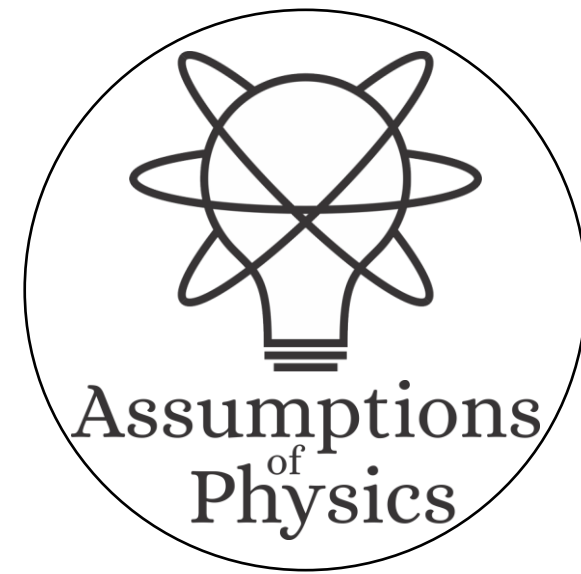
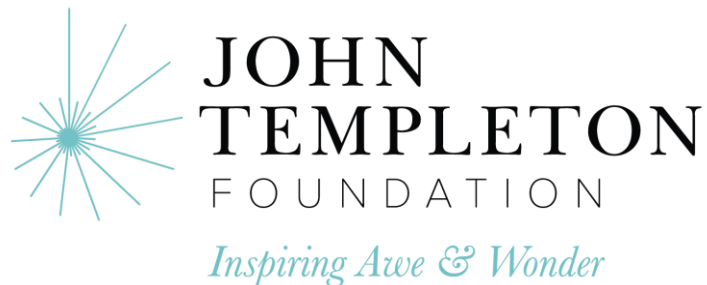


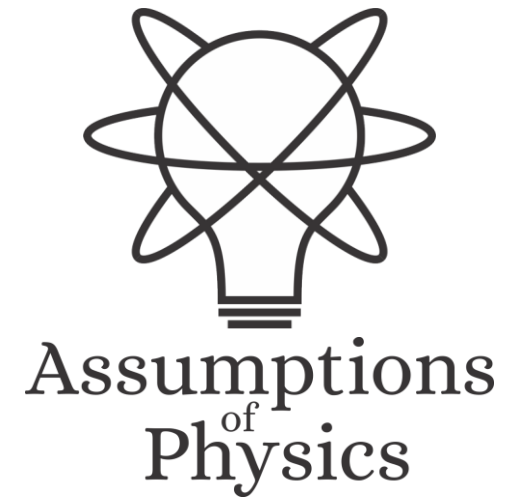
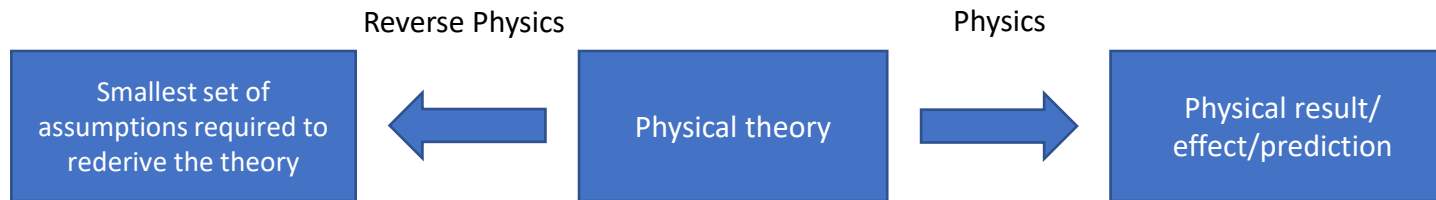
Non-additive measures for quantum probability?

Gabriele Carcassi and Christine Aidala
University of Michigan



Goal of the overall project

Identify a handful of physical starting points from which the basic laws can be rigorously derived



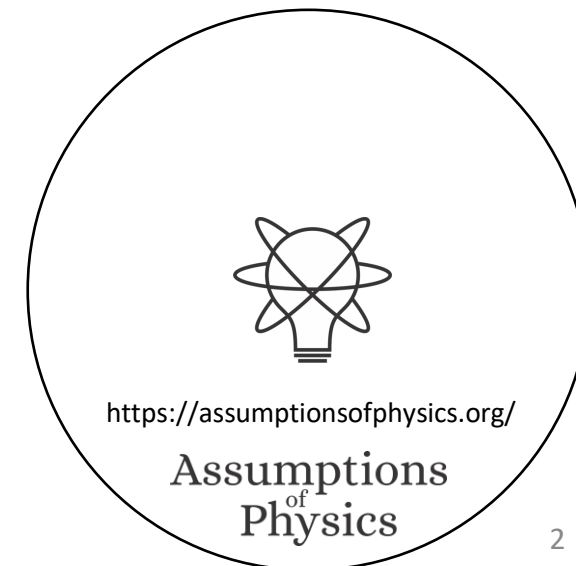
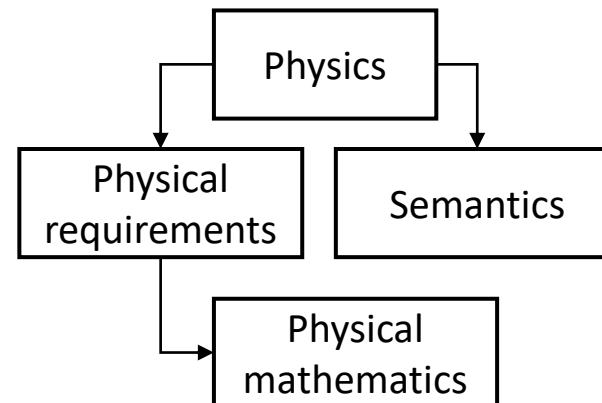
<https://assumptionsofphysics.org>

Reverse physics:

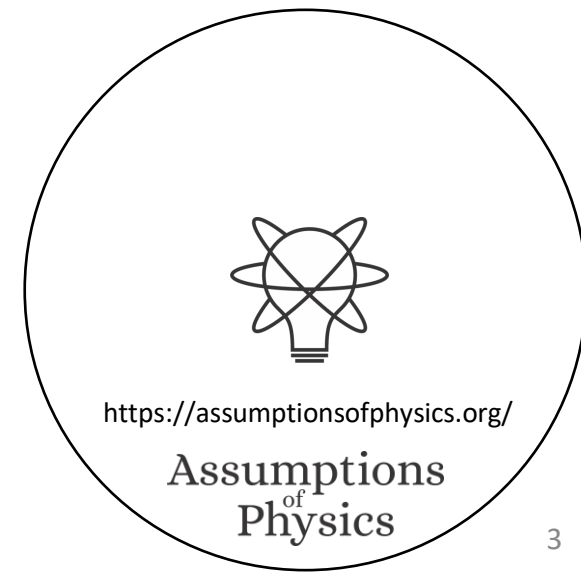
Start from the equations and identify principles and assumptions

Physical mathematics:

Start from scratch and rederive all mathematical structures from physical requirements



1. Quantum mechanics violates classical probability
2. Classical probability is additive measure theory
 \Rightarrow quantum mechanics violates additive measure theory
3. Can we use non-additive measure theory for a physically meaningful generalization?
4. What tools would (or would not) work?

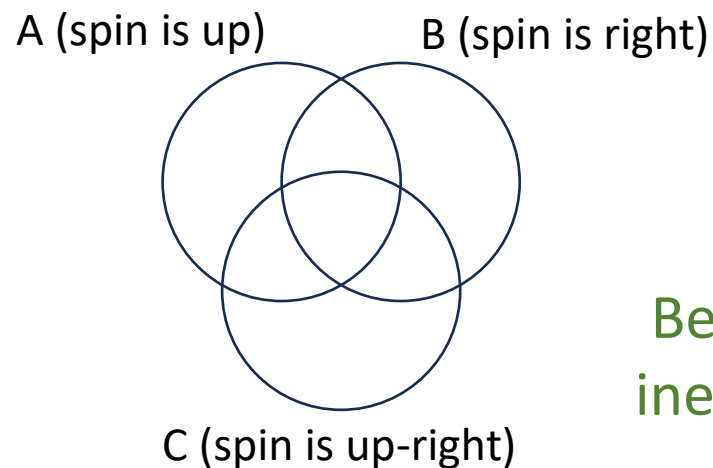


Classical probability and quantum mechanics

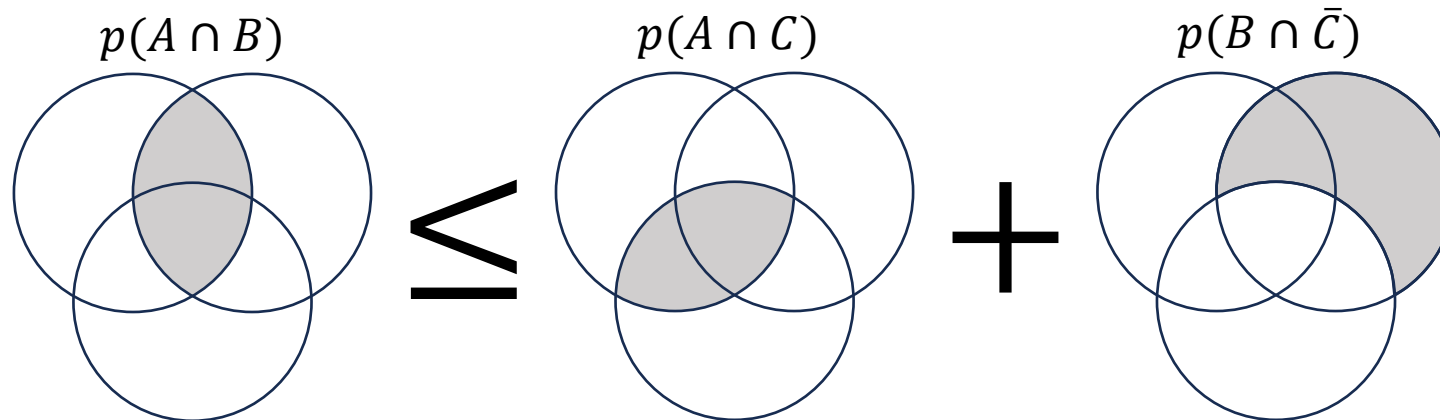


<https://assumptionsofphysics.org/>

Assumptions
of
Physics



Bell type
inequality



Violated by quantum mechanics

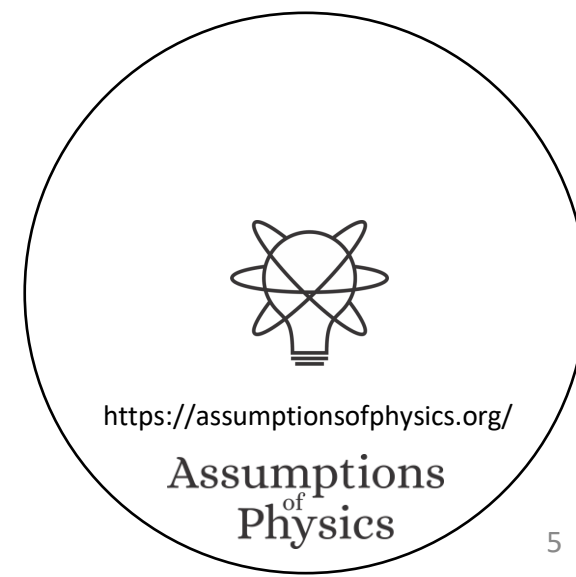
CHSH inequality

$$A, a, B, b \in \{-1, +1\} \quad (A - a, A + a) \in \{(0, \pm 2), (\pm 2, 0)\}$$

$$|\langle AB \rangle - \langle Ab \rangle - \langle aB \rangle - \langle ab \rangle| \leq 2$$

In quantum mechanics $2\sqrt{2}$

⇒ Failure of classical probability
(of standard measure theory)



Quasi-probability distribution

Wigner function

Sample space is not the state space

$$W(x, p) = \frac{1}{\pi\hbar} \int_X \psi^*(x + y) \psi(x - p) e^{2ipy/\hbar} dy$$

