

Comparative Quantum Monte Carlo studies of the two-dimensional bond-Su-Schrieffer-Heeger and Holstein models

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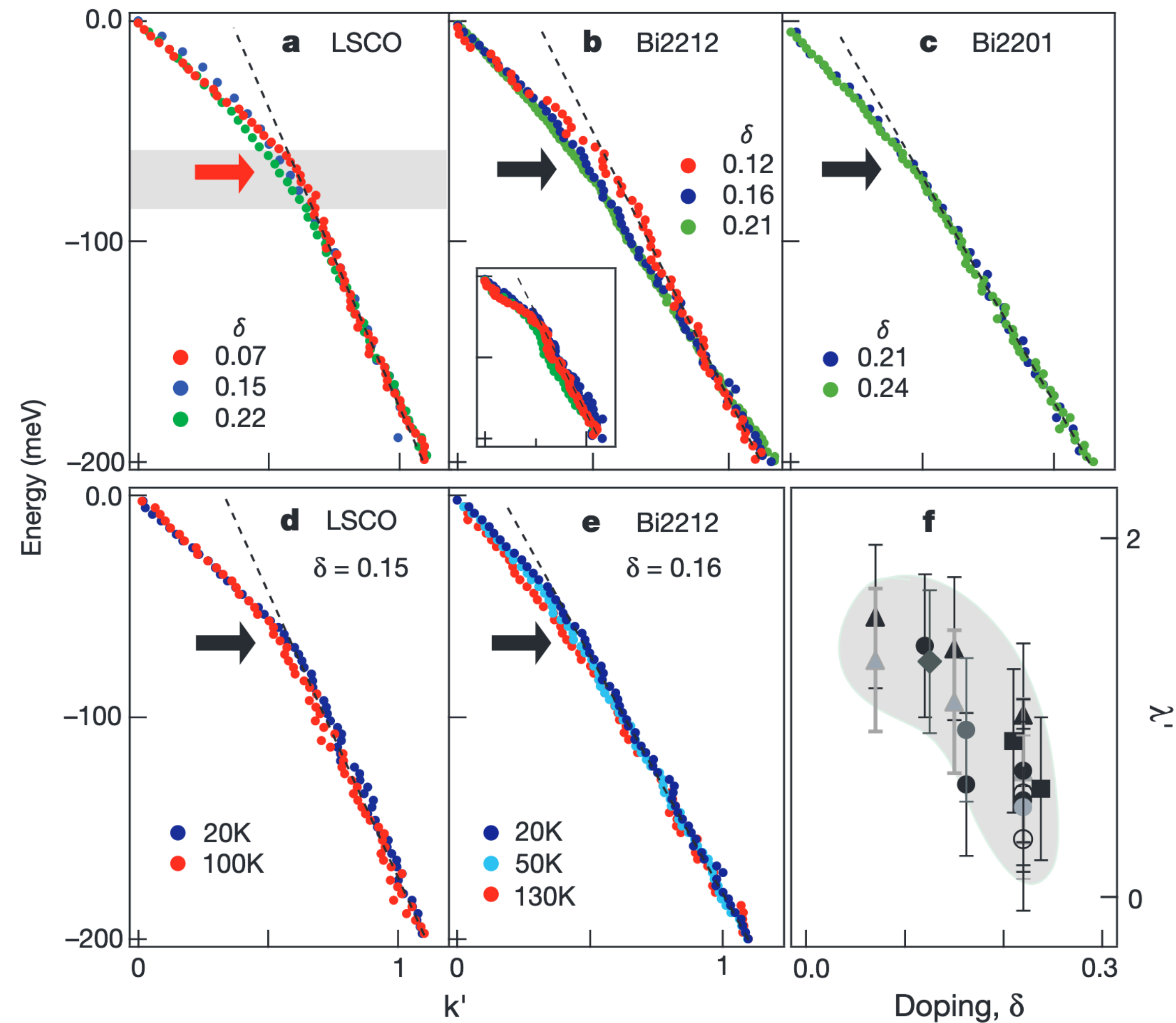
APS March Meeting

10 March 2023



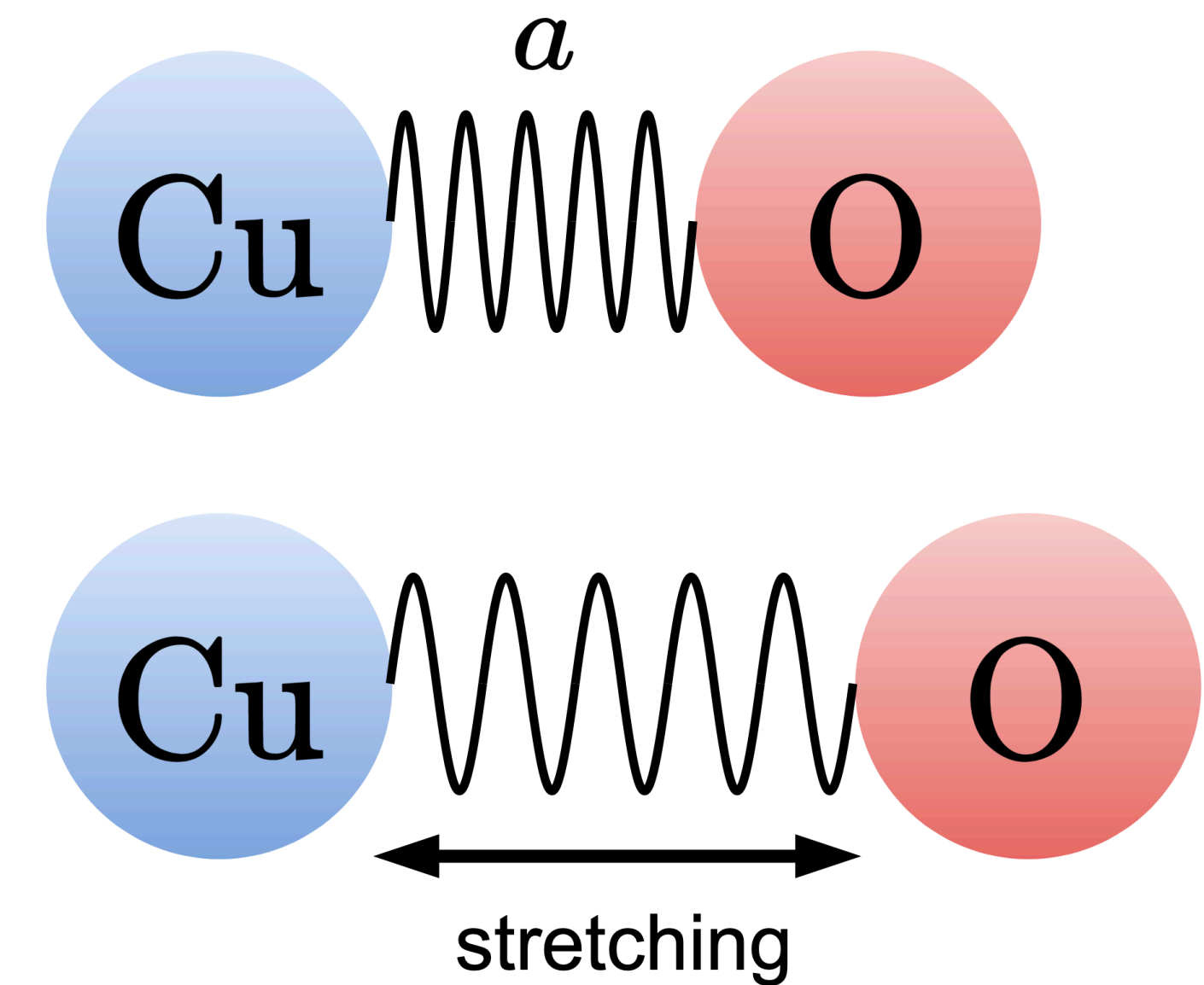
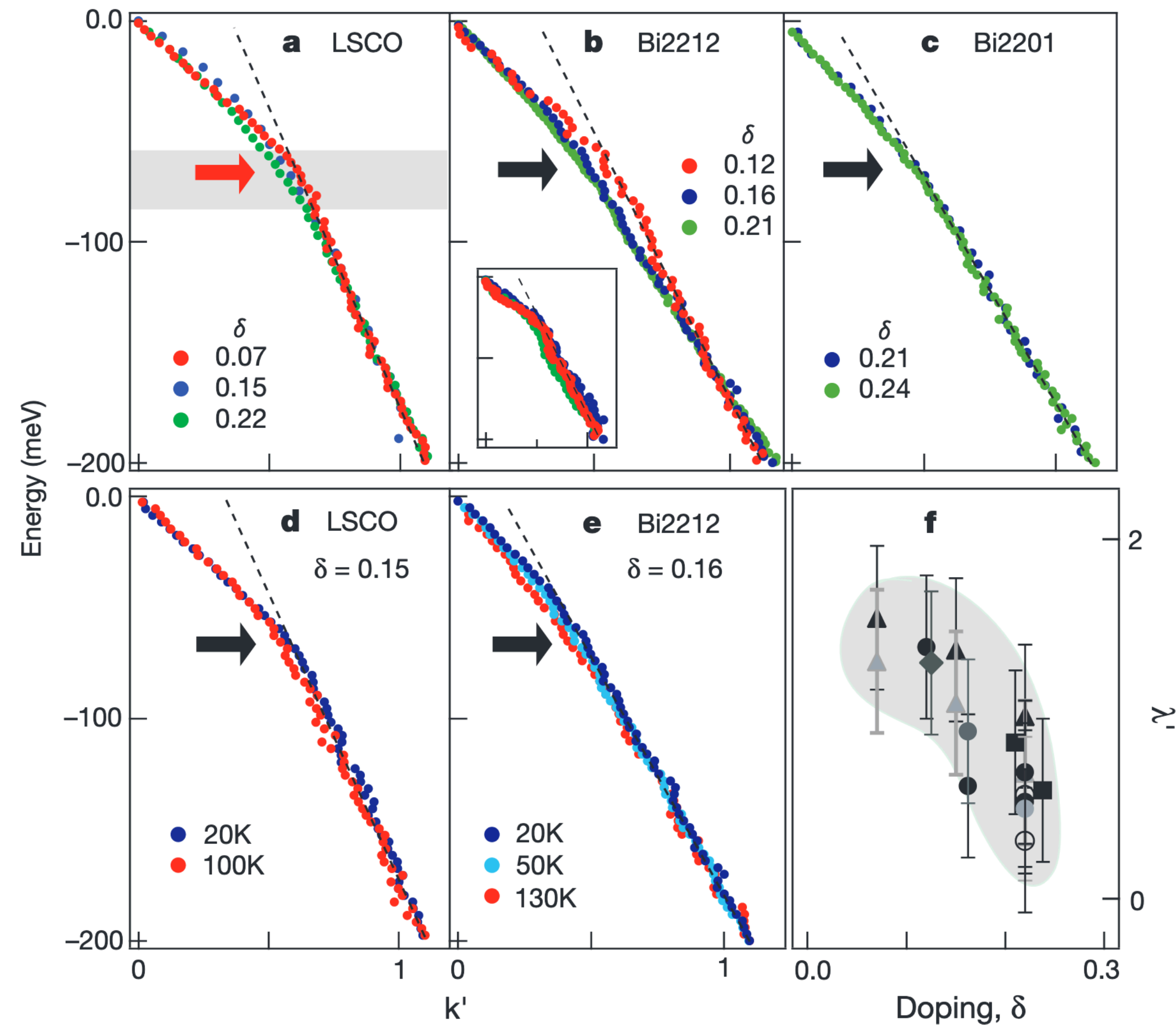
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Electron-phonon (e-ph) interactions in materials



