{T,2}/2 = 2000$  GeV and 1680 GeV, respectively. Due to some kinematical constraints, the minimum cut on  $H_{T,2}/2$  is 300 GeV (explained in Sec. 2.1.1). Hence the differential cross sections are studied in the range  $300 \text{ GeV} \leq H_{T,2}/2 < 2000$  GeV for  $n_j \geq 2$  and  $300 \text{ GeV} \leq H_{T,2}/2 < 1680$  GeV for  $n_j \geq 3$  events.

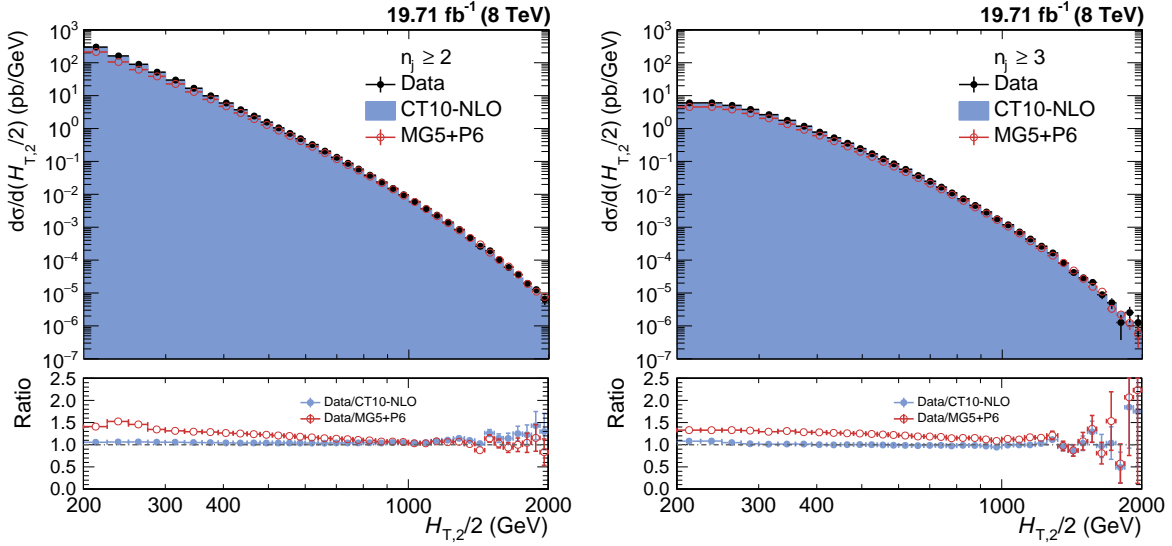


Figure 1.7: The differential cross sections are compared for data (black solid circles) and LO MADGRAPH5+PYTHIA6 (MG5+P6) Monte Carlo (red empty circles), at reconstructed level with CT10-NLO theory predictions (blue histogram), as a function of  $H_{T,2}/2$  for inclusive 2-jet (left) and 3-jet events (right). Ratios of data to the Monte Carlo predictions (red line) as well as to the CT10-NLO predictions (blue line) are shown in bottom panel of each plot.

The ratio of differential cross sections,  $R_{32}$  as a function of  $H_{T,2}/2$ , is extracted by dividing the cross section of selected inclusive 3-jet events to that of inclusive 2-jet events at any given bin size of  $H_{T,2}/2$ . In the cross section ratios, the numerator and denominator are not independent samples. So to calculate the

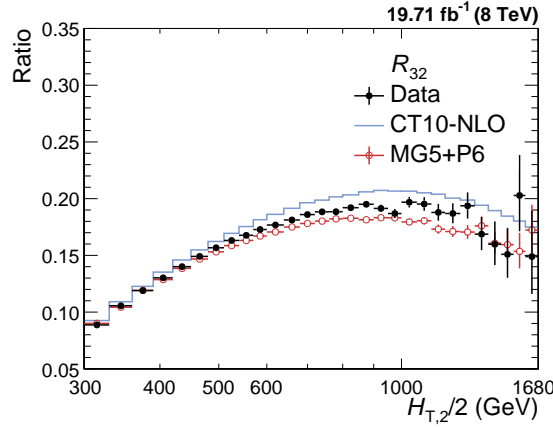


Figure 1.8: Comparison of the cross section ratio  $R_{32}$  as a function of  $H_{T,2}/2$ , for data (black solid circles) and LO MADGRAPH5+PYTHIA6 (MG5+P6) Monte Carlo (red empty circles), at reconstructed level with CT10-NLO theory predictions (blue line). The error bars give the asymmetrical statistical uncertainty, calculated by the Wilson score interval method which takes into the account the correlation between the numerator and denominator.

statistical uncertainty for the cross section ratios at reconstructed level, the Wilson score interval method is used which takes into account the correlation between the numerator and the denominator and give asymmetric errors. Figure 1.8 shows the comparison of the cross section ratio  $R_{32}$  as a function of  $H_{T,2}/2$ , for data (black solid circles) and LO MADGRAPH5+PYTHIA6 (MG5+P6) Monte Carlo (red empty circles), at reconstructed level with CT10-NLO theory predictions (blue line). Since in  $n_j \geq 3$  events, the enough statistics for differential cross section is available upto 1680 GeV of  $H_{T,2}/2$  only,  $R_{32}$  is also studied in the range  $300 \text{ GeV} \leq H_{T,2}/2 < 1680 \text{ GeV}$ . The bin-wise inclusive 2-jet and 3-jet events differential cross sections as well as their ratio  $R_{32}$ , calculated at detector level, along with statistical uncertainty (in %) are tabulated in Table A.1.

## 1.4 Jet Energy Resolution (JER)

In an ideal experiment, the value of a physical quantity would be determined exactly with an infinite precision. For e.g. whenever a particle with energy  $E$  passes an ideal calorimeter having infinite resolution, the measured energy should always be equal

to  $E$ . But in real world, the measured energy of the above mentioned particle might differ from the value  $E$ . This difference of the measured quantity from its true value may be due to detector noise, uncertainties in the calibration, non-linearity of the response etc. Hence this results in the finite value of the resolution of the detector known as jet energy resolution (JER). In such case, the measured values of energy of different particles, passing through the same detector with same energy  $E$ , will be different. Such measurements are described by a Gaussian distribution which is centered around the true value of the measured quantity and its width is generally interpreted as detector resolution. Hence the importance of the detector resolution lies in the fact that it indicates how much the measured value of the observable differs from the true one i.e. how precisely a physical observable can be measured. The narrower the distribution, the higher the resolution is and hence the more efficient is the detector.

Due to finite resolution of the CMS detector, the measured transverse momentum of jets gets smeared. Since the observable in this study i.e.  $H_{T,2}/2$  is the average sum of transverse momentum of leading and sub-leading jets, the resolution of the detector has to be studied in terms of the observable. CMS detector simulation based on MG5+P6 MC event generators is used to determine the resolution as both the particle and reconstructed level information is available. The jets clustered from stable generator particles called Gen jets as well as from particle flow candidates reconstructed from the simulated detector output called Reco jets, are used. The studies of the JETMET working group at CMS has shown that the jet energy resolution in data is actually worse than in simulation [7]. So the reconstructed jet transverse momentum needs to be smeared additionally to match the resolution in data. Table 1.5 shows the scaling factors ( $c$ ) which need to be applied on the transverse momentum of simulated reconstructed jets. The scaling factors depend on the absolute  $\eta$  of the jet. The uncertainty on these measured scaling factors ( $c_{central}$ ) needs to be taken into account in a physics analysis. This is done by smearing the reconstructed jets with two additional sets of scaling factors,  $c_{up}$

and  $c_{down}$ , that correspond to varying the factors up and down respectively, by one sigma and evaluating the impact of these new sets.

Table 1.5: JETMET working group at CMS has shown that the jet energy resolution in data is actually worse than in simulation [7]. To match the resolution in data, the reconstructed jet transverse momentum in simulated events need to be smeared by applying the scale factors. The uncertainty on the resolution is given by an upwards and downwards variation  $c_{up}$  and  $c_{down}$  of the measured scaling factor  $c_{central}$ .

$\eta$	0.0 – 0.5	0.5 – 1.1	1.1 – 1.7	1.7 – 2.3	2.3 – 2.8
$c_{central}$	1.079	1.099	1.121	1.208	1.254
$c_{down}$	1.053	1.071	1.092	1.162	1.192
$c_{up}$	1.105	1.127	1.150	1.254	1.316

The reconstructed jet  $p_T$  is smeared randomly using a Gaussian width widened by the scaling factor ( $c_{central}$ )

$$p_T \rightarrow Gauss\left(\mu = p_T, \sigma = \sqrt{c_{central}^2 - 1} \cdot JER(p_T)\right) \quad (1.6)$$

where  $JER(p_T)$  is the resolution determined as a function of jet  $p_T$  using MG5+P6 MC simulated events. After smearing transverse momentum of each reco jet,  $H_{T,2}/2$  is calculated from both generator particle jets (Gen  $H_{T,2}/2$ ) as well as the particle flow or reconstructed jets (Reco  $H_{T,2}/2$ ). Then the response is calculated as defined in the Eq. 1.7.

$$R = \frac{\text{Reco } H_{T,2}/2}{\text{Gen } H_{T,2}/2} \quad (1.7)$$

The width of the response distribution in a given Gen  $H_{T,2}/2$  bin is interpreted as the resolution which in good approximation can be described by the  $\sigma$  of a Gaussian fit of the response distribution. A double-sided Crystal-Ball function takes into account the non-Gaussian tails of the jet response distribution. The resolution as a function of  $H_{T,2}/2$  is calculated separately for both  $n_j \geq 2$  and  $n_j \geq 3$  events. A

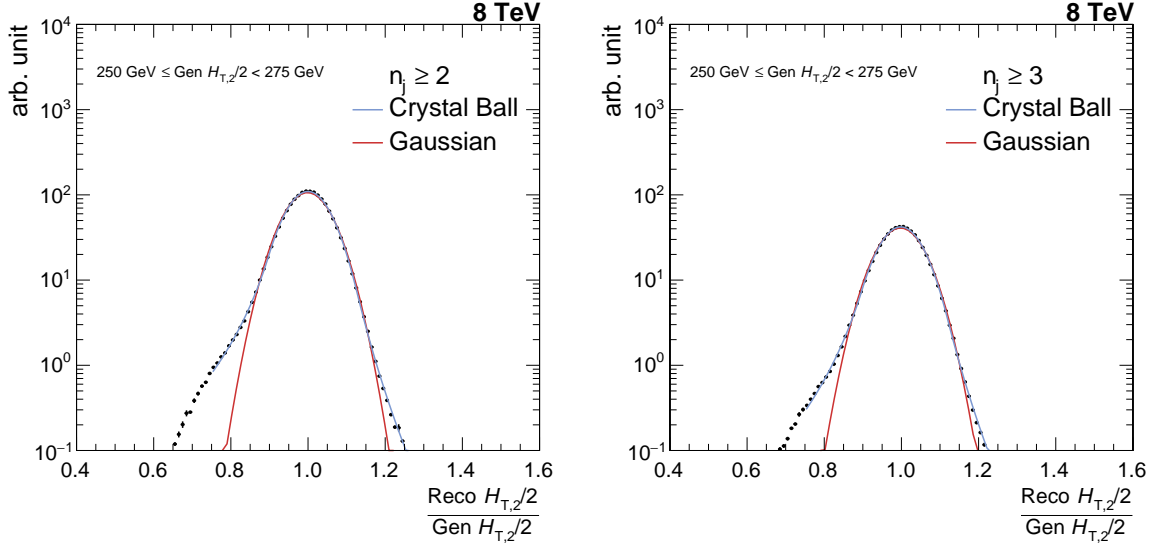


Figure 1.9: Fitting of the resolution distribution, obtained using LO MADGRAPH5+PYTHIA6 (MG5+P6) Monte Carlo simulated events, as a function of  $H_{T,2}/2$  for inclusive 2-jet (left) and inclusive 3-jet events (right). The blue line shows the double-sided Crystal Ball function fit of  $\frac{\text{Reco } H_{T,2}/2}{\text{Gen } H_{T,2}/2}$  in each  $\text{Gen } H_{T,2}/2$  bin, overlayed by Gaussian fitting the core of the resolution (red line).

fit example for one  $\text{Gen } H_{T,2}/2$  bin is shown in Fig. 1.9 for  $n_j \geq 2$  (left) and inclusive 3-jet events (right). Here the black dots represent the jet response distribution and the double-sided Crystal-Ball fit (blue line) is overlayed by the Gaussian fit (red line). The resolution in each  $\text{Gen } H_{T,2}/2$  bin is then plotted as a function of  $\text{Gen } H_{T,2}/2$ .

As expected, it has been observed from Fig. 1.10 that the Crystal Ball function (blue solid circles) describes the measured distributions better as compared to Gaussian function fit (red solid circles), especially in the low- $H_{T,2}/2$  region where the non-Gaussian tails are more pronounced. Hence JER is determined using Crystal Ball function fit. Figure 1.11 shows the final relative jet energy resolution (JER) which is described by a modified version of the NSC formula (blue solid line) [8], as mentioned in Equation 1.8. To consider the migration to lower as well higher bins and to obtain the resolution with reasonable statistics over the full range of  $\text{Gen } H_{T,2}/2$ , the fit function is extrapolated to 80 GeV and upto 2000 GeV which is

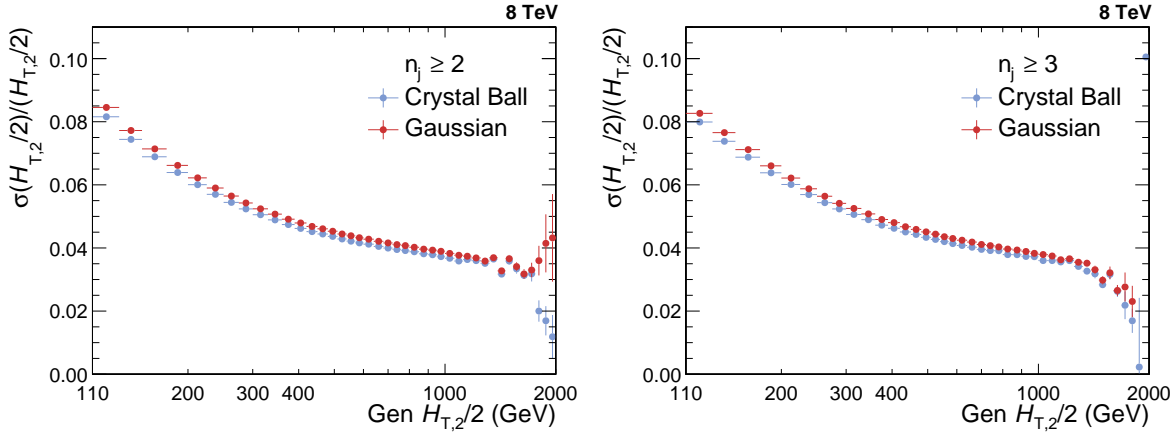


Figure 1.10: Comparison of jet energy resolution calculated using Crystal-Ball fit function (blue solid circles) and Gaussian fit function (red solid circles) for inclusive 2-jet (left) and inclusive 3-jet events (right).

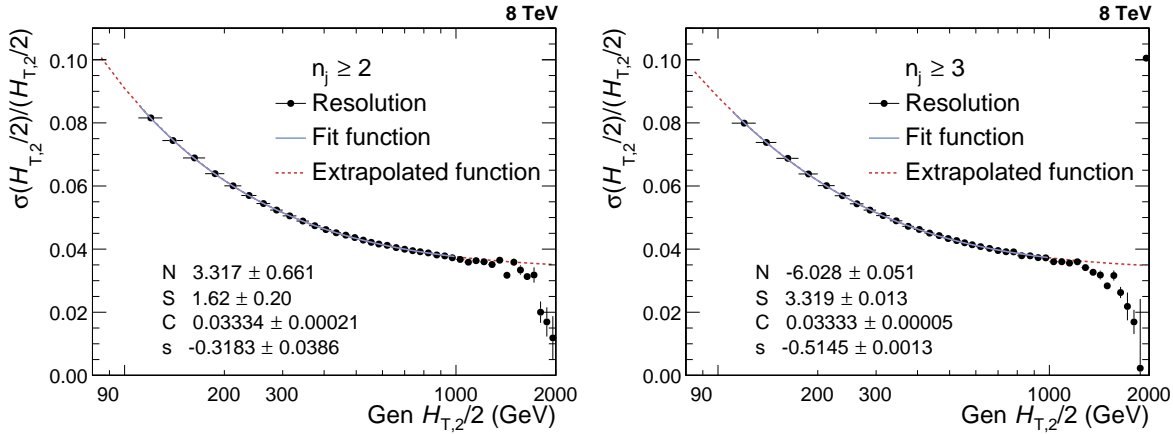


Figure 1.11: Jet energy resolution (JER) is shown as a function of  $\text{Gen } H_{T,2}/2$  for inclusive 2-jet (left) and inclusive 3-jet events (right). JER (black solid circles) is fitted by using the modified NSC-formula (blue solid line) which is extrapolated to 80 GeV and upto 2000 GeV (red dashed line) to consider the migration into lower as well as higher bins.

shown by red dashed line. The fit formula used here is basically the usual NSC formula which describes the resolution in terms of noise  $N$  originating due to electronic and pileup noise and is independent of  $H_{T,2}/2$ ; a stochastic component  $S$  due to sampling fluctuation and EM fraction fluctuation per hadrons; and a constant term  $C$  because of dead material, magnetic field and calorimeter cell to cell fluctuation. In the low  $H_{T,2}/2$  region the tracking has a non-negligible influence on the resolution due to the particle flow algorithm, so the additional parameter  $s$  is introduced to ob-



tain slightly better fits. The parameters obtained after fitting the relative resolution using the above mentioned NSC formula are tabulated in Table 1.6 for  $n_j \geq 2$  and  $n_j \geq 3$  events. This calculated JER is used in unfolding procedure to smear the generated truth spectrum which is used as input in getting the response matrices and is explained in details in Sec. 1.5.1. Since JER in  $n_j \geq 2$  events is similar to that one in  $n_j \geq 3$  events, so N, S and C fit parameters obtained for  $n_j \geq 3$  events are used for unfolding  $R_{32}$ .

$$\frac{\sigma(x)}{x} = \sqrt{\text{sgn}(N) \cdot \frac{N^2}{x^2} + S^2 \cdot x^{s-1} + C^2} \quad (1.8)$$

Table 1.6: The parameters obtained by fitting the relative resolution as a function of  $H_{T,2}/2$ , using the modified NSC formula, for inclusive 2-jet and inclusive 3-jet events.

	N	S	C	s
Inclusive 2-jet	3.32	1.62	0.0333	-0.318
Inclusive 3-jet	-6.03	3.32	0.0333	-0.515

Since the JER is calculated using MG5+P6 Reco and Gen  $H_{T,2}/2$  distributions, so it is expected that if Gen  $H_{T,2}/2$  is smeared using this JER, it should match the Reco  $H_{T,2}/2$ . But this extracted JER in one large rapidity bin, smears the Gen  $H_{T,2}/2$  too much because Smeared Gen/Gen ratio (red line) shows a discrepancy from simulated Reco/Gen ratio (blue line), as observed in Fig. 1.12 for  $n_j \geq 2$  (left) and  $n_j \geq 3$  events (right). Some shortcomings in the detector simulation of the theory spectra leads to these small nonclosures. When the 30% reduced JER is used to smear Gen, then the ratio Smeared Gen/Gen (pink line) matches with simulated Reco/Gen ratio (blue line) within the statistical fluctuations. Hence an additional unfolding uncertainty is attributed by comparison to 30% reduced JER for both  $n_j \geq 2$  and  $n_j \geq 3$  events. Due to high statistical fluctuations at high  $H_{T,2}/2$ , range is presented upto 1680 GeV only.

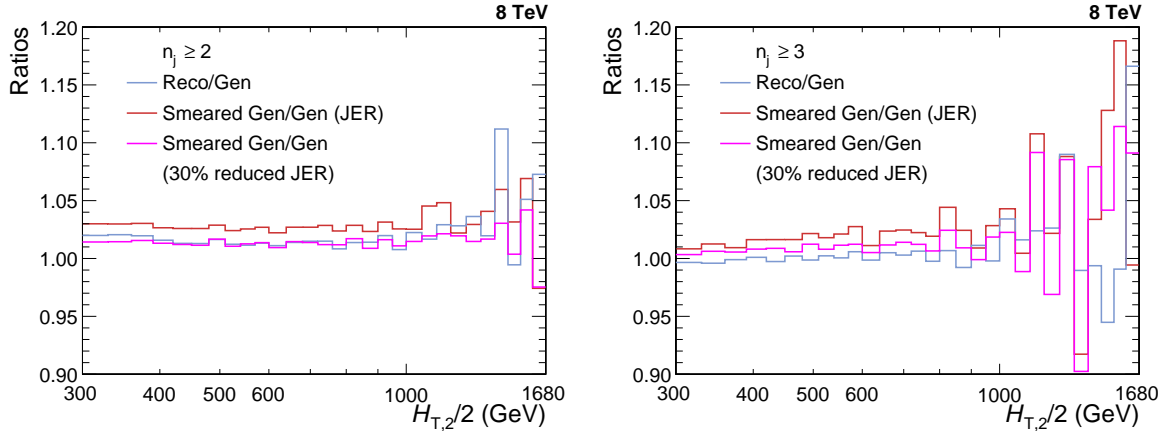


Figure 1.12: MADGRAPH5+PYTHIA6 (MG5+P6) Gen smeared using extracted jet energy resolution (JER) shows a discrepancy from simulated Reco as Smeared Gen/Gen ratio (red line) does not match with Reco/Gen ratio (blue line), for both inclusive 2-jet (left) and inclusive 3-jet events (right). Smeared Gen/Gen ratio (pink line) where Gen is smeared using 30% reduced JER matches with simulated Reco/Gen ratio (blue line) within the statistical fluctuations. Hence an additional unfolding uncertainty is attributed by comparison to 30% reduced JER.

