# On the integrability of hybrid Hamiltonian systems

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# Symplectic geometry

- Symplectic geometry is the natural framework for classical mechanics.
- Recall that a symplectic form  $\omega$  on M is a 2-form such that  $d\omega = 0$  and  $T_x M \ni v \mapsto \omega_x(v, \cdot) \in T_x^* M$  is an isomorphism of vector spaces.
- Given a function f on M, its its Hamiltonian vector field  $X_f$  is given by

$$\omega(X_f,\cdot)=\mathrm{d}f.$$

• The Poisson bracket  $\{\cdot,\cdot\}$  is given by

$$\{f,g\} := \omega(X_f,X_g) = X_g(f) = -X_f(g).$$

## Theorem (Liouville – Arnol'd theorem)

Let  $f_1, \ldots, f_n$  be independent functions in involution (i.e.,  $\{f_i, f_j\} = 0 \ \forall i, j$ ) on a symplectic manifold  $(M^{2n}, \omega)$ . Let  $M_{\Lambda} = \{x \in M \mid f_i = \Lambda_i\}$ .

- **1** Any compact connected component of  $M_{\Lambda}$  is diffeomorphic to  $\mathbb{T}^n$ .
- **2** On a neighborhood of  $M_{\Lambda}$  there are coordinates  $(\varphi^i, J_i)$  such that

$$\omega = \mathrm{d}\varphi^i \wedge \mathrm{d}J_i,$$

and the Hamiltonian dynamics are given by

$$\frac{\mathrm{d}\varphi^i}{\mathrm{d}t} = \Omega^i(J_1,\ldots,J_n),$$
$$\frac{\mathrm{d}J_i}{\mathrm{d}t} = 0.$$

# Hybrid systems

#### Definition

A **hybrid system** is a 4-tuple  $\mathcal{H} = (M, X, S, \Delta)$ , formed by

- $\mathbf{0}$  a manifold M,
- **2** a vector field  $X \in \mathfrak{X}(M)$ ,
- **3** a submanifold  $S \subset M$  of codimension 1 or greater,
- **4** an embedding  $\Delta: S \to M$ .

The dynamics generated by  $\mathscr{H}$  are the curves  $c: I \subseteq \mathbb{R} \to M$  such that

$$\dot{c}(t) = X(c(t)), \quad \text{if } c(t) \notin S,$$
  
 $c^+(t) = \Delta(c^-(t)), \quad \text{if } c(t) \in S,$ 

where

$$c^{\pm}(t) = \lim_{\tau \to t^{\pm}} c(\tau).$$

# Hybrid Hamiltonian systems

#### Definition

A hybrid dynamical system  $(M, X, S, \Delta)$  is said to be a **hybrid Hamiltonian system** and denoted by  $\mathcal{H}_h$  if

- **1**  $M \subseteq T^*Q$  is a zero-codimensional submanifold of the cotangent bundle  $\pi_Q \colon T^*Q \to Q$  of a manifold Q,
- **2** S projects onto a codimension-one submanifold  $\pi_Q(S)$  of Q,
- **4**  $X = X_h$  is the Hamiltonian vector field of  $h \in C^{\infty}(T^*Q)$  w.r.t. the canonical symplectic form  $\omega_Q$ , namely,

$$\omega_{\mathcal{O}}(X_h) = \mathrm{d}h$$
.

# Hybrid Hamiltonian systems

## Physically,

- Q represents the space of positions,
- T\*Q the phase space,
- $X_h$  the dynamics between the impacts,
- $\pi_Q(S)$  the hypersurface where impacts occur, and
- ullet  $\Delta$  the change of momenta on the impacts.

# Hybrid Lie group action

### Definition

A Lie group action  $\Phi \colon G \times Q \to Q$  is called a **hybrid action for**  $\mathscr{H}_h$  if its cotangent lift  $\Phi^{T^*} \colon G \times T^*Q \to T^*Q$  satisfies the following conditions:

- $oldsymbol{0}$  h is  $\Phi^{\mathrm{T}^*}$ -invariant, namely,  $h \circ \Phi_g^{\mathrm{T}^*} = h$  for all  $g \in G$ ,
- 2 the restriction  $\Phi^{T^*}\Big|_{G\times S}$  is a Lie group action of G on S,
- 3 the impact map is equivariant w.r.t. this action, i.e.,

$$\left. \Delta \circ \Phi_g^{\mathrm{T}^*} \right|_S = \Phi_g^{\mathrm{T}^*} \circ \Delta \,, \quad \forall \, g \in \textit{G} \,.$$