A. Lanzara et al. Nature **412**, 510 (2001)

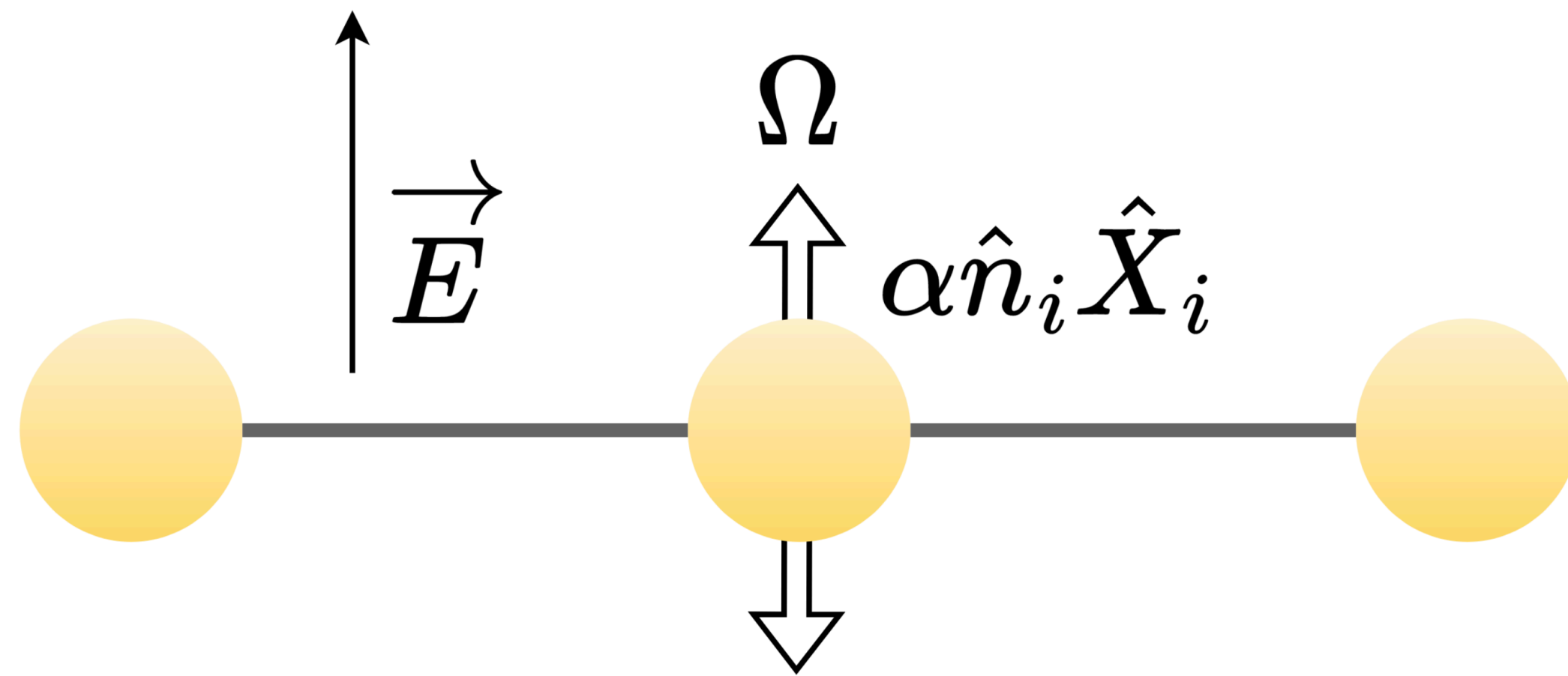
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Models of e-ph interactions

