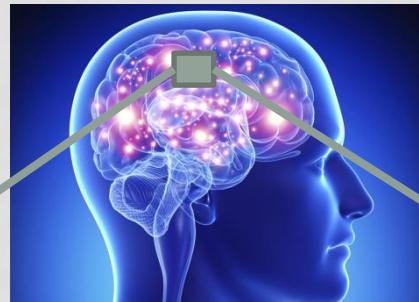




COMPUTING WITH RHYTHMS: THE SEARCH FOR OSCILLATORY DEEP NEURAL NETWORKS

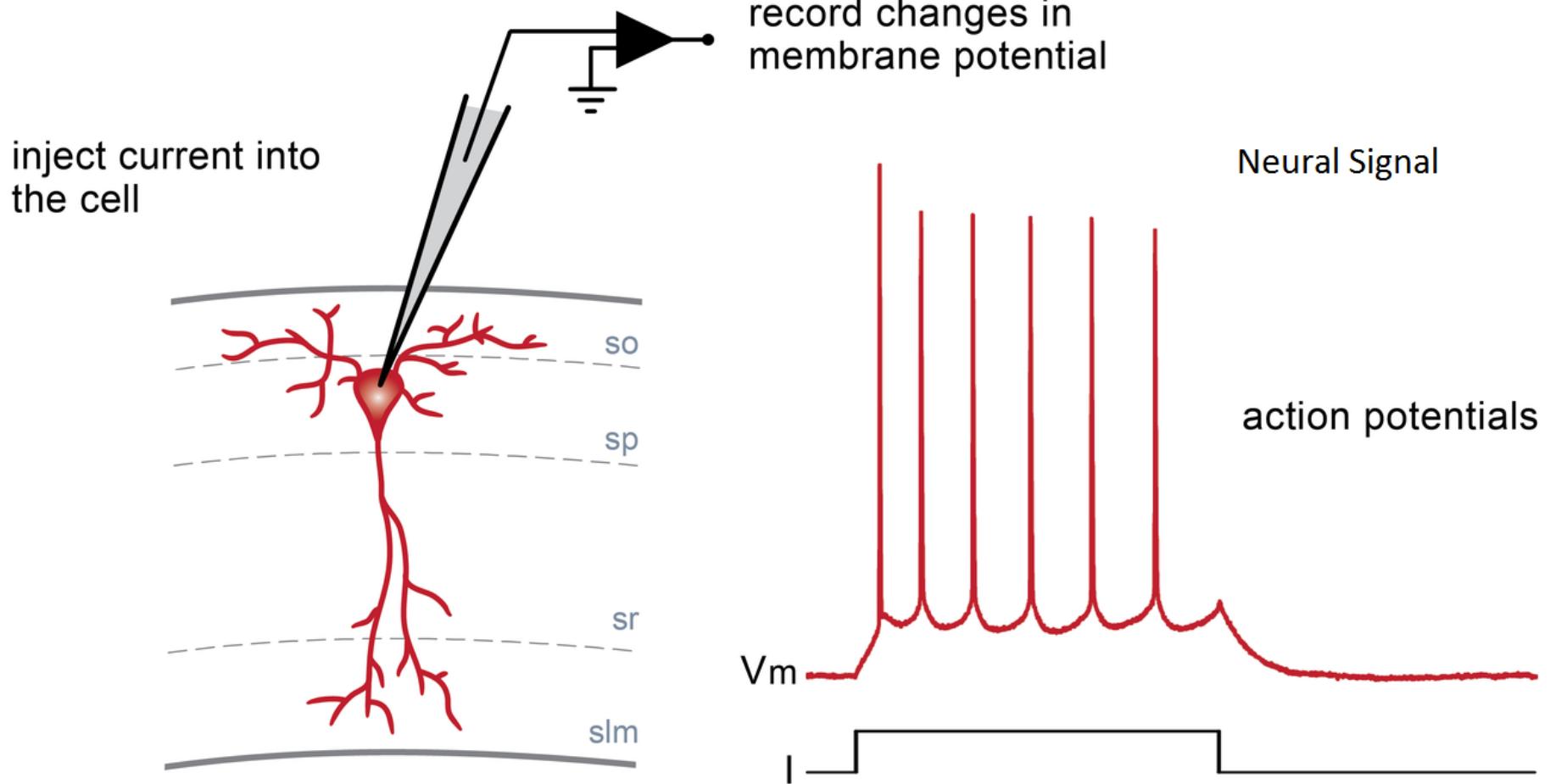
V. SRINIVASA CHAKRAVARTHY
IIT MADRAS.

PhD Student
Dipayan Biswas

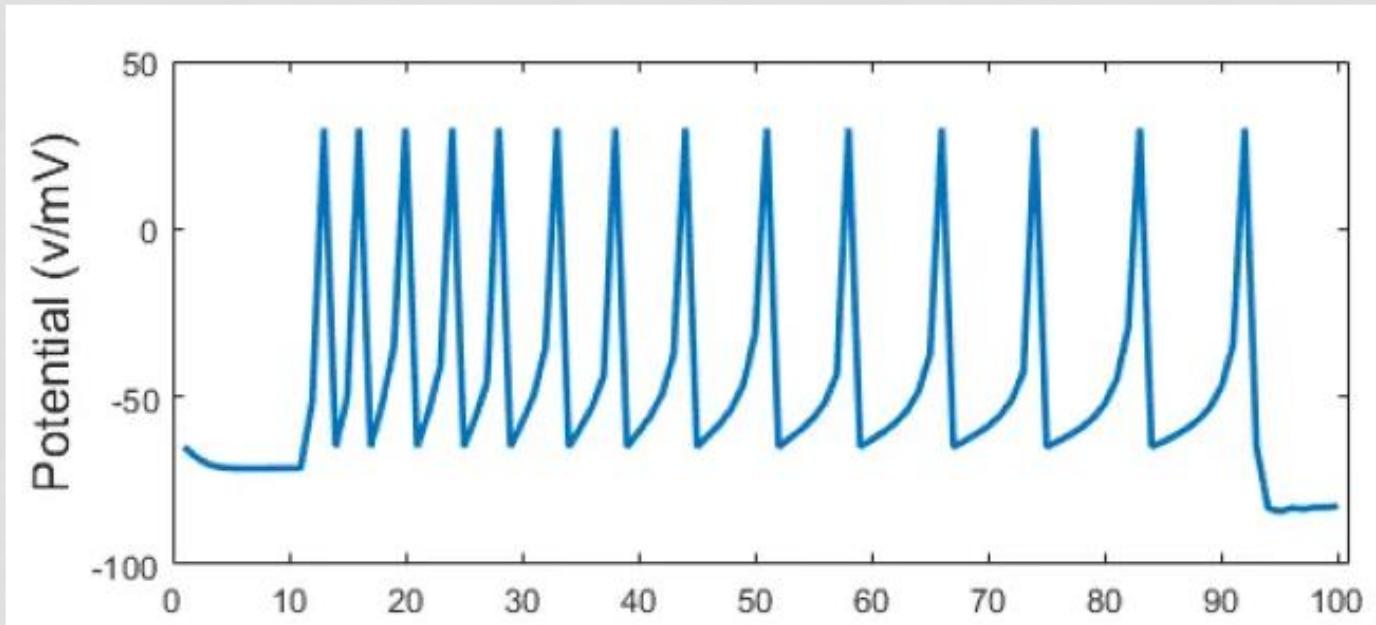


The Brain is a network of Neurons

Neurons communicate with each other by exchanging
electrical signals



WHAT IS THE NEURAL CODE?



How do you mathematically describe the information present in the electrical spiking activity of a neuron?

TWO THEORIES

Spike code

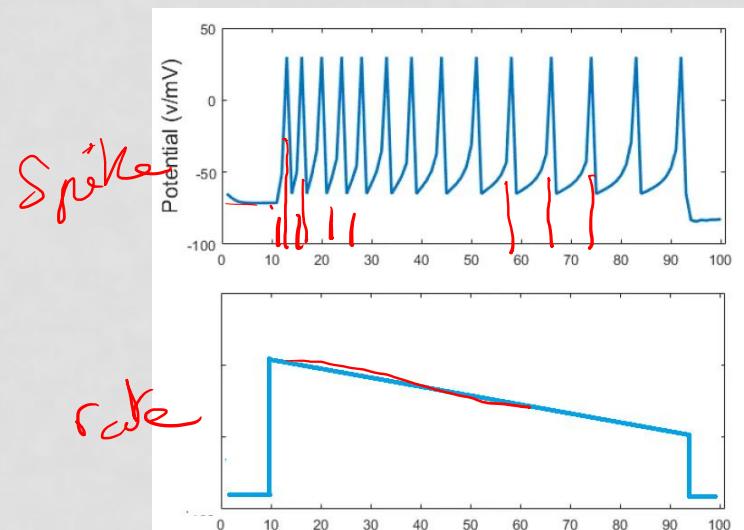
- Information is present in the times of occurrence of the spikes

$$s(t) = \sum_i \delta(t - t_i)$$



Rate Code

- Information is present in the frequency (rate) of the spikes



TWO KINDS OF NEURAL NETWORK MODELS

Spiking neuron networks

- Used often to model brain function

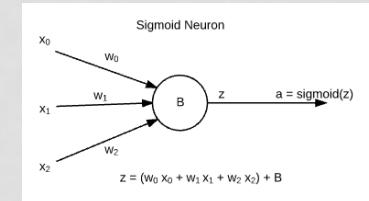
Rate coded neural networks

- Used more in artificial, engineering domains
- Eg. Deep neural networks

RATE-CODED NEURON MODELS

Used in
Deep
Neural
Networks

Sigmoidal neuron - static, rate-coded



$$\begin{aligned} y &= x, & x > 0 \\ &= 0, & \text{otherwise} \end{aligned}$$

Relu neuron - static, rate-coded

Long Short-Term Memory block:

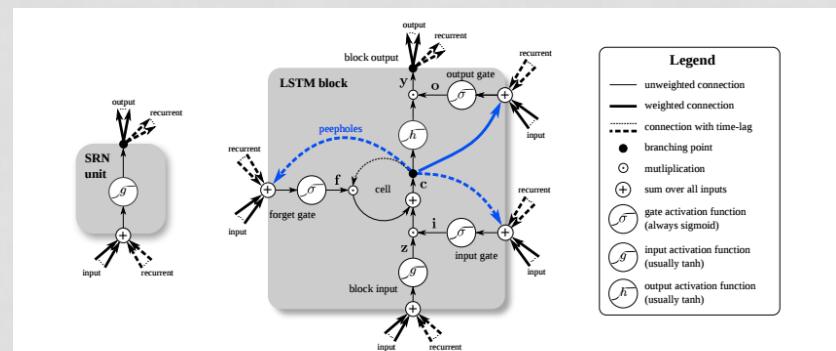


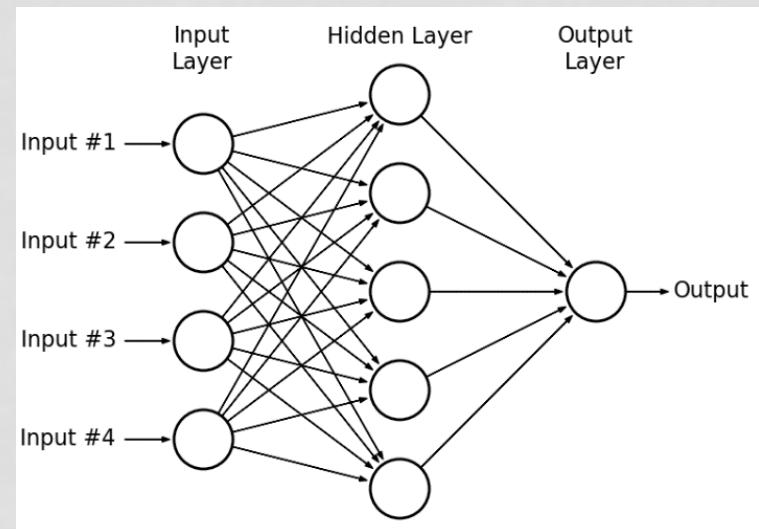
Figure 1. Detailed schematic of the Simple Recurrent Network (SRN) unit (left) and a Long Short-Term Memory block (right) as used in the hidden layers of a recurrent neural network.

DEEP LEARNING/AI

- The aim of Deep learning/AI is to reproduce human intelligence.
- The approach is to achieve it through an implementation of the human brain
- Deep Neural networks are offered as an answer

NNS AND DEEP NNS

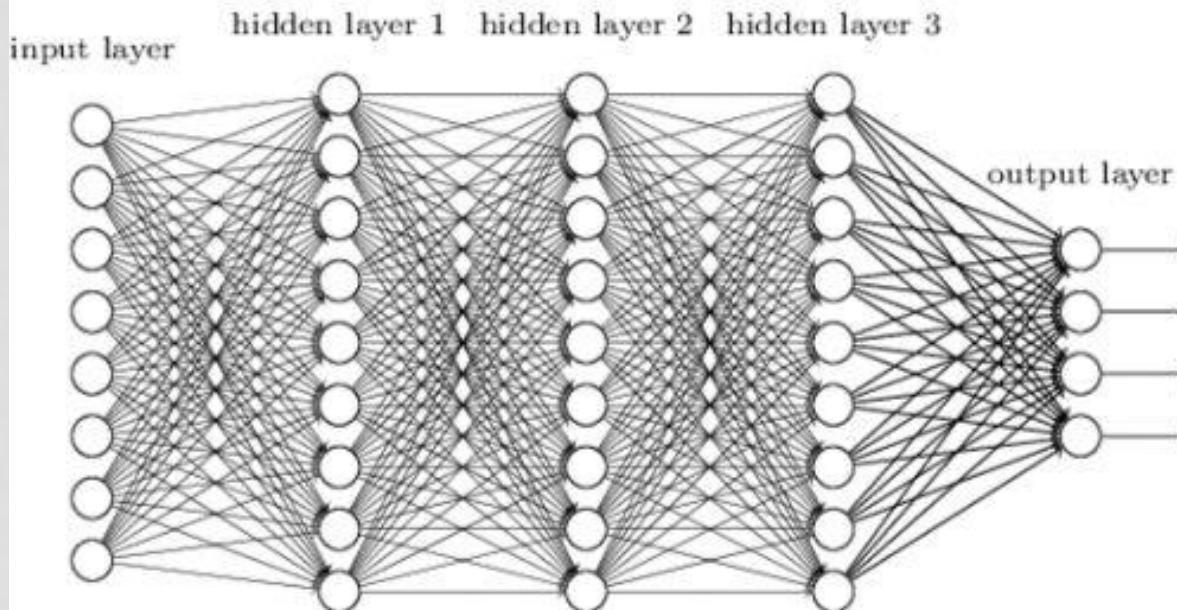
- 1980s-1990s: Multilayer perceptrons and backpropagation
- Deep learning – 2006
- Powerful computational properties.
- Universal approximation results



Multi-layer Perceptron

DEEP NEURAL NETWORK

Deep neural network

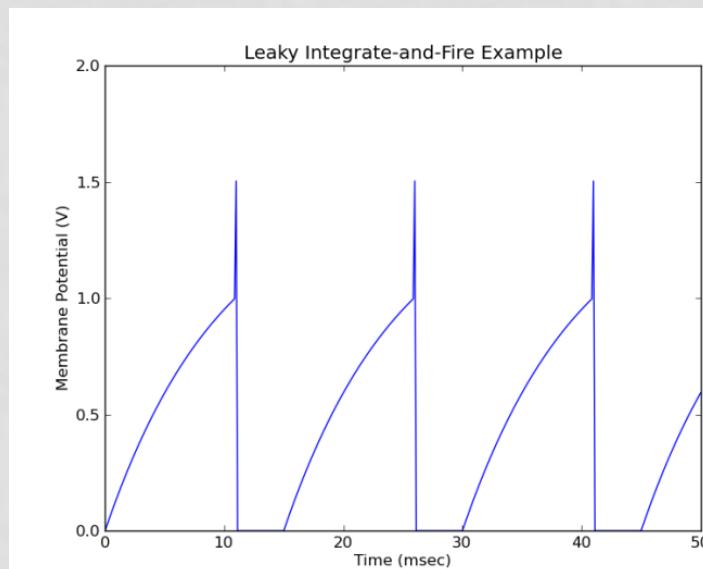
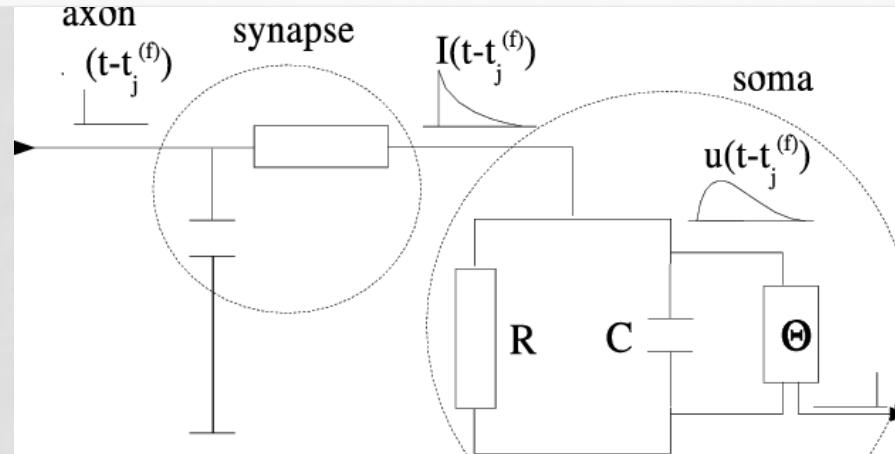


Has Universal Approximation properties
Can learn to map arbitrary input/output
Excellent applications in image pattern recognition,
Signal processing, robotics...

