



The CLAS12 Forward Time-of-Flight system

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ABSTRACT

The Forward Time-of-Flight system for the large-acceptance CLAS12 spectrometer in Hall B at the Thomas Jefferson National Accelerator Facility is described. The system is positioned at distances in the range from 6.2 m to 7.2 m from the beam–target interaction point and spans laboratory polar angles from 5° to 45° and nearly the full azimuth. The system consists of 540 individual scintillation counters with double-ended readout that range in length from 17 cm to 426 cm of discrete widths of 6 cm, 15 cm, and 22 cm, and of discrete thicknesses of 5 cm and 6 cm. The effective counter time resolution for passing charged particles varies from 50 ps for the shortest counters at small angles to 200 ps for the longest counters at large angles. The detectors are part of the forward-angle particle identification system for CLAS12 during offline event reconstruction and are a component of the online data acquisition trigger to select final state event topologies with forward-going charged particles.

1. Overview of CLAS12

The Thomas Jefferson National Accelerator Facility (JLab) recently completed a project to double the maximum energy of its electron accelerator from 6 GeV to 12 GeV. The experimental equipment in Hall B forms the large-acceptance CLAS12 spectrometer that is designed to operate with beam energies up to 11 GeV at a beam–target luminosity of up to $10^{35} \text{ cm}^{-2} \text{ s}^{-1}$ to allow for precision measurements of exclusive reactions with polarized beams and both unpolarized and polarized targets. This spectrometer is based on two superconducting magnets, a solenoid in the central region around the target and a toroid at forward angles.

The CLAS12 torus magnet has a six-fold symmetry that divides the forward azimuthal acceptance in the polar angle range from 5° to 35° into six 60°-wide sectors. The torus produces a field primarily in the azimuthal direction of strength $\int B d\ell$ at its nominal full current of 2.8 Tm at 5° and 0.5 Tm at 35°. A set of three multi-layer drift chambers in each sector (before the field, within the field, and after the field) and a forward micromegas vertex tracker are used for charged particle tracking to measure momenta. Downstream of the torus each sector is instrumented with a Cherenkov counter for π/K separation (four sectors are instrumented with low threshold gas Cherenkov counters, one sector is instrumented with a ring-imaging Cherenkov counter, and the final sector will eventually be instrumented with a second

ring-imaging Cherenkov counter), three planar layers of scintillation counters for charged particle time measurements called the Forward Time-of-Flight (FTOF) system, and an electromagnetic calorimeter system for electron and neutral particle identification. Just upstream of the torus is a large-volume high-threshold gas Cherenkov counter for electron identification and a tagging system to detect electrons and photons at polar angles below 5°.

The CLAS12 solenoid spans the central angular range from 35° to 125° and has a uniform 5 T central field at its nominal full current. The solenoid serves to focus the low-energy Møller background down the beam pipe to the beam dump away from the acceptance of the spectrometer. The detectors mounted within the solenoid include a thick scintillation counter for neutron identification, a barrel of thin scintillation counters for charged particle timing measurements, and a set of tracking detectors around the target.

Fig. 1 shows a model representation of CLAS12 to highlight its overall layout and scale. See Ref. [1] (and references therein) for more complete information on CLAS12 and its individual subsystems. CLAS12 was installed and instrumented in Hall B in the period from 2012 to 2017 and took the place of the original CLAS spectrometer [2] that operated in Hall B in the period from 1997 to 2012 when it was decommissioned.

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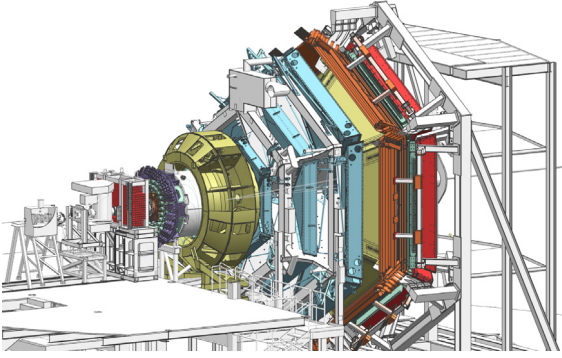


Fig. 1. Model representation of the CLAS12 spectrometer in Hall B at Jefferson Laboratory. The electron beam is incident from the left side of this figure. The CLAS12 detector is roughly 20 m in scale along the beam axis.

This paper focuses on the CLAS12 FTOF detector system and is organized as follows: Section 2 reviews the scheme for particle identification in the CLAS12 Forward Detector, Section 3 provides a high-level overview of the FTOF system and its overall design requirements, Section 4 provides a technical description of the system design, and Section 5 highlights the performance of the system through both bench testing with cosmic rays, as well as during the 2017 commissioning run and 2018 first production running with electrons. Finally, Section 6 provides a summary of the FTOF system for CLAS12.

2. CLAS12 forward detector particle identification

Particle identification in the CLAS12 Forward Detector relies on input from each of the different Forward Detector subsystems. A reconstructed track in the drift chambers (DC) [3] identifies the presence of a charged particle and is used as a veto for forward-going neutral particles. The curvature of the particle tracks in the magnetic field of the torus provides the electric charge and momentum. The other detector subsystems in the forward direction are used to identify the particle type. These subsystems include the different Cherenkov counters (HTCC [4], LTCC [5], and RICH [6]), the electromagnetic calorimeters (ECAL [7]), and the FTOF. These systems are used as part of the overall CLAS12 particle identification scheme to separate the different particle species as a function of momentum. See Ref. [1] for details on the different CLAS12 detector subsystems used for forward-going charged particle identification and the range of momenta for which they are responsible for the separation of the different particle species.

The FTOF is the primary system for particle identification in CLAS12 for forward-going charged particles for momenta up to ~ 5 GeV. Due to the requirements of the physics program to separate charged pions and kaons up to ~ 3 GeV, the FTOF was designed to measure the flight time of charged particles emerging from the target with an average time resolution of 80 ps. Given this nominal time resolution for the counters, the momentum threshold for particle identification can be defined. For our purposes, thresholds are given at the 4σ level for FTOF, which amounts to the momenta where particle identification can occur with up to an order of magnitude difference in the relative yields of the different species. The time resolution is illustrated by computing the flight time differences between different charged particle species, pions, kaons, and protons, for tracks normally incident on the detector. Fig. 2 shows the computed time differences as a function of momentum. Where the 4σ line crosses the computed time difference curves defines the momentum limit for particle identification for each particle species. These limits are given as 2.8 GeV for π/K separation, 4.8 GeV for K/p separation, and 5.4 GeV for π/p separation.

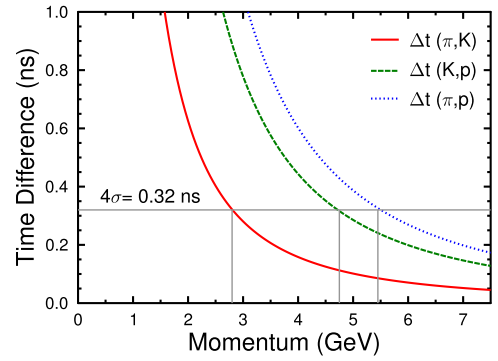


Fig. 2. Flight time differences (ns) between protons and pions, protons and kaons, and kaons and pions (as indicated) for a 7 m path length from the target to the FTOF vs. particle momentum (GeV). The horizontal line indicates a time difference four times larger than the average FTOF counter design resolution of $\sigma_{TOF} \approx 80$ ps. The vertical lines that meet each curve represent the momentum limit for 4σ particle species separation.

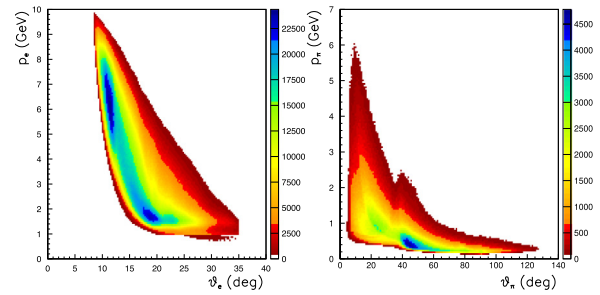


Fig. 3. Plot of momentum vs. lab polar angle from beam data for a 10.6 GeV electron beam incident upon a liquid-hydrogen target in CLAS12 for scattered electrons (left) and π^\pm (right). The discontinuity at $\theta = 35^\circ$ is due to the small acceptance gap between the Forward and Central Detectors. The typical momentum of charged hadron tracks in the Forward Detector in these kinematics is between 0.5 GeV and 6 GeV.

Fig. 3 illustrates the momentum vs. polar angle coverage in CLAS12 from beam data of a 10.6 GeV electron beam incident upon a liquid-hydrogen target. The plots show the kinematic phase space for scattered electrons (left) and charged pions (right) for the semi-inclusive reactions $ep \rightarrow e'\pi^\pm X$ (X represents all other possible reaction products). For these reactions the typical charged hadron track momenta accepted by FTOF are in the range from 0.5 GeV to 6 GeV.

3. Overview of the FTOF system

The Forward Time-of-Flight System (FTOF) is a major component of the CLAS12 Forward Detector used to measure the time-of-flight of charged particles emerging from interactions in the target. The requirements for FTOF include excellent time resolution for charged particle identification and good segmentation to minimize count rates and to provide for flexible trigger options (for details on FTOF in the CLAS12 trigger, see Ref. [8]). The system specifications call for an average time resolution of $\sigma_{TOF} = 80$ ps at the more forward angles of CLAS12 and 150 ps at angles larger than 35° . The system must also be capable of operating in a high-rate environment where the maximum count rate for each FTOF scintillator at an operating luminosity of $1 \times 10^{35} \text{ cm}^{-2} \text{ s}^{-1}$ is up to 1 MHz.

In each of the six 60° -wide sectors of the CLAS12 Forward Detector, the FTOF system is comprised of three arrays of counters, referred to as panels, named panel-1a, panel-1b, and panel-2. Each panel consists of a set of rectangular scintillators with a photomultiplier tube (PMT) on each end. Panel-1 refers to the counters located at forward angles (5° to 35°) (where two panels are employed to meet the 80 ps average

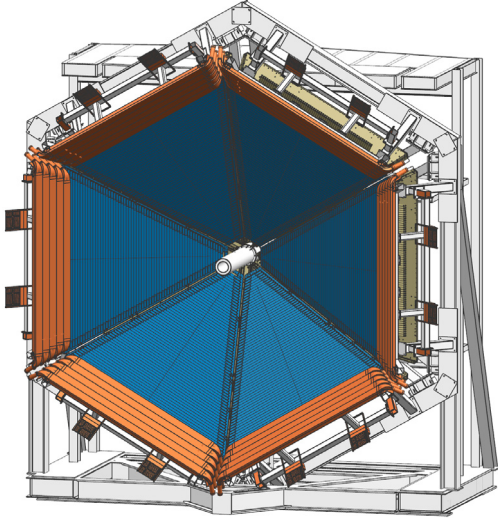


Fig. 4. View of the FTOF counters for CLAS12 highlighting the location of the counters. The panel-1b counters are shown in dark blue and the panel-2 counters, mounted around the perimeter of the Forward Carriage, are shown in light orange. The panel-1a counters, mounted just downstream of the panel-1b counters, are not visible in this picture. The Forward Carriage is roughly 10 m in diameter. (For interpretation of the references to color in this figure legend, the reader is referred to the web version of this article.)

Table 1

Parameters for the scintillators, PMTs, and counters for the FTOF panel-1a, panel-1b, and panel-2 arrays in each of the six sectors of the CLAS12 Forward Carriage.

Parameter	Design Value
Panel-1a (23 counters per sector)	
Angular Coverage	$\theta = 5^\circ \rightarrow 35^\circ$, ϕ : 50% at $5^\circ \rightarrow 85^\circ$ at 35°
Counter Dimensions	$L = 32.3 \text{ cm} \rightarrow 376.1 \text{ cm}$, $w \times h = 15 \text{ cm} \times 5 \text{ cm}$
Scintillation Material	BC-408
PMTs	EMI 9954A, Philips XP2262
Counter Time Resolution	90 ps \rightarrow 180 ps
Panel-1b (62 counters per sector)	
Angular Coverage	$\theta = 5^\circ \rightarrow 35^\circ$, ϕ : 50% at $5^\circ \rightarrow 85^\circ$ at 35°
Counter Dimensions	$L = 17.3 \text{ cm} \rightarrow 407.9 \text{ cm}$, $w \times h = 6 \text{ cm} \times 6 \text{ cm}$
Scintillation Material	BC-404 (#1 \rightarrow #31), BC-408 (#32 \rightarrow #62)
PMTs	Hamamatsu R9779
Counter Time Resolution	60 ps \rightarrow 110 ps
Panel-2 (5 counters per sector)	
Angular Coverage	$\theta = 35^\circ \rightarrow 45^\circ$, ϕ : 85% at $35^\circ \rightarrow 95^\circ$ at 45°
Counter Dimensions	$L = 371.3 \text{ cm} \rightarrow 426.1 \text{ cm}$, $w \times h = 22 \text{ cm} \times 5 \text{ cm}$
Scintillation Material	BC-408
PMTs	Photonis XP4312B, EMI 4312KB
Counter Time Resolution	170 ps \rightarrow 180 ps

time resolution requirement) and panel-2 refers to the sets of counters at larger angles (35° to 45°). The positioning and attachment of the FTOF detectors on their Forward Carriage supports are shown in Fig. 4.

The FTOF counters in the angular range from 5° to 35° consist of two sets of six triangular arrays. Just upstream of the electromagnetic calorimeter (ECAL) detectors, the panel-1a arrays are mounted. These detector sets are the refurbished panel-1 TOF counters from the decommissioned CLAS spectrometer [9]. Upstream of the panel-1a arrays the new panel-1b arrays are mounted. In the event reconstruction the hit times for panel-1a and panel-1b are combined together to determine the charged particle hit time (see Ref. [10] for details). In the angular range from 35° to 45° the panel-2 arrays are mounted. These counters are refurbished panel-2 counters of the CLAS TOF system. A detailed summary of the FTOF technical parameters is given in Table 1.

The panel-1 arrays consist of the old CLAS panel-1 TOF counters (called panel-1a) and a new set of panel-1 counters (called panel-1b).