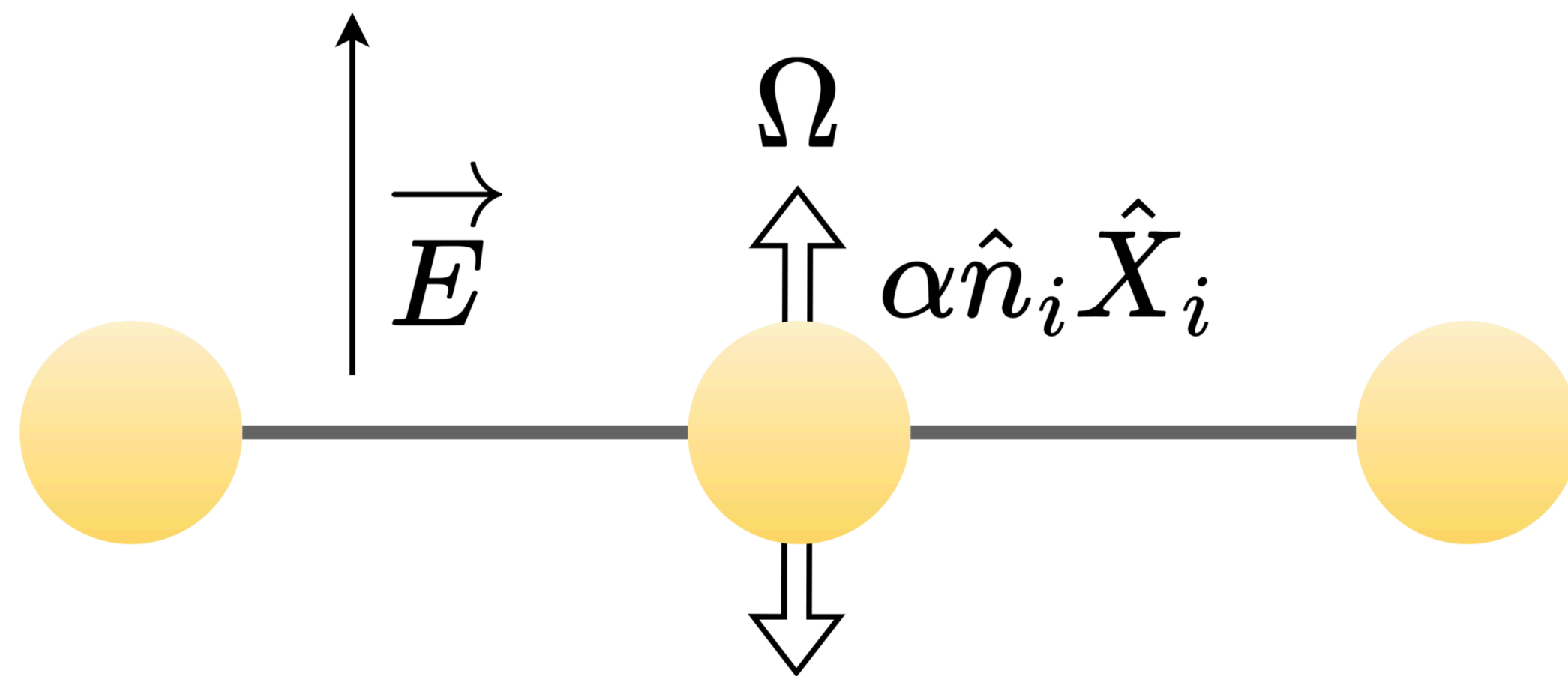
“Buckling” modes



Holstein model

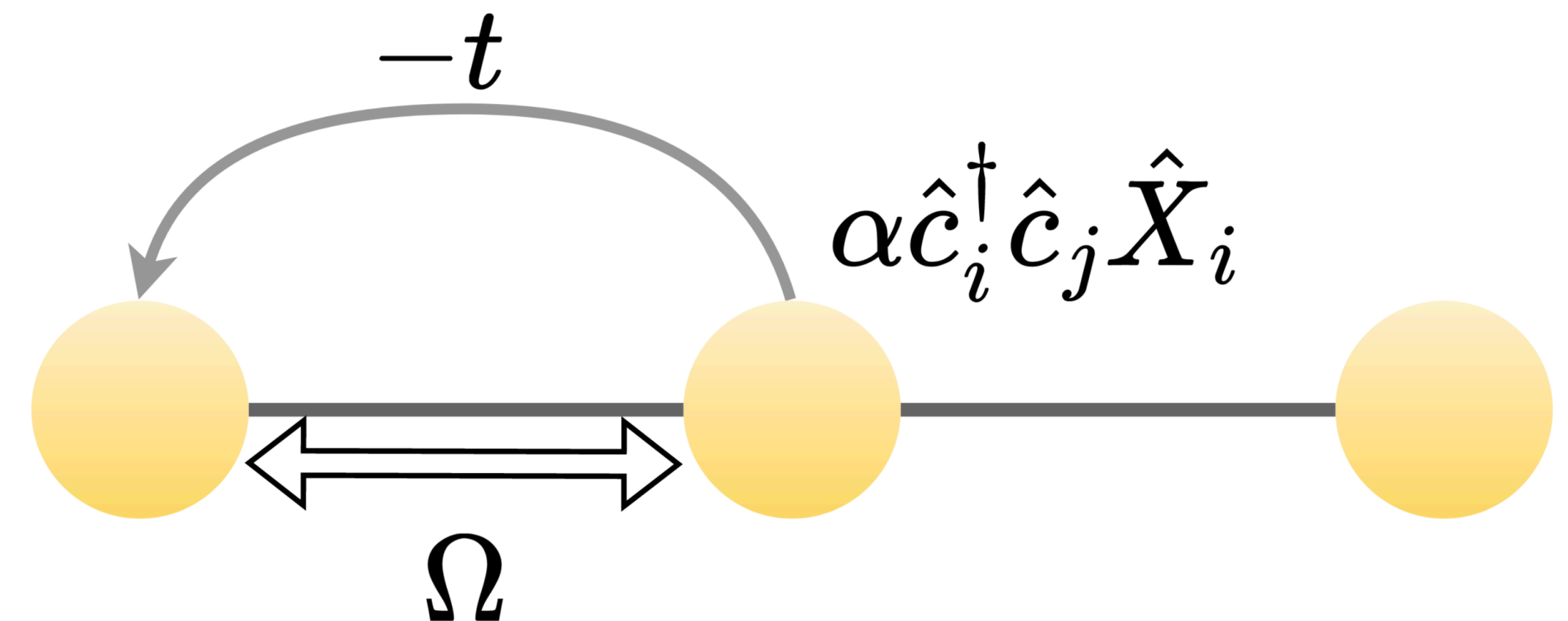
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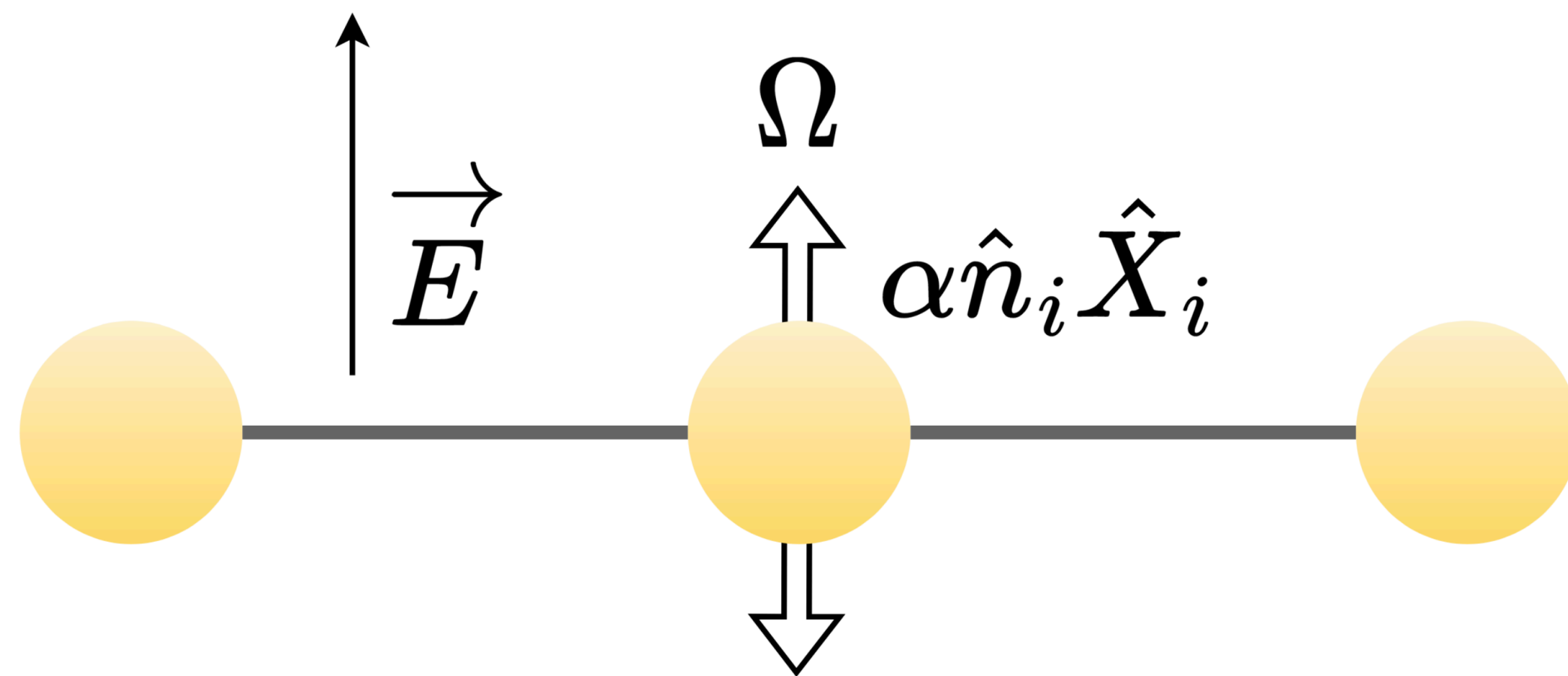
“Breathing” modes



**bond-Su-Schrieffer-Heeger
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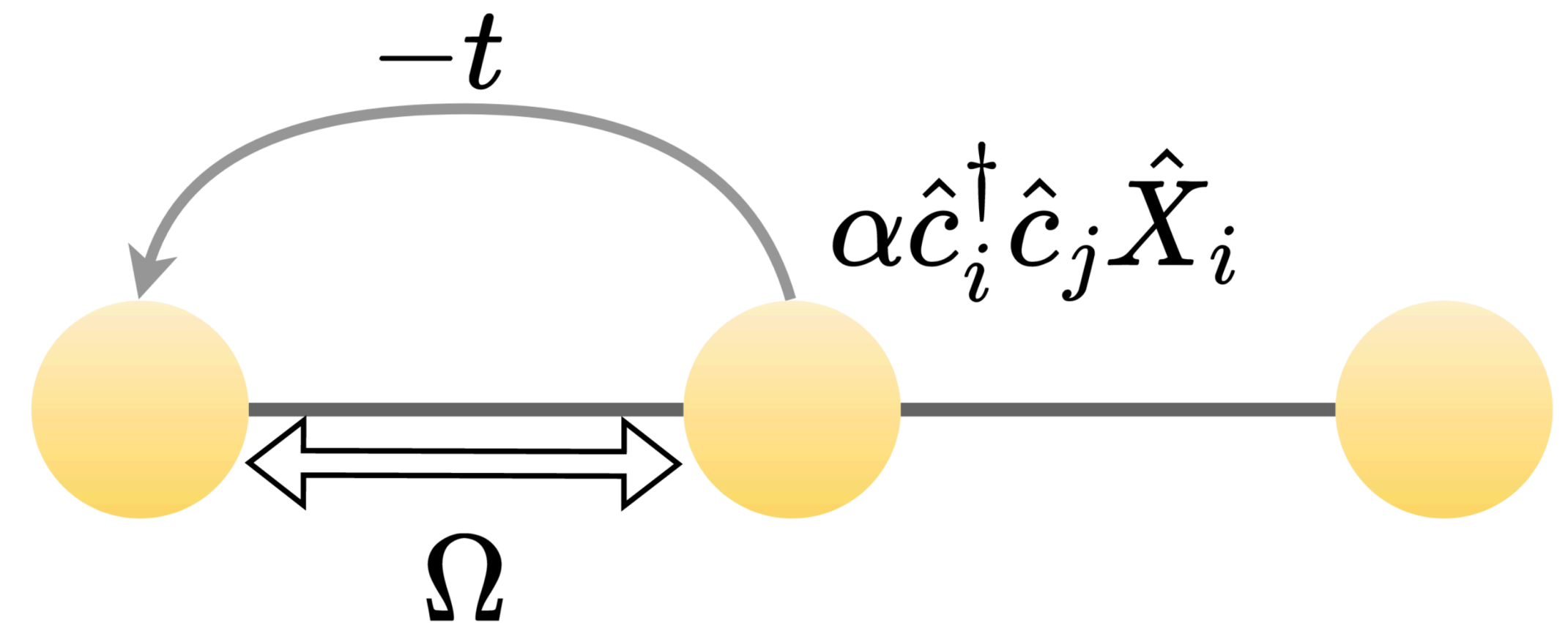
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**bond-Su-Schrieffer-Heeger
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Proposal by C. Zhang et al., Phys. Rev. X 13, 011010 (2023)

Methods

Tight-binding model and phonon modes

$$\hat{H}_{\text{el}} = -t \sum_{\langle i,j \rangle, \sigma} \left(\hat{c}_{i\sigma}^\dagger \hat{c}_{j\sigma} + \text{h.c.} \right) - \mu \sum_{i\sigma} \hat{n}_{i\sigma}$$

$$\hat{H}_{\text{lat}} = \sum_i \left(\frac{1}{2} M \Omega^2 \hat{X}_i^2 + \frac{1}{2M} \hat{P}_i^2 \right)$$

e-ph coupling

$$\hat{H}_{\text{H}} = \alpha \sum_{i\sigma} \hat{n}_{i\sigma} \hat{X}_i$$

$$\hat{H}_{\text{bSSH}} = \alpha \sum_{\langle i,j \rangle, \sigma} \hat{X}_{ij} \left(\hat{c}_{i\sigma}^\dagger \hat{c}_{j\sigma} + \text{h.c.} \right)$$