SPIKING NEURON MODELS

- Leaky Integrate and fire neuron (1 variable)
- Izhikevich neuron (2 variables)
- Hodgkin-Huxley neuron model (4 variables)
- Biophysical neuron models (large number of variables)

LEAKY INTEGRATE AND FIRE NEURON MODEL

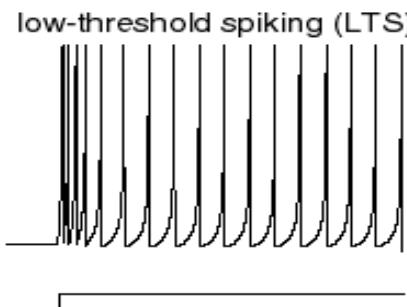
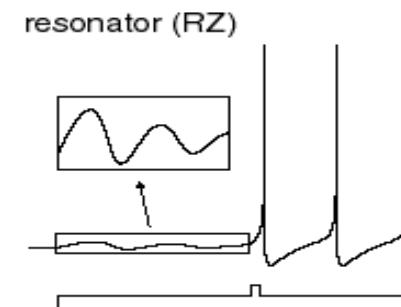
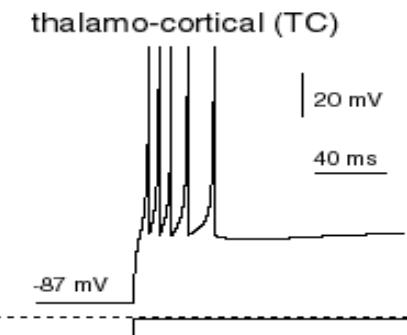
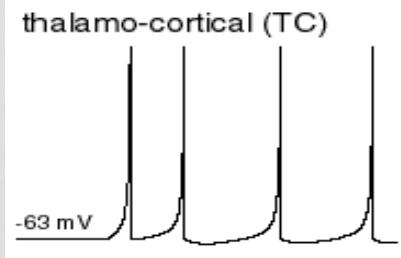
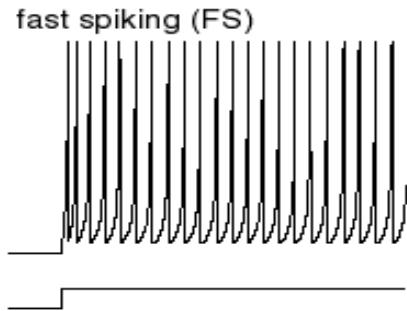
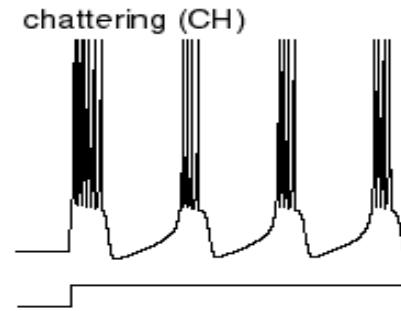
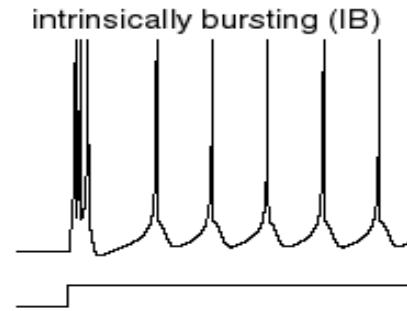
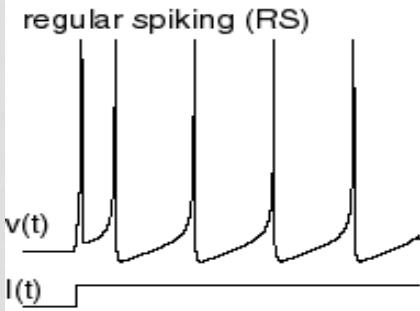
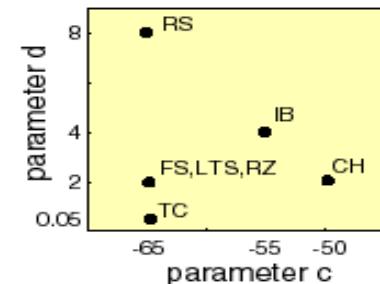
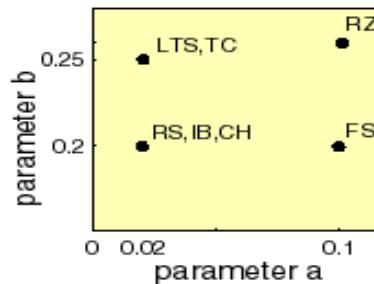
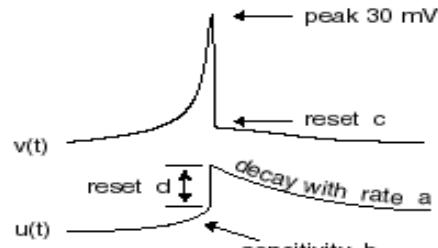


IZHIKEVICH NEURON MODEL

$$v' = 0.04v^2 + 5v + 140 - u + I$$

$$u' = a(bv - u)$$

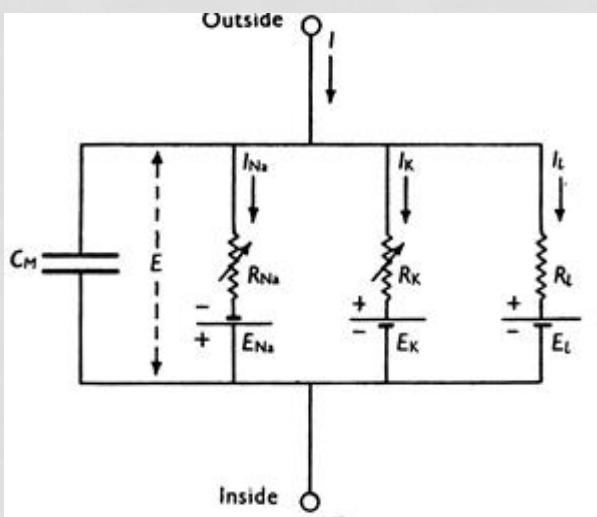
if $v = 30 \text{ mV}$,
then $v \leftarrow c$, $u \leftarrow u + d$



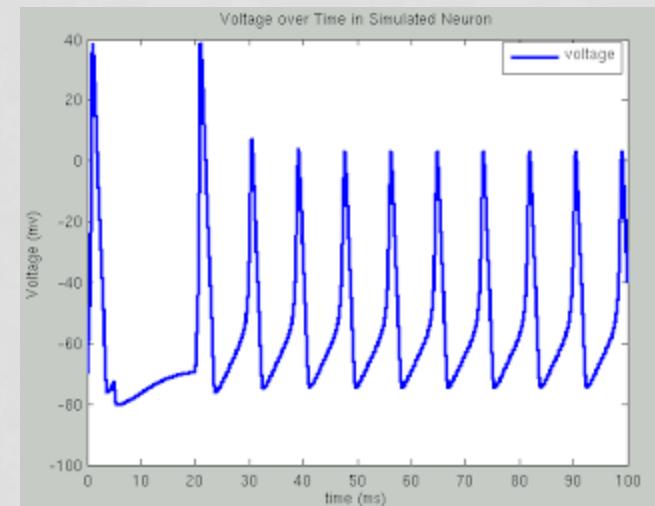
HODGKIN-HUXLEY NEURON MODEL

Hodgkin-Huxley Model (Example of a “point neuron” model)

$$\begin{aligned}\frac{dv}{dt} &= \frac{1}{C_m}[I - g_{Na}m^3h(v - E_{Na}) - g_Kn^4(v - E_K) \\ &\quad - g_L(v - E_L)] \\ \frac{dm}{dt} &= \alpha_m(v)(1 - m) - \beta_m(v)m \\ \frac{dn}{dt} &= \alpha_m(v)(1 - m) - \beta_m(v)m \\ \frac{dh}{dt} &= \alpha_h(v)(1 - h) - \beta_h(v)h\end{aligned}\tag{1}$$

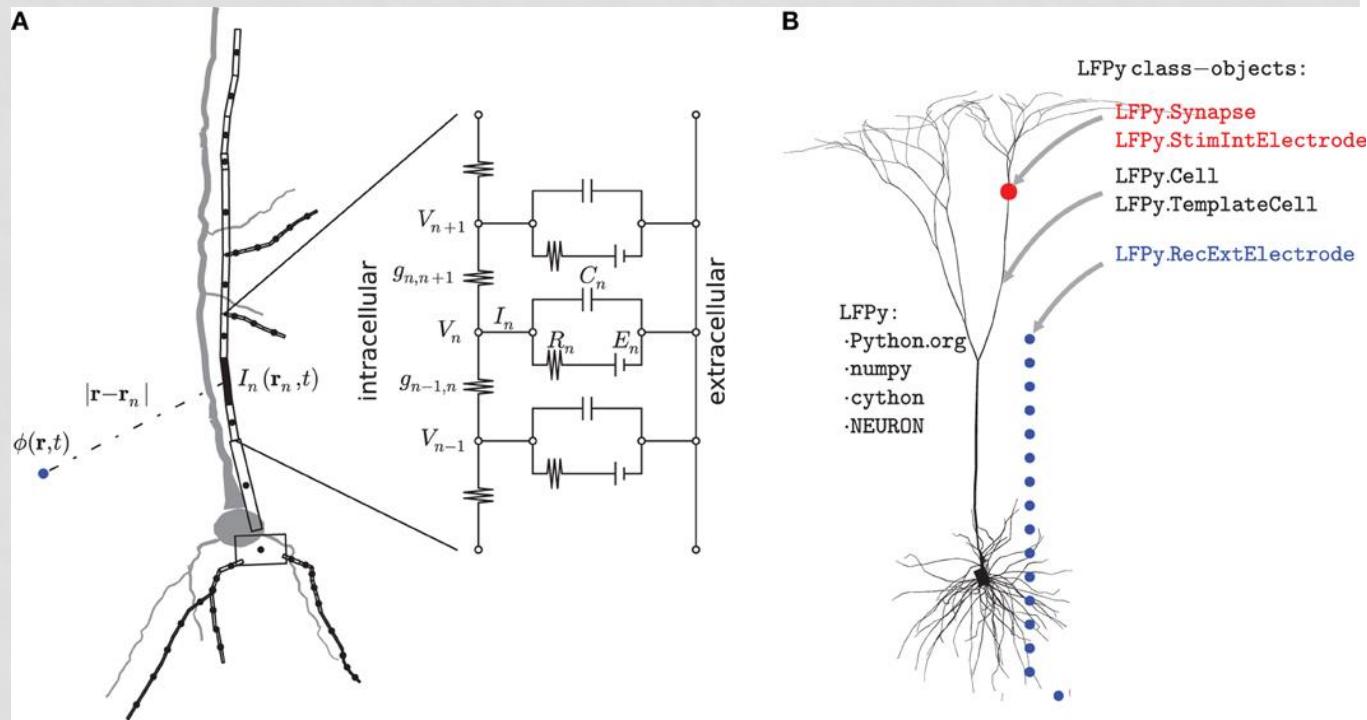


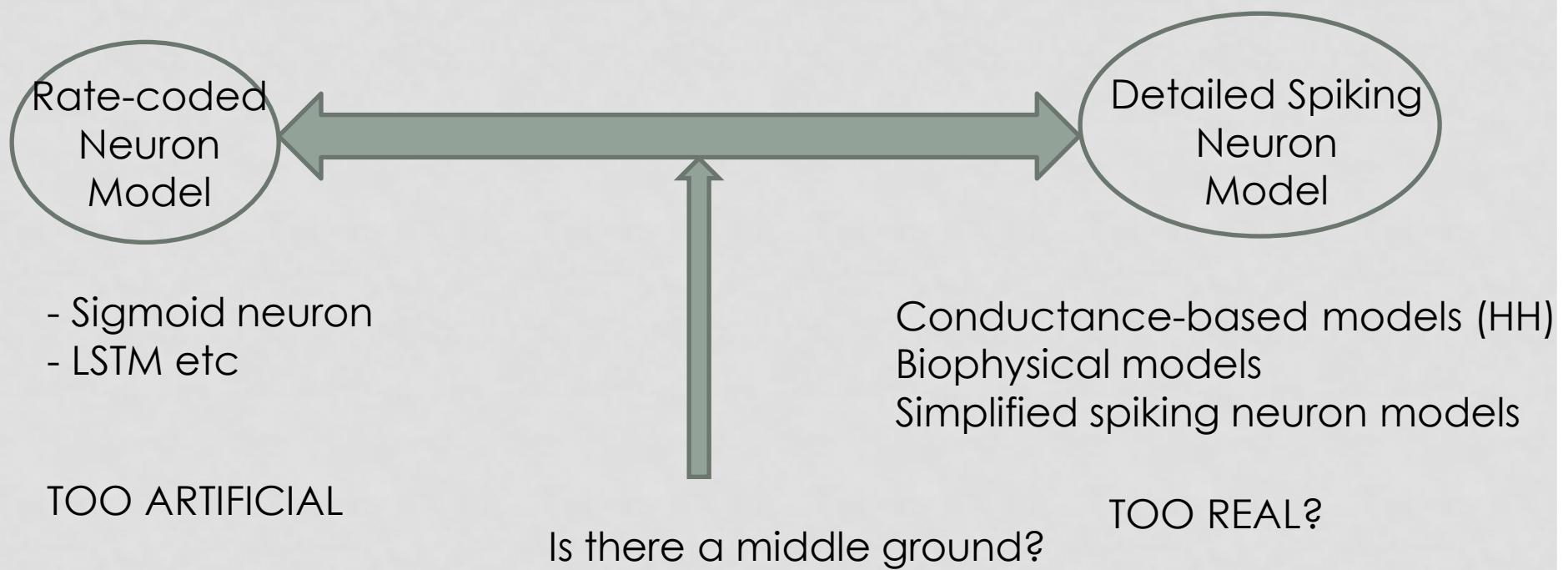
Hodgkin & Huxley
Were awarded the
Nobel Prize in 1963



BIOPHYSICAL NEURON MODEL

Simulate the entire neuronal arbor

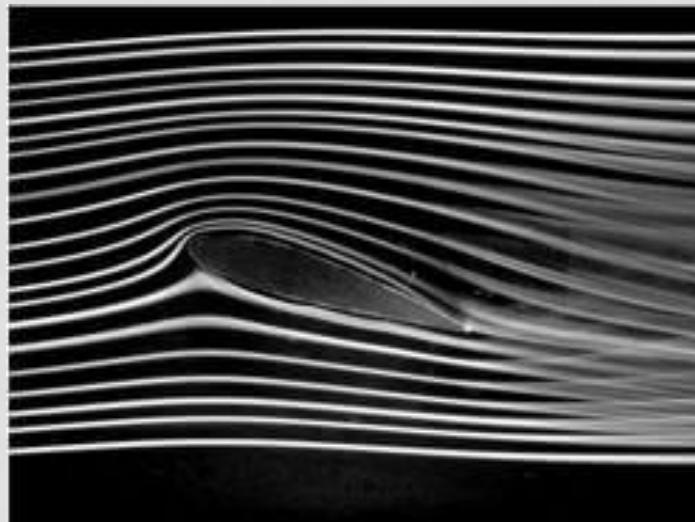




THE QUESTION OF THE RIGHT LEVEL

Example:

Aircraft wing in a wind tunnel

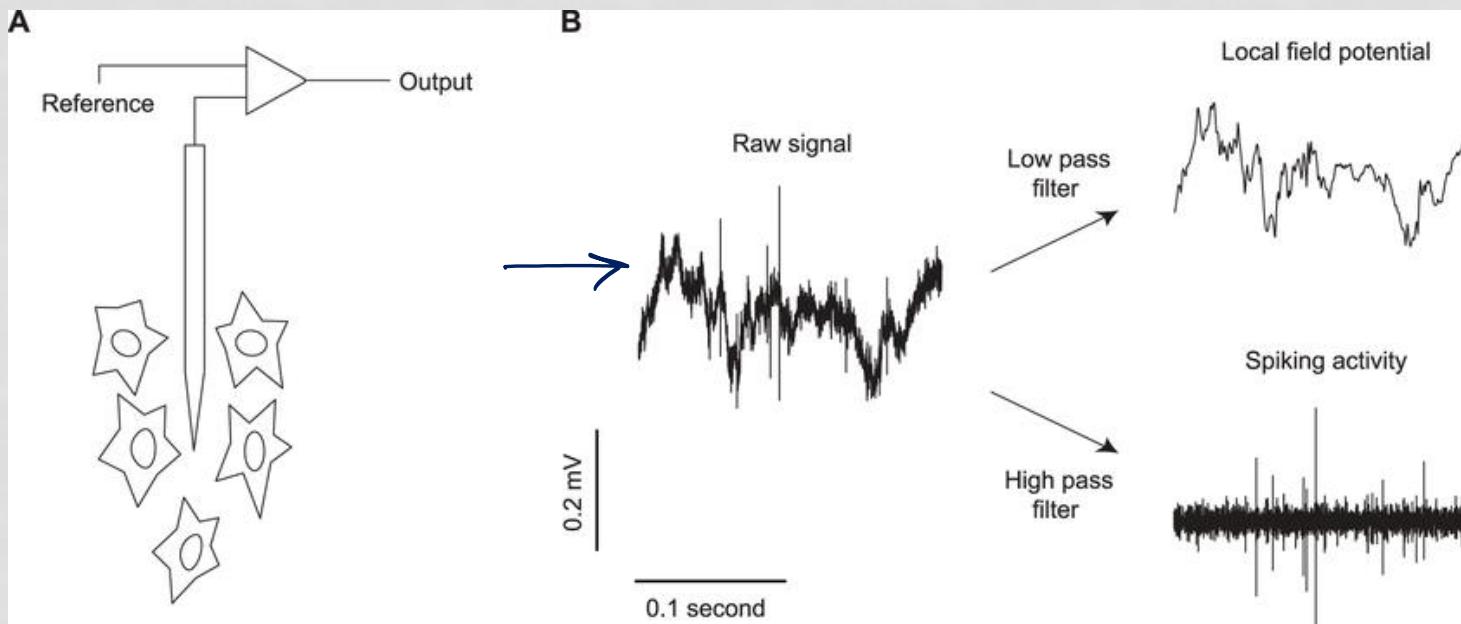


Study fluid flow using Navier Stokes Equation
FLUID is an abstraction

FLUID is NOT molecules

FROM SINGLE NEURON TO NEURAL ENSEMBLE

- Local field potential
 - activity of a local neural population
- Neural field models:
 - Coarse-grained models of spatio-temporal evolution of neural tissue (eg cortex)



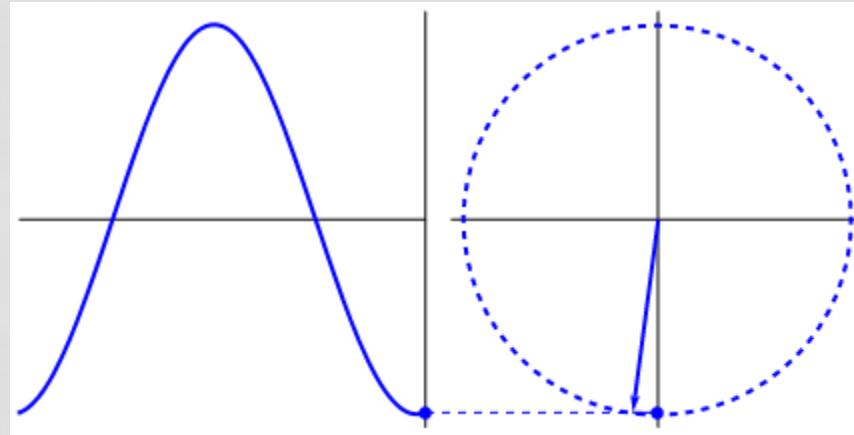
NEURAL "WAVES"

Electro Encephalo Gram (EEG)

Wave	Frequency (hertz)	Properties	EEG Trace
Delta	0.2-3	Lowest in frequency and arising in deep, dreamless sleep	
Theta	3-8	Important in strengthening synapses during learning	
Alpha	8-12	Predominant in a relaxed state when the brain is at rest and the eyes are closed	
Beta	12-30	Common when the brain is in an alert state, attentive and concentrating	
Gamma	30-120	Associated with information processing in the cerebral cortex, thinking and learning	

DO MODELS OF NEURAL OSCILLATORS OCCUPY THAT IDEAL MIDDLE GROUND?

“Synchronous activity of oscillating networks is now viewed as the critical ‘middle ground’ linking single-neuron activity to behavior” - (Buzsaki and Draghun, Science 2004)

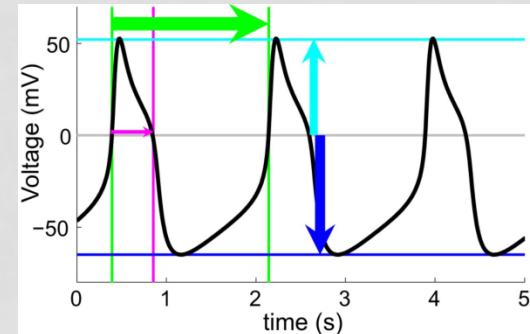


Ex
single ion channel
SINUSOID
Population of I.C.s

NEURAL OSCILLATOR MODELS (TYPICALLY 2 VARIABLE MODELS)

Wilson-Cowan model

$$\begin{aligned}\dot{x} &= -\alpha x + f(ax - by + \rho_x) \\ \dot{y} &= -\beta y + f(cx - dy + \rho_y).\end{aligned}$$

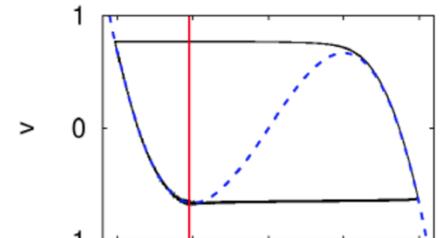


FitzHugh - Nagumo model

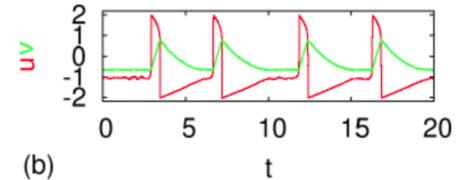
$$\frac{\partial V}{\partial t} = \frac{\partial^2 V}{\partial x^2} + (a - V)(V - 1)V - v \quad (1)$$

$$\frac{\partial v}{\partial t} = \epsilon(\beta V - \gamma v - \delta) \quad (2)$$

V = Voltage (fast variable)
 v = v -gate (slow variable)

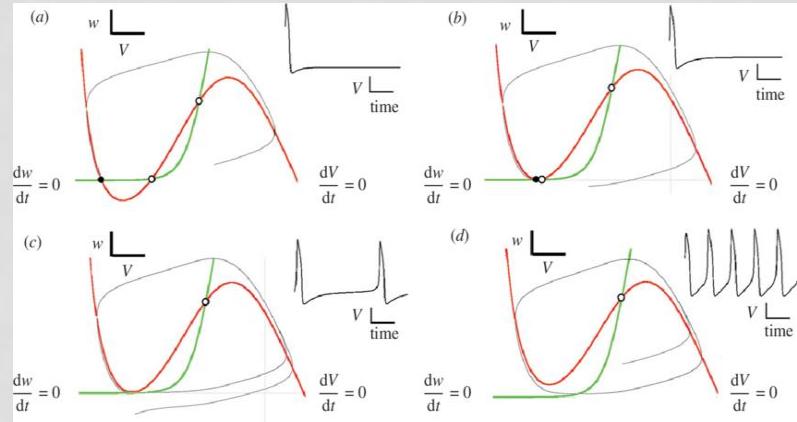


(a)



(b)

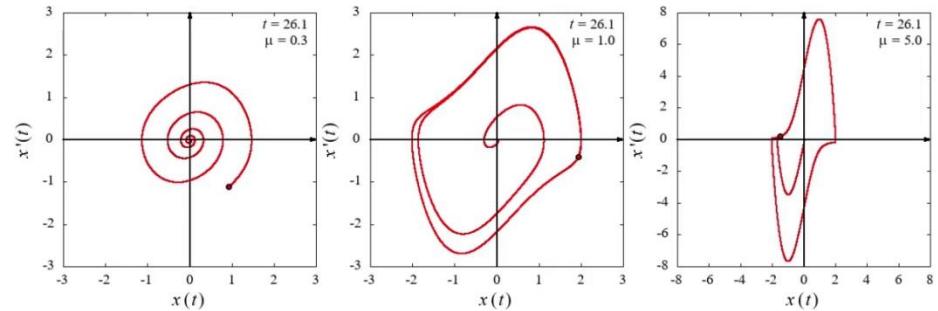
Morris-Lecarr model



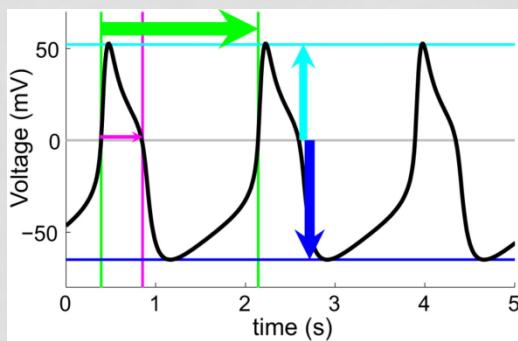
Van Der Pol model

Van der Pol equation

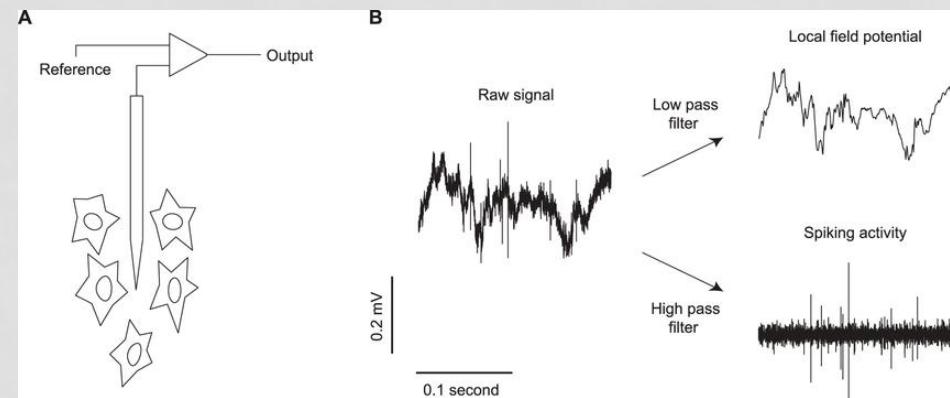
$$\frac{d^2x}{dt^2} - \mu(1-x^2) \frac{dx}{dt} + x = 0$$



Oscillation as a model of an Action Potential



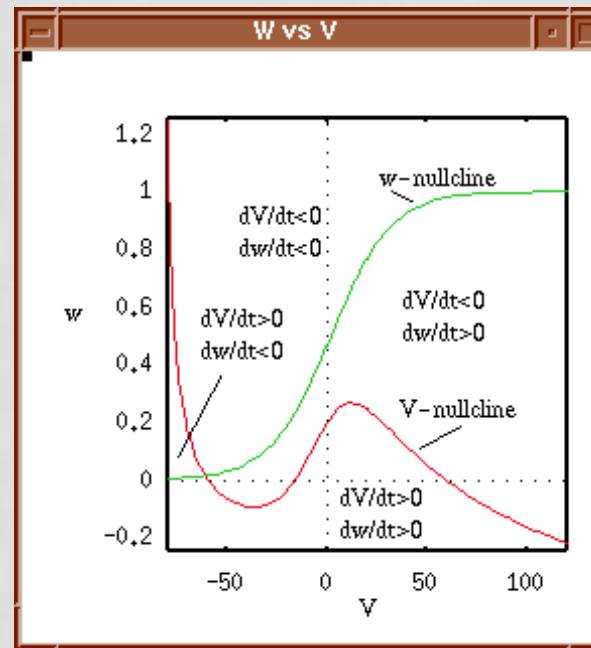
Oscillation as a model of Local Field Potential



The same models have been used for both

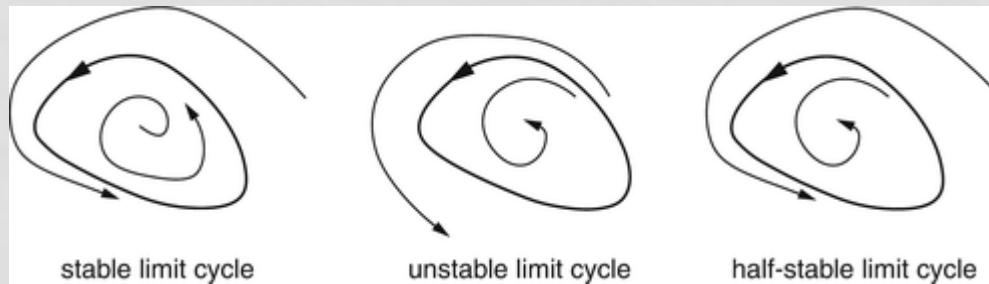
PHASE PLANE ANALYSIS

Two-variable form:
 $dx/dt = f(x,y)$
 $dy/dt = g(x,y)$



Can exhibit limit cycle oscillations

LIMIT CYCLES



Definition: A *limit cycle* is a closed trajectory in phase space having the property that at least one other trajectory spirals into it either as time approaches infinity or as time approaches negative infinity.

Can only be exhibited by nonlinear systems.

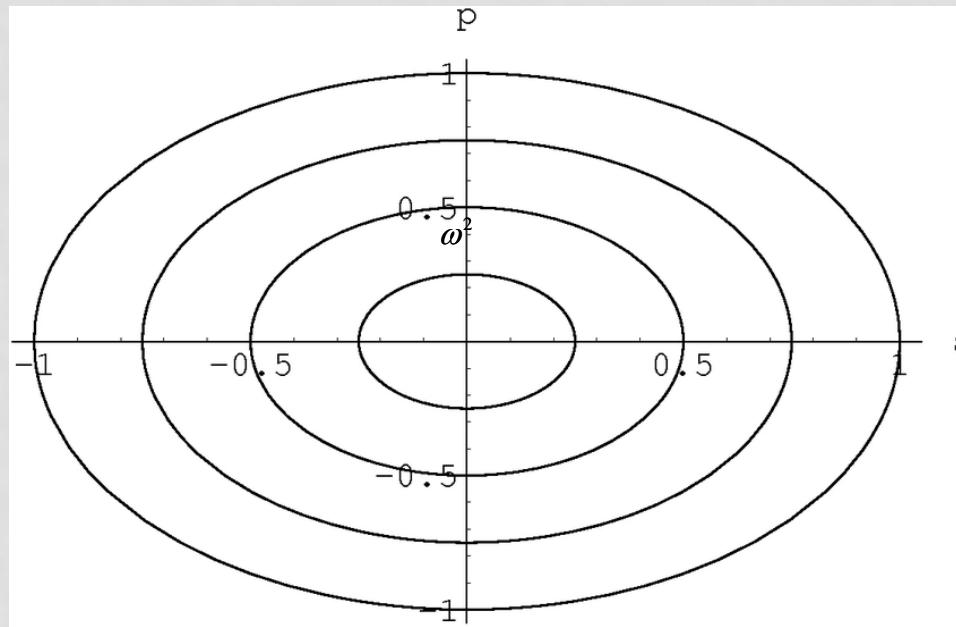
HARMONIC OSCILLATOR

Two-variable form:

$$\frac{dx}{dt} = f(x,y) = wy$$

$$\frac{dy}{dt} = -wx$$

$$\frac{d^2x}{dt^2} = -w^2 x$$



Solution is strictly dependent on the initial condition

Neural oscillations are Limit Cycle type

HOPF OSCILLATOR

$$\dot{x} = -y + \mu x(1 - x^2 - y^2)$$

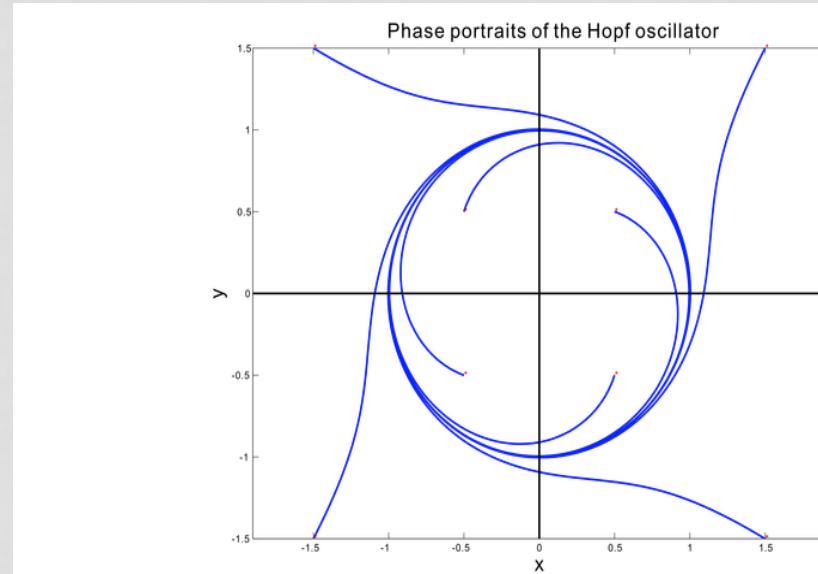
$$\dot{y} = x + \mu y(1 - x^2 - y^2)$$

In polar coordinates the equations have an elegant form.

