

On the integrability of hybrid Hamiltonian systems

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Symplectic geometry

- Symplectic geometry is the natural framework for classical mechanics.
- Recall that a symplectic form ω on M is a 2-form such that $d\omega = 0$ and $T_x M \ni v \mapsto \omega_x(v, \cdot) \in T_x^* M$ is an isomorphism of vector spaces.
- Given a function f on M , its Hamiltonian vector field X_f is given by

$$\omega(X_f, \cdot) = df.$$

- The Poisson bracket $\{\cdot, \cdot\}$ is given by

$$\{f, g\} := \omega(X_f, X_g) = X_g(f) = -X_f(g).$$

Theorem (Liouville–Arnold theorem)

Let f_1, \dots, f_n be independent functions in involution (i.e., $\{f_i, f_j\} = 0 \ \forall i, j$) on a symplectic manifold (M^{2n}, ω) . Let $M_\Lambda = \{x \in M \mid f_i = \Lambda_i\}$.

- ① Any compact connected component of M_Λ is diffeomorphic to \mathbb{T}^n .
- ② On a neighborhood of M_Λ there are coordinates (φ^i, J_i) such that

$$\omega = d\varphi^i \wedge dJ_i,$$

and the Hamiltonian dynamics are given by

$$\begin{aligned} \frac{d\varphi^i}{dt} &= \Omega^i(J_1, \dots, J_n), \\ \frac{dJ_i}{dt} &= 0. \end{aligned}$$

Hybrid systems

Definition

A **hybrid system** is a 4-tuple $\mathcal{H} = (M, X, S, \Delta)$, formed by

- ① a manifold M ,
- ② a vector field $X \in \mathfrak{X}(M)$,
- ③ a submanifold $S \subset M$ of codimension 1 or greater,
- ④ an embedding $\Delta: S \rightarrow M$.

The dynamics generated by \mathcal{H} are the curves $c: I \subseteq \mathbb{R} \rightarrow M$ such that

$$\begin{aligned}\dot{c}(t) &= X(c(t)), & \text{if } c(t) \notin S, \\ c^+(t) &= \Delta(c^-(t)), & \text{if } c(t) \in S,\end{aligned}$$

where

$$c^\pm(t) = \lim_{\tau \rightarrow t^\pm} c(\tau).$$

Hybrid Hamiltonian systems

Definition

A hybrid dynamical system (M, X, S, Δ) is said to be a **hybrid Hamiltonian system** and denoted by \mathcal{H}_h if

- ① $M \subseteq T^*Q$ is a zero-codimensional submanifold of the cotangent bundle $\pi_Q: T^*Q \rightarrow Q$ of a manifold Q ,
- ② S projects onto a codimension-one submanifold $\pi_Q(S)$ of Q ,
- ③ $\pi_Q \circ \Delta = \pi_Q$,
- ④ $X = X_h$ is the Hamiltonian vector field of $h \in C^\infty(T^*Q)$ w.r.t. the canonical symplectic form ω_Q , namely,

$$\omega_Q(X_h) = dh.$$

Hybrid Hamiltonian systems

Physically,

- Q represents the space of positions,
- T^*Q the phase space,
- X_h the dynamics between the impacts,
- $\pi_Q(S)$ the hypersurface where impacts occur, and
- Δ the change of momenta on the impacts.

Hybrid Lie group action

Definition

A Lie group action $\Phi: G \times Q \rightarrow Q$ is called a **hybrid action for \mathcal{H}_h** if its cotangent lift $\Phi^{T^*}: G \times T^*Q \rightarrow T^*Q$ satisfies the following conditions:

- ① h is Φ^{T^*} -invariant, namely, $h \circ \Phi_g^{T^*} = h$ for all $g \in G$,
- ② the restriction $\Phi^{T^*}|_{G \times S}$ is a Lie group action of G on S ,
- ③ the impact map is equivariant w.r.t. this action, i.e.,

$$\Delta \circ \Phi_g^{T^*}|_S = \Phi_g^{T^*} \circ \Delta, \quad \forall g \in G.$$