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Determinant Quantum Monte Carlo^{1,2} (DQMC) and Hybrid Monte Carlo³ (HMC)

- Auxiliary field Monte Carlo
- DQMC: $\Omega \geq t$, HMC: $\Omega \leq t$
- Finite temperature
- 8 x 8 square clusters

¹S. Johnston, E. A. Nowadnick, Y. F. Kung, B. Moritz, R. T. Scalettar, and T. P. Devereaux, Phys. Rev. B **87**, 235133 (2013)

²C. Feng, B. Xing, D. Poletti, R. Scalettar, and G. Batrouni, Phys. Rev. B **106**, L081114 (2022)

³B. Cohen-Stead, O. Bradley, C. Miles, G. Batrouni, R. Scalettar, and K. Barros, Phys. Rev. E, **105** 065302 (2022)

Electron-phonon coupling constant

- Dimensionless coupling: $\lambda = N_F \left\langle \left\langle \frac{\alpha_{\mathbf{k}-\mathbf{k}'}}{\Omega_{\mathbf{k}-\mathbf{k}'}}^2 \right\rangle \right\rangle_{FS}$

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$$\lambda_{bSSH} \approx \frac{4\alpha^2}{\Omega^2 W}$$

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Momentum averaged:

$$\begin{aligned} \lambda_{bSSH} &\approx \frac{4\alpha^2}{\Omega^2 W} \int_{-\pi}^{\pi} \int_{-\pi}^{\pi} dk dq \cos^2[(k + q/2)/2] \\ &= \frac{4\alpha^2}{\Omega^2 W} \cdot \frac{1}{2} = \frac{2\alpha^2}{\Omega^2 W} \end{aligned}$$

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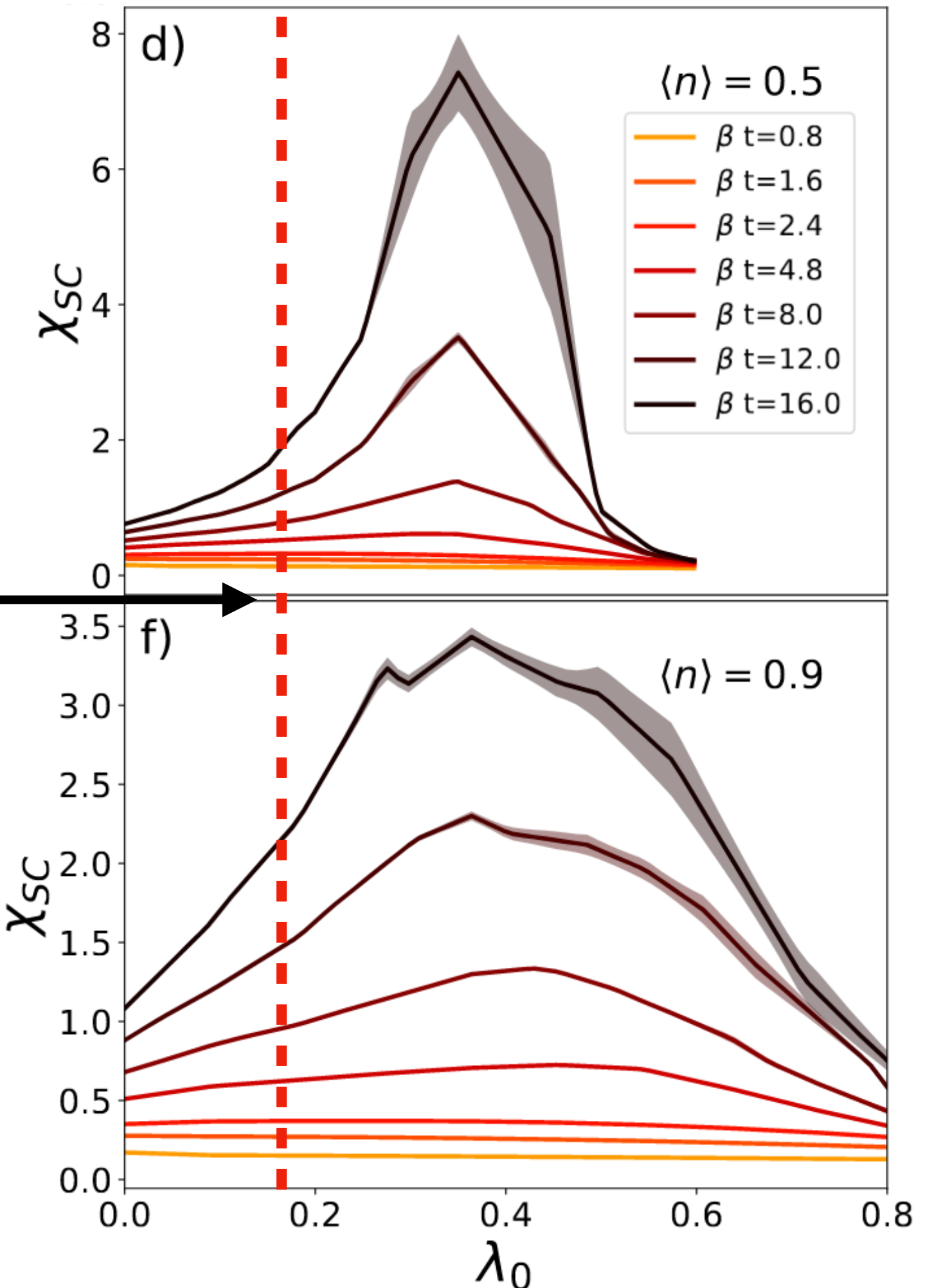
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$$\lambda = 0.15$$

Holstein $L = 8$ $\Omega = 2.0t$



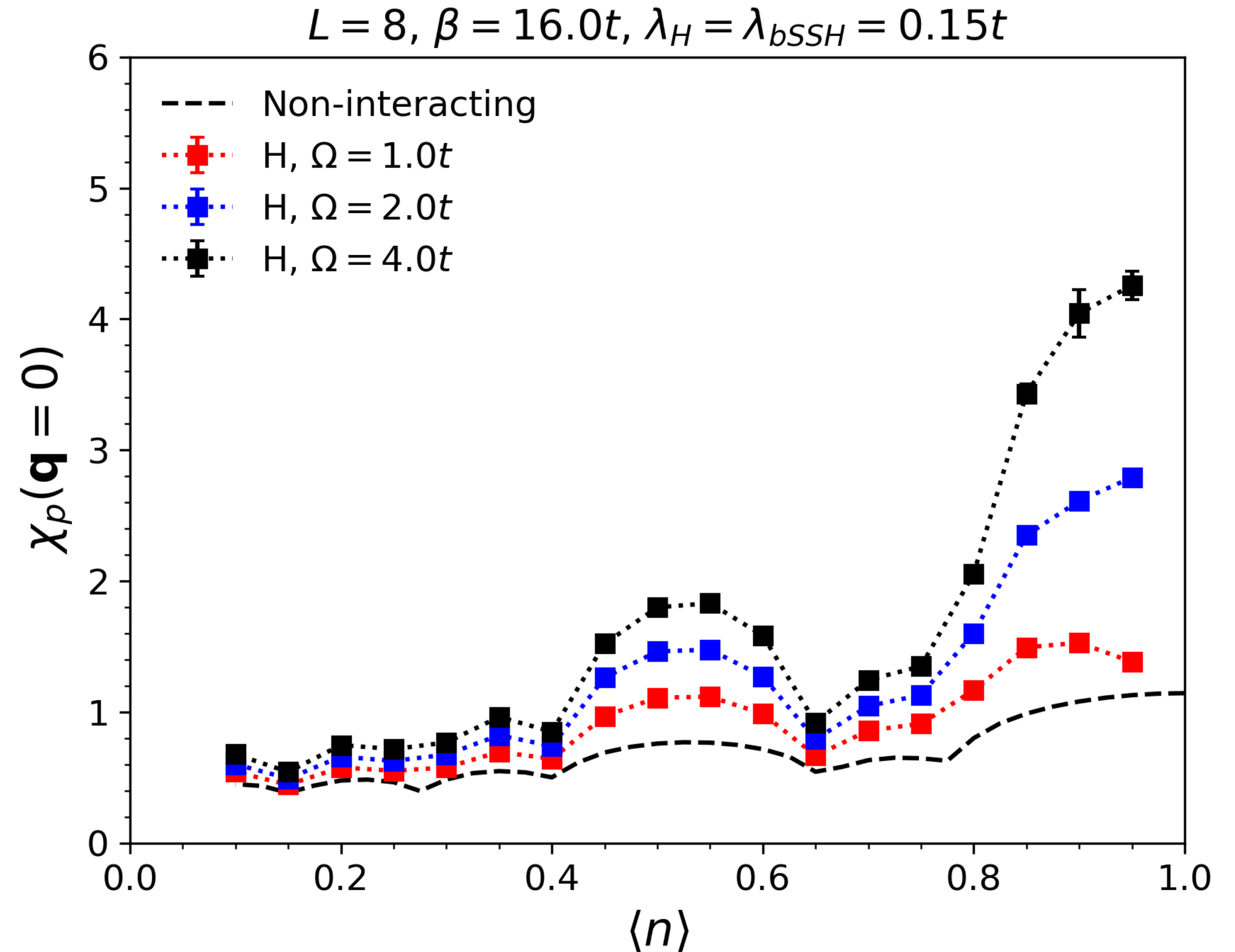
Density dependent pairing correlations

Pair-field operator

$$\hat{\Delta}(\tau) = \sum_i \hat{c}_{i\uparrow}(\tau) \hat{c}_{i\downarrow}(\tau)$$