Let $x = r \cos(\theta)$, $y = r \sin(\theta)$

$$\dot{r} = \mu r(1 - r^2)$$

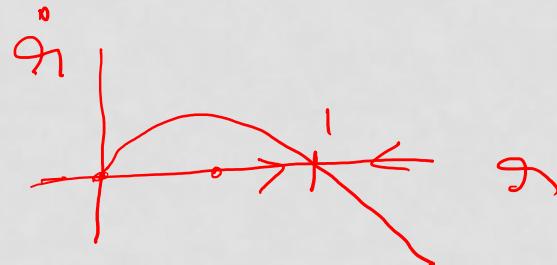
$$\dot{\theta} = 1$$



$$\begin{aligned}x\dot{x} + y\dot{y} &= \mu(x^2 + y^2)(1 - x^2 - y^2) \\&= \mu r^2(1 - r^2)\end{aligned}$$

$$\frac{1}{2} \frac{d(r^2)}{dt} = \mu r^2(1 - r^2)$$

$$\frac{dr}{dt} = \mu r(1 - r^2)$$



Now consider, $\theta = \arctan(y/x)$

- Differentiating both sides,

$$\dot{\theta} = \frac{1}{1 + \frac{y^2}{x^2}} \frac{x\dot{y} - y\dot{x}}{x^2}$$

$$x\dot{y} - y\dot{x} = x^2 + y^2 = r^2 \quad \dot{\omega} = 1$$

$$\dot{\theta} = \omega$$

IN COMPLEX VARIABLE FORM

$$\frac{dz}{dt} = i\omega z + (1 - |z|^2)z$$

$$z = x + iy$$

Let

$$z = re^{i\theta}$$

$$\frac{dz}{dt} = e^{i\theta} \frac{dr}{dt} + ire^{i\theta} \frac{d\theta}{dt}$$

$$e^{i\theta} \frac{dr}{dt} + ire^{i\theta} \frac{d\theta}{dt} = i\omega r e^{i\theta} + (1 - r^2) r e^{i\theta}$$

cancel $e^{i\theta}$,

$$\frac{dr}{dt} + ir \frac{d\theta}{dt} = i\omega r + (1 - r^2)r$$

$$\frac{dr}{dt} = (1 - r^2)r$$

$$\frac{d\theta}{dt} = \omega$$

Equating real and
imaginary parts:

NO NONLINEARITY → SHM

$$\frac{dz}{dt} = i\omega z + (1 - |z|^2)z$$

Eliminating nonlinearity...

$$\frac{dz}{dt} = i\omega z$$

$$\frac{dx}{dt} + i \frac{dy}{dt} = i\omega(x + iy) = \omega(ix - y)$$

$$\frac{dx}{dt} = \omega y$$

$$\frac{dy}{dt} = -\omega x$$

$$A = \begin{bmatrix} 0 & \omega \\ -\omega & 0 \end{bmatrix}$$

$$\lambda = \pm i\omega$$

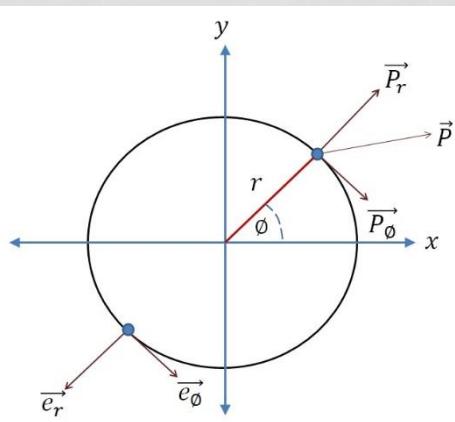
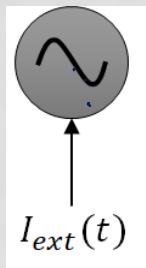
RESEARCH AGENDA

- To construct general oscillatory neural networks using Hopf Oscillators:
 - Oscillatory Hopfield networks (done) ✓
 - Oscillatory Auto Encoders (done) ✓
 - Oscillatory Self-organizing Maps (SOMs) ✓
 - Oscillatory Deep networks ✓
 - Oscillatory Convolutional Networks etc etc ✓

COMPLEX ADAPTIVE HOPF OSCILLATOR:

$$\dot{z} = z(\mu + i\omega - |z|^2) + \varepsilon I_{ext}(t)$$

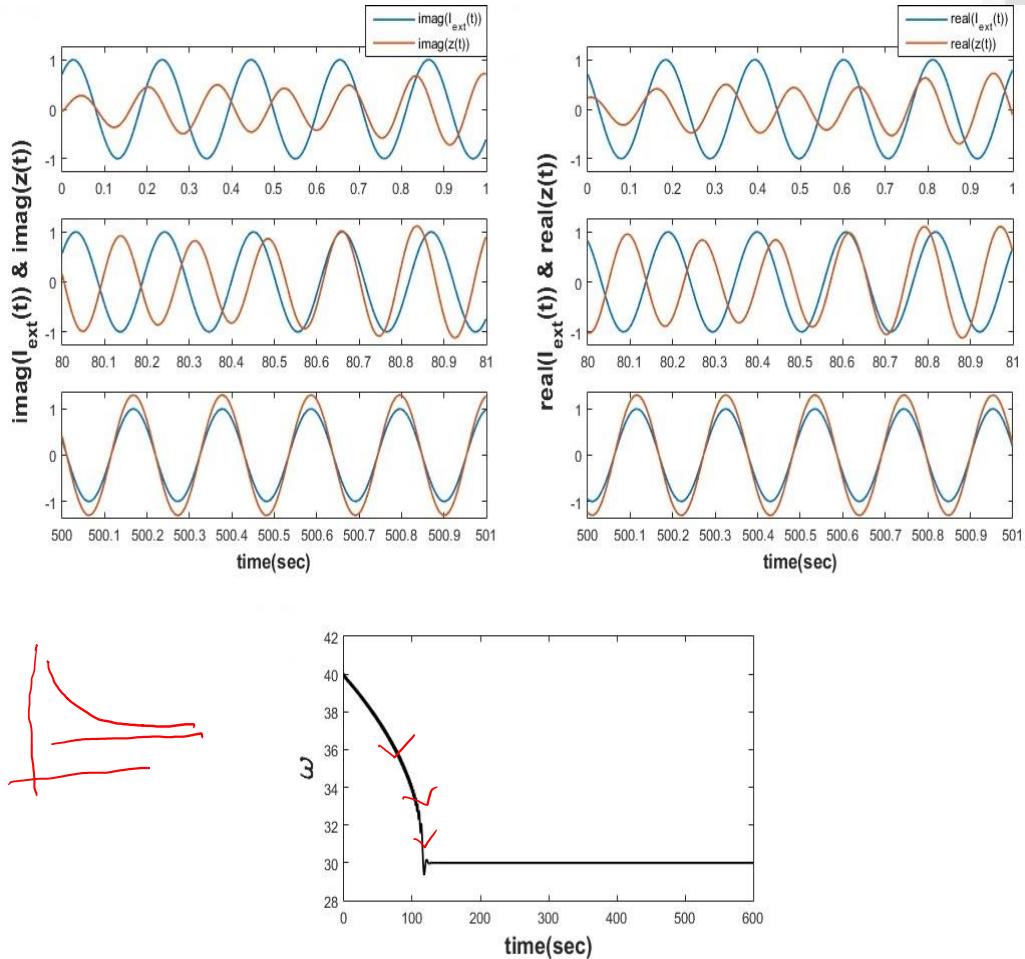
$$\dot{\omega} = -\varepsilon I_{ext}(t) \sin \phi$$



$$I_{ext}(t) = I_0 e^{i(\omega_0 t + \varphi)}$$

$$\dot{\omega} = -\varepsilon (\text{real}(I_{ext}(t)) \sin \phi - \text{img}(I_{ext}(t)) \cos \phi)$$

$$\dot{\omega} = -\varepsilon I_0 \sin(\phi - \omega_0 t - \varphi)$$



Righetti et al., 2005

(Biswas & Chakravarthy, 2020)

In this scenario, $I_{ext}(t)$ is a complex sinusoidal signal. It is straight forward to derive the learning rule for the natural frequency of the oscillator if eq. 2a is represented in the Cartesian and polar coordinate forms, respectively, as follows:

$$\begin{aligned}\dot{x} &= (\mu - r^2)x - \omega y + \varepsilon I_0 \cos(\omega_0 t + \varphi) \\ \dot{y} &= (\mu - r^2)y + \omega x + \varepsilon I_0 \sin(\omega_0 t + \varphi)\end{aligned}\tag{2a1}$$

$$\begin{aligned}\dot{r} &= (\mu - r^2)r + \varepsilon I_0 \cos(\omega_0 t + \varphi - \emptyset) \\ \dot{\emptyset} &= \omega + \underbrace{\frac{\varepsilon I_0}{r} \sin(\omega_0 t + \varphi - \emptyset)}\end{aligned}\tag{2a2}$$

In the phase plane representation, it can be observed from eq. 2a2 that the influence caused by the input perturbation on the oscillator phase is $\frac{\varepsilon I_0}{r} \sin(\omega_0 t - \emptyset)$. Whereas from the Cartesian coordinate system



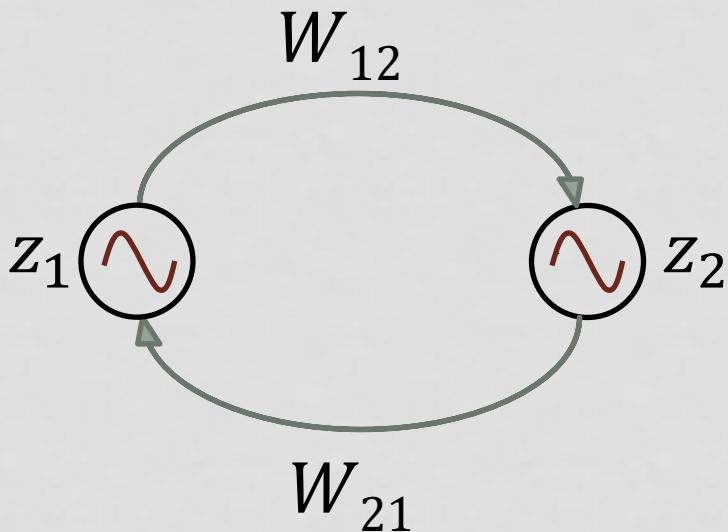
$$\phi = \omega t + \xi$$

$$i = -\epsilon \sin(\omega t + \xi - \omega_0 t - \varphi)$$

PROBLEM#1

A PAIR OF COUPLED HOPF OSCILLATORS

$$\begin{aligned}\dot{z}_1 &= (\mu - |z_1|^2)z_1 + i\omega_1 z_1 + W_{21} \text{real}(z_2) \\ \dot{z}_2 &= (\mu - |z_2|^2)z_2 + i\omega_2 z_2 + W_{12} \text{real}(z_1)\end{aligned}$$



What is the effect of the Coupling Strength (W) on the Phase Difference (ψ)?



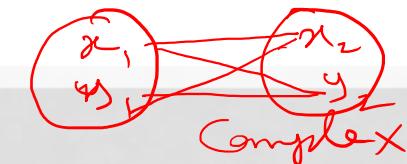
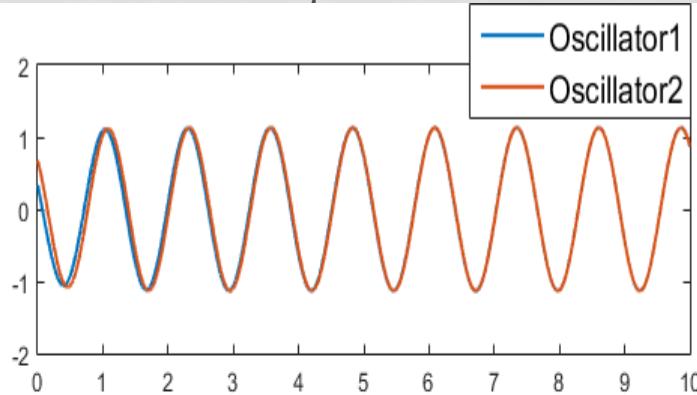
TWO HOPF OSCILLATORS

SAME FREQUENCY ($\omega_1 = \omega_2$), REAL COUPLING

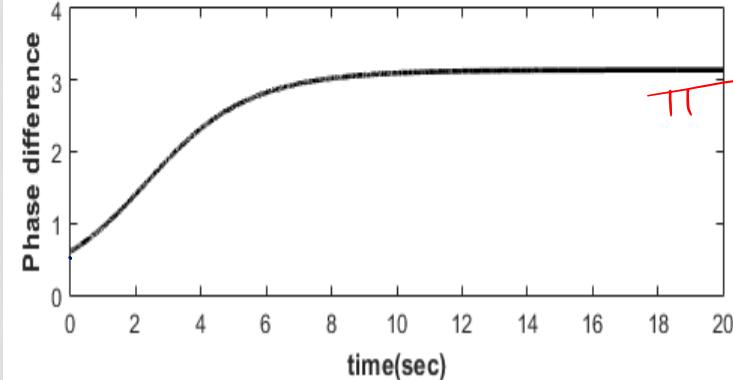
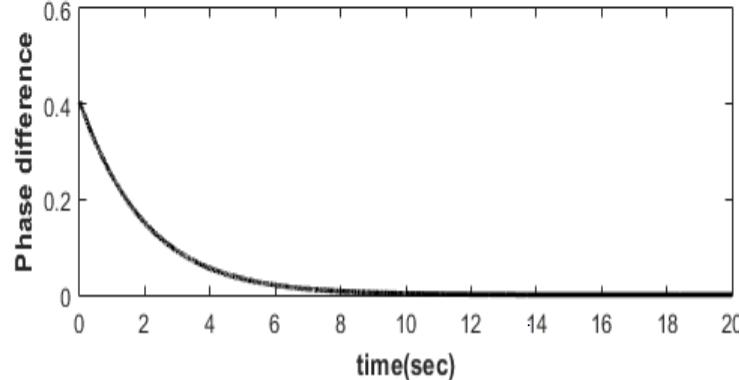
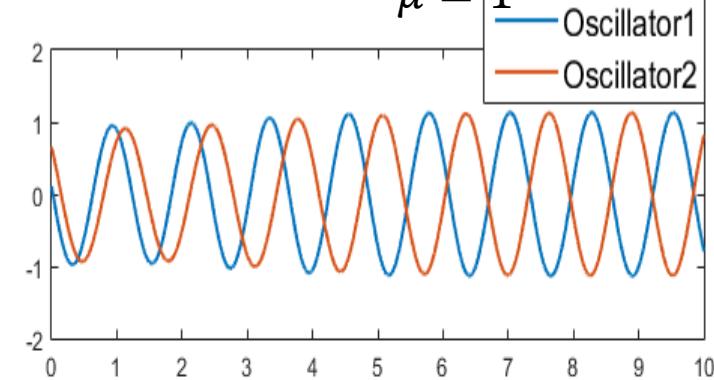


$$\begin{aligned}\omega_1 &= \omega_2 = 5 \\ W_{21} &= W_{12} = 0.5 \\ \mu &= 1\end{aligned}$$

$\omega_1 = \omega_2$
Symmetric coupling



$$\begin{aligned}\omega_1 &= \omega_2 = 5 \\ W_{21} &= W_{12} = -0.5 \\ \mu &= 1\end{aligned}$$



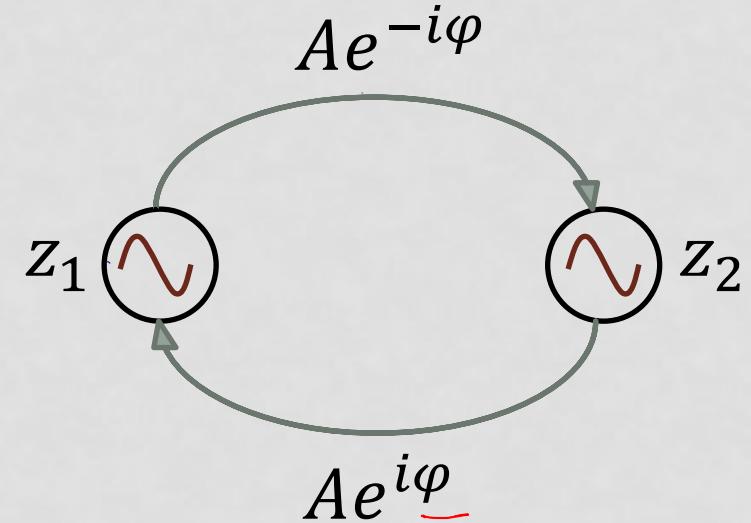
TWO COUPLED HOPF OSCILLATORS: SAME FREQUENCIES ($\omega_1 = \omega_2$), COMPLEX COUPLING

- Complex variable representation of Hopf oscillator.

$$\dot{z} = (\mu - |z|^2)z + i\omega z$$

Where, $z = x + iy = re^{i\theta}$

- Two oscillators coupled through complex lateral weights.



$$\dot{z}_1 = (\mu - |z_1|^2)z_1 + i\omega_1 z_1 + \underline{Ae^{i\varphi}} z_2$$

$$\dot{z}_2 = (\mu - |z_2|^2)z_2 + i\omega_2 z_2 + \underline{Ae^{-i\varphi}} z_1$$

THEORETICAL EXPLANATION

- Polar coordinate representation:

$$\dot{r}_1 = (\mu - r_1^2)r_1 + A \underline{r_2} \cos(\theta_2 - \theta_1 + \varphi)$$

$$\dot{\theta}_1 = \omega_1 + A \frac{\underline{r_2}}{r_1} \sin(\theta_2 - \theta_1 + \varphi)$$

$$\dot{r}_2 = (\mu - r_2^2)r_2 + A \underline{r_1} \cos(\theta_1 - \theta_2 - \varphi)$$

$$\psi = \theta_1 - \theta_2$$

$$\dot{\theta}_2 = \omega_2 + A \frac{\underline{r_1}}{r_2} \sin(\theta_1 - \theta_2 - \varphi)$$

$$\theta_1 - \theta_2 =$$

- From which:

$$\underline{\omega_1 = \omega_2}$$

arbitrary?

$$\dot{\psi} = \dot{\theta}_1 - \dot{\theta}_2 = (\omega_1 - \cancel{\omega_2}) - \frac{A \sin(\psi - \varphi)}{r_1 r_2} (r_1^2 + r_2^2)$$

- Assuming $\omega_1 = \omega_2$, at steady state:

$$\dot{\psi} = 2A \sin(\varphi - \psi) = 0$$

$$r_1 = r_2 = \sqrt{\mu}$$

- Which ensures the phase difference $(\theta_1 - \theta_2)$ between the two oscillators to be φ : the angle of the complex coupling weight.