The panel-1a arrays consist of 23 scintillators, each measuring 5.08-cm thick and 15-cm wide. The lengths of these counters range from 32 cm at the smallest scattering angles to 376 cm at the largest scattering angles. The scintillators are constructed from Bicron BC-408 and are coupled to short acrylic light guides read out with 2-in Thorn EMI-9954A PMTs. The new panel-1b arrays consist of 62 scintillators 6-cm wide by 6-cm thick constructed from Bicron BC-404 scintillator for the shortest 31 counters and BC-408 for the longest 31 counters. The lengths of these counters range from 17 cm at the smallest scattering angles to 408 cm at the largest scattering angles. The scintillators are read out by 2-in Hamamatsu R9779 PMTs coupled directly to the scintillation bars. These new panels are mounted to the Forward Carriage in front of the panel-1a counter arrays. The design and bench testing results for these counters are described in detail in Ref. [11].

The panel-2 arrays consist of selected counters from the old CLAS panel-2 TOF counters, and include 5 22-cm wide, 5.08-cm thick scintillators in each sector. The length of these counters ranges from roughly 370 cm to 440 cm. These scintillators are constructed from Bicron BC-408 and are read out through curved acrylic light guides coupled to 3-in Philips XP4312B PMTs. These scintillators are included to give complete acceptance for forward-going, low-momentum, outbending charged particles.

4. FTOF system design

In order to meet the performance and mechanical requirements for the FTOF system, the major considerations in its design included the system geometry and areal coverage, the counter time resolutions, the system components, and the design and materials associated with its mechanical support structure in the active area of the spectrometer. These system design elements are described in the following subsections. In addition, this section also includes information on the readout electronics and the high voltage system used for the FTOF.

4.1. Geometry

The projected space behind the coils of the CLAS12 torus as defined by straight lines projecting radially outward from the center of the nominal target position is referred to downstream of the torus as its “shadow” region. This region is inactive and defines the space available for locating the light guides, PMTs, voltage dividers, and signal and high voltage cables. The remaining area in the forward direction is the sensitive fiducial region of the detector and must be covered by scintillation counters. The design specification for FTOF called for a minimum of 50% azimuthal acceptance at 5° increasing to 95% at 45° .

Fig. 5 provides an illustration of the shadow region projected onto the Forward Carriage created primarily by the torus cryostats and the drift chamber endplates as projected onto the face of the FTOF system. Fig. 5 (top) shows a picture of the shadow bands on the Forward Carriage that defines a uniform gap of ~ 40 cm between each sector. Fig. 5 (bottom) shows the defined active region in one sector of CLAS12 on the face of the FTOF. The azimuthal width of this area at the position of the Forward Carriage in Hall B essentially defined the length of the scintillation counters. The final limits of the shadow region at the location of FTOF are actually defined by the endplates of the drift chamber system [3] located within the torus coils. The drift chamber systems upstream of the torus and downstream of the torus have their endplates, on-board electronics, and readout cables located mainly in the shadow of the torus cryostats.

The FTOF panel-1a and panel-1b arrays in each sector are triangular in shape with the shortest counters located closest to the beamline and the longest counters furthest from the beamline. The length of each counter for a given counter number N_C is as follows:

- Panel-1a:
 $L [\text{cm}] = 15.85 \times N_C + 16.43 \quad (N_C = 1 \rightarrow 5),$
 $L [\text{cm}] = 15.85 \times N_C + 11.45 \quad (N_C = 6 \rightarrow 23),$

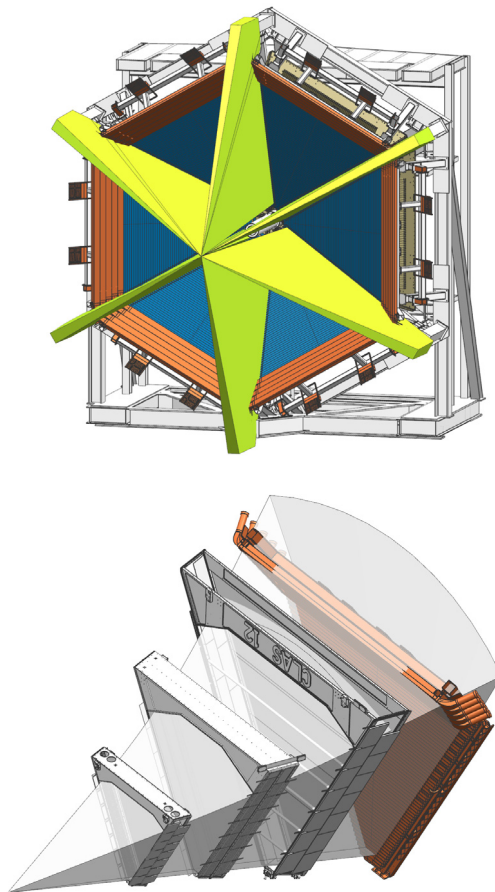


Fig. 5. (Top) View of the shadows created by the main torus cryostats and drift chamber endplates as projected on the face of the FTOF system. (Bottom) The defined active area between the shadow projections through the three regions of drift chambers projected onto the face of the FTOF in a representative Forward Carriage sector.

- Panel-1b:

$$L \text{ [cm]} = 6.40 \times N_C + 10.84 \quad (N_C = 1 \rightarrow 62).$$

The panel-1a and panel-1b arrays are tilted toward the target at an angle of 25° consistent with the other subsystems in the CLAS12 Forward Detector (DC, LTCC, RICH, ECAL). The panel-1a counters are located at a radial distance from the target in the range from $R = 724.21$ cm for $N_C = 1$ to $R = 691.74$ cm for $N_C = 23$. The panel-1b counters are located at a radial distance from the target in the range from $R = 716.15$ cm for $N_C = 1$ to $R = 677.97$ cm for $N_C = 62$. The gap between the coplanar panel-1b and panel-1a arrays in each sector is 10.72 cm. The minimum angle covered by panel-1a based on a straight line from the target is 5.453° . The corresponding minimum angle covered by panel-1b is 3.667° . Each of the panel-1a arrays covers an area of 7.0 m^2 and each of the panel-1b arrays covers an area of 7.9 m^2 . Fig. 6 shows a two-dimensional schematic of the layout and positioning of the arrays defining the key geometry parameters, which are listed in Table 2. See Ref. [12] for more information.

The panel-2 arrays are mounted outside of the panel-1a and panel-1b arrays at larger polar angles as shown in Fig. 6. The length of each counter for a given counter number N_C is as follows:

- Panel-2:

$$L \text{ [cm]} = 13.73 \times N_C + 357.55 \quad (N_C = 1 \rightarrow 5).$$

The panel-2 arrays are tilted toward the target at an angle of 58.11° . The minimum angle covered by panel-2 based on a straight line from the target is 34.698° . Each of the six panel-2 arrays covers an area of

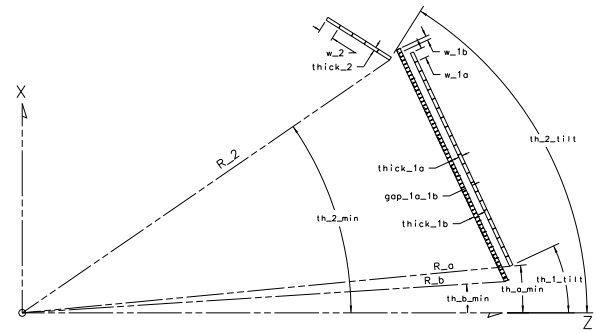


Fig. 6. View of the FTOF scintillators for panel-1a, panel-1b, and panel-2 in the sector mid-plane for one representative sector of the CLAS12 Forward Detector with the key parameters indicated. See Table 2 for the nominal parameter values.

Table 2
Nominal geometry parameters for the CLAS12 FTOF detector system.

Parameter	Panel-1a	Panel-1b	Panel-2
R_min	726.689 cm	717.236 cm	659.71 cm
th_min	5.453°	3.667°	34.698°
th_tilt	25.00°	25.00°	58.11°
thick	5.08 cm	6.00 cm	5.08 cm
width	15.00 cm	6.00 cm	22.00 cm
gap_1a_1b	10.717 cm		–

4.4 m². Note that the panel-2 arrays have no direct line of sight to the target due to the solenoid. However, due to the presence of the toroidal magnetic field, they provide additional acceptance for outbending, low momentum tracks and for tracks associated with strange particles that decay in-flight after emerging from the target (e.g. $\Lambda \rightarrow N\pi$ with $c\tau = 7.89$ cm).

Given the active area coverage requirements for the FTOF system within each CLAS12 sector on the Forward Carriage, another key aspect of the geometry associated with the FTOF system design is the width of the individual scintillation counters. An essential optimization was made to select the counter width to minimize the number of readout channels, while considering the overall count rates per bar at the nominal luminosity associated with incident charged and neutral particles (including photons). These rates must allow for reasonable PMT anode currents that do not affect the stability of the PMT response in terms of pulse shape or saturation effects, nor lead to unreasonably short PMT lifetimes. In addition, the width of the scintillation bars determines the granularity of the scattering angle definition in the trigger and its matching to the projected tracks from the drift chambers to the electromagnetic calorimeters. Note that the 15-cm widths of panel-1a and the 22-cm widths of panel-2 of the existing refurbished counters of the CLAS TOF system were optimized for nominal beam-target luminosities a factor of 10 lower than for CLAS12. The 6-cm widths of the newly constructed panel-1b counters were optimized for the higher rate operating conditions of CLAS12.

4.2. Counter hit time resolution

The FTOF counters are designed to provide an output signal for the CLAS12 data acquisition system [13] that reflects the time a charged particle passed through the scintillation counter. As the particle passes through the scintillation material, it causes ionization that subsequently generates scintillation light. The photons that are created travel on various paths inside of the scintillator and light guide (if present), and may get absorbed or reflected (internally or on outer wrapping materials) before they can impinge on the photocathode of the PMT. This interaction produces a photoelectron signal that is amplified within the stages of the PMT and the generated pulse is then input into the readout electronics, which includes an analog-to-digital converter (ADC), a

Table 3

Parameters determined for the CLAS TOF panel-1a and panel-2 counters in Ref. [9] and used for a parameterization of the CLAS12 FTOF counters using the functional form for σ_{TOF} in Eq. (1) and for N_{pe} in Eq. (2).

Parameter	Nominal value
σ_0	0.062 ns (CLAS TOF); 0.040 ns (CLAS12 FTOF)
σ_1	2.1 ns (panel-1a/1b); 2.0 ns (panel-2)
σ_2	2.0 ns/m
N_{pe}^0	918
λ	$0.358 \cdot L + 81.725$ cm

discriminator, and a time-to-digital converter (TDC) (see Section 4.4 for details). The net effect of these different processes accounts for the time resolution of the counter.

The contributions to the time resolution of time-of-flight systems have been parameterized in Ref. [14] by:

$$\sigma_{TOF} = \sqrt{\sigma_0^2 + \frac{\sigma_1^2 + (\sigma_2 L/2)^2}{N_{pe}}}. \quad (1)$$

Here σ_{TOF} represents the timing resolution for a scintillation counter with double-sided PMT readout. σ_0 represents the intrinsic electronic resolution of the measurement system, a floor-term contribution that is independent of the light level. The remaining terms σ_1 and σ_2 are directly dependent on the average photo-statistics seen at either of the PMT photocathodes N_{pe} (see Eq. (2)). The term σ_1 models the jitter in the combined single photoelectron response of the scintillation counter and its PMTs and the term σ_2 accounts for path length variations in the light collection. These path length variations in the scintillator scale with the distance from the source to the PMT, which we take to be half the length of the counter ($L/2$), since the scintillators are read out at either side. The statistical behavior of the last two terms is indicated by scaling the single-photoelectron responses by $\sqrt{N_{pe}}$. For scintillators that are several meters long, the dominant contribution to the timing resolution comes from transit time variations of photon paths along the scintillator to the PMT due to the counter geometry.

The values of the parameters σ_0 , σ_1 , and σ_2 for the panel-1a and panel-2 counters are given in Ref. [9], where the above functional form with the parameters listed in Table 3 was found to describe the measured data. A direct extension of these parameters is assumed to be reasonable for estimating the time resolution for the panel-1b counters. A summary of all parameters employed are listed in Table 3. Note that due to improvements in the resolution of the readout electronics for the CLAS12 FTOF system compared to the CLAS TOF system, the floor-term σ_0 has been reduced from 62 ps to 40 ps.

The number of photoelectrons N_{pe} in Eq. (1) for panel-1a and panel-2 at the PMT photocathode was determined in Ref. [9] by:

$$N_{pe} = N_{pe}^0 \exp\left(\frac{L_0}{2\lambda_0} - \frac{L}{2\lambda}\right) \cdot F, \quad (2)$$

where N_{pe} for all counters was referenced to the average value measured for the response of the shortest panel-1a counter N_{pe}^0 of length $L_0 = 32$ cm with attenuation length λ_0 . The attenuation length of the scintillation bars represents the distance λ into the material where the probability that the photon has been absorbed is $1/e$. For the panel-2 counters, N_{pe} is further scaled by the factor $F = 0.9$ to account for light collection efficiencies at the end of the larger panel-2 counters with their 3-in PMTs and longer light guides compared to the smaller panel-1a PMTs with their relatively short light guides [9]. For the panel-1b counters, N_{pe} is determined as for panel-1a using Eq. (2) by scaling by the ratio of the cross sectional areas of the scintillation bars (15 cm \times 5 cm vs. 6 cm \times 6 cm).

Fig. 7 shows the parameterized resolution for the counters in panel-1a, panel-1b, and panel-2 as a function of counter length. The Forward Detector event reconstruction and particle identification uses time information from both panel-1a and panel-1b. For tracks that

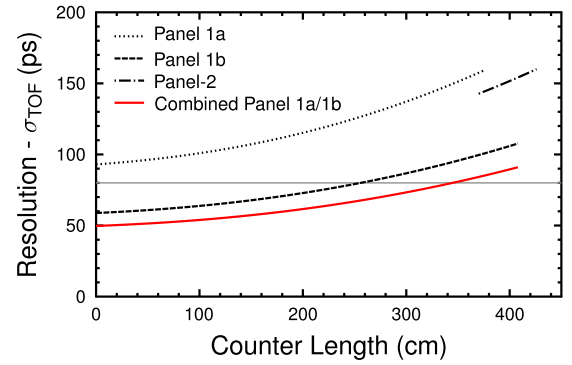


Fig. 7. Parameterized expectation of the counter hit time resolution for the FTOF panel-1a (dotted), panel-1b (dashed), and panel-2 (dot-dashed) counters as a function of length. The solid (red) line indicates the final expected resolution in the forward direction by combining the hit time information from the panel-1a and panel-1b counters. The horizontal line indicates the 80 ps average time resolution specification for the FTOF system.