## 1.5 Unfolding

One of the main goals in an experimental measurement is to do the comparison of data with theory predictions or with the results obtained from other experiments. But the finite resolution of a detector and the steeply falling jet  $p_T$  spectrum distorts the physical quantities. As a result, the measured observables are different from their corresponding true values. Each  $p_T$  bin content contains the migrated events from neighbouring bins along with the original events. So an unfolding process of the data should be followed in order to remove detector effects. In this analysis, the measurements are corrected for detector smearing effects and unfolded to stable particle level by using the iterative D'Agostini Bayesian algorithm [9, 10] as implemented in RooUnfold software package [11]. In this algorithm, the number of iterations regularize the unfolding process. The obtained distribution in one iteration is taken as the input in the next one.  $\chi^2$  between two successive iterations is given by Eq. 1.9. The number of iterations stop when  $\chi^2/N_{bins}$  is  $< 1$ . A reduced  $\chi^2$  is obtained by a higher number of iterations but this will also increase the uncertainty and there

are larger bin-by-bin fluctuations and correlations. So the optimization of number of iterations is very important. In the current analysis, unfolding done with “four” iterations gives the best results with low  $\chi^2$  and low bin-by-bin correlations.

$$\chi^2 = \sum_{i=1}^{N_{bins}} \left( \frac{n_i^{j+1} - n_i^j}{\sqrt{n_i^j}} \right)^2 \quad (1.9)$$

where  $n_i^j$  number of events in  $i$ -th bin for  $j$ -th iteration

The measured differential cross sections as a function of  $H_{T,2}/2$ , are unfolded separately for  $n_j \geq 2$  and  $n_j \geq 3$  events. The measured cross section ratio  $R_{32}$  is also corrected for detector smearing effects and unfolded to particle level. There can be two ways to obtain unfolded cross section ratio :

- **Method I** : First unfold separately the inclusive 2-jet and 3-jet measured cross sections and then construct the ratio  $R_{32}$
- **Method II** : Unfold directly the cross section ratio  $R_{32}$

In further analysis, unfolded cross section ratio  $R_{32}$  and its systematic uncertainties are calculated using Method I, whereas Method II is used only to propagate the statistical uncertainties including bin-by-bin correlations and statistical correlations between the inclusive 3-jet and 2-jet events cross sections. Unfolding takes the response matrix as an input which are explained in the next section.

### 1.5.1 Response Matrices

The response matrix is a two dimensional mapping between the true and measured distributions. It is usually derived from simulated Monte Carlo (MC) samples, which takes the true distribution from MC as an input and smears it by taking into account the detector resolution. Then this response matrix is used to unfold the measured

data spectrum. But there are several drawbacks of constructing response matrix using this method. In some phase space regions, the shape of the distribution is not well described by the LO predictions. Also, the limited number of events in the MC samples at high transverse momenta introduces high statistical fluctuations in the response matrix.

However, there is an indirect way of constructing the response matrix which uses a custom Toy Monte Carlo method. In this method, the particle level or true  $H_{T,2}/2$  spectrum is obtained by fitting the theoretically predicted NLO spectrum. Then this distribution is smeared with forward smearing technique, using the extracted jet energy resolution (JER) to obtain the reconstructed level or measured  $H_{T,2}/2$  spectrum. After that, the response matrix is constructed from these two distributions is used for the unfolding procedure.

#### 1.5.1.1 Inclusive Cross Sections

The NLO spectrum of the differential cross sections for  $n_j \geq 2$  and  $n_j \geq 3$  events obtained using CT10-NLO PDF set are fitted with the following two different functions defined in Eq. 1.10 and 1.13. These functions describes the shape as well as normalization of the distribution.

- **Function I :**

$$f(H_{T,2}/2) = N[x_T]^{-a}[1 - x_T]^b \times \exp[-c/x_T] \quad (1.10)$$

where N is normalization factor and a, b, c are fit parameters.

This function is derived from the below function [12] :

$$f(p_T; \alpha, \beta, \gamma) = N_0[p_T]^{-\alpha} \left[ 1 - \frac{2 p_T \cosh(y_{min})}{\sqrt{s}} \right]^\beta \times \exp[-\gamma/p_T] \quad (1.11)$$

using

$$\alpha = a, \quad \beta = b, \quad \gamma = c * \sqrt{s}/2, \quad x_T = \frac{2 * H_{T,2}/2 * \cosh(y_{min})}{\sqrt{s}} = \frac{2 * H_{T,2}/2}{\sqrt{s}} \quad (1.12)$$

where transverse scaling variable  $x_T$  corresponds to the proton fractional momentum  $x$  for dijets with rapidity  $y = 0$ ,  $\sqrt{s} = 8000$  GeV and  $y_{min}$  is low-edge of the rapidity bin  $y$  under consideration (here  $y_{min}$  is taken equal to 0)

• **Function II :**

$$f(H_{T,2}/2) = A_0 \left(1 - \frac{H_{T,2}/2}{A_6}\right)^{A_7} \times 10^{F(H_{T,2}/2)}, \text{ where } F(x) = \sum_{i=1}^5 A_i \left(\log\left(\frac{x}{A_6}\right)\right)^i \quad (1.13)$$

where the parameter  $A_6$  is fixed to  $\frac{\sqrt{s}}{2 \cosh(y_{min})}$ , where  $\sqrt{s} = 8000$  GeV and  $y_{min}$  is the minimum rapidity. The other parameters are derived from the fitting.

Figure 1.13 shows the fitted CT10-NLO spectrum of differential cross section as a function of  $H_{T,2}/2$  (green solid circles) using Function I (top) and using Function II (bottom) : for inclusive 2-jet (left) and inclusive 3-jet events (right). Function I is used primarily to generate response matrices and perform the closure tests and Function II is used as an alternative function to calculate unfolding uncertainty, described in Sec. 1.6.3. To include the migration to lower bins, the fit functions described by red lines are extrapolated to 80 GeV (blue dashed lines).

A flat  $H_{T,2}/2$  spectrum is generated by using toy Monte Carlo events and the fit parameters obtained from the NLO spectrum using function I (as shown in Fig. 1.13) provides weights to the flat spectrum. A total of ten million events are generated randomly (in  $H_{T,2}/2$  range 80-2000). These generated values are then smeared with a Gaussian function, where  $\sigma$  of the Gaussian is determined from the relative resolution

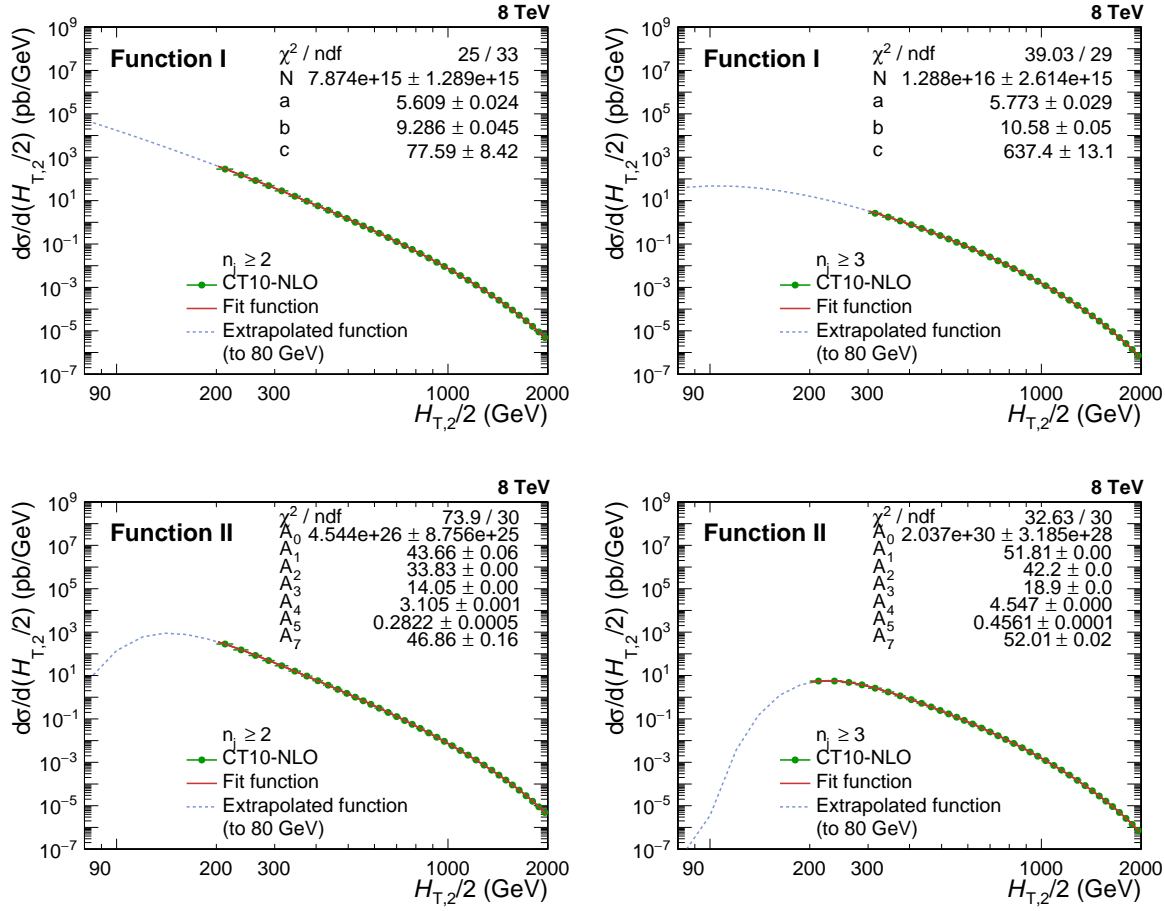


Figure 1.13: Fitted CT10-NLO spectrum of differential cross section as a function of  $H_{T,2}/2$  (green solid circles) using Function I (top) defined in Eq. 1.10 and using Function II (bottom) given by Eq. 1.13, for inclusive 2-jet (left) and inclusive 3-jet events (right). To consider the migration to lower  $H_{T,2}/2$  bins, the fit functions described by red lines are extrapolated to 80 GeV (blue dashed lines).

parametrization as a function of  $H_{T,2}/2$  calculated from NSC formula mentioned in equation 1.8. The parameters N, S, C used for smearing are taken from Table 1.6. These randomly generated ( $\text{Gen}_{\text{Toy}}$ ) and smeared ( $\text{Measured}_{\text{Toy}}$ ) values are used to fill the response matrices. Figure 1.14 shows the response matrices derived using the Toy MC for  $n_j \geq 2$  (left) and  $n_j \geq 3$  events (right). The matrices are normalized to the number of events in each column. The response matrices are diagonal as the migrations in off-diagonal bins are much smaller than the bins along the diagonal.

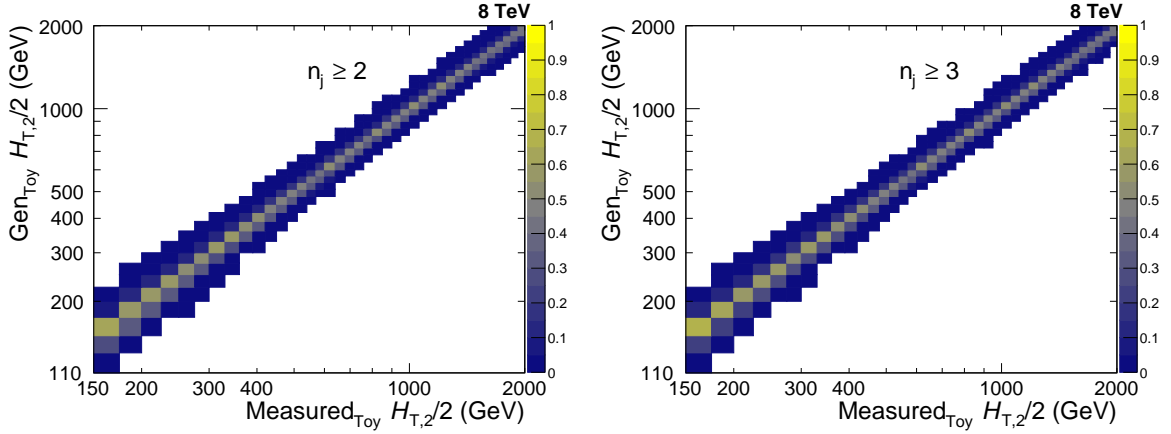


Figure 1.14: The response matrices are derived using the Toy Monte Carlo and forward smearing method, for inclusive 2-jet (left) and inclusive 3-jet events (right). The matrices are normalized to the number of events in each column and are diagonal with small off-diagonal migrations between close-by  $H_{T,2}/2$  bins.

### 1.5.1.2 Cross Section Ratio, $R_{32}$

To obtain the statistical uncertainty on the unfolded cross section ratio  $R_{32}$ , Method II is used. In this method, the response matrix is constructed using Toy MC method as done in Sec. 1.5.1.1 for differential cross sections. To obtain the true spectrum for  $R_{32}$ , the ratio of cross section spectrum described by Eq. 1.10 for inclusive 3-jet to that of 2-jet events is taken. This ratio is shown by green solid circles in Fig. 1.15 (left) which is fitted using a polynomial function of degree 8 (red line). Then as explained in above section, response matrix is derived for  $R_{32}$  using the Toy Monte Carlo and forward smearing method which is shown in Fig. 1.15 (right). The matrix is normalized to the number of events in each column and is diagonal with small off-diagonal migrations between close-by  $H_{T,2}/2$  bins.

## 1.5.2 Closure Test

A closure test has been performed to confirm the working of the unfolding procedure. In this test,  $\text{Measured}_{\text{Toy}}$  spectrum is unfolded using the constructed response matrices shown in Figure 1.14. It is expected that the same  $\text{Gen}_{\text{Toy}}$  spectrum should be re-obtained after unfolding. Figure 1.16 confirms that the unfolded  $\text{Measured}_{\text{Toy}}$

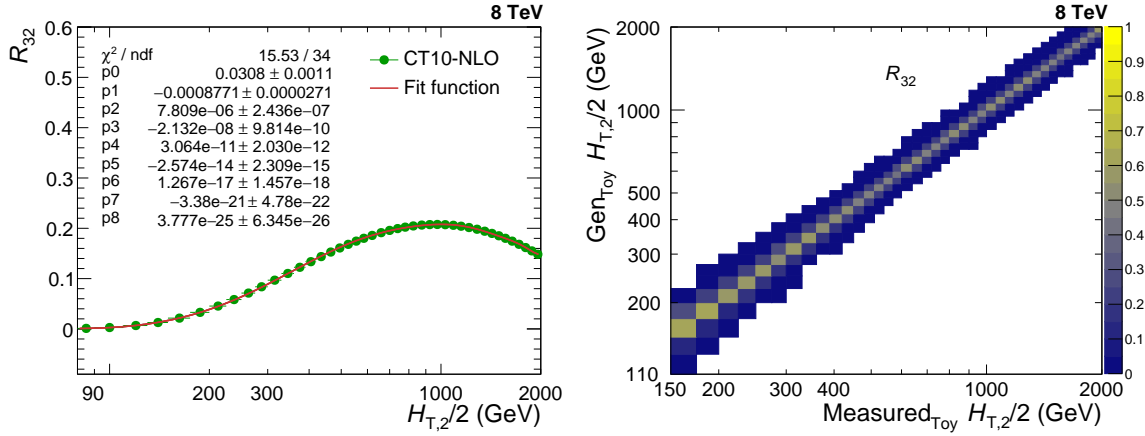


Figure 1.15: Left : The ratio of cross sections described by Eq. 1.10 for inclusive 3-jet to that of 2-jet events is shown as a function of  $H_{T,2}/2$  (green solid circles). It is fit using a polynomial function of degree 8 (red line). Right : The response matrix is derived using the Toy Monte Carlo and forward smearing method, for the cross section ratio  $R_{32}$ . The matrix is normalized to the number of events in each column and is diagonal with small off-diagonal migrations between close-by  $H_{T,2}/2$  bins.

spectrum matches exactly with  $\text{Gen}_{\text{Toy}}$  spectrum as the ratio of these distributions is perfectly flat at one for both  $n_j \geq 2$  (top left) and  $n_j \geq 3$  events (top right) cross sections as well as the cross section ratio  $R_{32}$  (bottom).

For another closure test, Reco MG5+P6 MC differential cross section distribution is unfolded using the above constructed response matrices using JER for forward smearing the randomly generated spectrum. While taking ratio of the unfolded distribution to that of Gen MG5+P6 MC, it is observed that a well closure is not obtained. This is represented by blue line in Fig. 1.17 for  $n_j \geq 2$  (top left) and  $n_j \geq 3$  events (top right). As observed in Fig. 1.12 in Sec. 1.4, if Reco MG5+P6 MC is unfolded using the response matrices obtained using 30% reduced JER, then the good closure is obtained as shown by red line in Fig. 1.17. Since unfolded cross section ratio  $R_{32}$  is the ratio of unfolded differential cross sections (Method I), same behaviour is observed for  $R_{32}$  (bottom).



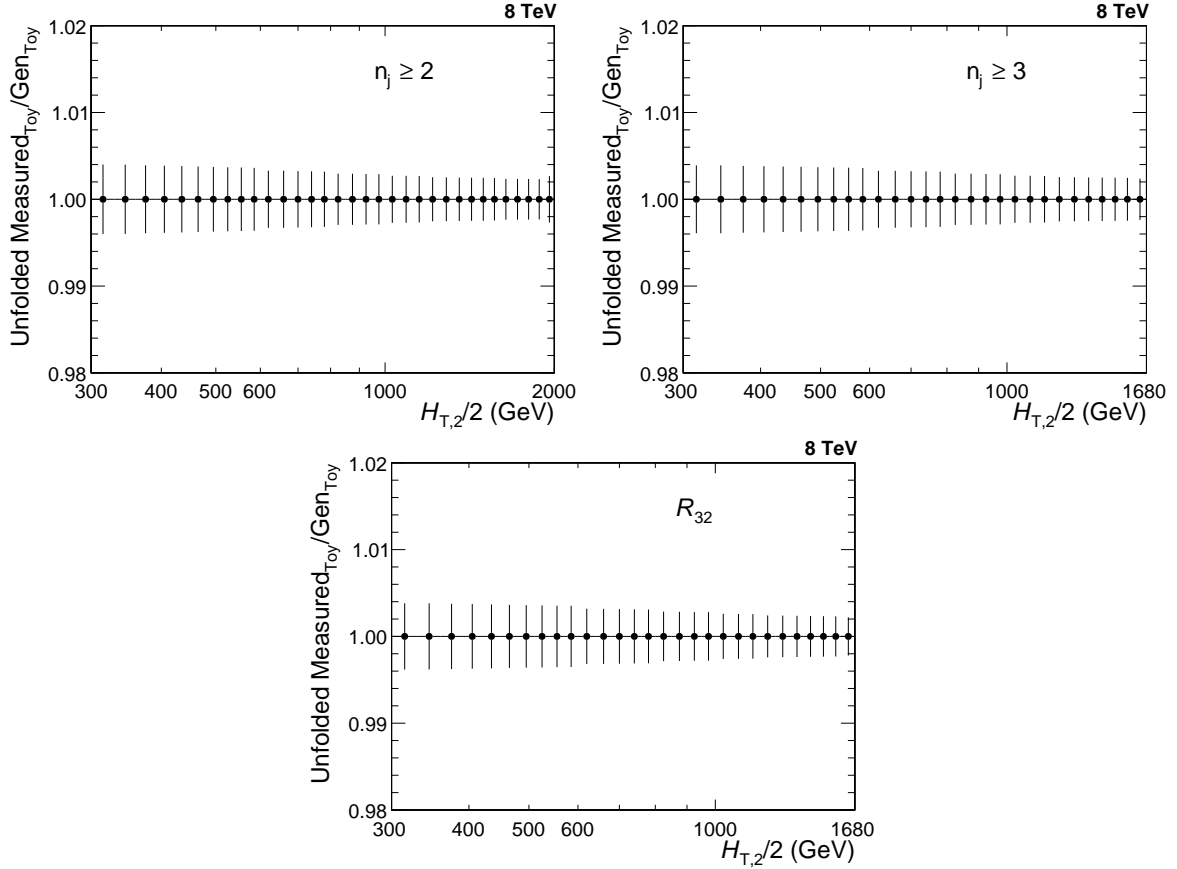


Figure 1.16: Closure test of the unfolding technique where the smeared spectrum obtained from Toy Monte Carlo method ( $\text{Measured}_{\text{Toy}}$ ), is unfolded using the constructed response matrices (obtained by forward smearing the randomly generated spectrum ( $\text{Gen}_{\text{Toy}}$ ) using extracted jet energy resolution (JER)). As expected, the unfolded  $\text{Measured}_{\text{Toy}}$  spectrum matches exactly with  $\text{Gen}_{\text{Toy}}$  spectrum as the ratio of these distributions is perfectly flat at one for both inclusive 2-jet (top left) and inclusive 3-jet events (top right) cross sections as well as the cross section ratio  $R_{32}$  (bottom).

