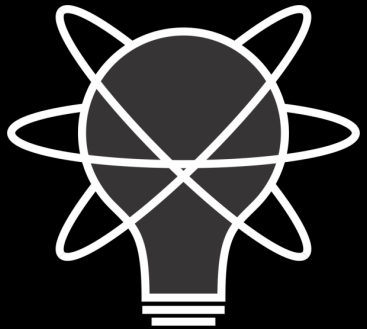


What is classical? What is quantum?

Insights from the Assumptions of Physics program

Gabriele Carcassi
University of Michigan



Assumptions
of
Physics



UNIVERSITY OF
MICHIGAN

International Year of Quantum Science and Technology 2025



International Year of Quantum logo

Date 1 January – 31 December 2025
Type Exhibitions
Website quantum2025.org

wikipedia

nature

Quantum mechanics 100 years on: an unfinished revolution

A century ago, physics had its Darwinian moment – a change in perspective that was as consequential for the physical sciences as the theory of evolution by natural selection was for biology.

ScienceNews

SIGN IN

NEWS

QUANTUM PHYSICS

Quantum mechanics was born 100 years ago. Physicists are celebrating

The International Year of Quantum marks a century of scientific developments



A century of quantum mechanics

On 9 July 1925, in a letter to Wolfgang Pauli, Werner Heisenberg revealed his new ideas, which were to revolutionise physics



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Assumptions
of
Physics

A brief history of quantum mechanics

There were problems
with classical physics

Blackbody
radiation

Photoelectric
effect

Stability of atoms

People futzed around with matrices and waves,
to find something new that has some
correspondence with classical mechanics

⇒ Quantum mechanics!



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Assumptions
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Physics

How to quantize a classical theory

Classical mechanics

$$\frac{df}{dt} = \{f, H\}$$



Quantum mechanics

$$\frac{df}{dt} = \frac{[f, H]}{i\hbar}$$

Quantization!

Dirac's correspondence principle



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Assumptions
of
Physics

$$\{x, p\} = 1 \qquad \frac{[x, p]}{i\hbar} = 1$$

$$\{L_x, L_y\} = L_z \qquad \frac{[L_x, L_y]}{i\hbar} = L_z$$

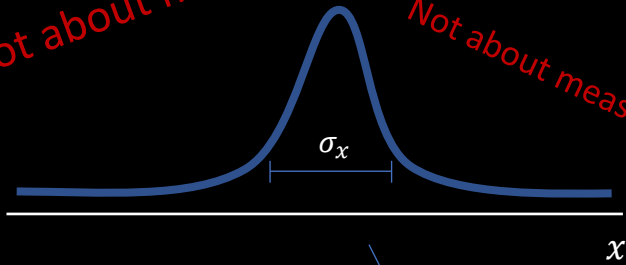
Quantization is a “formal” operation
(i.e. substitute some symbols with some other symbols)



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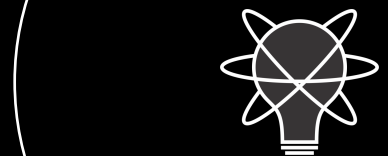
Assumptions
of
Physics

Uncertainty principle



$$\sigma_x \sigma_p \geq \frac{1}{2} |\langle [x, p] \rangle| = \frac{\hbar}{2}$$

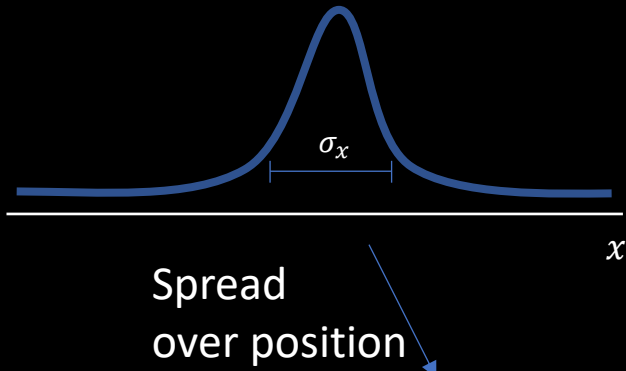
Not about measurement



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Assumptions
of
Physics

Indetermination ~~Uncertainty~~ principle



$$\sigma_x \sigma_p \geq \frac{1}{2} |\langle [x, p] \rangle| = \frac{\hbar}{2}$$

About preparation
not measurement!



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Assumptions
of
Physics

Quantum contextuality

Measurements cannot be thought as revealing pre-existing values

Quantum complementarity

For every quantity we prepare, another quantity is indeterminate



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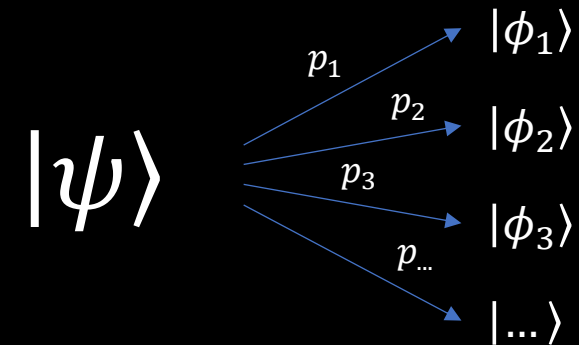
Assumptions
of
Physics

During physical processes
Schrödinger equation

$$i\hbar \frac{d}{dt} |\psi\rangle = H |\psi\rangle$$

smooth, deterministic and reversible

During measurements
Projections



jumpy, non-deterministic, irreversible

Two laws of evolution

Quantum measurement problem



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Assumptions
of
Physics

What is quantum mechanics about?

Hidden variables?

Consistent histories?

Interpretations

Many worlds?

Agent information?

Consciousness?

Quantum reconstructions



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Assumptions
of
Physics

What is the ultimate reason classical mechanics fails, and how is it fixed by quantum mechanics?

Since some of the classical failures are related to thermodynamics,
let's look at the relationship between entropy and uncertainty in
classical physics



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Assumptions
of
Physics

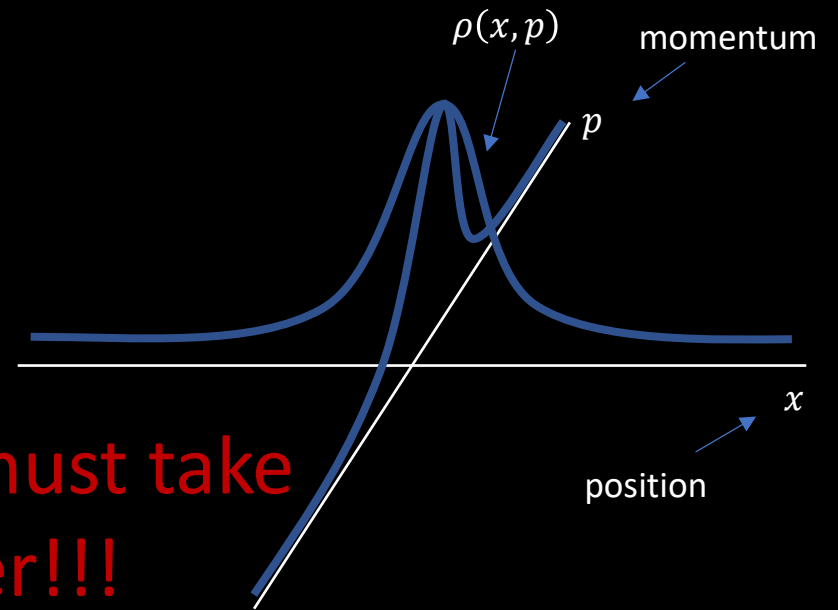
Entropy in classical mechanics

$$S(\rho) = -\int \rho \log \rho \, dx dp$$

Entropy

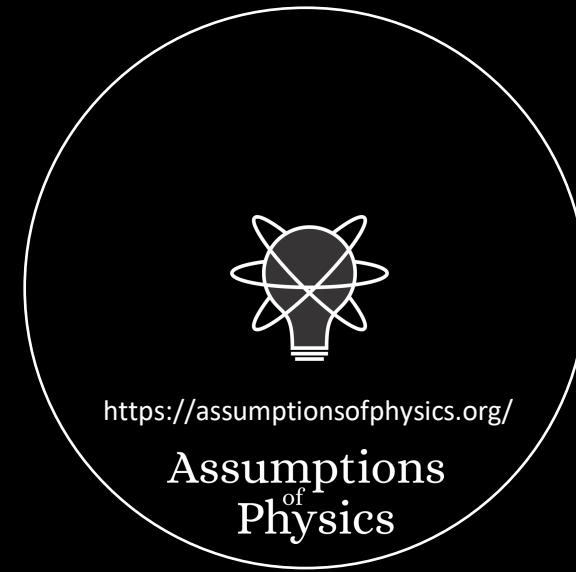
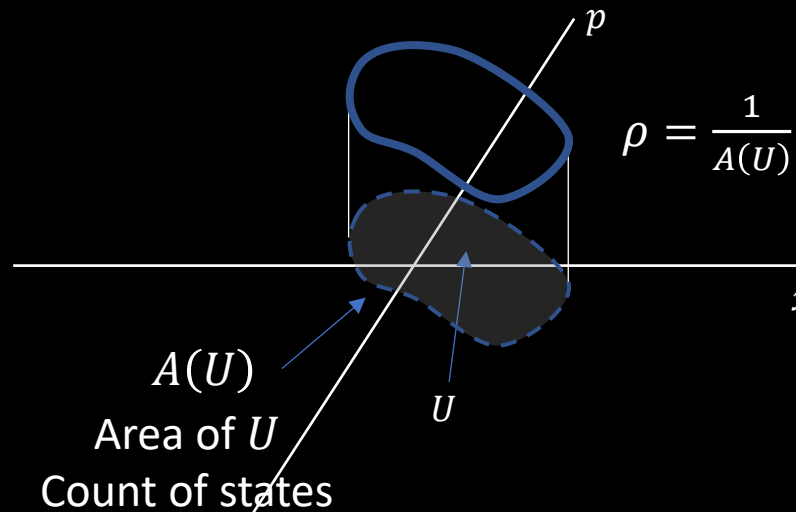
Probability
distribution

Logarithm must take
pure number!!!



$$S(\rho_U) = \log A(U)$$

Uniform
distribution

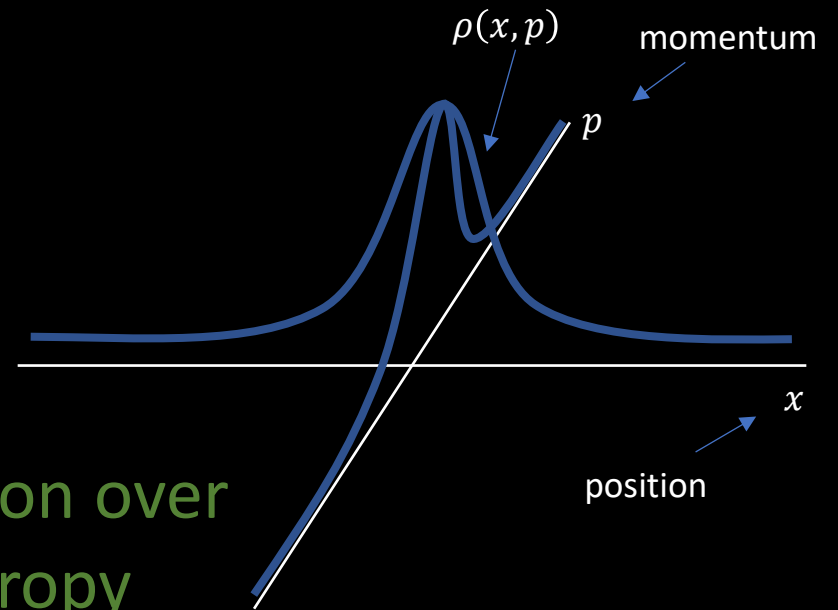


Entropy in classical mechanics

$$S(\rho) = -\int \rho \log h \rho \, dx dp$$

Entropy

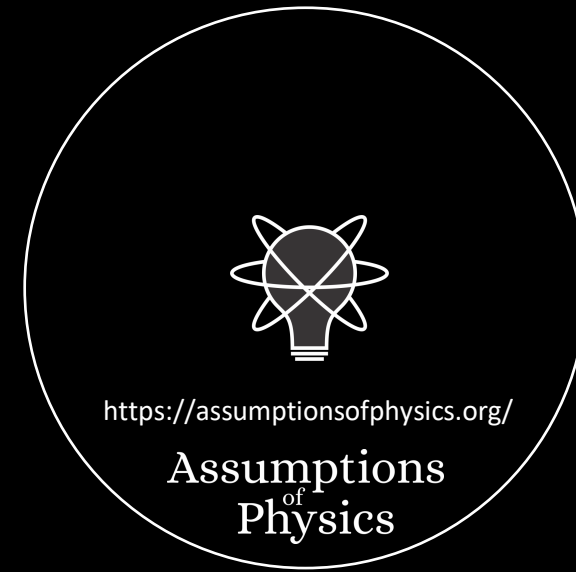
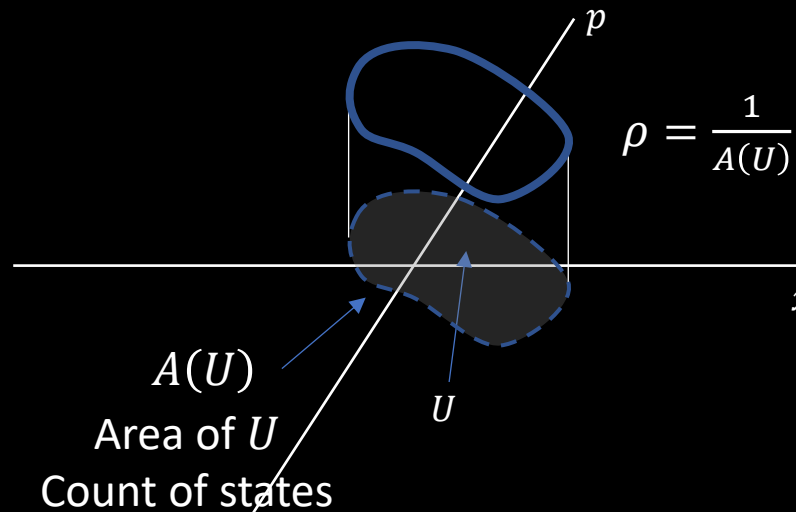
Probability
distribution