Pair-field susceptibility

$$\chi_p = \frac{1}{N} \int_0^\beta d\tau \langle \hat{\Delta}(\tau) \hat{\Delta}(0) \rangle$$



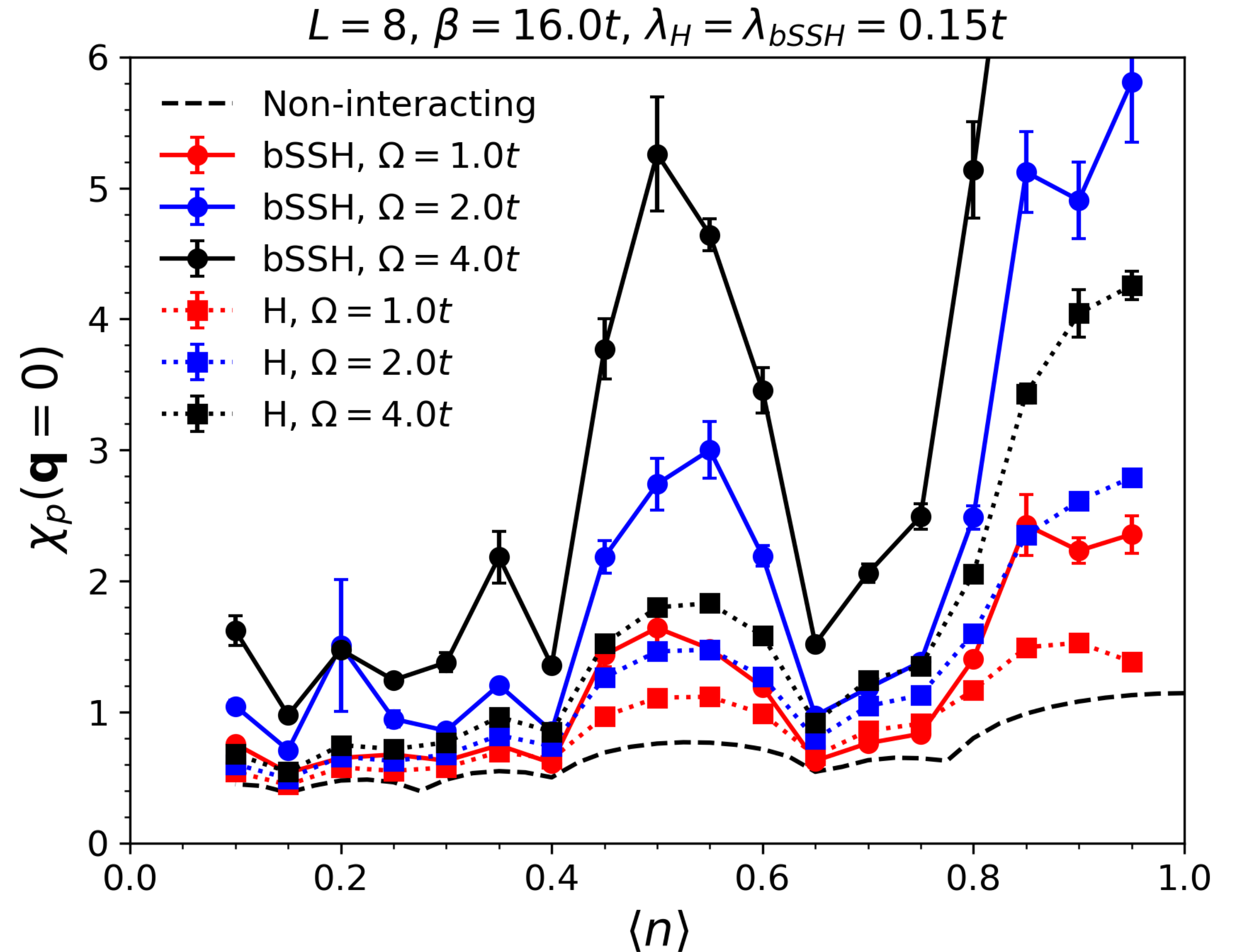
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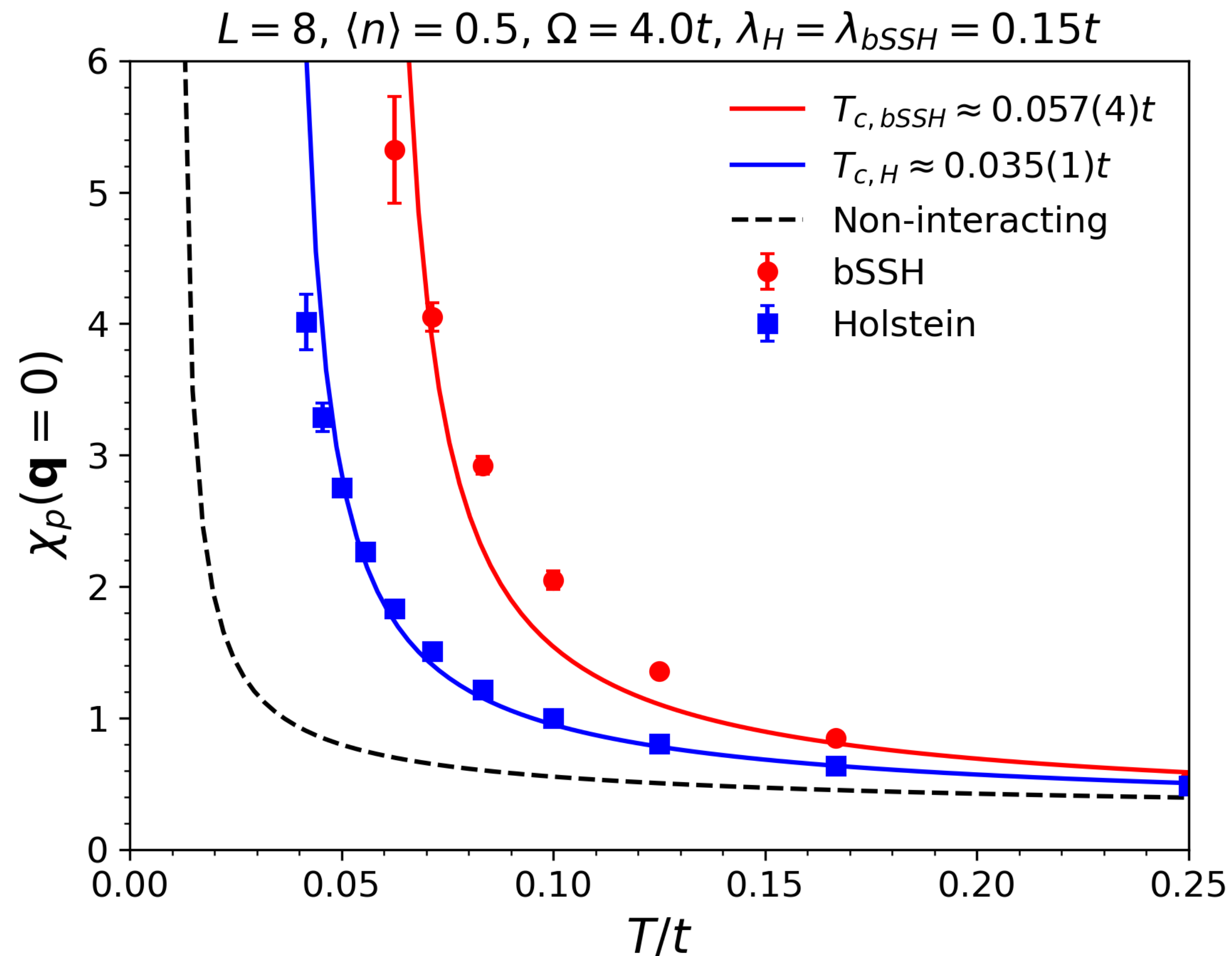
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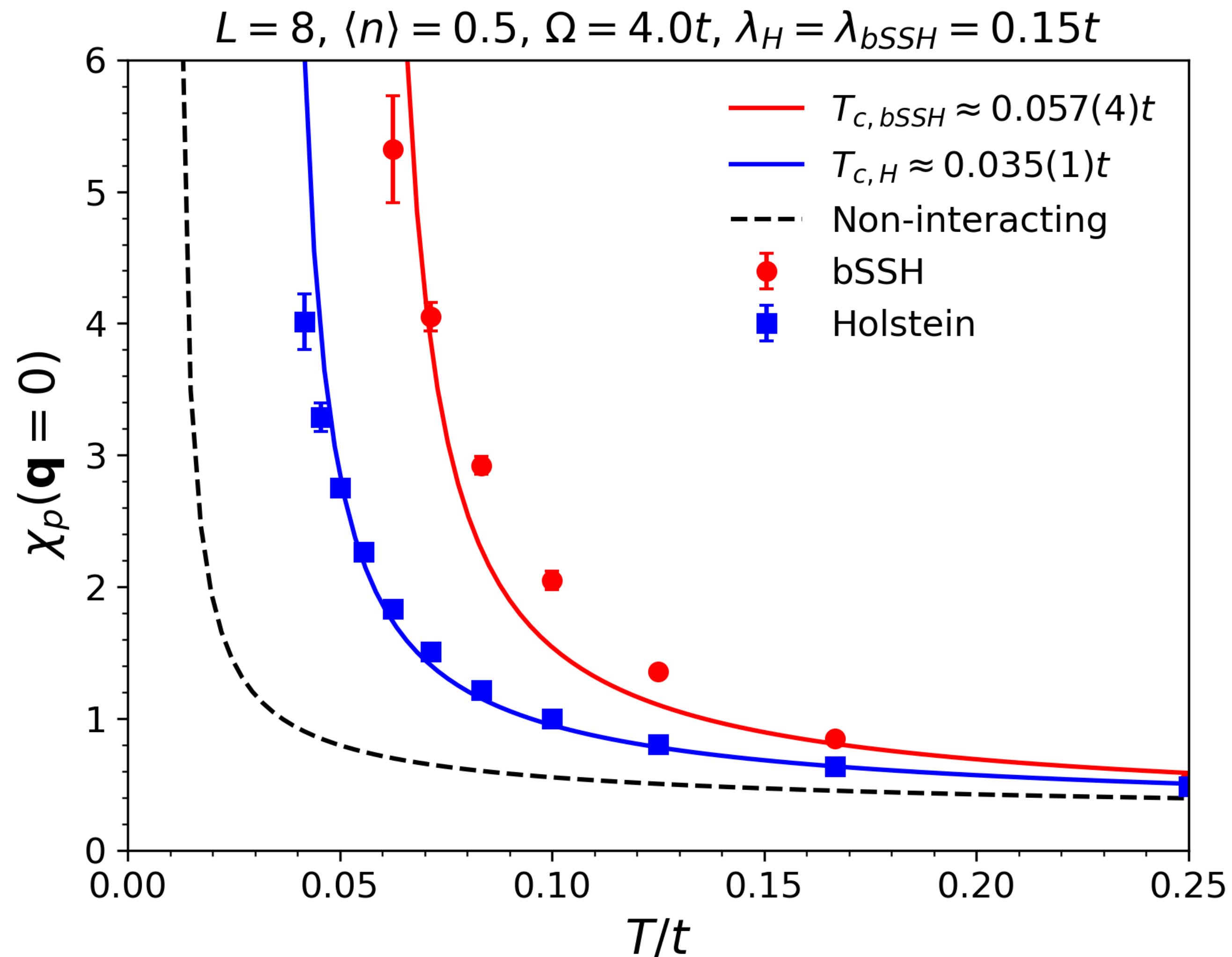
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Temperature dependent pairing correlations