### 1.5.3 Unfolding of the Measurement

After validating the unfolding method, the measured differential cross sections as well as  $R_{32}$  are unfolded using the above reconstructed response matrices. The unfolded data spectrum is compared to that of measured one in Fig. 1.18 for  $\eta_j \geq 2$  (top left) and  $\eta_j \geq 3$  events (top right) cross sections and for the cross section ratio  $R_{32}$  (bottom). As already discussed that 30% reduced JER gives better closures than JER, so the unfolding of data is done with response matrices using JER (blue solid

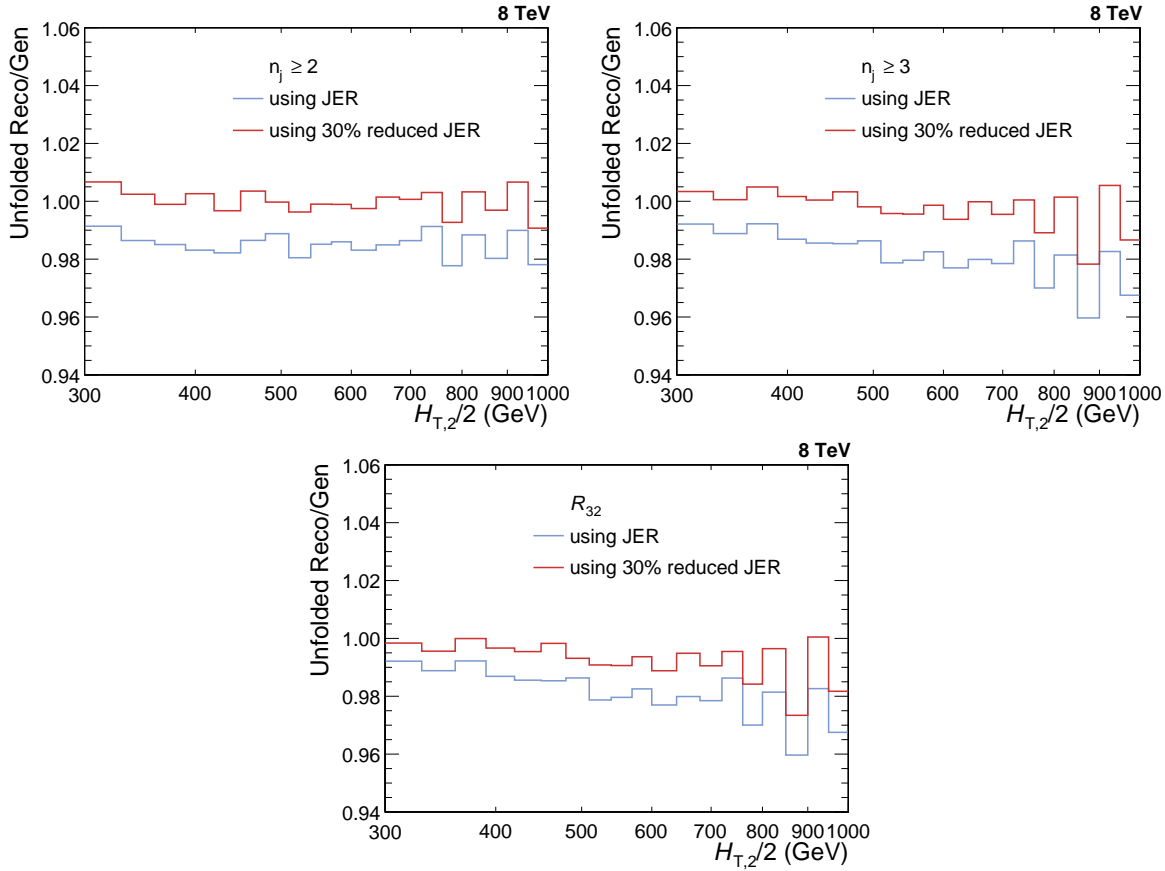


Figure 1.17: Reco MADGRAPH5+PYTHIA6 Monte Carlo (MG5+P6 MC) differential cross section distributions unfolded with the response matrices (obtained by forward smearing the randomly generated spectrum (Gen) using extracted jet energy resolution (JER)), does not give a good closure with Gen MG5+P6 MC (blue line), for inclusive 2-jet (top left) and inclusive 3-jet events (top right). After performing the unfolding using 30% reduced JER, a good closure is obtained (red line). Since unfolded the cross section ratio  $R_{32}$  is the ratio of unfolded differential cross sections, same behaviour is observed for  $R_{32}$  (bottom).

circles) as well as 30% reduced JER (red solid circles) for smearing. The difference between both is taken as an additional uncertainty on the unfolded measurement.

## 1.6 Experimental Uncertainties

In an experimental measurement of any physical observable, the uncertainties play a key role and hence are important to study in a physics analysis. The uncertainties can be categorized into two types : statistical and systematic. The statistical

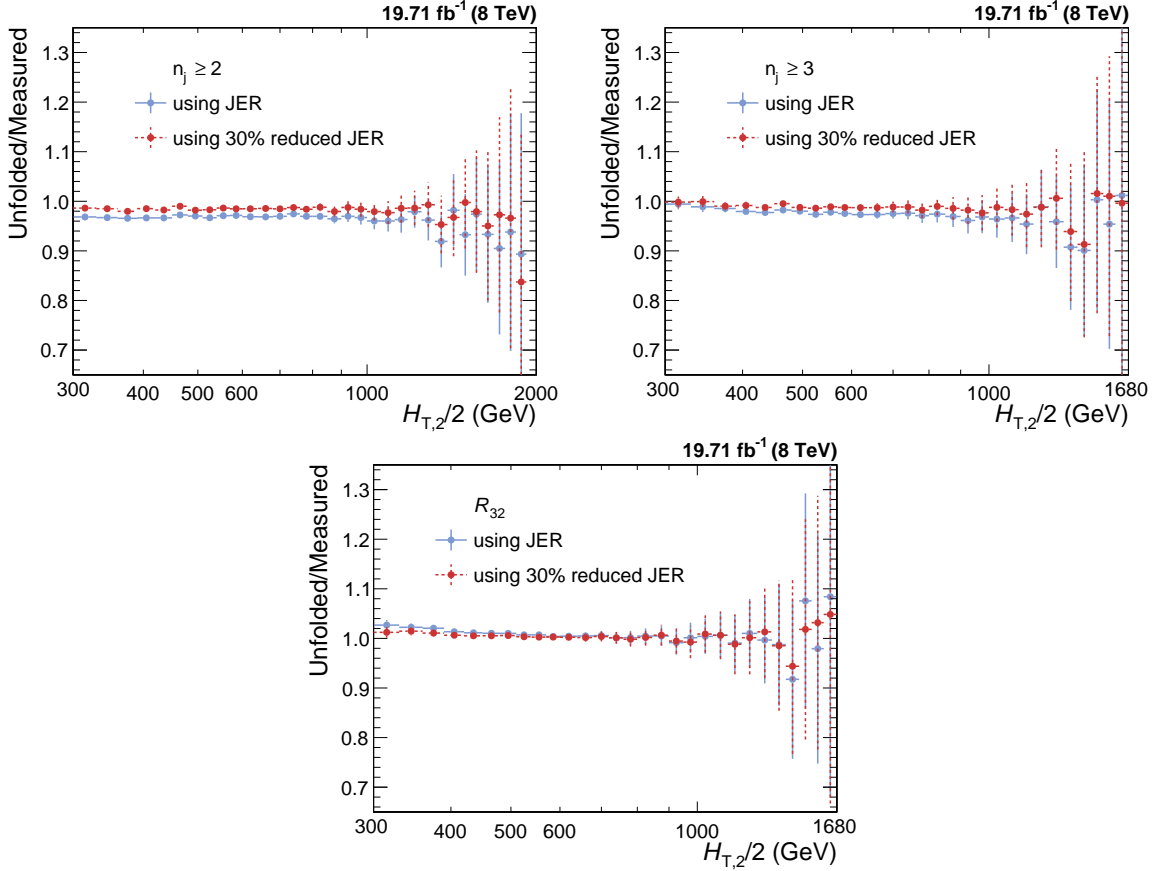


Figure 1.18: The measured differential cross sections as well as the cross section ratio  $R_{32}$  are unfolded as a function of  $H_{T,2}/2$  using the response matrices derived using the Toy Monte Carlo and forward smearing method. The unfolded spectrum are compared with that of the measured one for inclusive 2-jet (top left) and inclusive 3-jet events cross sections (right) as well as for  $R_{32}$  (bottom). The unfolding is done with response matrices using JER (blue solid circles) as well as 30% reduced JER (red solid circles) for smearing. The difference between both is taken as an additional uncertainty on the unfolded measurement.

uncertainties arise due to random fluctuations depending on the number of events. The more the number of events, less is the statistical uncertainty. The systematic uncertainties may be due to known detector effects, model dependence, assumptions made or various corrections applied. In general, if the statistical and systematic uncertainties are uncorrelated, these can be added in quadrature to obtain the total uncertainty on the measurement. In this section, all the experimental uncertainties affecting the measurement of cross sections and the cross section ratio  $R_{32}$  are described. The systematic experimental uncertainties for  $R_{32}$  are propagated from the cross sections to the ratio taking into account correlations. Due to this, the

systematic uncertainties may cancel for  $R_{32}$  completely or partially as compared to those for the individual cross sections.

### 1.6.1 Statistical Uncertainty

Statistical uncertainty on the measurement is obtained through the unfolding procedure using a toy MC method. The measured data points are smeared within their statistical uncertainties to get the smeared spectrum. Such smeared spectra are produced million in number and the unfolding is performed multiple times for each smeared spectra. The differences between the unfolded spectra and the measured one give the statistical uncertainty. The unfolding process introduces more statistical fluctuations which can be observed in Fig. 1.19. Here the fractional statistical uncertainties of the unfolded data (red line) are compared with those of the measured one (blue line) for  $n_j \geq 2$  (top left) and  $n_j \geq 3$  events cross sections (top right) as well as for the cross section ratio  $R_{32}$  (bottom).

After the unfolding, the final statistical uncertainties become correlated among the bins such that the size of these correlations varies between 10 and 20%. The correlation (anti-) is more significant for neighbouring bins in  $H_{T,2}/2$  as compared to the far off ones. In Fig. 1.20, the correlations of the statistical uncertainty after the unfolding can be seen for  $n_j \geq 2$  (top left) and  $n_j \geq 3$  events cross sections (top right) and for the cross section ratio  $R_{32}$  (bottom). These correlations must be considered while performing the fits to extract the value of the strong coupling constant,  $\alpha_S$ .

### 1.6.2 Jet Energy Corrections Uncertainty

As explained in Sec. ?, the measured jet energy is corrected for a variety of detector effects by applying jet energy corrections (JEC) [13]. There are 25 mutually independent sources which contribute to JEC. Each source presents a  $1\sigma$  shift and is

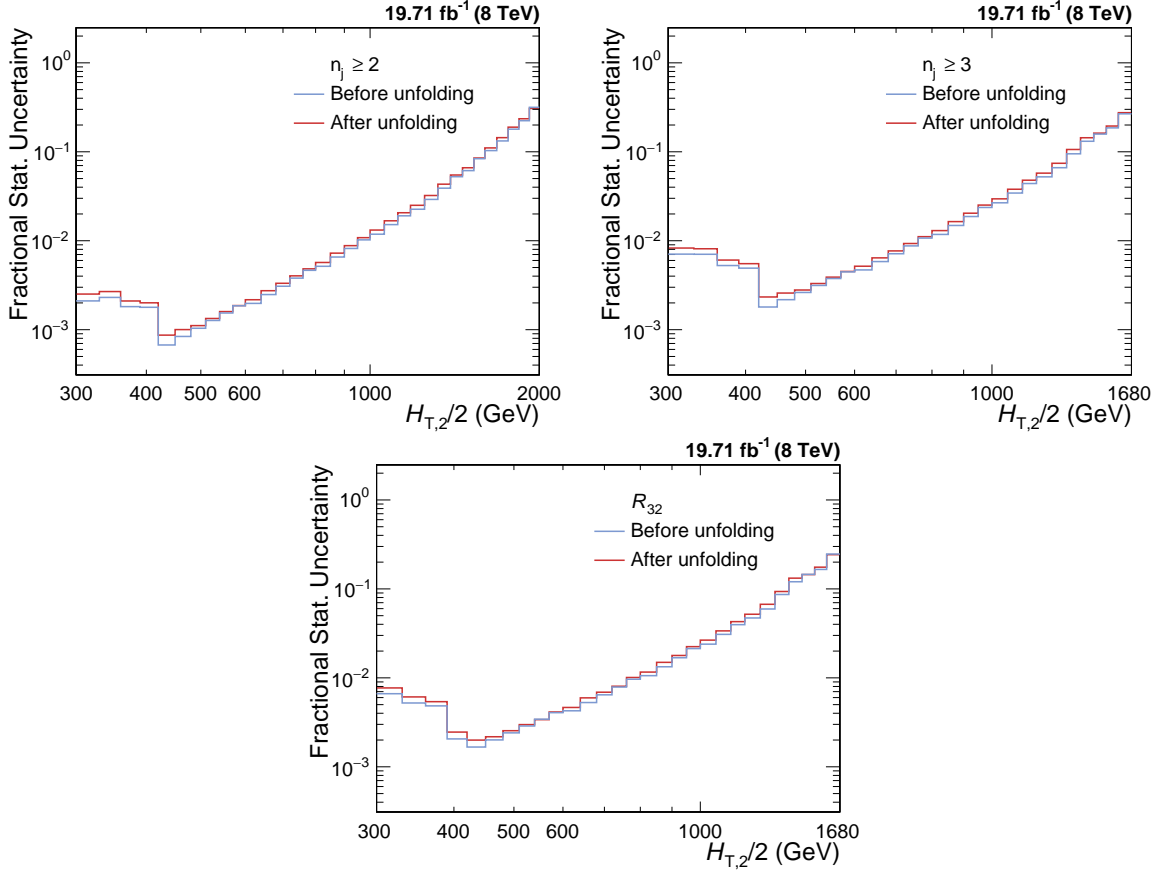


Figure 1.19: The fractional statistical uncertainties of the unfolded data (red line) are compared with those of the measured one (blue line) for inclusive 2-jet (top left) and inclusive 3-jet events cross sections (top right) as well as for the cross section ratio  $R_{32}$  (bottom). After unfolding, the statistical uncertainty increases slightly.

fully correlated in  $p_T$  and  $\eta$  but uncorrelated to all other sources. The observable is studied with the nominal values of the jet energy which gives nominal distributions as well as by varying up and down the energy of all jets by the uncertainty. The differences between the nominal distributions and the ones obtained by varying the jet energy gives the uncertainties from each source. The JEC uncertainties can be asymmetric in nature which leads to separate treatment of upwards and downwards variation of each source. The sum in quadrature of uncertainties from all sources gives the total JEC uncertainty. In the current analysis, JEC uncertainties are a dominant source of experimental uncertainty at low  $H_{T,2}/2$ . The JEC uncertainty ranges from 3% to 10% for  $n_j \geq 2$  and from 3% to 8% for  $n_j \geq 3$  events cross sections

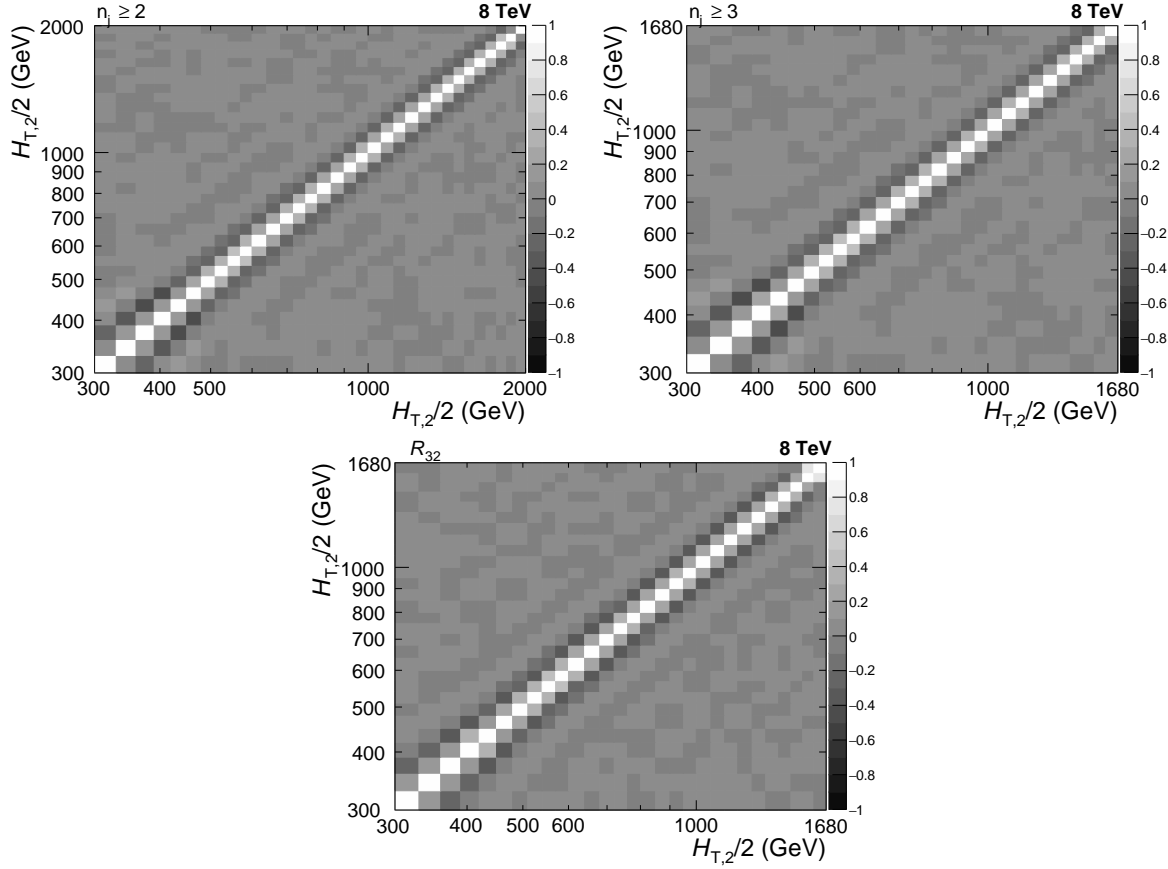


Figure 1.20: The unfolding procedure introduces the correlations of the statistical uncertainty through bin migrations which are shown here for inclusive 2-jet (top left) and 3-jet events cross sections (top right) as well as for the cross section ratio  $R_{32}$  (bottom). The correlation (anti-) is more significant between neighbouring bins than far-off ones.

measurement. To calculate JEC uncertainty for ratio  $R_{32}$ , the inclusive 2-jet and 3-jet events cross sections are measured as a function of  $H_{T,2}/2$  by shifting the jet  $p_T$  according to the JEC uncertainty for each source of JEC separately. Then the ratio of these cross sections is taken and the difference of these from the central ratio  $R_{32}$ , gives the JEC uncertainty for  $R_{32}$ . As expected, JEC uncertainty for  $R_{32}$  is small as compared to that for individual cross sections and is about 1 to 2% over all  $H_{T,2}/2$  bins.

The sources of JEC considered in this analysis are : AbsoluteStat, AbsoluteScale, AbsoluteFlavMap, AbsoluteMPFBias, Fragmentation, SinglePionECAL, SinglePionHCAL, FlavorQCD, RelativeJEREC1, RelativeJEREC2, Rel-

ativeJERHF, RelativePtBB, RelativePtEC1, RelativePtEC2, RelativePtHF, RelativeFSR, RelativeStatFSR, RelativeStatEC2, RelativeStatHF, PileUpDataMC, PileUpPtRef, PileUpPtBB, PileUpPtEC1, PileUpPtEC2 and PileUpPtHF. The AbsoluteFlavMap uncertainty is exactly zero for the 8 TeV and can be ignored. For the four sources : RelativeJERHF, RelativePtHF, RelativeStatHF, PileUpPtHF, the JEC uncertainty is exactly zero because of  $|y| < 2.5$  cut used in the analysis. So only 20 sources contribute to the total JEC uncertainty. The Figs. A.1-A.3 show the JEC uncertainty from each source separately for inclusive 2-jet (top) and 3-jet events cross sections (middle) as for cross section  $R_{32}$  (bottom). Depending on the origin of sources, they are categorized into four groups which are described below in brief :

1. **Pileup** : This uncertainty originates from the differences in the transverse momentum between the true offset and the Random Cone method (i.e. essentially difference of pile-up inside and outside of jets), in simulated events. This uncertainty is derived from  $Z/\gamma$ +jet, dijet and multijet data using fit procedure to estimate the residual pileup uncertainty after the calibration.
2. **Relative** : The forward jets are calibrated by the relative  $\eta$ -dependent corrections using dijet events. The main contribution to the uncertainty comes from jet energy resolution (JER), derived by varying JER scale factors up and down by quoted uncertainties and the initial and final state radiation bias corrections.
3. **Absolute** : A global fit to  $Z/\gamma$ +jet and multi-jet events gives the absolute calibration of the jet energy scale. The uncertainties are related to the lepton momentum scale for muons in  $Z (\rightarrow \mu\mu)$ +jet and the single pion response in the HCAL.
4. **Flavor** : Flavor response differences are studied from simulation by cross-checking the results with quark- and gluon-tagged  $\gamma$ +jet and  $Z$ +jet events.