Can be negative:
no clear meaning

Recovers marginal distributions

$$|\psi(x)|^2 = \int W(x, p) dp \quad |\psi(p)|^2 = \int W(x, p) dx$$

$$\langle G \rangle = \int g(x, p) W(x, p) dx dp$$

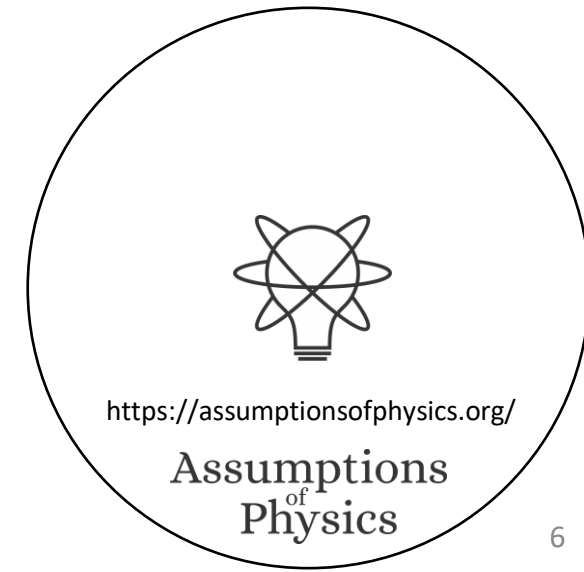
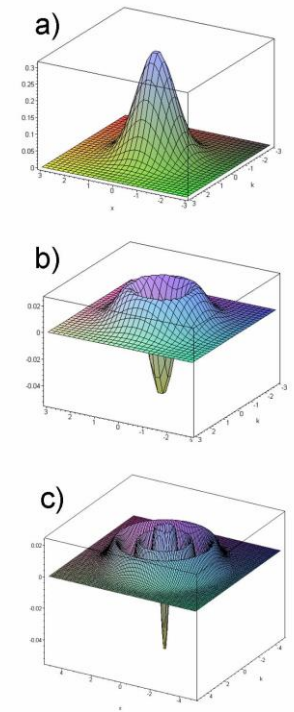
Acts as a probability density to recover probabilities and expectation values

Additive under statistical mixing

Density first (not measure first)

Not clear what densities are allowed or not

Wigner function for first two solutions of the harmonic oscillator



Can we find a generalization of probability theory that works in the same way in both classical and quantum mechanics?

Sample space is the respective state space

Probability defined as a set function on the sample space

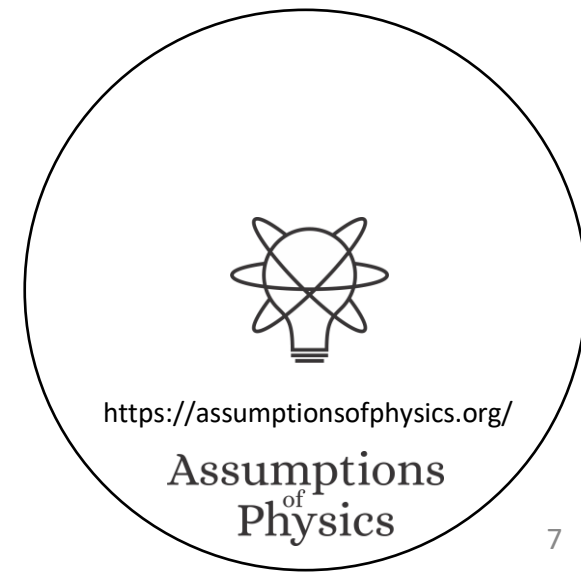
Expectations recovered as integral of some function of the sample space

...

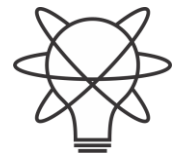
What set function for probability?

What integral for expectation?

What derivative for probability density?



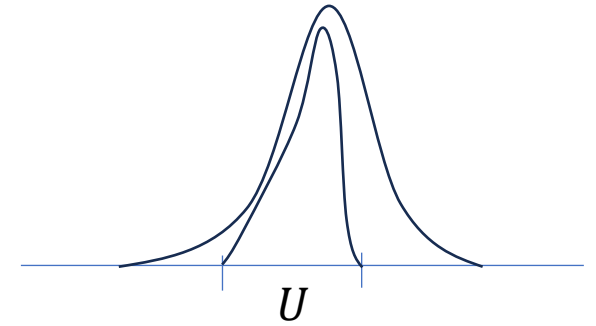
Ensemble spaces and mixing probability



<https://assumptionsofphysics.org/>

Assumptions
of
Physics

Probability as biggest coefficient when mixing measures



Measure over
whole space

$$p(A) = p(A|U)p(U) + p(A|U^c)p(U^c)$$

Measure over U

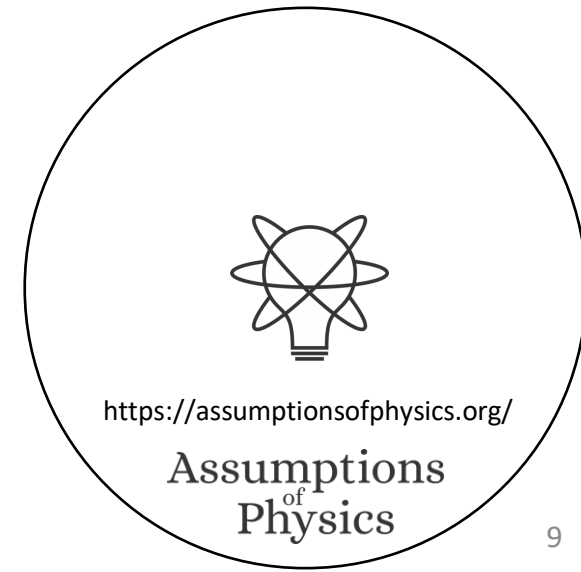
Coefficient of
convex combination

\mathcal{E} convex space of probability
measures

\mathcal{E}_U convex subspace of \mathcal{E}
with support U

Given $\mu \in \mathcal{E}$, $\mu(U) = \sup(\{p \in [0,1] \mid \exists x \in \mathcal{E}_U, y \in \mathcal{E}$
 $\mu = px + (1 - p)y\})$

This we can generalize to quantum mechanics
(to any convex set)



General structure for space of ensembles

Convex structure $px + (1 - p)y \in \mathcal{E} \quad p \in [0,1] \quad x, y \in \mathcal{E}$