Hybrid momentum map

Definition

Let $\Phi: G \times Q \rightarrow Q$ be a hybrid action for \mathcal{H}_h . A momentum map $\mathbf{J}: T^*Q \rightarrow \mathfrak{g}^*$ for the cotangent lift action Φ^{T^*} is called a **generalized hybrid momentum map** if, for each connected component $C \subseteq S$ and for each regular value μ_- of \mathbf{J} , there is another regular value μ_+ such that

$$\Delta(\mathbf{J}|_C^{-1}(\mu_-)) \subset \mathbf{J}^{-1}(\mu_+).$$

In particular, if $\mu_- = \mu_+$ it is called a **hybrid momentum map**. A **hybrid regular value** of \mathbf{J} is a regular value of both \mathbf{J} and $\mathbf{J}|_S$.

Hybrid momentum map

In other words, \mathbf{J} is a generalized hybrid momentum map if, for every point in the connected component C of the switching surface S such that the momentum before the impact takes a value of μ_- , the momentum will take a value μ_+ after the impact; and it is a hybrid momentum map if its value does not change with the impacts.

Hybrid reduction

Proposition

If μ_- and μ_+ are regular values of \mathbf{J} such that $\Delta \left(\mathbf{J}|_S^{-1}(\mu_-) \right) \subset \mathbf{J}^{-1}(\mu_+)$, then the isotropy subgroups in μ_- and μ_+ coincide, that is, $G_{\mu_-} = G_{\mu_+}$.

Hybrid reduction

Theorem (Colombo, de León, Eyrea Irazú, L. G., 2022)

Let $\Phi: G \times Q \rightarrow Q$ be a hybrid action on \mathcal{H}_h . Assume that G is connected and that $\Phi^{T^*}: G \times T^*Q \rightarrow T^*Q$ is free and proper. Consider a sequence $\{\mu_i\}_{i \in \mathbb{N}}$ of hybrid regular values of \mathbf{J} , such that $\Delta \left(\mathbf{J}|_S^{-1}(\mu_i) \right) \subset \mathbf{J}^{-1}(\mu_{i+1})$. Let $G_{\mu_i} = G_{\mu_0}$ be the isotropy subgroup in μ_i under the co-adjoint action. Then, the reduction leads to a sequence of reduced hybrid forced Hamiltonian systems

$$\mathcal{H}_h^{\mu_i} = \left(\mathbf{J}^{-1}(\mu_i)/G_{\mu_0}, X_{h_{\mu_i}}, \mathbf{J}|_S^{-1}(\mu_i)/G_{\mu_0}, (\Delta)_{\mu_i} \right).$$

Hybrid reduction

$$\begin{array}{ccccccc}
 \dots & \longrightarrow & \mathbf{J}^{-1}(\mu_i) & \longleftrightarrow & \mathbf{J}|_S^{-1}(\mu_i) & \xrightarrow{\Delta|_{\mathbf{J}^{-1}(\mu_i)}} & \mathbf{J}^{-1}(\mu_{i+1}) & \longleftrightarrow & \dots \\
 & & \downarrow & & \downarrow & & \downarrow & & \\
 \dots & \longrightarrow & \frac{\mathbf{J}^{-1}(\mu_i)}{G_{\mu_0}} & \longleftrightarrow & \mathbf{J}|_S^{-1}(\mu_i)/G_{\mu_0} & \xrightarrow{(\Delta)_{\mu_i}} & \frac{\mathbf{J}^{-1}(\mu_{i+1})}{G_{\mu_0}} & \longleftrightarrow & \dots
 \end{array}$$

Integrable hybrid Hamiltonian systems

- A particular case is when we have the Abelian Lie group action $\Phi: \mathbb{R}^n \times T^*Q \rightarrow T^*Q$ generated by the Hamiltonian flows of n functions f_1, \dots, f_n in involution.
- In that case, we can identify the momentum map with $F = (f_1, \dots, f_n): T^*Q \rightarrow \mathbb{R}^n$.
- We may obtain action-angle coordinates for each time interval between impacts. The action-angle coordinates before and after the impact will be related by Δ .