Uniform distribution over
area $h \Rightarrow$ zero entropy

$$S(\rho_U) = \log \frac{A(U)}{h}$$

Uniform
distribution



Entropy vs uncertainty

Gaussian maximizes entropy for a given uncertainty

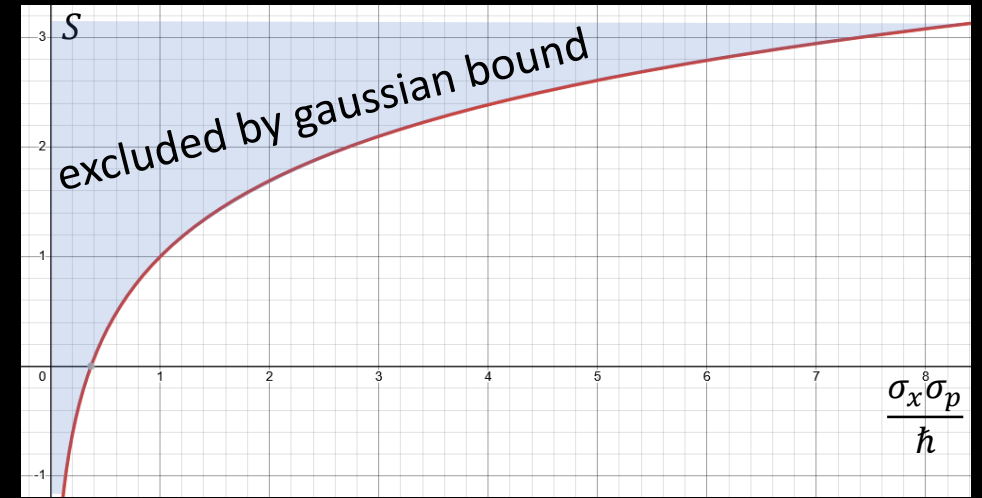
$$S(\rho) \leq \log 2\pi e \frac{\sigma_x \sigma_p}{h}$$

$$\sigma_x \sigma_p \geq \frac{h}{2\pi e} e^{S(\rho)} = \frac{\hbar}{e} e^{S(\rho)}$$

Entropy puts a lower bound
on the uncertainty

Is there a universal lower bound to the entropy?

Let's plot one against the other



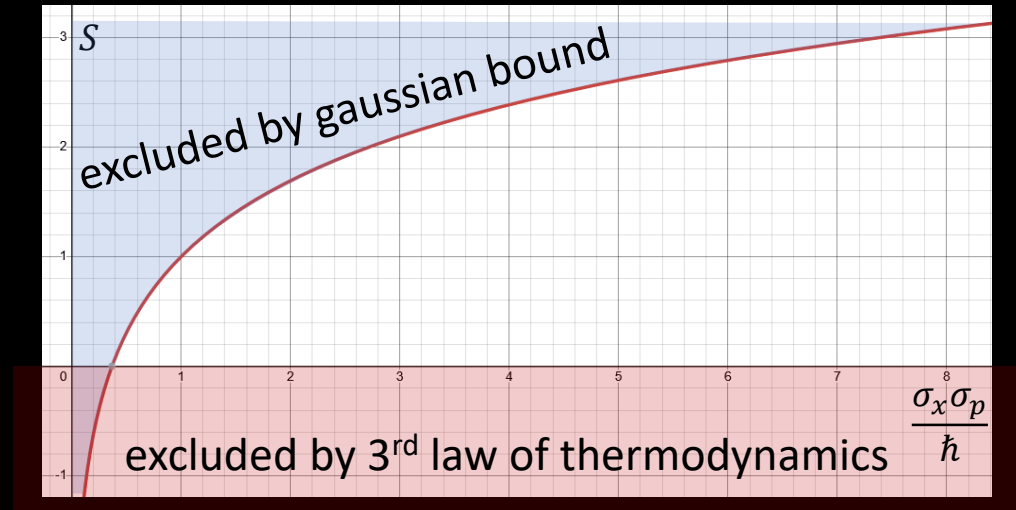
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Assumptions
of
Physics

Enter the 3rd law

Every substance has a finite positive entropy, but at the absolute zero of temperature the entropy may become zero, and does so become in the case of perfect crystalline substances.

G. N. Lewis and M. Randall, Thermodynamics and the free energy of chemical substances (McGraw-Hill, 1923)



$$S \geq 0 \Rightarrow \sigma_x \sigma_p \geq \frac{\hbar}{e}$$

~~Classical uncertainty principle~~
indetermination



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Assumptions
of
Physics

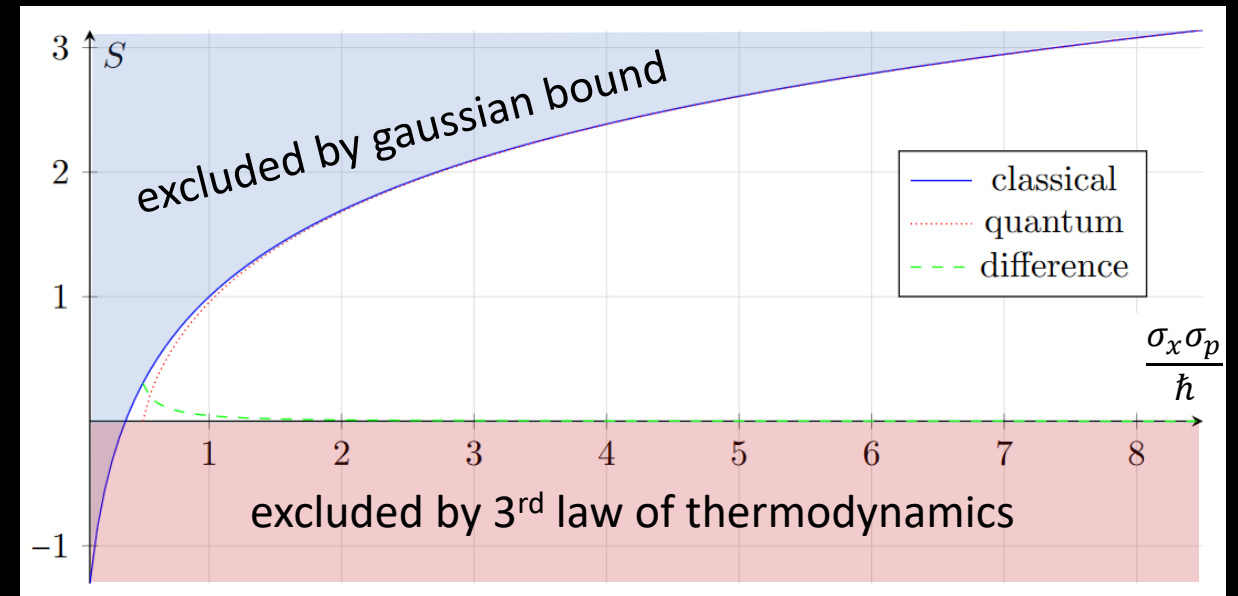
Comparing theories

$$\overset{\text{classical}}{\sigma_x \sigma_p} \geq \frac{\hbar}{e} \quad \overset{\text{quantum}}{\sigma_x \sigma_p} \geq \frac{\hbar}{2}$$

2.71828...

The gaussian bound quickly becomes very similar across theories

Entropy of quantum states is already non-negative



$$S_C = \ln e \sigma$$

$$S_Q = \left(\sigma + \frac{1}{2}\right) \ln \left(\sigma + \frac{1}{2}\right) - \left(\sigma - \frac{1}{2}\right) \ln \left(\sigma - \frac{1}{2}\right)$$

Quantum mechanics incorporates the third law
Classical mechanics does not

Conjecture: does quantum mechanics recover
classical mechanics at high entropy?



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Assumptions
of
Physics



greetings

ML

Manuele Landini <manulando@gmail.com>
To carcassi@umich.edu

↩ Reply

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⋮

Wed 7/10/2024 6:15 AM

📄

You replied to this message on 7/10/2024 10:00 AM.

Caro Gabriele,

Mi chiamo Manuele Landini e lavoro a Innsbruck (Austria) come senior scientist in un gruppo di fisica atomica sperimentale. Puoi vedere di cosa ci occupiamo sul nostro sito: <https://quantummatter.at>.

Ho visto un po' dei tuoi video su youtube. Mi sembra un progetto molto ambizioso, ma promettente. Mi farebbe piacere riuscire a spiegare agli studenti in futuro in termini piu' fisici concetti come le sovrapposizioni o il teorema spin-statistica.

Per la storia della metrica, da quel che ho capito hai bisogno di una metrica che non sia basata sull'entropia, visto che vuoi definire una distanza a entropia costante. Ci sono varie opzioni, ma la trace distance [Trace distance - Wikipedia](#) funziona perche' ha una proprieta' fondamentale che puoi usare. Chiamala: $T(\rho, \sigma)$

Se parti da stati puri, si riduce a $(1 - \langle \psi | \phi \rangle)^{1/2}$. Quindi per massimizzarla, scegli due stati ortogonali (non importa quali). Il massimo e' $T_0 = 1$. Una volta che hai questi stati, che hanno entropia 0, li puoi trasformare in stati con entropia finita (in particolare quelli con massima distanza) tramite una trace preserving map M .

Siccome T si contrae, hai che $T(M(\rho), M(\sigma)) \leq T(\rho, \sigma)$. L'uguale vale se la mappa e' unitaria. Così definisci un serie di step in cui la distanza massima decresce $T_{n+1} < T_n$, fino ad arrivare a 0 per stati fully mixed.

arXiv

> quant-ph > arXiv:2411.00972

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All fields

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Quantum Physics

[Submitted on 1 Nov 2024 (v1), last revised 3 Dec 2024 (this version, v2)]

Classical mechanics as the high-entropy limit of quantum mechanics

Gabriele Carcassi, Manuele Landini, Christine A. Aidala

We show that classical mechanics can be recovered as the high-entropy limit of quantum mechanics. That is, the high entropy masks quantum effects, and mixed states of high enough entropy can be approximated with classical distributions. The mathematical limit $\hbar \rightarrow 0$ can be reinterpreted as setting the zero entropy of pure states to $-\infty$, in the same way that non-relativistic mechanics can be recovered mathematically with $c \rightarrow \infty$. Physically, these limits are more appropriately defined as $S \gg 0$ and $v \ll c$. Both limits can then be understood as approximations independently of what circumstances allow those approximations to be valid. Consequently, the limit presented is independent of possible underlying mechanisms and of what interpretation is chosen for both quantum states and entropy.

Comments: 14 pages, 3 figures

Subjects: Quantum Physics (quant-ph)

Cite as: arXiv:2411.00972 [quant-ph]
(or arXiv:2411.00972v2 [quant-ph] for this version)
<https://doi.org/10.48550/arXiv.2411.00972>

Submission history

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[v2] Tue, 3 Dec 2024 13:52:45 UTC (20 KB)

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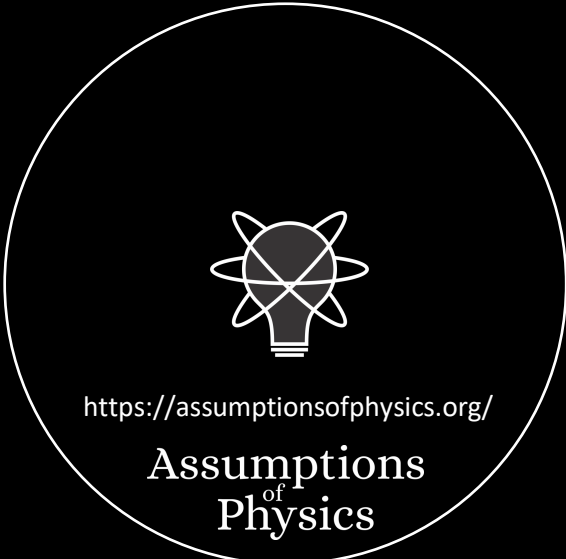
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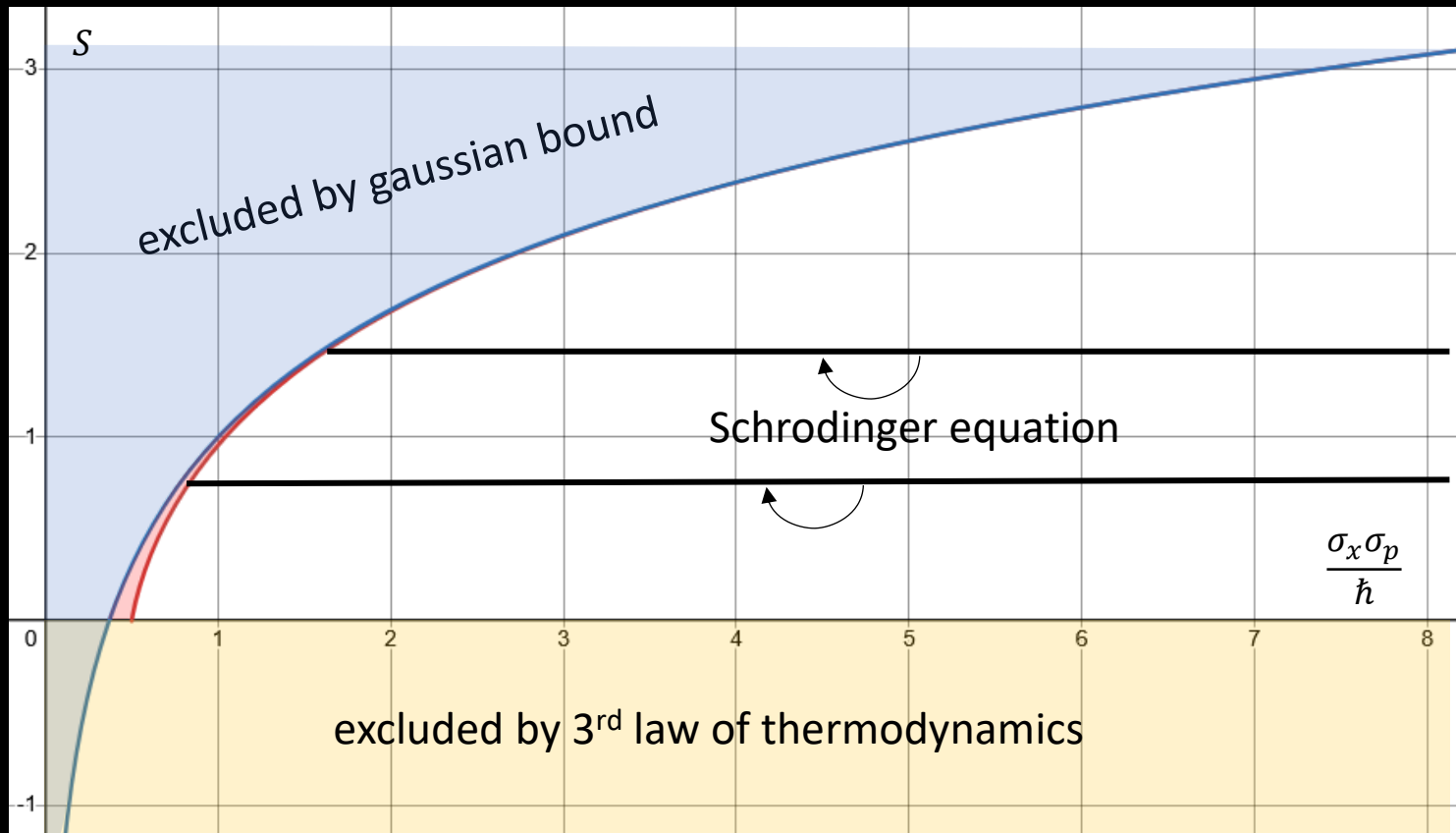
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Bookmark





