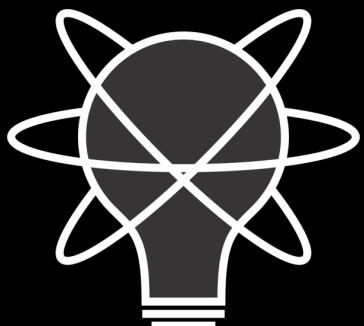


# What is classical? What is quantum?

Insights from the Assumptions of Physics program



Assumptions  
of  
Physics

Gabriele Carcassi  
University of Michigan



UNIVERSITY OF  
MICHIGAN

## International Year of Quantum Science and Technology 2025



International Year of Quantum logo

Date	1 January – 31 December 2025
Type	Exhibitions
Website	<a href="https://quantum2025.org">quantum2025.org</a>

wikipedia

nature

## Quantum mechanics 100 years on: an unfinished revolution

A century ago, physics had its Darwinian moment – a change in perspective that was as consequential for the physical sciences as the theory of evolution by natural selection was for biology.



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QUANTUM PHYSICS



## Quantum mechanics was born 100 years ago. Physicists are celebrating

The International Year of Quantum marks a century of scientific developments

## A century of quantum mechanics

On 9 July 1925, in a letter to Wolfgang Pauli, Werner Heisenberg revealed his new ideas, which were to revolutionise physics



<https://assumptionsofphysics.org/>

Assumptions  
of  
Physics

# A brief history of quantum mechanics

There were problems  
with classical physics

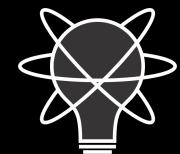
Blackbody  
radiation

Stability of atoms

Photoelectric  
effect

People futzed around with matrices and waves,  
to find something new that has some  
correspondence with classical mechanics

⇒ Quantum mechanics!



# How to quantize a classical theory

Classical mechanics

$$\frac{df}{dt} = \{f, H\}$$

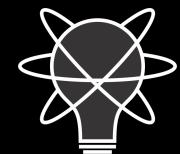


Quantum mechanics

$$\frac{df}{dt} = \frac{[f, H]}{i\hbar}$$

## Quantization!

Dirac's correspondence principle



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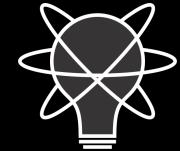
$$\{x, p\} = 1$$

$$\frac{[x, p]}{i\hbar} = 1$$

$$\{L_x, L_y\} = L_z$$

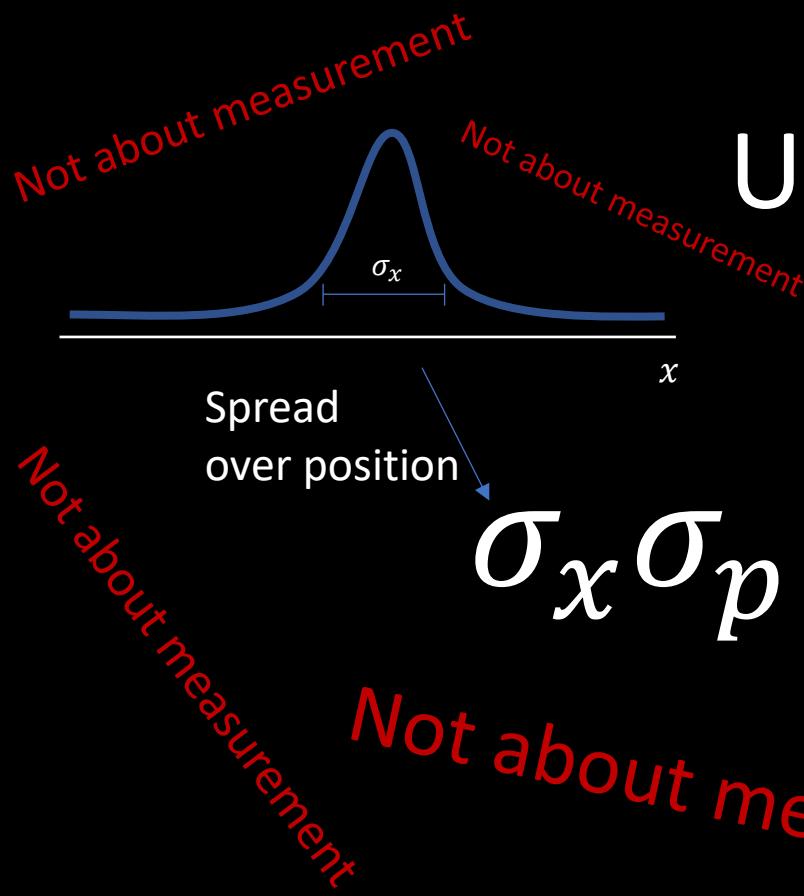
$$\frac{[L_x, L_y]}{i\hbar} = L_z$$

Quantization is a “formal” operation  
(i.e. substitute some symbols with some other symbols)



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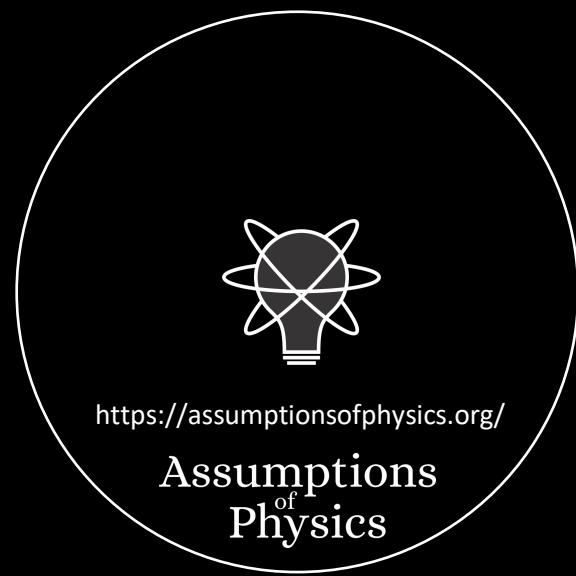
Assumptions  
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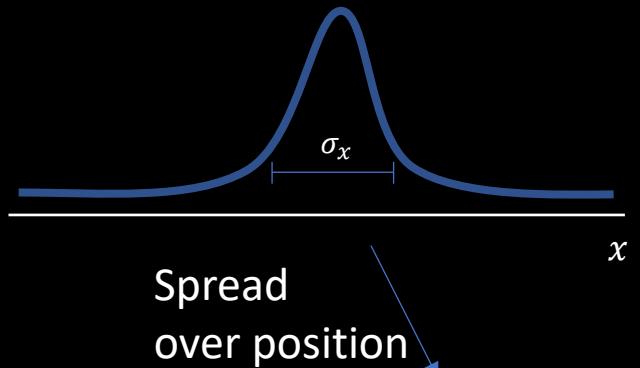
# Uncertainty principle

$$\sigma_x \sigma_p \geq \frac{1}{2} |\langle [x, p] \rangle| = \frac{\hbar}{2}$$

# Not about measurement

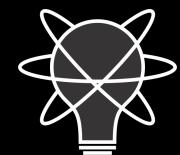


# Indetermination ~~Uncertainty~~ principle



$$\sigma_x \sigma_p \geq \frac{1}{2} |\langle [x, p] \rangle| = \frac{\hbar}{2}$$

About preparation  
not measurement!

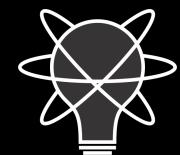


# Quantum contextuality

Measurements cannot be thought as revealing pre-existing values

# Quantum complementarity

For every quantity we prepare, another quantity is indeterminate



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During physical processes

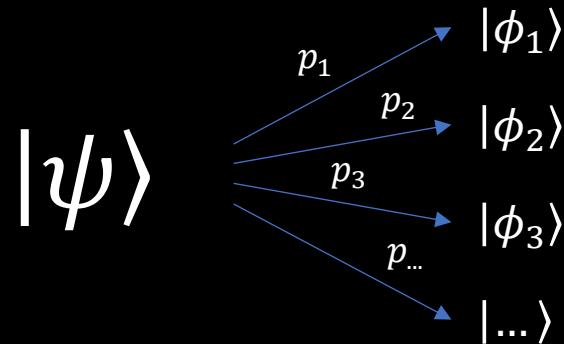
Schrödinger equation

$$i\hbar \frac{d}{dt} |\psi\rangle = H|\psi\rangle$$

smooth, deterministic and reversible

During measurements

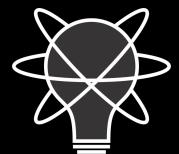
Projections



jumpy, non-deterministic, irreversible

# Two laws of evolution

Quantum measurement problem



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Assumptions  
of  
Physics

# What is quantum mechanics about?

Hidden variables?

Consistent histories?

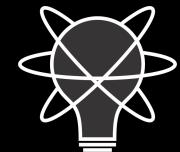
## Interpretations

Many worlds?

Agent information?

Consciousness?

## Quantum reconstructions

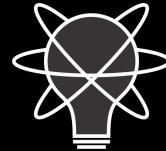


<https://assumptionsofphysics.org/>

Assumptions  
of  
Physics

# What is the ultimate reason classical mechanics fails, and how is it fixed by quantum mechanics?

Since some of the classical failures are related to thermodynamics,  
let's look at the relationship between entropy and uncertainty in  
classical physics



# Entropy in classical mechanics

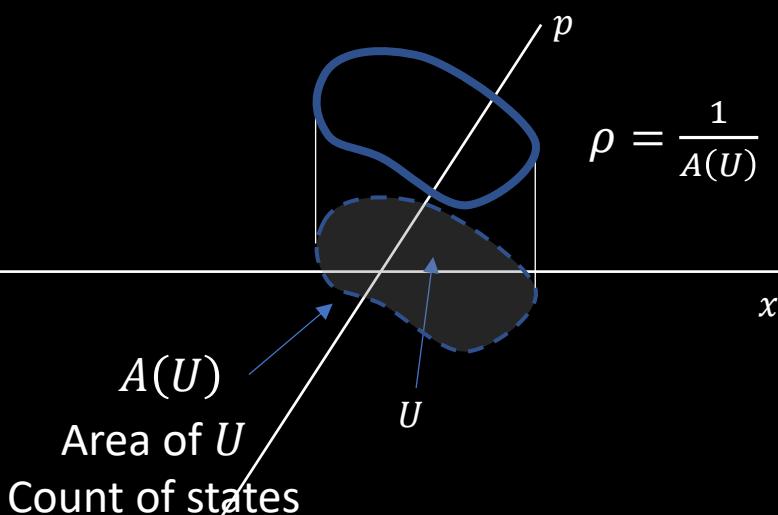
$$S(\rho) = - \int \rho \log \rho dx dp$$

Entropy

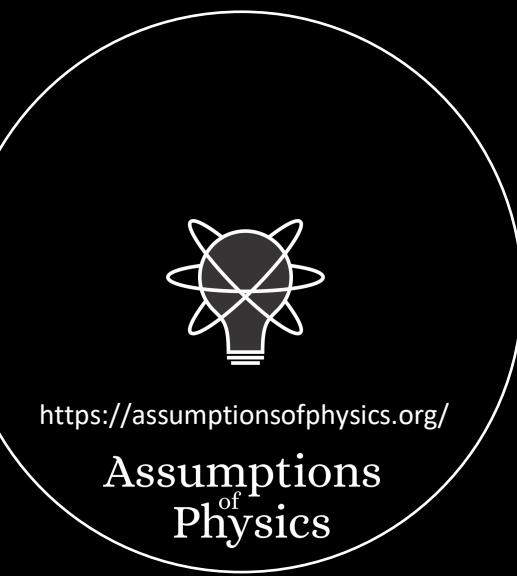
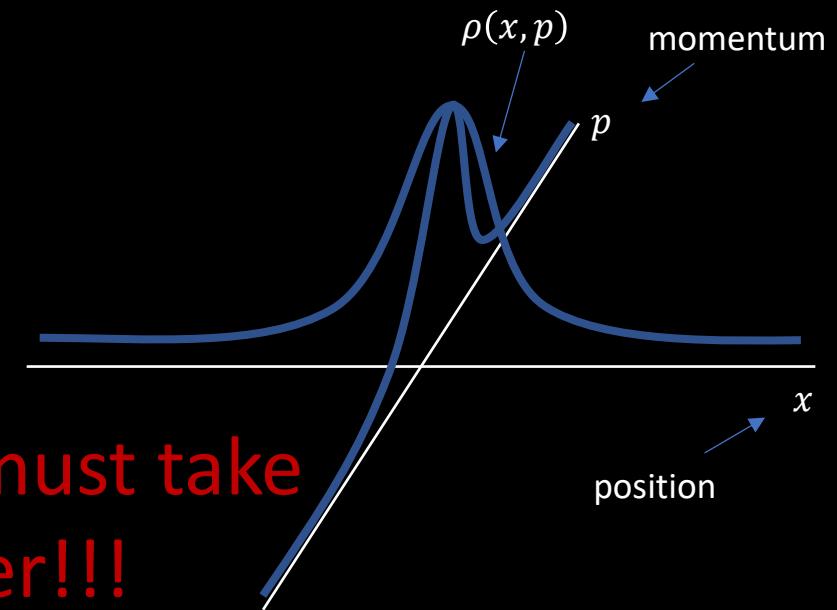
Probability distribution

$$S(\rho_U) = \log A(U)$$

Uniform distribution



Logarithm must take  
pure number!!!



# Entropy in classical mechanics

$$S(\rho) = - \int \rho \log h \rho \, dx dp$$

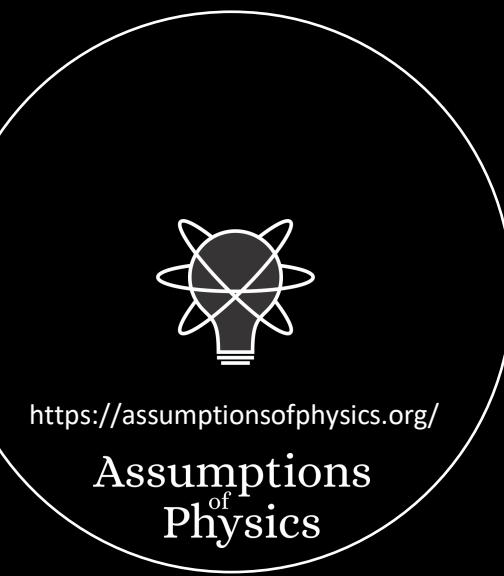
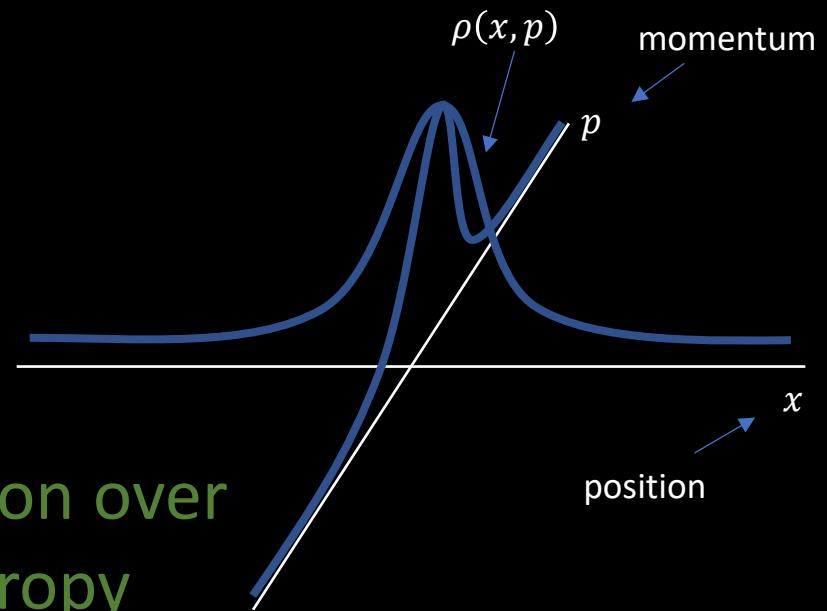
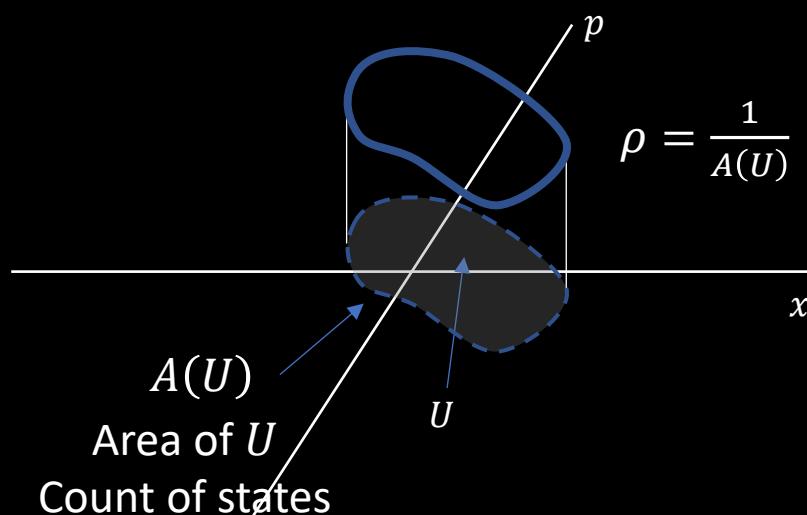
Entropy