These uncertainties are based on PYTHIA6.4 and HERWIG++2.3 differences propagated through the data-based calibration method.

The details of the jet energy corrections and uncertainties can be found in [14].

### 1.6.3 Unfolding Uncertainty

The unfolding uncertainty is comprised of three uncertainties which are explained as follows :

1. **Jet Energy Resolution :** The calculation of the jet energy resolution (JER) using simulated MG5+P6 Monte Carlo events is already explained in Sec. 1.4. As mentioned before, the measured jet transverse momentum ( $p_T$ ) in simulated MC events needs to be smeared additionally to match the resolution in data. This smearing is done by using measured scale factors ( $c_{central}$ ) mentioned in Table 1.5. It is recommended by JETMET group that the uncertainty on these measured scaling factors must be taken into account in a physics analysis. Since JER is used in constructing the response matrix which is an input in unfolding procedure, so the uncertainty on scale factors accounts for the unfolding uncertainty. To calculate JER uncertainty,  $p_T$  is smeared with two additional sets of scale factors corresponding to varying the factors up and down by one sigma, and corresponding  $H_{T,2}/2$  is calculated. Then again JER is calculated as a function of  $H_{T,2}/2$  using these upwards ( $c_{up}$ ) and downwards ( $c_{down}$ ) variations of the scaling factors. Alternative response matrices are built using the JER with above variations and the unfolding is performed again. The differences of the obtained unfolded spectrums to the nominal ones accounts for a systematic JER uncertainty.
2. **Model Dependence :** It is explained in Sec. 1.5.1 that to obtain the true  $H_{T,2}/2$  spectrum to be used in constructing response matrix using Toy MC



method, the fitting of the CT10-NLO predictions is performed with the Function I described in Eq. 1.10. Using the alternative function, Function II given by Eq. 1.13, for this fitting and then constructing different response matrix, gives the model dependence of the true  $H_{T,2}/2$  spectrum. The differences in unfolded distributions using the above mentioned two different response matrices gives the model dependence uncertainty.

3. **Additional Uncertainty :** Small nonclosures observed in Fig. 1.12 introduces a supplementary uncertainty which is attributed by comparison of distributions unfolded using response matrices constructed using JER from simulation with that obtained with a 30% reduced JER.

All the three above mentioned uncertainties are added in quadrature to get the total unfolding uncertainty which increases from about 1% at low  $H_{T,2}/2$  up to 2% at the high  $H_{T,2}/2$  ends of the cross sections for both  $n_j \geq 2$  and  $n_j \geq 3$  events. This uncertainty account for about less than 1% for  $R_{32}$ .

#### 1.6.4 Luminosity Measurement Uncertainty

The measurement of the luminosity delivered to CMS detector by LHC in the proton-proton collisions in the year of 2012 is done by using the silicon pixel cluster counting method [15]. The uncertainty related to the integrated luminosity measurement is estimated to be 2.5% (syst.) and 0.5% (stat.). This uncertainty propagates directly to any absolute cross section measurement. Hence, a total systematic uncertainty of 2.6% is considered across all the  $H_{T,2}/2$  bins. At low  $H_{T,2}/2$ , it is similar in size as the one from JEC. This uncertainty cancels completely for  $R_{32}$ .

### 1.6.5 Residual Uncertainty

The small trigger and jet identification inefficiencies account for smaller than 1% uncertainties on the cross section measurements [16, 17]. Hence, an uncorrelated residual uncertainty of 1% is assumed across all  $H_{T,2}/2$  bins for both  $n_j \geq 2$  and  $n_j \geq 3$  events cross sections whereas for  $R_{32}$ , it gets cancel completely.

### 1.6.6 Total Experimental Uncertainty

After calculating the uncertainties from all the above mentioned sources, the total experimental uncertainty on measurement of cross sections as well as cross section ratio  $R_{32}$ , is obtained by adding in quadrature the uncertainties from individual sources. The values of uncertainties (in %) from each source as well as total uncertainty, for each  $H_{T,2}/2$  bin, are tabulated in Tables A.2, A.3 and A.4 for  $n_j \geq 2$  and  $n_j \geq 3$  events cross sections and cross section ratio  $R_{32}$ , respectively. Figure 1.21 shows the experimental uncertainties, from different sources as well as the total uncertainty, affecting the measurement of  $n_j \geq 2$  (top left) and  $n_j \geq 3$  events cross sections (top right) and cross section ratio  $R_{32}$  (bottom). The error bars represent the statistical uncertainty obtained after unfolding. The systematic uncertainties due to jet energy corrections (JEC by blue line), luminosity (red dashed line), unfolding (green dashed line) and residual effects (light purple line) are also presented. The uncertainties due to luminosity and residual effects cancel completely in  $R_{32}$ . The total uncertainty (black dashed line) on the measurements is asymmetric in nature and dominated by the uncertainty due to the jet energy corrections (JEC) at lower  $H_{T,2}/2$  values and by statistical uncertainty at higher  $H_{T,2}/2$  values. The experimental uncertainties from each source as well as total uncertainty are also quoted in Table 1.7.

The complete data analysis of the differential inclusive 2-jet and inclusive 3-jet events cross sections as well as their ratio  $R_{32}$  has been presented as a function of

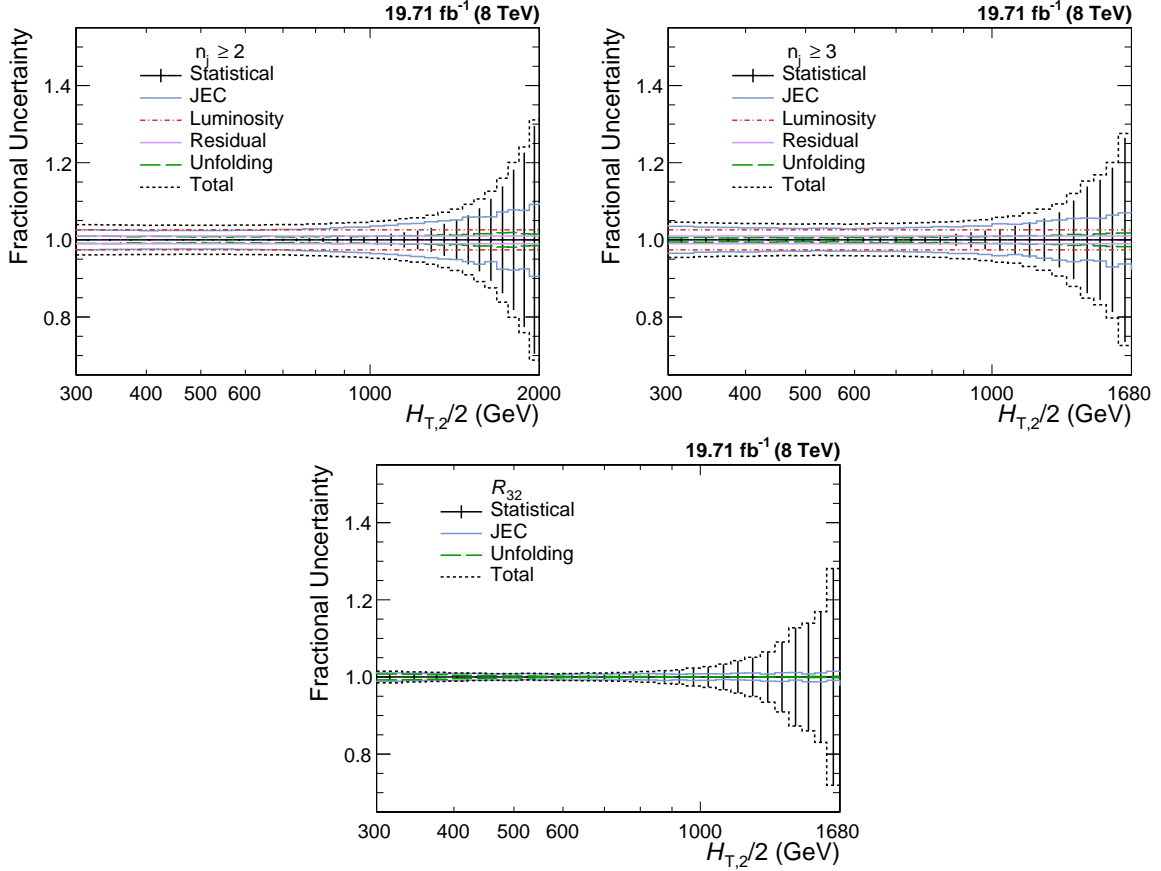


Figure 1.21: Experimental uncertainties from different sources affecting the measurement of cross sections for inclusive 2-jet (top left) and inclusive 3-jet events (top right) and cross section ratio  $R_{32}$  (bottom). The error bars represent the statistical uncertainty after unfolding. The systematic uncertainties due to jet energy corrections (JEC by blue line), luminosity (red dashed line), unfolding (green dashed line) and residual effects (light purple line) are also presented. The uncertainties due to luminosity and residual effects cancel completely in  $R_{32}$ . The total uncertainty (black dashed line) is calculated by adding in quadrature the individual sources of uncertainty.

$H_{T,2}/2$ . The measured spectrums after correcting for detector effects through the unfolding procedure, are compared with the next-to-leading order (NLO) pQCD calculations in the next chapter.

Table 1.7: Overview of all experimental uncertainties affecting the measurement of cross sections for inclusive 2-jet (left) and inclusive 3-jet events (middle) and cross section ratio  $R_{32}$  (right). The uncertainties due to luminosity and residual effects cancel completely in  $R_{32}$ . The total uncertainty is calculated by adding in quadrature the individual sources of uncertainty.

Uncertainty Source	Inclusive 2-jet	Inclusive 3-jet	$R_{32}$
Statistical	< 1 to 30%	< 1 to 27%	< 1 to 28%
Jet energy corrections (JEC)	3 to 10%	3 to 8%	1 to 2%
Unfolding	1 to 2%	1 to 2%	< 1%
Luminosity	2.6%	2.6%	cancels
Residual	1%	1%	cancels
Total	4 to 32%	4 to 28%	1 to 28%

# Chapter 2

## Theoretical Calculations

In an experiment, the measurements are validated by doing the comparison with the perturbative QCD (pQCD) theoretical calculations. The lowest order (LO) calculations describe well the shapes of the measured distributions but not the normalization due to the dependence on the unphysical renormalization ( $\mu_r$ ) and factorization ( $\mu_f$ ) scales. The next-to-leading order calculations (NLO) improves the precision by reducing the dependence on  $\mu_r$  and  $\mu_f$  scales and become an essential feature in the determination of fundamental parameters such as  $\alpha_S$  and the parton density distributions. In this chapter, the next-to-leading order pQCD calculations are described in details. NLO pQCD calculations are corrected for the multiparton interactions (MPI) and hadronization effects by applying non-perturbative (NP) corrections and also corrected for the electroweak interactions (EW).

### 2.1 Fixed Order NLO Calculations

The predictions of the inclusive differential jet event cross section at NLO accuracy in pQCD are computed with the NLOJET++ program version 4.1.3 [18, 19]. The results are provided within the framework of FASTNLO version 2.3 [20, 21] for use within fits. The parton distribution functions (PDFs) are accessed through the

LHAPDF6 library [22, 23]. The FASTNLO is preferred over the direct calculation with NLOJET++ as the calculations of the cross sections can be repeated several times with different PDFs as well as scale choices required for the calculating PDF and scale uncertainties. The renormalization and factorization scales are chosen equal to  $H_{T,2}/2$ , i.e.  $\mu_r = \mu_f = H_{T,2}/2$ .

In the current study, different PDF sets available for a series of different assumptions on  $\alpha_s(M_Z)$  are used for NLO calculations. In Table 2.1, already existing PDF sets in LHC Run 1 (upper rows) and newer ones for Run 2 (lower rows) are listed together with the corresponding number of flavours  $N_f$ , the assumed masses  $M_t$  and  $M_Z$  of the top quark and the  $Z$  boson, respectively, the default values of  $\alpha_s(M_Z)$ , and the range in  $\alpha_s(M_Z)$  variation available for fits. All sets employ a variable-flavour number scheme with at most five or six flavours apart from the ABM11 PDFs, which use a fixed-flavour number scheme with  $N_F = 5$ . Mainly CT10 PDF set is considered for comparison between data and theory predictions as well as for calculating theoretical uncertainties. Out of these eight PDF sets the following three are not considered further because of the below mentioned reasons :

- At NLO, predictions based on ABM11 do not describe LHC jet data at small jet rapidity [24–27].
- The HERAPDF2.0 set exclusively fits HERA DIS data with only weak constraints on the gluon PDF.
- The range in values available for  $\alpha_s(M_Z)$  is too limited for the NNPDF3.0 set.

### 2.1.1 NLO Correction Factors

The differences between LO predictions and NLO predictions give the impact of the higher-order contributions to the pQCD predictions. These are described by a NLO

Table 2.1: NLO PDF sets are available via LHAPDF6 with various assumptions on the value of  $\alpha_s(M_Z)$ . The already existing sets in LHC Run 1 (upper rows) and newer ones for Run 2 (lower rows) are listed here with the corresponding number of flavours  $N_f$ , the assumed masses  $M_t$  and  $M_Z$  of the top quark and the  $Z$  boson, respectively, the default values of  $\alpha_s(M_Z)$ , and the range in  $\alpha_s(M_Z)$  variation available for fits. A \* behind the  $\alpha_s(M_Z)$  values signifies that the parameter was fixed, not fitted.

Base set	$N_F$	$M_t$ ( GeV)	$M_Z$ ( GeV)	$\alpha_s(M_Z)$	$\alpha_s(M_Z)$ range
ABM11 [28]	5	180	91.174	0.1180	0.110–0.130
CT10 [29]	$\leq 5$	172	91.188	0.1180*	0.112–0.127
MSTW2008 [30, 31]	$\leq 5$	$10^{10}$	91.1876	0.1202	0.110–0.130
NNPDF2.3 [32]	$\leq 6$	175	91.1876	0.1180*	0.114–0.124
CT14 [33]	$\leq 5$	172	91.1876	0.1180*	0.113–0.123
HERAPDF2.0 [34]	$\leq 5$	173	91.1876	0.1180*	0.110–0.130
MMHT2014 [35]	$\leq 5$	$10^{10}$	91.1876	0.1180*	0.108–0.128
NNPDF3.0 [36]	$\leq 5$	173	91.2	0.1180*	0.115–0.121

correction factor, k-factor, which is defined as the ratio of cross sections at NLO accuracy to that at LO i.e.

$$\text{k-factor} = \frac{\sigma_{\text{NLO}}}{\sigma_{\text{LO}}} \quad (2.1)$$

The size of k-factor determine the effect of the higher-order corrections. The small size of k-factor indicates that the cross section predictions are precisely described at the LO whereas the larger size hints the contributions from NLO. Figure 2.1 shows the k-factors of the NLOJET++ calculations, for inclusive 2-jet and 3-jet events cross sections and their ratio  $R_{32}$ , using five different PDF sets. k-factor for  $R_{32}$  is obtained by taking the ratio of k-factors for inclusive 3-jet events cross sections to that of inclusive 2-jet. The k-factors are similar for all the PDF sets in the lower region, but the differences increase in regions with larger  $H_{T,2}/2$ . It is observed that for inclusive 3-jet events cross sections, k-factor jumps at lowest  $H_{T,2}/2$ . This is because some jet configurations are kinematically forbidden near the  $p_T$  cut bin i.e. 150 GeV. Since the first few bins in  $H_{T,2}/2$  (below 225 GeV) still suffer from these

kinematical effects, the minimum value of  $H_{T,2}/2$  studied is 300 GeV.

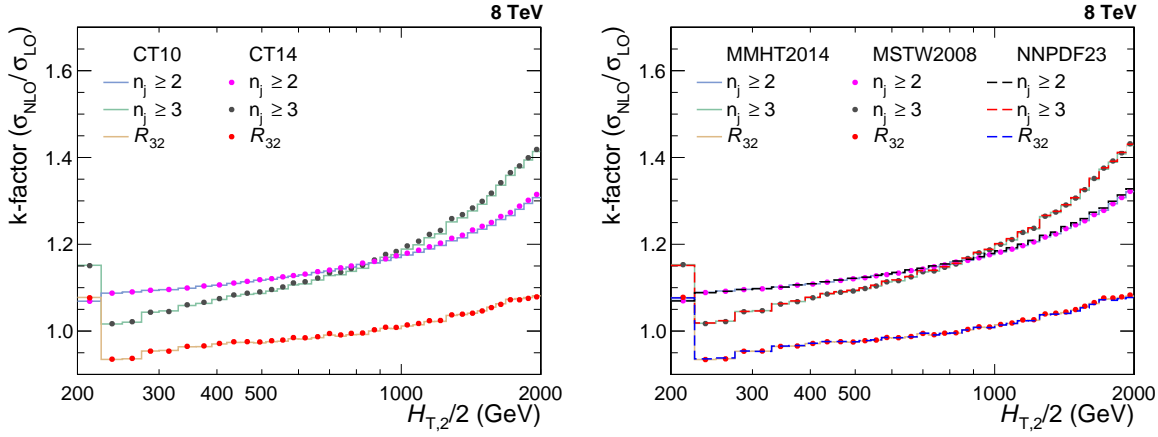


Figure 2.1: The k-factors of the NLOJET++ calculations, for inclusive 2-jet and 3-jet events cross sections and their ratio  $R_{32}$ , using five different PDF sets.

### 2.1.2 Non-Perturbative Corrections

The fixed-order pQCD calculations predict the parton-level cross section but lacks accuracy due to several effects. The partons which are emitted close to each other in phase space are not handled well in lower order perturbation theories and hence requires a parton shower (PS) correction. The scattering phenomena between partons within a colliding proton, other than the hard scattering, are known as multiparton interactions (MPI). The partons of the hard scattering forms colorless bound states called hadrons through a process of hadronization (HAD). The MPI and hadronization cannot be modelled well within the perturbative framework. Since the fixed-order NLO calculations do not include these additional soft QCD effects, these calculations cannot be compared directly to unfolded data. So these calculations should be corrected for non-perturbative effects (NP). The cross section ratio between a nominal event generation interfaced to the simulation of UE contributions and a sample without hadronization and MPI effects are taken as correction factors and are defined as :

$$C^{\text{NP}} = \frac{\sigma^{\text{PS+HAD+MPI}}}{\sigma^{\text{PS}}} \quad (2.2)$$



In this analysis, the NP effects are evaluated by using samples obtained from different MC event generators with a simulation of parton shower and underlying-event (UE) contributions. The leading order (LO), HERWIG++ [37] with the default tune of version 2.3 and PYTHIA6 [2] with tune Z2\*, and the NLO, POWHEG [38–40], MC event generators are considered. The matrix-element calculation performed with POWHEG is interfaced to PYTHIA8 with tune CUETS1 [41] for the UE simulation. The ratio, defined in Eq. 2.2, is obtained for each MC generator and is fitted by a power-law function defined in Eq. 2.3. Since this ratio obtained from different MC generators have large differences, so the average of the envelope, which covers all the differences, is taken as the correction factor which is then applied as bin-by-bin multiplicative factor to the parton-level NLO cross section. The half of the envelope it is taken as the uncertainty on the NP correction factor.

$$f(H_{T,2}/2) = a \cdot (H_{T,2}/2)^b + c \quad (2.3)$$

The NP correction factors,  $C_{3\text{-jet}}^{\text{NP}}$  and  $C_{2\text{-jet}}^{\text{NP}}$  are calculated for  $n_j \geq 2$  and  $n_j \geq 3$  event cross sections respectively and then their ratio gives the correction factor for  $R_{32}$ . These are shown in Fig. 2.2 for the inclusive 2-jet (top left) and 3-jet event cross sections (top right), as well as for cross section ratio  $R_{32}$  (bottom). The NP corrections amount to  $\sim 4\text{--}5\%$  for inclusive 2-jet and 3-jet events cross section and  $\sim 1\%$  for ratio  $R_{32}$ , for  $H_{T,2}/2 \sim 300$  GeV and decrease rapidly for increasing  $H_{T,2}/2$ . On comparing the NP correction factors of cross section ratio with that for individual cross sections, it has been observed that the non-perturbative effects get reduced in the cross section ratio.