Table 4

Properties of the plastic scintillation material BC-404 and BC-408 employed for the counters of the FTOF system [15].

Property	BC-404	BC-408
Light Output, % Anthracene	68	64
Rise Time (ns)	0.7	0.9
Decay Time (ns)	1.8	2.1
Pulse Width, FWHM (ns)	2.2	2.5
Wavelength of maximum emission (nm)	408	425
Light attenuation length (cm)	140	210
Bulk attenuation length (cm)	160	380
Polymer base	Polyvinyltoluene	
Refractive index	1.58	
Density (g/cm ³)	1.023	

pass through both arrays the combined time information (described in Ref. [10]) is used and results in a 20% improvement compared to using the hit information from panel-1b alone.

4.3. System components

4.3.1. Scintillator material

To optimize the time resolution over the full volume of the FTOF counters, scintillation materials with fast time response and long attenuation length are essential. For the panel-1a and panel-2 FTOF counters that were refurbished from the older CLAS TOF system, Bicron BC-408 was selected. For the panel-1b counters constructed for the new CLAS12 FTOF system, a different design approach was considered that optimized the overall system time resolution. For counters less than 2 m in length, the overall performance is improved by the use of a faster scintillator with a small decay time τ_{decay} , whereas for the longer counters, a material with a longer attenuation length is the better choice. The final decision for the panel-1b counters was to use BC-404 for counters 1 \rightarrow 31 (lengths from 17.3 cm to 209.4 cm) and BC-408 for counters 32 \rightarrow 62 (lengths from 215.8 cm to 407.9 cm). Table 4 lists the properties of the FTOF scintillation materials.

The bulk attenuation length of the scintillation material is stated by its manufacturer to be 160 cm for BC-404 and 380 cm for BC-408. However, the practical attenuation length of the prepared bars is smaller than the nominal bulk value as the actual path length of photons from the charged particle intersection point to the ends of the bar is increased and reflections occur due to the finite geometry of the bar. For optimal response, this practical attenuation length should typically be longer than the bar to ensure sufficient photon statistics. Measurements of the practical attenuation length of the FTOF counters are given in Section 5.2.2.

Table 5

Properties of the PMTs used for the readout of the FTOF panel-1a, panel-1b, and panel-2 counters. All of these PMTs have a borosilicate glass window and employ green-sensitive bialkali photocathodes.

	9954A	R9779	XP4312B/D1
Property	Panel-1a	Panel-1b	Panel-2
Diameter	2 in	2 in	3 in
Photocathode area	16.6 cm ²	16.6 cm ²	36.3 cm ²
Dynode stages	12	8	12
Spectral response	290 → 680 nm	300 → 650 nm	290 → 650 nm
Max. wavelength emission	400 nm	420 nm	420 nm
Gain	1.8×10^7	5.0×10^5	3×10^7
Quantum eff. @ λ_{max}	28%	28%	28%
Max. anode current rating	100 μ A	100 μ A	100 μ A
Anode dark current	2 nA	10 nA	10 nA
Anode pulse rise time	2 ns	1.8 ns	2.1 ns
Electron transit time	41 ns	20 ns	31 ns
Transition time spread	0.4 ns	0.25 ns	0.4 ns

4.3.2. Photomultiplier tubes and voltage dividers

The panel-1a counters are read out at either end through 2-in diameter Thorn EMI 9954A PMTs (later manufactured by Electron Tubes). The PMTs are coupled to the scintillation bars using 12-cm-long acrylic light guides that match the 15 cm \times 5 cm scintillator on one end and the 2-in diameter PMT on the other end. For the panel-2 counters, 3-in diameter Philips XP4312B/D1 3-in PMTs (later manufactured by Photonis) are employed. The PMTs are coupled to the scintillators through acrylic light guides that match the 22 cm \times 5 cm scintillator on one end and the 3-in diameter PMT on the other end. Both the 9954A and XP4312B/D1 PMTs have 12 linear-focused dynode stages. For both the panel-1a and panel-2 counters the PMTs are glued on the light guides using BC-600 optical glue. See Ref. [9] for full details on the PMT selection criteria and the light guide designs.

The voltage dividers employed for the panel-1a and panel-2 readout are custom units built specifically for the CLAS TOF project [9]. The dividers use high voltage field-effect transistors to fix the PMT gain by stabilizing the voltage and to protect the PMT against high light levels by shutting down the circuit in case of over-current. The grid voltage for both types of dividers followed the manufacturer's specifications.

The photomultipliers employed for the panel-1b counters are 2-in diameter Hamamatsu R9779 PMTs with a high gain selection (minimum 0.5×10^5 , average 1.0×10^6) that have been integrated with a voltage divider to form the R9779-20MOD assembly. These PMTs include 8 linear-focused dynode stages. This high time resolution PMT was selected due to its particularly compact overall assembly length of 113 mm. The length restriction was necessary to fit within the defined shadow region of the torus cryostats at the location of the PMTs (see Section 4.1 for details). These PMTs are coupled directly to the ends of the scintillation bars using BC-600 optical glue. The performance specifications for all FTOF PMTs are listed in Table 5.

4.3.3. PMT magnetic shielding

The FTOF PMTs are located in the range from 6.2 m to 7.2 m from the target in a region where the stray magnetic field from the torus is computed to be less than 30 G when the torus is operated at full field. Custom magnetic shields for the PMTs are included to reduce both the axial and transverse components of the field along the full accelerating structure of the PMT to a level of less than 0.2 G. For the panel-1a counters, the PMT magnetic shields consist of 7.5-in long, 2-in diameter, 0.040-in thick μ -metal cylinders. For the panel-2 counters, the PMT magnetic shields consist of a 9.5-in long, 3-in diameter, 1.5-mm thick μ -metal cylinders. These shields extend from the voltage divider attachment point to 2 in beyond the front face of the PMT. They are held in place by the PMT itself and the counter support structure (see Ref. [9] for details). For the panel-1b counters, the magnetic shields are composed of 2-mm thick μ -metal boxes, 6.7-in long by 2.36-in wide. These boxes contain the entire voltage divider assembly and

extend 1.6 in beyond the front face of the PMT. Note that the last 4 cm of each end of the panel-1b scintillation bars were machined down to a width of 5.6 mm (with a diamond-tool finish) to allow the shield boxes to fit over the ends of the bars. The shields are screwed down to the counter support structure to hold them into position. The face of the shield box opposite the PMT side has a small penetration to allow the signal and HV cables and connectors to pass through. Further details on the FTOF magnetic shielding and the field tests that were conducted are included in Ref. [16]. All FTOF μ -metal shields are made from 80% nickel high permeability alloy with hydrogen annealing.

4.3.4. Counter assembly and support

Each scintillation counter is individually wrapped first with a reflective layer and then an opaque outer layer. For panel-1a and panel-2 the scintillation counter wrapping materials include:

- 1 layer of 9.4 mil thick black Kapton,
- 2 layers of 1 mil thick aluminum foil.

For panel-1b, the scintillation counter wrapping materials include:

- 3 layers of 1.5 mil thick Tedlar,
- 1 layer of 0.3 mil thick aluminized polyester film.

After wrapping, each of the FTOF scintillation counters was attached to a support structure that runs along the full length of the scintillation bar. The attachment was made using multiple bands of 4-in-wide fiberglass tape that wrapped around the counter and support structure at discrete locations along its length. These support structures are necessary to reduce the gravitational sagging of the scintillation bars. This structure consists of a composite sandwich structure of thin stainless steel skins over structural foam that is attached to the detector frame at the ends of each counter. The composite structure, which mounts on the scintillator side facing away from the target, provides uniform material thickness to the detected particles. The support was undersized across the counter width so they could be placed as close together as allowed by the wrapping material and the scintillation bar manufacturing tolerances. See Ref. [9] for more details.

Each panel-1a counter is mounted on a 1-in thick support to minimize the thickness of the package from the standpoint of Coulomb multiple scattering and energy loss considerations. The maximum deflection for the installed scintillators is 4.4 mm, as estimated from deflection tests and the compound angle of each detector, which relieves the overall support requirements. The space behind the panel-2 counters allowed for 3-in thick sandwich supports, which are mechanically much stiffer and result in no appreciable deflection. Again, each panel-2 counter is mounted to its own support structure. For the panel-1b counters, the backing structures are 2-in thick and designed to support two panel-1b counters. The maximum deflection for the installed scintillators is less than 5 mm, which occurs at the middle of the longest counters.

The support structures onto which the scintillator counters are attached are bolted to box-beam support frames (steel for panel-1a and panel-2, aluminum for panel-1b) that reside in the torus shadow regions. The support frames are triangular in shape for panel-1a and panel-1b, and form a rhombus shape for panel-2. The panel-1a frames are bolted directly to the upstream faces of the electromagnetic calorimeters in each sector of the Forward Carriage. The panel-1b frames are bolted directly to the panel-1a frames. The panel-2 frames are attached to the steel super-structure of the Forward Carriage.

4.4. Electronics

The outputs from the FTOF PMTs include both an anode and a dynode signal. The anode signals are sent first to a discriminator and then to a TDC. The dynode signals are sent to a flash ADC (FADC). A

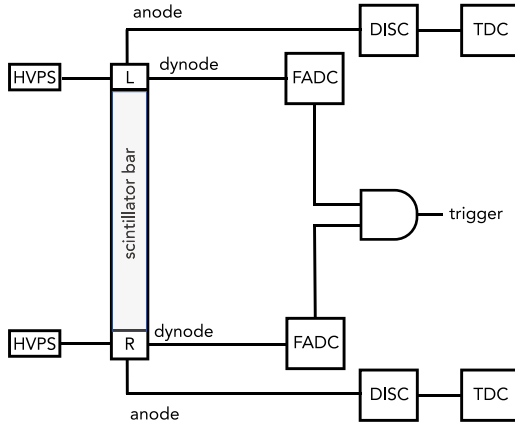


Fig. 8. Schematic of the electronics for each counter in the CLAS12 FTOF system.

block diagram of the electronics layout for each FTOF counter is shown in Fig. 8.

For the FTOF PMTs the anode signal is roughly three times larger in amplitude than the dynode signal. For the panel-1a and panel-2 PMTs, the dynode signals are bipolar with a negative polarity primary pulse with a long tail that overshoots the baseline. This tail is not included in the determination of the pulse charge. For the panel-1b PMTs, the anode signal has negative polarity and the dynode signal has positive polarity. To ensure compatibility with the negative polarity input requirements of the FADC, the dynode signal is inverted before the readout electronics using an inline Phillips Scientific 460 IT inverting transformer.

The output from the panel-1b FADCs is also used as part of the CLAS12 level-1 trigger to select charged particles. Signals in panel-1b above the FADC threshold are geometrically matched to hits in the electromagnetic calorimeter and to found track candidates in the drift chambers (see Ref. [8]). Signals from the FTOF system are also used to provide an effective charged particle veto for the detection of neutrals in the electromagnetic calorimeters. While high resolution time measurements are the primary role of the FTOF system for charged particle identification in the forward direction of CLAS12, the pulse height information from the FADCs is also employed for energy loss measurements to provide an independent means for identification of slow particles. In addition, pulse fitting techniques are employed using the FADC pulse shape to determine the hit time of the track that can be compared to the TDC time to better ensure matching of the ADC and TDC information in the high rate operating environment of CLAS12 (see Section 5.2.8). In order to minimize the amount of data collected from the FTOF readout to reduce data sizes, a deposited energy threshold is applied to the FADC readout corresponding to approximately 1 MeV.

The intrinsic resolution of the electronics system (σ_0) must be optimized to ensure that it does not become a limitation to the effective counter timing resolution. There are several contributions to this term and each electronic component was studied to understand its effect. From our measurements on the bench and from CLAS12, a reasonable approximation for the floor term in the counter hit time resolution is $\sigma_0 = 40$ ps (see Section 4.2). The PMT anode outputs are connected to JLab-designed VME leading-edge discriminators. A leading-edge rather than a constant-fraction discriminator was chosen for the FTOF system. Although a constant-fraction discriminator delivers better timing initially, off-line time-walk corrections to leading-edge times give comparable results at a significantly lower cost since the off-line analysis can use the measured charge. Time walk is an instrumental shift in the measured hit time that arises due to the finite rise time of the analog pulse. For a given event time, pulses of different amplitude cross the leading-edge discriminator threshold at slightly different times. The time-walk correction algorithm is described in Section 5.2.4. The

Table 6

Key performance specifications of the FTOF CAEN V1190A and VX1290A pipeline TDCs and the JLab FADC250 flash ADCs.

TDC specs (V1190A/VX1290A)	ADC specs
No. Channels: 128/32	16
RMS resolution 100 ps/25 ps	Sampling 250 MHz
Resolution: 19 bit/21 bit	Resolution: 12-bit
Inter-channel isolation ≤ 3 LSB	Clock jitter 350 fs
Double-hit resolution 5 ns	Data memory 8 μ s
Full-scale range 52 μ s	Trigger/Data latency 8 μ s / 32 ns
Integral/Differential non-linearity	
<2.5 LSB / <3 LSB	± 0.5 LSB / ± 0.8 LSB
Inter-channel isolation <3 LSB	SNR 56.8 dB @ 100 MHz input

discriminator threshold was set at -25 mV, significantly above the 2 mV noise level. This threshold corresponds to roughly 1 MeV of deposited energy, consistent with the deposited energy threshold on the FADC readout as mentioned above. The discriminator signal output width was set to 35 ns in order to prevent multiple outputs from the same input pulse.

The output of the discriminator goes to a CAEN VME TDC. Both high resolution TDCs (25 ps LSB CAEN VX1290A) and lower resolution TDCs (100 ps LSB CAEN V1190A) are employed, where the lower resolution TDCs are associated with the longer counters at large polar angles for panel-1a ($N_C = 17 \rightarrow 23$), panel-1b ($N_C = 49 \rightarrow 62$), and panel-2 ($N_C = 1 \rightarrow 5$). These multi-hit pipeline TDCs were chosen in order to allow for readout capability in the operating luminosity of $10^{35} \text{ cm}^{-2} \text{ s}^{-1}$. The TDC readout window was set to 250 ns to ensure the full dynamic range of the data was safely in time with the trigger. The key performance specifications of these TDC units are given in Table 6. Note that the integral non-linearity of each TDC channel in the system was measured and corrected for using a look-up table stored in the system memory. See Ref. [13] for more details.