NUMERICAL RESULTS

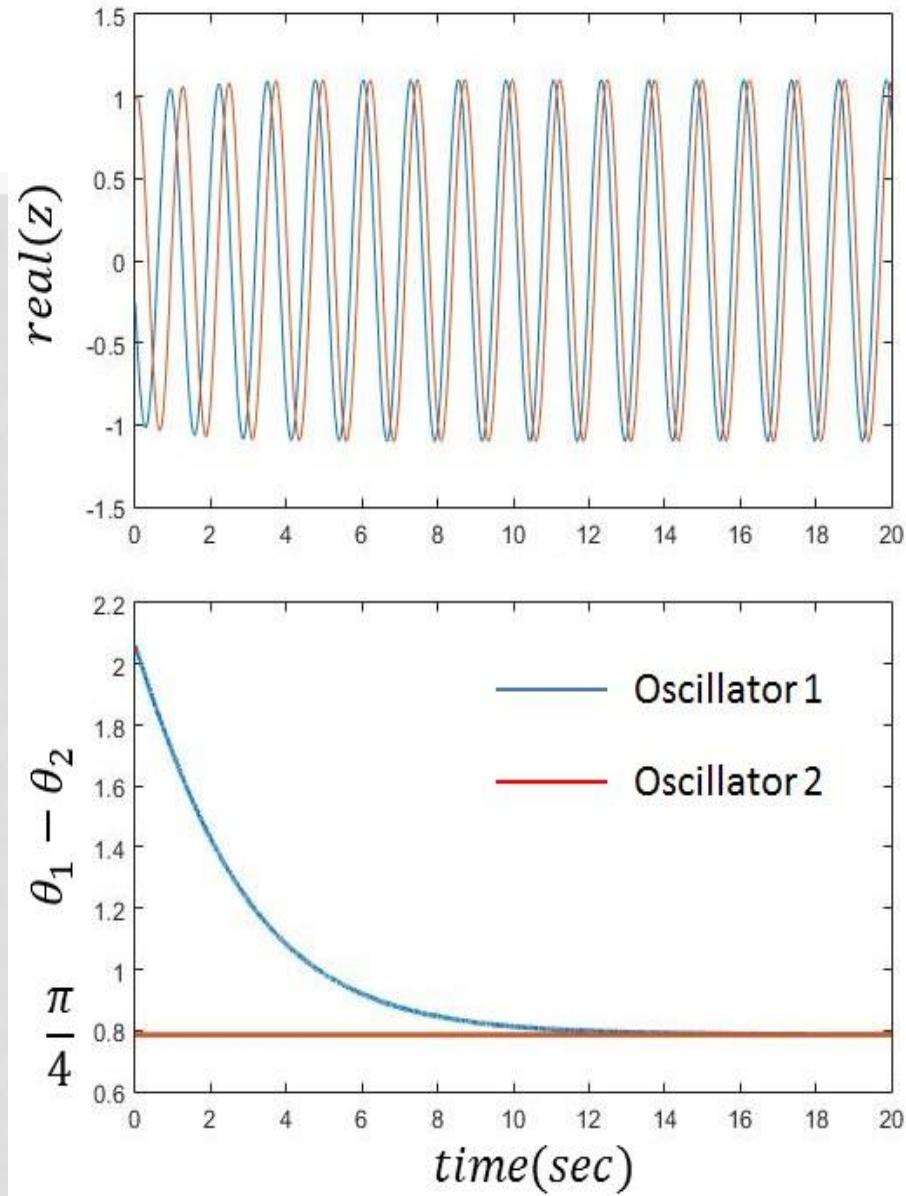
- Parameters:

$$A = 0.2, \varphi = \frac{\pi}{4}$$

$$\omega_1 = \omega_2 = 5$$

Where: $z_1 = r_1 e^{i\theta_1}$ and

$$z_2 = r_2 e^{i\theta_2}$$



COUPLING TWO HOPF OSCILLATORS NEARBY FREQUENCIES, REAL COUPLING

- When two Hopf oscillators, with slightly different frequencies ($\omega_1 \approx \omega_2$) are coupled, for a sufficiently strong coupling strength, the two oscillators come to a common frequency (ω) such that ($\omega_1 < \omega < \omega_2$), with oscillator #2 leading oscillator #1.

$$\omega_1 = \omega_2 + \epsilon \quad \epsilon \ll 1$$

$$\theta_1 > \theta_2$$

$$\begin{matrix} \omega_1 \rightarrow \omega \\ \omega_2 \end{matrix}$$



$$\omega_1 > \omega > \omega_2$$

$$\begin{matrix} \delta \cdot 1 Hz \\ \theta \\ L \\ P \end{matrix}$$

$$\begin{matrix} R \\ 10^0 Hz \end{matrix}$$

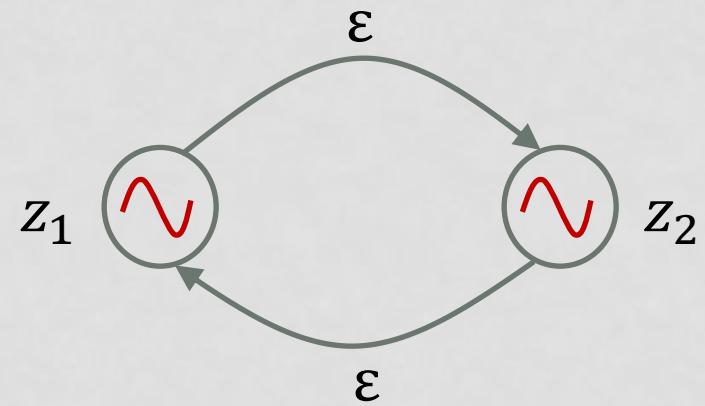
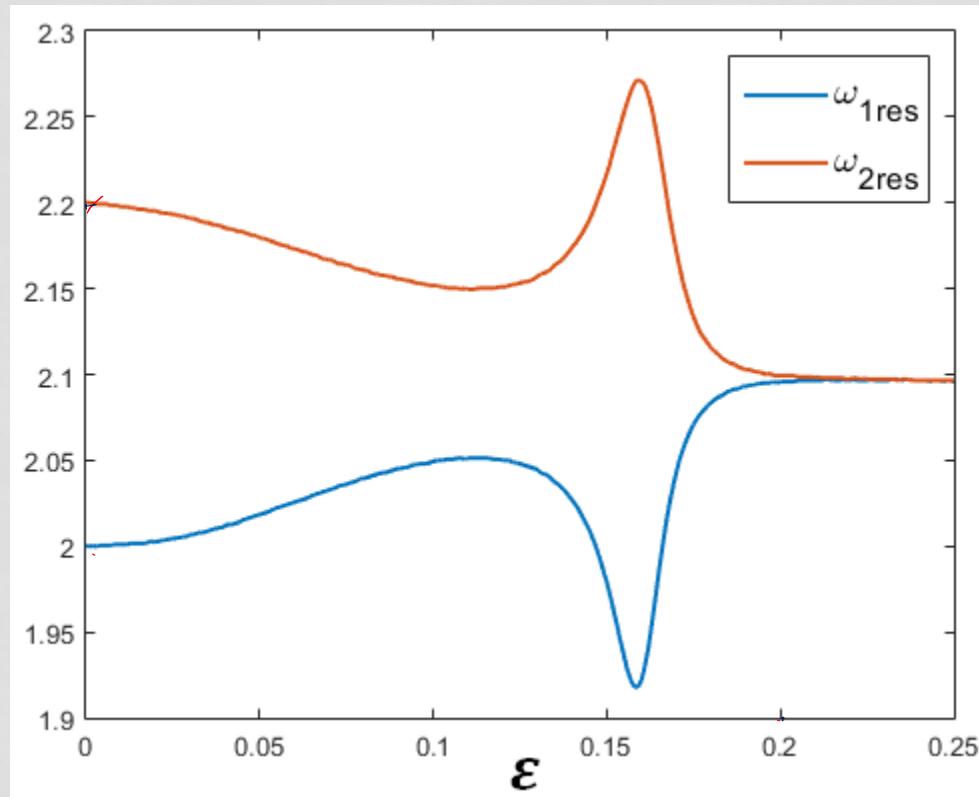


$$\dot{z}_1 = (\mu - |z_1|^2)z_1 + i\omega_1 z_1 + \varepsilon \times \text{real}(z_1)$$

$$\dot{z}_2 = (\mu - |z_2|^2)z_2 + i\omega_2 z_2 + \varepsilon \times \text{real}(z_2)$$

Simulation parameters:

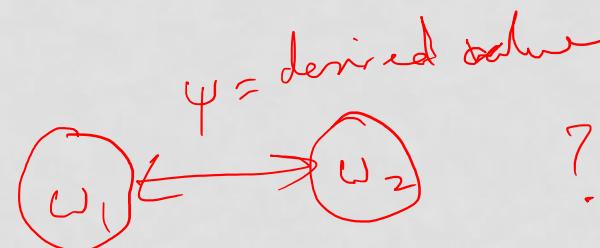
$$\mu = 1, \omega_1 = 2, \omega_2 = 2.2$$



COUPLING TWO HOPF OSCILLATORS WITH ARBITRARILY DIFFERENT FREQUENCIES

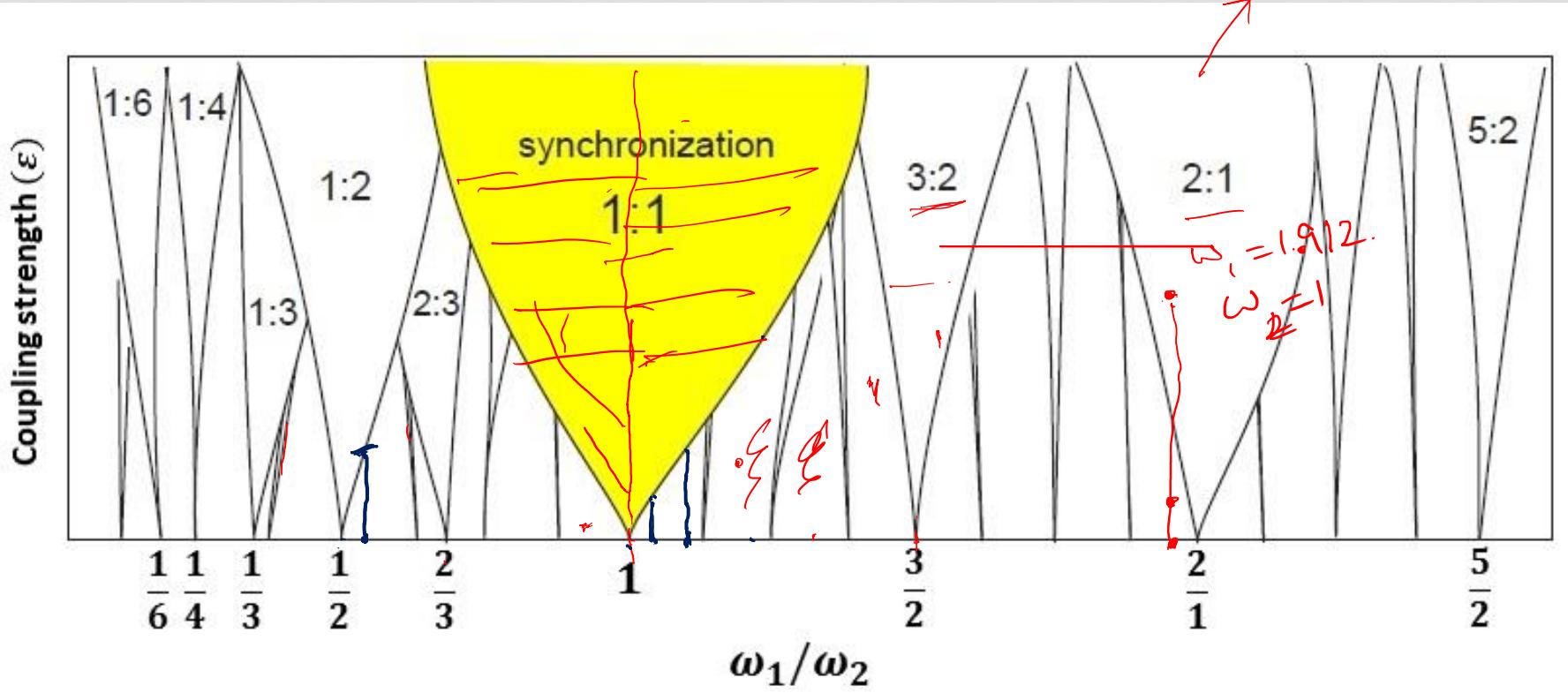
- When two oscillators with different intrinsic frequencies are coupled through real coupling weights a new behavior emerges while analyzing their synchronization behavior, called Arnold Tongues.

$$\begin{aligned}\dot{z}_1 &= (\mu - |z_1|^2)z_1 + i\omega_1 z_1 + \varepsilon \times \text{real}(z_1) \\ \dot{z}_2 &= (\mu - |z_2|^2)z_2 + i\omega_2 z_2 + \varepsilon \times \text{real}(z_2)\end{aligned}$$



COUPLING TWO HOPF OSCILLATORS: ARBITRARILY DIFFERENT FREQUENCIES, REAL COUPLING GENERAL SCENARIO

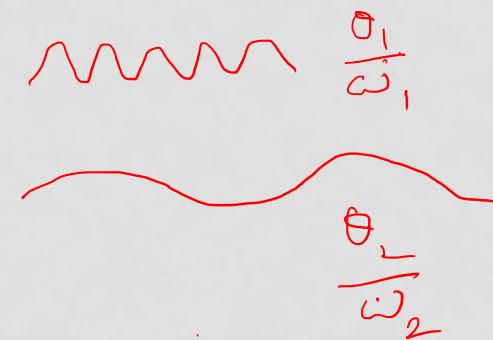
ARNOLD TONGUES



NORMALIZED PHASE DIFFERENCE

- When the frequencies are different, the usual phase difference cannot be constant
- But normalized phase difference can be constant
- normalized phase difference:

$$\psi = \frac{\theta_1}{\omega_1} - \frac{\theta_2}{\omega_2}$$



- This is what is constant when two oscillators are entrained in a simple integral ratio

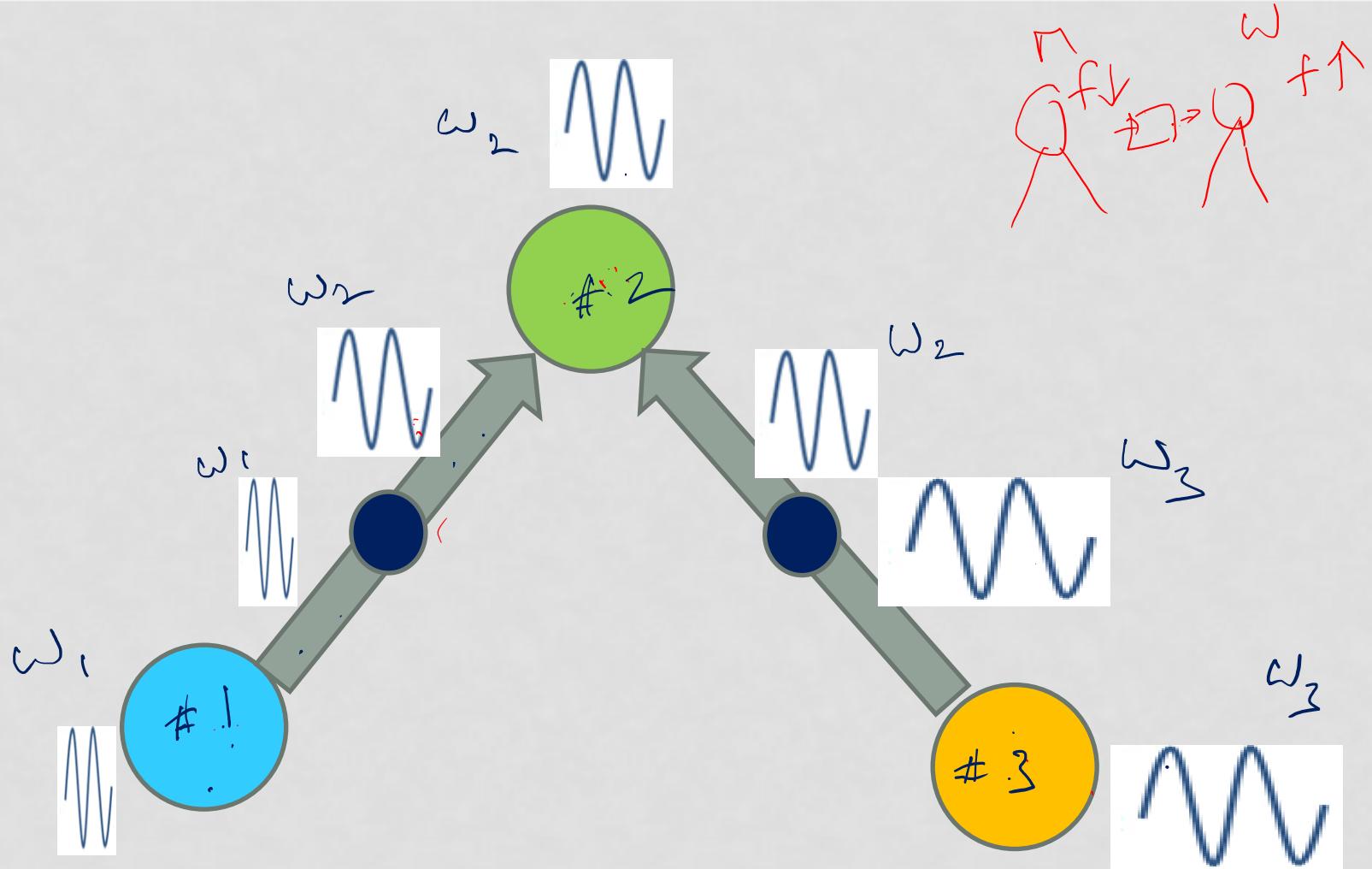
- How to achieve constant normalized phase difference for arbitrary frequencies?