# Hybrid momentum map

#### Definition

Let  $\Phi\colon G\times Q\to Q$  be a hybrid action for  $\mathscr{H}_h$ . A momentum map  $\mathbf{J}\colon \mathrm{T}^*Q\to \mathfrak{g}^*$  for the cotangent lift action  $\Phi^{\mathrm{T}^*}$  is called a **generalized** hybrid momentum map if, for each connected component  $C\subseteq S$  and for each regular value  $\mu_-$  of  $\mathbf{J}$ , there is another regular value  $\mu_+$  such that

$$\Delta(\mathbf{J}|_{C}^{-1}(\mu_{-})) \subset \mathbf{J}^{-1}(\mu_{+}).$$

In particular, if  $\mu_- = \mu_+$  it is called a **hybrid momentum map**. A **hybrid regular value** of **J** is a regular value of both **J** and  $\mathbf{J}|_{\mathcal{S}}$ .

# Hybrid momentum map

In other words,  ${\bf J}$  is a generalized hybrid momentum map if, for every point in the connected component C of the switching surface S such that the momentum before the impact takes a value of  $\mu_-$ , the momentum will take a value  $\mu_+$  after the impact; and it is a hybrid momentum map if its value does not change with the impacts.

# Hybrid reduction

## **Proposition**

If  $\mu_-$  and  $\mu_+$  are regular values of  $\mathbf{J}$  such that  $\Delta\left(\mathbf{J}|_{\mathcal{S}}^{-1}(\mu_-)\right)\subset\mathbf{J}^{-1}(\mu_+)$ , then the isotropy subgroups in  $\mu_-$  and  $\mu_+$  coincide, that is,  $G_{\mu_-}=G_{\mu_+}$ .

# Hybrid reduction

## Theorem (Colombo, de León, Eyrea Irazú, L. G., 2022)

Let  $\Phi \colon G \times Q \to Q$  be a hybrid action on  $\mathscr{H}_h$ . Assume that G is connected and that  $\Phi^{T^*} \colon G \times T^*Q \to T^*Q$  is free and proper. Consider a sequence  $\{\mu_i\}_{i \in I \subseteq \mathbb{N}}$  of hybrid regular values of  $\mathbf{J}$ , such that  $\Delta \left(\mathbf{J}|_{S}^{-1}(\mu_i)\right) \subset \mathbf{J}^{-1}(\mu_{i+1})$ . Let  $G_{\mu_i} = G_{\mu_0}$  be the isotropy subgroup in  $\mu_i$ 

under the co-adjoint action. Then, the reduction leads to a sequence of reduced hybrid forced Hamiltonian systems

$$\mathscr{H}_{h}^{\mu_{i}} = \left(\mathbf{J}^{-1}(\mu_{i})/G_{\mu_{0}}, X_{h_{\mu_{i}}}, \mathbf{J}|_{S}^{-1}(\mu_{i})/G_{\mu_{0}}, (\Delta)_{\mu_{i}}\right).$$

# Hybrid reduction

$$\cdots \longrightarrow \mathbf{J}^{-1}(\mu_{i}) \longleftarrow \mathbf{J}|_{S}^{-1}(\mu_{i}) \xrightarrow{\Delta|_{\mathbf{J}^{-1}(\mu_{i})}} \mathbf{J}^{-1}(\mu_{i+1}) \longleftarrow \cdots$$

$$\downarrow \qquad \qquad \downarrow \qquad \qquad \downarrow$$

$$\cdots \longrightarrow \frac{\mathbf{J}^{-1}(\mu_{i})}{G_{\mu_{0}}} \longleftarrow \mathbf{J}|_{S}^{-1}(\mu_{i})/G_{\mu_{0}} \xrightarrow{(\Delta)_{\mu_{i}}} \mathbf{J}^{-1}(\mu_{i+1}) \longleftarrow \cdots$$

# Integrable hybrid Hamiltonian systems

- A particular case is when we have the Abelian Lie group action
   Φ: ℝ<sup>n</sup> × T\*Q → T\*Q generated by the Hamiltonian flows of n functions f<sub>1</sub>,..., f<sub>n</sub> in involution.
- In that case, we can identify the momentum map with  $F = (f_1, \dots, f_n) \colon \mathrm{T}^* Q \to \mathbb{R}^n$ .
- We may obtain action-angle coordinates for each time interval between impacts. The action-angle coordinates before and after the impact will be related by  $\Delta$ .