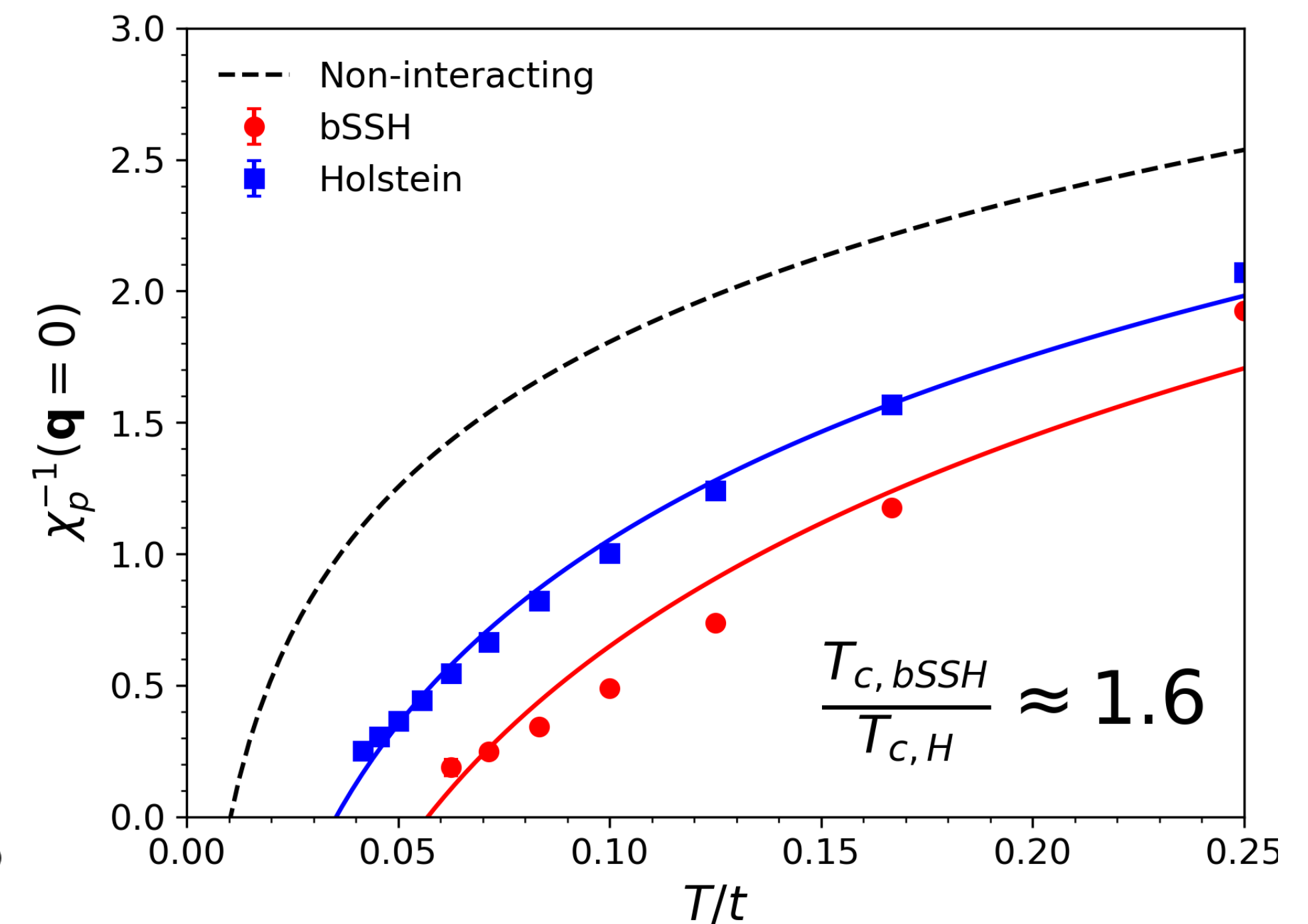


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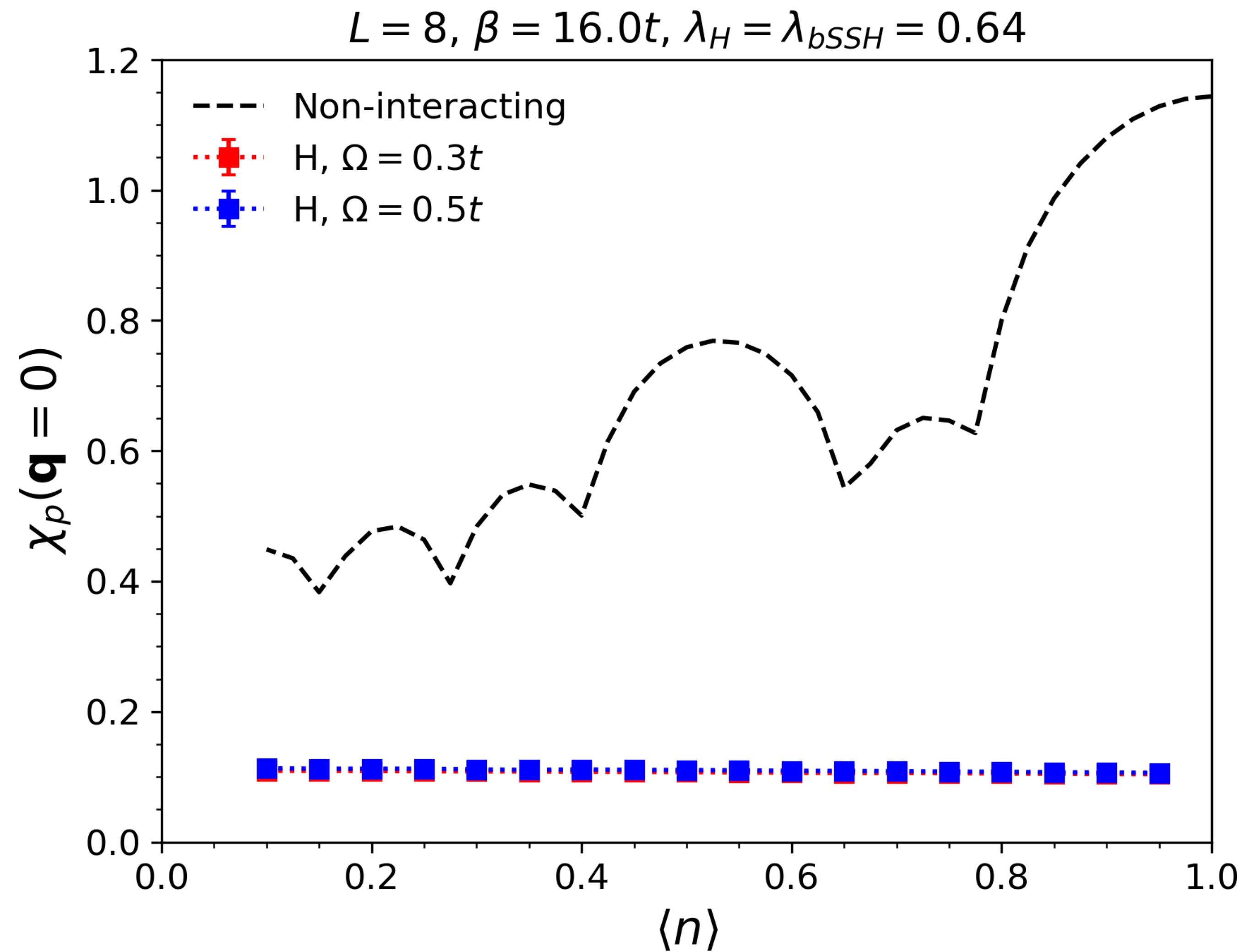