Probability distribution

$$S(\rho_U) = \log \frac{A(U)}{h}$$

Uniform distribution

Uniform distribution over area  $h \Rightarrow$  zero entropy



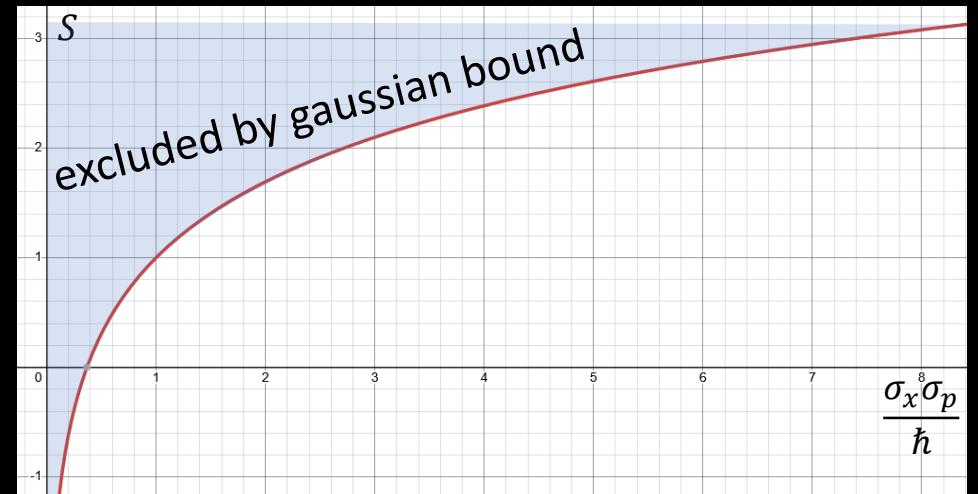
# Entropy vs uncertainty

Gaussian maximizes entropy for a given uncertainty

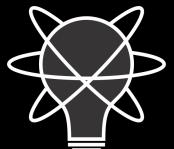
$$S(\rho) \leq \log 2\pi e \frac{\sigma_x \sigma_p}{h}$$

$$\sigma_x \sigma_p \geq \frac{h}{2\pi e} e^{S(\rho)} = \frac{\hbar}{e} e^{S(\rho)}$$

Entropy puts a lower bound  
on the uncertainty



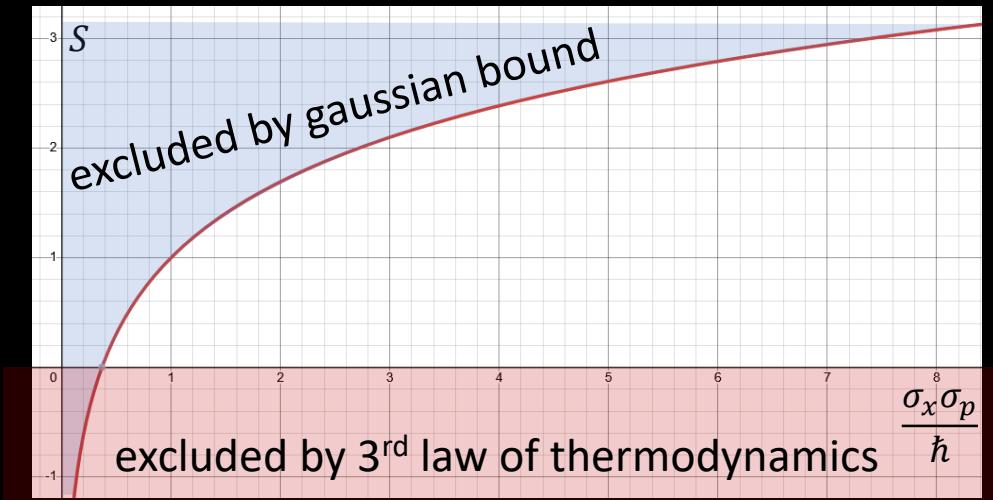
Is there a universal lower bound to the entropy?



# Enter the 3<sup>rd</sup> law

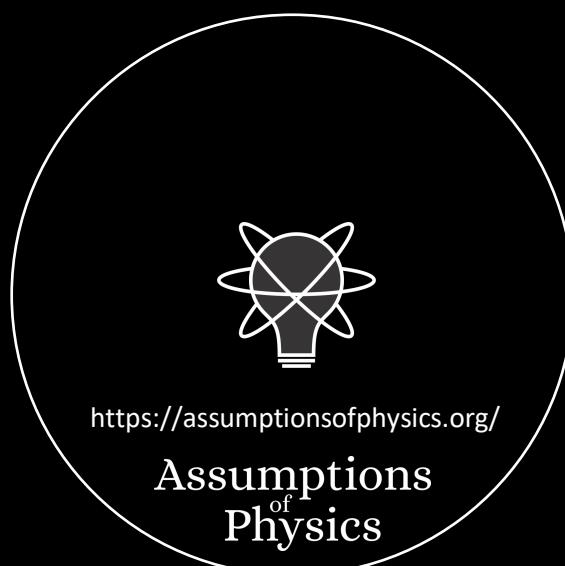
Every substance has a finite positive entropy, but at the absolute zero of temperature the entropy may become zero, and does so become in the case of perfect crystalline substances.

G. N. Lewis and M. Randall, Thermodynamics and the free energy of chemical substances (McGraw-Hill, 1923)



$$S \geq 0 \Rightarrow \sigma_x \sigma_p \geq \frac{\hbar}{e}$$

Classical ~~uncertainty~~ principle  
indetermination



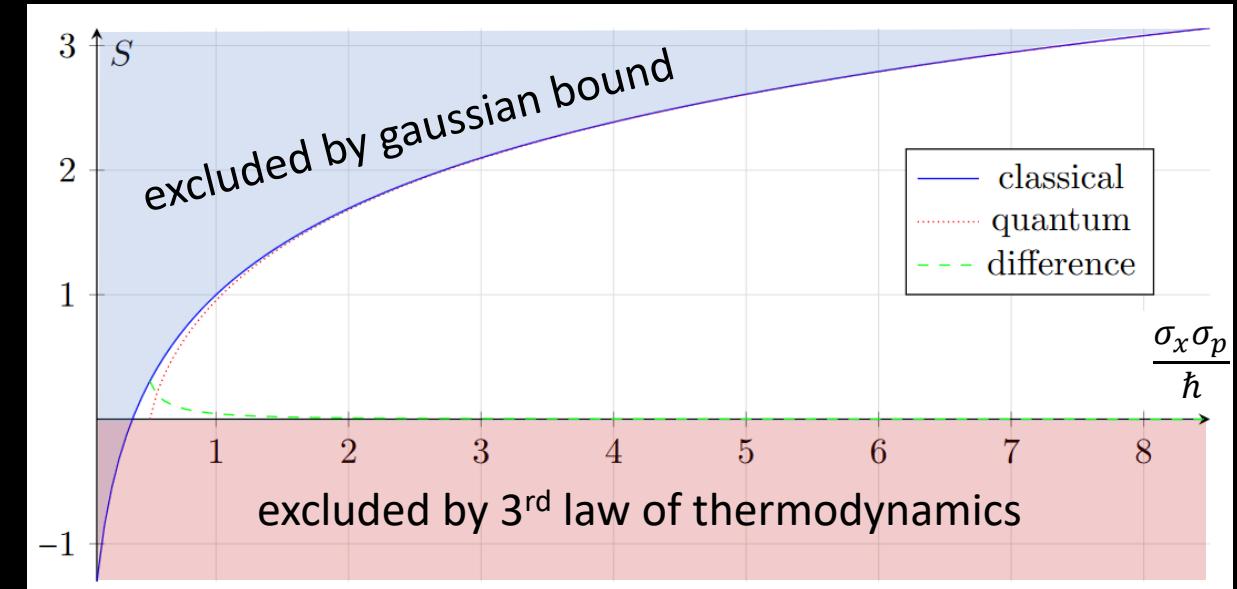
# Comparing theories

$$\text{classical} \quad \sigma_x \sigma_p \geq -\frac{\hbar}{e}$$
$$\text{quantum} \quad \sigma_x \sigma_p \geq -\frac{\hbar}{2}$$

2.71828...

The gaussian bound quickly becomes very similar across theories

Entropy of quantum states is already non-negative

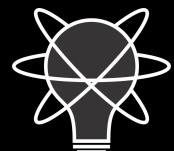


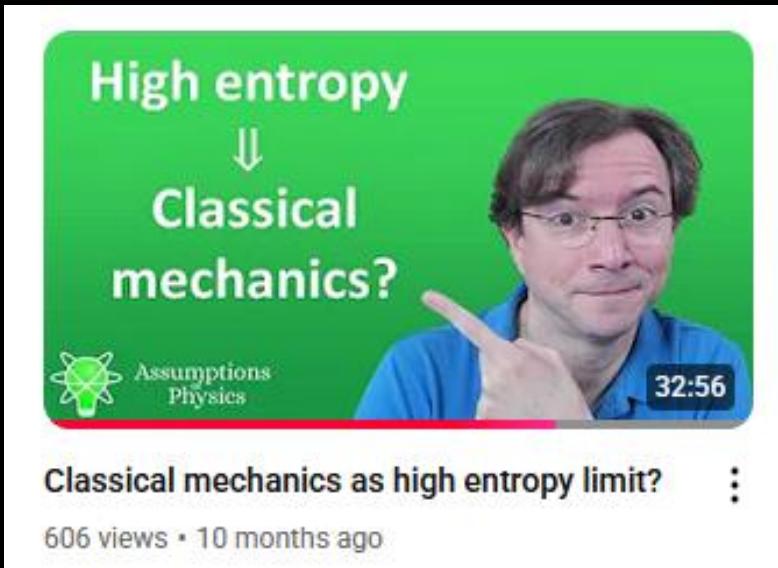
$$S_C = \ln e\sigma$$

$$S_Q = \left(\sigma + \frac{1}{2}\right) \ln \left(\sigma + \frac{1}{2}\right) - \left(\sigma - \frac{1}{2}\right) \ln \left(\sigma - \frac{1}{2}\right)$$

Quantum mechanics incorporates the third law  
Classical mechanics does not

Conjecture: does quantum mechanics recover  
classical mechanics at high entropy?





greetings

ML Manuele Landini <manulando@gmail.com>  
To: carassi@umich.edu

(i) You replied to this message on 7/10/2024 10:00 AM.

Caro Gabriele,

Mi chiamo Manuele Landini e lavoro a Innsbruck (Austria) come senior scientist in un gruppo di fisica atomica sperimentale. Puoi vedere di cosa ci occupiamo sul nostro sito: <https://quantummatter.at>.

Ho visto un po' dei tuoi video su youtube. Mi sembra un progetto molto ambizioso, ma promettente. Mi farebbe piacere riuscire a spiegare agli studenti in futuro in termini piu' fisici concetti come le sovrapposizioni o il teorema spin-statistica.

Per la storia della metrica, da quel che ho capito hai bisogno di una metrica che non sia basata sull'entropia, visto che vuoi definire una distanza a entropia costante. Ci sono varie opzioni, ma la trace distance [Trace distance - Wikipedia](#) funziona perche' ha una propriet'a fondamentale che puoi usare. Chiamala:  $T(\rho, \sigma)$

Se parti da stati puri, si riduce a  $(1 - \langle \psi | \phi \rangle)^{1/2}$ . Quindi per massimizzarla, scegli due stati ortogonali (non importa quali). Il massimo e'  $T_0=1$ . Una volta che hai questi stati, che hanno entropia 0, li puoi trasformare in stati con entropia finita (in particolare quelli con massima distanza) tramite una trace preserving map  $M$ .

Siccome  $T$  si contrae, hai che  $T(M(\rho), M(\sigma)) \leq T(\rho, \sigma)$ . L'uguale vale se la mappa e' unitaria. Cosi' definisci un serie di step in cui la distanza massima decresce  $T_{n+1} < T_n$ , fino ad arrivare a 0 per stati fully mixed.

arXiv > quant-ph > arXiv:2411.00972

Quantum Physics

[Submitted on 1 Nov 2024 (v1), last revised 3 Dec 2024 (this version, v2)]

## Classical mechanics as the high-entropy limit of quantum mechanics

Gabriele Carassi, Manuele Landini, Christine A. Aidala

We show that classical mechanics can be recovered as the high-entropy limit of quantum mechanics. That is, the high entropy masks quantum effects, and mixed states of high enough entropy can be approximated with classical distributions. The mathematical limit  $\hbar \rightarrow 0$  can be reinterpreted as setting the zero entropy of pure states to  $-\infty$ , in the same way that non-relativistic mechanics can be recovered mathematically with  $c \rightarrow \infty$ . Physically, these limits are more appropriately defined as  $S \gg 0$  and  $v \ll c$ . Both limits can then be understood as approximations independently of what circumstances allow those approximations to be valid. Consequently, the limit presented is independent of possible underlying mechanisms and of what interpretation is chosen for both quantum states and entropy.