Answer:

It requires a trick called

Power coupling

TRANSFORMING TRANSMITTER FREQUENCIES TO RECEIVER FREQUENCY





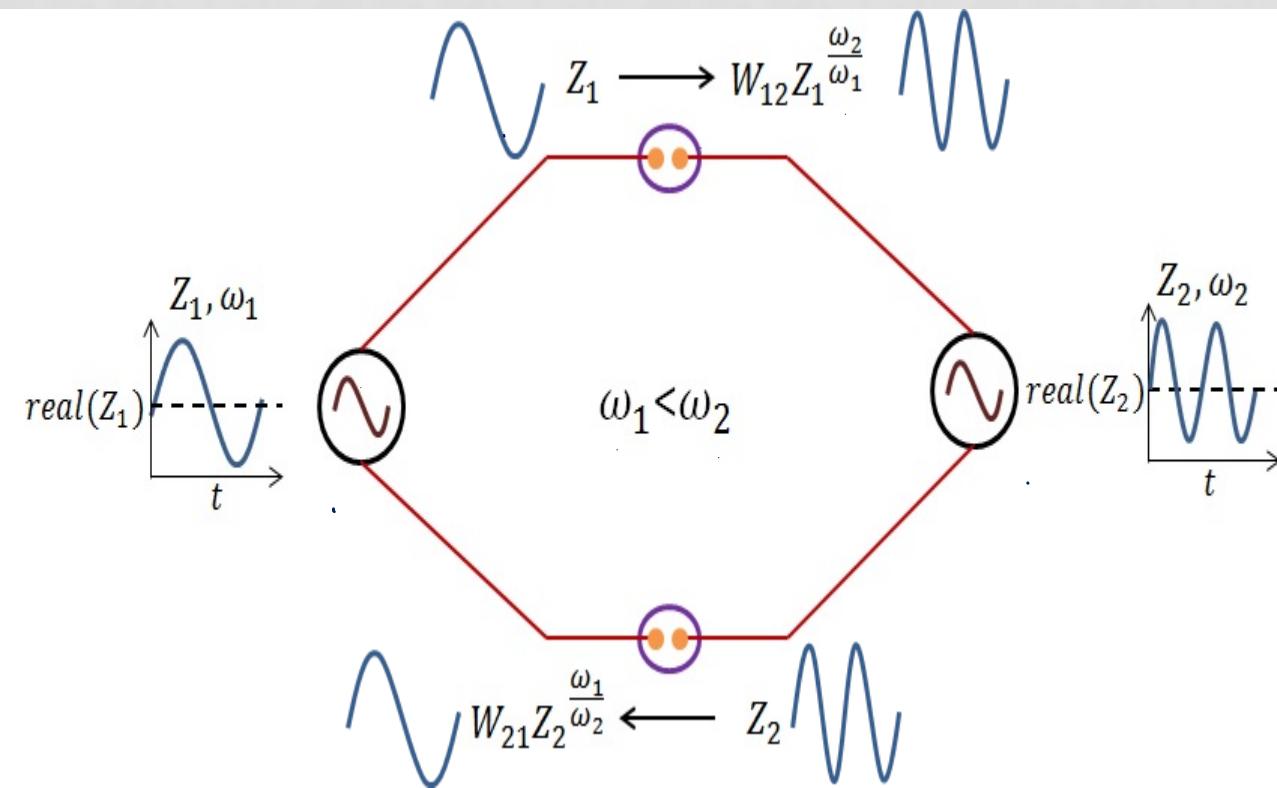
$$z_1 = \underline{A_1 e^{i\omega_1 t}}$$

How to transform ω_1 to ω_2 ?

$$(z_1)^{\circlearrowleft \omega_2 / \omega_1} = (A_1 e^{i\omega_1 t})^{\omega_2 / \omega_1} = (A_1 e^{i\omega_2 t})$$

POWER COUPLING

- Frequency of oscillator 1 is transformed into frequency of oscillator 2, as the signal passes over the connection



THEORETICAL ANALYSIS

- Two oscillators coupled through “power coupling”:

$$\begin{aligned}\dot{z}_1 &= (\mu - |z_1|^2)z_1 + i\omega_1 z_1 + Ae^{t\frac{\varphi}{\omega_2}} z_2 \frac{\omega_1}{\omega_2} \\ \dot{z}_2 &= (\mu - |z_2|^2)z_2 + i\omega_2 z_2 + Ae^{-i\frac{\varphi}{\omega_1}} z_1 \frac{\omega_2}{\omega_1}\end{aligned}$$

- The angular part of the polar form:

$$\begin{aligned}\dot{\theta}_1 &= \omega_1 + A \frac{r_2 \frac{\omega_1}{\omega_2}}{r_1} \sin \omega_1 \left(\frac{\theta_2}{\omega_2} - \frac{\theta_1}{\omega_1} + \frac{\varphi}{\omega_1 \omega_2} \right) \\ \dot{\theta}_2 &= \omega_2 + A \frac{r_1 \frac{\omega_2}{\omega_1}}{r_2} \sin \omega_2 \left(\frac{\theta_1}{\omega_1} - \frac{\theta_2}{\omega_2} - \frac{\varphi}{\omega_1 \omega_2} \right)\end{aligned}$$

- We denote the normalized phase difference:

$$\psi = \frac{\theta_1}{\omega_1} - \frac{\theta_2}{\omega_2}$$

- The derivative of normalized phase difference w.r.t time:

$$\dot{\psi} = \frac{\dot{\theta}_1}{\omega_1} - \frac{\dot{\theta}_2}{\omega_2}$$

$$\dot{\psi} = \frac{Ar_2 \frac{\omega_1}{\omega_2}}{\omega_1 r_1} \sin \omega_1 \left(\frac{\varphi}{\omega_1 \omega_2} - \psi \right) + \frac{Ar_1 \frac{\omega_2}{\omega_1}}{\omega_2 r_2} \sin \omega_2 \left(\frac{\varphi}{\omega_1 \omega_2} - \psi \right)$$

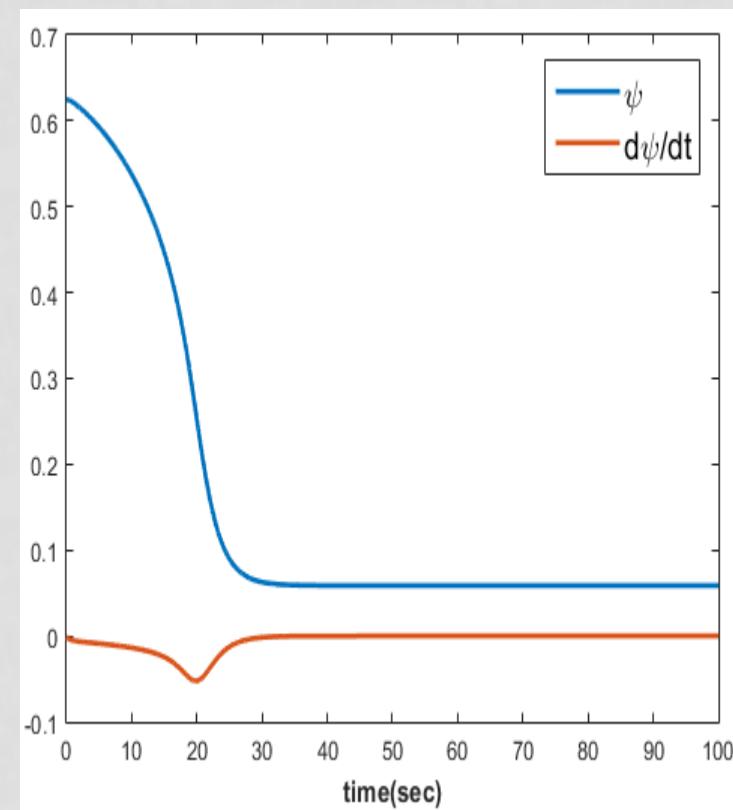
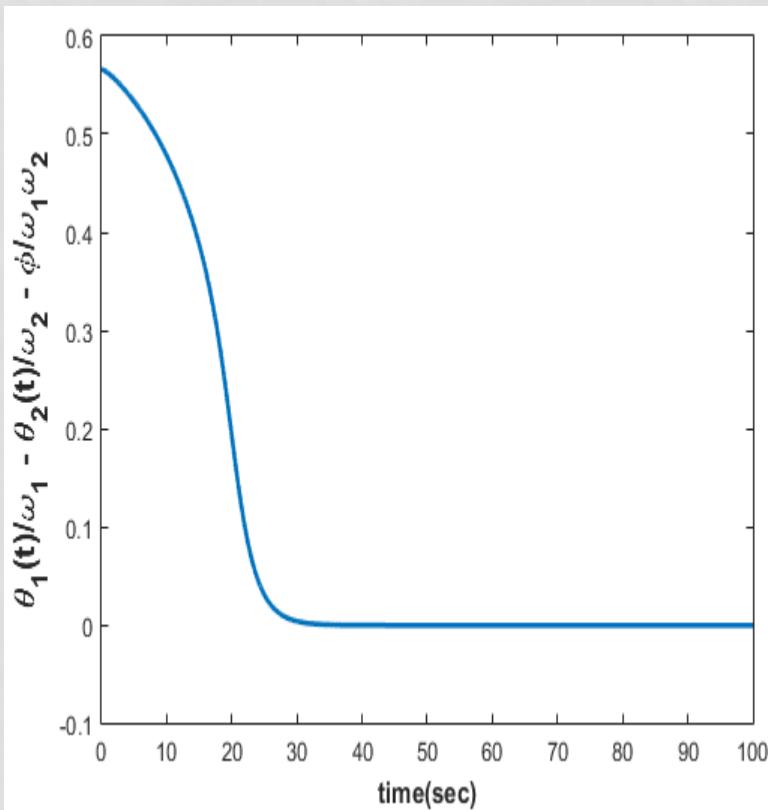
- At steady state; $\psi = \text{constant}$, $\dot{\psi} = 0$ and $r_1 = r_2 = 1$:

$$\dot{\psi} = \frac{A}{\omega_1} \sin \omega_1 \left(\frac{\varphi}{\omega_1 \omega_2} - \psi \right) + \frac{A}{\omega_2} \sin \omega_2 \left(\frac{\varphi}{\omega_1 \omega_2} - \psi \right) = 0$$

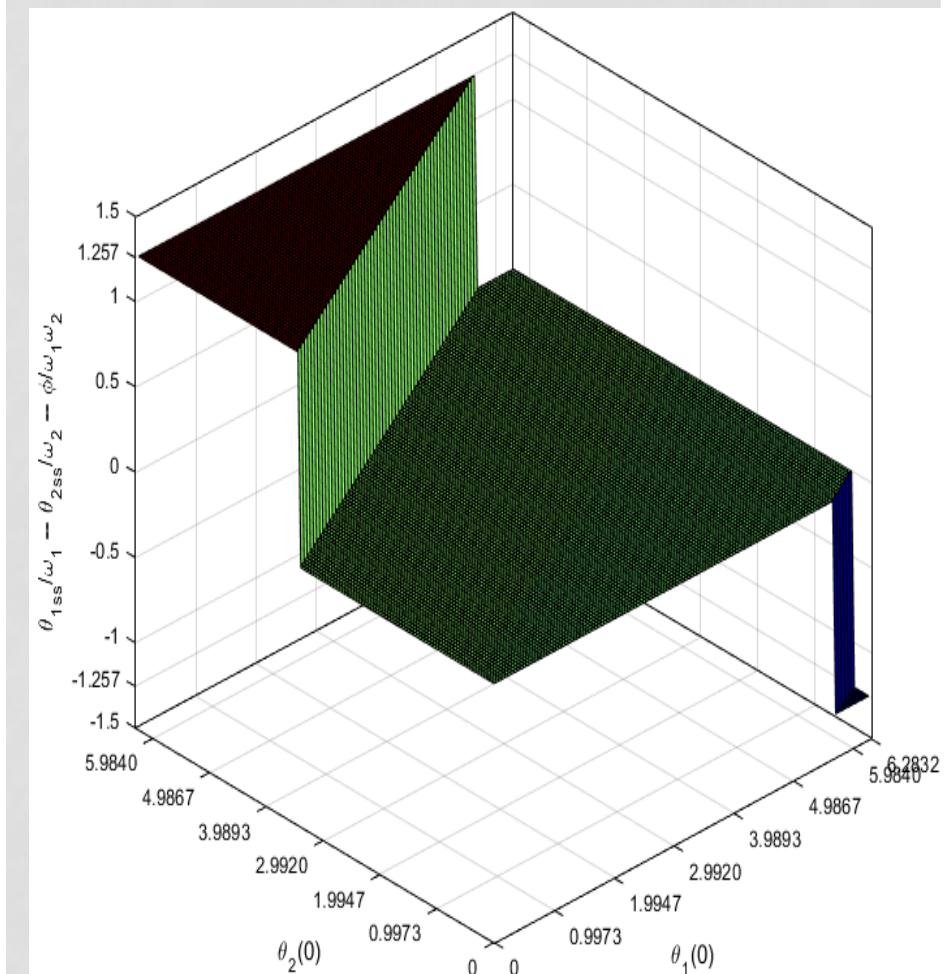
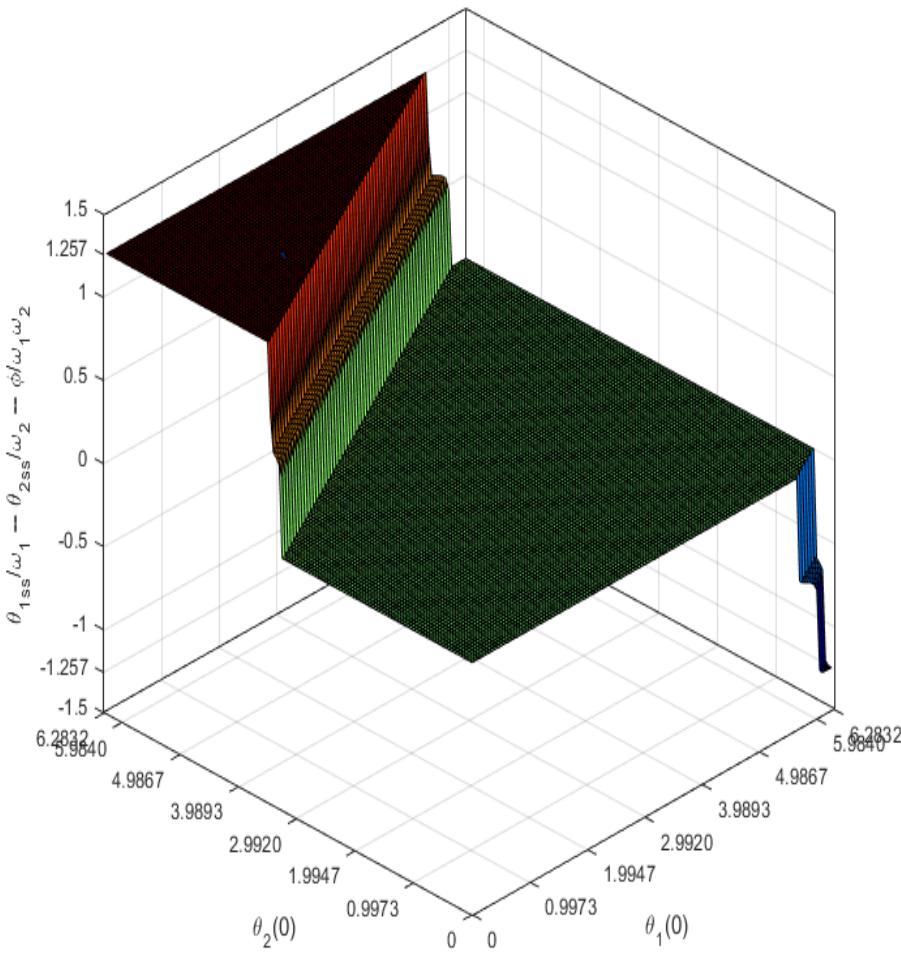
- The steady state normalized phase difference satisfies the equation:

$$\psi = \frac{\theta_1}{\omega_1} - \frac{\theta_2}{\omega_2} = \frac{\varphi}{\omega_1 \omega_2}$$