### 2.1.3 Electroweak Corrections

In LHC, the centre-of-mass energy of proton-proton collisions is well beyond the electroweak (EW) scale  $\sim O(100 \text{ GeV})$ . At such a high energy, the impact of higher

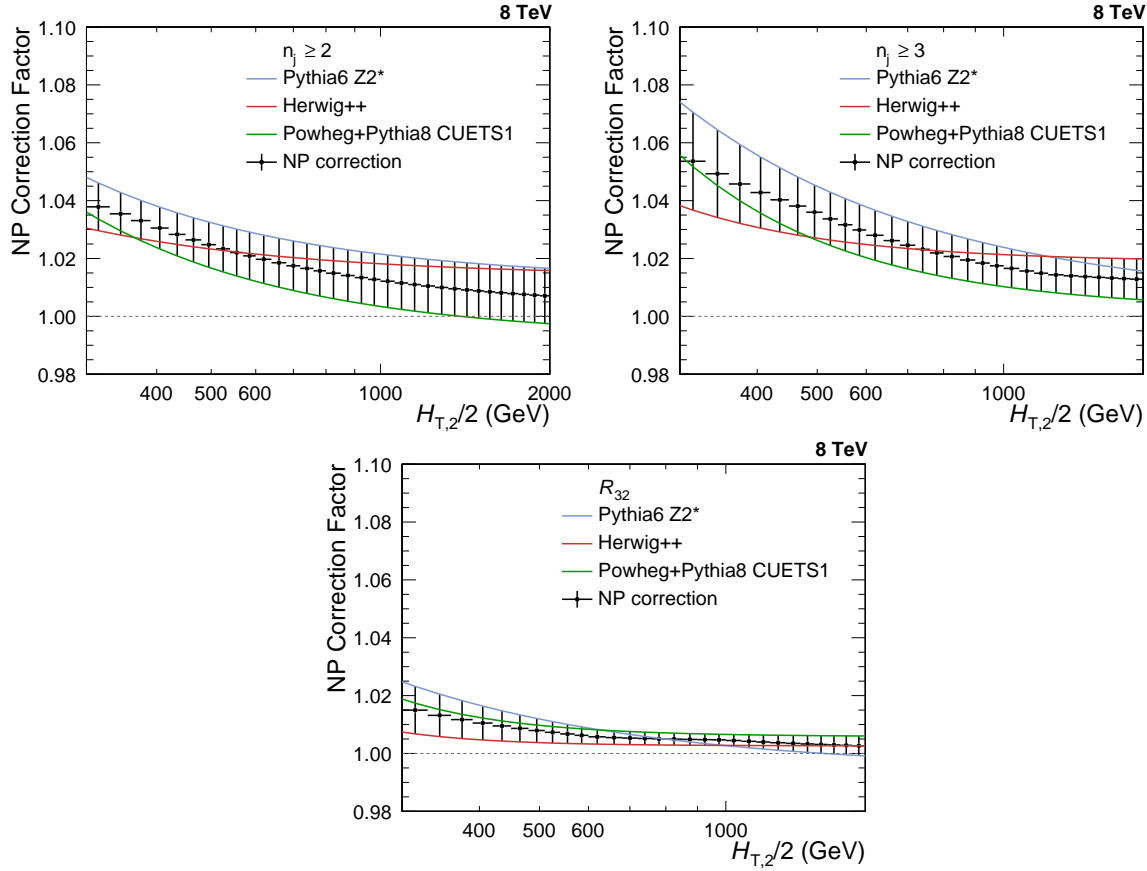


Figure 2.2: The nonperturbative (NP) corrections are presented as a function of  $H_{T,2}/2$  for inclusive 2-jet (top left) and 3-jet (top right) event cross sections, as well as their ratio  $R_{32}$ . These corrections are calculated from the leading order HERWIG++ with the default tune of version 2.3 (red line) and PYTHIA6 with tune Z2\* (blue line); and the next-to-leading order POWHEG interfaced to PYTHIA8 with tune CUETS1 (green line) Monte Carlo event generators. The black solid circles give the average NP correction factor along with the uncertainty shown by the error bars.

order EW corrections is much more with respect to QCD effects [42] and affect jet cross sections at large jet  $p_T$ . The quark-quark scattering processes involving virtual exchanges of massive W and Z bosons contribute to electroweak (EW) corrections. The fixed-order QCD calculations do not include EW corrections and hence the NLO theory calculations are corrected for EW effects. The EW corrections have been calculated for inclusive 1-jet and 2-jet case, in Ref. [43] but are not available yet for inclusive 3-jet production. The EW correction factors in the phase space of the measurement are shown as a function of  $H_{T,2}/2$  in Fig. 2.3 for inclusive 2-jet

events cross sections. These correction factor increases up to 13% at high ends of  $H_{T,2}/2$  and significantly improves the agreement between data and prediction. Since the guess from theory side is that EW for inclusive 2-jet and 3-jet will be similar, so for  $R_{32}$ , it is assumed to be equal to the factor of 1. These corrections are applied as a bin-by-bin correction factor to the fixed-order calculation of NLOJET++ as well as to the Monte Carlo predictions .

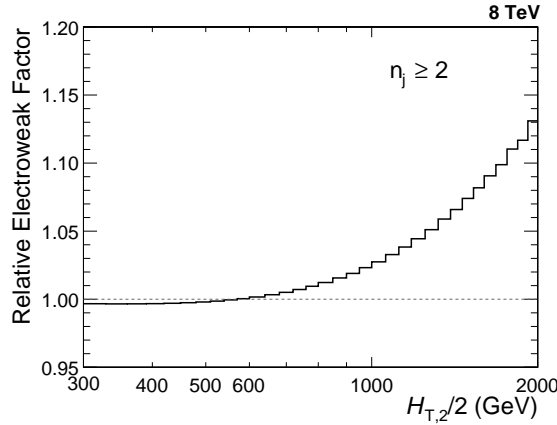


Figure 2.3: The electroweak (EW) corrections [43] in the phase space of the measurement are shown as a function of  $H_{T,2}/2$  for inclusive 2-jet events cross sections. These corrections are applied as a bin-by-bin correction factor to the fixed-order calculation of NLOJET++ as well as the MC predictions of MADGRAPH5+PYTHIA6. The EW correction factor increases up to 13% at high ends of  $H_{T,2}/2$  and significantly improves the agreement between data and prediction.

## 2.2 Theoretical Uncertainties

The measurements are not only sensitive to experimental uncertainties but also to the theoretical uncertainties. The various sources contributing to theoretical uncertainties are described below :

### 2.2.1 Scale Uncertainty

In perturbative QCD calculations of cross sections, one has to choose a renormalization ( $\mu_r$ ) and factorization ( $\mu_f$ ) scale. The dependence on scales is negligible

if these calculations are performed for all orders of the perturbative series. But the perturbative series is truncated at NLO, so there is a scale dependence of the measurement which is covered by systematic uncertainty known as scale uncertainty. The scale uncertainty is evaluated with the conventional recipe of varying the default scale  $H_{T,2}/2$  chosen for  $\mu_r$  and  $\mu_f$  independently in the following six combinations:  $(\mu_r/H_{T,2}/2, \mu_f/H_{T,2}/2) = (1/2, 1/2), (1/2, 1), (1, 1/2), (1, 2), (2, 1)$  and  $(2, 2)$ . The maximal upwards and downwards deviations in cross section from the central prediction give the scale uncertainty. To calculate the scale uncertainty for cross section ratio  $R_{32}$ , first  $R_{32}$  is obtained for each above mentioned scale choice and then its difference from central  $R_{32}$  is taken. The scale uncertainty calculated using CT10-NLO PDF set ranges for inclusive 2-jet events from 5% to 13%, for inclusive 3-jet events from 11% to 17% and for  $R_{32}$  from 6% to 8%.

### 2.2.2 PDF Uncertainty

The calculation of the jet production cross sections in proton-proton collisions relies upon the knowledge of parton distribution functions (PDF). These PDF sets are best determined by global fits to all the available deep inelastic scattering (DIS) and related hard scattering data from different experiments. The various sources affect the PDFs such as theory model, input parameters like the strong coupling constant  $\alpha_S$ , the quark masses and the statistical and systematic uncertainty sources of the data included in the PDF fit. These sources contribute to PDF uncertainty which is evaluated according to the prescriptions given for each PDF set. The CT10 NLO PDF set [29, 44] employ the eigenvector method to evaluate the PDF uncertainties. The CT10 PDF set consists of  $N_{\text{ev}} = 26$  eigenvectors with two PDF members per eigenvector  $k$ , which are varied upwards and downwards to generate a set of eigenvector pairs. The asymmetric uncertainties,  $\Delta X^+$  and  $\Delta X^-$ , of a quantity  $X$  are given by Eq. 2.4 where  $X_0$  is the central prediction,  $X_k^+$  and  $X_k^-$  are

the predictions using the upwards and downwards variation of each eigenvector  $k$ .

$$\begin{aligned}\Delta X^+ &= \sqrt{\sum_{k=1}^{N_{\text{ev}}} [\max(X_k^+ - X^0, X_k^- - X^0, 0)]^2} \\ \Delta X^- &= \sqrt{\sum_{k=1}^{N_{\text{ev}}} [\min(X_k^+ - X^0, X_k^- - X^0, 0)]^2}\end{aligned}\tag{2.4}$$

The symmetric uncertainty ( $\Delta X^\pm$ ) is given by half the difference of the upwards and downwards variations :

$$\Delta X^\pm = \sqrt{\sum_{k=1}^{N_{\text{ev}}} \left[ \frac{X_k^+ - X_k^-}{2} \right]^2}\tag{2.5}$$

The CT10 uncertainties are downscaled by a factor of 1.64 in order to have the uncertainties at the 68.3% confidence level CL( $1\sigma$ ) instead of 90% CL( $2\sigma$ ). The PDF uncertainty as derived with the CT10 PDF set is the dominant source of uncertainty and ranges from 3% to 30% for inclusive 2-jet and from 4% to 32% for 3-jet cross sections. For  $R_{32}$ , the ratio of predictions for inclusive 3-jet to that of 2-jet is taken for each eigen vector with upwards and downwards variations separately and then PDF uncertainty is calculated as done for individual cross sections. The PDF uncertainty ranges and from 2% to 10% for cross section ratio  $R_{32}$ .

### 2.2.3 Non-perturbative Uncertainty

As discussed in 2.1.2, the differences in the non-perturbative (NP) corrections calculated from different Monte Carlo event generators introduce the NP uncertainty which is of the order of 1% and 1 to 2% for inclusive 2-jet and 3-jet events cross sections respectively, and  $< 1\%$  for cross section ratio  $R_{32}$ .

### 2.2.4 Total Theoretical Uncertainty

The total systematic theoretical uncertainties are evaluated as the quadratic sum of the scale, PDF and NP uncertainties. Figure 2.4 presents the systematic theoretical uncertainties affecting the cross section measurement for inclusive 2-jet (top left) and 3-jet events (top right) and the cross section ratio  $R_{32}$  (bottom), using CT10 PDF set. The scale (red dashed line), PDF (green line) and NP (blue dashed line) uncertainties as well as total theoretical uncertainty (black dashed line) are shown. The total theoretical uncertainty is asymmetric and is dominated by PDF uncertainty. Table 2.2 quotes the values of the theoretical uncertainty from each source as well as total uncertainty affecting the measurements. The bin-wise values of uncertainties (in %) from each source as well as total uncertainty are shown in Tables A.5, A.6 and A.7 for  $n_j \geq 2$  and  $n_j \geq 3$  events cross sections and cross section ratio  $R_{32}$ , respectively. The computation of the NLO predictions with NLOJET++ is also subject to statistical fluctuations from the complex numerical integrations. For the inclusive 2-jet event cross sections this uncertainty is smaller than about a per mille, while for the inclusive 3-jet event cross section it amounts to 1-9 per mille. Hence the statistical uncertainty is not considered in the total theoretical uncertainty.

Table 2.2: Overview of all systematic theoretical uncertainties, obtained using CT10-NLO PDF set, affecting the measurement of cross sections for inclusive 2-jet (left) and inclusive 3-jet events (middle) and cross section ratio  $R_{32}$  (right).

Uncertainty Source	Inclusive 2-jet	Inclusive 3-jet	$R_{32}$
Scale	5 to 13%	11 to 17%	6 to 8%
PDF	3 to 30%	4 to 32%	2 to 10%
Non-perturbative (NP)	1%	1 to 2%	< 1%
Total	3 to 30%	5 to 34%	3 to 11%

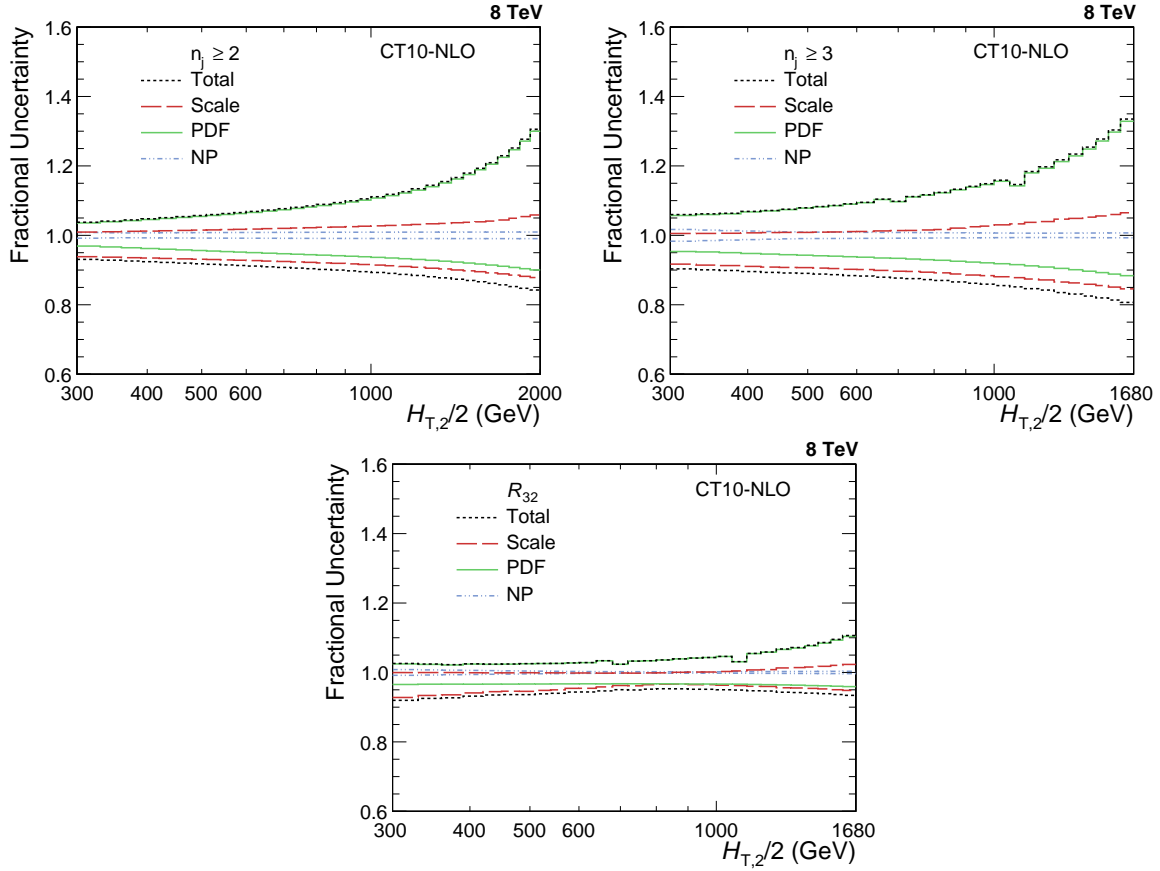


Figure 2.4: The systematic theoretical uncertainties affecting the cross section measurement for inclusive 2-jet (top left) and 3-jet events (top right) and their ratio  $R_{32}$  (bottom). The scale (red dashed line), PDF (green line) and NP (blue dashed line) uncertainties as well as total uncertainty (black dashed line) obtained using CT10-NLO PDF set are shown. The total theoretical uncertainty is asymmetric and is dominated by PDF uncertainty.

## 2.3 Comparison of theory to data

After correcting the measurement for detector effects as well as NLO pQCD calculations for non-perturbative (NP) and electroweak (EW) effects, it is now possible to compare the measured cross sections with the theory predictions. Figure 2.5 shows the measured differential inclusive 2-jet and 3-jet event cross sections as a function of  $H_{T,2}/2$  after unfolding for detector effects. On the left, the measurements (points) are compared to the NLOJET++ predictions using the CT10 PDF set (line), corrected for NP effects and in addition for EW effects in the 2-jet case. On the right, the comparison is made to the predictions from MADGRAPH5+PYTHIA6 (MG+P6)

with tune Z2<sup>\*</sup> (line), corrected for EW effects in the 2-jet case. The error bars correspond to the total experimental uncertainty, for which the statistical and systematic uncertainties are added in quadrature. On a logarithmic scale, the data are in agreement with the NLO predictions over the whole range of  $H_{T,2}/2$  from 300 GeV up to 2000 (2-jet) and 1680 GeV (3-jet) respectively.

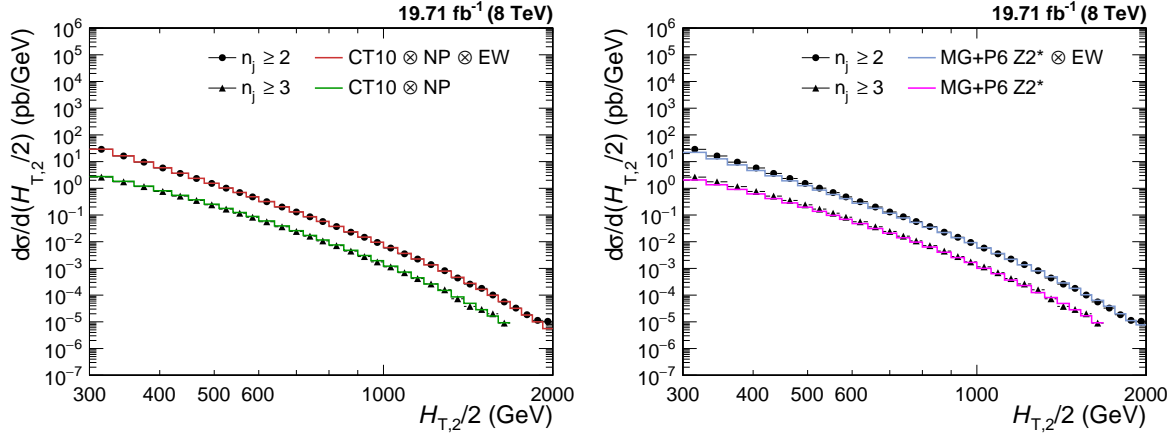


Figure 2.5: Comparison of the measured differential inclusive 2-jet and 3-jet event cross sections as a function of  $H_{T,2}/2$  to theoretical predictions. On the left, the data (points) are shown together with NLOJET++ predictions (line) using the CT10 PDF set, corrected for non-perturbative (NP) and electroweak (EW) effects (2-jet) or only NP effects (3-jet). On the (right), the data (points) are compared to predictions from MADGRAPH5+PYTHIA6 (MG+P6) with tune Z2<sup>\*</sup> (line), corrected for EW effects in the 2-jet case. The error bars correspond to the total experimental uncertainty, for which the statistical and systematic uncertainties are added in quadrature.

Figure 2.6 presents the cross section ratio  $R_{32}$  as obtained from unfolded data data (solid circles) in comparison to that from NLO pQCD predictions obtained using the CT10 NLO PDF set corrected with NP corrections (line). The error bars here correspond to the total experimental uncertainty derived as quadratic sum from all uncertainty sources. The deviations of measured  $R_{32}$  from the theoretical predicted value can be explained by the electroweak effects which are not considered yet because of their unavailability for inclusive 3-jet events cross sections.

For better visibility, the ratios of data over the theory at NLO are also studied in details. In Fig. 2.7, the ratios of data over NLOJET++ predictions using the CT10 PDF set are shown for inclusive 2-jet (top left) and 3-jet event cross sections



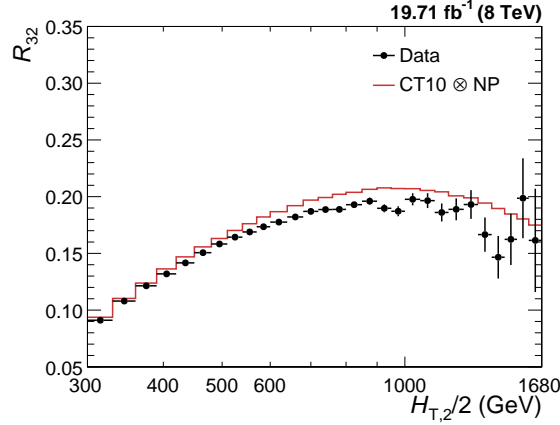


Figure 2.6: Cross section ratio  $R_{32}$  as a function of  $H_{T,2}/2$  calculated from data (solid circles) in comparison to that from NLO pQCD predictions obtained using the CT10 NLO PDF set corrected with non-perturbative (NP) corrections (line). The error bars correspond to the total experimental uncertainty derived as quadratic sum from all uncertainty sources.

(top right) as well as their ratio  $R_{32}$  (bottom). The data are well described by the predictions within their uncertainty, which is dominated at large  $H_{T,2}/2$  by PDF effects in the upwards and by scale variations in the downwards direction. A trend towards an increasing systematic excess of the 2-jet data with respect to theory, starting at about 1 TeV in  $H_{T,2}/2$ , is remedied by the inclusion of EWK corrections. In the 3-jet case the statistical precision of the data and the reach in  $H_{T,2}/2$  is insufficient to observe any effect. The alternative PDF sets MSTW2008 and NNPDF2.3 exhibit a small underestimation of the cross sections at high  $H_{T,2}/2$ .