— classical
— quantum

$$\frac{[x, p]}{i\hbar} = 1 \Rightarrow \frac{[T(x), T(p)]}{i\hbar} = \lambda$$

$$\sigma_{T(x)} \sigma_{T(p)} \geq \frac{1}{2} |\langle [T(x), T(p)] \rangle| = \frac{\hbar}{2} \lambda$$

1. Increase the entropy of all states

T - Entropy increasing map

Spread increase in position and momentum



<https://assumptionsofphysics.org/>

Assumptions
of
Physics

Mathematical limit

$$\frac{[x,p]}{i\hbar} \xrightarrow{\hbar \rightarrow 0} \{x,p\}$$

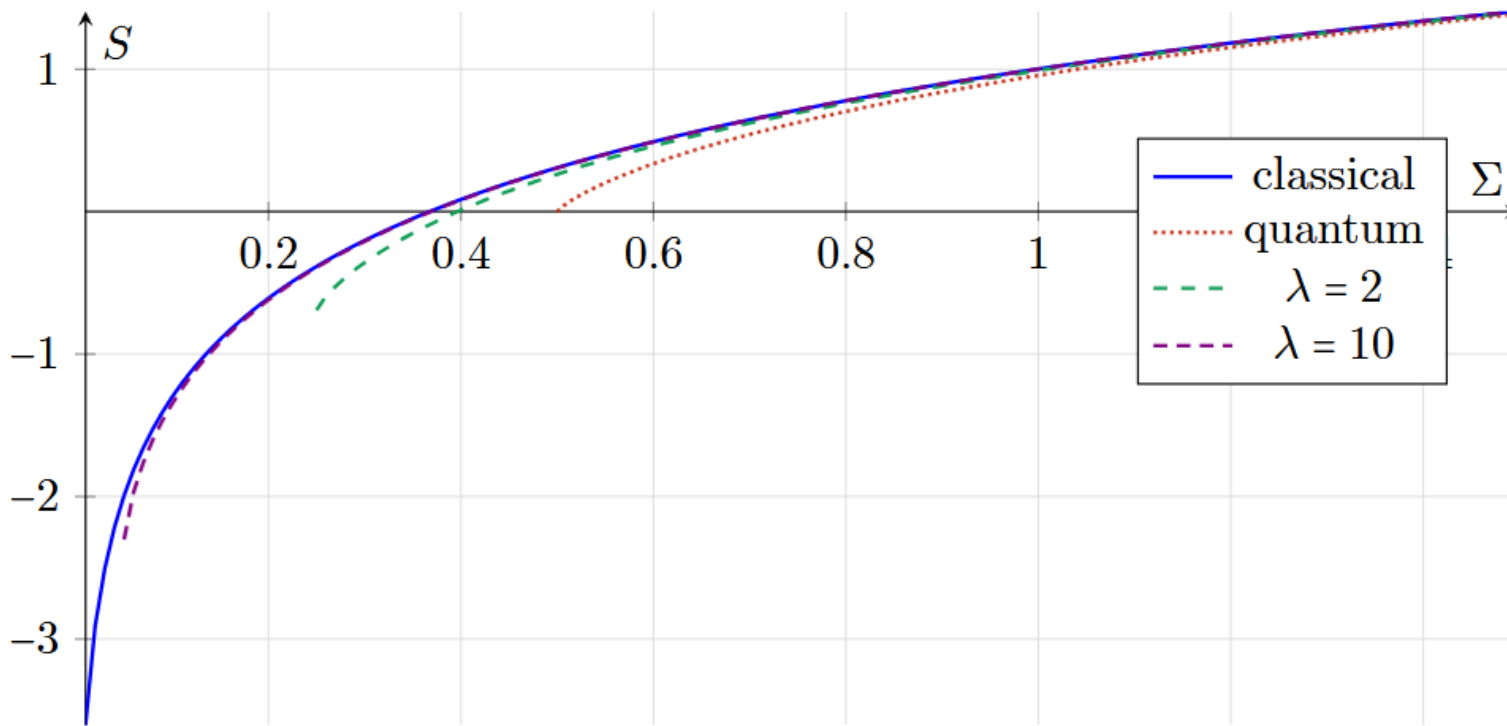
Makes no physical sense!

\hbar is a constant: it doesn't change



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Assumptions
of
Physics



2. Decrease the entropy of pure states

$$\frac{[X, P]}{i\hbar} = \frac{1}{\lambda} \quad \frac{[X, P]}{i(\hbar/\lambda)} = 1 \xrightarrow[\hbar/\lambda \rightarrow 0]{\lambda \rightarrow \infty} \{x, p\} = 1$$

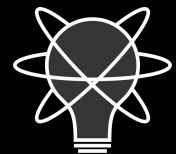


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Assumptions
of
Physics

Low speed " $c \rightarrow \infty$ "		Speed	
Entropy High entropy " $\hbar \rightarrow 0$ "	Classical Mechanics	Relativistic Mechanics	
	Quantum Mechanics	Quantum Field Theory	

No-mechanism limit
(same as non-relativistic limit)



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Assumptions
of
Physics

Quantum to classical

$$\frac{[A, B]}{i\hbar} \longrightarrow \{A, B\}$$

High entropy limit

Classical to quantum

$$\{A, B\} \longrightarrow \frac{[A, B]}{i\hbar}$$

???

Up to isomorphisms, the Moyal bracket is the unique one-parameter Lie-algebraic deformation of the Poisson bracket

If I asked for a theory that puts a hard lower bound on the entropy,
and recovers classical mechanics at high entropy....

there is only one way to do it



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Assumptions
of
Physics

Quantum to classical

$$\frac{[A, B]}{i\hbar} \longrightarrow \{A, B\}$$

High entropy limit

Classical to quantum

$$\{A, B\} \longrightarrow \frac{[A, B]}{i\hbar}$$

Entropic lower bound
that recovers classical mechanics

Quantizing a classical theory
means putting a lower bound
on the entropy



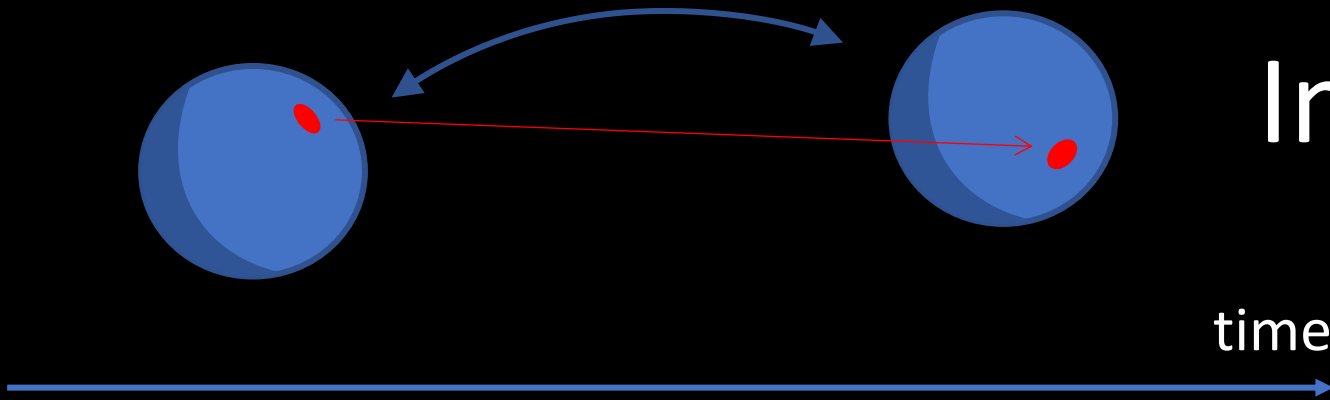
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Assumptions
of
Physics

Classical system

Infinitesimally reducible

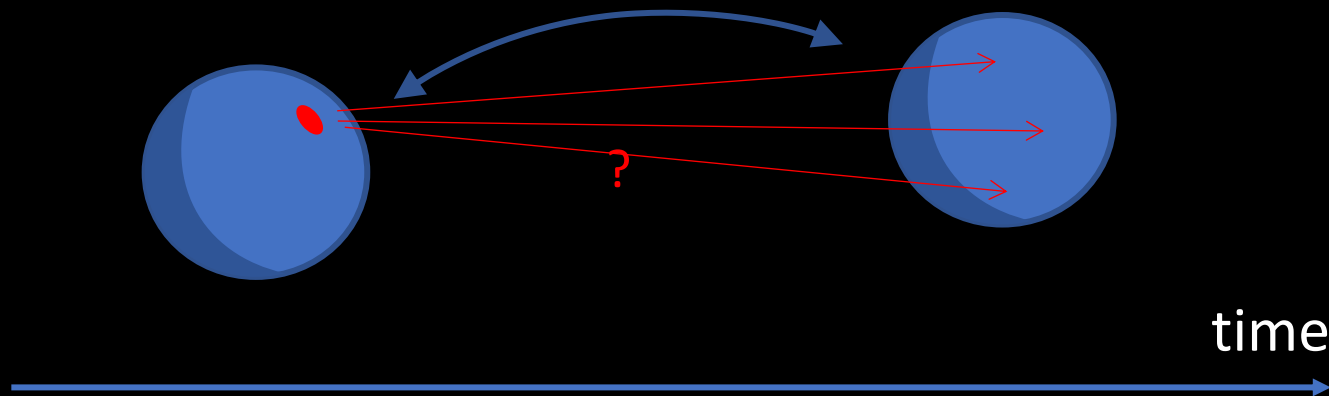
Internal dynamics
perfectly accessible



Quantum system

Irreducible

Internal dynamics
inaccessible

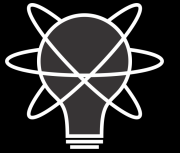


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Assumptions
of
Physics

What is entropy?

Why a lower bound?



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Assumptions
of
Physics

What is entropy?

You can think of it as disorder

You can think of it as information

You can think of it as
lack of information

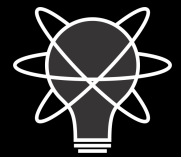


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Assumptions
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Physics

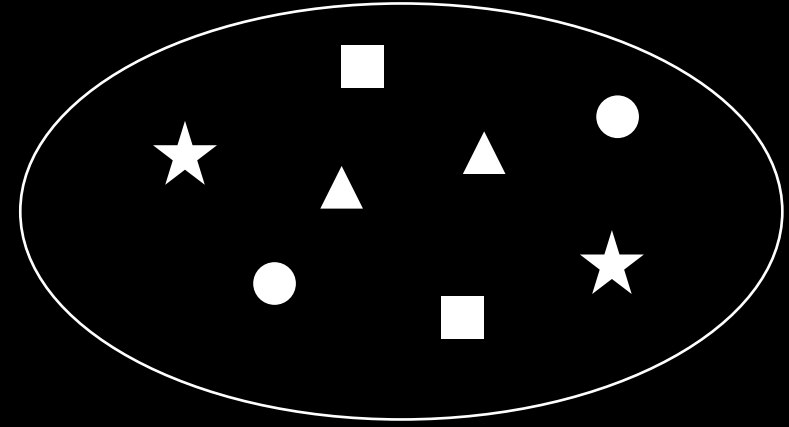
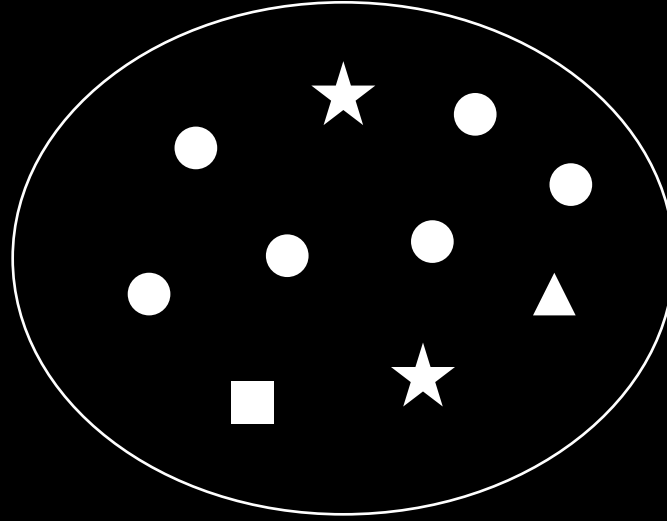
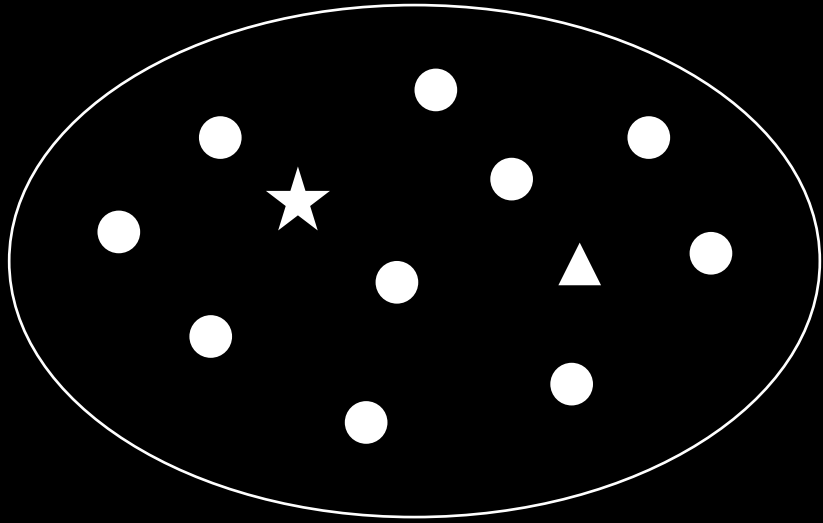


Is or is not.
There is no “you can think of it as.”