Definition

Let (M, S, X, Δ) be a hybrid dynamical system. A function $f: M \rightarrow \mathbb{R}$ is called a **generalized hybrid constant of the motion** if

- ① $Xf = 0$,
- ② For each connected component $C \subseteq S$ and each $a \in \text{Im } f$, there exists a $b \in \text{Im } f$ such that

$$\Delta \left(f|_C^{-1}(a) \right) \subseteq f^{-1}(b).$$

In particular, f is called a **hybrid constant of the motion** if, in addition, $b = a$ for each $a \in \text{Im } f$.

Definition

Let Q be an n -dimensional manifold. A **completely integrable hybrid Hamiltonian system** is a 5-tuple

$(T^*Q, S, X_H, \Delta, F)$, formed by a hybrid Hamiltonian system (T^*Q, S, X_H, Δ) , together with a function $F = (f_1, \dots, f_n): T^*Q \rightarrow \mathbb{R}^n$ such that:

- ① $\text{rank } T_x F = n$ a.e.,
- ② the functions f_1, \dots, f_n are generalized hybrid constant of the motion
- ③ $\{f_i, f_j\} = X_{f_j}(f_i) = 0 \quad \forall i, j \in \{1, \dots, n\}$.

Theorem (L. G., Colombo, 2024)

Consider a completely integrable hybrid Hamiltonian system (T^*Q, S, X_H, Δ) , with $F = (f_1, \dots, f_n)$, where $n = \dim Q$. Let M_Λ be a regular level set of F . Then:

- 1 For each regular level set M_Λ and each connected component $C \subseteq S$, there exists a $\Lambda' \in \mathbb{R}^n$ such that $\Delta(M_\Lambda \cap C) \subset M_{\Lambda'} = F^{-1}(\Lambda')$.
- 2 On a neighbourhood U_λ of M_Λ there are coordinates (φ^i, s_i) s.t.
 - 1 $\omega_Q = d\varphi^i \wedge ds_i$,
 - 2 the action coordinates s_i are functions depending only on the integrals f_1, \dots, f_n ,
 - 3 the continuous part hybrid dynamics are given by

$$\dot{\varphi}^i = \Omega^i(s_1, \dots, s_n), \quad \dot{s}_i = 0.$$

- 4 In these coordinates, for each connected component $C \subseteq S$, the impact map reads $\Delta: (\varphi_-^i, s_i^-) \in M_\Lambda \cap C \mapsto (\varphi_+^i, s_i^+) \in M_{\Lambda'}$, where s_1^+, \dots, s_n^+ are functions depending only on s_1^-, \dots, s_n^- .

Rolling disk with a harmonic potential hitting fixed walls

- Consider a homogeneous circular disk of radius R and mass m moving in the plane.
- The configuration space is $Q = \mathbb{R}^2 \times \mathbb{S}^1$, with canonical coordinates (x, y, θ) .
- The coordinates (x, y) represent then position of the center of the disk, while the coordinate θ represents the angle between a fixed reference point of the disk and the y -axis.

Rolling disk with a harmonic potential hitting fixed walls

- The Hamiltonian function $H: T^*Q \rightarrow \mathbb{R}$ of the system is

$$H = \frac{1}{2m}(p_x^2 + p_y^2) + \frac{1}{2mk^2}p_\theta^2 + \frac{1}{2}\Omega^2(x^2 + y^2),$$

where $(x, y, \theta, p_x, p_y, p_\theta)$ are the bundle coordinates in $T^*(\mathbb{R}^2 \times \mathbb{S}^1)$.