The PMT dynode outputs are connected to the FADCs for the pulse charge measurement. The readout employs JLab-designed FADC250 16-channel VME 250-MHz flash ADCs [17]. Fig. 9 shows a raw ADC pulse from a representative FTOF PMT. The pedestal is determined event-by-event and subtracted offline. Our procedure determines the pedestal over the first 15 channels. This average is used to determine and subtract off the baseline noise in our pulse signal region, which lies between channels 35 and 65. The measured ADC values for each counter PMT that are referred to in this paper represent the pedestal-subtracted pulse integral in our defined signal region. A pulse fitting algorithm, which fits the leading edge of the pulse down to the baseline, is used to determine the hit time from the FADC signal. The readout window for the FTOF FADCs is set to 48 samples (192 ns). The applied readout threshold is set to 1 MeV to ensure that the hit cluster energy can be determined with a reasonable accuracy. Details on the hit clustering for FTOF are described in Ref. [10]. The key performance specifications of these FADC units are given in Table 6.

The signal cables used for the FTOF system to connect from the PMT anodes and dynodes to the Forward Carriage patch panels are RG-58C/U fire-retardant coaxial cables. This type of cable is appropriate for moderate length cable runs for fast signals with low signal-distortion requirements. The cable runs vary from 47 ft to 59 ft. The connections from the patch panels to the readout electronics are made with a final 5 ft run of low loss RG-174 coaxial cable. The inline signal inverting transformers for the panel-1b dynodes (see Section 4.4) were attached directly to the Forward Carriage patch panels.

4.4.1. High-voltage supplies

The PMTs for the FTOF counters typically operate in the range from 1500 V to 2000 V with negative polarity. The typical dark current drawn by the PMTs is measured to be <20 nA. The system is powered by a single high voltage mainframe for each sector. These mainframes are either CAEN 1527LC or CAEN 4527 units outfitted with negative

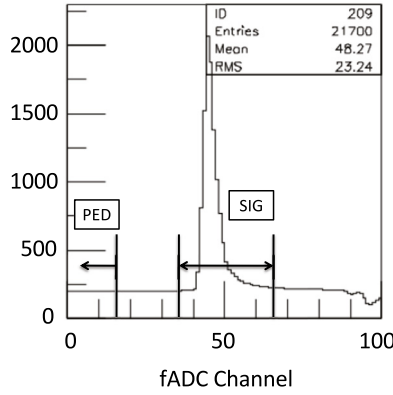


Fig. 9. Typical FADC pulse for a representative FTOF counter from the JLab FADC250. The “PED” region is used to determine the average pedestal in the “SIG” region shown about the PMT pulse. This plot shows ADC counts vs. FADC sample number (1 sample = 4 ns).

polarity 24-channel A1535N modules that can supply up to 3.5 kV per channel with a maximum current of 3 mA. The power supply has a voltage ripple specification of <20 mV peak-to-peak. Each channel consumes less than 1 W during counter operation with typical supply currents per channel between 300 μ A to 500 μ A.

The mainframe is controlled remotely through the Hall B Slow Controls system. A graphical user interface using EPICS [18] running on a UNIX system communicates with the mainframe via Ethernet. The mainframe settings enable basic protection of the PMTs in terms of maximum voltage and current settings, and channel ramp rates.

The high voltage cables for each PMT are fire-retardant RG-59 coaxial cables that run from the PMT voltage divider to a local disconnect HV distribution box located behind the panel-2 arrays in each sector. There are four 48-channel HV distribution boxes for each sector, two for the left PMTs and two for the right PMTs. The output of each HV distribution box is a pair of 35-ft long multi-conductor cables, each containing 24-channels, with a Radiall connector to mate with the HV A1535N board input connector. Each multi-conductor high voltage cable contains individual conductors wrapped in Tefzel insulation, an outer wire shield, and a PVC insulation wrap. Each conductor is rated at 5 kV.

5. FTOF performance

This section highlights the performance of the FTOF system both on the test bench and in Hall B during the first beam runs for CLAS12. The bench test timing performance is important to ensure that the refurbished counters that make up the panel-1a and panel-2 arrays from the CLAS TOF system still meet their original performance specifications as detailed in Table 1 and Ref. [9]. These bench performance studies are even more important for the newly constructed panel-1b arrays for CLAS12 as they are primarily responsible for the limits of the particle identification separation for CLAS12 in the forward direction. Full details on the bench test performance results for the panel-1a and panel-2 counters are provided in Ref. [19] and for the panel-1b counters in Ref. [11].

In this section the essential performance results from the bench testing studies are presented in terms of the counter photoelectron statistics and benchmark time calibrations. Then the calibration algorithms are introduced to provide details on how the in-beam FTOF time resolution performance was quantified. Finally, this section provides the current status of the particle identification capabilities of the FTOF system in relation to the design specifications.

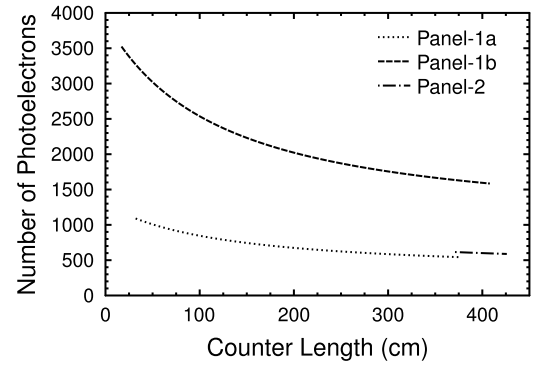


Fig. 10. Parameterized distribution of the number of photoelectrons vs. counter length (cm) for the panel-1a (dotted), panel-1b (dashed), and panel-2 (dot-dashed) counters based on direct measurements with the shortest panel-1a counter.

5.1. Bench measurements

5.1.1. Counter photoelectron statistics

The primary approach to determine the average number of photoelectrons at the photocathode of a PMT generated by minimum-ionizing particles passing through the scintillation bars employs the ratio of the integral of the signal pulse to the integral of the pulse for a single photoelectron [20]. For these measurements we used a 350 MHz (4 GSa/s) oscilloscope with a pulse averaging mode and averaged over 1000 pulses. The minimum-ionizing particle signals were analyzed by connecting the scope to a PMT mounted on one of the shortest FTOF panel-1a counters. For the single photoelectron signal, we took data using just a bare PMT on the bench using the same HV setting. For both measurements the oscilloscope threshold was adjusted appropriately. For the minimum-ionizing peak analysis the threshold had to be set high enough (>200 mV) to eliminate tracks that did not pass through the full thickness of the bar. For the single photoelectron peak the threshold had to be set low enough (1 mV) to pick out the single-electron emission noise pulses from the photocathode that are the dominant source of the PMT intrinsic dark current. This somewhat crude measurement scheme yielded $N_{pe} = 1000 \pm 100$, a value consistent with that found during the initial characterization of the number of photoelectrons seen by the PMTs for these counters for the CLAS TOF system [9]. Based on the parameterization given in Eq. (2), Fig. 10 shows the number of photoelectrons at the PMT photocathodes for the different FTOF counters.

A second method to estimate the number of photoelectrons produced at the PMT photocathode, which accounts for the quantum efficiency at the photocathode, can be estimated from cosmic ray data using the form of Ref. [21]:

$$\langle N_{pe} \rangle = \left(\frac{M_{ADC}}{\sigma_{ADC}} \right)^2, \quad (3)$$

where M_{ADC} is the ADC mean for the minimum-ionizing peak in the ADC spectrum and σ_{ADC} is the width of the ADC distribution. The form of Eq. (3) assumes that a finite σ_{ADC} arises solely due to statistical variations in the number of photoelectrons created at the photocathode for an event sample with a fixed energy loss per track, which we can assume to be a good approximation for perpendicularly incident minimum-ionizing tracks. From the measured data averaged across the counters in panel-1a and panel-1b it was found that $N_{pe}^{1a} = 373 \pm 39$ and $N_{pe}^{1b} = 1158 \pm 77$ [22]. These results, while roughly a factor of two below the parameterized estimates, also show the same factor of three difference in the expected number of photoelectrons for panel-1b relative to panel-1a. The estimates from the first approach are considered to be more reliable not only because they are connected to a more direct measurement of the number of photoelectrons, but also because this parameterization for N_{pe} used in Eq. (1) agrees reasonably well with the measured counter resolutions shown in Section 5.2.7.

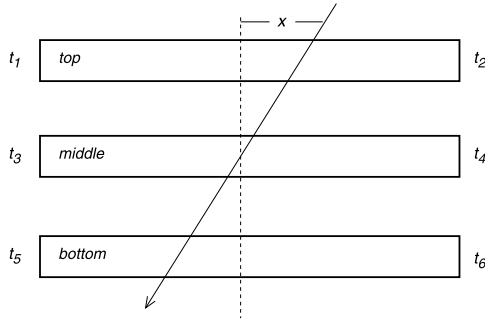


Fig. 11. Schematic representation of a triplet of counters (labeled top - t , middle - m , and bottom - b) with a cosmic ray track traversing the stack. The geometry of the triplet was configured such that the counters were equally spaced.

5.1.2. Bench time resolution performance

The basic algorithm used on the test bench for the panel-1a and panel-2 counters to determine the time resolution of a given reference counter was to use cosmic ray tracks to compare the measured time for a reference counter to the time measured by two other identical counters in a triplet counter configuration (see Fig. 11). For a triplet measurement, where the track passes through all three counters with double-sided readout, six times are measured ($t_1 \rightarrow t_6$). Each time measurement actually represents the difference between the discriminated PMT signal (TDC start) and the trigger time (TDC stop) generated from the six-fold PMT coincidence. These time measurements are then translated into three counter hit times $t_{t,m,b} = \frac{1}{2}(t_{1,3,5} + t_{2,4,6})$.

For incident tracks that pass fully through each counter of the triplet, we can define a time residual

$$t_r = t_m - \frac{1}{2}(t_t + t_b), \quad (4)$$

where the time t_m of the middle scintillator hit should be the average of the measured times t_t and t_b for the top and bottom scintillator hits, respectively. Thus the measured residual t_r should nominally be zero. However, due to the smearing of the measured times t_t , t_m , and t_b due to the finite time resolution of the measurements, the residual time t_r will also be smeared. The width of the t_r distribution can be used to determine the average time resolution of the counters in the triplet.

The average time resolution of each of the identical counters was computed from the variance δt_r of the measured time residual t_r summed over all the full event sample (see Fig. 12 for a representative result). Assuming the average time resolution for each PMT in the triplet is identical and taking into account that each counter is read out using two PMTs, we can write the final expression for the average counter time resolution as:

$$\sigma_{\text{counter}} = \frac{2}{\sqrt{6}} \delta t_r. \quad (5)$$

Thus a measure of the variance of the time residual distribution provides a measure of the average resolution of a counter in the triplet.

Fig. 13 shows the average time resolution measured in the triplet configurations for the panel-1a and panel-2 FTOF counters. For these measurements the fully assembled counter arrays were stacked one above the other in the cosmic ray test stand. The triplets t , m , and b were formed from the counters for sectors 1, 6, and 5 and separately for the counters for sectors 2, 4, and 3. This analysis included a minimum PMT ADC cut to remove events that did not pass through the full thickness of the counter (“corner-clippers”) and also included a coordinate cut of ± 10 cm about the center of the scintillation bar. Due to the use of leading edge discriminators, the measured PMT times were corrected with a power-law time-walk function of the form:

$$t_{\text{walk}}^{L,R} = \frac{A_0}{1 + A_1 \sqrt{(ADC - PED)_{L,R}}}. \quad (6)$$

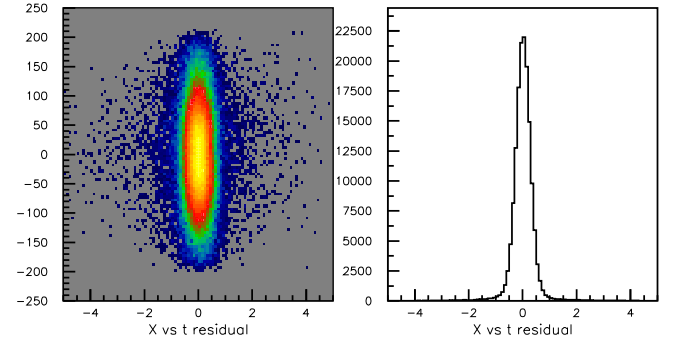


Fig. 12. Data from the S2-S4-S3 panel-1a triplet for counter #23. (Left) Plot of the triplet hit coordinate (cm) (defined from the hit in the middle counter of the triplet) vs. the time residual t_r (ns). (Right) The triplet time residual t_r distribution (ns) that is fit to determine the average counter timing resolution using Eq. (5).

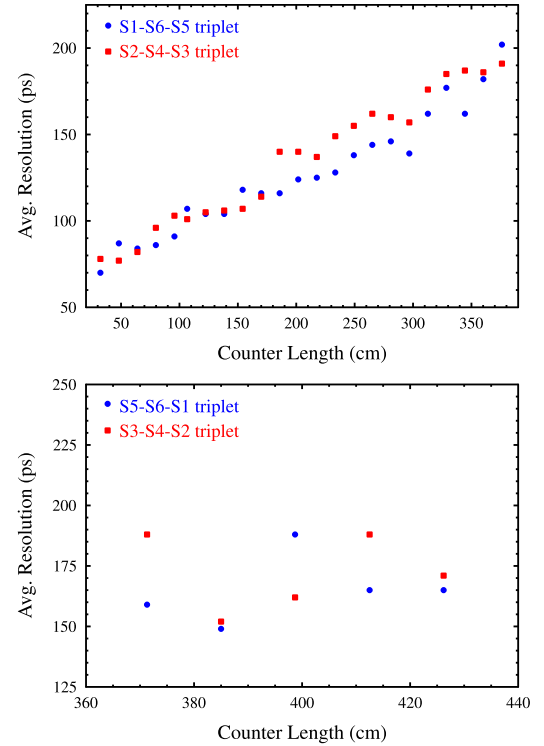


Fig. 13. Average bench measurement resolutions (ps) vs. counter length (cm) using cosmic rays for the refurbished FTOF panel-1a (top) and panel-2 (bottom) counters. The different sets of data points correspond to the two different cosmic ray test stands used for calibration. The data points for the counters in sectors 1, 5, and 6 were averaged together, as well as those for the counters in sectors 2, 3, and 4.