Comments: 14 pages, 3 figures  
 Subjects: Quantum Physics (quant-ph)  
 Cite as: arXiv:2411.00972 [quant-ph]  
 (or arXiv:2411.00972v2 [quant-ph] for this version)  
<https://doi.org/10.48550/arXiv.2411.00972> ⓘ

Submission history

From: Christine Aidala [view email]  
 [v1] Fri, 1 Nov 2024 18:48:04 UTC (19 KB)  
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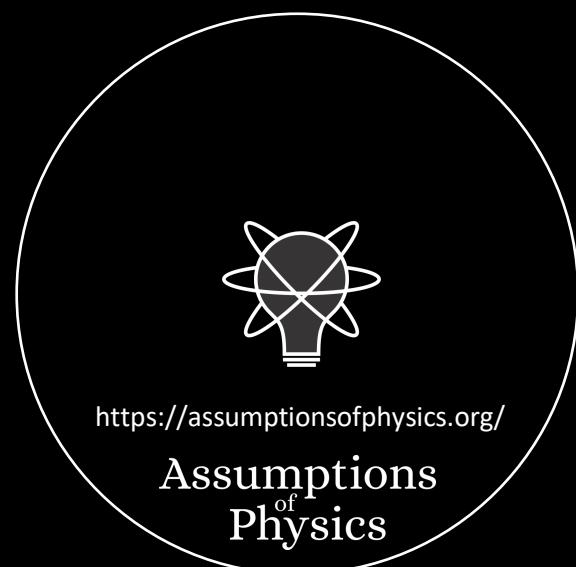
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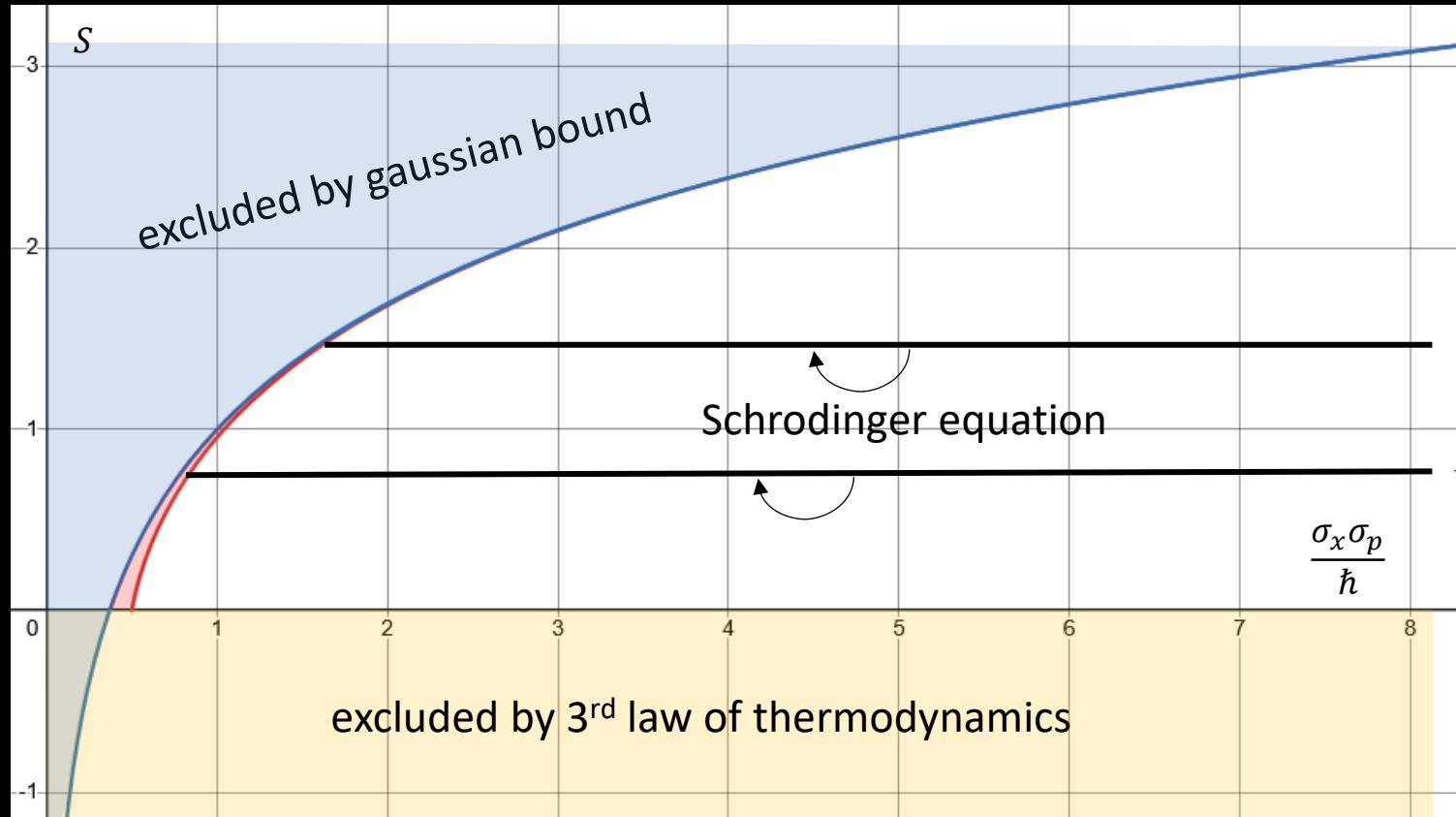
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— classical  
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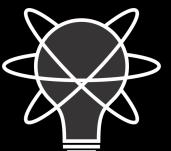
$$\frac{[x, p]}{i\hbar} = 1 \Rightarrow \frac{[T(x), T(p)]}{i\hbar} = \lambda$$

$$\sigma_{T(x)} \sigma_{T(p)} \geq \frac{1}{2} |\langle [T(x), T(p)] \rangle| = \frac{\hbar}{2} \lambda$$

1. Increase the entropy of all states

$T$  - Entropy increasing map

Spread increase in position and momentum

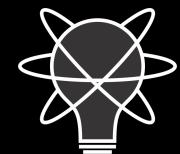


# Mathematical limit

$$\frac{[x,p]}{i\hbar} \xrightarrow{\hbar \rightarrow 0} \{x, p\}$$

Makes no physical sense!

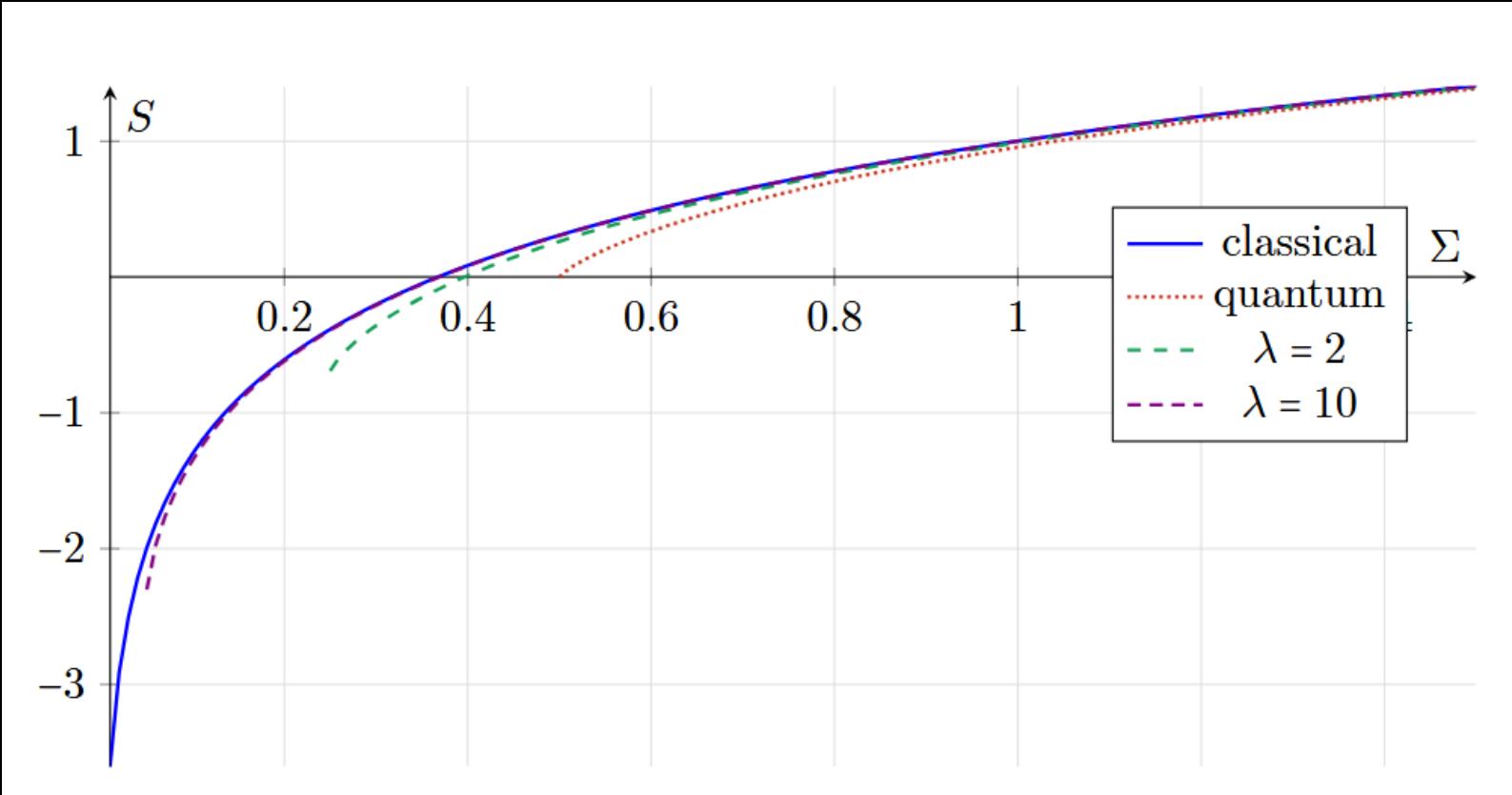
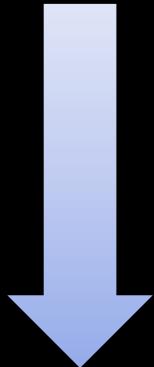
$\hbar$  is a constant: it doesn't change



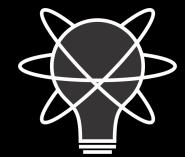
<https://assumptionsofphysics.org/>

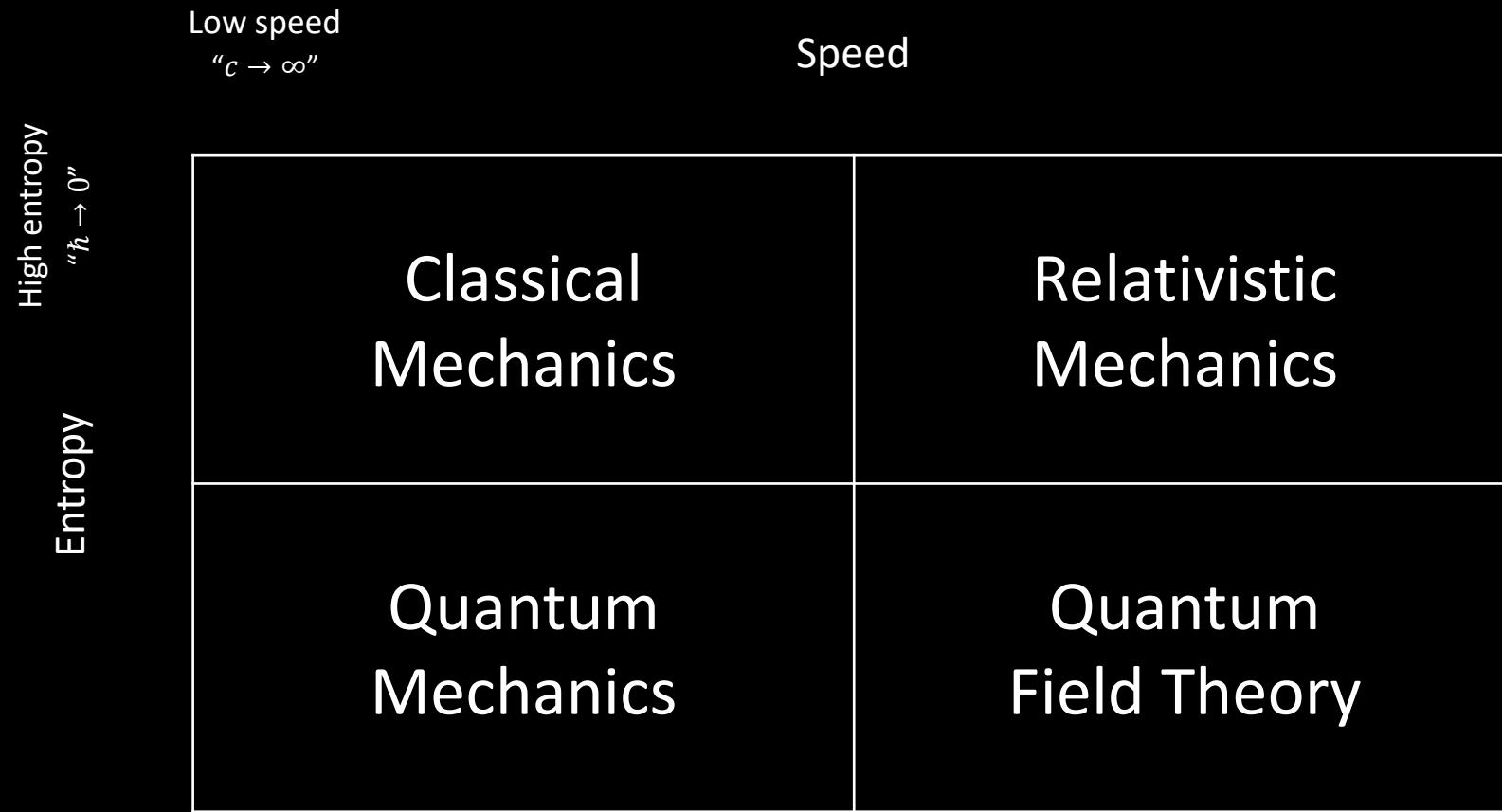
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## 2. Decrease the entropy of pure states

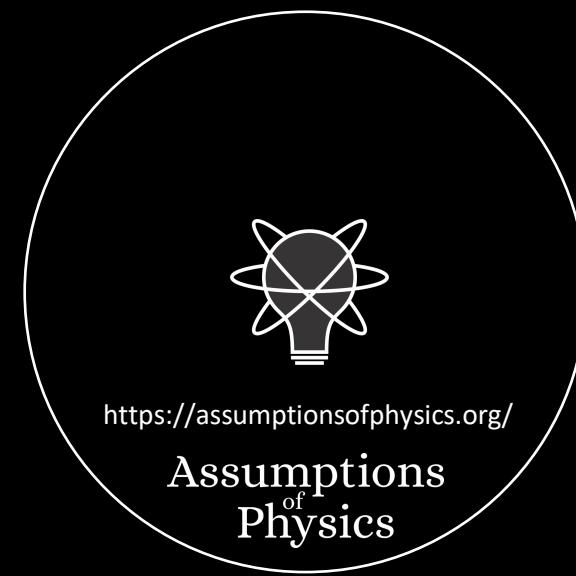


$$\frac{[X,P]}{\imath\hbar} = \frac{1}{\lambda} \quad \frac{[X,P]}{\imath(\hbar/\lambda)} = 1 \xrightarrow[\hbar/\lambda \rightarrow 0]{\lambda \rightarrow \infty} \{x,p\} = 1$$





No-mechanism limit  
(same as non-relativistic limit)



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Physics

Quantum to classical

$$\frac{[A, B]}{\imath\hbar} \rightarrow \{A, B\}$$

High entropy limit

Classical to quantum

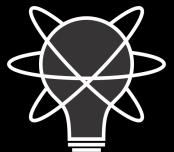
$$\{A, B\} \rightarrow \frac{[A, B]}{\imath\hbar}$$

???