BCS instability:

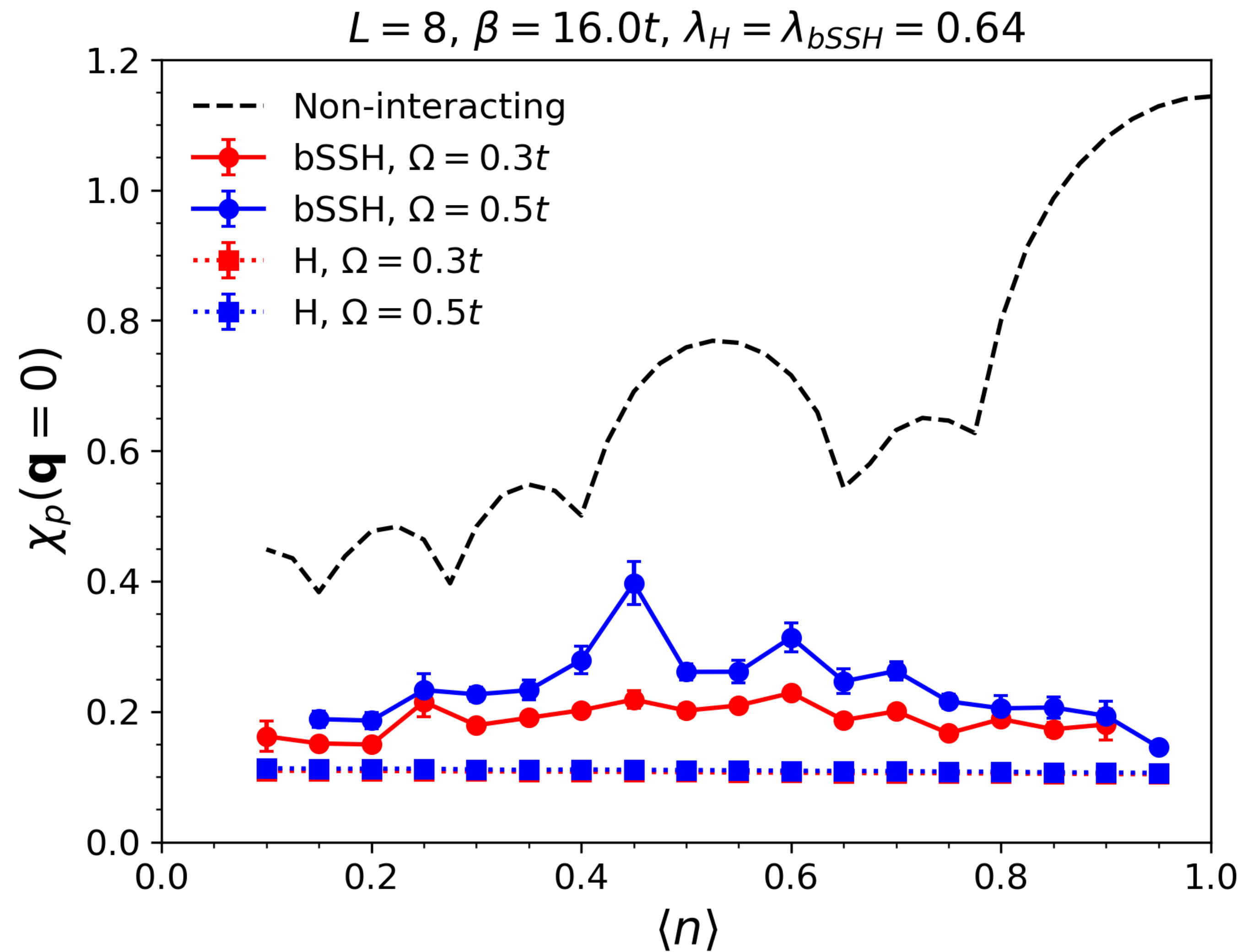
$$\chi_p^{-1} = A \log \left(\frac{T}{T_c} \right)$$



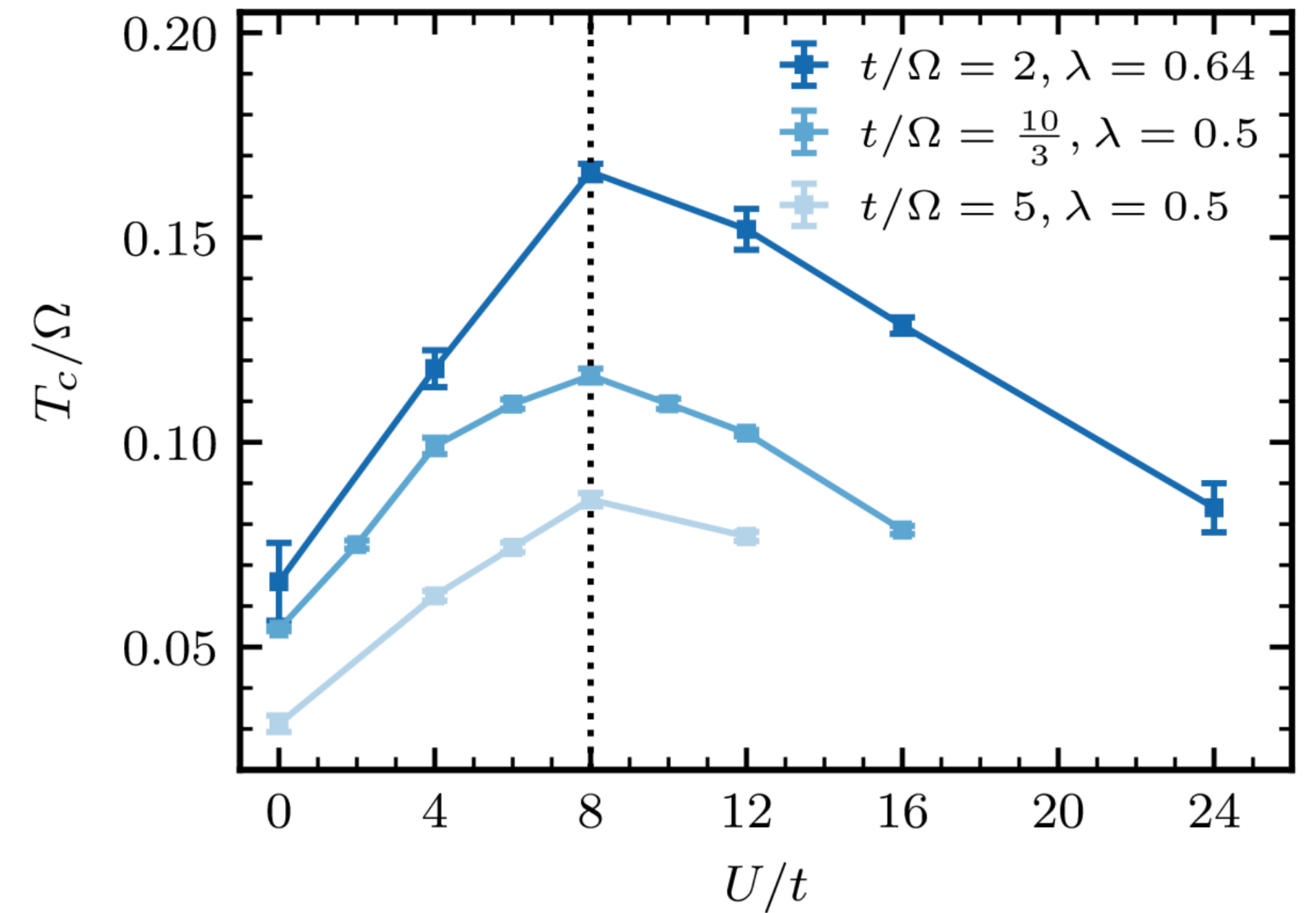
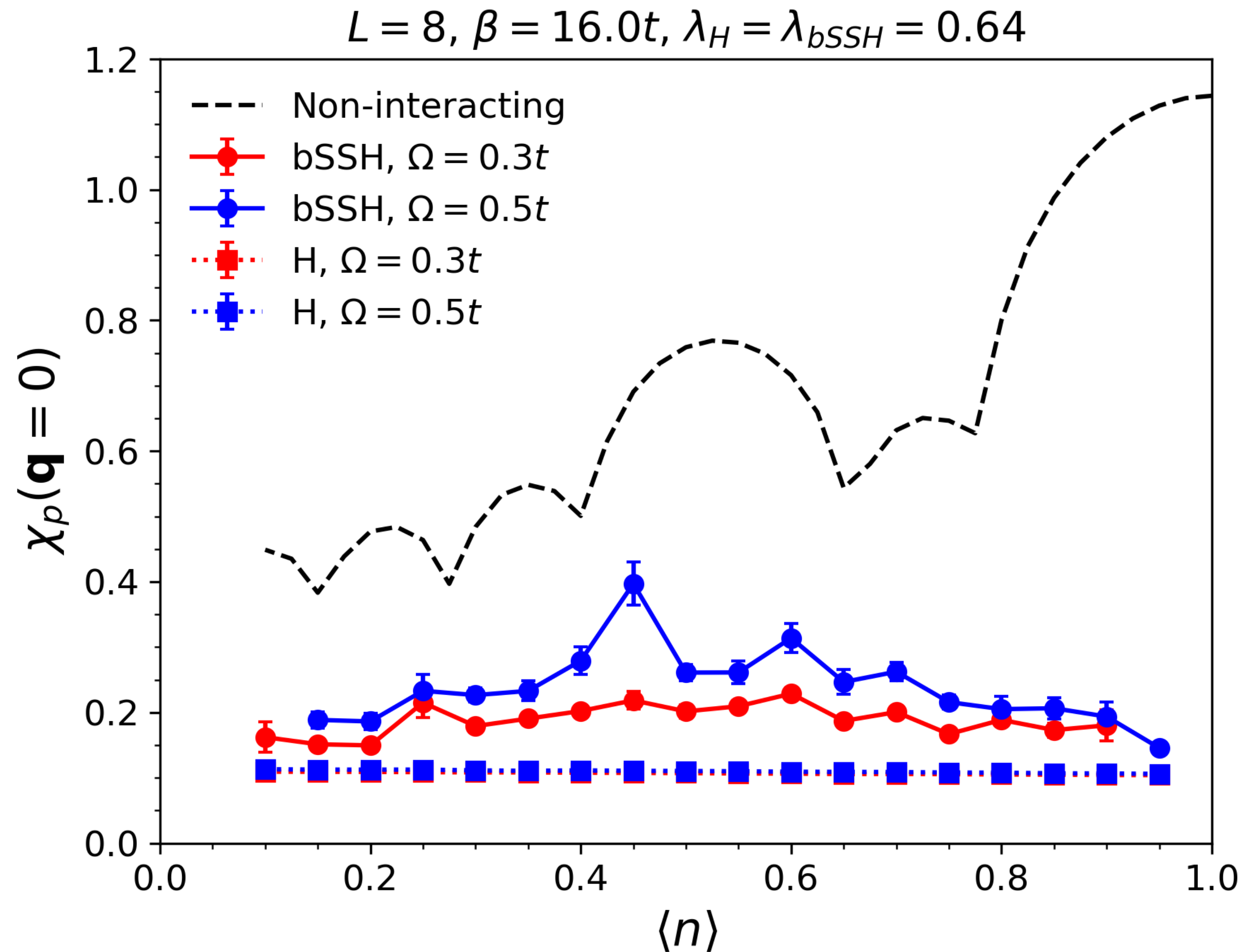
Strong coupling