Concave entropy $S: \mathcal{E} \rightarrow \mathbb{R}$

Increase of entropy due to mixing is bound

$$0 \leq S(px + (1 - p)y) - (pS(x) + (1 - p)S(y)) \leq I(p, 1 - p)$$

No increase

Equality

Semi-metric

Max increase = Shannon entropy

Orthogonality \equiv No overlap



<https://assumptionsofphysics.org/>

Assumptions
of
Physics

Mixing probability

Hull

$$\text{hull}(U) = \{\sum p_i e_i, e_i \in U\}$$

All convex combinations

\approx All probability distributions over U

Probability

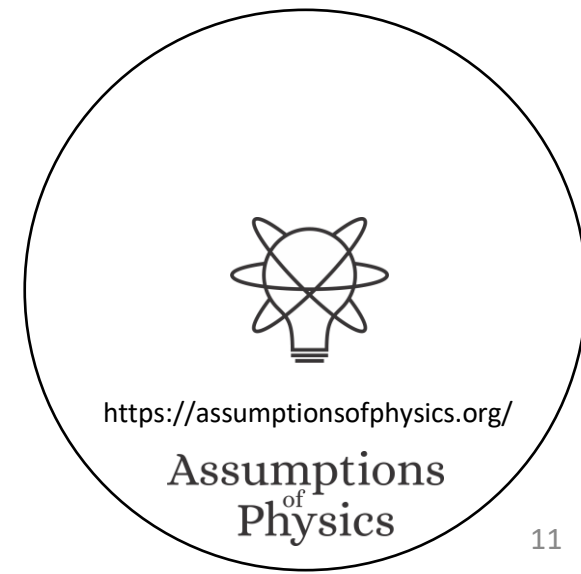
$$p_e(U) = \sup(\{p \in [0,1] \mid \exists x \in \text{hull}(U), y \in \mathcal{E} \\ e = px + (1-p)y\})$$

Biggest coefficient reachable by a mixture of elements of U

Set function

1. non-negative and unit bounded: $p_e(U) \in [0,1]$
2. monotone: $U \subseteq V \Rightarrow p_e(U) \leq p_e(V)$
3. subadditive: $p_e(U \cup V) \leq p_e(U) + p_e(V)$

Additive on subspaces that can be orthogonally decomposed



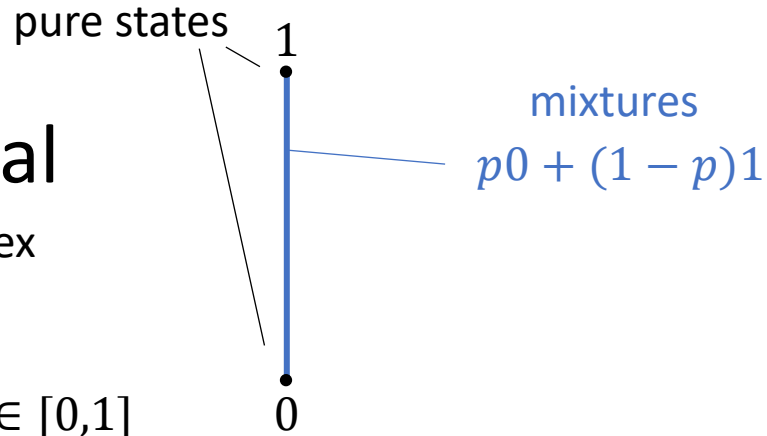
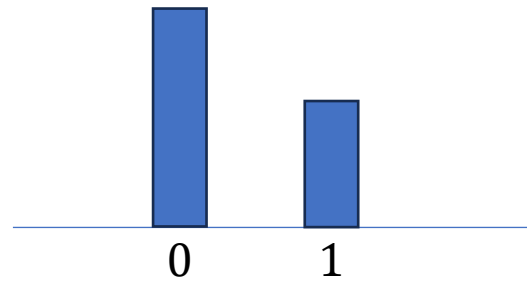
Two-state system (bit and qubit)

Classical

2D simplex

$$e \sim (p) \in [0,1]$$

$$\begin{aligned} p_e(\emptyset) &= 0 & p_e(\{0,1\}) &= 1 \\ p_e(0) &= p & p_e(1) &= 1 - p \end{aligned}$$



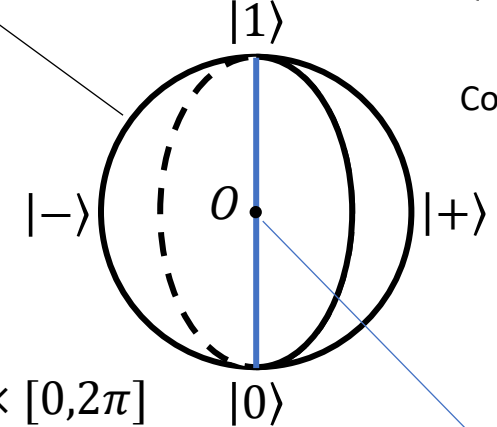
pure states (superpositions)

$$|\psi\rangle = \cos \hat{\theta}/2 |0\rangle + \sin \hat{\theta}/2 e^{i\hat{\phi}} |1\rangle$$

Quantum

Bloch ball

Complex projective
Hilbert space



$$e \sim (r, \theta, \phi) \in [0,1] \times [0,\pi] \times [0,2\pi]$$

$$p_e(\emptyset) = 0 \quad p_e(X) = 1$$

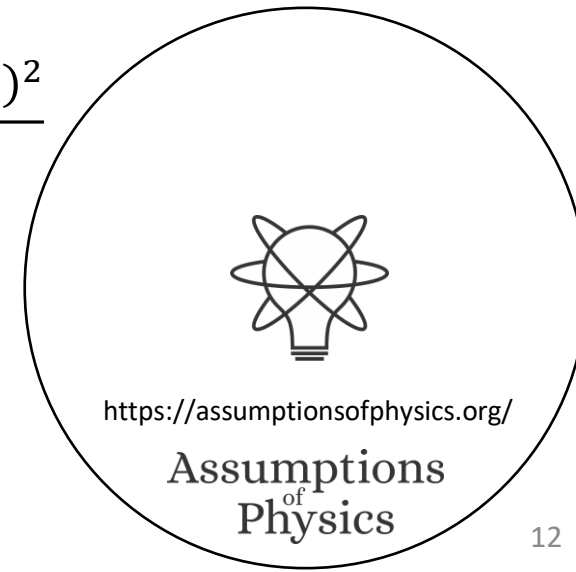
$$p_e(\{|\psi\rangle\}) = 1 - \frac{r^2 \sin^2 \Delta\theta + (1 - r \cos \Delta\theta)^2}{2(1 - r \cos \Delta\theta)}$$