The POWHEG framework providing a NLO dijet calculation matched to the parton showers of PYTHIA8 employed with the CUETS1 and CUETM1 tunes [41] is also used for a comparison. The ratios of data over theory from POWHEG+PYTHIA8 with tune CUETS1 are shown for inclusive 2-jet (top left) and 3-jet event cross sections (top right) as well as their ratio  $R_{32}$  (bottom) in Fig. 2.8. For comparison, the LO prediction from PYTHIA6 with tune Z2\*, the tree-level multi-leg improved prediction by MADGRAPH5+PYTHIA6 with tune Z2\*, and the matched NLO prediction from POWHEG+PYTHIA8 with tune CUETM1 are shown as well. EW corrections have been accounted for in this comparison in the 2-jet case only. Significant dis-

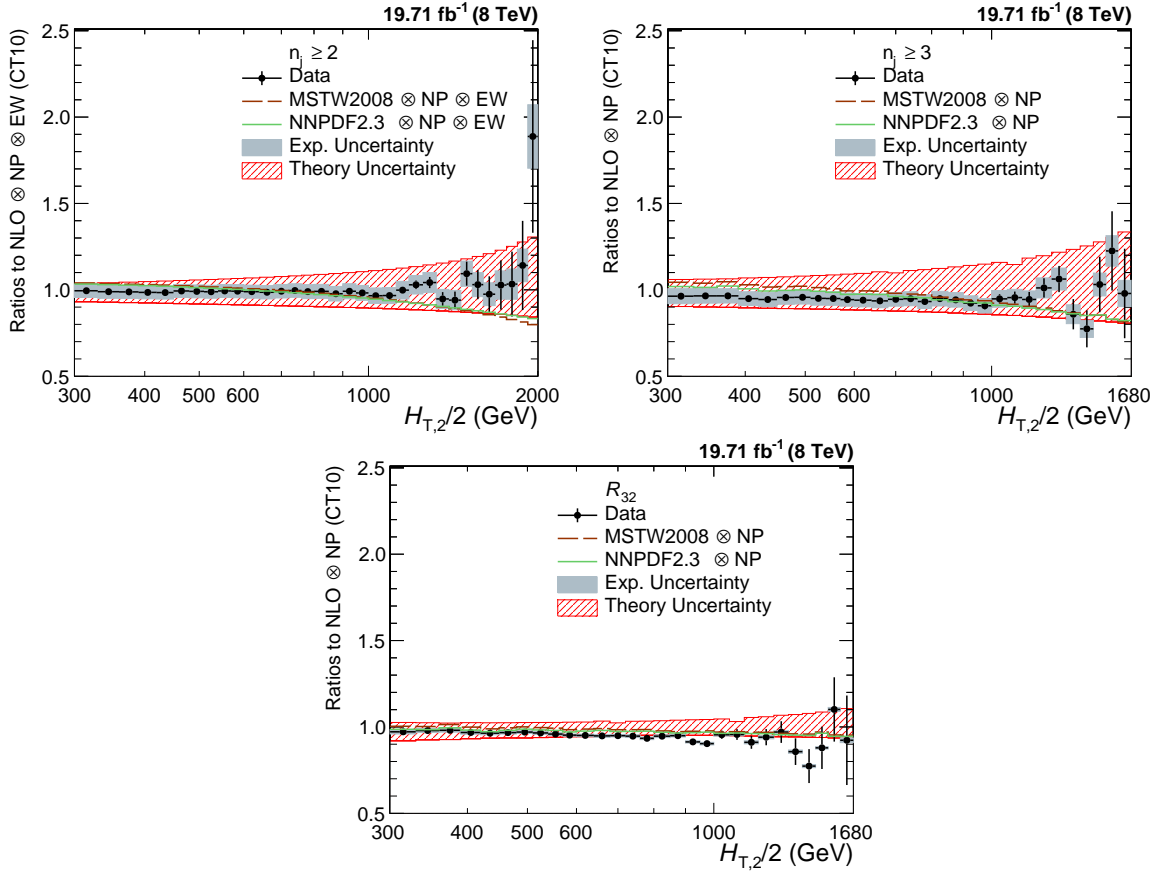


Figure 2.7: Ratio of data over theory using the CT10 PDF set for inclusive 2-jet (top left) and inclusive 3-jet event cross sections (top right) and their ratio  $R_{32}$  (bottom). For comparison predictions employing two other PDF sets, MSTW2008 and NNPDF2.3, are also shown. The error bars correspond to the statistical uncertainty of the data and the shaded rectangles to the total experimental systematic uncertainty. The shaded band around unity represents the total uncertainty of the theory.

crepancies, which are cancelled to a large extent in the ratio  $R_{32}$ , are visible in the comparison with the LO prediction from MADGRAPH5+PYTHIA6 with tune Z2\*, in particular for small  $H_{T,2}/2$ . In contrast, the employed dijet MC POWHEG+PYTHIA8 better describe the 2-jet event cross section, but fail for the 3-jet case.

The jet measurements at hadron colliders can be used to determine the strong coupling constant  $\alpha_S$ , which are discussed in the next chapter.

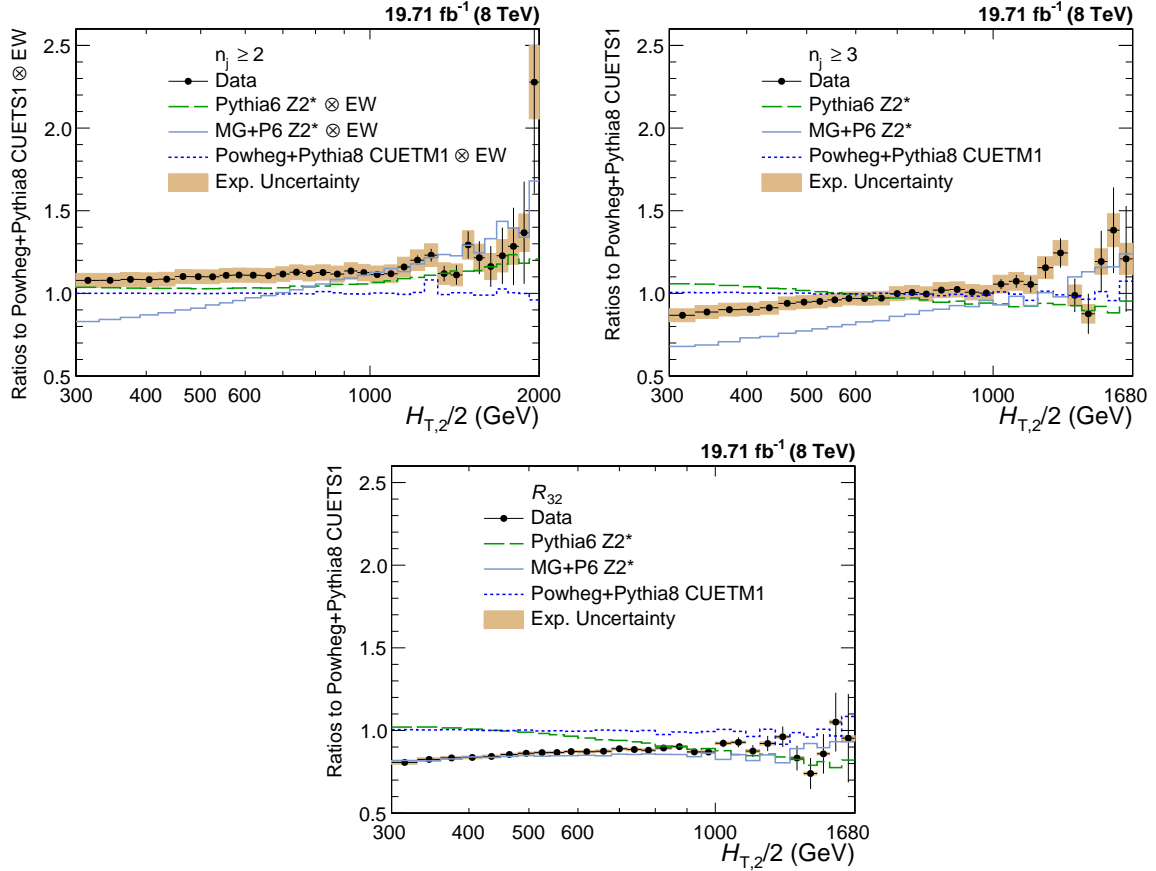


Figure 2.8: Ratio of data over the prediction from POWHEG+PYTHIA8 with tune CUETS1 are presented for inclusive 2-jet (top left) and 3-jet event cross sections (top right) as well as their ratio  $R_{32}$  (bottom). For comparison the alternative tune CUETM1 of POWHEG+PYTHIA8, the tree-level multi-leg improved prediction by MADGRAPH5+PYTHIA6 with tune Z2\*, and the the LO MC predictions from PYTHIA6 tune Z2\* are shown as well. The error bars correspond to the statistical uncertainty of the data and the shaded rectangles to the total experimental systematic uncertainty. EW corrections have been accounted for in this comparison in the 2-jet case only.

# Chapter 3

## Fits of the strong coupling constant

As discussed in the previous section, the measured inclusive 2-jet and 3-jet event cross sections and their ratio  $R_{32}$  can be used for a determination of the strong coupling constant  $\alpha_s(M_Z)$ . The value of  $\alpha_s(M_Z)$  is determined by minimizing the  $\chi^2$  between the experimental measurement and the theoretical predictions. The fit procedure here follows closely the one previously used in Refs. [45] and [46]. The  $\chi^2$  is defined as:

$$\chi^2 = M^T C^{-1} M, \quad (3.1)$$

where  $M$  is the vector of the differences between the data ( $D^i$ ) and the theoretical values ( $T^i$ ) in each bin  $i$ ,

$$M^i = D^i - T^i \quad (3.2)$$

and  $C$  is the covariance matrix including all experimental uncertainties as described in Section ?? and some theoretical uncertainties. More precisely,  $C = C_{\text{exp}} + C_{\text{theo}}$  is defined as the sum of covariances of experimental and theoretical

sources of uncertainty as follows

$$C_{\text{exp}} = \text{Cov}^{\text{ExpStat}} + \sum \text{Cov}^{\text{JEC}} + \text{Cov}^{\text{Unfolding}} + \text{Cov}^{\text{Lumi}} + \text{Cov}^{\text{Uncor}}, \quad (3.3)$$

$$C_{\text{theo}} = \text{Cov}^{\text{TheoStat}} + \text{Cov}^{\text{NP}} + \text{Cov}^{\text{PDF}}, \quad (3.4)$$

where the labelled covariance matrices account for the following effects:

- $\text{Cov}^{\text{ExpStat}}$ : the statistical uncertainty of the data including correlations introduced by the unfolding,
- $\text{Cov}^{\text{JEC}}$ : the JEC systematic uncertainty,
- $\text{Cov}^{\text{Unfolding}}$ : the unfolding systematic uncertainty including the JER,
- $\text{Cov}^{\text{Lumi}}$ : the luminosity uncertainty,
- $\text{Cov}^{\text{Uncor}}$ : a residual uncorrelated systematic uncertainty summarizing individual causes such as trigger and identification inefficiencies, time dependence of the jet  $p_{\text{T}}$  resolution, and uncertainty on the trigger prescale factors,
- $\text{Cov}^{\text{TheoStat}}$ : the statistical uncertainty caused by numerical integrations in the cross section computations,
- $\text{Cov}^{\text{NP}}$ : the systematic uncertainty of the NP corrections, and
- $\text{Cov}^{\text{PDF}}$ : the PDF uncertainty.

In fits of the ratio  $R_{32}$ , the luminosity and residual uncorrelated uncertainties cancel completely. Partial cancellations between the other sources of uncertainty are taken into account in the fit. The JEC, unfolding, and luminosity uncertainties are treated as multiplicative to avoid the statistical bias that arises when estimating uncertainties from data.

The derivation of PDF uncertainties depends on each PDF set. The CT10 PDF set consists of  $N_{\text{ev}} = 26$  eigenvectors with two PDF members per eigenvector  $k$ , which lead to the predictions  $S_k^\pm$  that follow from PDF variations with respect to the plus and minus directions of eigenvector  $k$ . Symmetric uncertainties as required by the use of covariance matrices are then computed by [44]:

$$(\Delta X)^2 = \frac{1}{4} \sum_{k=1}^{N_{\text{ev}}} [X(S_k^+) - X(S_k^-)]^2, \quad (3.5)$$

where  $\Delta X$  is the uncertainty of the cross section and  $X(S_k^\pm)$  is the predicted cross section for each eigenvector orientation,  $+$  or  $-$ .

Scale uncertainties of the pQCD predictions are taken into account employing the offset method, i.e. by performing separate fits with varying scale factors as described in the previous section. The largest upwards and downwards deviations from the default factors are defined as the uncertainty. At NLO such scale variations predominantly lead to smaller cross sections and also a smaller ratio  $R_{32}$  as visible in Fig. 2.4. As a consequence the scale uncertainty in fits is equally asymmetric, where smaller cross sections or ratios are compensated by an increase in the fitted value for  $\alpha_s(M_Z)$ .

First, fits to the cross sections are performed, where the range in  $H_{T,2}/2$  is restricted to be between 300 GeV and 1 TeV to avoid the region close to the minimal  $p_T$  threshold of 150 GeV for each jet at low  $p_T$  and the onset of electroweak effects at high  $p_T$ , which are available for the dijet case only. The results are reported in Table 3.1 for the 2-jet and 3-jet event cross sections. For comparison, a simultaneous fit to both cross sections ignoring any correlations, and a fit to their ratio fully accounting for correlations are given in Table 3.2. Also, EWK effects are assumed to cancel in the ratio as do the luminosity and the uncorrelated uncertainty.

All cross section fits give compatible values for  $\alpha_s(M_Z)$  in the range of 0.115–

0.118; for the ratio  $R_{32}$  somewhat smaller values are obtained. A common issue, except for the ratio fits, is the rather small  $\chi^2/n_{\text{dof}}$ . A possible explanation is an overestimation of the residual uncorrelated uncertainty of 1% that is cancelled for  $R_{32}$ . If the fits are repeated with an assumed uncertainty of 0.25% instead, the  $\chi^2/n_{\text{dof}}$  values lie around unity while the  $\alpha_s(M_Z)$  values are still compatible with the previous results but with slightly reduced uncertainties.

Table 3.1: Determination of  $\alpha_s(M_Z)$  from the inclusive 2-jet and 3-jet event cross sections using five PDF sets at NLO. Only total uncertainties without scale variations are quoted. The results are obtained from a simultaneous fit to all 19  $H_{T,2}/2$  bins in the restricted range of  $0.3 < H_{T,2}/2 < 1.0$  TeV.

PDF set	2-jets			3-jets		
	$\alpha_s(M_Z)$	$\pm\Delta\alpha_s(M_Z)$	$\chi^2/n_{\text{dof}}$	$\alpha_s(M_Z)$	$\pm\Delta\alpha_s(M_Z)$	$\chi^2/n_{\text{dof}}$
CT10	0.1174	0.0032	3.0/18	0.1169	0.0027	5.4/18
CT14	0.1160	0.0035	3.5/18	0.1159	0.0031	6.1/18
MSTW2008	0.1159	0.0025	5.3/18	0.1161	0.0021	6.7/18
MMHT2014	0.1165	0.0034	5.9/18	0.1166	0.0025	7.1/18
NNPDF2.3	0.1183	0.0025	9.7/18	0.1179	0.0021	9.1/18

Table 3.2: Determination of  $\alpha_s(M_Z)$  from the inclusive 2-jet and 3-jet event cross sections simultaneously and from their ratio  $R_{32}$  using five PDF sets at NLO. Only total uncertainties without scale variations are quoted. The results are obtained from a simultaneous fit to all 38 (19)  $H_{T,2}/2$  bins in the restricted range of  $0.3 < H_{T,2}/2 < 1.0$  TeV. For comparison, correlations between the two cross sections are neglected in the simultaneous fit on the left, but fully taken into account in the ratio fit on the right.

PDF set	2- & 3-jets			$R_{32}$		
	$\alpha_s(M_Z)$	$\pm\Delta\alpha_s(M_Z)$	$\chi^2/n_{\text{dof}}$	$\alpha_s(M_Z)$	$\pm\Delta\alpha_s(M_Z)$	$\chi^2/n_{\text{dof}}$
CT10	0.1170	0.0026	8.2/37	0.1141	0.0028	19./18
CT14	0.1161	0.0029	9.1/37	0.1139	0.0032	15./18
MSTW2008	0.1161	0.0021	11./37	0.1150	0.0023	21./18
MMHT2014	0.1168	0.0025	11./37	0.1142	0.0022	19./18
NNPDF2.3	0.1188	0.0019	15./37	0.1184	0.0021	12./18

To investigate how the EWK corrections affect the fit results for  $\alpha_s(M_Z)$ , the range in  $H_{T,2}/2$  is extended to  $0.3 < H_{T,2}/2 < 1.68$  TeV. Table 3.3 reports the values obtained for  $\alpha_s(M_Z)$  from fits to the 2-jet event cross section in this range with or without EWK correction factors. The largest impact is a reduction

in  $\chi^2/n_{\text{dof}}$ , which indicates a better agreement when EWK effects are included. In addition, a tendency to slightly smaller  $\alpha_s(M_Z)$  values is observed without the EWK corrections. For the ratio  $R_{32}$  it is expected that these effects are much reduced.

Table 3.3: Determination of  $\alpha_s(M_Z)$  from the inclusive 2-jet event cross section using five PDF sets at NLO with (right) and without (left) EWK corrections. Only total uncertainties without scale variations are quoted. The results are obtained from a simultaneous fit to all 29  $H_{T,2}/2$  bins in the range of  $0.3 < H_{T,2}/2 < 1.68$  TeV.

PDF set	2-jets, without EWK			2-jets, with EWK		
	$\alpha_s(M_Z)$	$\pm\Delta\alpha_s(M_Z)$	$\chi^2/n_{\text{dof}}$	$\alpha_s(M_Z)$	$\pm\Delta\alpha_s(M_Z)$	$\chi^2/n_{\text{dof}}$
CT10	0.1163	0.0034	15./28	0.1165	0.0032	14./28
CT14	0.1137	0.0033	24./28	0.1144	0.0033	17./28
MSTW2008	0.1093	0.0028	27./28	0.1133	0.0023	19./28
MMHT2014	0.1127	0.0032	32./28	0.1141	0.0032	21./28
NNPDF2.3	0.1162	0.0024	31./28	0.1168	0.0024	23./28

From Fig. ?? follows that only the PDF sets MSTW2008 and MMHT2014 provide a large enough range in  $\alpha_s(M_Z)$  values to ensure fits without extrapolation. The other three PDF sets are at the limit such that reliable fits cannot be performed for all scale settings and/or bins in scale  $Q = H_{T,2}/2$ . Tables 3.4–3.6 give the complete results for MSTW2008 and MMHT2014 for the full range in  $H_{T,2}/2$  of 300 GeV up to 1.68 TeV, for scale variations in this range, and for subranges in  $H_{T,2}/2$ .

Using the MSTW2008 PDF set, which dates from before the LHC start, the strong coupling constant finally is determined to

$$\begin{aligned}
 \alpha_s(M_Z) &= 0.1150 \pm 0.0010 (\text{exp}) \pm 0.0013 (\text{PDF}) \pm 0.0015 (\text{NP})^{+0.0050}_{-0.0000} (\text{scale}) \quad (3.6) \\
 &= 0.1150 \pm 0.0023 (\text{all except scale})^{+0.0050}_{-0.0000} (\text{scale}) .
 \end{aligned}$$

The MMHT2014 PDF set, although using LHC jet data to determine the PDF



parameters, leads to a very similar result of

$$\begin{aligned}\alpha_s(M_Z) &= 0.1142 \pm 0.0010 (\text{exp}) \pm 0.0013 (\text{PDF}) \pm 0.0014 (\text{NP}) \stackrel{+0.0049}{-0.0006} (\text{scale}) \quad (3.7) \\ &= 0.1142 \pm 0.0022 (\text{all except scale}) \stackrel{+0.0049}{-0.0006} (\text{scale}) .\end{aligned}$$

In contrast to fits at NLO using cross sections, where the scale uncertainty recipe usually leads to a very asymmetric behaviour with the larger uncertainty towards smaller values of  $\alpha_s(M_Z)$ , this is inverted for the fits to the cross section ratio.

Table 3.7 provides in addition to the extracted  $\alpha_s(M_Z)$  value for each range in  $H_{T,2}/2$  the  $\alpha_s(Q)$  values with total uncertainty as evolved to the respective cross-section averaged scale  $\langle Q \rangle$  in that range. The evolution is performed for five flavours at 2-loop order with the RUNDEC program [47, 48]. The obtained  $\alpha_s(Q)$  points are illustrated in Fig. 3.1 together with the world average [49] and results from other measurements of the CMS [17, 26, 45, 46, 50], ATLAS [51], D0 [52, 53], H1 [54, 55], and ZEUS [56] experiments.

Table 3.4: Determination of  $\alpha_s(M_Z)$  from the ratio  $R_{32}$  using the two most compatible PDF sets MSTW2008 and MMHT2014 at NLO. The results are obtained from a simultaneous fit to all 29  $H_{T,2}/2$  bins in the full range of  $0.3 < H_{T,2}/2 < 1.68$  TeV.

PDF set	$R_{32}: \Delta\alpha_s(M_Z) \times 1000$						
	$\alpha_s(M_Z)$	exp	PDF	NP	all exc.	scale	$\chi^2/n_{\text{dof}}$
MSTW2008	0.1150	$\pm 10$	$\pm 13$	$\pm 15$	$\pm 23$	$^{+50}_{-0}$	26./28
MMHT2014	0.1142	$\pm 10$	$\pm 13$	$\pm 14$	$\pm 22$	$^{+49}_{-6}$	24./28

Table 3.5: Fitted values of  $\alpha_s(M_Z)$  using  $R_{32}$  in the  $H_{T,2}/2$  range from 0.3 up to 1.68 TeV at the central scale and for the six scale factor combinations for the two PDF sets MSTW2008 and MMHT2014.