Up to isomorphisms, the Moyal bracket is the unique one-parameter Lie-algebraic deformation of the Poisson bracket

If I asked for a theory that puts a hard lower bound on the entropy, and recovers classical mechanics at high entropy....

there is only one way to do it



Quantum to classical

$$\frac{[A, B]}{i\hbar} \rightarrow \{A, B\}$$

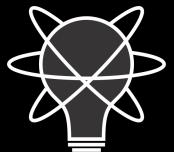
High entropy limit

Classical to quantum

$$\{A, B\} \rightarrow \frac{[A, B]}{i\hbar}$$

Entropic lower bound  
that recovers classical mechanics

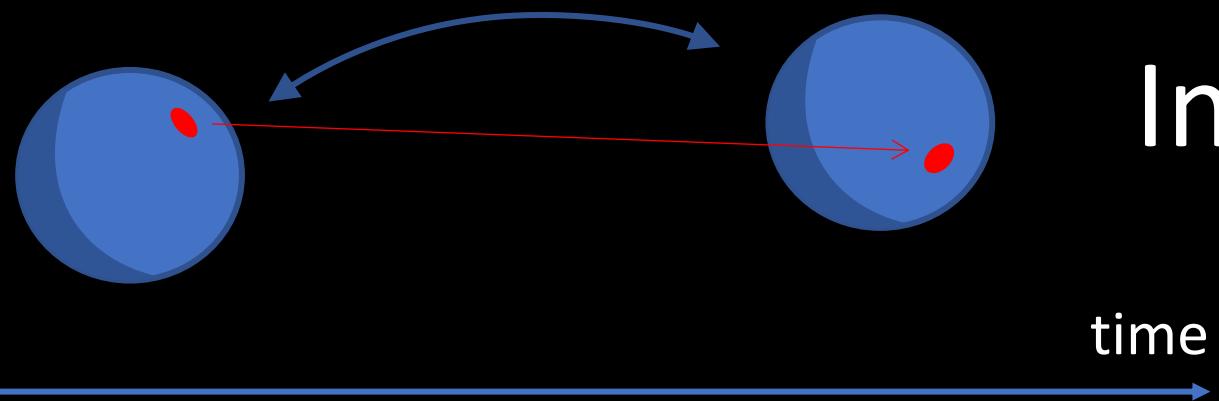
Quantizing a classical theory  
means putting a lower bound  
on the entropy



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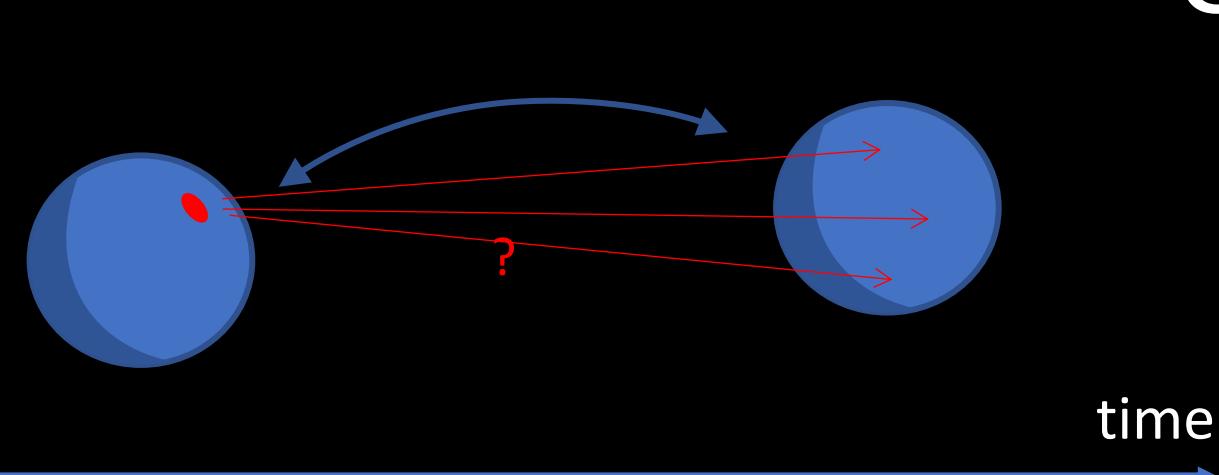
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# Classical system



Infinitesimally reducible

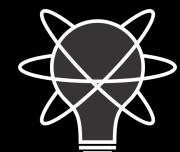
Internal dynamics  
perfectly accessible



# Quantum system

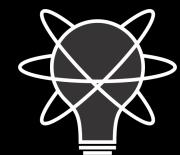
Irreducible

Internal dynamics  
inaccessible



# What is entropy?

## Why a lower bound?



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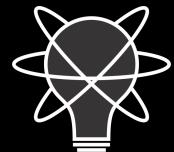
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# What is entropy?

You can think of it as disorder

You can think of it as information

You can think of it as  
lack of information



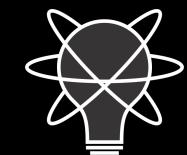
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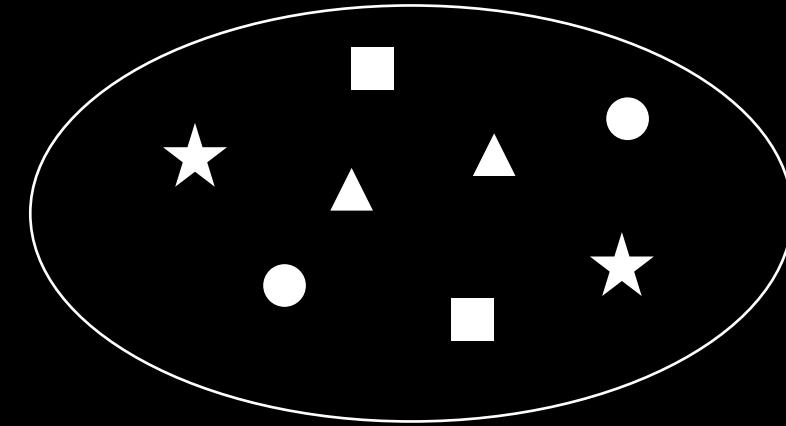
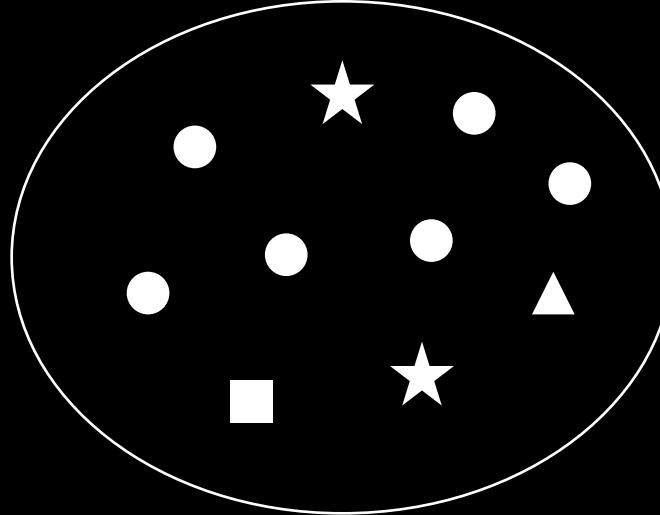
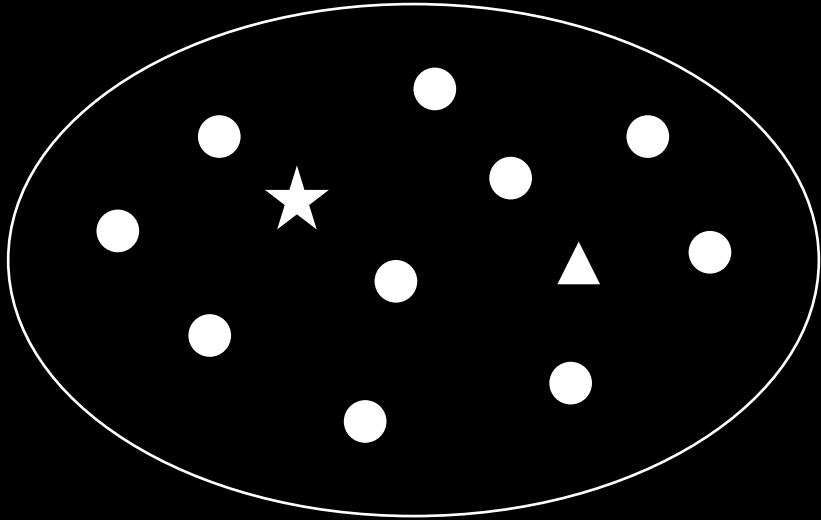
Is or is not.

There is no “you can think of it as.”



<https://assumptionsofphysics.org/>

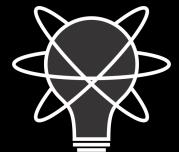
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Suppose you have a (possibly infinite) collection of things

**What is the variability of the collection?**

How different are the elements within the collection?



Given a discrete distribution  $p_i$   
we want an indicator  $I(p_i)$  of variability such that:

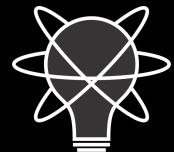
1. Continuous in  $p_i$
2. Increases if the number of cases increases
3. Linear in probability

$$I(p_1 \hat{p}_i, p_2 \bar{p}_j) = I(p_1, p_2) + p_1 I(\hat{p}_i) + p_2 I(\bar{p}_j)$$

⇒ Then  $I(p_i) = -\sum_i p_i \log p_i$

Proof can be generalized to classical continuum and quantum mechanics

Entropy is variability!



# Variability as a better characterization of Shannon entropy

Gabriele Carcassi, Christine A Aidala and Julian Barbour

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[European Journal of Physics, Volume 42, Number 4](#)

[Machine learning to detect schedules using spatiotemporal data of behavior: A proof of concept](#)

Marc J. Lanovaz, Varsovia Hernandez, Alejandro León

2025, Journal of the Experimental Analysis of Behavior - Article

[Why so slow? Models of parkinsonian bradykinesia](#)

David Williams

2024, Nature Reviews Neuroscience - Article

[Using Entropy Metrics to Analyze Information Processing Within Production Systems: The Role of Organizational Constraints](#)

Frits van Merode, Henri Boersma, Fleur Tournois, Windi Winasti, Nelson Aloysio Reis de Almeida Passos, Annelies van der Ham

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[Internal Versus Forced Variability Metrics for General Circulation Models Using Information Theory](#)

Aakash Sane, Baylor Fox-Kemper, David S. Ullman

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[Parameterizing Vertical Mixing Coefficients in the Ocean Surface Boundary Layer Using Neural Networks](#)

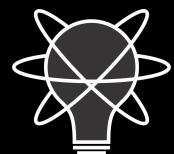
Aakash Sane, Brandon G. Reichl, Alistair Adcroft, Laure Zanna

2023, Journal of Advances in Modeling Earth Systems - Article

[Spatiotemporal characteristics of future precipitation variability in the Tianshan Mountain region of China](#)

Xianglin Lyu, Junkai Du, Yaqin Qiu, Yangwen Jia, Chunfeng Hao, Hao Dong

2025, Journal of Hydrology Regional Studies - Article



<https://assumptionsofphysics.org/>

Assumptions  
of  
Physics

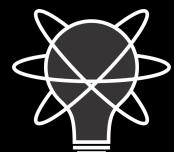
Entropy is not the property  
of an element, but of a distribution

A distribution of what?

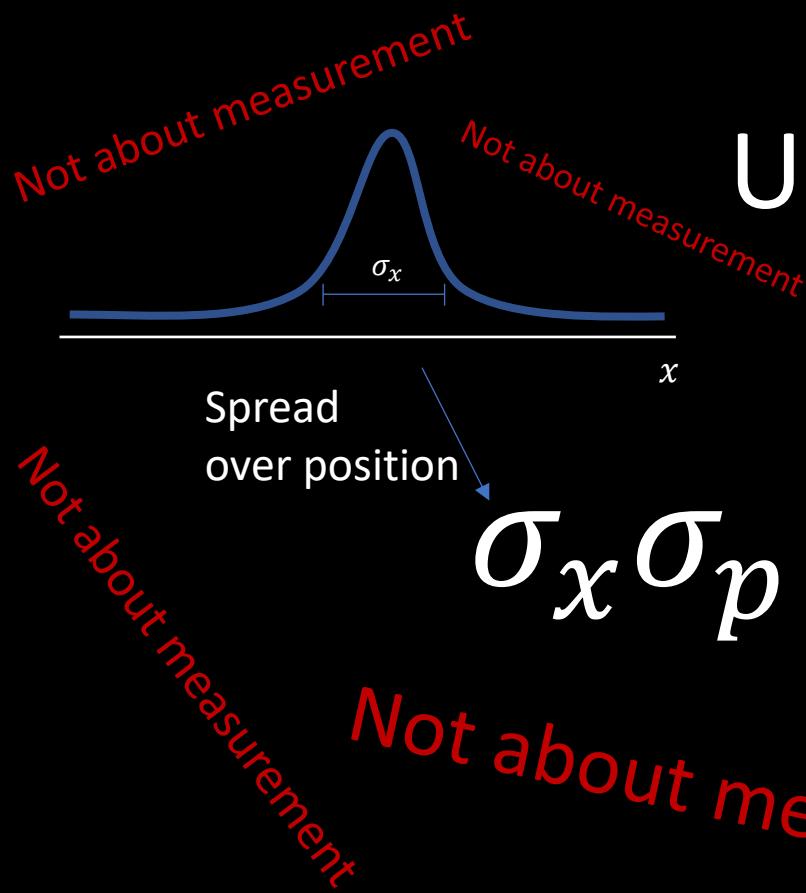
A statistical distribution of  
the molecules of a gas?

A credence distribution of what  
an agent may think is there?

A probability distribution of the  
outcomes of a measurement?

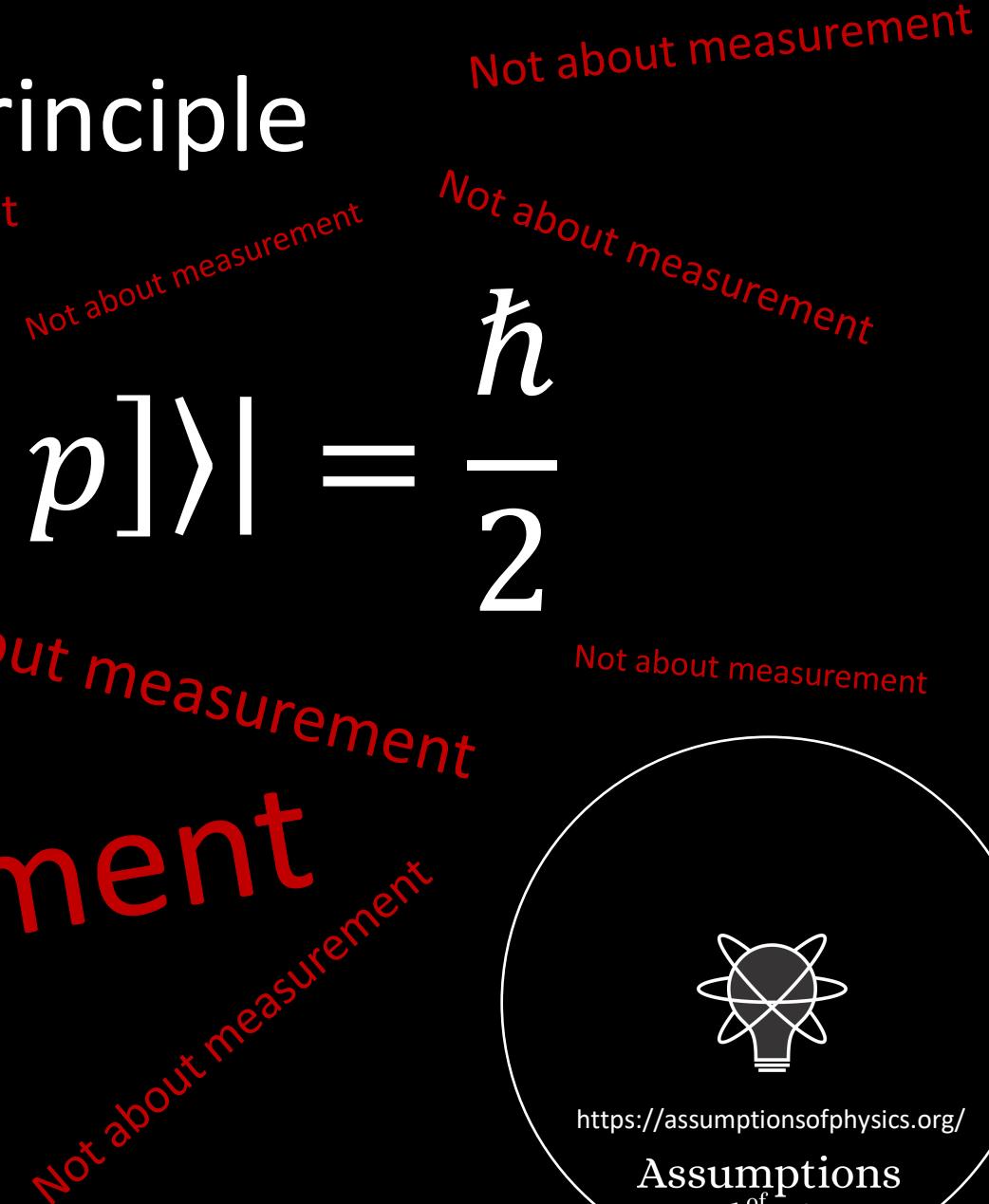


# Uncertainty principle

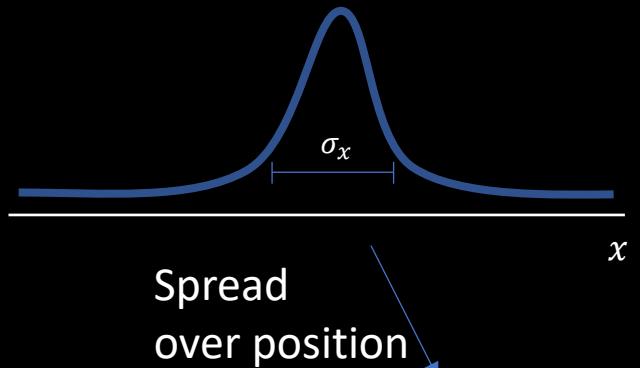


$$\sigma_x \sigma_p \geq \frac{1}{2} |\langle [x, p] \rangle| = \frac{\hbar}{2}$$