Here, $ADC - PED$ is the pedestal-subtracted ADC value for each PMT (i.e. the pedestal-subtracted integral of the SIG region in Fig. 9). The parameters A_0 and A_1 were determined by fitting the residual times for the left and right PMTs of a given counter from Eq. (4) vs. ADC_L and ADC_R after the PMTs of all counters were gain-matched. See Ref. [19] for full details on these measurements.

The average time resolutions for the counters in the panel-1a and panel-2 triplets were found to be within 20% of those achieved for the original CLAS TOF baseline measurements [9]. However, when accounting for the different TDC LSB contribution to the floor term (25 ps/100 ps for CLAS12 and 50 ps for CLAS) the results agree remarkably well.

The bench measurements for the panel-1b FTOF counters were carried out using a stack of six equidistant counters of a given counter

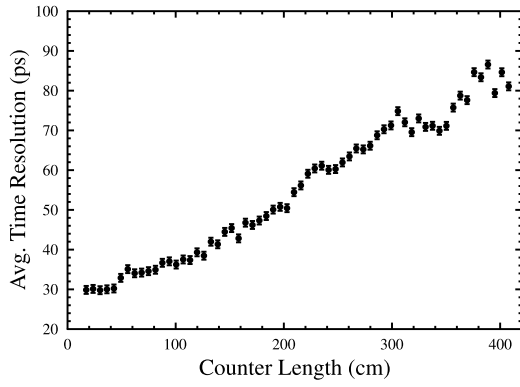


Fig. 14. Measurements of the time resolution (ps) vs. counter length (cm) achieved for the FTOF panel-1b system averaged over the six counters of a given length belonging to each CLAS12 Forward Carriage sector. These data were acquired on the bench using cosmic rays. Full details are included in Ref. [11].

number from $N_C = 1 \rightarrow 62$. Accounting for the necessary path length corrections to relate the individual counter times to each other, six simultaneous triplet counter measurements were analyzed and the resulting time-residual system of equations was then solved for the individual counter time resolution. Also in distinction to the simple time-walk correction employed for the panel-1a and panel-2 measurements shown in Eq. (6), a more sophisticated position-dependent time-walk correction was employed that generalizes the simpler position-independent form. Precision measurements of the time-walk amplitude (A_0 in Eq. (6)) vs. the distance from the PMT showed a nearly linear fall-off of the amplitude with increasing distance from the PMT. On average the time-walk amplitude is $\sim 30\%$ larger at the PMT compared to the far end of the bar, although this near to far end ratio of the amplitude decreases linearly with the length of the bar. Our measurements showed this ratio varies between 20% for the shortest bars to 40% for the longest bars. After accounting for this correction vs. hit position along the bar, a final baseline for the average panel-1b time resolutions was extracted averaging over the six counters of a given number from $N_C = 1 \rightarrow 62$. The resulting time resolutions for the panel-1b counters are shown in Fig. 14 and range from 30 ps for the shortest counters (17 cm long) to 80 ps for the longest counters (408 cm long). Full details describing the measurements are provided in Ref. [11].

For our purposes in quoting counter time resolution values, it is essential to distinguish between two different quantities. The first is the *intrinsic* counter time resolution that reflects the resolution parameterized in Eq. (1). This includes the resolution contributions that mainly depend on the photon statistics at the PMT photocathode and hence the counter geometry, surface quality, scintillation material and bulk quality, wrapping preparations, etc., the transit time spread of the PMT, and the readout electronics noise (the floor term of the resolution). However, when calibrating the counter timing in situ in Hall B with beam interactions in the experimental target, an *effective* time resolution is extracted that includes not only the intrinsic resolution contributions, but also contributions from the angle-dependent uncertainty in the path length determined by the CLAS12 forward tracking system and the resolution spread in the accelerator RF signal that is used as a comparison reference time. The results quoted in this section represent the intrinsic time resolutions of the counters. The effective in situ time resolutions are discussed in Section 5.2.7.

5.2. FTOF beam-data calibrations

In the nominal data taking mode for CLAS12, the FTOF ADCs and TDCs for all PMTs with a signal above the readout threshold are recorded. For the FADCs, the charge of the pulse is integrated over the extent of the pulse region and the pedestal is subtracted event by event

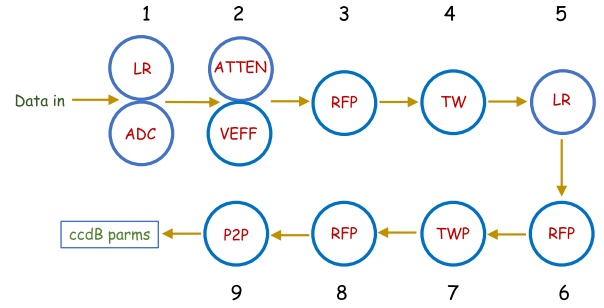


Fig. 15. Schematic representation of the different steps in the FTOF calibration sequence and their order. See Section 5.2 for full details.

during offline data processing as discussed in Section 4.4. For the TDCs the time recorded is relative to the trigger. To determine the flight time of the charged track from the target to the FTOF, the TDC time must be correlated with the time of the electron beam bunch initiating the trigger that is defined by the accelerator radio frequency (RF) pulse. The RF signal from the accelerator has a period of 2.004 ns. The RF bunch length itself corresponds to about 2 ps. Although the signal timing is very accurate (with a resolution of < 20 ps), the determination of which beam bunch produced a given interaction must be determined by the experiment. Note that for the beam operations with CLAS12 the electron beam was actually delivered in every other RF bucket, resulting in an effective T_{RF} of 4.008 ns.

The full calibration of each of the FTOF counters involves a number of discrete steps that are carried out sequentially for a given data run (where a run typically lasts for about two hours of data collection). The associated calibration constants for each run are stored in the CLAS12 calibration database (CCDB) [10]. After the calibration of a given data reference run is completed, the calibrations for subsequent data runs are only carried out if there is a response shift outside of our allowed timing or energy tolerances (which are typically 5%). The steps to complete the FTOF calibration are carried out in a particular sequence as detailed in Ref. [23] and shown schematically in the calibration flowchart of Fig. 15. The individual steps include:

1. Left/right PMT time offsets (LR): This time offset accounts for the difference in the time recorded between the left side and right side PMTs in a given counter due mainly to the different PMT transit times. These time offsets are determined from the centroid of the difference between the left/right TDC time difference and the left/right side hit times computed using the counter hit point from the forward tracking system divided by the effective speed of light in the counter. These time offsets range between ± 5 ns. This step is carried out initially in order to compute a hit coordinate from the FTOF information for the effective velocity determination and then a second time to account for the fact that the time-walk correction shifts the measured left and right PMT times.
2. ADC Calibration (ADC): Determine the ADC value to energy deposition calibration factor for each counter using minimum-ionizing events; see Section 5.2.1.
3. Attenuation Length Calibration (ATTN): This property of the counter quantifies the light absorption length in the scintillation bars and is determined by relating the measured ADC as a function of hit coordinate along the bar; see Section 5.2.2.
4. Effective Velocity Calibration (VEFF): Determine the effective speed of light propagation along the counter; see Section 5.2.3.
5. Time-Walk Amplitude Calibration (TW, TWP): Compare the measured hit time with respect to the measured ADC to determine the time-walk correction; see Section 5.2.4.

6. Counter-to-Counter Time Offset Calibration (RFP, P2P): In order to measure the absolute flight time of a charged particle from the target to the FTOF counter and to be able to reconstruct exclusive events when hits are associated with multiple FTOF counters, the relative time offsets of each counter relative to all of the other counters in the system need to be determined. This is done in two steps. The first step is to correlate each counter hit time to the RF time, which amounts to a precision time alignment in bins of the TDC LSB. The second step is a coarse alignment of each counter hit time in bins of the RF period T_{RF} ; see Section 5.2.5. During this step the effective counter time resolutions are extracted; see Section 5.2.7.
7. TDC Calibration (TDC): After calibrating the integral nonlinearities of each TDC channel in the system (see Section 4.4), the TDC channel to time calibration is completed using beam events; see Section 5.2.6 for details and a note of why this step is not included on Fig. 15 flowchart.

The calibration flowchart of Fig. 15 shows that the calibrations are completed in nine separate calibration steps that proceed in series. The data run is analyzed to complete a given step and the determined parameters are then used in the subsequent steps. Due to dependencies of the steps on each other, several calibration steps (LR, RFP) have to be completed multiple times. As the FTOF calibration relies on accurate path length measurements for the forward-going charged tracks, the drift chamber calibrations are completed before the FTOF calibrations in the overall CLAS12 subsystem calibration sequence.

To calibrate the FTOF system, events are selected that have a good electron reconstructed in the forward direction as determined by the CLAS12 Event Builder (see Ref. [10] for details). From these events the panel-1a and panel-1b counters are calibrated using forward-going charged leptons and pions. For the panel-2 calibration it is necessary to select charged pions and protons, as CLAS12 cannot cleanly identify leptons when there is no electromagnetic calorimeter signal. Due to the increased energy loss and Coulomb multiple scattering of the proton sample, the effective counter time resolutions derived for the panel-2 counters are noticeably worse than for the bench test results using cosmic rays.

The average hit time resolution for the FTOF from the TDCs is about 80 ps and that from the FTOF FADCs, given the rapid rise time of the fast PMT signals that provide for only 2–3 samples on the rising edge, is only about 1 ns. A matching requirement of 10 ns between the TDC time and the FADC time is employed during event reconstruction. While this matching requirement still needs to be tuned further, it is already reasonably effective in allowing the FADC hits to be matched with their corresponding TDC hits. This is important, as due to the slightly different thresholds on the discriminators and the FADCs, the number of entries in the hit lists can be up to a factor of two different. The matching criteria is also essential in order to assign the correct ADC information to the hit not only for the time-walk correction that directly uses the measured ADC, but also for the energy loss computation.

5.2.1. PMT gain matching

One of the purposes of gain matching the FTOF PMTs is to equalize the detector response to tracks that cross the FTOF arrays and pass through two neighboring counters. This is a necessary procedure because each counter must contribute equally to the trigger for a common-threshold discriminator level [8]. Gain matching, so that the minimum-ionizing particle peak response appears at the same ADC value for all counters, also allows for easier data monitoring during online and offline analyses.

The FTOF PMT high voltage settings were determined using calibration runs employing minimum-ionizing tracks. These tracks deposit roughly 10 MeV (12 MeV) as they pass through the 5-cm (6-cm) thick FTOF scintillation bars, as $dE/dx = 2 \text{ MeV/g/cm}^2$ for minimum-ionizing particles. The high voltage settings were initially based on runs

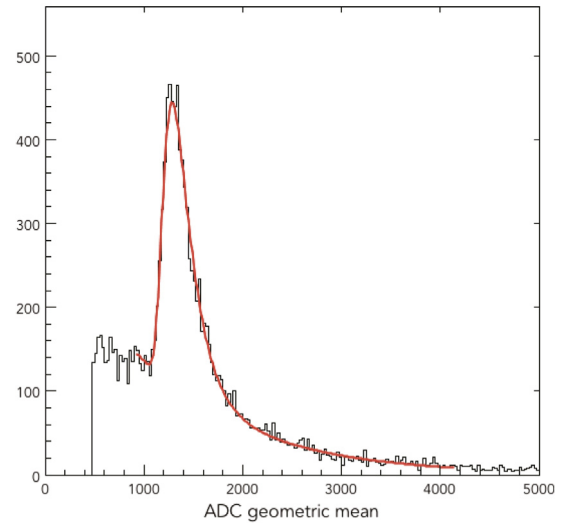


Fig. 16. Geometric ADC mean spectrum for one representative FTOF counter from beam data. The recorded events are pedestal subtracted. The red curve is a fit function that includes a Landau shape for the peak and an exponential for the background.

using cosmic rays with the readout based on a calorimeter pixel trigger (see Ref. [7]) that effectively selects tracks approximately perpendicular to the face of the FTOF counters in panel-1a and panel-1b. Currently the calibrations are carried out using minimum-ionizing tracks from beam data coming from the target. In this case the integral of the ADC pulse is scaled by a path length correction given by t/p , where t is the counter thickness and p is the path-length of the track in the counter as determined by extrapolation of the drift chamber track to the location of the FTOF counter. The energy deposited in the scintillation bars is recorded by the ADCs, which show Landau peaks above the pedestal. Minimum-ionizing tracks that do not pass through the full counter thickness and more heavily ionizing tracks give rise to a background beneath the Landau peak.

In the determination of the high voltage settings, to avoid issues with the attenuation of light for tracks that pass near the ends of the bars and with unbalanced light entering the left and right PMTs, we combine the information from the left and right PMTs to produce an average ADC spectrum for the counter through the quantity known as the geometric ADC mean:

$$\overline{ADC} = \sqrt{(ADC - PED)_L \cdot (ADC - PED)_R}. \quad (7)$$

Given the finite dynamic range of the ADC, we have chosen to position the minimum-ionizing peak in a particular ADC value that is different for the panel-1a, panel-1b, and panel-2 counters. For all counters this value is selected so that it is safely above the pedestal, but leaves sufficient range for the more highly ionizing charged tracks of our typical physics events. To minimize PMT aging effects that result in loss of PMT gain with time correlated with the total charge collected at the anode of the PMT, the gains are set as low as possible.

The position of the minimum-ionizing peak in the \overline{ADC} spectrum is set by the PMT HV values. For a given scintillation bar, a typical \overline{ADC} spectrum is shown in Fig. 16. Given that the same ADC geometric mean value is chosen to position the minimum-ionizing peak for all counters in either panel-1a or panel-1b, the required PMT gain increases linearly from the short to the long bars to compensate for the attenuation losses in the longer bars.