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Assumptions
of
Physics



Suppose you have a (possibly infinite) collection of things

What is the variability of the collection?

How different are the elements within the collection?



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Assumptions
of
Physics

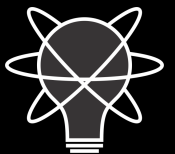
Given a discrete distribution p_i
we want an indicator $I(p_i)$ of variability such that:

1. Continuous in p_i
2. Increases if the number of cases increases
3. Linear in probability $I(p_1\hat{p}_i, p_2\bar{p}_j) = I(p_1, p_2) + p_1I(\hat{p}_i) + p_2I(\bar{p}_j)$

$$\Rightarrow \text{Then } I(p_i) = -\sum_i p_i \log p_i$$

Proof can be generalized to classical continuum and quantum mechanics

Entropy is variability!



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Assumptions
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Physics

Variability as a better characterization of Shannon entropy

Gabriele Carcassi, Christine A Aidala and Julian Barbour

Published 19 May 2021 • © 2021 European Physical Society

[European Journal of Physics](#), [Volume 42](#), [Number 4](#)

Machine learning to detect schedules using spatiotemporal data of behavior: A proof of concept

Marc J. Lanovaz, Varsovia Hernandez, Alejandro León

2025, Journal of the Experimental Analysis of Behavior - Article

Why so slow? Models of parkinsonian bradykinesia

David Williams

2024, Nature Reviews Neuroscience - Article

Using Entropy Metrics to Analyze Information Processing Within Production Systems: The Role of Organizational Constraints

Frits van Merode, Henri Boersma, Fleur Tournois, Windi Winasti, Nelson Aloysio Reis de Almeida Passos, Annelies van der Ham

2025, Logistics - Article

Internal Versus Forced Variability Metrics for General Circulation Models Using Information Theory

Aakash Sane, Baylor Fox-Kemper, David S. Ullman

2024, Journal of Geophysical Research - Oceans - Article

■ ■ ■

Parameterizing Vertical Mixing Coefficients in the Ocean Surface Boundary Layer Using Neural Networks

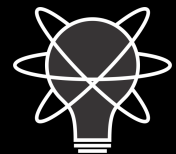
Aakash Sane, Brandon G. Reichl, Alistair Adcroft, Laure Zanna

2023, Journal of Advances in Modeling Earth Systems - Article

Spatiotemporal characteristics of future precipitation variability in the Tianshan Mountain region of China

Xianglin Lyu, Junkai Du, Yaqin Qiu, Yangwen Jia, Chunfeng Hao, Hao Dong

2025, Journal of Hydrology Regional Studies - Article



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Assumptions
of
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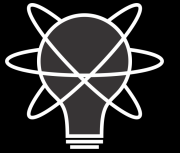
Entropy is not the property
of an element, but of a distribution

A distribution of what?

A statistical distribution of
the molecules of a gas?

A credence distribution of what
an agent may think is there?

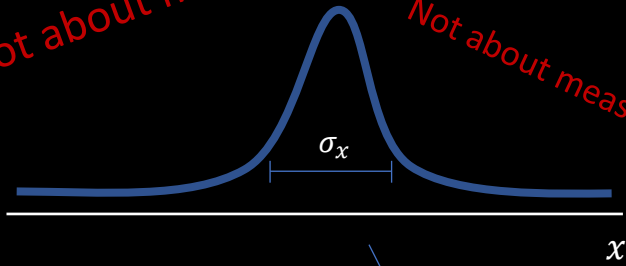
A probability distribution of the
outcomes of a measurement?



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Assumptions
of
Physics

Uncertainty principle



$$\sigma_x \sigma_p \geq \frac{1}{2} |\langle [x, p] \rangle| = \frac{\hbar}{2}$$

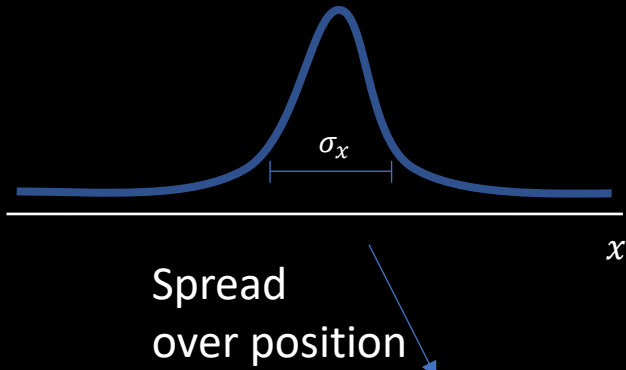
Not about measurement



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Assumptions
of
Physics

Indetermination ~~Uncertainty~~ principle



$$\sigma_x \sigma_p \geq \frac{1}{2} |\langle [x, p] \rangle| = \frac{\hbar}{2}$$

About preparation
not measurement!

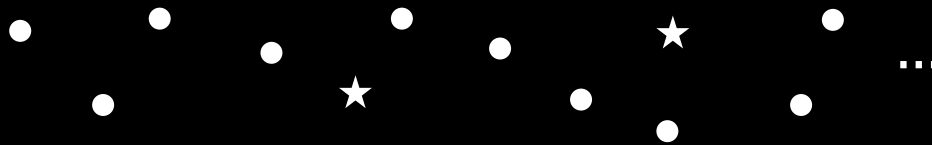


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Assumptions
of
Physics

Entropy is the variability of an ensemble:
the collection of all outputs
of a reliable preparation procedure

Preparation
procedure



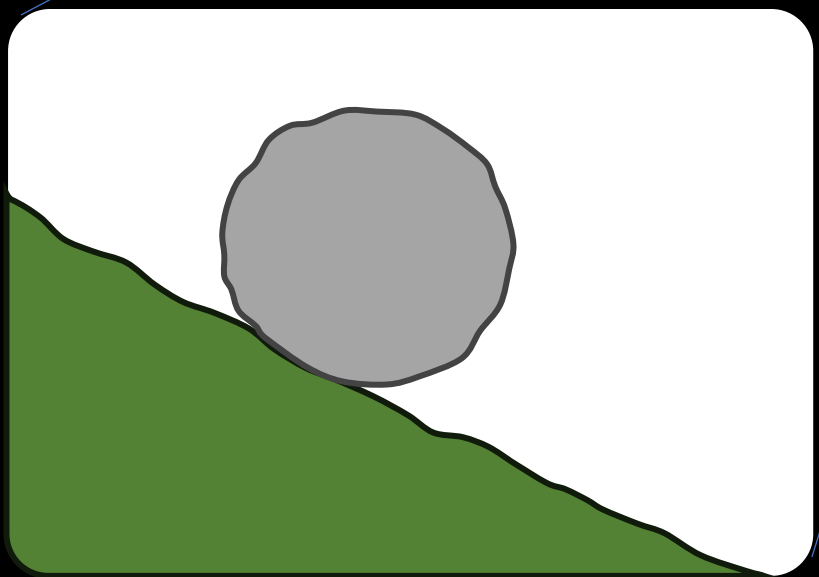
Ensemble of outputs



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Assumptions
of
Physics

→ t



Physics describes the evolution
of a particular system like a movie

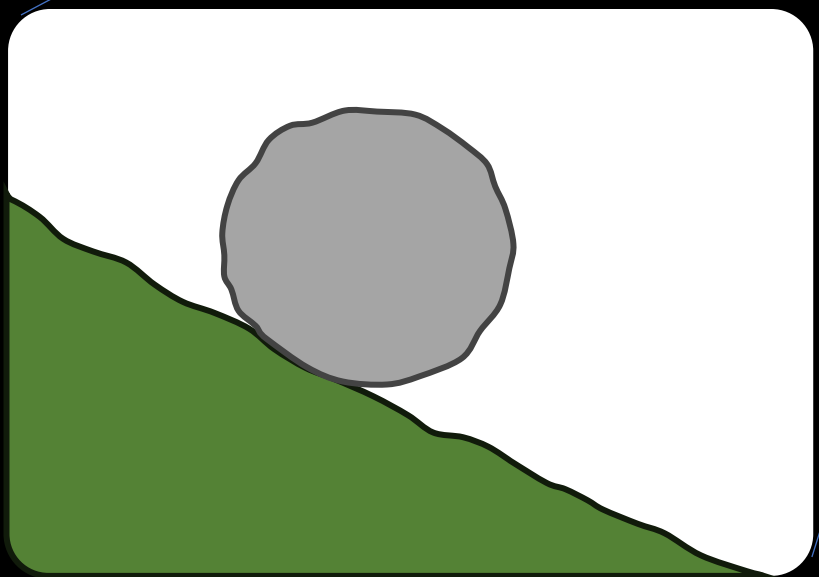
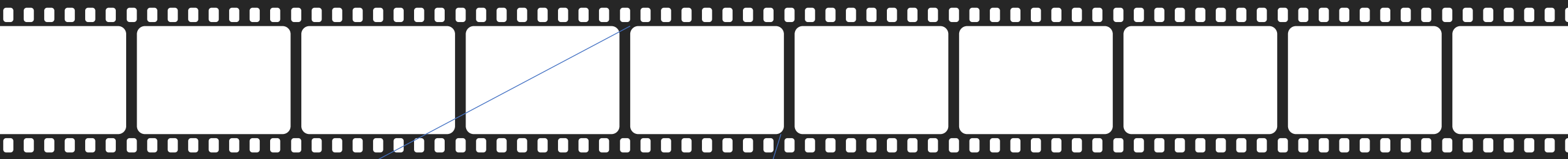
$$x(t)$$



<https://assumptionsofphysics.org/>

Assumptions
of
Physics

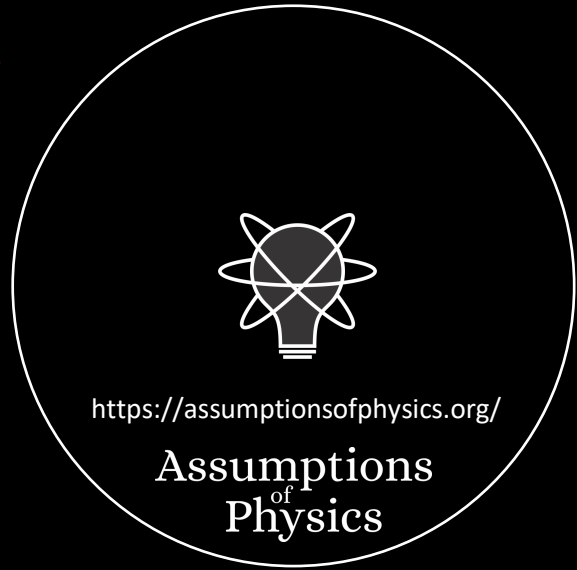
→ t



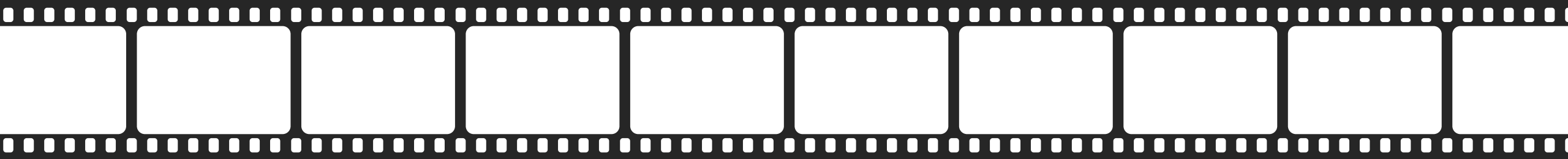
~~Physics describes the evolution
of a particular system like a movie~~