Not about measurement

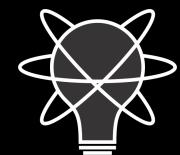


# Indetermination ~~Uncertainty~~ principle

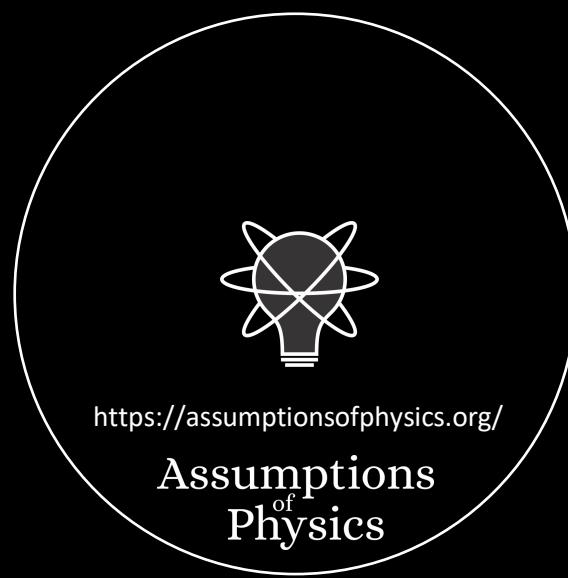
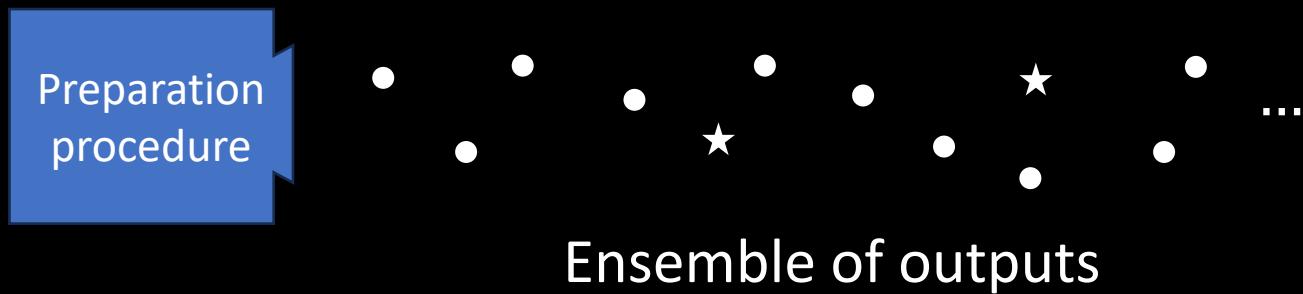


$$\sigma_x \sigma_p \geq \frac{1}{2} |\langle [x, p] \rangle| = \frac{\hbar}{2}$$

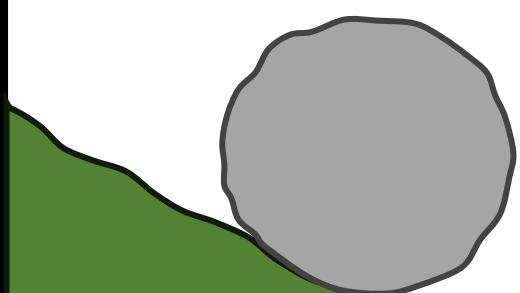
About preparation  
not measurement!



Entropy is the variability of an ensemble:  
the collection of all outputs  
of a reliable preparation procedure

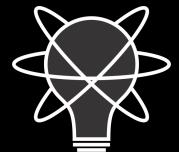


$t$



Physics describes the evolution  
of a particular system like a movie

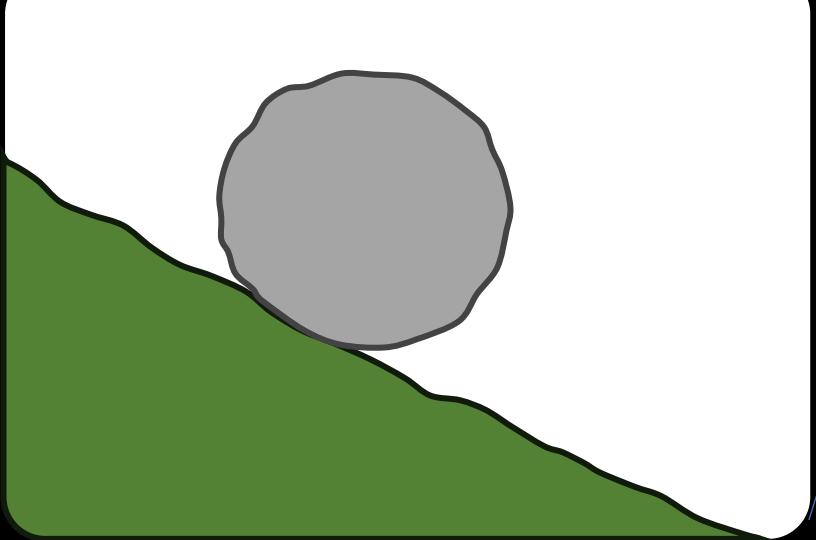
$x(t)$



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Assumptions  
of  
Physics

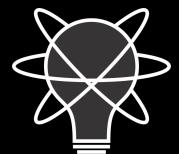
$t$



~~Physics describes the evolution  
of a particular system like a movie~~

-100 points

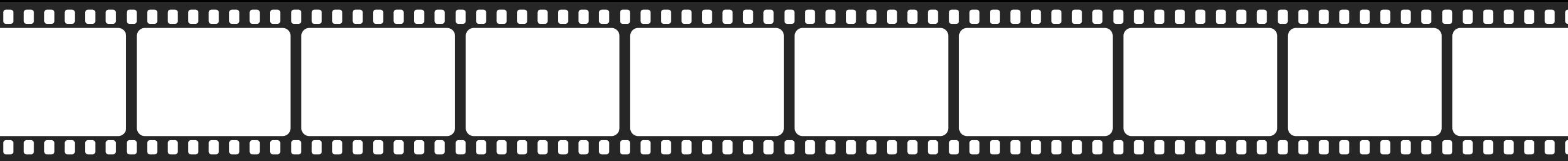
$x(t)$



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Assumptions  
of  
Physics

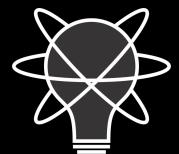
$\rightarrow t$



Physics describes the evolution  
of similarly prepared systems

Every time I do this...

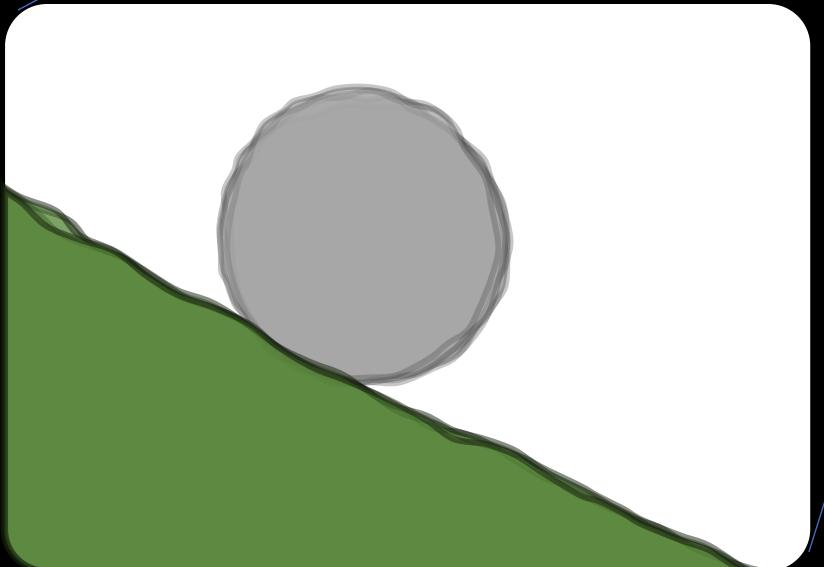
$x(t)$



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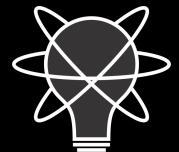
Assumptions  
of  
Physics

$t$



Physics describes the evolution  
of similarly prepared systems

$x(t)$



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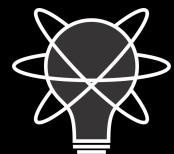
Assumptions  
of  
Physics

Statistical description of preparations  
and measurement outcomes

$$x(t)$$

Statistical description  
of all clocks synchronized  
at a particular initial event

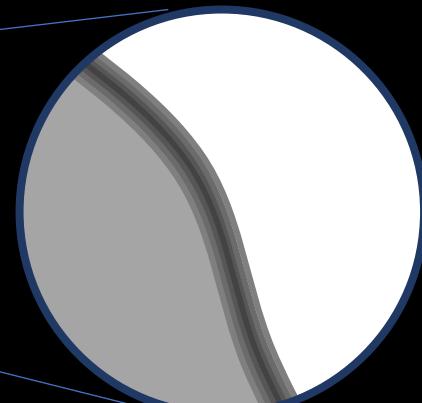
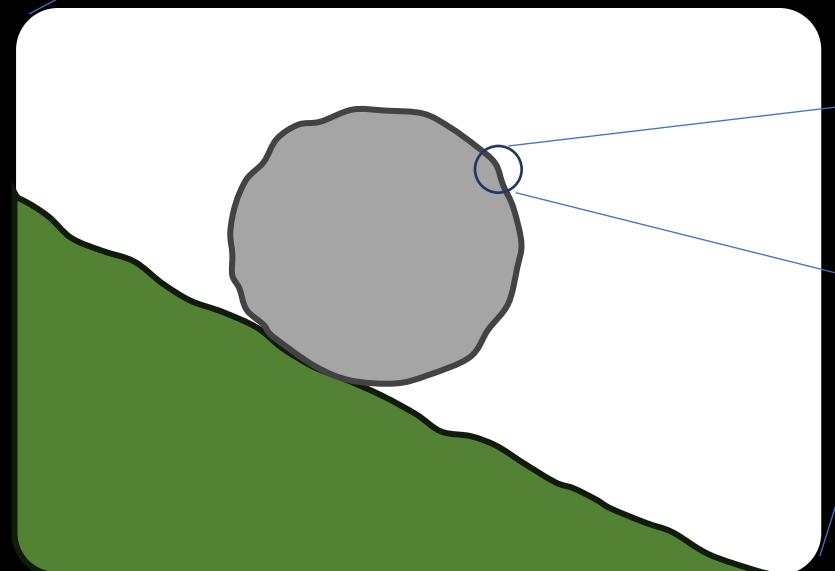
In physics, we can only  
study relationships that  
can be reproduced:  
relationships between  
ensembles



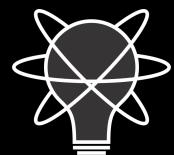
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Assumptions  
of  
Physics

$\rightarrow t$



Even a single measurement is an average over a finite time



To be able to define a probability distribution over heads or tails, you need first to define what heads and tails are

We can do that because we can leave the coin be, and it remains heads or tails

If it kept changing more than 100 times a second, we wouldn't perceive heads or tails: just a blur

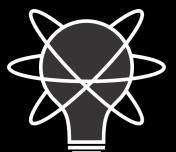


wikipedia

wikipedia

The state of a system is what remains “stable” for at least a short amount of time

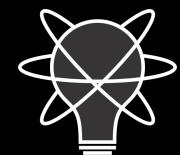
It's ensembles all the way down



# What is entropy?

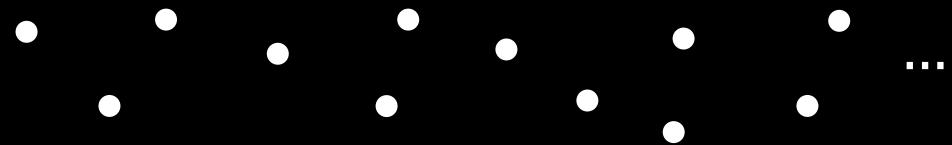
Variability within a preparation!

# Why a lower bound?



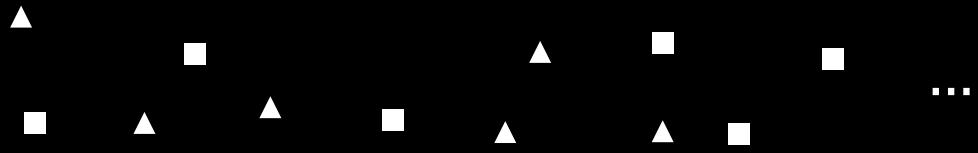
# Classical mechanics

Preparation procedure



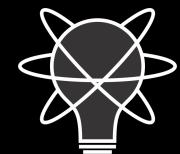
Assumes we can have perfect preparations

Preparation procedure



Concedes that there are no perfect preparations

# Quantum mechanics



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Assumptions  
of  
Physics

# For classical systems

$$S(\rho_U) = \log \frac{A(U)}{h}$$

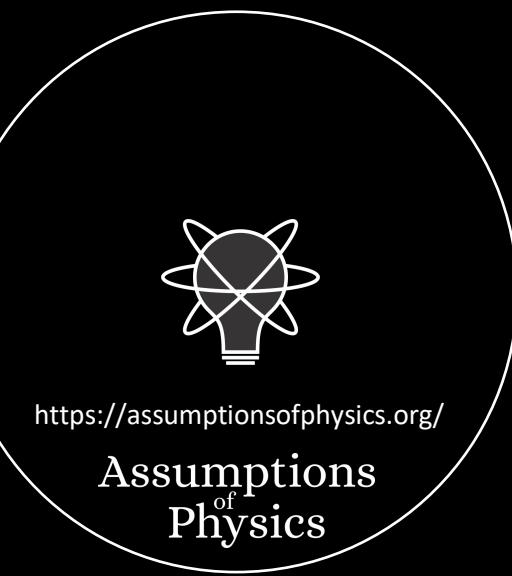
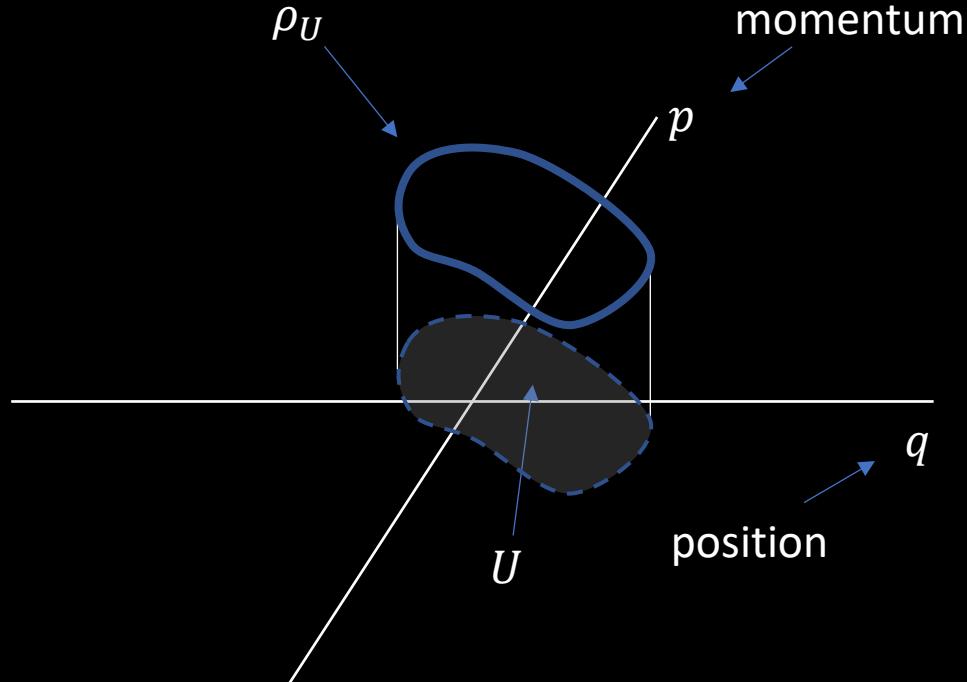
Entropy for a uniform distribution over  $U$

Logarithm of the count of states in  $U$

$$S(\rho_U) < 0$$

$$\frac{A(\rho_U)}{h} < 1$$

Less than one state!