The gain G of a PMT can be related to the high voltage setting V using $G \propto V^\alpha$, which represents a basic power law form with α the power law factor. This expression governs the gain change for a given change in voltage. With this expression, the PMT gain G_1 at a given

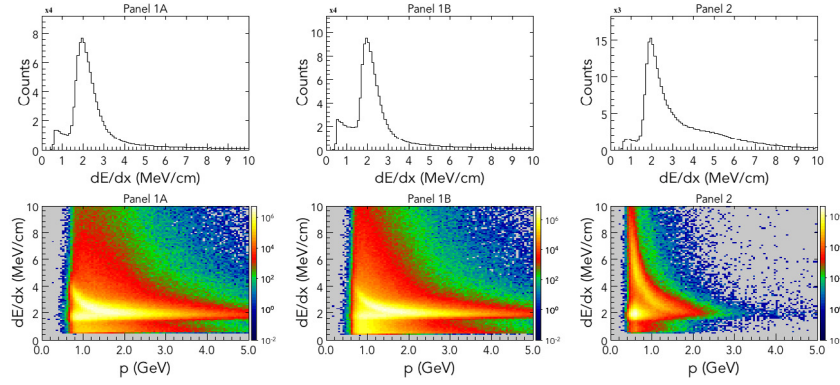


Fig. 17. Measured FTOF counter energy loss for positively charged particles from 10.6-GeV electrons incident upon a liquid-hydrogen target normalized by the extrapolated path length from the projection of the forward track through the counter array. The normalized dE/dx (MeV/cm) spectrum shows the separation of minimum-ionizing particles from more heavily ionizing particles summed over the counters in panel-1a (left), panel-1b (middle), and panel-2 (left). The top row of plots show the normalized dE/dx and the bottom row of plots show dE/dx vs. track momentum (GeV).

voltage V_1 can be related to the gain G_2 at a different voltage V_2 using:

$$\frac{G_1}{G_2} = \left(\frac{V_1}{V_2} \right)^\alpha, \quad (8)$$

which can be rewritten in a slightly different form as:

$$\frac{\Delta G}{G} = \alpha \frac{\Delta V}{V}. \quad (9)$$

For our purposes we assume that the gain G is directly proportional to the measured ADC value. With an expression that relates the measured ADC value at two different voltage settings, we have a relation that forms the basis for relating the position of the minimum-ionizing peak in the \overline{ADC} spectrum (see Eq. (7)) to the PMT HV setting. The gain-matching procedure then amounts to adjusting the HV settings of all PMTs to the values required to position the minimum-ionizing peak for each counter in the desired ADC location. At the same time the algorithm uses the individual left and right PMT ADC spectra for a given counter to ensure that the PMT gains for any given counter are balanced.

The power law factor α in Eq. (8) for each PMT type can be determined by looking at data with two different high voltage settings. In this manner the average α factors for the FTOF PMTs were determined to be 13.4 for panel-1a, 4.7 for panel-1b, and 8.6 for panel-2. With these values the calibrations converge within just a few iterations such that all of the minimum-ionizing particle peak locations are within \overline{ADC} values of ± 25 of their set targets and the left and right PMT ADC values are similarly gain matched.

The gain matching procedure is carried out before the beginning of each experiment to determine the high voltage settings for the PMTs. The PMT response is monitored throughout the run period. If the average PMT gains shift by more than $\sim 5\%$, new high voltage settings are determined to optimize the gain balance and to restore the ADC geometric means to their nominal values. In the first two years of operations of FTOF in CLAS12, the high voltage settings have been updated after every four to six weeks of beam operations.

The energy loss in a counter for a passing charged particle track is determined after the minimum-ionizing peak centroids are aligned. The energy loss in each counter is computed for each PMT as:

$$E_{L,R} = ADC_{L,R} \cdot \left[\frac{\left(\frac{dE}{dx} \right)_{MIP} \cdot t}{ADC_{MIP}} \right] \exp \left(\frac{d_{L,R}}{\lambda} \right), \quad (10)$$

where ADC_{MIP} is the centroid of the minimum-ionizing peak in the geometric mean distribution, $\left(\frac{dE}{dx} \right)_{MIP}$ is the energy loss for minimum-ionizing particles in the scintillation bars (2.001 MeV/cm), t is the counter thickness ($t = 5$ cm for panel-1a and panel-2, and $t = 6$ cm for panel-1b), d is the distance along the bar from the track hit position

to the PMT, and λ is the counter attenuation length. The energy loss used in the event reconstruction is the geometric mean of the separate measures $E_{L,R}$ (see Ref. [10] for details).

Fig. 17 shows the reconstructed energy loss normalized by the track path length through the bar for different panels from a data run with a 10.6-GeV electron beam incident upon a liquid-hydrogen target. The path length through the bar is determined from extrapolating the track from the forward tracking system through the FTOF system and determining the track entrance and exit points on each scintillation bar. The data allow for the separation of minimum-ionizing particles from more heavily ionizing particles. The minimum-ionizing particles lose a constant amount of energy as a function of path length. At low momentum the more heavily ionizing particles have energy loss that increases linearly with distance until they can pass through the counter. At that point their energy loss follows the Bethe–Bloch formula. However, the minimum proton track momenta seen by the FTOF is more than 0.5 GeV, so no protons are actually stopped in the FTOF scintillation bars.

5.2.2. Attenuation length measurements

The measured ADC values for each PMT can be written in terms of the attenuation length as:

$$(ADC - PED) = A_0 e^{-d/\lambda}, \quad (11)$$

where A_0 is a constant, d is the distance along the counter with respect to the PMT location, and λ is the counter attenuation length. Using the definition:

$$d_{L/R} = \frac{L}{2} \pm coor, \quad (12)$$

where L is the counter length, $coor$ is the FTOF hit coordinate along the bar (with the middle of the bar at $coor = 0$) defined as:

$$coor = \frac{v_{eff}}{2} \cdot (t_L - t_R - C_{LR}), \quad (13)$$

v_{eff} is the effective velocity of light in the scintillation bars (see Section 5.2.3 for details), and C_{LR} is the offset that centers the time difference distribution about 0.

The logarithmic ratio of the ADCs of the left and right PMTs from a given counter as a function of hit coordinate along the bar can be written as:

$$\log \left(\frac{(ADC - PED)_R}{(ADC - PED)_L} \right) = C + \frac{2 \cdot coor}{\lambda}. \quad (14)$$

This expression can be used to determine the effective counter attenuation length using a linear fit of the logarithmic ADC ratio vs. $coor$. The slope of this correlation is $2/\lambda$. In this expression, the y -intercept C is a constant given by $\log(A_0^R/A_0^L)$.

Fig. 18 shows the measured attenuation lengths for the FTOF counters in one sector of the CLAS12 Forward Detector extracted from data with a 10.6 GeV electron beam incident upon a liquid-hydrogen target.

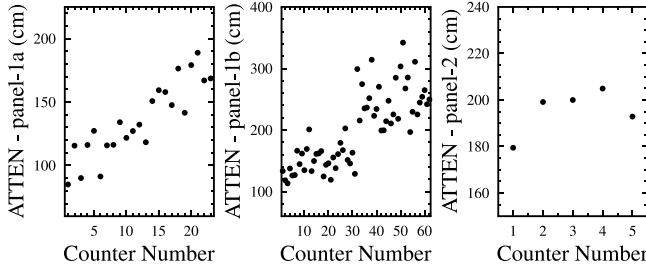


Fig. 18. Counter attenuation lengths (cm) vs. counter number for all FTOF counters in one sector of the CLAS12 Forward Detector determined from beam data.

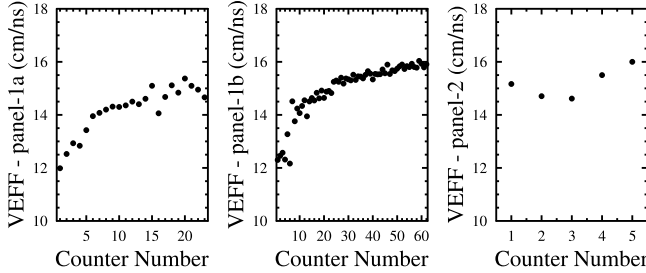


Fig. 19. Counter effective velocities (cm/ns) vs. counter number for all FTOF counters in one sector of the CLAS12 Forward Carriage determined from beam data.

5.2.3. Effective velocity determination

The effective velocity of light in each counter employs a calculation based on the comparison of the reconstructed coordinate information along the scintillation bar from the left and right PMT TDC times with the track hit coordinate determined from the extrapolation of the track beyond the drift chambers to the location of the FTOF counters. Fig. 19 shows the measured effective velocity for each counter in one sector of the CLAS12 Forward Detector using data with a 10.6-GeV electron beam incident on a liquid-hydrogen target.

As the counter length increases, so does v_{eff} because more reflected light with smaller v_{eff} is lost due to attenuation along the bar compared to direct light with higher v_{eff} . The effective velocity is used in the FTOF analysis to determine the hit time for each event from the measured TDC times. The intrinsic position resolution is given by $v_{eff} \times \sigma(t_L - t_R)$ for each counter, which is most relevant for the interactions of neutral particles. The position for charged particles at the location of the FTOF counters can be measured more precisely with the forward tracking system (the track hit coordinate resolution at the FTOF counters is 1–2 cm, while the FTOF hit coordinate resolution is 1–4 cm, depending on the counter timing resolution). However, the hit coordinate of charged particles along the length of the FTOF counters determined from FTOF information alone (see Eq. (13)) can be compared to that from forward tracking projected to the FTOF location to ensure the FTOF hits are properly matched to the tracks.

5.2.4. Time-walk corrections

The approach that we have adopted to correct the FTOF TDC times for time-walk effects is different from the one employed for our bench test studies of the counters in their cosmic ray test stands described in Section 5.1.2. We ultimately settled on an approach that first accounts for an average hit position-independent correction (called TW in Fig. 15) with a power-law functional form:

$$t_{L,R}^{corr} = t_{L,R} - \frac{tw0_{L,R}}{\sqrt{(ADC - PED)_{L,R}}}. \quad (15)$$

The time-walk amplitude parameters $tw0_{L,R}$ are determined by defining the following vertex time residuals for each PMT:

$$t_{L,R}^{res} = \left(t_{L,R} - \frac{d_{L,R}^{FT}}{v_{eff}} - \frac{P_L}{\beta c} \right) - \left(t_{RF} + \frac{z_{vert}}{\beta_e c} \right), \quad (16)$$

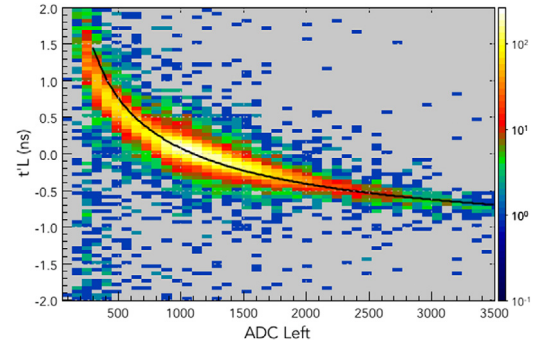


Fig. 20. Plot of t' (ns) vs. ADC for one representative left PMT from panel-1b from beam data. As t' is defined using the modulus of T_{RF} , its limits span $\pm T_{RF}/2$. The overlaid curve represents the time-walk functional fit from Eq. (15).

where $t_{L,R}$ are the measured TDC times after the left/right PMT time offset correction, $d_{L,R}^{FT}/v_{eff}$ corrects the time measured at the PMT to the time at the track hit point on the counter determined from the forward tracking information, and $P_L/(\beta c)$ is the track flight time from the reaction vertex to the FTOF. The track path length P_L and β are defined via forward tracking and the particle identification is determined from the Event Builder [13]. The term $z_{vert}/(\beta_e c)$ corrects the RF time t_{RF} for the actual electron beam event vertex location along the z-axis of the extended target. This vertex time residual represents the FTOF hit time from a single PMT traced back to the reaction vertex and compared to a precise time reference t_{RF} given by the RF signal from the accelerator. As t_{RF} represents a reference time for the arrival of the electron beam bunch at a fixed position along the beamline in Hall B assigned as the center of the target, the time must be corrected for the displacement of the reaction vertex along the extended length of the target.

As the beam bucket that was associated with the event is not determined at this point, the time walk for each PMT is actually determined using the modulus of $t_{L,R}^{res}$ with the RF beam bucket period $T_{RF} = 1/(\text{RF frequency})$ by fitting:

$$t'_{L,R} = \text{mod} \left[t_{L,R}^{res}, T_{RF} \right] \text{ vs. } (ADC - PED)_{L,R}. \quad (17)$$

Fig. 20 shows the t' vs. ADC distribution for a representative FTOF PMT in panel-1b from beam data using a 10.6-GeV electron beam incident upon a 5-cm long liquid-hydrogen target. Note that all distributions that employ $t'_{L,R}$ are sorted in 25 ps bins (consistent with the TDC LSB). The overall scale of the time-walk effects spanning the full dynamic range of the ADC is 2 ns.

The second part of the full time-walk correction accounts for additional position-dependent effects as discussed in Section 5.1.2 (called TWP in Fig. 15). After the position-independent time-walk correction is determined, in a second step we then fit a second-order polynomial to the counter hit time vs. hit position along the counter defined from the projection of the charged particle track on the FTOF from the forward tracking system. The time employed for this step is the track hit time at the vertex (relative to the vertex-corrected RF time) averaging the left and right PMT hit times (see definition in Section 5.2.5). Fig. 21 shows the distribution before and after the second correction. The before distribution where only the position-independent time-walk correction is applied, reveals a characteristic “smile” pattern, which reflects the unaccounted for position-dependent time-walk effects on the measured times incorporating both the left and right PMTs each with their own linearly falling time-walk parameters when moving away from each PMT as discussed in Section 5.1.2. Effectively this approach actually accounts for all remaining position dependences in the calibration parameters. Specifically, it also takes care of the effective velocity changes with position along the bar moving away from the PMT. However,

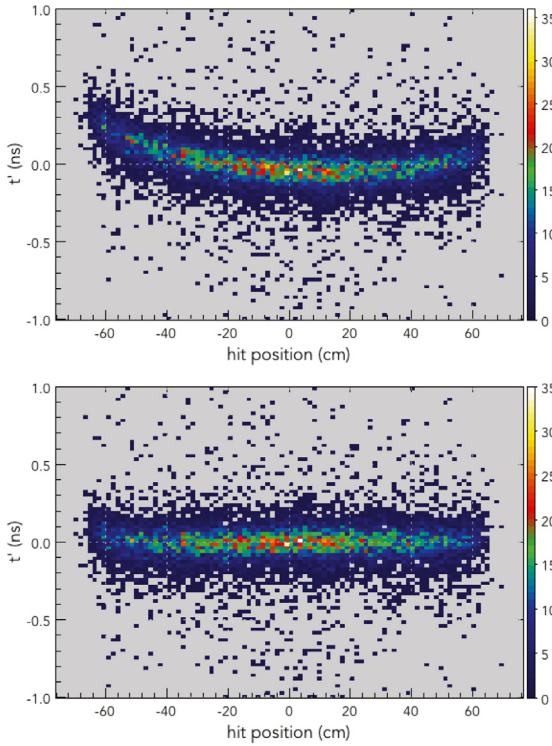


Fig. 21. Plot of t' (ns) vs. hit position (cm) along the bar from beam data after the position-independent time-walk correction (top) and after the ad hoc second-order polynomial correction to remove the residual coordinate dependence (bottom).

the dominant position-dependence accounted for here is related to the time-walk.