-100 points

$$x(t)$$



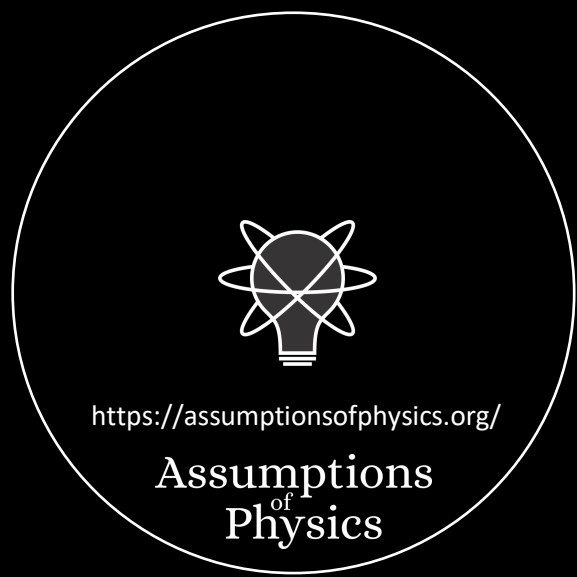
→ t



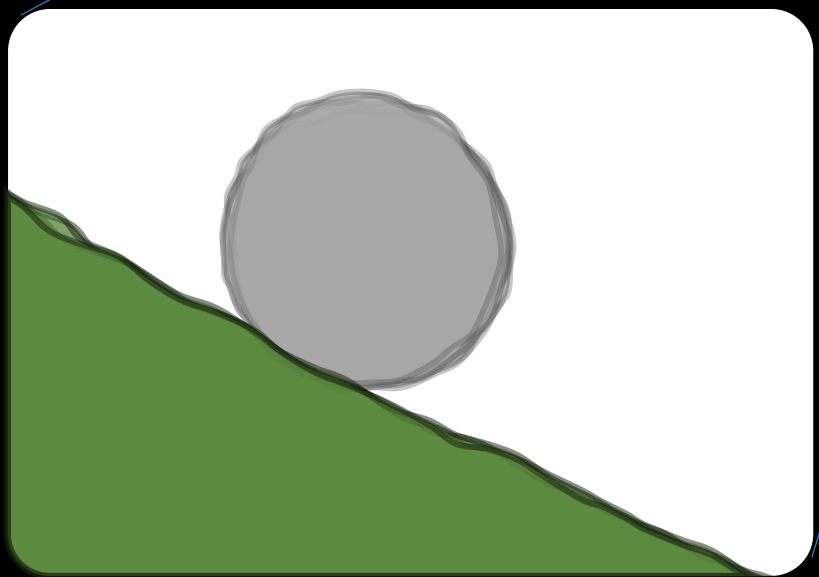
Physics describes the evolution
of similarly prepared systems

Every time I do this...

$$x(t)$$

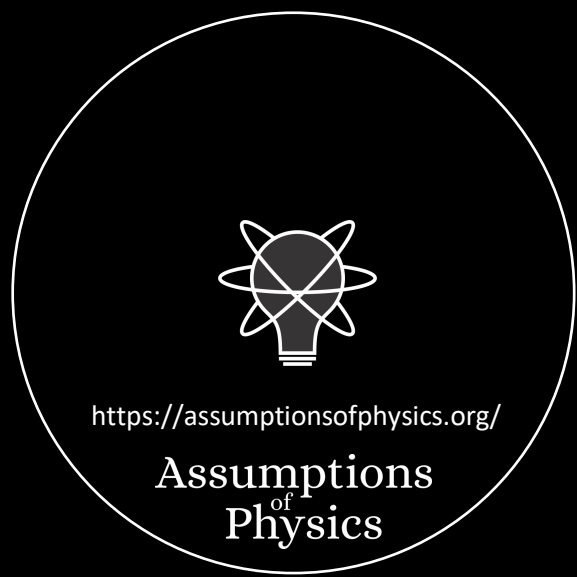


→ t



Physics describes the evolution
of similarly prepared systems

$$x(t)$$



Statistical description of preparations
and measurement outcomes

$x(t)$

In physics, we can only
study relationships that
can be reproduced:
relationships between
ensembles

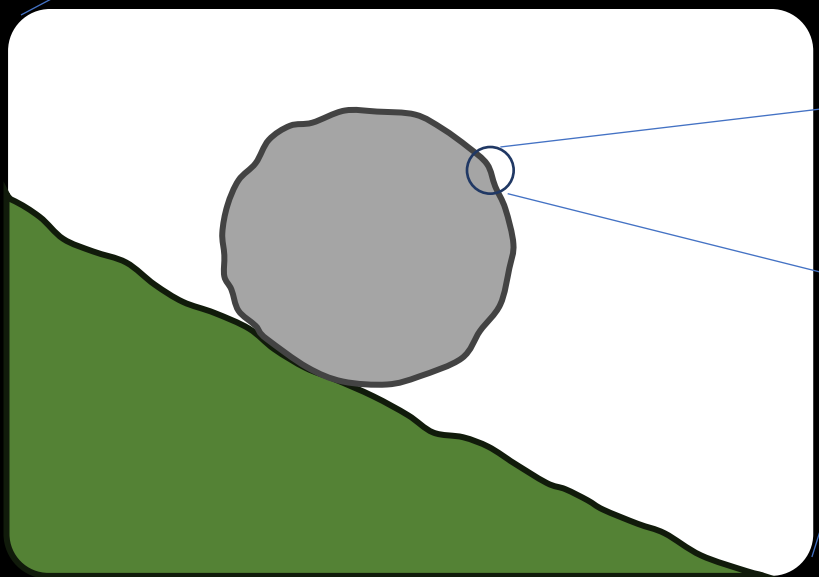
Statistical description
of all clocks synchronized
at a particular initial event



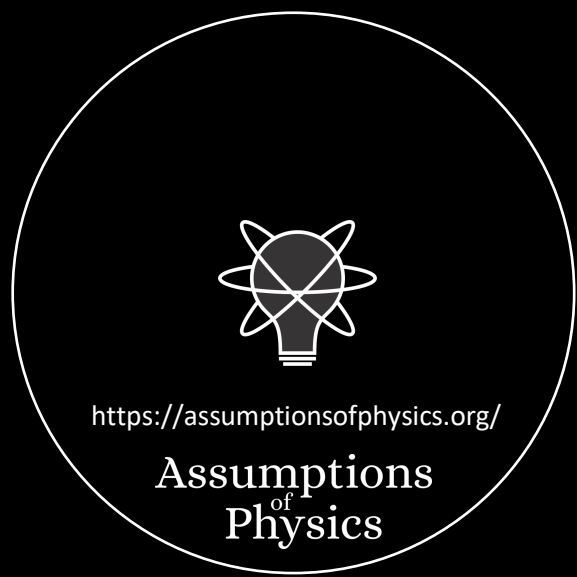
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Assumptions
of
Physics

→ t



Even a single measurement is an average over a finite time



To be able to define a probability distribution over heads or tails, you need first to define what heads and tails are

We can do that because we can leave the coin be, and it remains heads or tails

If it kept changing more than 100 times a second, we wouldn't perceive heads or tails: just a blur



wikipedia

wikipedia

The state of a system is what remains
“stable” for at least a short amount of time

It's ensembles all the way down



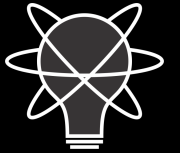
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Assumptions
of
Physics

What is entropy?

Variability within a preparation!

Why a lower bound?

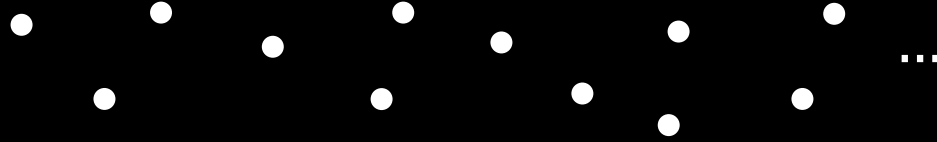


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Assumptions
of
Physics

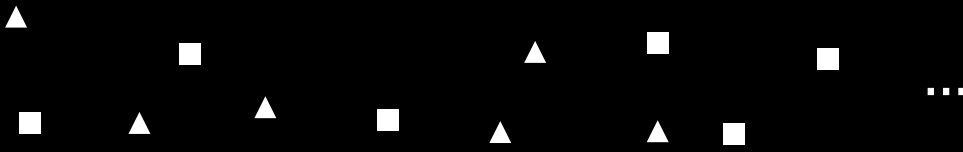
Classical mechanics

Preparation
procedure



Assumes we can have
perfect preparations

Preparation
procedure



Concedes that there are
no perfect preparations

Quantum mechanics



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Assumptions
of
Physics

For classical systems

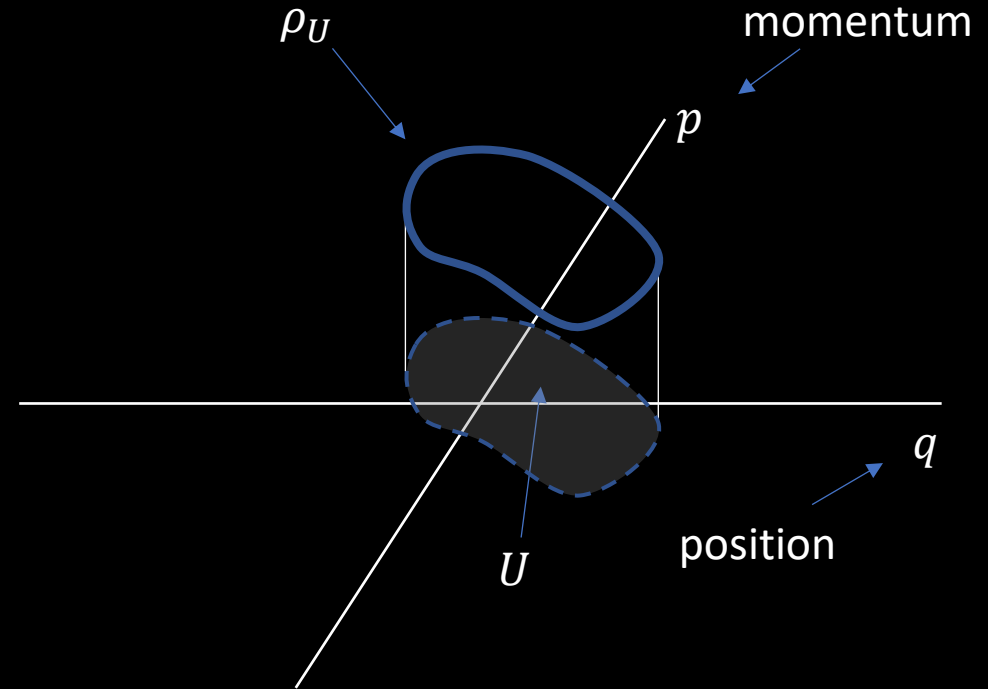
$$S(\rho_U) = \log \frac{A(U)}{h}$$

Entropy for a uniform
distribution over U

Logarithm of the count
of states in U

$$S(\rho_U) < 0 \quad \frac{A(\rho_U)}{h} < 1$$

Less than one state!



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Assumptions
of
Physics

What is entropy?

Variability within a preparation!

Why a lower bound?

Our preparations cannot be perfect

Regions with “less than one state”
make no sense

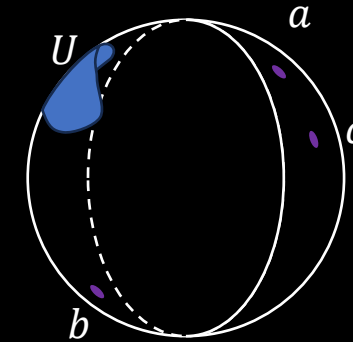


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Assumptions
of
Physics

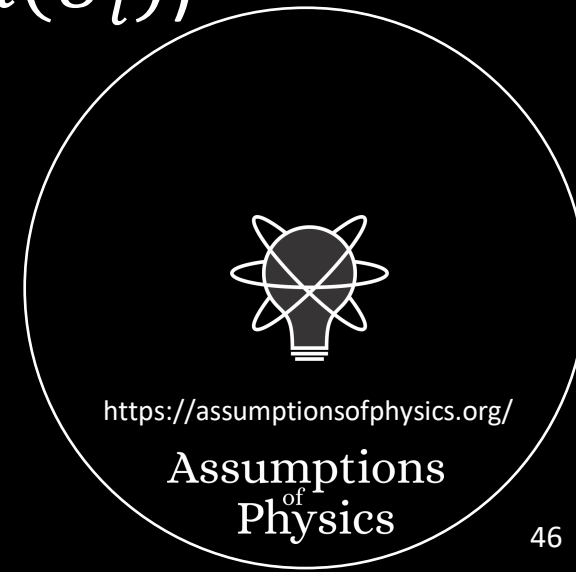
The problem with counting (on the continuum)

measure μ (a set of states) = how many states there are in the set



1. Every state counts as one (i.e. $\mu(\{c\}) = 1$)
2. Finite regions have infinitely many states (i.e. $\mu(U) < \infty$)
3. Count is additive for disjoint sets (i.e. $\mu(\cup U_i) = \sum \mu(U_i)$)

On a continuum, a finite region
has infinitely many states!