#### Definition

Let  $(M, S, X, \Delta)$  be a hybrid dynamical system. A function  $f: M \to \mathbb{R}$  is called a **generalized hybrid constant of the motion** if

- **2** For each connected component  $C \subseteq S$  and each  $a \in \operatorname{Im} f$ , there exists a  $b \in \operatorname{Im} f$  such that

$$\Delta\left(f|_{C}^{-1}(a)\right)\subseteq f^{-1}(b)$$
.

In particular, f is called a **hybrid constant of the motion** if, in addition, b = a for each  $a \in \text{Im } f$ .

#### Definition

Let Q be an n-dimensional manifold. A completely integrable hybrid Hamiltonian system is a 5-tuple

 $(\mathrm{T}^*Q,S,X_H,\Delta,F)$ , formed by a hybrid Hamiltonian system  $(\mathrm{T}^*Q,S,X_H,\Delta)$ , together with a function  $F=(f_1,\ldots,f_n)\colon \mathrm{T}^*Q\to\mathbb{R}^n$  such that:

- 2 the functions  $f_1, \ldots, f_n$  are generalized hybrid constant of the motion

## Theorem (L. G., Colombo, 2024)

Consider a completely integrable hybrid Hamiltonian system  $(T^*Q, S, X_H, \Delta)$ , with  $F = (f_1, \ldots, f_n)$ , where  $n = \dim Q$ . Let  $M_{\Lambda}$  be a regular level set of F. Then:

- **1** For each regular level set  $M_{\Lambda}$  and each connected component  $C \subseteq S$ , there exists a  $\Lambda' \in \mathbb{R}^n$  such that  $\Delta(M_{\Lambda} \cap C) \subset M_{\Lambda'} = F^{-1}(\Lambda')$ .
- **2** On a neighbourhood  $U_{\lambda}$  of  $M_{\Lambda}$  there are coordinates  $(\varphi^{i}, s_{i})$  s.t.

  - **2** the action coordinates  $s_i$  are functions depending only on the integrals  $f_1, \ldots, f_n$ ,
  - 3 the continuous part hybrid dynamics are given by

$$\dot{\varphi}^i = \Omega^i(s_1, \ldots, s_n), \qquad \dot{s}_i = 0.$$

**4** In these coordinates, for each connected component  $C \subseteq S$ , the impact map reads  $\Delta$ :  $(\varphi_-^i, s_i^-) \in M_{\Lambda} \cap C \mapsto (\varphi_+^i, s_i^+) \in M_{\Lambda'}$ , where  $s_1^+, \ldots, s_n^+$  are functions depending only on  $s_1^-, \ldots, s_n^-$ .

- Consider a homogeneous circular disk of radius *R* and mass *m* moving in the plane.
- The configuration space is  $Q = \mathbb{R}^2 \times \mathbb{S}^1$ , with canonical coordinates  $(x, y, \theta)$ .
- The coordinates (x, y) represent then position of the center of the disk, while the coordinate  $\theta$  represents the angle between a fixed reference point of the disk and the y-axis.

• The Hamiltonian function  $H \colon \mathrm{T}^*Q \to \mathbb{R}$  of the system is

$$H = \frac{1}{2m}(p_x^2 + p_y^2) + \frac{1}{2mk^2}p_\theta^2 + \frac{1}{2}\Omega^2(x^2 + y^2),$$

where  $(x, y, \theta, p_x, p_y, p_\theta)$  are the bundle coordinates in  $\mathrm{T}^*(\mathbb{R}^2 \times \mathbb{S}^1)$ .