Note that the RFP fine-timing calibration step (see Section 5.2.5) is completed three separate times in the calibration sequence. The first time is to center the vertex time distributions in the T_{RF} window to avoid wrap-around in the RF period range when taking the modulus of the vertex time difference with the RF time. The RFP calibration step is repeated after the second LR step due to shifts of the left and right PMT times and before the position-dependent time-walk calibration (TWP). A final RFP calibration is carried out to account for the vertex time shifts that result from the position-dependent time-walk corrections.

5.2.5. Counter-to-counter time alignment

The flight time of a charged particle from the reaction vertex to an FTOF counter is given by:

$$t_p = \bar{t}_{hit} - t_{ST}, \quad (18)$$

where \bar{t}_{hit} is the average FTOF counter hit time from the left and right PMTs (see Section 5.2.8) and t_{ST} is the event start time. The event start time is associated with the RF but needs to be synchronized with the particular RF beam bucket associated with the event. The beam bunch width within the RF beam bucket is only about 2 ps and, therefore, represents a precise time marker. However, as the RF time signal has a period of T_{RF} , it is not a priori known which RF beam bucket is the one associated with the event that led to the hit in the FTOF counter.

The determination of the absolute flight time of charged particle tracks from the reaction vertex to an FTOF counter is performed in two steps. In the first step (called RFP in Fig. 15), fine timing offsets (binned in the 25 ps TDC LSB) are determined to align the FTOF hit times traced back to the reaction vertex for each counter within the RF time window. In the second step (called P2P in Fig. 15), coarse timing offsets binned in units of the RF period T_{RF} are determined to select the specific RF beam bucket associated with the event.

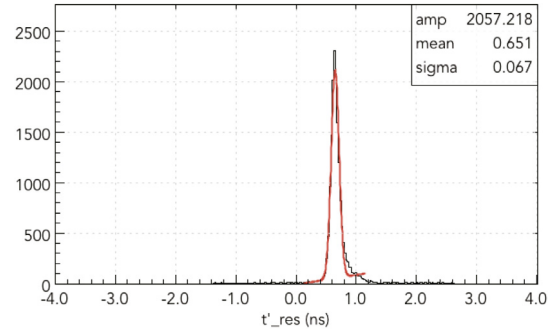


Fig. 22. Distribution of the FTOF hit times from beam data traced back to the vertex relative to the RF (ns) for one representative FTOF panel-1b counter with the Gaussian plus background fit overlaid to determine the counter RF offset and the effective counter time resolution. As t'_{res} is defined using the modulus of T_{RF} , this distribution is limited to span $\pm T_{RF}/2$ about the peak centroid.

The fine timing alignment algorithm uses the FTOF hit time traced back to the event vertex relative to the RF to align the vertex times of all FTOF hits (modulo T_{RF}). However, instead of using the separate left and right PMT hit times as in Eq. (16), this algorithm uses the average counter hit times,

$$t'_{res} = \text{mod} [t_{vtx}, T_{RF}], \quad (19)$$

$$t_{vtx} = \left(\bar{t}_{hit} - \frac{P_L}{\beta c} \right) - \left(t_{RF} + \frac{z_{vert}}{\beta c} \right).$$

Fig. 22 shows the t'_{res} distribution for one representative FTOF counter. The centroid of the Gaussian fit gives the fine timing offset. The width of the Gaussian fit is a measure of the effective time resolution of the counter. To display the full t'_{res} distribution avoiding any wrap-around effects near the edges of the T_{RF} range, the algorithm plots the t'_{res} distribution in a range of $\pm T_{RF}/2$ about the peak channel in the distribution.

After the fine timing offset calibration, the counter timing is precisely aligned modulo T_{RF} . The next step in the FTOF timing calibration is to fix the measured hit times for all counters to the specific RF bunch associated with the event. This is carried out using coincidences of charged particle tracks to link the hit times of all counters across the full FTOF system.

The coarse timing offset algorithm (called P2P for paddle-to-paddle) selects events with two forward-going charged tracks and computes the vertex time difference between any given FTOF counter relative to hits in all of the other FTOF counters,

$$t_{P2P} = t_{vert}^1 - t_{vert}^2, \quad (20)$$

where,

$$t_{vert}^i = \bar{t}_{hit}^i - \frac{P_L^i}{\beta c}. \quad (21)$$

At this point the counter timing has already been aligned to within a multiple of T_{RF} . Note that particle identification of each track is given by the Event Builder [13], and as both tracks are assumed to originate from the same reaction vertex, no vertex time corrections are necessary. The algorithm adjusts the vertex time differences between all counters to zero. The coarse time offsets represents a single parameter for each counter that is restricted to values of $n \cdot T_{RF}$, with $n = 0, \pm 1, \pm 2, \dots$

Fig. 23 shows the t_{P2P} distribution for one representative FTOF counter before and after the coarse timing alignment. As expected, the final histogram is dominated by events in a single channel (of width T_{RF}) centered at $T_{RF} = 0$. As these constants are predominantly determined by the fixed system cable lengths, of which there are four different lengths used to connect the panel-1a and the panel-1b counters, the constants primarily reflect the differences in the signal propagation times along the signal cables. Note that the algorithm specifically identifies two track events for the calibration and does not consider hits in panel-1b and panel-1a associated with the same track.

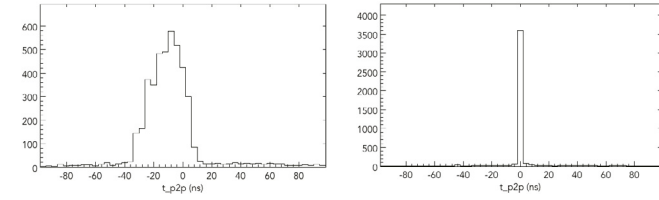


Fig. 23. Distribution of the vertex time differences (ns) for tracks in a single representative FTOF counter compared to tracks in all other FTOF counters using event samples with two forward-going charged particle tracks. (Left) Before P2P corrections and (right) after P2P corrections. The histogram is sorted in bins of T_{RF} .

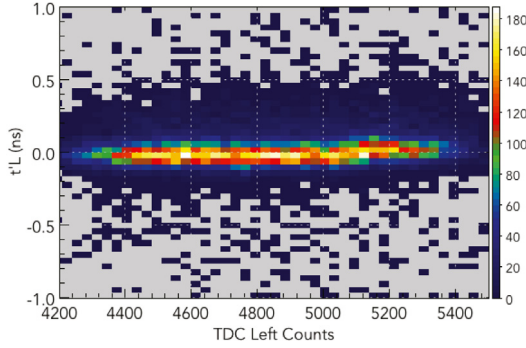


Fig. 24. Distribution of t_L^{res} (ns) vs. TDC digitization bin for one representative left FTOF panel-1b PMT. The TDC conversion constant for each input channel is that which forces the slope of the linear fit to be zero.

5.2.6. TDC calibration

The calibration sequence also allows for calibration of the TDCs. This calibration is a single constant for each TDC input channel in the system that converts the measured TDC digitization bin into time. The nominal TDC LSB is 25 ps for the CAEN VX1290A and V1190A TDC units employed for the FTOF readout (see Section 4.4).

The calibration is completed by fitting the PMT time residuals of Eq. (16) vs. TDC digitization bin using a linear function. The TDC calibration is the value that fixes the slope of t_{res}' to be zero. Fig. 24 shows the distribution of t_{res}' vs. TDC for a representative FTOF counter. Any bin-to-bin Δt variations reflect remaining integral non-linearities in the measured TDC compensation tables (see Section 4.4). At the present time a single conversion constant of $CONV = 23.45$ ps/bin is employed for all FTOF system TDC input channels. This value is derived using a TDC channel that digitizes the RF time (our most accurate time reference in Hall B). For this reason individual TDC channel calibrations are not shown as a separate step in Fig. 15 and the channel-by-channel TDC calibrations done for FTOF serve only as a cross-check.

5.2.7. Counter time resolutions

The effective time resolutions for each counter determined after complete calibrations of the FTOF system are shown in Fig. 25. These measurements are from a beam data run with 10.6-GeV electrons incident on a 5-cm long liquid-hydrogen target. These time resolutions represent the current quality of the overall CLAS12 calibrations. The results are based on calibration procedures that are not yet fully optimized, as well as uncertainties in the reconstructed momentum and path length from the forward track reconstruction. Note that the time resolution floor-term σ_0 discussed in Section 4.2 and Eq. (1) does not include the contributions from the reconstructed path length uncertainties. These uncertainties are polar and azimuthal angle dependent. Near the torus coils the true magnetic field has different variations than accounted for in our conductor model used to generate the field for the event reconstruction. Furthermore, the path length uncertainties grow strongly for high momentum tracks at small angles, which represent

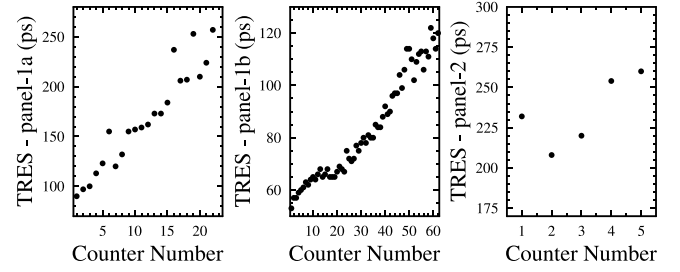


Fig. 25. The measured effective time resolution (ps) vs. counter number for each of the FTOF counters in sector 1 as determined using final state leptons and pions for panel-1a and panel-1b and pions and protons for panel-2.

the dominant part of our kinematic phase space at 10.6 GeV. It is also important to mention that studies of the CLAS12 subsystem detector alignment based on survey data and based on zero-field straight track data are in progress. Misalignments of the detector affect the quality and accuracy of the reconstruction. When all of these uncertainties and misalignments are accounted for their contribution to the floor-term of the resolution function will be reduced.

Nevertheless, the time resolutions already achieved meet the system design specifications in the forward direction outlined in Section 3 and shown in Table 1. For the panel-2 counters the time resolutions are 200–250 ps, but the calibrations are presently limited by statistics, by the use of low-momentum protons in the calibration sample (as discussed in Section 5.2), and by the use of the low-resolution 100 ps LSB TDCs for readout. With these FTOF counter resolutions, the quality of the particle identification in the Forward Detector of CLAS12 allows the experimental program in Hall B to reach its goals. As further operating experience with CLAS12 is gained, we expect to realize further modest but important improvements in the FTOF time resolution that will allow π/K , π/p , and K/p separation in the Forward Detector of CLAS12 to be pushed to higher momenta than currently seen.

5.2.8. Counter hit times

After completion of each of the timing calibration steps, the FTOF hit time associated with a matched charged particle track can be determined. Putting all of the timing corrections together, the track hit time reconstructed from the readout of the left and right PMTs are given by:

$$t_{L,R} = (CONV \cdot TDC_{L,R}) - t_{L,R}^{walk} \mp \frac{C_{LR}}{2} + C_{RF} + C_{p2p}, \quad (22)$$

where $CONV$ is the TDC digitization bin to time conversion factor, TDC is the measured TDC value relative to the trigger signal, t_{walk} is the time-walk correction (which includes both the position-independent and position-dependent time-walk corrections described in Section 5.2.4), C_{LR} is the time shift to center the TDC difference distribution relative to the track coordinate at zero, and C_{RF} and C_{p2p} are the time shifts to align all of the counter times with respect to the RF and to each other, respectively.

The actual hit time associated with the track has to be corrected for the propagation time of the light from the track hit point on the counter to the PMT. The final reported track hit time is then the average of the left and right corrected PMT time. Another aspect of the FTOF hit reconstruction is associated with tracks that cross through multiple scintillation bars as they pass through the FTOF system. These are referred to as hit clusters. If a track passes through both panel-1b and panel-1a, the hit time from panel-1a can be evolved back to the panel-1b hit location and the time information can be combined to give a hit time with improved precision. Full details on the FTOF reconstruction algorithms, including hit times, and the hit clustering and matching algorithms are provided in Refs. [10,24].

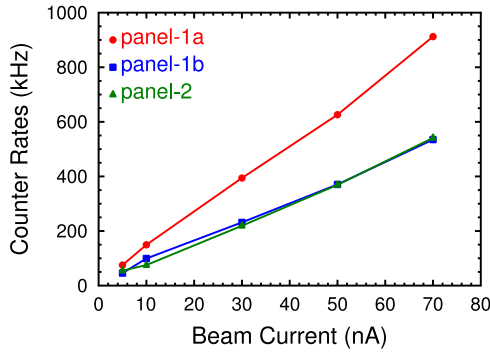


Fig. 26. Measured FTOF counter rates (kHz) for 10.6-GeV electrons on a liquid-hydrogen target as a function of beam current (nA) employing hardware thresholds of approximately 1 MeV. The nominal CLAS12 operating luminosity of $1 \times 10^{35} \text{ cm}^{-2} \text{ s}^{-1}$ corresponds to a beam current of $\sim 75 \text{ nA}$. The red circles correspond to the average panel-1a counter rates, the blue squares to the average panel-1b counter rates, and the green triangles to the average panel-2 counter rates.

5.3. Beam performance

The first in-beam characterization of the FTOF system took place during the Dec. 2017 to Feb. 2018 CLAS12 Engineering Run and subsequently during the first physics production running periods that took place from Mar. - May 2018 and Sep. - Dec. 2018. During these periods the performance of the FTOF system was tested at different beam energies (2.2, 6.5, 7.5, 10.6 GeV), different torus and solenoid magnetic field strengths and polarities (from 0 field to full field for both magnets), and over a range of beam-target luminosities up to the nominal planned CLAS12 luminosity of $1 \times 10^{35} \text{ cm}^{-2} \text{ s}^{-1}$. In this section the measured scaler rates and PMT currents as a function of beam current are presented, as well as the reconstruction results and particle identification capabilities relative to the system specifications based on the current system calibrations.