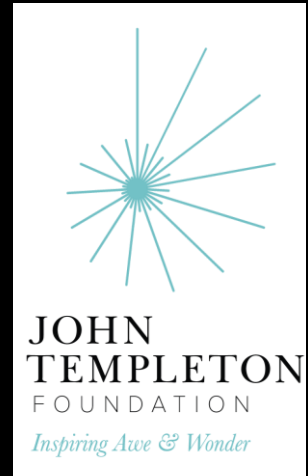
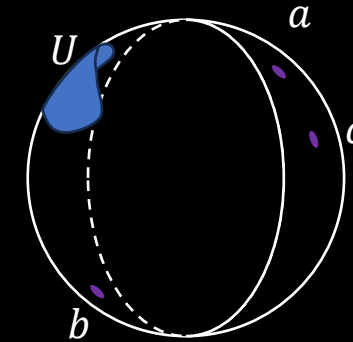


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Assumptions
of
Physics

The problem with counting on the continuum

measure μ (a set of states) = how many states are there in the set



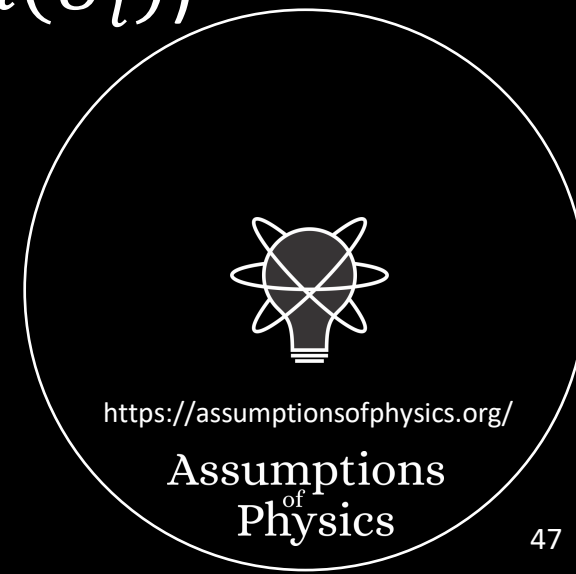
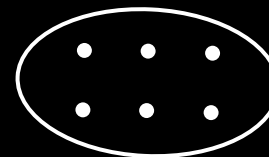
1. Every state counts as one (i.e. $\mu(\{a\}) = 1$)

~~2. Finite regions have finitely many states (i.e. $\mu(U) < \infty$)~~

3. Count is additive for disjoint sets (i.e. $\mu(\cup U_i) = \sum \mu(U_i)$)

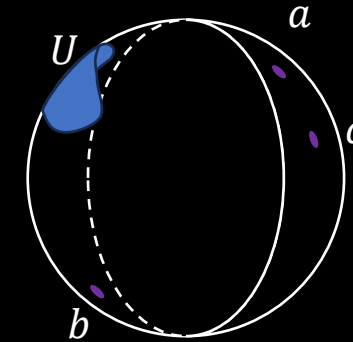
\Rightarrow Counting measure

What we use in classical discrete case



The problem with counting on the continuum

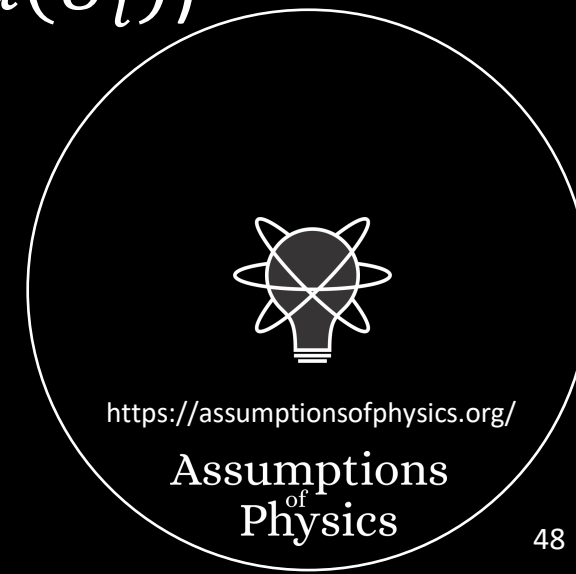
measure μ (a set of states) = how many states are there in the set



- ~~1. Every state counts as one (i.e. $\mu(\{a\}) = 1$)~~
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3. Count is additive for disjoint sets (i.e. $\mu(\cup U_i) = \sum \mu(U_i)$)

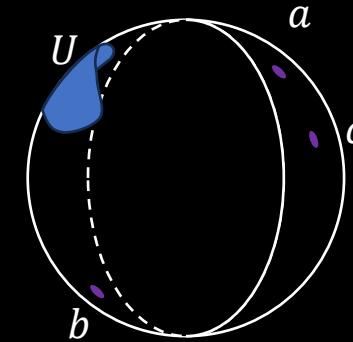
⇒ Lebesgue measure

What we use in classical continuum case (i.e. classical mechanics)



The problem with counting on the continuum

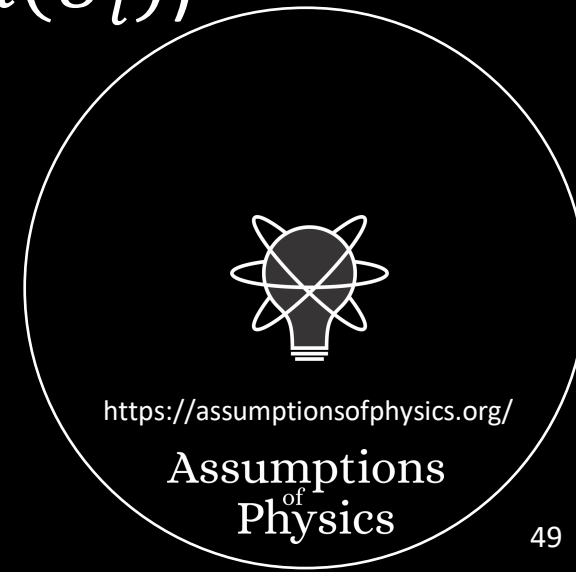
measure μ (a set of states) = how many states are there in the set



1. Every state counts as one (i.e. $\mu(\{a\}) = 1$)
2. Finite regions have finitely many states (i.e. $\mu(U) < \infty$)
- ~~3. Count is additive for disjoint sets (i.e. $\mu(\cup U_i) = \sum \mu(U_i)$)~~

⇒ “Quantum” measure

What we are implicitly using in quantum mechanics



We saw that thinking about entropy makes us understand how classical mechanics fails and gives rise to quantum mechanics

Can we use similar arguments to understand why quantum mechanics and general relativity are incompatible, and how to fix it?



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Assumptions
of
Physics

General relativity is a field theory

Particle mechanics

State defined by
position and momentum:

$$\begin{matrix} x^1 & x^2 & & x^n \\ p_1 & p_2 & \dots & p_n \end{matrix}$$

Finitely many variables

Field theory

State defined by
fields at each point:

$$A(x) \quad B(x) \quad \dots$$

Continuously (infinitely)
many variables

How do we count them?!?

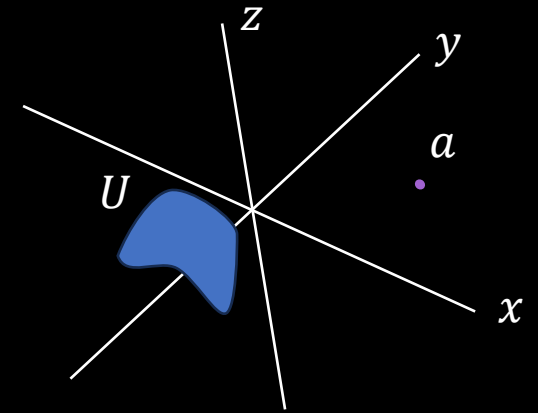


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Assumptions
of
Physics

The problem with counting (on the continuum)

$$\mu(\text{region of space}) = \text{Volume proportional to number of variables}$$



1. Every point has one variable (i.e. $\mu(\{a\}) = 1$)
2. Finite regions have finite volume (i.e. $\mu(U) < \infty$)
3. Count is additive for disjoint sets (i.e. $\mu(\cup U_i) = \sum \mu(U_i)$)

On a continuum, a finite region
has infinitely many points!



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Assumptions
of
Physics

Seems that “quantizing space-time” means putting a lower bound on the number of variables (degrees of freedom)

In the same way that a region of state space “with less than one state” makes no sense, a region of physical space with “less than one degree of freedom” makes no sense



I don't think we have the right math to solve this problem!



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Assumptions
of
Physics

But we don't develop theories by writing down assumptions and then derive observable consequences in a sequence of theorems and proofs. In physics, theories almost always start out as loose patchworks of ideas. Cleaning up the mess that physicists generate in theory development, and finding a neat set of assumptions from which the whole theory can be derived is often left to our colleagues in mathematical physics—a branch of mathematics, not of physics.

Sabine Hossenfelder – Lost in Math

-100,000 points

Math is not just a tool for calculation:
we use it to specify our theories



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Assumptions
of
Physics

Ultimately, it is up to the physicists to
develop a clear physical model

Mathematicians can help find holes... and give options

Clear physical model \Leftrightarrow clear math

Clarity, meaningfulness, consistency, ...
cannot be added after the fact



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Assumptions
of
Physics

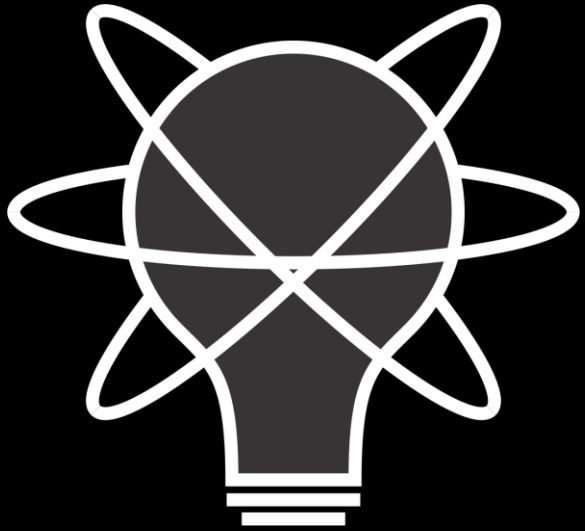
Big systems that work always
come from small systems that
work, they never come from big
systems that don't work



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Assumptions
of
Physics

The only way forward is to go back and
reconstruct everything from a minimal set of
physically meaningful assumptions



Assumptions
of
Physics

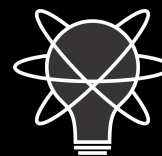
<https://assumptionsofphysics.org>

For more information, papers, presentations, ...

<https://www.youtube.com/@gcarcassi>

Videos with results and insights from the research

<https://assumptionsofphysics.org/foundation>



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Assumptions
of
Physics

The problem with counting on the continuum

We'd like to say:

1. Every state is a single case (i.e. $\mu(\{\psi\}) = 1$)
2. Finite continuous range carries finite information (i.e. $\mu(U) < \infty$)
3. Count is additive for disjoint sets (i.e. $\mu(\cup U_i) = \sum \mu(U_i)$)

Incompatible!

Pick two!

Discard 1 \Rightarrow Lebesgue measure

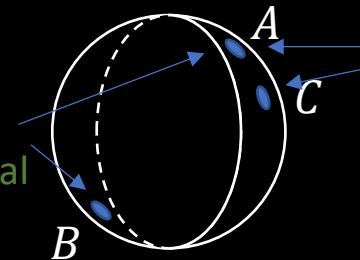
Discard 2 \Rightarrow counting measure

Discard 3 \Rightarrow "Quantum measure"

$$\mu(U) = 2^{\sup(s(\text{hull}(U)))}$$

Exponential of the maximum entropy reachable with convex combinations (statistical mixtures) of U (reduces to counting/Liouville measure)

Orthogonal states: additive
different states all else equal



$$\mu(\{A\}) = 2^0 = 1$$

$$\mu(\{A, B\}) = 2^1 = 2$$

$$\mu(\{A, C\}) < 2 = \mu(\{A\}) + \mu(\{C\})$$

Non-orthogonal states: different states but in different contexts
sub-additive

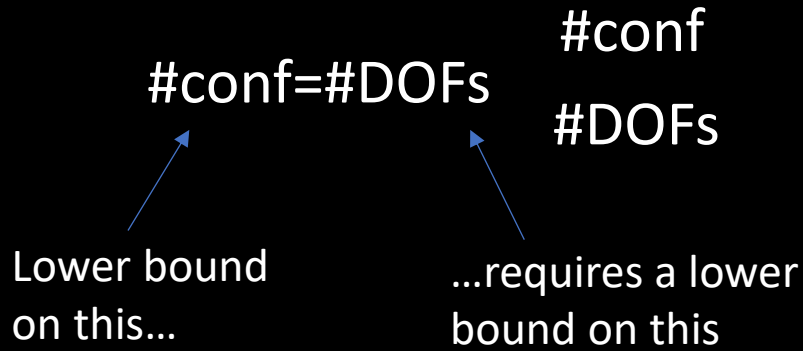
Quantum mechanics \Rightarrow lower bound
on #conf (entropy) on continuous DOF



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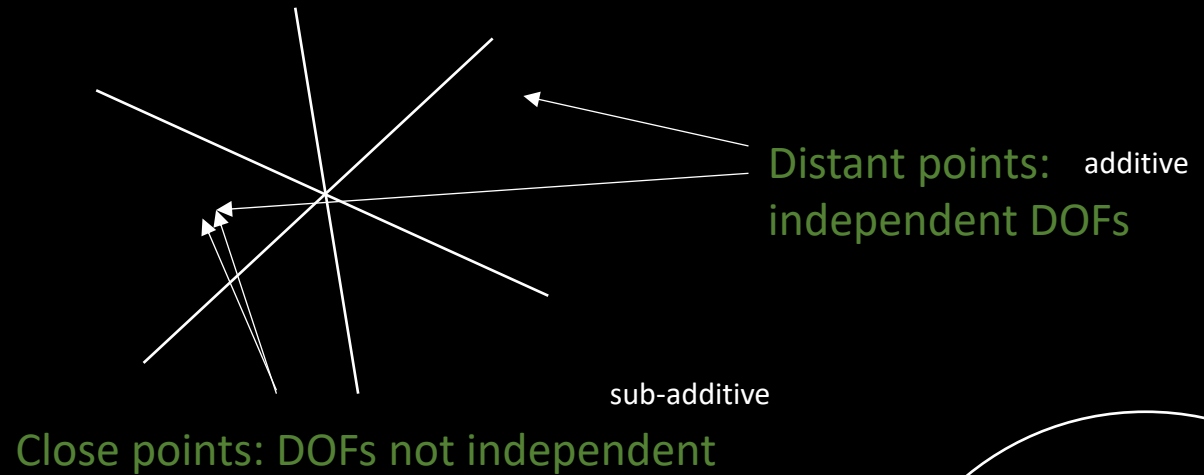
Assumptions
of
Physics

Conjecture: quantum gravity \Rightarrow lower bound on DOF count



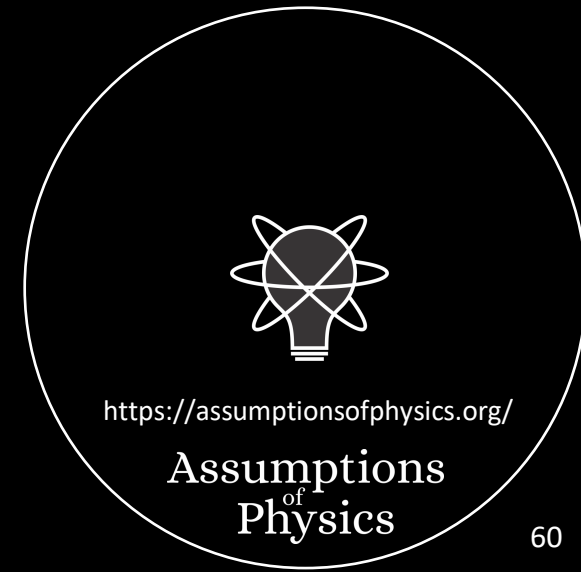
1. Every point is a single DOF (i.e. $\mu(\{x\}) = 1$)
2. Finite volume carries finitely many DOFs (i.e. $\mu(U) < \infty$)
3. Count is additive for disjoint regions (i.e. $\mu(\cup U_i) = \sum \mu(U_i)$)

Same problem!