- The $\sigma = \frac{\theta_1(t)}{\omega_1} - \frac{\theta_2(t)}{\omega_2} - \frac{\varphi}{\omega_1 \omega_2}$, $\psi(t)$ and $\dot{\psi}(t)$ variation w.r.t time for two 'power coupled' Hopf oscillators with the following parameters.
- $\omega_1 = 5$, $\omega_2 = 10$, $A_{12} = A_{21} = 0.2$, $\varphi = 2.9644$. $\varphi/50 = \underline{\psi}$



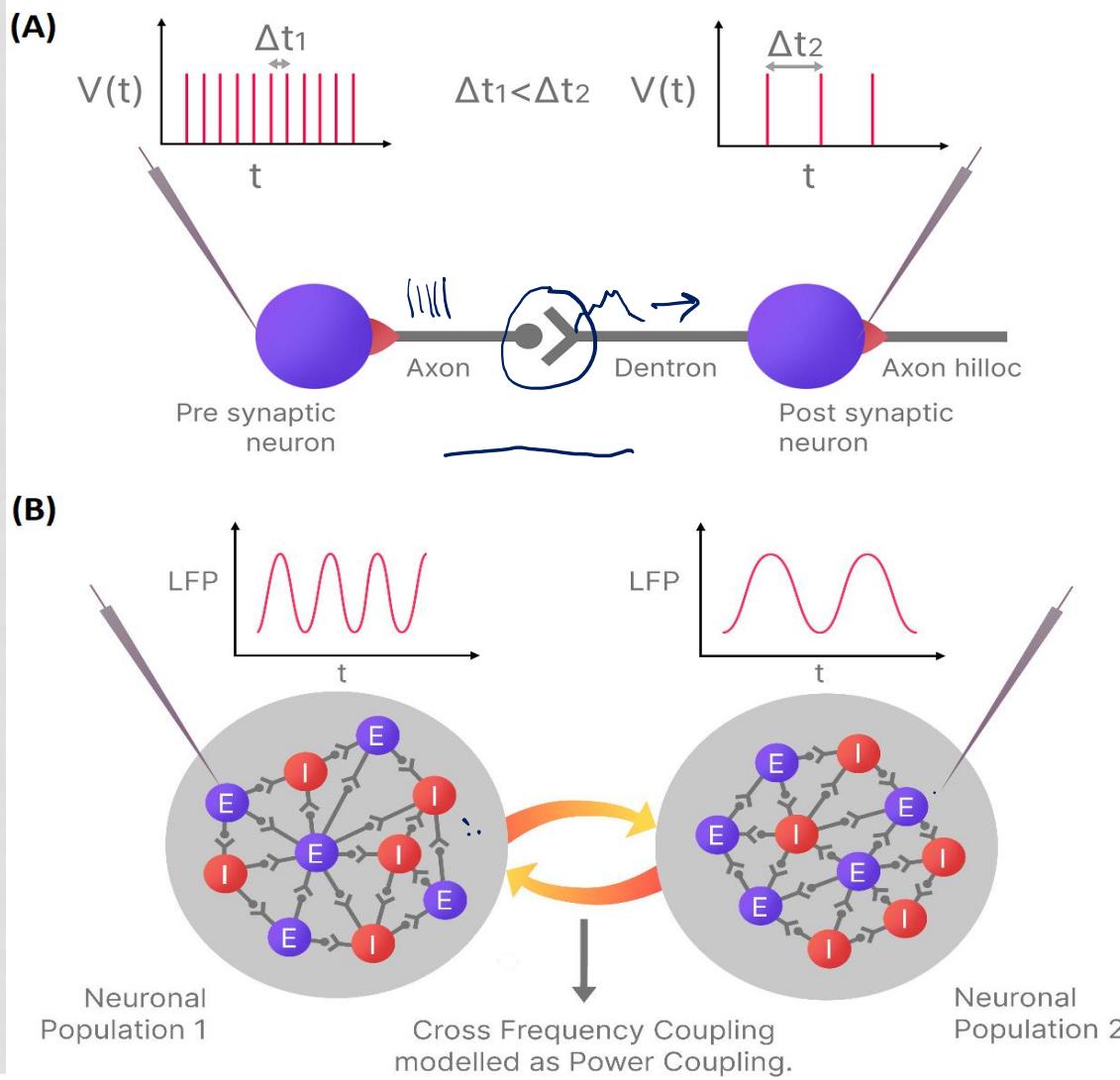
- Depending on initial values of θ_1 and θ_2 , σ_{ss} can attain any of the solutions of $\dot{\psi}_{ss} = 0$ equation.
- The solutions other than $\sigma_{ss} = 0$ are: $\sigma_{ss} = \frac{2\pi n_1}{\omega_1} = \frac{2\pi n_2}{\omega_2}$, where $\frac{n_1}{n_2} = \frac{\omega_1}{\omega_2}$
- In the simulated scenario the stable solutions are: $\sigma_{ss} = 0, \frac{2n_1\pi}{\omega_1} = \frac{2n_2\pi}{\omega_2}$



$\frac{\omega_1}{\omega_2}$ = rational #

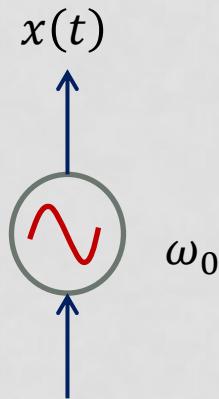
$$\sim \sin(\varphi_1 \left(\frac{2\pi n_1}{\omega_1} \right)) + \sim \sin(\varphi_2 \left(\frac{2\pi n_2}{\omega_2} \right)) = 0$$

BIOLOGICAL INTERPRETATION



FREQUENCY TRACKING

The intrinsic frequency of the oscillator can track the input frequency



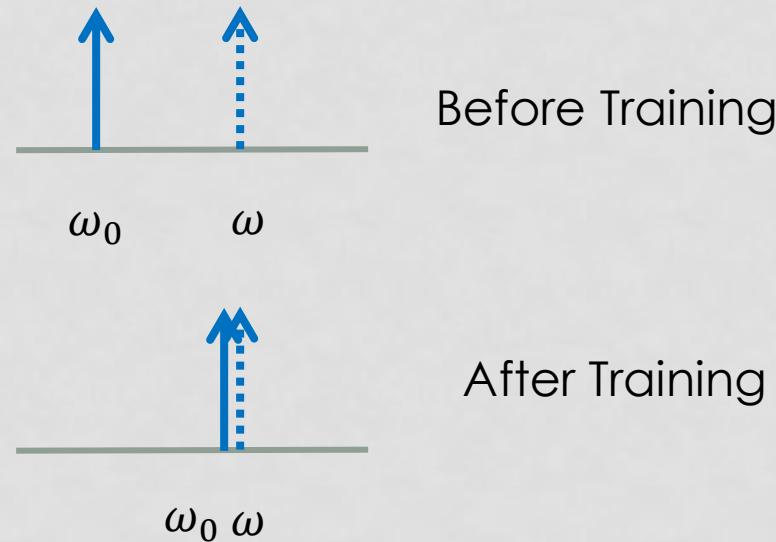
$$F(t) = Ae^{i(\omega t)}$$

$$\dot{x} = (1 - r^2)x - \omega y + \varepsilon F(t)$$

$$\dot{y} = (1 - r^2)y + \omega x$$

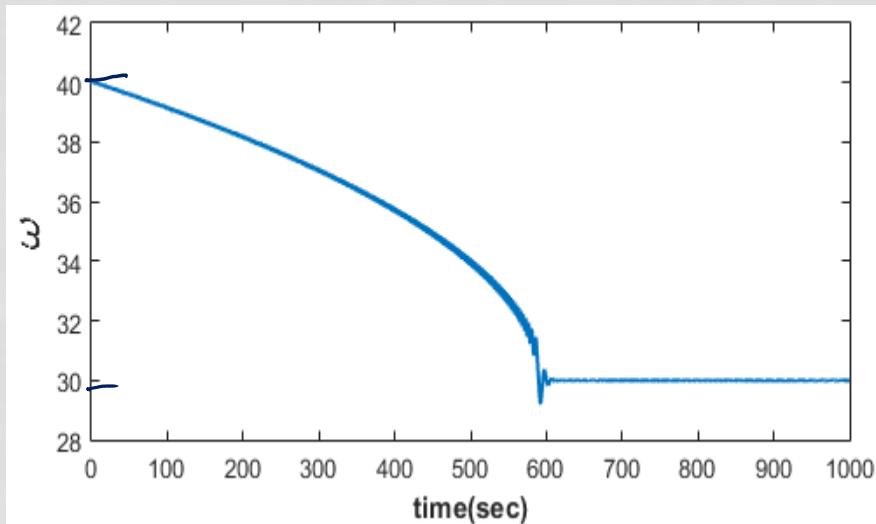
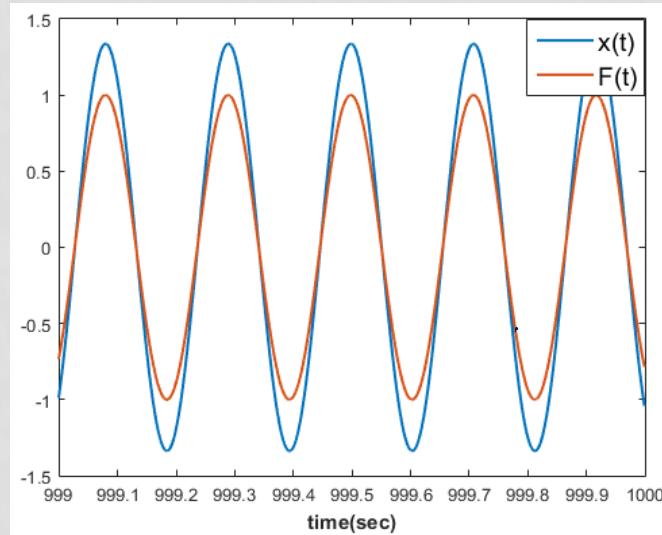
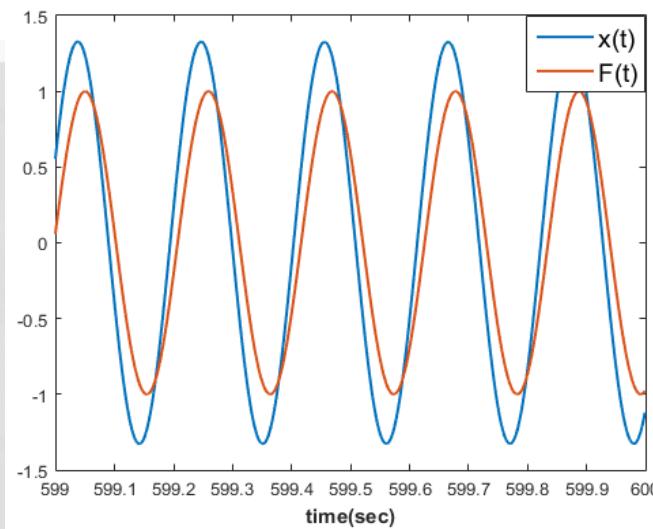
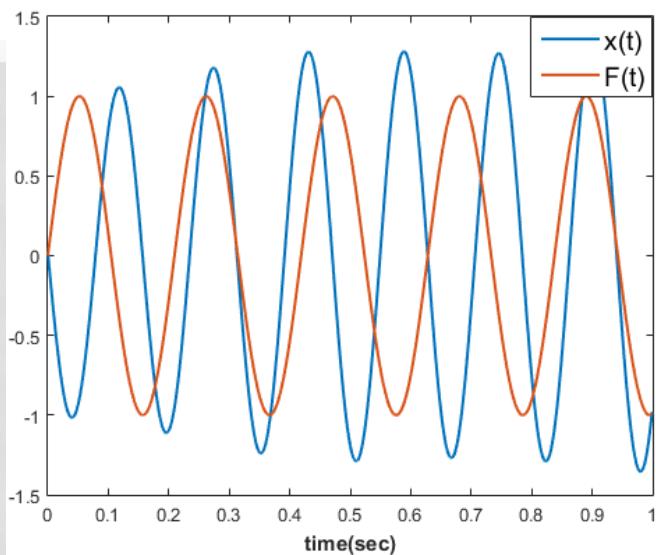
$$\dot{\omega} = -\varepsilon F(t) \frac{y}{r}$$

(Righetti et al 2005)

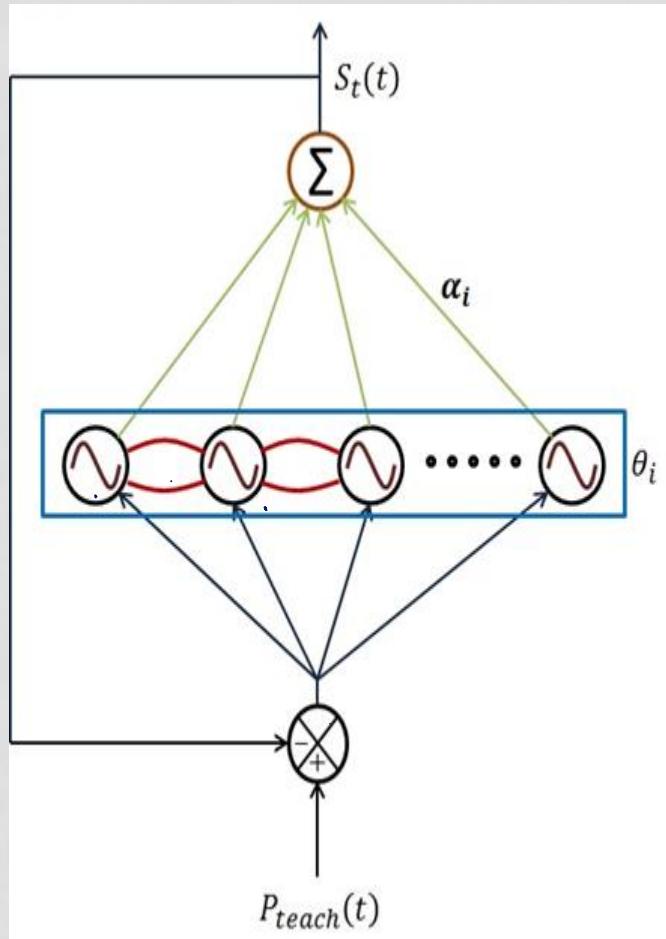


$$\begin{aligned}\dot{x} &= (1 - r^2)x - \omega y + \varepsilon F(t) \\ \dot{y} &= (1 - r^2)y + \omega x \\ \dot{\omega} &= -\varepsilon F(t) \frac{y}{r}\end{aligned}$$

$$\omega(0) = 40, \quad \varepsilon = 0.9, \quad F(t) = \sin 30t$$



COUPLED N-OSCILLATOR NETWORK



Reconstructing a time series
By taking a weighted sum of outputs
of N-oscillator network



RECONSTRUCTING A TIME SERIES

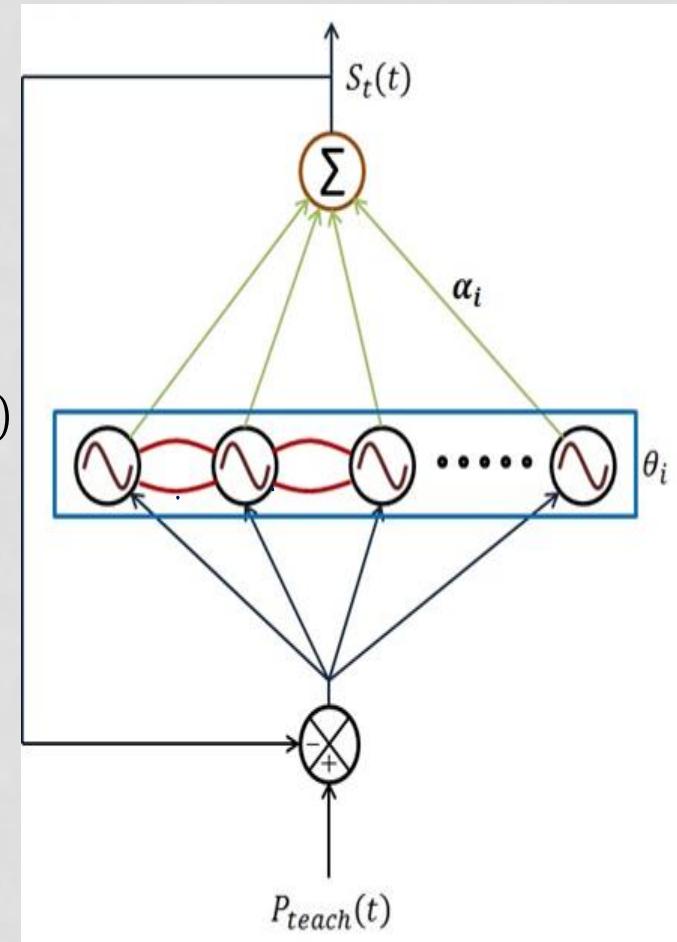
- Network of n no of Kuramoto oscillators can learn Some sort of Fourier decomposition of any $P_{teach}(t)$ signal.