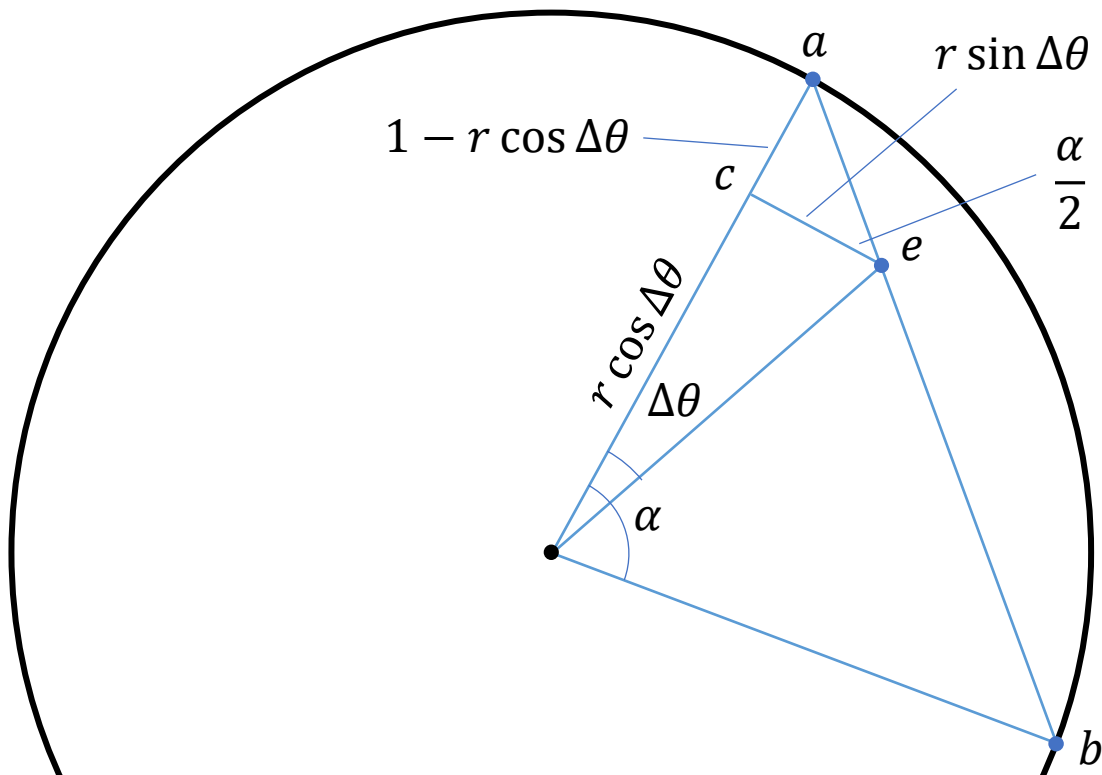
$$p_{|\psi\rangle}(|\phi\rangle) = \begin{cases} 1 & \psi = \phi \\ 0 & \psi \neq \phi \end{cases}$$

lowest entropy

$$p_o(|\phi\rangle) = \frac{1}{2}$$

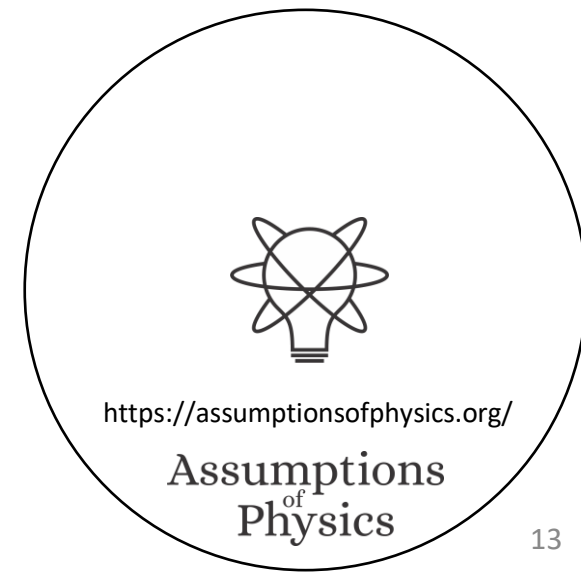
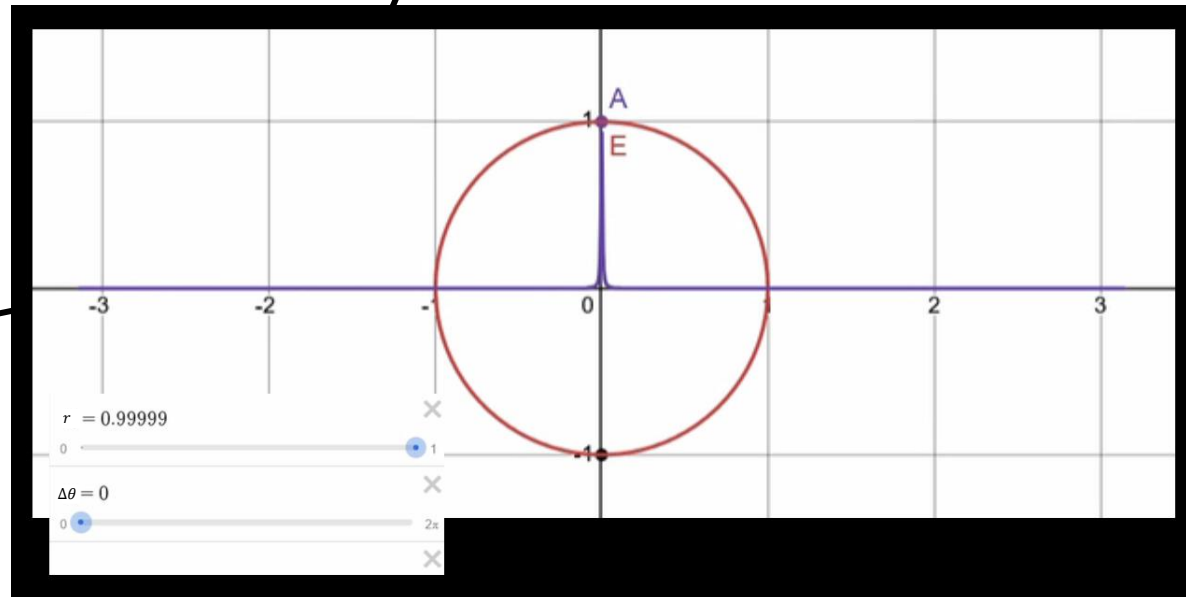
highest entropy





$$p = 1 - \frac{\overline{ae}}{\overline{ab}} = 1 - \frac{\overline{ae}}{2 \sin \frac{\alpha}{2}} = 1 - \frac{\overline{ae}}{2 \frac{\overline{ac}}{\overline{ae}}} = 1 - \frac{\overline{ae}^2}{2 \overline{ac}}$$