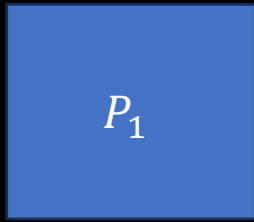


From QM: Lower bound on state count requires a severe revisitation of particle state space

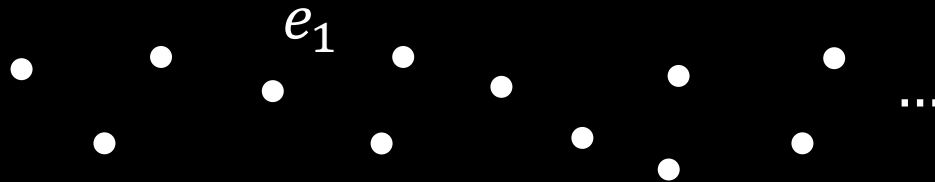
Does lower bound on DOF count require an equally severe revisitation of space-time?



Preparation

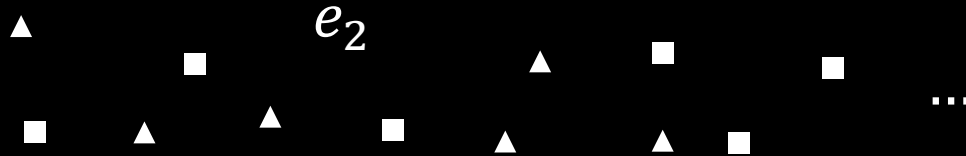
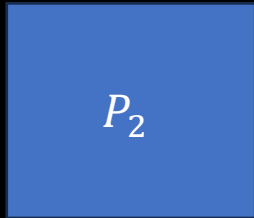


Ensemble



Identically
prepared
ensemble

There is no variability
in the ensemble



Pure
ensemble

No process can
reduce the variability
of the ensemble

There is no a such that $e_2 = pa + \bar{p}b$ for some $p \in (0,1)$ and $b \neq a$

In classical theories, all pure ensembles are identically prepared

No real distinction between ensembles and instances

In non-classical theories, no identically prepared ensembles

Zero entropy does NOT correspond to identically prepared ensembles

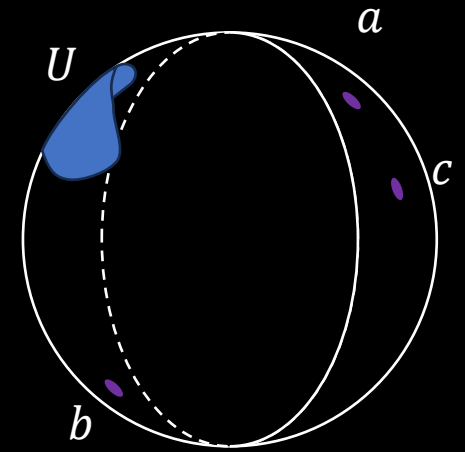


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Assumptions
of
Physics

The problem with counting on the continuum

$$\mu(\text{a set of states}) = \text{how many states are there in the set}$$



1. Every state counts as one (i.e. $\mu(\{e\}) = 1$)
2. Finite regions have finitely many states (i.e. $\mu(U) < \infty$)
3. Count is additive for disjoint sets (i.e. $\mu(\cup U_i) = \sum \mu(U_i)$)

On a continuum, a finite region
has infinitely many states!

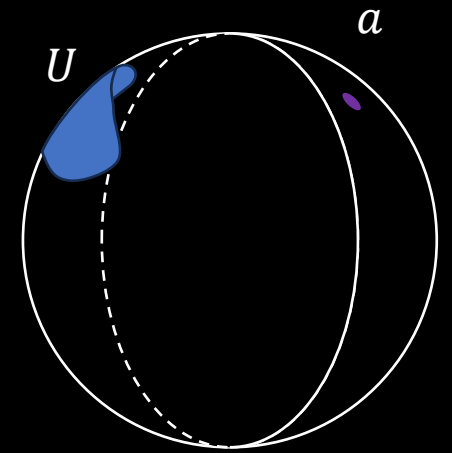


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Assumptions
of
Physics

The problem with counting on the continuum

measure μ (a set of states) = how many states are there in the set



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On a continuum, a finite region
has infinitely many states!



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Assumptions
of
Physics

How does the quantum measure work?

$$\mu(\{a\}) = \mu(\{b\}) = \mu(\{c\}) = 1$$

$$\mu(\{a, b\}) = 2 = \mu(\{a\}) + \mu(\{b\})$$

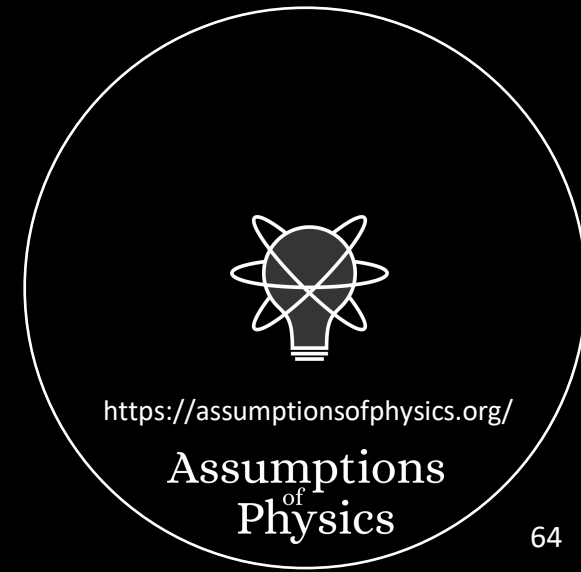
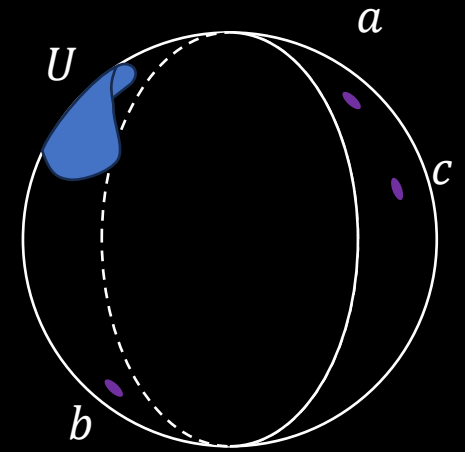
Additive when states are distinguishable (i.e. mutually exclusive)

$$\mu(\{a, c\}) < 2 = \mu(\{a\}) + \mu(\{c\})$$

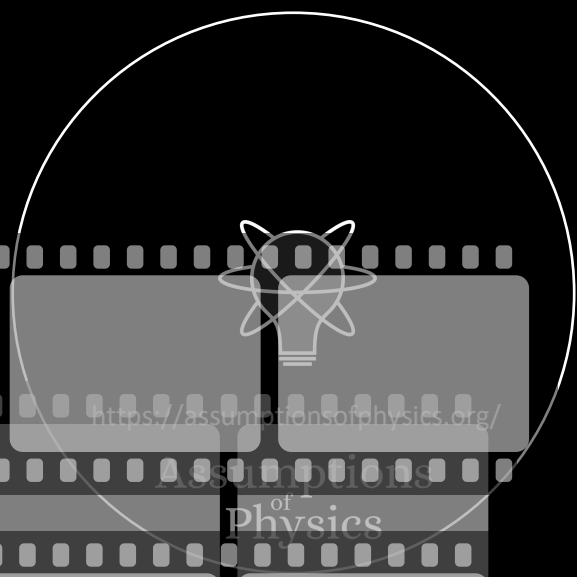
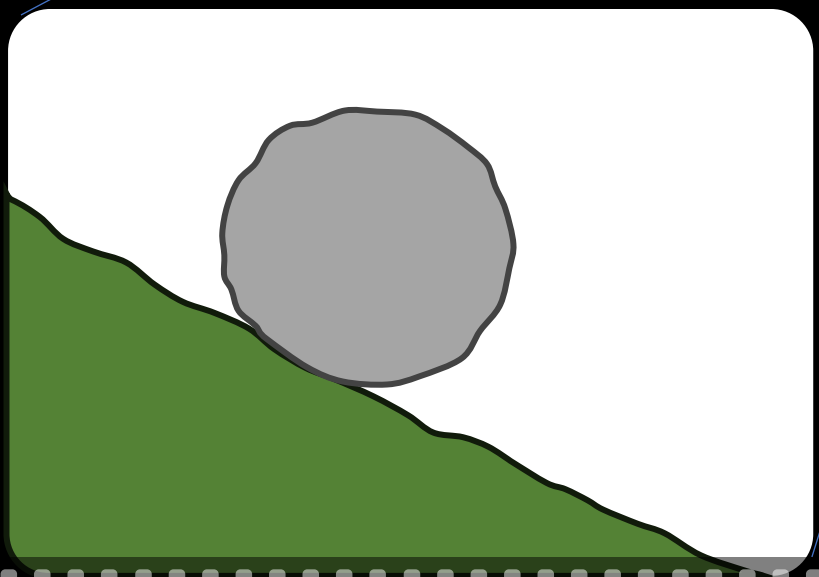
Sub-additive when states are not distinguishable (i.e. mutually exclusive)

**Non-additivity coincides
with non-distinguishability**

The measure counts
distinguishable (i.e. mutually
exclusive) states



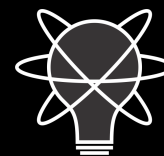
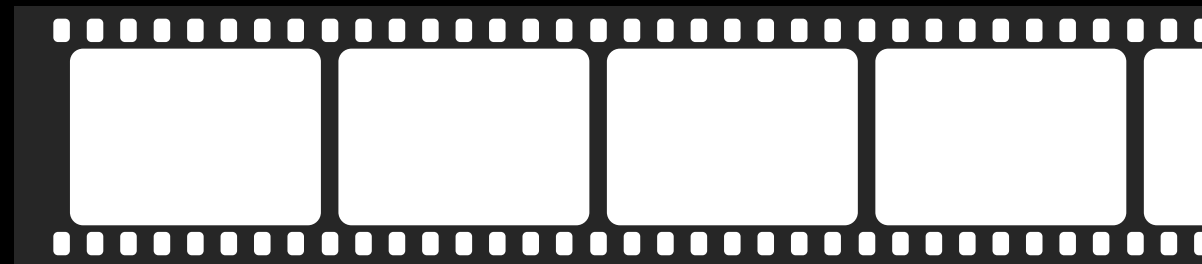
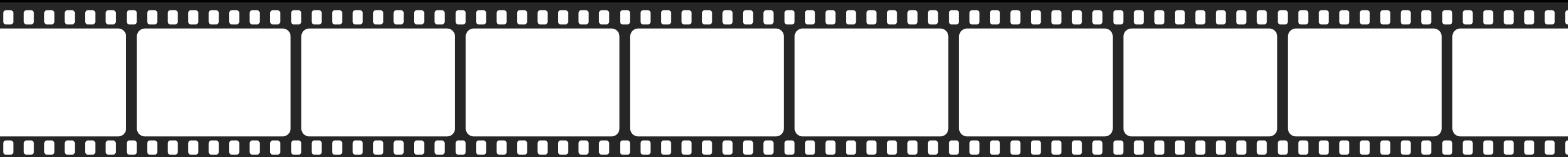
→ t



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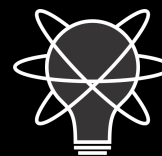
Assumptions
of
Physics

→ t



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Assumptions
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Physics



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Assumptions
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Physics

How should we count states?

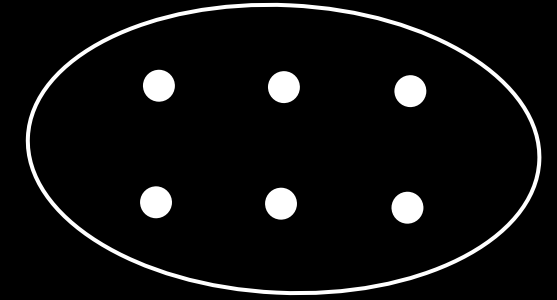


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Assumptions
of
Physics

For discrete classical systems

You simply count the points



For continuous classical systems

You use areas of position/momentum



A blue square is positioned within a white coordinate system defined by two perpendicular lines. The square is labeled with the expression $\Delta x \Delta p$.

$$\Delta x \Delta p$$

For quantum systems?