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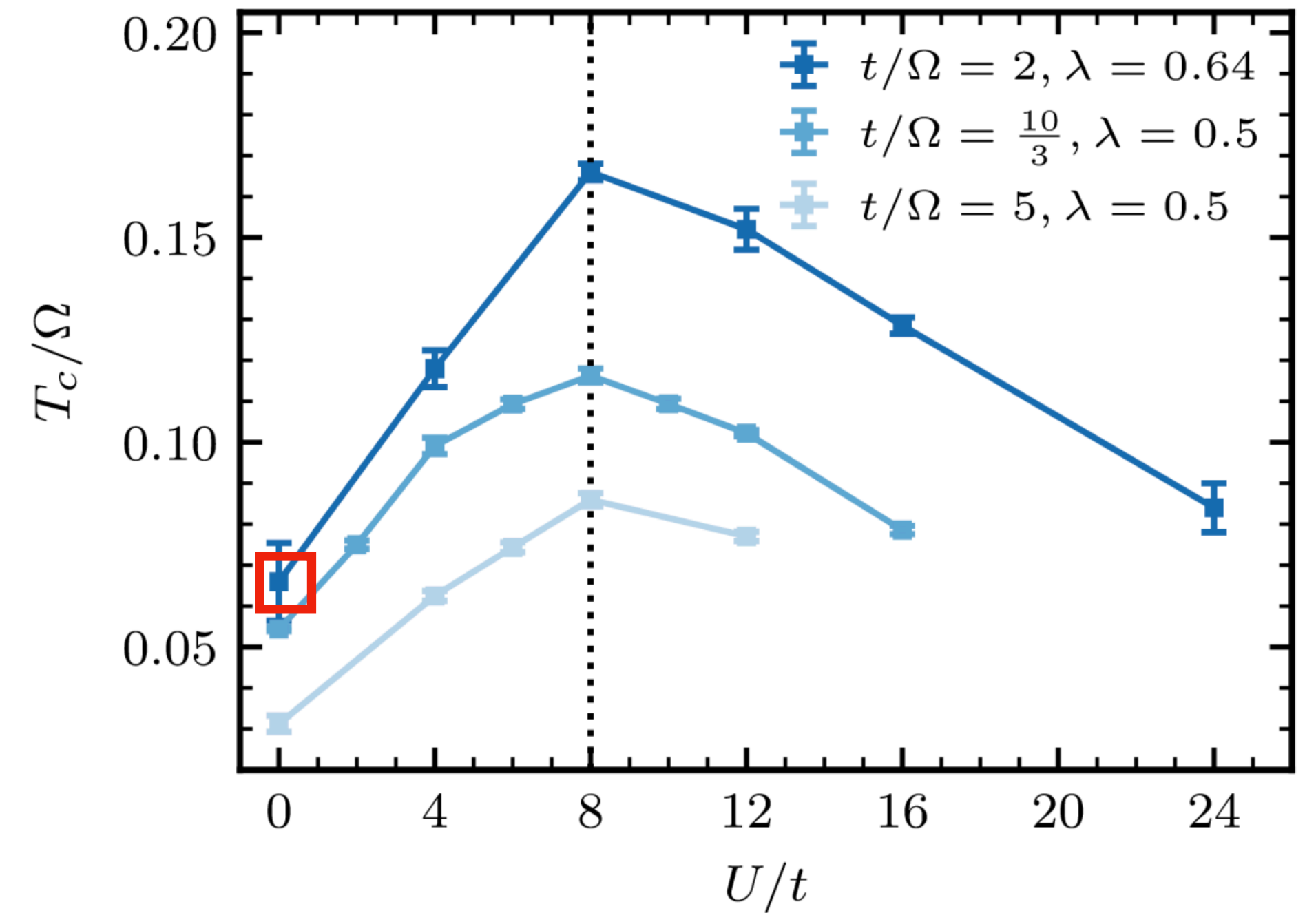
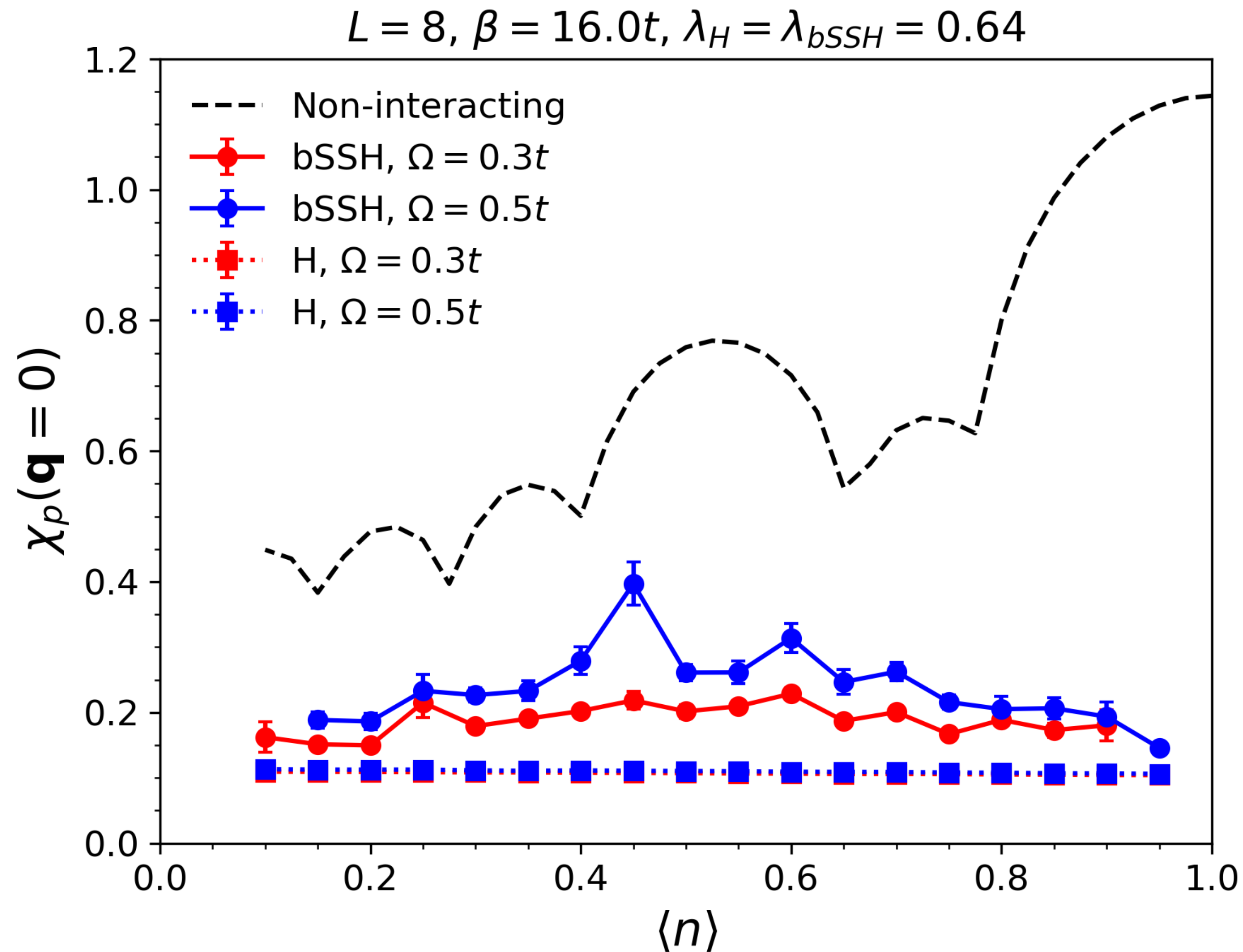


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C. Zhang et al., Phys. Rev. X 13, 011010 (2023)

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$$\beta^{-1} = 0.0625t < T_c$$

Summary

- Holstein (density coupled) versus bond-SSH models (bond modulated)
- Significant enhancements to s-wave pairing in bSSH model the anti-adiabatic limit
- Could potentially be applicable to explain “flat band” systems
- Little pairing enhancement for smaller frequencies, strong coupling