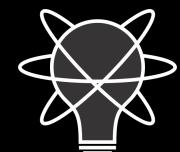
# What is entropy?

Variability within a preparation!

# Why a lower bound?

Our preparations cannot be perfect

Regions with “less than one state”  
make no sense



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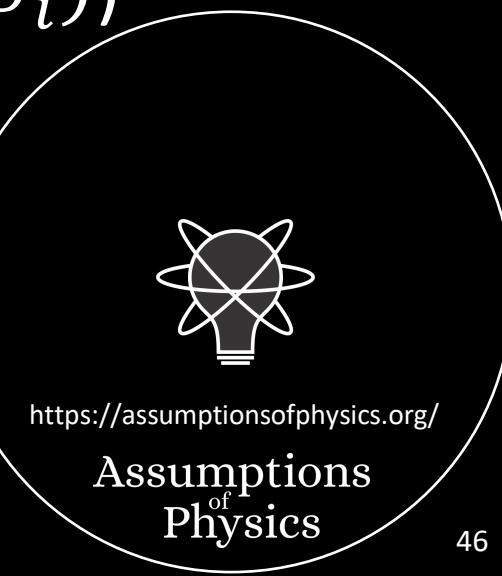
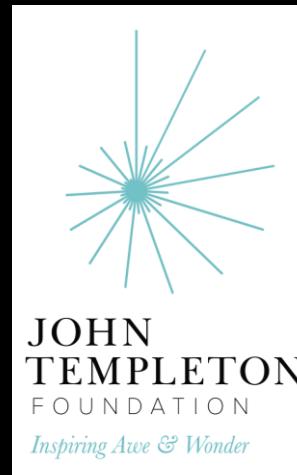
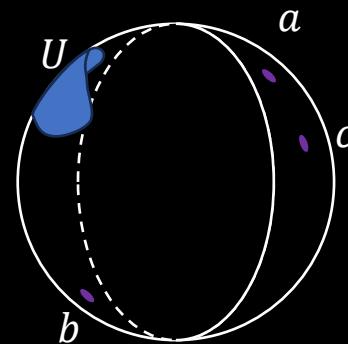
Assumptions  
of  
Physics

# The problem with counting (on the continuum)

measure  
→  $\mu(\text{ a set of states}) = \text{ how many states there are in the set}$

1. Every state counts as one (i.e.,  $\mu(\{c\}) = 1$ )
  2. Finite regions have finitely many states (i.e.,  $\mu([a, b]) < \infty$ )
  3. Count is additive for disjoint sets (i.e.,  $\mu(\bigcup U_i) = \sum \mu(U_i)$ )
- Incompatible! Pick two!**

On a continuum, a finite region has infinitely many states!



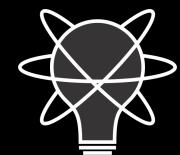
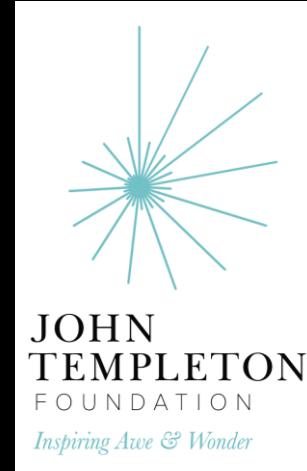
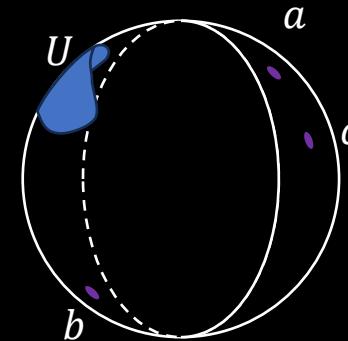
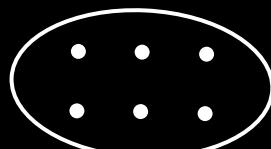
# The problem with counting on the continuum

$\mu$  ( a set of states ) = how many states are there in the set

1. Every state counts as one (i.e.  $\mu(\{a\}) = 1$ )
2. ~~Finite regions have finitely many states (i.e.  $\mu(U) < \infty$ )~~
3. Count is additive for disjoint sets (i.e.  $\mu(\cup U_i) = \sum \mu(U_i)$ )

⇒ Counting measure

What we use in classical discrete case



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Assumptions  
of Physics

# The problem with counting on the continuum

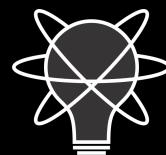
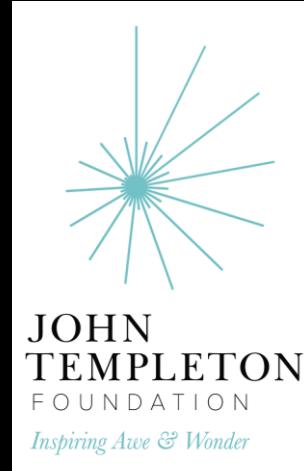
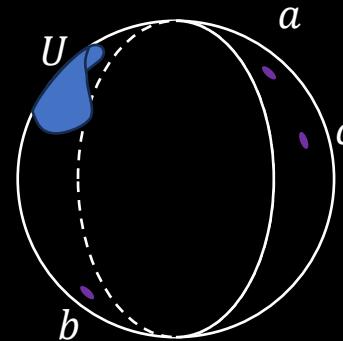
$\mu$  ( a set of states ) = how many states are there in the set

- ~~1. Every state counts as one (i.e.  $\mu(\{a\}) = 1$ )~~
- 2. Finite regions have finitely many states (i.e.  $\mu(U) < \infty$ )
- 3. Count is additive for disjoint sets (i.e.  $\mu(\cup U_i) = \sum \mu(U_i)$ )

⇒ Lebesgue measure

What we use in classical continuum case (i.e. classical mechanics)

$\Delta x \Delta p$



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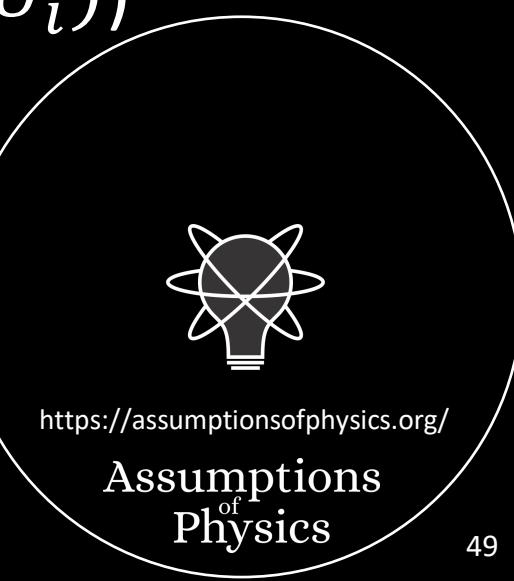
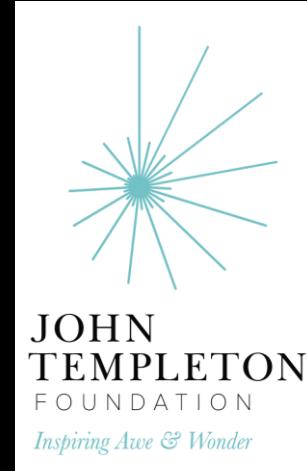
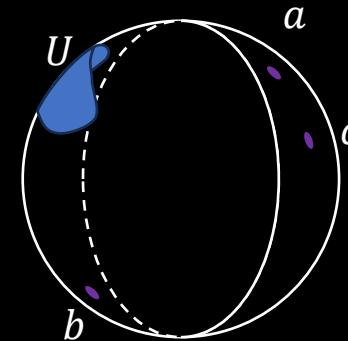
# The problem with counting on the continuum

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1. Every state counts as one (i.e.  $\mu(\{a\}) = 1$ )
2. Finite regions have finitely many states (i.e.  $\mu(U) < \infty$ )
3. Count is additive for disjoint sets (i.e.  $\mu(\bigcup U_l) = \sum \mu(U_l)$ )

⇒ “Quantum” measure

What we are implicitly using in quantum mechanics

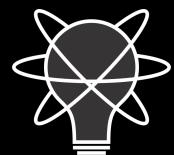


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Assumptions  
of Physics

We saw that thinking about entropy makes us understand how classical mechanics fails and gives rise to quantum mechanics

Can we use similar arguments to understand why quantum mechanics and general relativity are incompatible, and how to fix it?



# General relativity is a field theory

## Particle mechanics

State defined by  
position and momentum:

$$\begin{matrix} x^1 & x^2 & \dots & x^n \\ p_1 & p_2 & \dots & p_n \end{matrix}$$

Finitely many variables

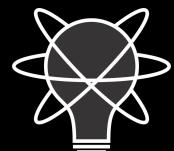
## Field theory

State defined by  
fields at each point:

$$A(x) \ B(x) \ \dots$$

Continuously (infinitely)  
many variables

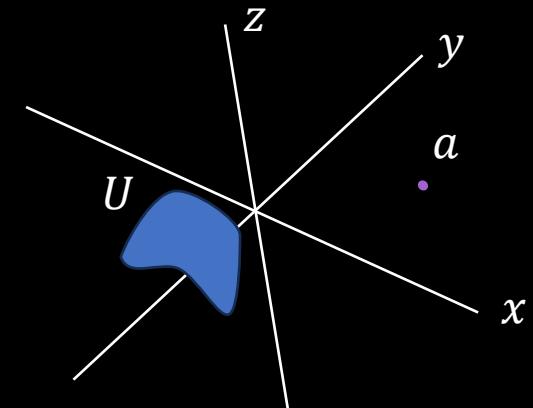
How do we count them?!?



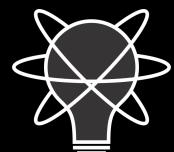
# The problem with counting (on the continuum)

$$\mu(\text{region of space}) = \text{Volume proportional to number of variables}$$

1. Every point has one variable (i.e.  $\mu(\{a\}) = 1$ )
  2. Finite regions have finite volume (i.e.  $\mu(U) < \infty$ )
  3. Count is additive for disjoint sets (i.e.  $\mu(\cup U_i) = \sum \mu(U_i)$ )
- Incompatible!**

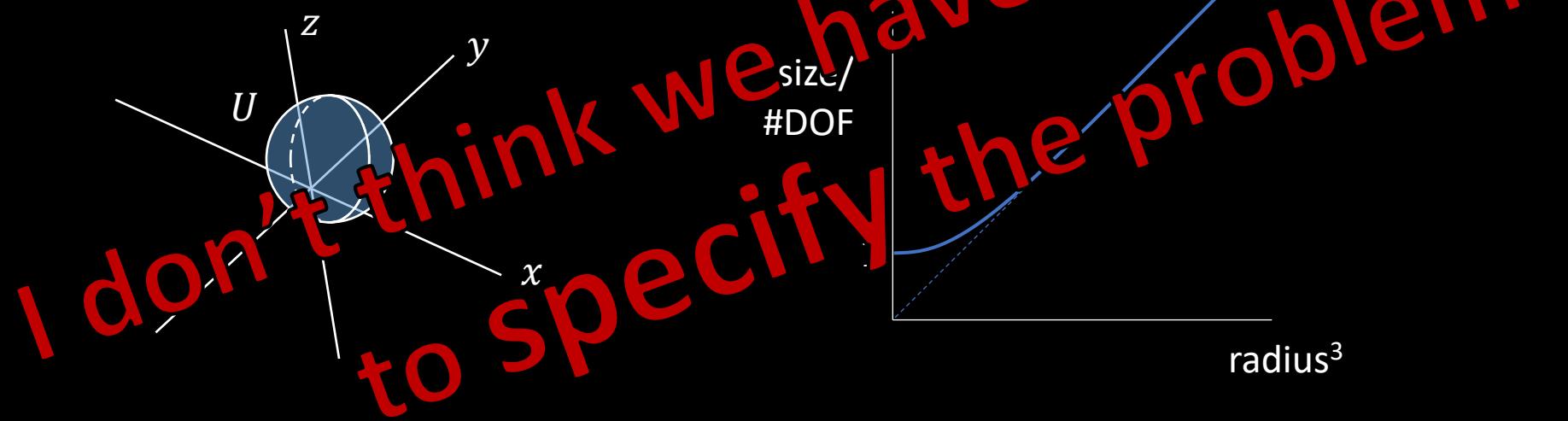


On a continuum, a finite region has infinitely many points!

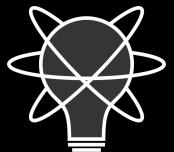


Seems that “quantizing space-time” means putting a lower bound on the number of variables (degrees of freedom)

In the same way that a region of state space “with less than one state” makes no sense, a region of physical space with “less than one degree of freedom” makes no sense



I don't think we have the right math to solve this problem!

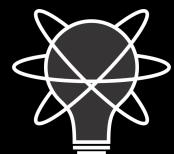


But we don't develop theories by writing down assumptions and then derive observable consequences in a sequence of theorems and proofs. In physics, theories almost always start out as loose patchworks of ideas. Cleaning up the mess that physicists generate in theory development, and finding a neat set of assumptions from which the whole theory can be derived is often left to our colleagues in mathematical physics—a branch of mathematics, not of physics.

Sabine Hossenfelder – Lost in Math

-100,000 points

Math is not just a tool for calculation:  
we use it to specify our theories



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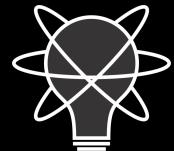
Assumptions  
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Physics

Ultimately, it is up to the physicists to develop a clear physical model

Mathematicians can help find holes... and give options

Clear physical model  $\Leftrightarrow$  clear math

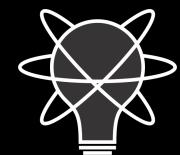
Clarity, meaningfulness, consistency, ...  
cannot be added after the fact



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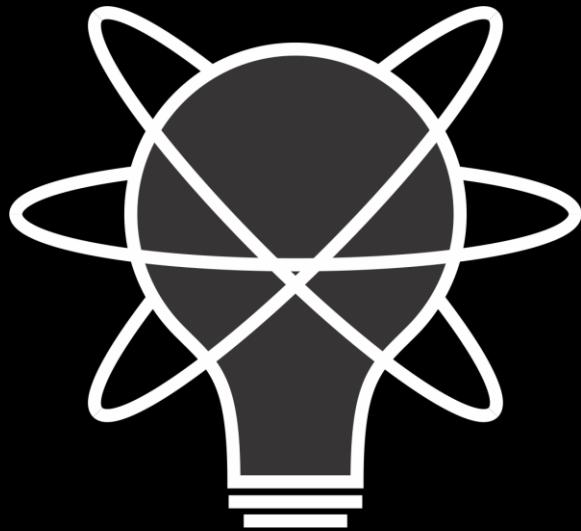
Big systems that work always  
come from small systems that  
work, they never come from big  
systems that don't work



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Assumptions  
of  
Physics

The only way forward is to go back and reconstruct everything from a minimal set of physically meaningful assumptions



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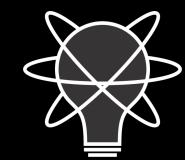
<https://assumptionsofphysics.org>

For more information, papers, presentations, ...

<https://www.youtube.com/@gcarcassi>

Videos with results and insights from the research

<https://assumptionsofphysics.org/foundation>



<https://assumptionsofphysics.org/>

Assumptions  
of  
Physics

# The problem with counting on the continuum

We'd like to say:

1. Every state is a single case (i.e.  $\mu(\{\psi\}) = 1$ )
2. Finite continuous range carries finite information (i.e.  $\mu(U) < \infty$ )
3. Count is additive for disjoint sets (i.e.  $\mu(\bigcup U_i) = \sum \mu(U_i)$ )

Incompatible!

Pick two!