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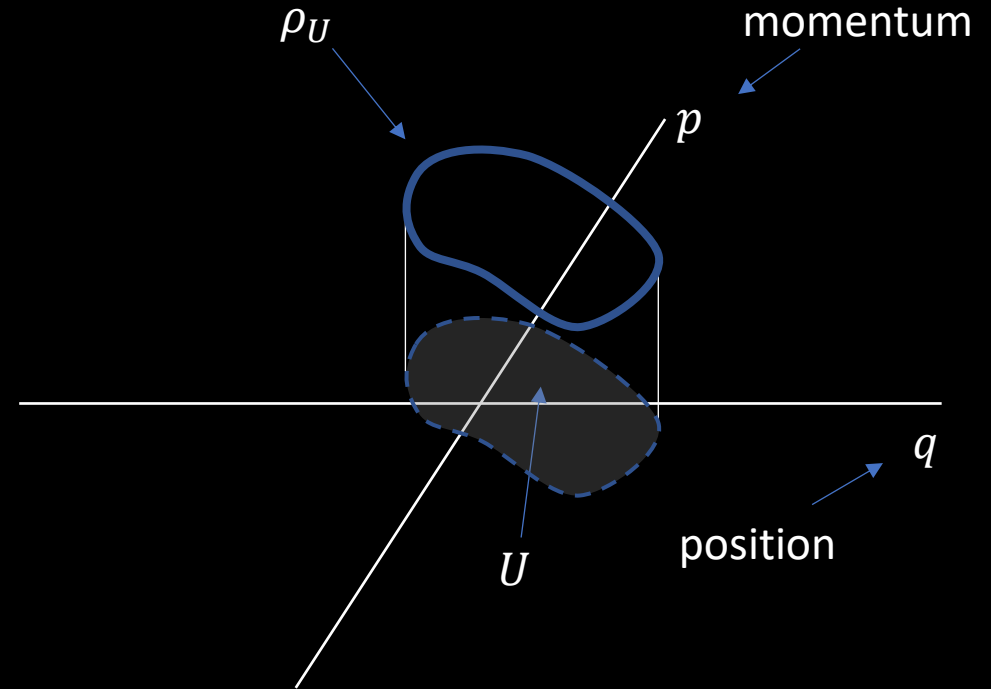
Assumptions
of
Physics

For classical systems

$$2^{S(\rho_U)} = \mu(U)$$

Exponential of the entropy
for a uniform
distribution over U

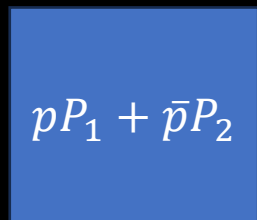
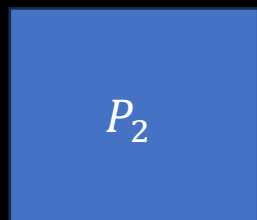
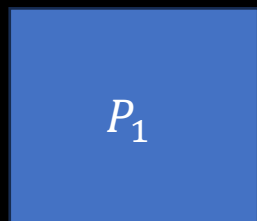
Count of states in U



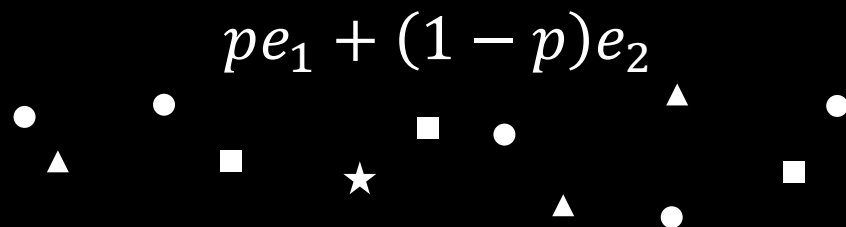
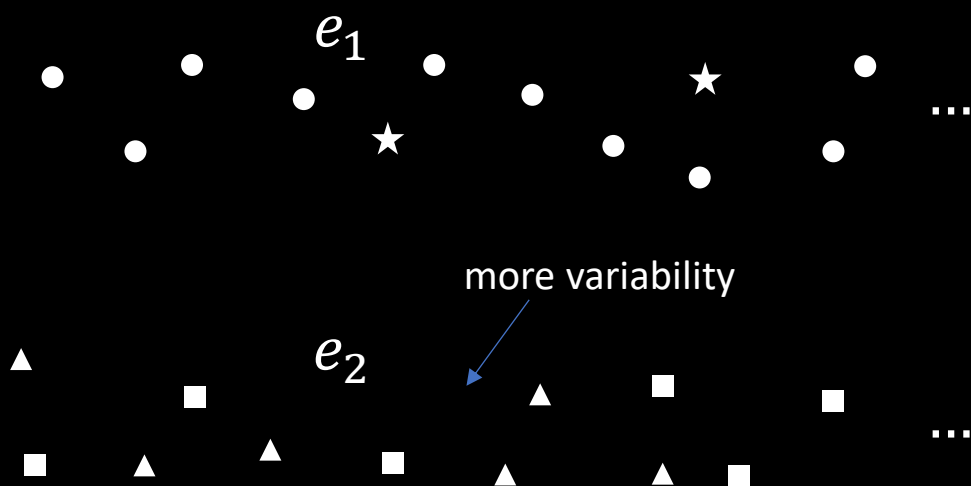
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Assumptions
of
Physics

Preparation



Ensemble



Entropy



Variability within an ensemble

$$S(e_1)$$

One instance is enough to tell e_1 and e_2 apart

Mutually exclusive

$$S(e_2)$$

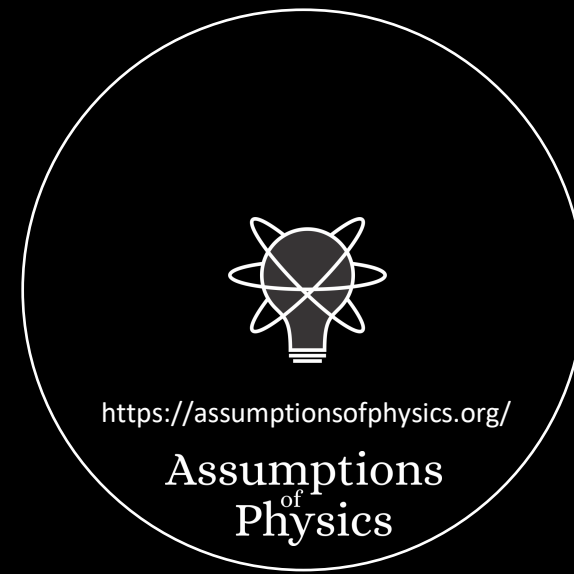
Orthogonal



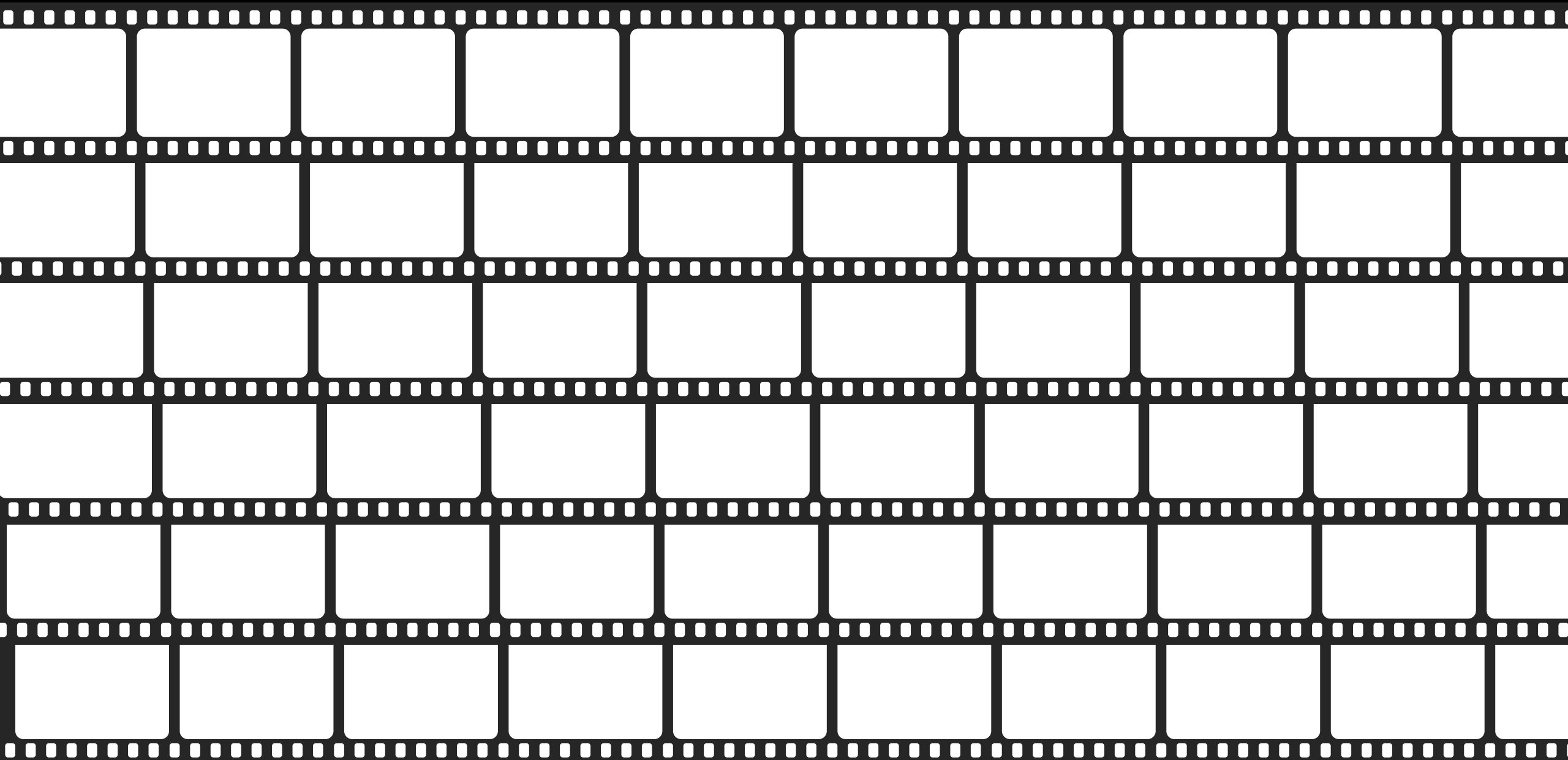
Maximal entropy increase if orthogonal

$$pS(e_1) + \bar{p}S(e_2) + I(p, \bar{p})$$

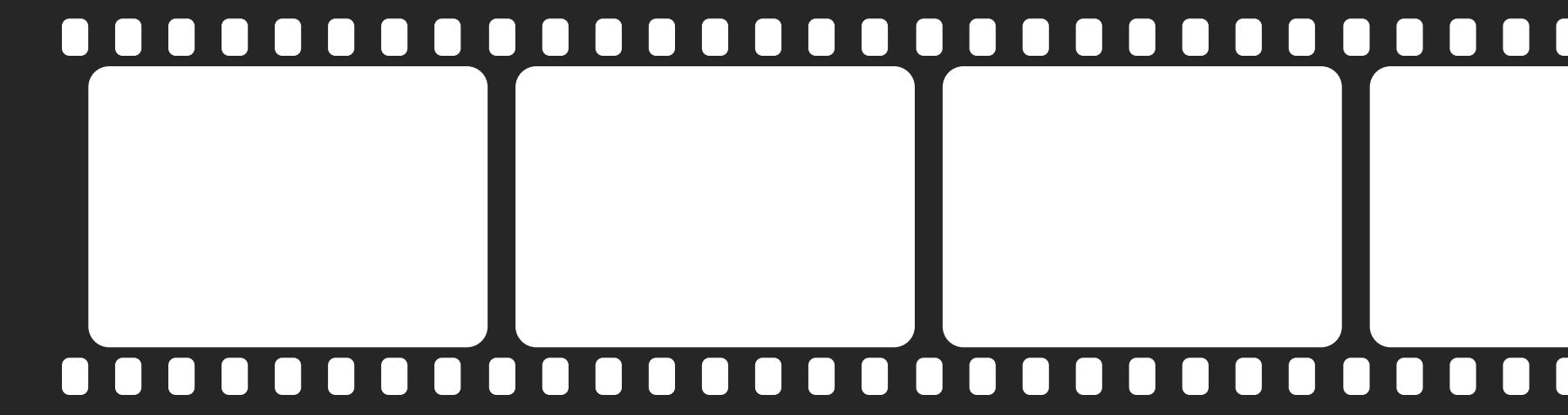
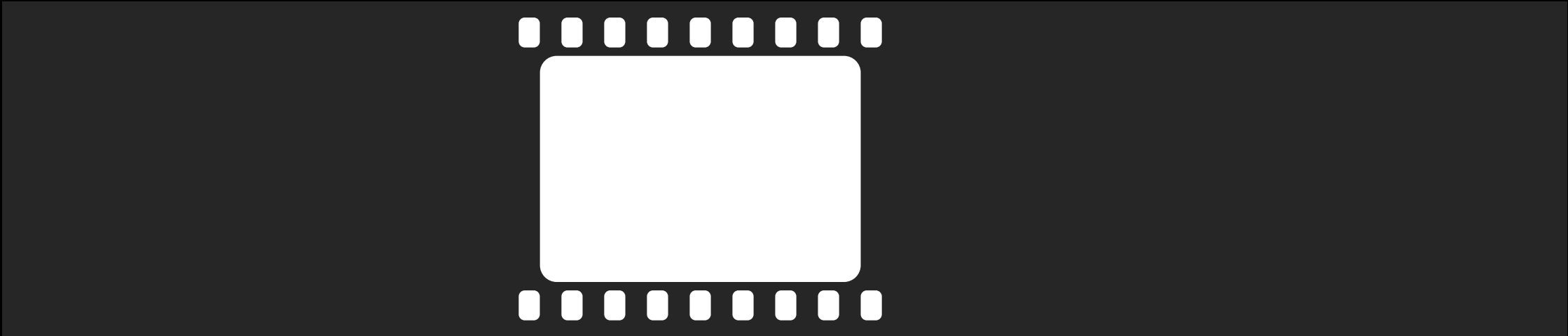
recovers Shannon entropy
 $I(p, \bar{p}) = -p \log p - \bar{p} \log \bar{p}$



→ t



Assumptions
of
Physics



Assumptions
of
Physics

- We don't know what it is
- What is a quantum system? What is a classical system?
- Electron/Proton/... all described by same equation. What makes them the same? Contextual. Inner dynamics is decoupled (it does not matter). Irreducibility?
- Can we reorganize all of physics through simple concepts? Most obvious?
- What



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Assumptions
of
Physics

- 100 years of QM: fake history, we got some math but we don't know what it means
- Uncertainty principle, two modes of evolution, contextuality, Dirac's correspondence principle (quantization)
- Interpretations (hidden variables) and reconstructions don't work
- Why does classical mechanics fail? Classical uncertainty principle.
- Motivations for the third law. Entropy of nothing. Entropy of one state.
- Classical mechanics recovered at high entropy, quantization is putting a lower bound on the entropy
- Entropy is about ensembles; physics is about ensembles; this is the fundamental mischaracterization of physics (film strip)
- Classical vs quantum: reducibility vs irreducibility; wave particle duality (electron a point particle, state probability distr; or electron a spread out blotch of energy? To tell the difference, needs lower entropy description)
- Back to counting states. 3 out of 2
- Space-time quantization



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Assumptions
of
Physics