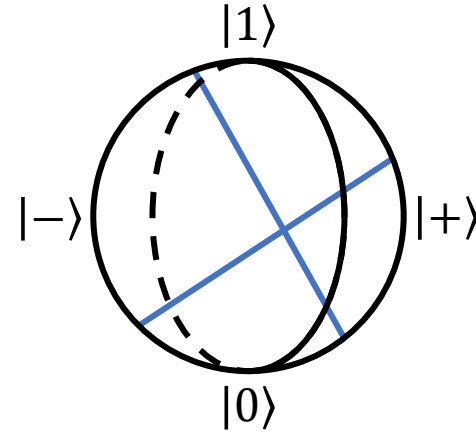
$$= 1 - \frac{(r \sin \Delta\theta)^2 + (1 - r \cos \Delta\theta)^2}{2(1 - r \cos \Delta\theta)}$$



The non-additivity is caused by the possible multiple decompositions



Classical probability:
single decomposition
in pure states



Quantum probability:
multiple decompositions
in pure states

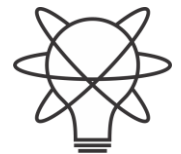
Finite probability on each point of a continuum

High degree of non-additivity

Uncountably many possible decompositions

$$p_o(X) = 1 \leq \sum_X p_o(\{x\}) = \frac{1}{2} |\mathbb{R}|$$

E.g. for maximally mixed state



<https://assumptionsofphysics.org/>

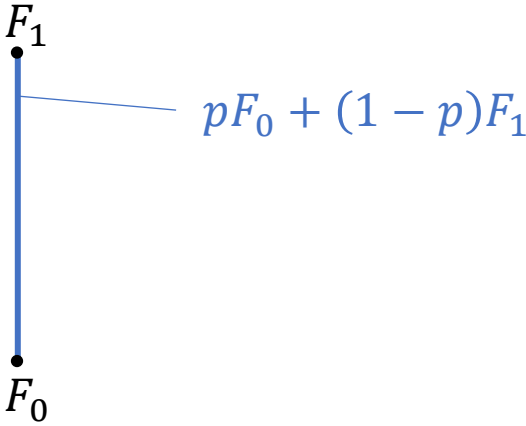
Assumptions
of
Physics

Random variables

$$F: \mathcal{E} \rightarrow \mathbb{R}$$
$$F(px + (1 - p)y) = pF(x) + (1 - p)F(y)$$

Classical

2D simplex



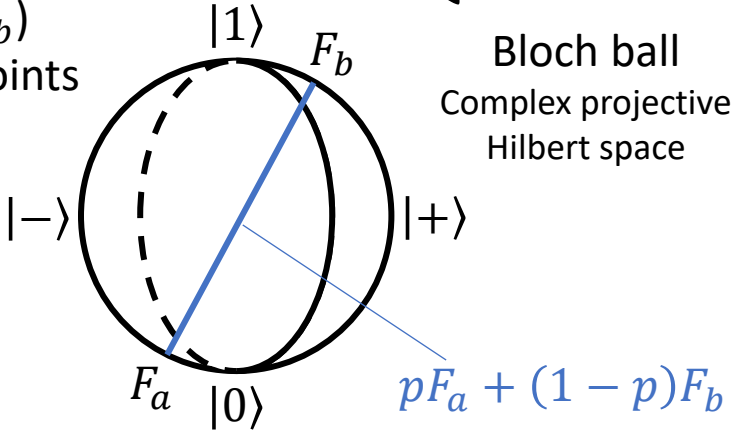
$$f(x) = F(\{x\})$$

$$E[F] = \sum f(x)p(x) = F(e)$$

Quantum

Bloch ball
Complex projective
Hilbert space

Max and min (F_a and F_b)
must be on opposite points

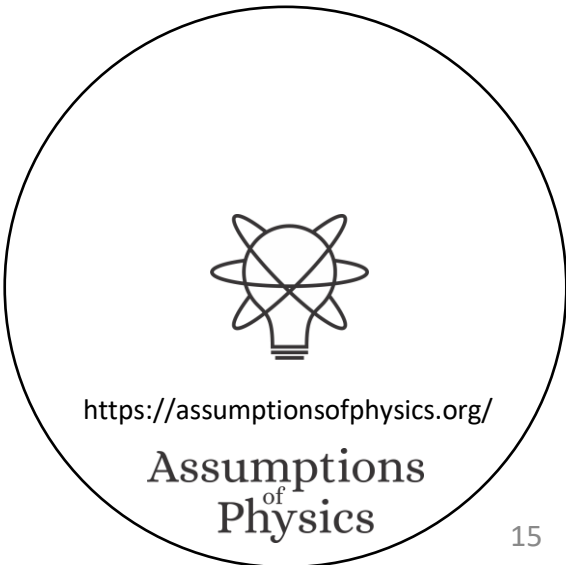


$$f(x) = F(\{x\})$$

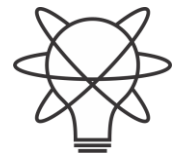
$$E[F] = \widehat{\sum} f(x)p(x) = F(e)$$

???