- Suppose that there are two rough walls situated at y = 0 and at y = h > R.
- Assume that the impact with a wall is such that the disk rolls without sliding and that the change of the velocity along the y-direction is characterized by an elastic constant e

• Then, the switching surface is  $S = C_1 \cup C_2$ , where

$$C_{1} = \left\{ \left( x, R, \theta, p_{x}, p_{y}, \frac{k^{2}}{R} p_{x} \right) \mid x, p_{x}, p_{y} \in \mathbb{R}, \ \theta \in \mathbb{S}^{1} \right\},$$

$$C_{2} = \left\{ \left( x, h - R, \theta, p_{x}, p_{y}, \frac{k^{2}}{R} p_{x} \right) \mid x, p_{x}, p_{y} \in \mathbb{R}, \ \theta \in \mathbb{S}^{1} \right\},$$

and the impact map  $\Delta \colon S o \mathrm{T}^*Q$  is given by

$$\left(p_{x}^{-},p_{y}^{-},p_{\theta}^{-}\right)\mapsto\left(\frac{R^{2}p_{x}^{-}+k^{2}Rp_{\theta}^{-}}{k^{2}+R^{2}},-ep_{y}^{-},\frac{Rp_{x}^{-}+k^{2}p_{\theta}^{-}}{k^{2}+R^{2}}\right)$$

- For simplicity's sake, let us hereafter take  $m=R=k=\Omega=1$ .
- The functions

$$f_1 = \frac{p_x^2 + x^2}{2}$$
,  $f_2 = \frac{p_y^2 + y^2}{2}$ ,  $f_3 = \frac{p_\theta^2}{2}$ ,

are conserved quantities with respect to the Hamiltonian dynamics of H.

- Moreover,  $\{f_i, f_i\} = 0$  and  $df_1 \wedge df_2 \wedge df_3 \neq 0$  a.e.
- Let  $F = (f_1, f_2, f_3) \colon \mathrm{T}^*(\mathbb{R}^2 \times \mathbb{S}) \to \mathbb{R}^3$ .
- It is clear that, for  $\Lambda \neq 0$ , the level sets  $F^{-1}(\Lambda)$  are diffeomorphic to  $\mathbb{S} \times \mathbb{S} \times \mathbb{R}$ .

• In the intersection of their domains of definition, the functions

$$\phi^1 = \arctan\left(\frac{x}{p_{\scriptscriptstyle X}}\right)\,,\quad \phi^2 = \arctan\left(\frac{y}{p_{\scriptscriptstyle Y}}\right)\,,\quad \phi^3 = \frac{\theta}{p_\theta}$$

are coordinates on each level set  $F^{-1}(\Lambda)$  for  $\Lambda \neq 0$ .

- Additionally,  $\omega_Q = \mathrm{d}\phi^i \wedge \mathrm{d}f_i$ .
- In these coordinates, the Hamiltonian function reads

$$H = f_1 + f_2 + f_3$$
.

• Hence, its Hamiltonian vector field is simply

$$X_H = \frac{\partial}{\partial \phi^1} + \frac{\partial}{\partial \phi^2} + \frac{\partial}{\partial \phi^3}.$$

• In the action-angle coordinates  $(\phi^i, f_i)$ , the impact surface reads

$$S = \left\{ \left( \phi^i, f_i \right) \mid 2f_2 \sin^2 \phi^2 = R^2 \text{ and } f_3 = \frac{2k^4 f_1 \cos^2 \phi^1}{R^2} \right\}$$

$$\cup \left\{ 2f_2 \sin^2 \phi^2 = (h - R)^2 \text{ and } f_3 = \frac{2k^4 f_1 \cos^2 \phi^1}{R^2} \right\}.$$

• The relations between the coordinates before,  $(\phi_-^i, f_i^-)$ , and after,  $(\phi_+^i, f_i^+)$ , are

$$\phi_{+}^{1} = \phi_{-}^{1}, \qquad \phi_{+}^{2} = -\arctan\left(\frac{\tan\phi_{-}^{2}}{e}\right), \qquad \phi_{+}^{3} = \phi_{-}^{3},$$
  $f_{+}^{1+} = f_{-}^{1-}, \qquad f_{2}^{+} = e^{2}f_{2} + \frac{1-e^{2}}{2}a^{2}, \qquad f_{3}^{+} = f_{3}^{-},$ 

 $f_1^+ = f_1^-, \qquad f_2^+ = e^2 f_2 + \frac{1 - e^2}{2} a^2, \qquad \qquad f_3^+ = f_3^-,$ 

where a = R or a = h - R depending on the wall where the impact takes place.

# Merci pour votre attention!

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