Discard 1  $\Rightarrow$  Lebesgue measure

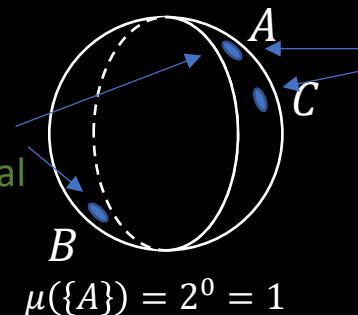
Discard 2  $\Rightarrow$  counting measure

Discard 3  $\Rightarrow$  “Quantum measure”

$$\mu(U) = 2^{\sup(S(\text{hull}(U)))}$$

Exponential of the maximum entropy reachable with convex combinations (statistical mixtures) of  $U$  (reduces to counting/Liouville measure)

Orthogonal states: additive  
different states all else equal



$$\mu(\{A, B\}) = 2^1 = 2$$

$$\mu(\{A, C\}) < 2 = \mu(\{A\}) + \mu(\{C\})$$

Non-orthogonal states: different states  
but in different contexts

sub-additive

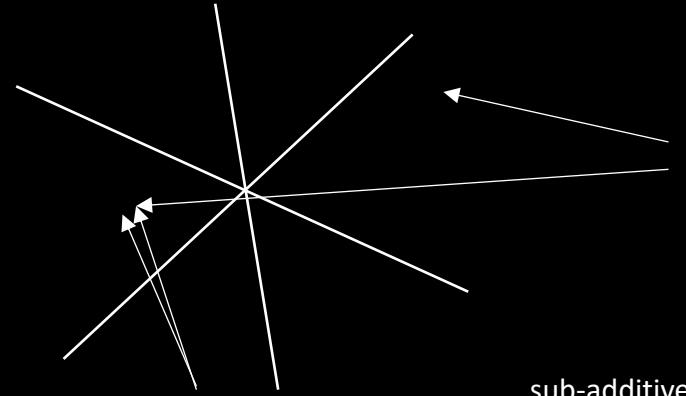
Quantum mechanics  $\Rightarrow$  lower bound  
on #conf (entropy) on continuous DOF

# Conjecture: quantum gravity $\Rightarrow$ lower bound on DOF count

#conf = #DOFs  
Lower bound on this...  
...requires a lower bound on this

From QM: Lower bound on state count requires a severe revisit of particle state space

1. Every point is a single DOF (i.e.  $\mu(\{x\}) = 1$ )
2. Finite volume carries finitely many DOFs (i.e.  $\mu(U) < \infty$ )
3. Count is additive for disjoint regions (i.e.  $\mu(\cup U_i) = \sum \mu(U_i)$ )



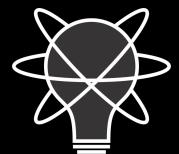
Close points: DOFs not independent

sub-additive

*Same problem!*

Distant points: additive independent DOFs

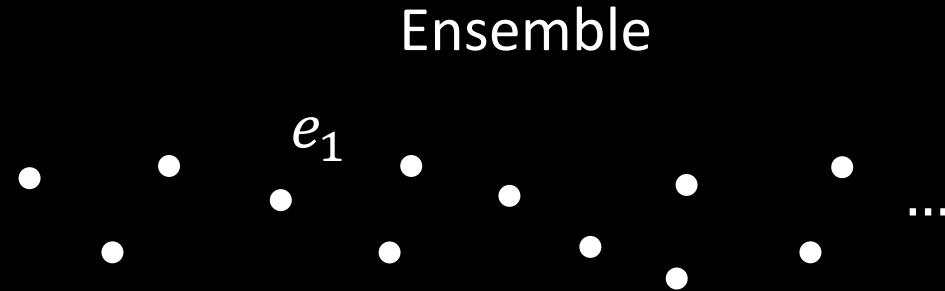
Does lower bound on DOF count require an equally severe revisit of space-time?



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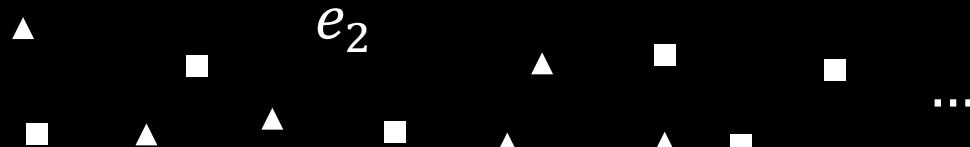
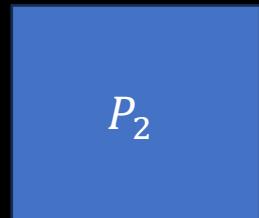
Assumptions  
of  
Physics

## Preparation



Identically prepared ensemble

There is no variability in the ensemble



Pure ensemble

No process can reduce the variability of the ensemble

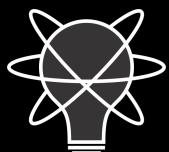
There is no  $a$  such that  $e_2 = pa + \bar{p}b$  for some  $p \in (0,1)$  and  $b \neq a$

In classical theories, all pure ensembles are identically prepared

No real distinction between ensembles and instances

In non-classical theories, no identically prepared ensembles

Zero entropy does NOT correspond to identically prepared ensembles



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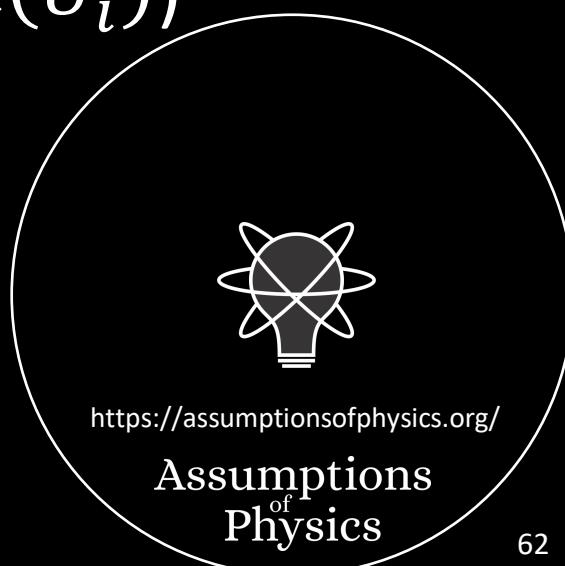
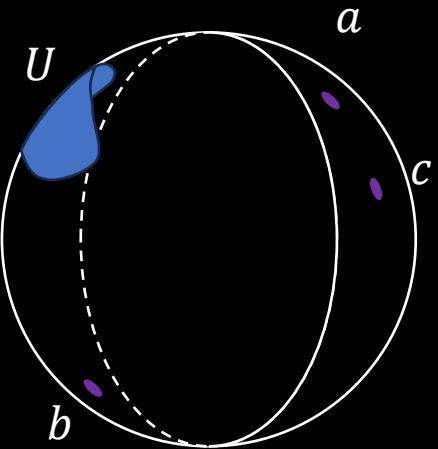
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# The problem with counting on the continuum

$\mu(\text{ a set of states}) = \text{ how many states are there in the set}$

1. Every state counts as one (i.e.  $\mu(\{c\}) = 1$ )
  2. Finite regions have finitely many states (i.e.  $\mu([a, b]) < \infty$ )
  3. Count is additive for disjoint sets (i.e.  $\mu(\bigcup U_i) = \sum \mu(U_i)$ )
- Incompatible! Pick two!**

On a continuum, a finite region has infinitely many states!

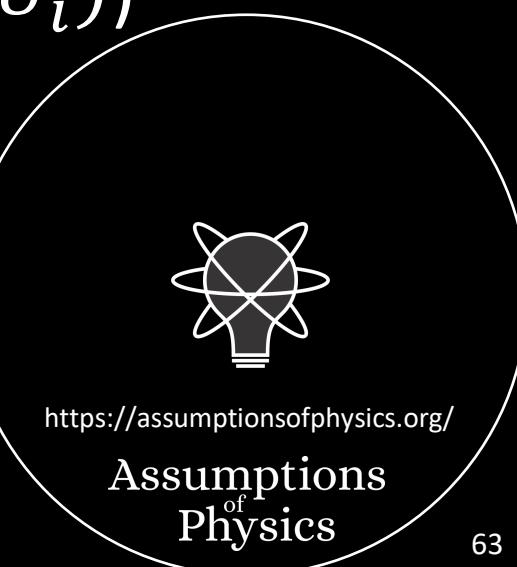
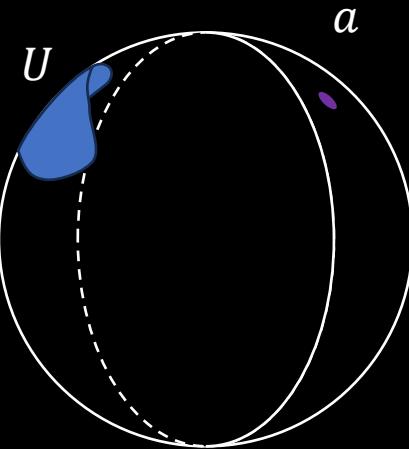


# The problem with counting on the continuum

measure  
→  $\mu(\text{ a set of states}) = \text{ how many states are there in the set}$

1. Every state counts as one (i.e.  $\mu(\{c\}) = 1$ )
  2. Finite regions have finitely many states (i.e.  $\mu([a, b]) < \infty$ )
  3. Count is additive for disjoint sets (i.e.  $\mu(\bigcup U_i) = \sum \mu(U_i)$ )
- Incompatible! Pick two!**

On a continuum, a finite region has infinitely many states!



# How does the quantum measure work?

$$\mu(\{a\}) = \mu(\{b\}) = \mu(\{c\}) = 1$$

$$\mu(\{a, b\}) = 2 = \mu(\{a\}) + \mu(\{b\})$$

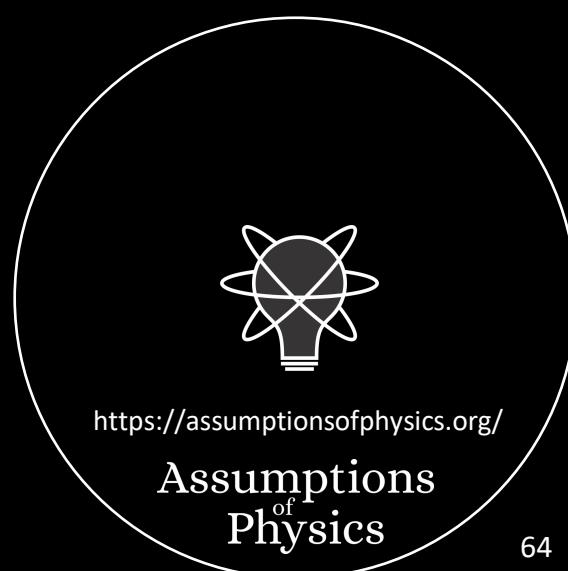
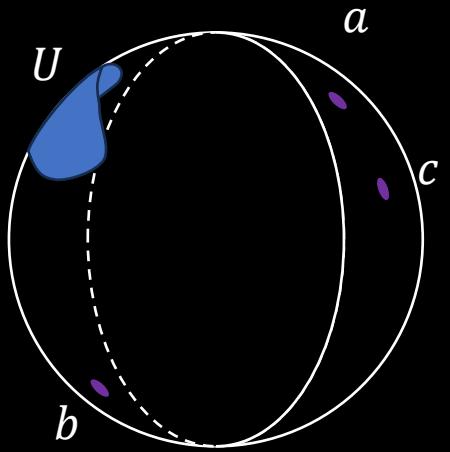
Additive when states are distinguishable (i.e. mutually exclusive)

$$\mu(\{a, c\}) < 2 = \mu(\{a\}) + \mu(\{c\})$$

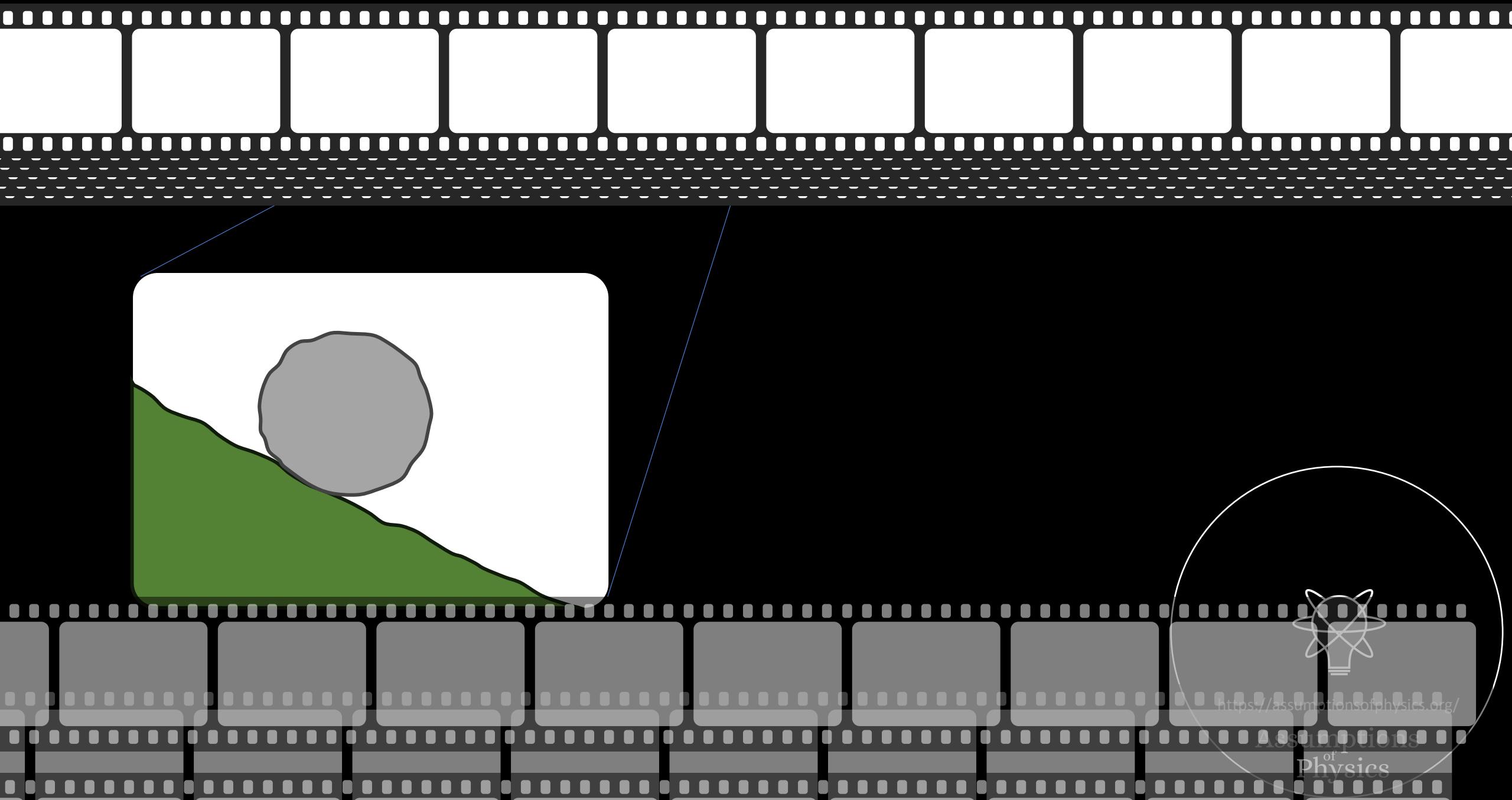
Sub-additive when states are not distinguishable (i.e. mutually exclusive)

Non-additivity coincides  
with non-distinguishability

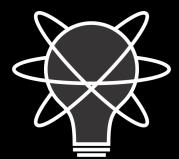
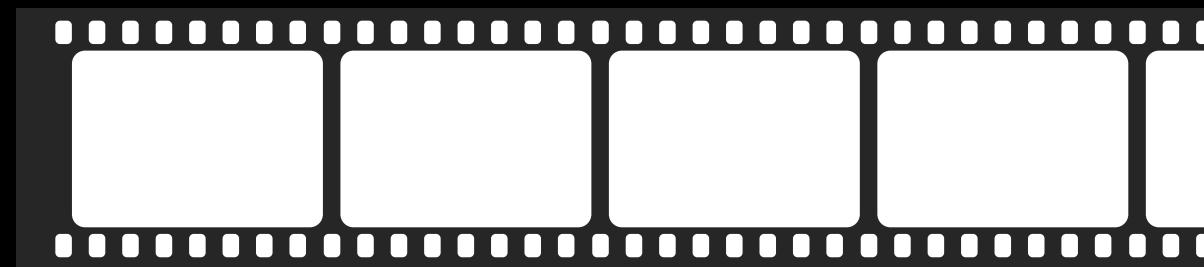
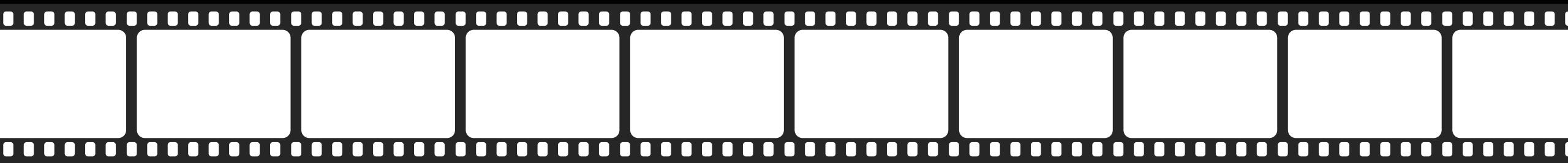
The measure counts  
distinguishable (i.e. mutually  
exclusive) states



$\rightarrow t$

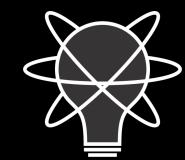


$\rightarrow t$



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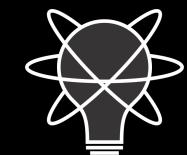
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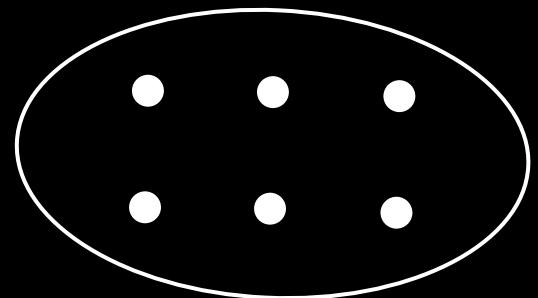
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# How should we count states?