This is what needs to be generalized
to a continuum for integral/derivative



Entropy and count of states



<https://assumptionsofphysics.org/>

Assumptions
of
Physics

Classical statistical mechanics links count of states and entropy

$$S(\rho_U) = \log \mu(U)$$

Shannon/Gibbs entropy

Uniform distribution over U

Count of states

Fundamental postulate of statistical mechanics

Quantum statistical mechanics has a somewhat related expression

$$S(\rho_U) \leq \log(\dim_{\mathbb{C}}(\text{span}(U)))$$

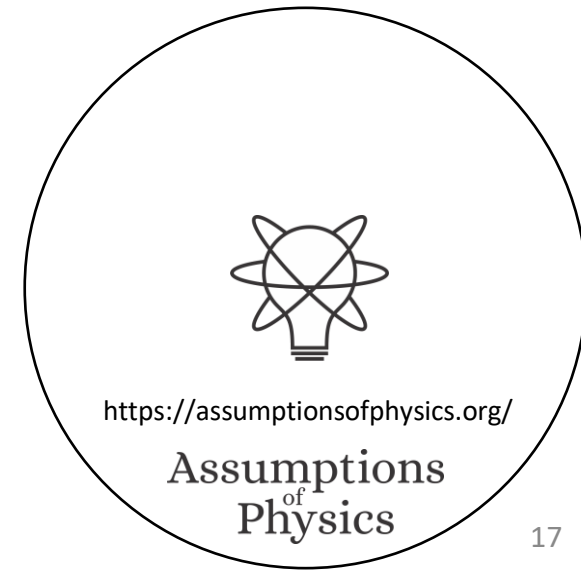
Von Neumann entropy

Uniform distribution over U

Dimensionality of subspace

Equal for a uniform distribution over $\text{span}(U)$

Generalized expression?



Proposition 1.24 (Exponential entropy subadditivity). Let $\mathbf{e}_1, \mathbf{e}_2 \in \mathcal{E}$. Let $S_1 = S(\mathbf{e}_1)$ and $S_2 = S(\mathbf{e}_2)$. Let $\mathbf{e} = p\mathbf{e}_1 + \bar{p}\mathbf{e}_2$ for some $p \in [0, 1]$ and $S = S(\mathbf{e})$. Then $2^S \leq 2^{S_1} + 2^{S_2}$, with the equality if and only if \mathbf{e}_1 and \mathbf{e}_2 are disjunct and $p = \frac{2^{S_1}}{2^{S_1} + 2^{S_2}}$.

Exponential of the entropy has key property

Proof. If p is fixed, the upper variability bound of entropy is saturated only if \mathbf{e}_1 and \mathbf{e}_2 are disjunct by definition. The entropy maximum for the mixed ensemble can only be achieved when the elements are disjunct, for some value of p .

Proof is mere calculation

$$0 = \frac{dS}{dp} = \frac{d}{dp} S(\mathbf{e}) = \frac{d}{dp} (-p \log p - \bar{p} \log \bar{p} + pS_1 + \bar{p}S_2)$$

$$= -\log p - 1 + \log \bar{p} + 1 + S_1 - S_2$$

$$\log \frac{p}{\bar{p}} = \log 2^{S_1} - \log 2^{S_2}$$

$$\log \frac{p}{1-p} = \log \frac{2^{S_1}}{2^{S_2}}$$

$$p2^{S_2} = (1-p)2^{S_1}$$

$$p(2^{S_1} + 2^{S_2}) = 2^{S_1}$$

$$p = \frac{2^{S_1}}{2^{S_1} + 2^{S_2}}$$

$$\bar{p} = 1 - \frac{2^{S_1}}{2^{S_1} + 2^{S_2}} = \frac{2^{S_2}}{2^{S_1} + 2^{S_2}}$$

$$S = S(\mathbf{e}) = -p \log p - \bar{p} \log \bar{p} + pS_1 + \bar{p}S_2$$

$$= -\frac{2^{S_1}}{2^{S_1} + 2^{S_2}} \log \frac{2^{S_1}}{2^{S_1} + 2^{S_2}} - \frac{2^{S_2}}{2^{S_1} + 2^{S_2}} \log \frac{2^{S_2}}{2^{S_1} + 2^{S_2}}$$

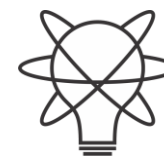
$$+ \frac{2^{S_1}}{2^{S_1} + 2^{S_2}} \log 2^{S_1} + \frac{2^{S_2}}{2^{S_1} + 2^{S_2}} \log 2^{S_2}$$

$$= \frac{2^{S_1}}{2^{S_1} + 2^{S_2}} \log (2^{S_1} + 2^{S_2}) + \frac{2^{S_2}}{2^{S_1} + 2^{S_2}} \log (2^{S_1} + 2^{S_2})$$

$$= \frac{2^{S_1} + 2^{S_2}}{2^{S_1} + 2^{S_2}} \log (2^{S_1} + 2^{S_2})$$

$$\log 2^S = \log (2^{S_1} + 2^{S_2})$$

$$2^S = 2^{S_1} + 2^{S_2}$$



<https://assumptionsofphysics.org/>

Assumptions
of
Physics

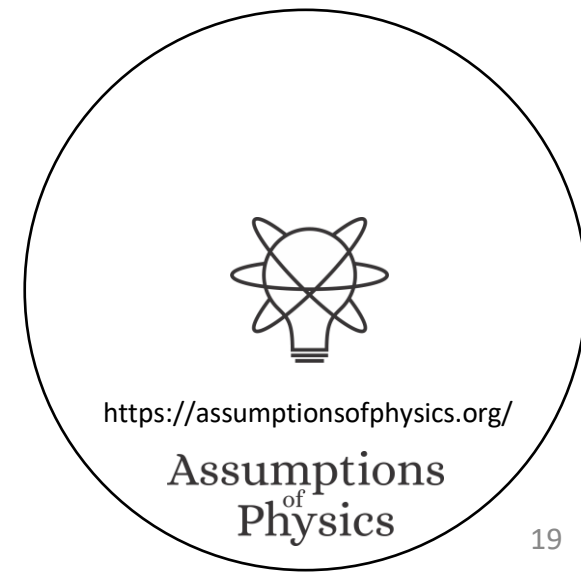
Define $\text{stateCount}(U) = \sup\left(2^{S(\text{hull}(U))}\right)$ ← Exponential of the highest entropy reachable through convex combinations

Set function

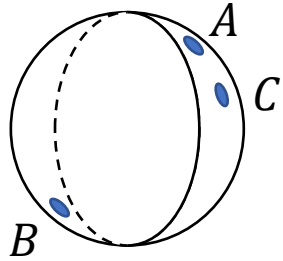
1. non-negative: $\text{stateCount}(U) \in [0, +\infty]$
2. monotone: $U \subseteq V \Rightarrow \text{stateCount}(U) \leq \text{stateCount}(V)$
3. subadditive: $\text{stateCount}(U \cup V) \leq \text{stateCount}(U) + \text{stateCount}(V)$