$\mu_r/H_{T,2}/2$	$\mu_f/H_{T,2}/2$	MSTW2008		MMHT2014	
		$\alpha_s(M_Z)$	$\chi^2/n_{\text{dof}}$	$\alpha_s(M_Z)$	$\chi^2/n_{\text{dof}}$
1	1	0.1150	26./28	0.1142	24./28
1/2	1/2	0.1165	77./28	0.1160	73./28
2	2	0.1200	18./28	0.1191	18./28
1/2	1	0.1150	53./28	0.1136	48./28
1	1/2	0.1150	30./28	0.1142	28./28
1	2	0.1155	23./28	0.1147	22./28
2	1	0.1180	19./28	0.1175	19./28

Table 3.6: Uncertainty composition for  $\alpha_s(M_Z)$  from the determination of  $\alpha_s$  from the jet event rate  $R_{32}$  in bins of  $H_{T,2}/2$ . The statistical uncertainty of the NLO computation is negligible in comparison to any of the other sources of uncertainty. Electroweak corrections, significant only at high  $H_{T,2}/2$ , are assumed to cancel between the numerator and denominator.

$H_{T,2}/2$ (GeV)	MSTW2008: $\Delta\alpha_s(M_Z) \times 1000$					MMHT2014: $\Delta\alpha_s(M_Z) \times 1000$				
	$\alpha_s(M_Z)$	exp	PDF	NP	scale	$\alpha_s(M_Z)$	exp	PDF	NP	scale
300–420	0.1157	$\pm 15$	$\pm 14$	$\pm 19$	$^{+53}_{-0}$	0.1158	$\pm 14$	$\pm 10$	$\pm 19$	$^{+52}_{-0}$
420–600	0.1153	$\pm 11$	$\pm 14$	$\pm 18$	$^{+57}_{-0}$	0.1154	$\pm 11$	$\pm 12$	$\pm 17$	$^{+56}_{-0}$
600–1000	0.1134	$\pm 13$	$\pm 16$	$\pm 19$	$^{+52}_{-0}$	0.1140	$\pm 12$	$\pm 12$	$\pm 18$	$^{+45}_{-0}$
1000–1680	0.1147	$\pm 29$	$\pm 17$	$\pm 18$	$^{+63}_{-11}$	0.1154	$\pm 25$	$\pm 14$	$\pm 15$	$^{+56}_{-11}$
300–1680	0.1150	$\pm 10$	$\pm 13$	$\pm 15$	$^{+50}_{-0}$	0.1142	$\pm 10$	$\pm 13$	$\pm 14$	$^{+49}_{-6}$

Table 3.7: Evolution of the strong coupling constant between the scale of the Z boson mass and the cross-section averaged  $H_{T,2}/2$  scale  $\langle Q \rangle$  for the separate determinations in each respective fit range. The evolution is performed for five flavours at 2-loop order with the RUNDEC program [47, 48].

$H_{T,2}/2$ (GeV)	$\langle Q \rangle$ (GeV)	$\alpha_s(M_Z)$	$\alpha_s(Q)$	No. of data points	$\chi^2/n_{\text{dof}}$
300–420	340	$0.1157^{+0.0060}_{-0.0030}$	$0.0969^{+0.0041}_{-0.0021}$	4	2.8/3
420–600	476	$0.1153^{+0.0062}_{-0.0025}$	$0.0928^{+0.0039}_{-0.0016}$	6	6.1/5
600–1000	685	$0.1134^{+0.0059}_{-0.0028}$	$0.0879^{+0.0035}_{-0.0017}$	9	7.1/8
1000–1680	1114	$0.1147^{+0.0074}_{-0.0040}$	$0.0841^{+0.0039}_{-0.0021}$	10	5.4/9

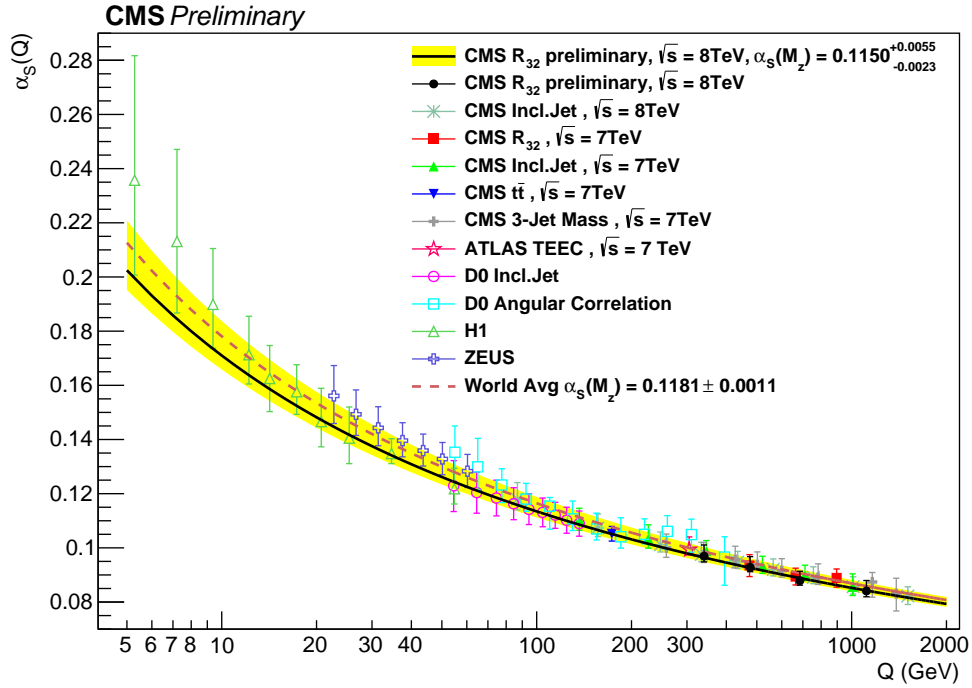


Figure 3.1: The running  $\alpha_s(Q)$  as a function of the scale  $Q$  is shown as obtained by using the MSTW2008 NLO PDF set. The solid line and the uncertainty band are drawn by evolving the extracted  $\alpha_s(M_Z)$  values using the 2-loop 5-flavour renormalization group equations as implemented in RUNDEC [47, 48]. The dashed line represents the evolution of the world average [49] and the black circles correspond to the  $\alpha_s(Q)$  determinations presented in Table 3.7. Results from other measurements of CMS [17, 26, 45, 46, 50], ATLAS [51], D0 [52, 53], H1 [54, 55], and ZEUS [56] are superimposed.



# Appendix A

## A.1 Cross Section Ratio, $R_{32}$

Table A.1: Differential cross sections and cross section ratio at detector level in each bin of  $H_{T,2}/2$ , along with statistical uncertainty (in %).

Bin	2-jet x-section (x $10^{-3}$ (pb/GeV))	Stat. unc.	3-jet x-section (x $10^{-3}$ (pb/GeV))	Stat. unc.	Ratio $R_{32}$	Stat. unc.
300 - 330	29772.726	0.211	2640.629	0.707	0.089	+0.665 -0.661
330 - 360	16792.917	0.231	1773.485	0.704	0.106	+0.523 -0.521
360 - 390	9889.326	0.182	1176.544	0.526	0.119	+0.485 -0.483
390 - 420	5976.777	0.179	778.034	0.492	0.130	+0.206 -0.206
420 - 450	3731.760	0.067	522.624	0.180	0.140	+0.167 -0.167
450 - 480	2398.741	0.084	357.622	0.217	0.149	+0.201 -0.200
480 - 510	1570.192	0.104	246.051	0.262	0.157	+0.241 -0.241
510 - 540	1048.665	0.127	171.080	0.314	0.163	+0.288 -0.287
540 - 570	713.042	0.154	119.566	0.376	0.168	+0.344 -0.343
570 - 600	490.776	0.186	84.798	0.447	0.173	+0.407 -0.406
600 - 640	325.046	0.198	57.463	0.470	0.177	+0.427 -0.426
640 - 680	205.727	0.248	37.282	0.583	0.181	+0.529 -0.527
680 - 720	133.674	0.308	24.859	0.714	0.186	+0.646 -0.643
720 - 760	87.911	0.380	16.560	0.875	0.188	+0.791 -0.786
760 - 800	58.657	0.465	11.056	1.071	0.188	+0.968 -0.961
800 - 850	38.106	0.516	7.318	1.178	0.192	+1.063 -1.054
850 - 900	23.587	0.656	4.600	1.485	0.195	+1.339 -1.326
900 - 950	15.130	0.819	2.896	1.872	0.191	+1.694 -1.672
950 - 1000	9.696	1.023	1.812	2.366	0.187	+2.151 -2.116
1000 - 1060	6.026	1.185	1.186	2.670	0.197	+2.414 -2.371
1060 - 1120	3.668	1.518	0.716	3.436	0.195	+3.118 -3.046
1120 - 1180	2.327	1.906	0.437	4.398	0.188	+4.024 -3.903
1180 - 1250	1.419	2.260	0.265	5.227	0.187	+4.798 -4.627
1250 - 1320	0.853	2.915	0.165	6.623	0.194	+6.080 -5.811
1320 - 1390	0.477	3.898	0.080	9.492	0.169	+8.951 -8.355
1390 - 1460	0.263	5.249	0.042	13.131	0.160	+12.619 -11.449
1460 - 1530	0.192	6.143	0.029	15.811	0.151	+15.437 -13.698
1530 - 1600	0.104	8.362	0.021	18.570	0.203	+17.571 -15.536
1600 - 1680	0.060	10.314	0.009	26.726	0.149	+27.132 -22.170

## A.2 Individual Sources of Jet Energy Correction Uncertainties

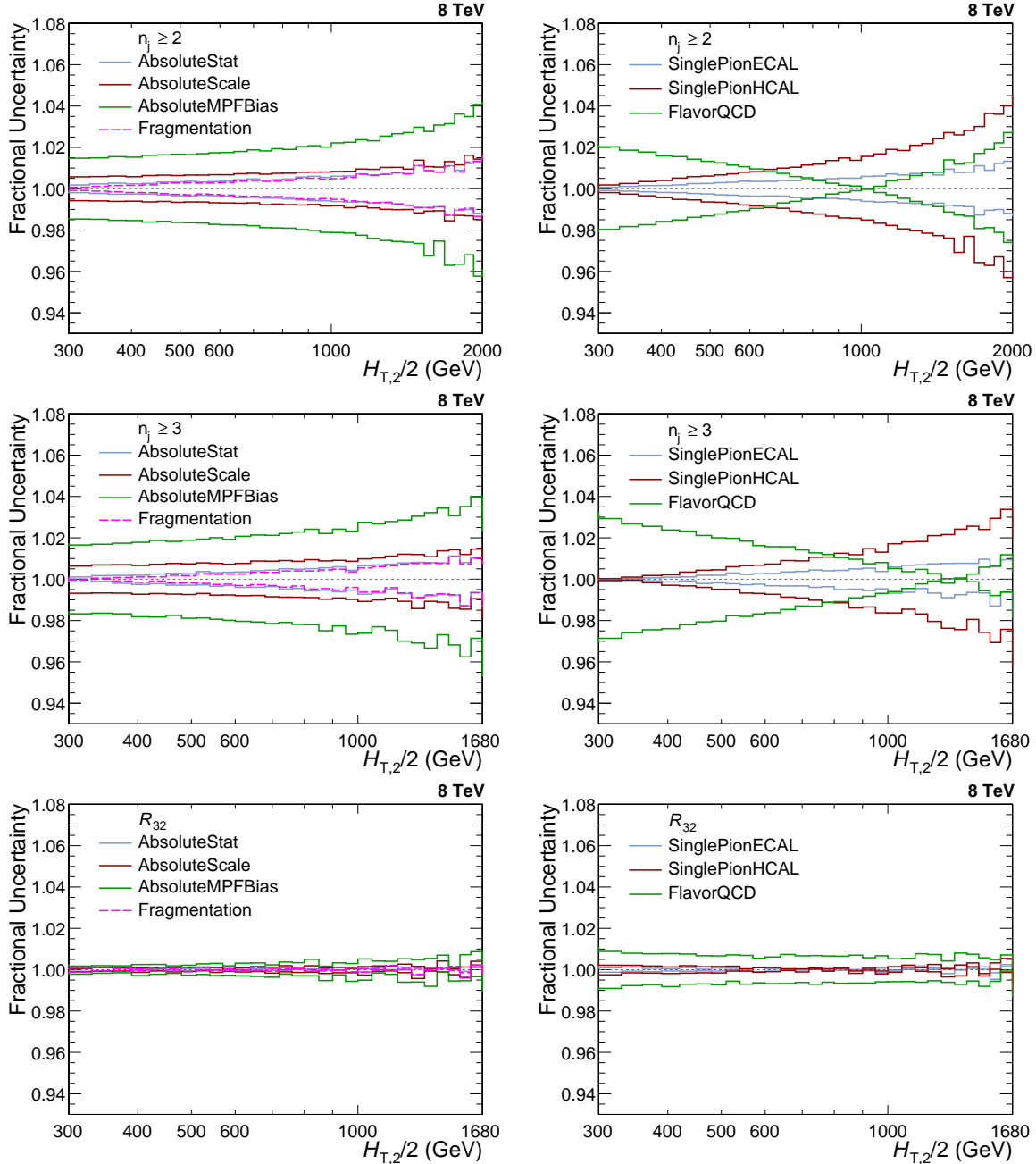


Figure A.1: The fractional jet energy correction (JEC) uncertainties from individual sources are shown for inclusive 2-jet (top) and 3-jet events cross sections (middle); and cross section ratio  $R_{32}$  (bottom). On left, JEC uncertainties are evaluated from AbsoluteStat (blue), AbsoluteScale (red), AbsoluteMPFBias (green) and Fragmentation (pink) sources whereas on right, these are evaluated from SinglePionECAL (blue), SinglePionHCAL (red) and FlavorQCD (green) sources.

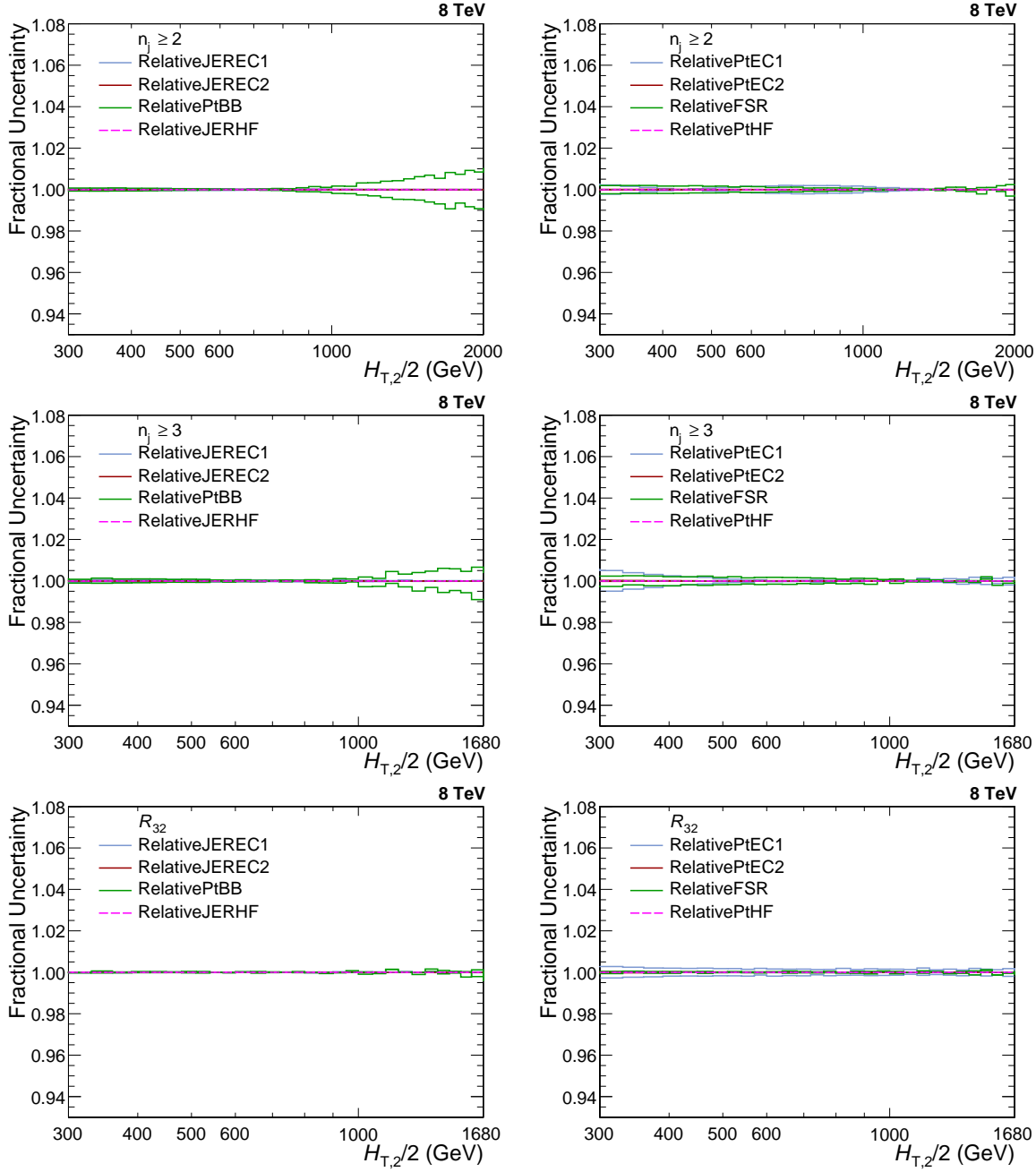


Figure A.2: The fractional jet energy correction (JEC) uncertainties from individual sources are shown for inclusive 2-jet (top) and 3-jet events cross sections (middle); and cross section ratio  $R_{32}$  (bottom). On left, JEC uncertainties are evaluated from RelativeJEREC1 (blue), RelativeJEREC2 (red), RelativePtBB (green) and RelativeJERHF (pink) sources whereas on right, these are evaluated from RelativePtEC1 (blue), RelativePtEC2 (red), RelativePFSR (green) and RelativePtHF (pink) sources.



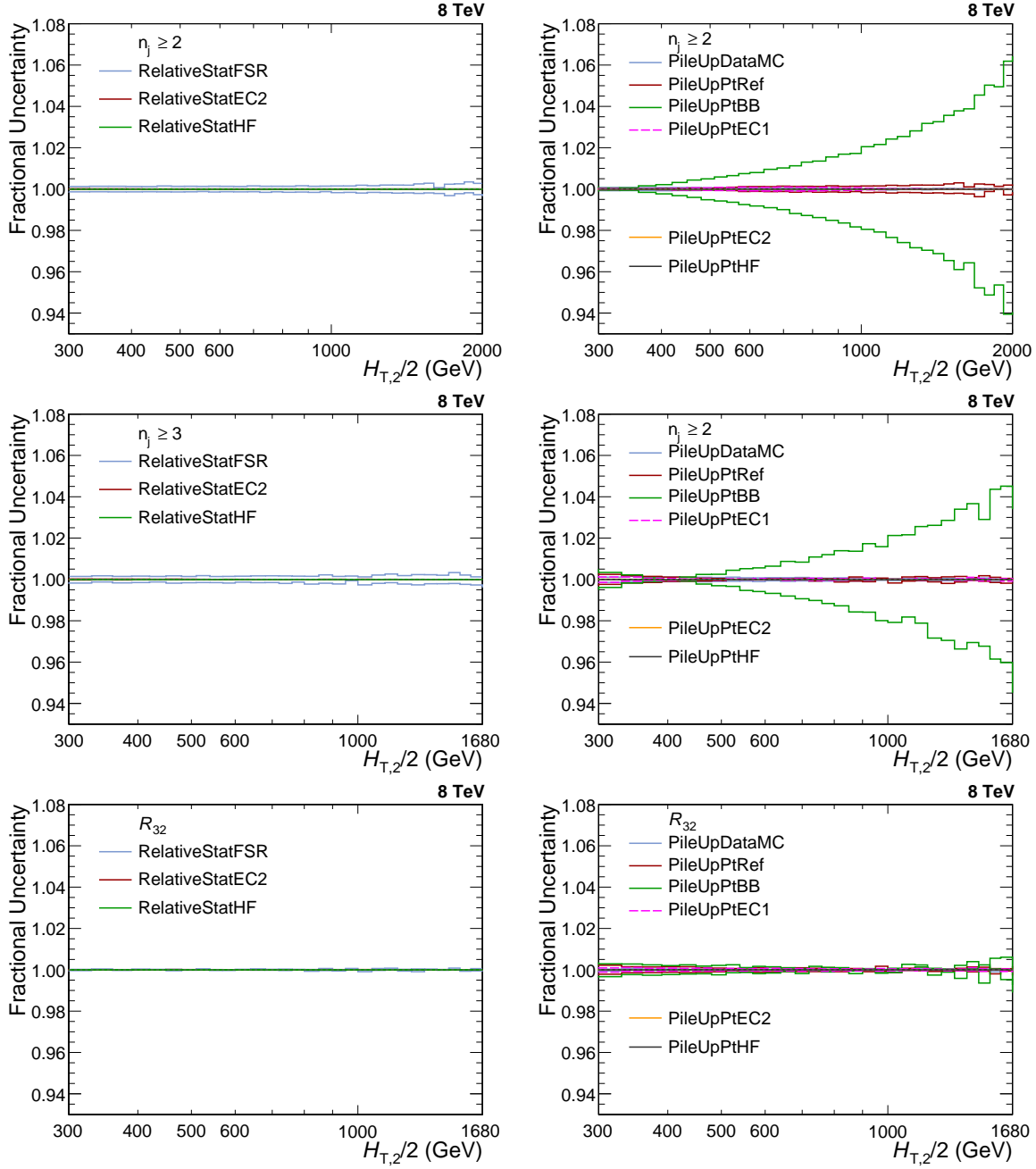


Figure A.3: The fractional jet energy correction (JEC) uncertainties from individual sources are shown for inclusive 2-jet (top) and 3-jet events cross sections (middle); and cross section ratio  $R_{32}$  (bottom). On left, JEC uncertainties are evaluated from RelativeStatFSR (blue), RelativeStatEC2 (red) and RelativeStatHF (green) sources whereas on right, these are evaluated from PileUpDataMC (blue), PileUpPtRef (red), PileUpPtBB (green), PileUpPtEC1 (pink), PileUpPtEC2 (orange) and PileUpPtHF (black) sources.