Rolling disk with a harmonic potential hitting fixed walls

- Suppose that there are two rough walls situated at $y = 0$ and at $y = h > R$.
- Assume that the impact with a wall is such that the disk rolls without sliding and that the change of the velocity along the y -direction is characterized by an elastic constant e

Rolling disk with a harmonic potential hitting fixed walls

- Then, the switching surface is $S = C_1 \cup C_2$, where

$$C_1 = \left\{ \left(x, R, \theta, p_x, p_y, \frac{k^2}{R} p_x \right) \mid x, p_x, p_y \in \mathbb{R}, \theta \in \mathbb{S}^1 \right\},$$

$$C_2 = \left\{ \left(x, h - R, \theta, p_x, p_y, \frac{k^2}{R} p_x \right) \mid x, p_x, p_y \in \mathbb{R}, \theta \in \mathbb{S}^1 \right\},$$

and the impact map $\Delta: S \rightarrow T^*Q$ is given by

$$(p_x^-, p_y^-, p_\theta^-) \mapsto \left(\frac{R^2 p_x^- + k^2 R p_\theta^-}{k^2 + R^2}, -e p_y^-, \frac{R p_x^- + k^2 p_\theta^-}{k^2 + R^2} \right)$$

Rolling disk with a harmonic potential hitting fixed walls

- For simplicity's sake, let us hereafter take $m = R = k = \Omega = 1$.
- The functions

$$f_1 = \frac{p_x^2 + x^2}{2}, \quad f_2 = \frac{p_y^2 + y^2}{2}, \quad f_3 = \frac{p_\theta^2}{2},$$

are conserved quantities with respect to the Hamiltonian dynamics of H .

- Moreover, $\{f_i, f_j\} = 0$ and $df_1 \wedge df_2 \wedge df_3 \neq 0$ a.e.
- Let $F = (f_1, f_2, f_3): T^*(\mathbb{R}^2 \times \mathbb{S}) \rightarrow \mathbb{R}^3$.
- It is clear that, for $\Lambda \neq 0$, the level sets $F^{-1}(\Lambda)$ are diffeomorphic to $\mathbb{S} \times \mathbb{S} \times \mathbb{R}$.

Rolling disk with a harmonic potential hitting fixed walls

- In the intersection of their domains of definition, the functions

$$\phi^1 = \arctan\left(\frac{x}{p_x}\right), \quad \phi^2 = \arctan\left(\frac{y}{p_y}\right), \quad \phi^3 = \frac{\theta}{p_\theta}$$

are coordinates on each level set $F^{-1}(\Lambda)$ for $\Lambda \neq 0$.

- Additionally, $\omega_Q = d\phi^i \wedge df_i$.
- In these coordinates, the Hamiltonian function reads

$$H = f_1 + f_2 + f_3.$$

- Hence, its Hamiltonian vector field is simply

$$X_H = \frac{\partial}{\partial \phi^1} + \frac{\partial}{\partial \phi^2} + \frac{\partial}{\partial \phi^3}.$$

Rolling disk with a harmonic potential hitting fixed walls

- In the action-angle coordinates (ϕ^i, f_i) , the impact surface reads

$$S = \left\{ (\phi^i, f_i) \mid 2f_2 \sin^2 \phi^2 = R^2 \text{ and } f_3 = \frac{2k^4 f_1 \cos^2 \phi^1}{R^2} \right\} \\ \cup \left\{ 2f_2 \sin^2 \phi^2 = (h - R)^2 \text{ and } f_3 = \frac{2k^4 f_1 \cos^2 \phi^1}{R^2} \right\}.$$

Rolling disk with a harmonic potential hitting fixed walls

- The relations between the coordinates before, (ϕ_-^i, f_i^-) , and after, (ϕ_+^i, f_i^+) , are

$$\phi_+^1 = \phi_-^1, \quad \phi_+^2 = -\arctan\left(\frac{\tan \phi_-^2}{e}\right), \quad \phi_+^3 = \phi_-^3,$$

$$f_1^+ = f_1^-, \quad f_2^+ = e^2 f_2 + \frac{1-e^2}{2} a^2, \quad f_3^+ = f_3^-,$$

where $a = R$ or $a = h - R$ depending on the wall where the impact takes place.

Merci pour votre attention!

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