Additive on subspaces that can be orthogonally decomposed

Recovers number of states for discrete classical spaces,
Liouville measure for continuous classical spaces,
dimension of the subspaces in quantum mechanics



Need for non-additive measure



$$\mu(\{A\}) = 2^0 = 1$$

$$\mu(\{A, B\}) = 2^1 = 2 \quad \text{not additive}$$

$$\mu(\{A, C\}) < 2 = \mu(\{A\}) + \mu(\{C\})$$

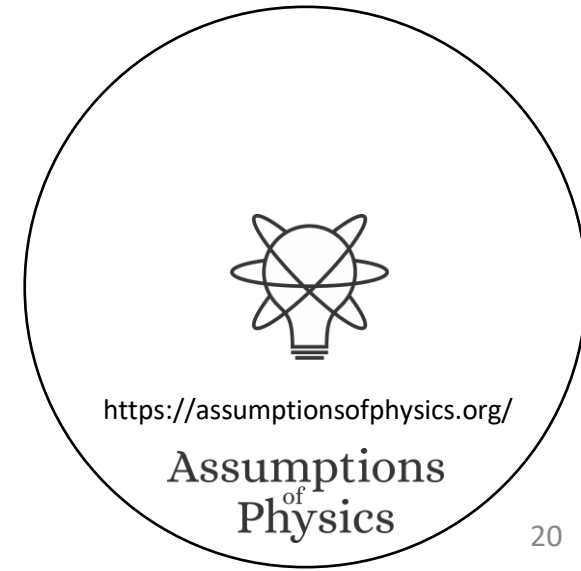
In quantum mechanics, literally $1 + 1 \leq 2$

Pick two!

1. Single point is a single case (i.e. $\mu(\{\psi\}) = 1$)
2. Finite range carries finite information (i.e. $\mu(U) < \infty$)
3. Measure is additive for disjoint sets (i.e. $\mu(\cup U_i) = \sum \mu(U_i)$)

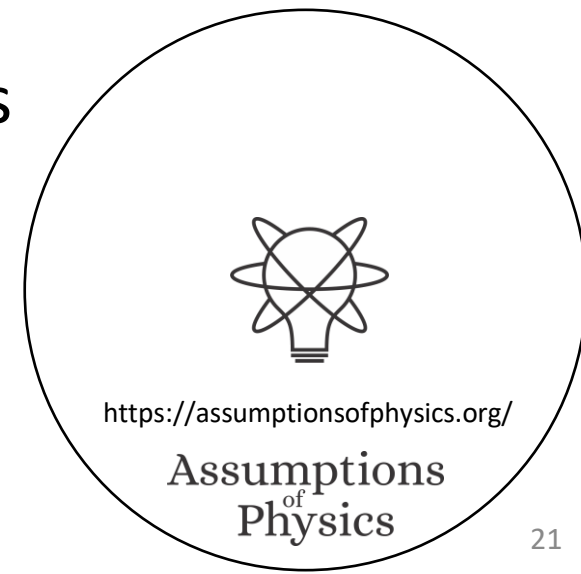
Physically, we count states all else equal
Contextuality \Leftrightarrow non-additive measure

	Single point		Finite continuous range	
	$\mu(U)$	$\log \mu(U)$	$\mu(U)$	$\log \mu(U)$
Counting measure $\mu(U) = \#U$ Number of points	1	0	$+\infty$	$+\infty$
Lebesgue measure $\mu([a, b]) = b - a$ Interval size	0	$-\infty$	$< \infty$	$< \infty$
"Quantized" measure $\mu(U) = \sup(2^{S(\text{hull}(U))})$ Entropy over uniform distribution	1	0	$< \infty$	$< \infty$



Summary

- In Assumptions of Physics (<https://assumptionsofphysics.org>) we are looking for a generalized setting for both classical and quantum probability
 - The setting is the convex space of ensembles at equilibrium (i.e. statistical distributions): can be mixed and have an entropy defined
- This generalization requires non-additive measures to deal with multiple decompositions (i.e. quantum contextuality)
 - Non-additive measures are not typically used in quantum foundations
 - Looking for discussion/advice/collaboration/...
- Same generalization connects other areas of math and physics
 - Information theory and information geometry
 - Functional analysis and spectral theory
 - Order theory and quantum logic



For more information

Gabriele Carcassi



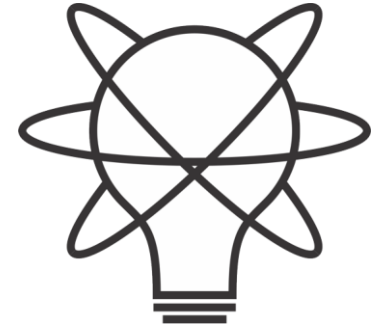
carcassi@umich.edu

Instructive videos for undergrads and above

<https://www.youtube.com/@gcarcassi>

<https://www.youtube.com/@AssumptionsofPhysicsResearch>

Videos about active research

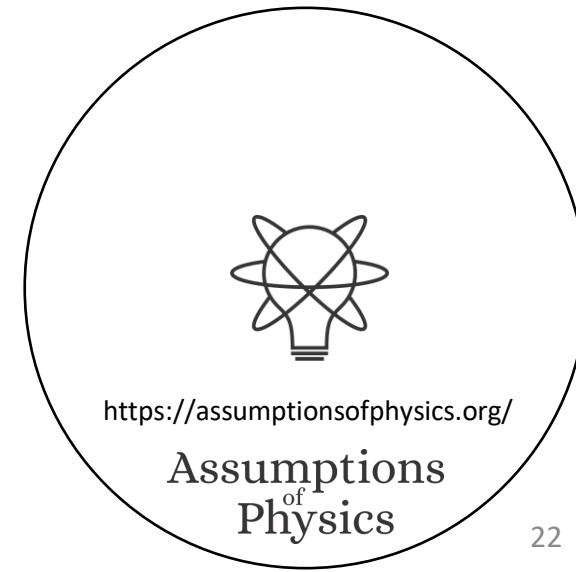


Assumptions
of
Physics

<https://assumptionsofphysics.org>

Thanks to all participants that have shared insights and ideas

Special thanks to Vicenç Torra and Zuzana Ontkovičová





<https://assumptionsofphysics.org/>

**Assumptions
of
Physics**