### A.3 Experimental Uncertainties

Table A.2: Experimental uncertainties (in %), from all sources as well as the total uncertainty, affecting the cross section measurement in each bin of  $H_{T,2}/2$  for inclusive 2-jet events.

Bin	Statistical	JEC	Unfolding	Lumi	Residual	Total
300 - 330	0.242	+2.612 -2.565	+0.948 -0.928	2.6	1.0	+3.942 -3.906
330 - 360	0.258	+2.507 -2.473	+0.976 -0.969	2.6	1.0	+3.882 -3.858
360 - 390	0.202	+2.504 -2.465	+0.779 -0.783	2.6	1.0	+3.831 -3.807
390 - 420	0.193	+2.363 -2.381	+0.905 -0.904	2.6	1.0	+3.768 -3.780
420 - 450	0.084	+2.448 -2.422	+0.904 -0.895	2.6	1.0	+3.818 -3.799
450 - 480	0.096	+2.440 -2.352	+0.797 -0.795	2.6	1.0	+3.789 -3.733
480 - 510	0.107	+2.427 -2.406	+0.728 -0.715	2.6	1.0	+3.767 -3.751
510 - 540	0.128	+2.425 -2.395	+0.835 -0.862	2.6	1.0	+3.789 -3.775
540 - 570	0.154	+2.425 -2.376	+0.687 -0.674	2.6	1.0	+3.760 -3.726
570 - 600	0.180	+2.497 -2.474	+0.839 -0.827	2.6	1.0	+3.838 -3.820
600 - 640	0.209	+2.495 -2.491	+0.744 -0.743	2.6	1.0	+3.819 -3.816
640 - 680	0.264	+2.582 -2.545	+0.912 -0.912	2.6	1.0	+3.915 -3.891
680 - 720	0.320	+2.691 -2.574	+0.763 -0.756	2.6	1.0	+3.961 -3.880
720 - 760	0.387	+2.690 -2.755	+0.705 -0.712	2.6	1.0	+3.955 -4.001
760 - 800	0.465	+2.858 -2.846	+0.859 -0.846	2.6	1.0	+4.109 -4.098
800 - 850	0.548	+2.889 -2.913	+0.783 -0.787	2.6	1.0	+4.126 -4.143
850 - 900	0.698	+3.145 -3.102	+0.961 -0.958	2.6	1.0	+4.366 -4.334
900 - 950	0.847	+3.298 -3.233	+0.828 -0.829	2.6	1.0	+4.476 -4.429
950 - 1000	1.041	+3.291 -3.330	+0.895 -0.872	2.6	1.0	+4.525 -4.549
1000 - 1060	1.268	+3.598 -3.569	+0.945 -0.956	2.6	1.0	+4.817 -4.798
1060 - 1120	1.611	+3.759 -3.756	+0.970 -0.967	2.6	1.0	+5.043 -5.040
1120 - 1180	1.985	+4.154 -4.053	+1.089 -1.080	2.6	1.0	+5.490 -5.413
1180 - 1250	2.406	+4.251 -4.313	+1.062 -1.070	2.6	1.0	+5.722 -5.770
1250 - 1320	3.101	+4.696 -4.624	+1.151 -1.144	2.6	1.0	+6.384 -6.330
1320 - 1390	4.157	+4.934 -4.979	+1.343 -1.341	2.6	1.0	+7.155 -7.186
1390 - 1460	5.270	+5.148 -5.104	+1.185 -1.177	2.6	1.0	+7.965 -7.936
1460 - 1530	6.360	+5.890 -5.652	+1.405 -1.406	2.6	1.0	+9.213 -9.063
1530 - 1600	8.183	+5.924 -6.311	+1.598 -1.590	2.6	1.0	+10.601 -10.821
1600 - 1680	10.630	+5.969 -5.655	+1.607 -1.592	2.6	1.0	+12.608 -12.461
1680 - 1760	13.864	+7.245 -7.603	+1.821 -1.839	2.6	1.0	+15.993 -16.161
1760 - 1840	18.192	+7.781 -7.820	+1.902 -1.906	2.6	1.0	+20.071 -20.087
1840 - 1920	22.612	+7.647 -7.537	+1.588 -1.590	2.6	1.0	+24.085 -24.050
1920 - 2000	29.530	+9.199 -9.469	+1.511 -1.505	2.6	1.0	+31.092 -31.172

Table A.3: Experimental uncertainties (in %), from all sources as well as the total uncertainty, affecting the cross section measurement in each bin of  $H_{T,2}/2$  for inclusive 3-jet events.

Bin	Statistical	JEC	Unfolding	Lumi	Residual	Total
300 - 330	0.796	+3.503 -3.475	+0.564 -0.552	2.6	1.0	+4.581 -4.558
330 - 360	0.781	+3.303 -3.186	+0.640 -0.633	2.6	1.0	+4.437 -4.350
360 - 390	0.583	+3.221 -3.094	+0.490 -0.496	2.6	1.0	+4.326 -4.233
390 - 420	0.531	+3.092 -3.149	+0.584 -0.584	2.6	1.0	+4.236 -4.278
420 - 450	0.224	+3.125 -2.996	+0.604 -0.592	2.6	1.0	+4.236 -4.140
450 - 480	0.248	+2.984 -2.890	+0.531 -0.528	2.6	1.0	+4.124 -4.056
480 - 510	0.269	+2.937 -2.963	+0.511 -0.512	2.6	1.0	+4.089 -4.108
510 - 540	0.318	+3.021 -2.797	+0.592 -0.612	2.6	1.0	+4.164 -4.007
540 - 570	0.375	+2.999 -2.935	+0.506 -0.500	2.6	1.0	+4.141 -4.094
570 - 600	0.434	+2.824 -2.906	+0.646 -0.620	2.6	1.0	+4.042 -4.096
600 - 640	0.497	+2.952 -2.956	+0.598 -0.604	2.6	1.0	+4.133 -4.136
640 - 680	0.617	+3.111 -3.001	+0.777 -0.786	2.6	1.0	+4.292 -4.215
680 - 720	0.739	+3.067 -2.984	+0.642 -0.611	2.6	1.0	+4.257 -4.194
720 - 760	0.895	+3.185 -3.111	+0.595 -0.607	2.6	1.0	+4.366 -4.313
760 - 800	1.068	+3.231 -3.166	+0.763 -0.774	2.6	1.0	+4.464 -4.419
800 - 850	1.250	+3.427 -3.295	+0.674 -0.687	2.6	1.0	+4.639 -4.544
850 - 900	1.578	+3.364 -3.540	+0.903 -0.898	2.6	1.0	+4.731 -4.857
900 - 950	1.961	+3.594 -3.524	+0.792 -0.793	2.6	1.0	+5.015 -4.965
950 - 1000	2.420	+3.603 -3.783	+0.846 -0.843	2.6	1.0	+5.226 -5.351
1000 - 1060	2.844	+4.164 -4.116	+0.916 -0.940	2.6	1.0	+5.834 -5.803
1060 - 1120	3.647	+4.038 -3.815	+0.963 -0.957	2.6	1.0	+6.188 -6.044
1120 - 1180	4.607	+4.278 -4.183	+1.084 -1.087	2.6	1.0	+6.961 -6.904
1180 - 1250	5.532	+4.894 -4.771	+1.074 -1.069	2.6	1.0	+7.967 -7.891
1250 - 1320	7.141	+5.144 -5.273	+1.222 -1.217	2.6	1.0	+9.312 -9.383
1320 - 1390	10.207	+5.542 -5.642	+1.414 -1.428	2.6	1.0	+12.027 -12.076
1390 - 1460	13.831	+5.630 -5.265	+1.257 -1.256	2.6	1.0	+15.242 -15.111
1460 - 1530	15.578	+5.576 -5.491	+1.546 -1.551	2.6	1.0	+16.850 -16.822
1530 - 1600	18.729	+6.409 -7.019	+1.718 -1.716	2.6	1.0	+20.063 -20.266
1600 - 1680	26.465	+7.017 -6.255	+1.775 -1.765	2.6	1.0	+27.578 -27.393

Table A.4: Experimental uncertainties (in %), from all sources as well as the total uncertainty, affecting the measurement of cross section ratio  $R_{32}$ , in each bin of  $H_{T,2}/2$ .

Bin	Statistical	JEC	Unfolding	Total
300 - 330	0.741	+1.059 -1.097	+0.754 -0.751	+1.496 -1.522
330 - 360	0.587	+0.954 -0.923	+0.685 -0.689	+1.313 -1.292
360 - 390	0.519	+0.902 -0.855	+0.594 -0.593	+1.199 -1.163
390 - 420	0.236	+0.907 -0.952	+0.439 -0.438	+1.035 -1.074
420 - 450	0.192	+0.900 -0.835	+0.360 -0.361	+0.988 -0.930
450 - 480	0.209	+0.788 -0.802	+0.307 -0.308	+0.872 -0.884
480 - 510	0.245	+0.795 -0.867	+0.254 -0.235	+0.870 -0.931
510 - 540	0.287	+0.852 -0.682	+0.264 -0.268	+0.937 -0.787
540 - 570	0.326	+0.807 -0.803	+0.193 -0.189	+0.891 -0.887
570 - 600	0.397	+0.656 -0.774	+0.199 -0.219	+0.792 -0.898
600 - 640	0.447	+0.763 -0.797	+0.150 -0.154	+0.897 -0.926
640 - 680	0.573	+0.861 -0.781	+0.153 -0.140	+1.045 -0.979
680 - 720	0.663	+0.766 -0.787	+0.147 -0.164	+1.024 -1.042
720 - 760	0.774	+0.842 -0.769	+0.118 -0.118	+1.149 -1.097
760 - 800	0.970	+0.800 -0.729	+0.115 -0.096	+1.263 -1.218
800 - 850	1.116	+0.873 -0.775	+0.115 -0.104	+1.422 -1.363
850 - 900	1.436	+0.770 -0.896	+0.069 -0.069	+1.631 -1.694
900 - 950	1.716	+0.704 -0.752	+0.050 -0.051	+1.855 -1.874
950 - 1000	2.156	+0.824 -0.897	+0.089 -0.045	+2.310 -2.336
1000 - 1060	2.554	+0.812 -0.870	+0.045 -0.040	+2.680 -2.698
1060 - 1120	3.244	+0.792 -0.658	+0.018 -0.027	+3.339 -3.310
1120 - 1180	4.121	+0.985 -0.757	+0.025 -0.043	+4.237 -4.191
1180 - 1250	4.990	+1.031 -0.848	+0.023 -0.041	+5.095 -5.062
1250 - 1320	6.456	+0.750 -1.087	+0.079 -0.079	+6.500 -6.548
1320 - 1390	8.990	+1.112 -1.144	+0.080 -0.099	+9.059 -9.063
1390 - 1460	12.699	+1.157 -0.815	+0.076 -0.078	+12.751 -12.725
1460 - 1530	13.926	+0.768 -1.235	+0.143 -0.145	+13.948 -13.981
1530 - 1600	16.903	+1.050 -1.258	+0.120 -0.127	+16.936 -16.950
1600 - 1680	28.070	+1.471 -0.859	+0.178 -0.177	+28.109 -28.084

## A.4 Theoretical Uncertainties

Table A.5: Theoretical uncertainties (in %) calculating using CT10-NLO PDF set from all sources as well as the total uncertainty, affecting the cross section measurement in each bin of  $H_{T,2}/2$  for inclusive 2-jet events.

Bin	Scale	PDF	NP	Total
300 - 330	+0.942 -6.149	+3.566 -3.090	0.825	+3.780 -6.931
330 - 360	+1.035 -6.289	+3.906 -3.342	0.736	+4.107 -7.159
360 - 390	+1.159 -6.438	+4.232 -3.573	0.696	+4.442 -7.396
390 - 420	+1.220 -6.536	+4.551 -3.794	0.723	+4.767 -7.592
420 - 450	+1.326 -6.660	+4.857 -3.997	0.745	+5.089 -7.802
450 - 480	+1.421 -6.776	+5.153 -4.186	0.765	+5.399 -8.001
480 - 510	+1.512 -6.888	+5.444 -4.365	0.782	+5.704 -8.192
510 - 540	+1.566 -6.967	+5.721 -4.527	0.797	+5.984 -8.347
540 - 570	+1.666 -7.082	+6.000 -4.682	0.810	+6.279 -8.528
570 - 600	+1.731 -7.172	+6.269 -4.825	0.822	+6.555 -8.683
600 - 640	+1.805 -7.271	+6.597 -4.979	0.833	+6.890 -8.852
640 - 680	+1.930 -7.416	+6.978 -5.143	0.845	+7.289 -9.064
680 - 720	+2.007 -7.527	+7.364 -5.295	0.856	+7.680 -9.243
720 - 760	+2.113 -7.663	+7.749 -5.437	0.865	+8.078 -9.436
760 - 800	+2.196 -7.781	+8.140 -5.569	0.873	+8.476 -9.609
800 - 850	+2.323 -7.945	+8.573 -5.706	0.881	+8.926 -9.822
850 - 900	+2.389 -8.062	+9.082 -5.863	0.889	+9.433 -10.008
900 - 950	+2.499 -8.227	+9.600 -6.018	0.896	+9.961 -10.232
950 - 1000	+2.631 -8.402	+10.134 -6.166	0.902	+10.509 -10.460
1000 - 1060	+2.738 -8.569	+10.747 -6.343	0.908	+11.127 -10.700
1060 - 1120	+2.853 -8.751	+11.431 -6.526	0.914	+11.817 -10.955
1120 - 1180	+2.992 -8.970	+12.183 -6.727	0.919	+12.579 -11.250
1180 - 1250	+3.135 -9.194	+13.019 -6.944	0.924	+13.423 -11.558
1250 - 1320	+3.324 -9.469	+14.004 -7.189	0.929	+14.423 -11.925
1320 - 1390	+3.434 -9.677	+15.080 -7.444	0.933	+15.494 -12.244
1390 - 1460	+3.629 -9.976	+16.223 -7.700	0.937	+16.650 -12.637
1460 - 1530	+3.760 -10.224	+17.505 -7.980	0.940	+17.929 -13.004
1530 - 1600	+3.894 -10.471	+18.891 -8.258	0.943	+19.311 -13.368
1600 - 1680	+4.107 -10.813	+20.496 -8.560	0.946	+20.925 -13.824
1680 - 1760	+4.421 -11.101	+22.481 -8.905	0.949	+22.931 -14.263
1760 - 1840	+4.921 -11.461	+24.654 -9.251	0.951	+25.158 -14.760
1840 - 1920	+5.404 -11.813	+27.143 -9.607	0.953	+27.692 -15.256
1920 - 2000	+5.867 -12.154	+29.986 -9.973	0.955	+30.570 -15.751

Table A.6: Theoretical uncertainties (in %) calculating using CT10-NLO PDF set from all sources as well as the total uncertainty, affecting the cross section measurement in each bin of  $H_{T,2}/2$  for inclusive 3-jet events.

Bin	Scale	PDF	NP	Total
300 - 330	+0.539 -8.294	+5.716 -4.657	1.692	+5.986 -9.662
330 - 360	+0.550 -8.577	+5.977 -4.779	1.516	+6.191 -9.935
360 - 390	+0.599 -8.709	+6.187 -4.987	1.363	+6.363 -10.128
390 - 420	+0.719 -8.948	+6.751 -5.223	1.228	+6.900 -10.433
420 - 450	+0.799 -9.145	+7.031 -5.395	1.110	+7.162 -10.676
450 - 480	+0.847 -9.247	+7.404 -5.578	1.005	+7.520 -10.845
480 - 510	+0.847 -9.294	+7.837 -5.717	0.937	+7.938 -10.951
510 - 540	+0.922 -9.436	+8.198 -5.884	0.921	+8.301 -11.158
540 - 570	+0.974 -9.566	+8.529 -6.000	0.904	+8.632 -11.328
570 - 600	+1.086 -9.786	+8.970 -6.156	0.886	+9.079 -11.595
600 - 640	+1.107 -9.852	+9.402 -6.297	0.866	+9.506 -11.724
640 - 680	+1.278 -10.101	+10.310 -6.526	0.842	+10.423 -12.055
680 - 720	+1.384 -10.342	+9.682 -6.618	0.820	+9.815 -12.305
720 - 760	+1.415 -10.404	+11.051 -6.826	0.798	+11.170 -12.469
760 - 800	+1.547 -10.615	+11.565 -7.009	0.777	+11.694 -12.744
800 - 850	+1.679 -10.804	+12.242 -7.185	0.755	+12.379 -12.997
850 - 900	+2.085 -11.134	+13.097 -7.461	0.731	+13.282 -13.422
900 - 950	+2.475 -11.432	+13.889 -7.703	0.709	+14.125 -13.804
950 - 1000	+2.655 -11.608	+14.614 -7.915	0.688	+14.869 -14.066
1000 - 1060	+3.025 -11.926	+15.576 -8.173	0.667	+15.881 -14.473
1060 - 1120	+3.299 -12.189	+14.250 -8.441	0.645	+14.641 -14.840
1120 - 1180	+3.741 -12.584	+17.984 -8.787	0.625	+18.380 -15.361
1180 - 1250	+3.969 -12.843	+19.324 -9.127	0.625	+19.737 -15.768
1250 - 1320	+4.663 -13.452	+21.246 -9.517	0.642	+21.761 -16.490
1320 - 1390	+4.878 -13.702	+22.884 -9.899	0.657	+23.407 -16.916
1390 - 1460	+5.242 -14.095	+24.854 -10.332	0.670	+25.410 -17.489
1460 - 1530	+5.582 -14.464	+27.170 -10.733	0.682	+27.746 -18.024
1530 - 1600	+6.003 -14.907	+29.741 -11.165	0.692	+30.349 -18.637
1600 - 1680	+6.503 -15.418	+32.855 -11.617	0.702	+33.500 -19.317

Table A.7: Theoretical uncertainties (in %) calculating using CT10-NLO PDF set from all sources as well as the total uncertainty, affecting the measurement of cross section ratio  $R_{32}$ , in each bin of  $H_{T,2}/2$ .

Bin	Scale	PDF	NP	Total
300 - 330	+0.038 -7.203	+2.458 -3.463	0.822	+2.592 -8.035
330 - 360	+0.027 -6.626	+2.317 -3.378	0.734	+2.431 -7.474
360 - 390	+0.024 -6.449	+2.149 -3.367	0.656	+2.247 -7.304
390 - 420	+0.084 -5.894	+2.411 -3.383	0.586	+2.482 -6.821
420 - 450	+0.113 -5.532	+2.345 -3.362	0.523	+2.405 -6.494
450 - 480	+0.109 -5.409	+2.390 -3.357	0.467	+2.438 -6.383
480 - 510	+0.073 -5.442	+2.506 -3.327	0.416	+2.541 -6.392
510 - 540	+0.107 -5.168	+2.559 -3.326	0.371	+2.588 -6.157
540 - 570	+0.112 -5.010	+2.586 -3.292	0.330	+2.609 -6.004
570 - 600	+0.163 -4.576	+2.729 -3.292	0.292	+2.750 -5.645
600 - 640	+0.146 -4.565	+2.824 -3.270	0.253	+2.839 -5.621
640 - 680	+0.198 -4.163	+3.368 -3.298	0.236	+3.382 -5.316
680 - 720	+0.155 -3.754	+2.352 -3.247	0.227	+2.368 -4.968
720 - 760	+0.196 -3.842	+3.267 -3.268	0.219	+3.280 -5.049
760 - 800	+0.126 -3.523	+3.366 -3.272	0.212	+3.375 -4.813
800 - 850	+0.110 -3.368	+3.596 -3.261	0.206	+3.604 -4.693
850 - 900	+0.048 -3.351	+3.909 -3.309	0.200	+3.915 -4.714
900 - 950	+0.116 -3.504	+4.148 -3.334	0.196	+4.154 -4.841
950 - 1000	+0.127 -3.511	+4.300 -3.335	0.192	+4.306 -4.846
1000 - 1060	+0.282 -3.683	+4.604 -3.357	0.204	+4.617 -4.988
1060 - 1120	+0.436 -3.779	+3.079 -3.375	0.224	+3.118 -5.071
1120 - 1180	+0.732 -3.982	+5.430 -3.452	0.241	+5.485 -5.276
1180 - 1250	+0.813 -4.031	+5.835 -3.511	0.258	+5.897 -5.352
1250 - 1320	+1.303 -4.414	+6.626 -3.591	0.275	+6.759 -5.697
1320 - 1390	+1.403 -4.471	+7.036 -3.659	0.290	+7.180 -5.785
1390 - 1460	+1.564 -4.590	+7.657 -3.778	0.304	+7.822 -5.953
1460 - 1530	+1.765 -4.738	+8.438 -3.853	0.316	+8.626 -6.115
1530 - 1600	+2.040 -4.972	+9.306 -3.962	0.328	+9.532 -6.366
1600 - 1680	+2.313 -5.179	+10.381 -4.075	0.339	+10.641 -6.599

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