5.3.1. FTOF rates and PMT currents

The count rates during beam operations can be viewed during data taking using the scalers associated with the discriminators or with the FADCs. The threshold applied for these scalers are set at 1 MeV. During a beam current scan from 5 nA to 70 nA with a 10.6-GeV electron beam incident upon the 5-cm long liquid-hydrogen target (where 75 nA corresponds roughly to the nominal CLAS12 design luminosity of $1 \times 10^{35} \text{ cm}^{-2} \text{ s}^{-1}$) the average count rate in the different FTOF counters was studied. Averaged over the three different arrays, the results shown in Fig. 26, display a reasonably linear behavior. The rates in panel-1a are about a factor of two larger than those for panel-1b. This is in agreement with the fact that the panel-1a counters are 2.5 times wider than the panel-1b counters. However, some portion of the incident radiation is absorbed in the panel-1b counters reducing the flux seen in panel-1a. At the nominal luminosity of CLAS12, the average measured rates in the panel-1b counters are about 500 kHz and those in panel-1a are about 1 MHz.

The response of the FTOF with an 11 GeV electron beam incident upon a 5-cm liquid-hydrogen target has been studied with a detailed simulation of the full CLAS12 detector and beamline based on our Geant4 Monte Carlo called GEMC [25]. Shown in Fig. 27 are the overall rates associated with hits above the readout threshold, including contributions from photons, neutrons, and charged particles. By far the dominant contribution to the overall measured FTOF rate is associated with low energy photons, whose energy deposition in the counters is significantly less than the contribution from minimum-ionizing hadrons. The results from Fig. 26 (data) and Fig. 27 (Monte Carlo) agree to within a factor of two. However, the measured rates depend very strongly on the hardware thresholds. The threshold set for

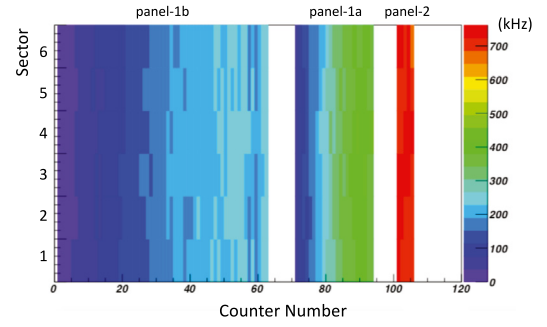


Fig. 27. Simulation results for the FTOF counter rates (kHz) for each sector for 11 GeV electrons on a liquid-hydrogen target at the nominal CLAS12 design luminosity of $1 \times 10^{35} \text{ cm}^{-2} \text{ s}^{-1}$. The rates are calculated for a 1-MeV deposited-energy threshold and expressed in kHz. Here the left-most group of 62 paddles corresponds to panel-1b, the middle group of 23 paddles corresponds to panel-1a, and the right-most group of 5 paddles corresponds to panel-2.

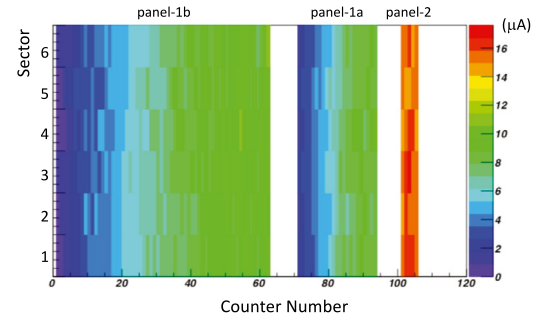


Fig. 28. Computed PMT currents (μA) from Monte Carlo studies for each sector at the nominal operating luminosity of CLAS12 of $1 \times 10^{35} \text{ cm}^{-2} \text{ s}^{-1}$ with 11 GeV electrons incident upon a liquid-hydrogen target. Here the left-most group of 62 paddles corresponds to panel-1b, the middle group of 23 paddles corresponds to panel-1a, and the right-most group of 5 paddles corresponds to panel-2.

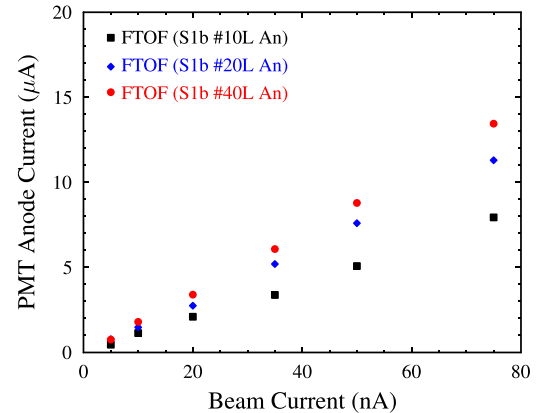


Fig. 29. Measurements of the PMT anode current for representative panel-1b left-side PMTs ($N_c = 10, 20, 40$) as a function of beam current with a 10.6-GeV electron beam incident upon a 5-cm long liquid-hydrogen target.

the FADC readout and on the discriminator are not set directly on the deposited energy, but on signal pulse height in the case of the discriminators or on a digitized amplitude in the FADC. The set threshold on the hardware mentioned in Section 4.4 are only approximately set to 1 MeV.

The average PMT current is directly proportional to the average number of photoelectrons ($\langle N_{pe} \rangle$) created at the photocathode by the scintillation light and the average incident charged particle event rate

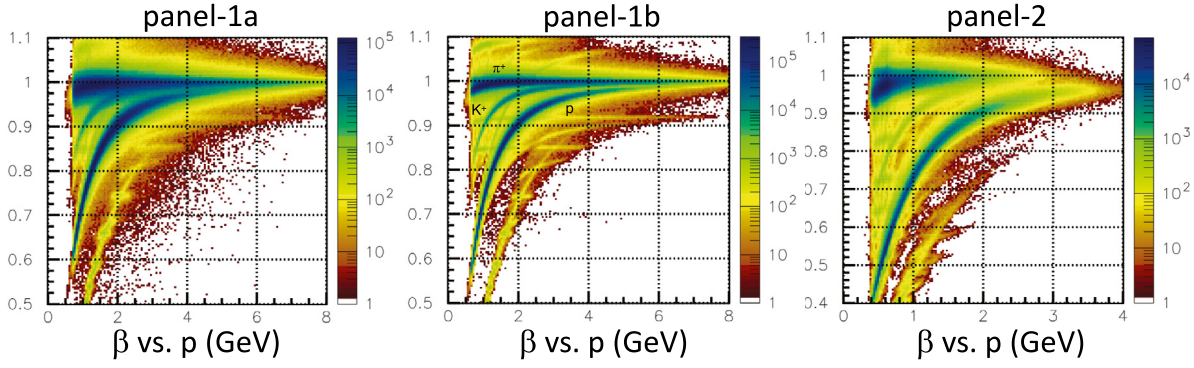


Fig. 30. Velocity of positive hadrons (β) vs. momentum (GeV) for all counters in panel-1a (left), panel-1b (middle), and panel-2 (right) from beam data with a 10.6 GeV electron beam incident on a liquid-hydrogen target.

$\langle R \rangle$. This current can be expressed as:

$$\langle i_{PMT} \rangle = \langle N_{pe} \rangle \cdot Q_e \cdot G \cdot \langle R \rangle, \quad (23)$$

where $Q_e = 1.6 \times 10^{-19}$ C/e is the electron charge and G is the PMT gain (assumed to be 1×10^6). Using the photoelectron statistics estimated in Section 4.2, Fig. 28 shows the predictions of the PMT anode currents for all of the FTOF counters at nominal operating luminosity at 11 GeV from our detailed Geant4 Monte Carlo studies [26]. These predictions show typical PMT anode currents in the panel-1a and panel-1b PMTs at the level of 5 μ A to 10 μ A increasing linearly with counter length. Direct in-beam measurements of the PMT anode currents (made connecting a picoammeter to the PMT anode output) are shown in Fig. 29. The in-beam measurements are in good accord with the simulation expectations. Note that the measured anode currents reflect the current integrated over all incident particles down to zero threshold. The comparison of data to Monte Carlo avoids the issues with the threshold uncertainty mentioned above when comparing scaler rates between data and Monte Carlo and serves to validate the Monte Carlo simulation and its modeling of the FTOF.

5.3.2. Reconstruction results

Particle identification in the Forward Detector of CLAS12 relies heavily on the combination of the measured charged particle momentum from the forward tracking system and the flight time from the target to the FTOF system. The vertex time is determined with respect to the accelerator RF, modulo the RF period T_{RF} . The beam bunch for each event is identified using the flight time of scattered electrons or high-momentum pions, traced back to the interaction point. The FTOF resolution of <200 ps allows clear selection of the correct RF beam bucket.

A plot of relative velocity (β) vs. momentum is shown in Fig. 30 for positively charged particles for the data taken with a 10.6-GeV electron beam incident upon a 5-cm long liquid-hydrogen target. These reconstructed beam data were based on initial calibrations of both the forward tracking system and the FTOF. Here the distributions are presented separately for the counters in panel-1a, panel-1b, and panel-2 summed over all CLAS12 Forward Detector sectors. Each distribution clearly shows primary bands corresponding to π^+ , K^+ , and p (see particle type labels on the panel-1b plot). These distributions show evidence of other bands that correspond to so-called “accidentals”. These are mainly due to particles from neighboring RF beam buckets separated by $\pm n T_{RF}$, where n is an integer. These distributions demonstrate the particle separation limits for π/K , π/p , and K/p vs. momentum.

To connect the current state of the particle identification limits from the FTOF system with those detailed in Section 3, Fig. 31 shows the reconstructed mass for positively charged particles in the Forward Detector of CLAS12 based on initial calibrations of the FTOF system. These plots are based only on the hit time reconstructed from the panel-1b counters. In order to avoid timing resolution effects that would

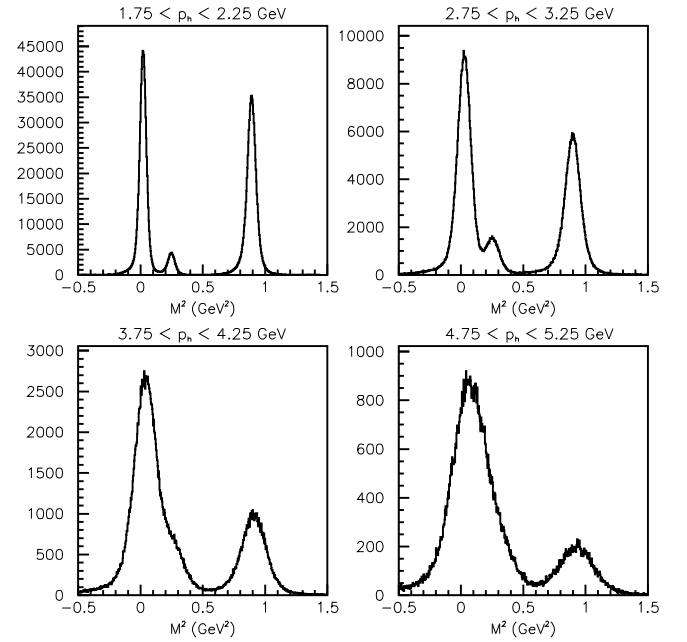


Fig. 31. Reconstructed mass squared (GeV^2) for positively charged particles from the timing in the panel-1b FTOF counters from beam data with a 10.6-GeV electron beam incident on a liquid-hydrogen target. The data are sorted into four bins in hadron momentum (UL) [1.75:2.25] GeV, (UR) [2.75:3.25] GeV, (LL) [3.75:4.25] GeV, and (LR) [4.75:5.25] GeV and are based on the current CLAS12 detector calibrations, detector alignments, and knowledge of the torus and solenoid field maps.

truncate the mass squared distribution when taking the square root, Fig. 31 plots the reconstructed hadron mass squared using:

$$M^2 = p_p^2 \cdot \frac{1 - \beta^2}{\beta^2}, \quad \beta = \frac{P_L}{t_p c}, \quad (24)$$

where p_p is the particle momentum determined by the forward tracking system, P_L is the hadron path length from the event vertex in the target to the FTOF system, $t_p = \bar{t}_{hit} - t_{ST}$ is the hadron flight time over P_L , and c is the speed of light. Fig. 31 shows the mass squared distributions for several bins in hadron momentum from 2 GeV to 5 GeV. Even with the current state of CLAS12 detector calibrations, detector alignment, and knowledge of the torus and solenoid magnetic field, these distributions show good separation between π/K up to about 3 GeV and separation of π and K from p to about 5 GeV. Improvements in the FTOF and forward tracking calibrations, detector alignment, and knowledge of the magnetic fields in the CLAS12 tracking volume that are still expected should give improved particle identification and separation of the different charged hadron species to higher momenta. In addition,

employing the combined timing measurement for the hits in panel-1a and panel-1b (see Section 5.2.8) will also result in improved particle identification using the FTOF system.

6. Summary

We have designed and assembled a time-of-flight system for the Forward Detector of the new CLAS12 Spectrometer in Hall B at Jefferson Lab known as the Forward Time-of-Flight or FTOF system, which consists of 90 scintillation bars in each of the six sectors of the CLAS12 Forward Detector for a total of 540 counters. This design is based on rectangular counters varying in length from 17 cm to 426 cm in three different counter arrays in each sector. In the polar angle range from 5° to 35° the FTOF system consists of two layers of counters referred to as panel-1a and panel-1b. The panel-1a counters were refurbished from the forward TOF counters that were part of the original CLAS spectrometer. The panel-1b counters were newly constructed for CLAS12. Together the timing measurements from these two arrays of counters currently provide effective time resolutions from 50 ps for the shortest counters at small polar angles to 100 ps for the longest counters at large polar angles. In the polar angle range from 35° to 45° the FTOF system consists of the panel-2 counters that were refurbished from the CLAS TOF system. These counters provide effective time resolutions of about 250 ps. With these time resolutions the FTOF system can separate π/K to 2.8 GeV, K/p to 4.8 GeV, and π/p to 5.4 GeV with 4σ separation with up to an order of magnitude difference in the relative yields. The specifications are sufficient to meet the design particle identification requirements in the forward direction for the full CLAS12 physics program. The performance of the FTOF system has been verified in extensive bench studies in our cosmic ray test stands, as well as after installation in the first beam runs with the CLAS12 system in the period from Dec. 2017 to Dec. 2018.

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