# For discrete classical systems

You simply count the points

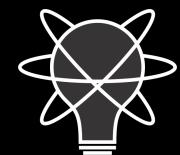


$$\Delta x \Delta p$$

# For continuous classical systems

You use areas of position/momentum

# For quantum systems?



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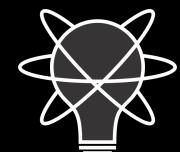
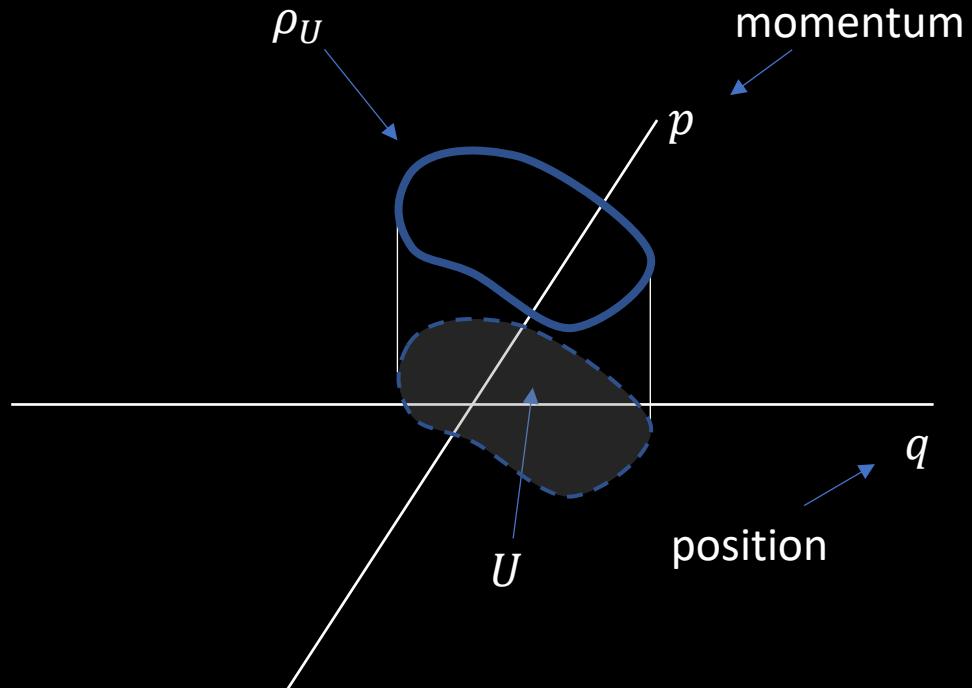
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# For classical systems

$$2^{S(\rho_U)} = \mu(U)$$

Exponential of the entropy  
for a uniform  
distribution over  $U$

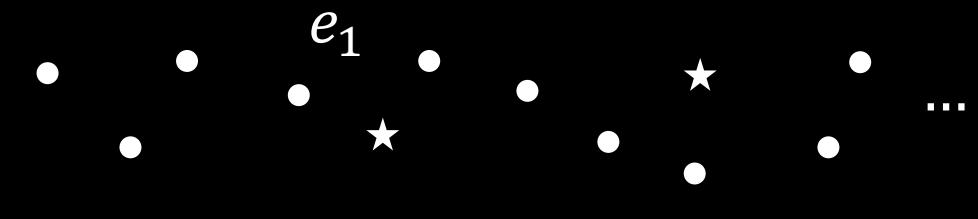
Count of states in  $U$



## Preparation



## Ensemble



## Entropy

$$S(e_1)$$

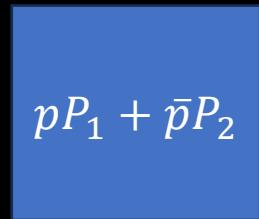
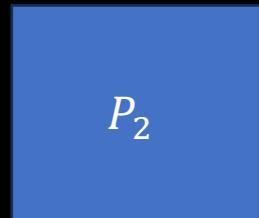
$$S(e_2)$$

Variability within  
an ensemble

One instance is enough  
to tell  $e_1$  and  $e_2$  apart

Mutually exclusive

Orthogonal

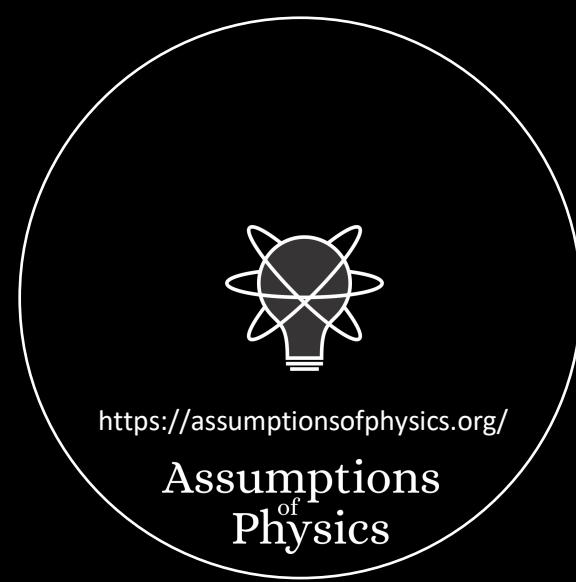


$$pe_1 + (1 - p)e_2$$

Maximal entropy increase  
if orthogonal

$$pS(e_1) + \bar{p}S(e_2)  
+ I(p, \bar{p})$$

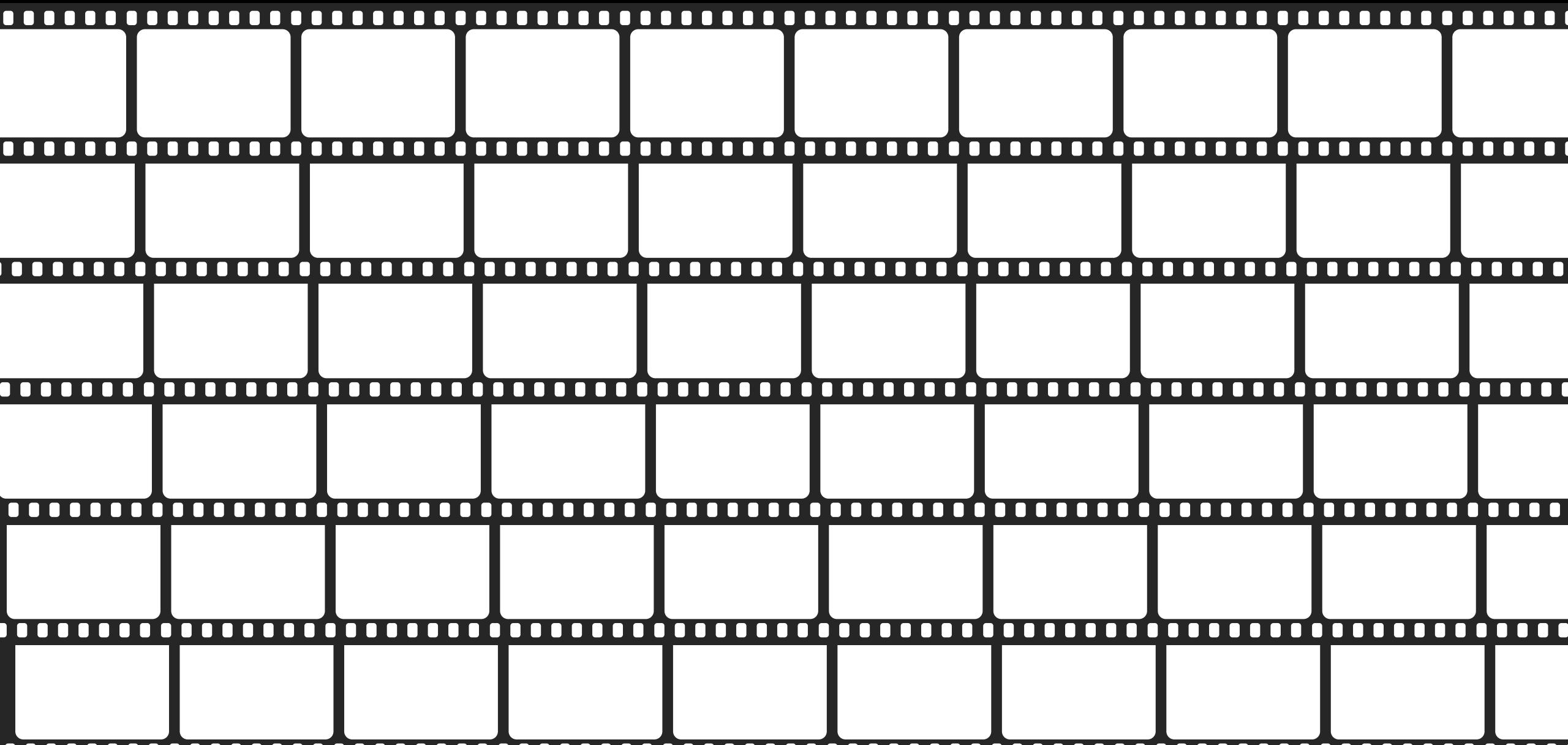
recovers Shannon entropy  
 $I(p, \bar{p}) = -p \log p - \bar{p} \log \bar{p}$

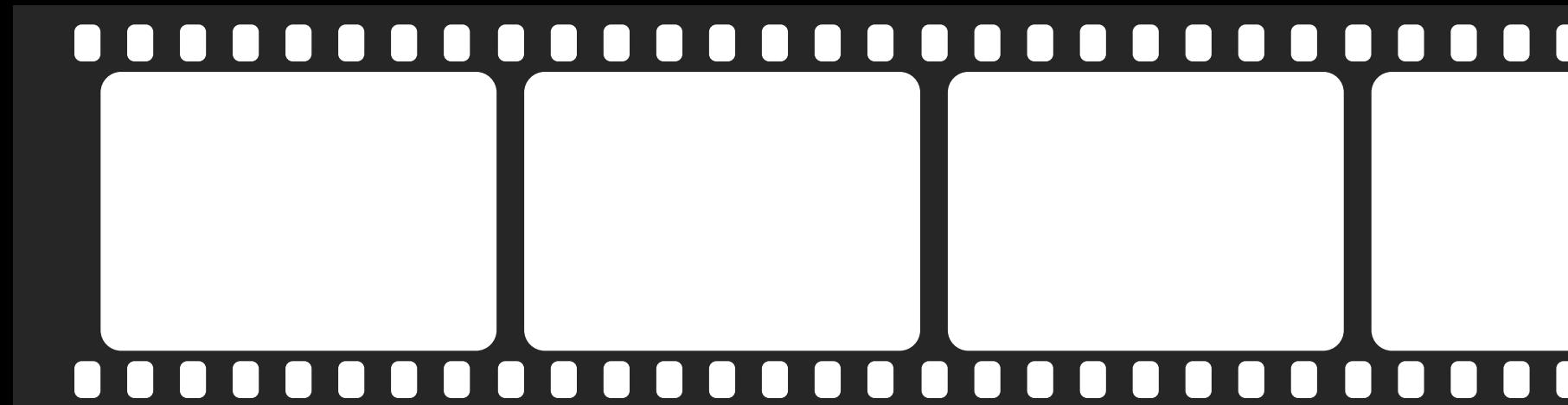
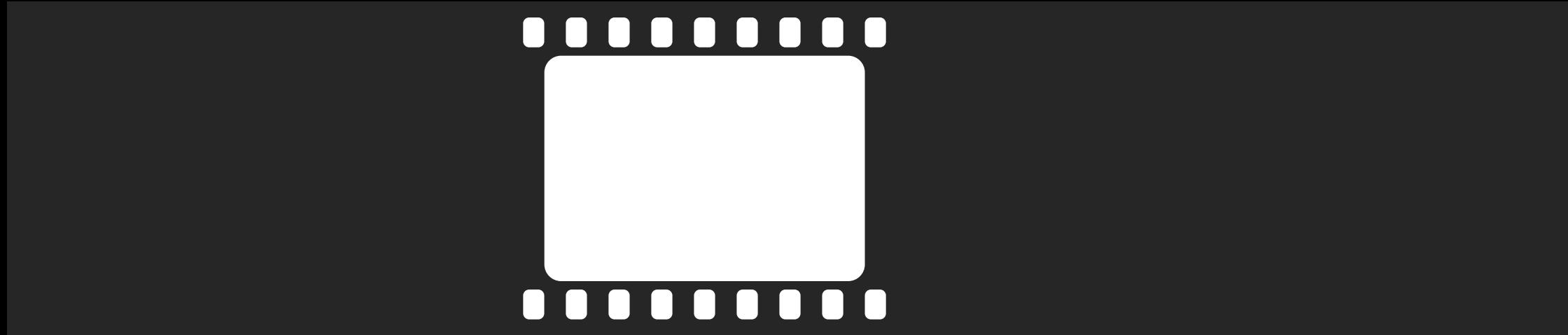


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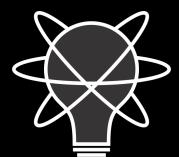
$\rightarrow t$





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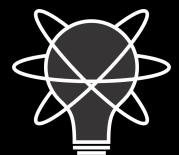
- We don't know what it is
- What is a quantum system? What is a classical system?
- Electron/Proton/... all described by same equation. What makes them the same? Contextual. Inner dynamics is decoupled (it does not matter). Irreducibility?
- Can we reorganize all of physics through simple concepts? Most obvious?
- What



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- 100 years of QM: fake history, we got some math but we don't know what it means
- Uncertainty principle, two modes of evolution, contextuality, Dirac's correspondence principle (quantization)
- Interpretations (hidden variables) and reconstructions don't work
- Why does classical mechanics fail? Classical uncertainty principle.
- Motivations for the third law. Entropy of nothing. Entropy of one state.
- Classical mechanics recovered at high entropy, quantization is putting a lower bound on the entropy
- Entropy is about ensembles; physics is about ensembles; this is the fundamental mischaracterization of physics (film strip)
- Classical vs quantum: reducibility vs irreducibility; wave particule duality (electron a point particle, state probability distr; or electron a spread out blotch of energy? To tell the difference, needs lower entropy description)
- Back to counting states. 3 out of 2
- Space-time quantization



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