$$\dot{Z}_i = (\mu - |Z_i|^2)Z_i + i\omega Z_i + \sum_j A_{ij} e^{i\frac{\phi_{ij}}{\omega_j} Z_j \frac{\omega_i}{\omega_j}} + \varepsilon(P_{teach}(t) - S(t))$$

$$\begin{aligned} \dot{\theta}_j &= \omega_j + \sum_k A_{jk} \sin \omega_j \left(\frac{\theta_k}{\omega_k} - \frac{\theta_j}{\omega_j} + \frac{\phi_{jk}}{\omega_j \omega_k} \right) \\ &+ \varepsilon(P_{teach}(t) - S(t)) \sin \theta_j \end{aligned}$$

$$\dot{\omega}_j = -\eta_\omega \underbrace{(P_{teach}(t) - S(t))}_{\text{sin } \theta_j} \sin \theta_j \quad \checkmark$$

$$S(t) = \sum_j \alpha_j \cos \theta_j \quad (\text{reconstructed signal})$$



TEST CASE#1

THE INTRINSIC FREQUENCIES OF THE OSCILLATORS LEARNED THE FREQUENCY COMPONENTS OF THE INPUT SIGNAL

$$P_{teach}(t) = A_1 \cos(\omega_1 t + \phi_1) + A_2 \cos(\omega_2 t + \phi_2) + A_3 \cos(\omega_3 t + \phi_3)$$

Where: $A_1 = 1.5909, A_2 = 1.377, A_3 = 1.3938$

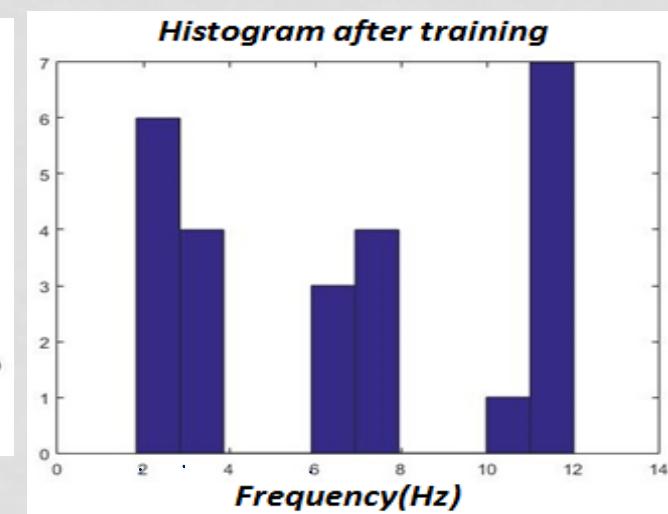
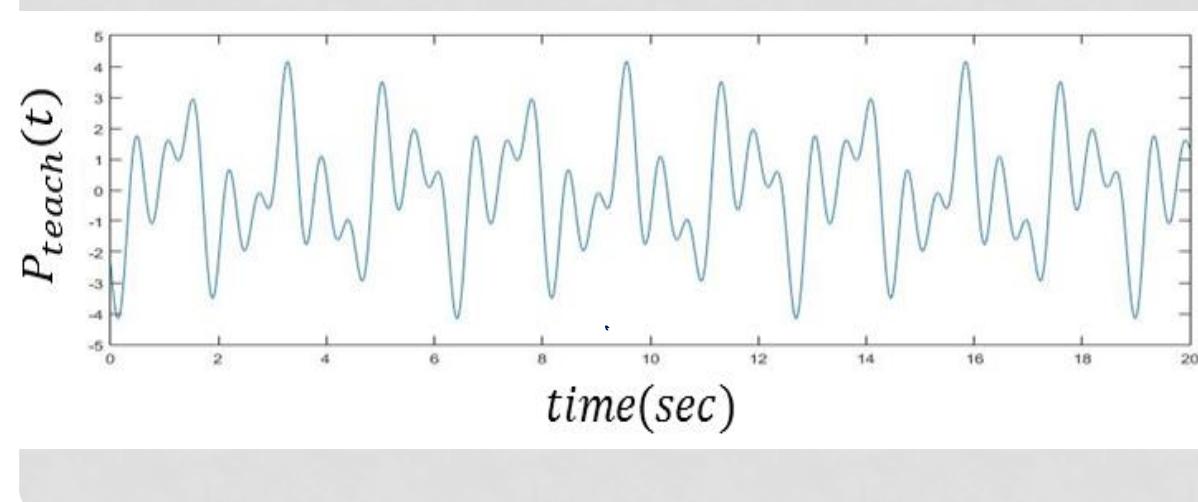
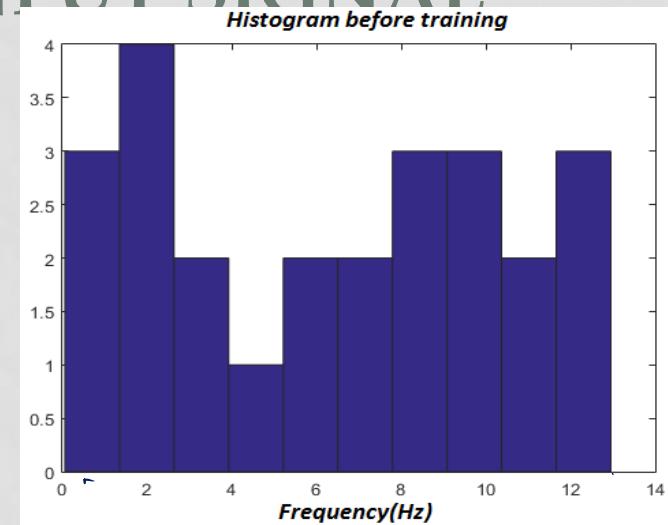
$$\omega_1 = 3, \omega_2 = 7, \omega_3 = 11$$

$$\phi_1 = 2.6911, \phi_2 = 2.6579, \phi_3 = 1.3755$$

$$n = 25$$

$$\eta_\omega = 0.001$$

$$\alpha_j = 0.2$$



GENERALIZED NETWORK#2
 FROM LEARNING FREQUENCY SPECTRUM REPRESENTATION TO
 MODELLING ARBITRARY M NO OF TIME SERIES SIGNALS HAVING SAME
 FREQUENCY COMPONENTS PRESENT IN THE SIGNAL USED TO TRAIN
 NETWORK#1

$$L = \frac{1}{2} \sum_{i=1}^m \sum_t \|Y_{d_i}(t) - Y_{p_i}(t)\|^2$$

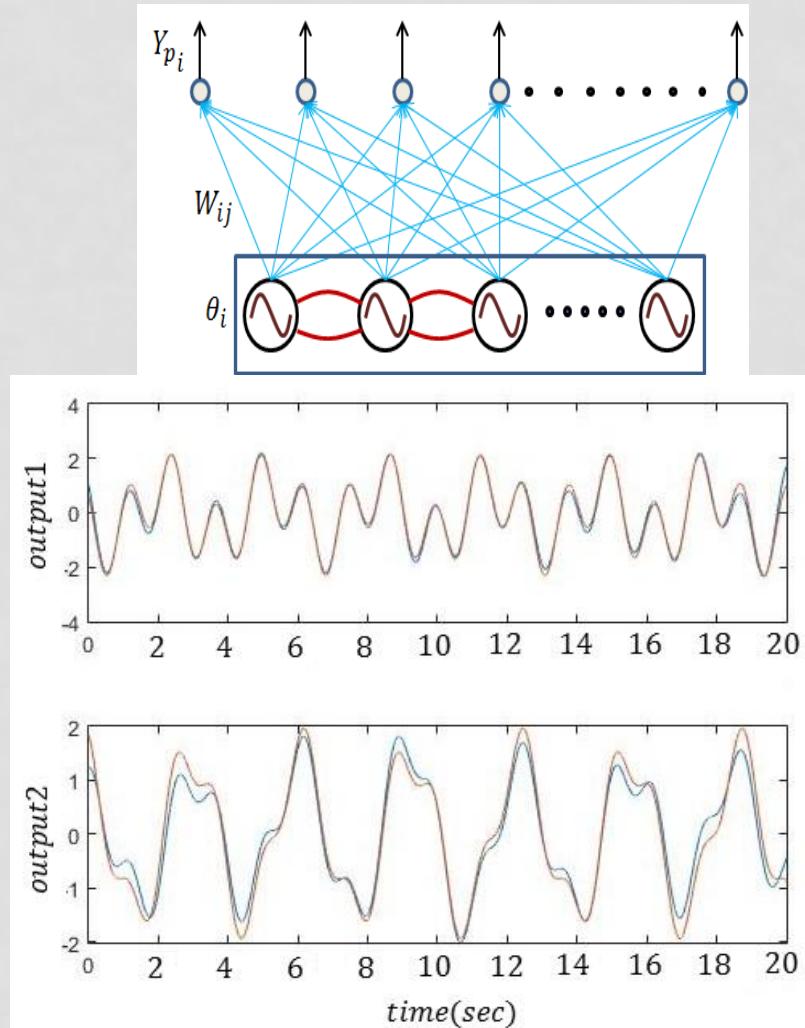
$$Y_{p_i}(t) = \operatorname{real} \left(\sum_{j=1}^n W_{ij} e^{i\theta_j} \right)$$

$$W_{ij} = K_{ij} e^{i\phi_{ij}}$$

$$\Delta K_{ij} = (-1) \eta_{K_{ij}} \sum_t (Y_{d_i}(t) - Y_{p_i}(t)) \cos(\theta_j(t) + \phi_{ij})$$

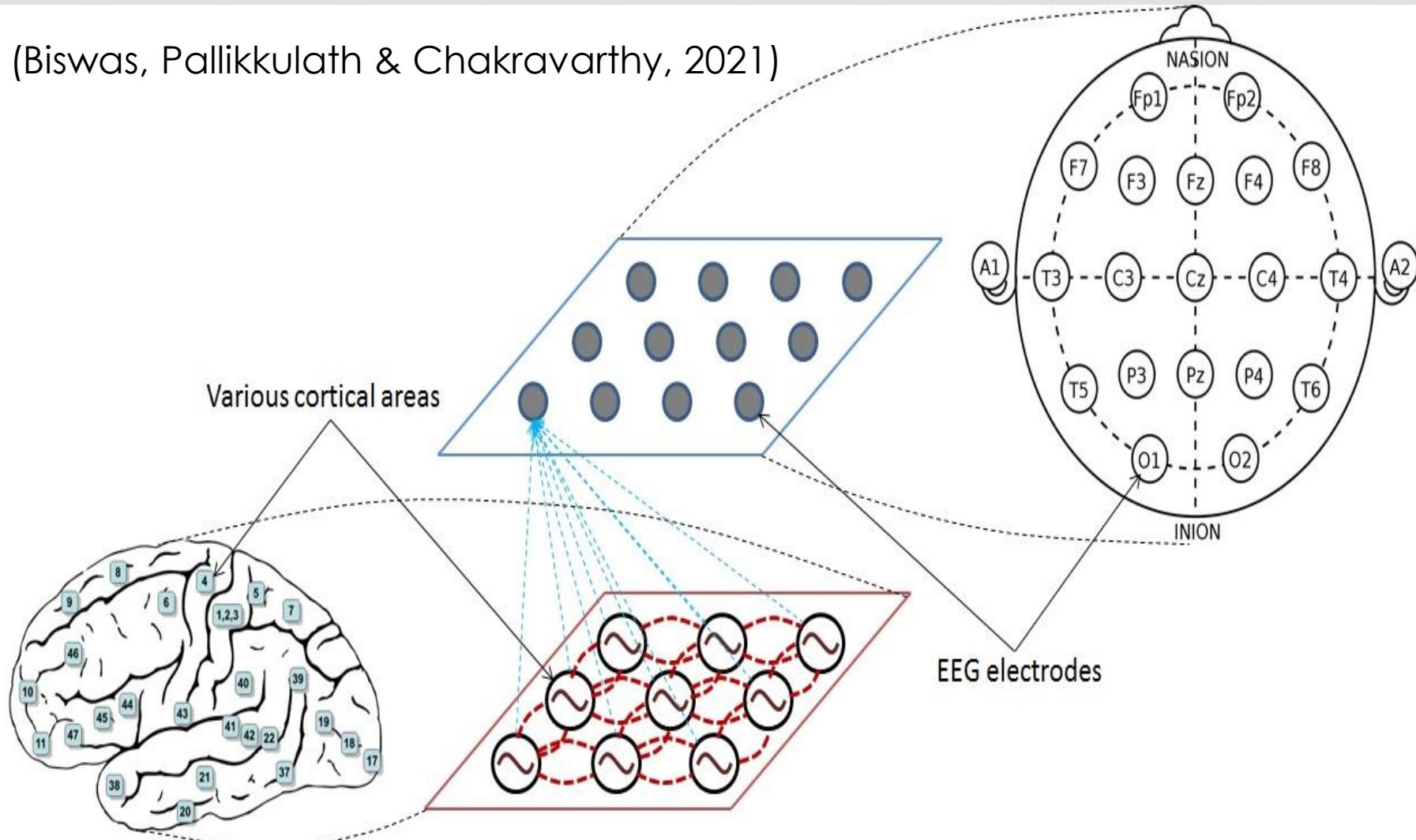
$$\Delta \phi_{ij}$$

$$= (-1) \eta_{\phi_{ij}} \sum_t (Y_{d_i}(t) - Y_{p_i}(t)) (-K_{ij} \sin(\theta_j(t) + \phi_{ij}))$$

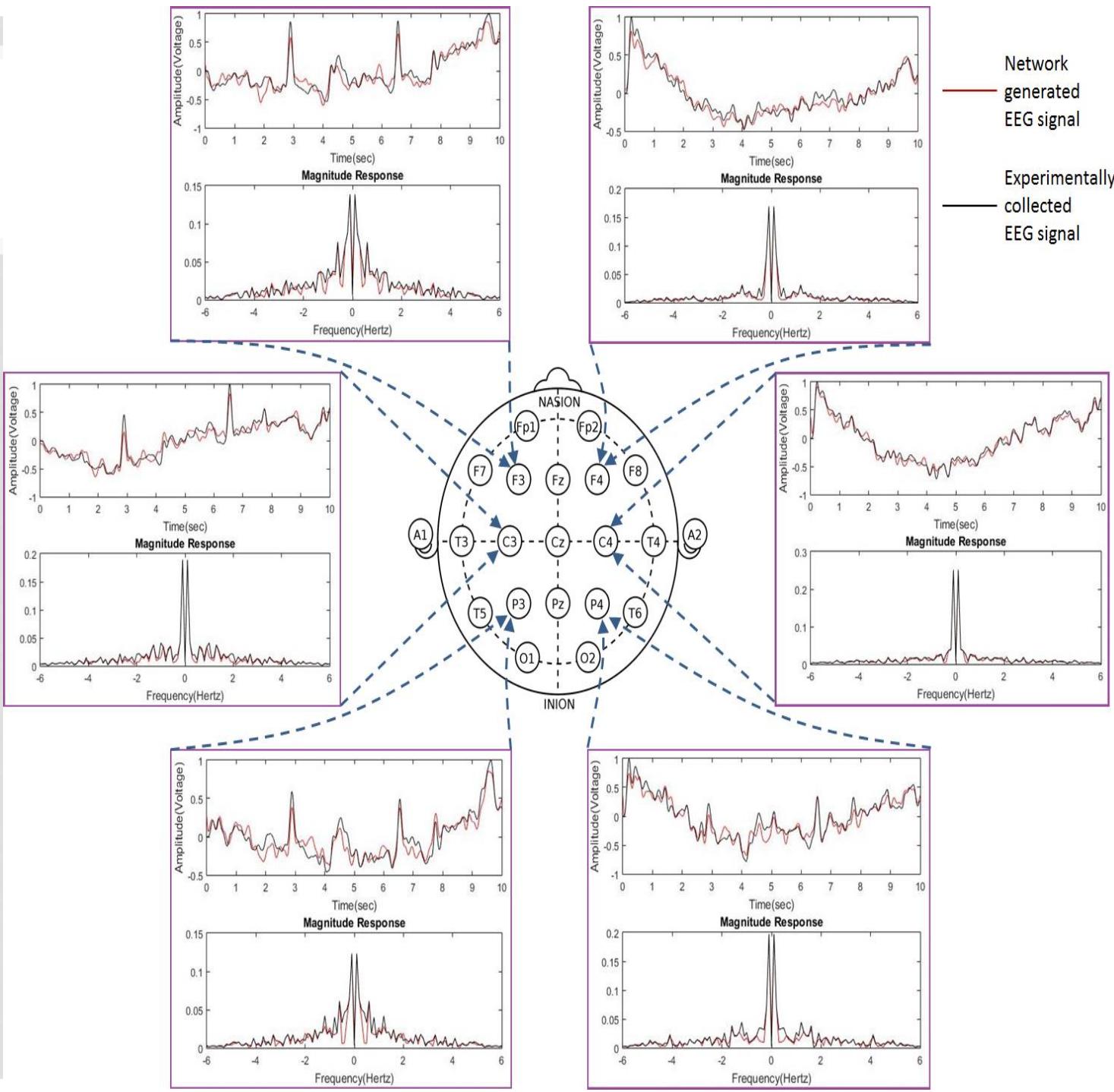


MODELLING ELECTROENCEPHALOGRAPHIC SIGNAL USING NETWORK#2

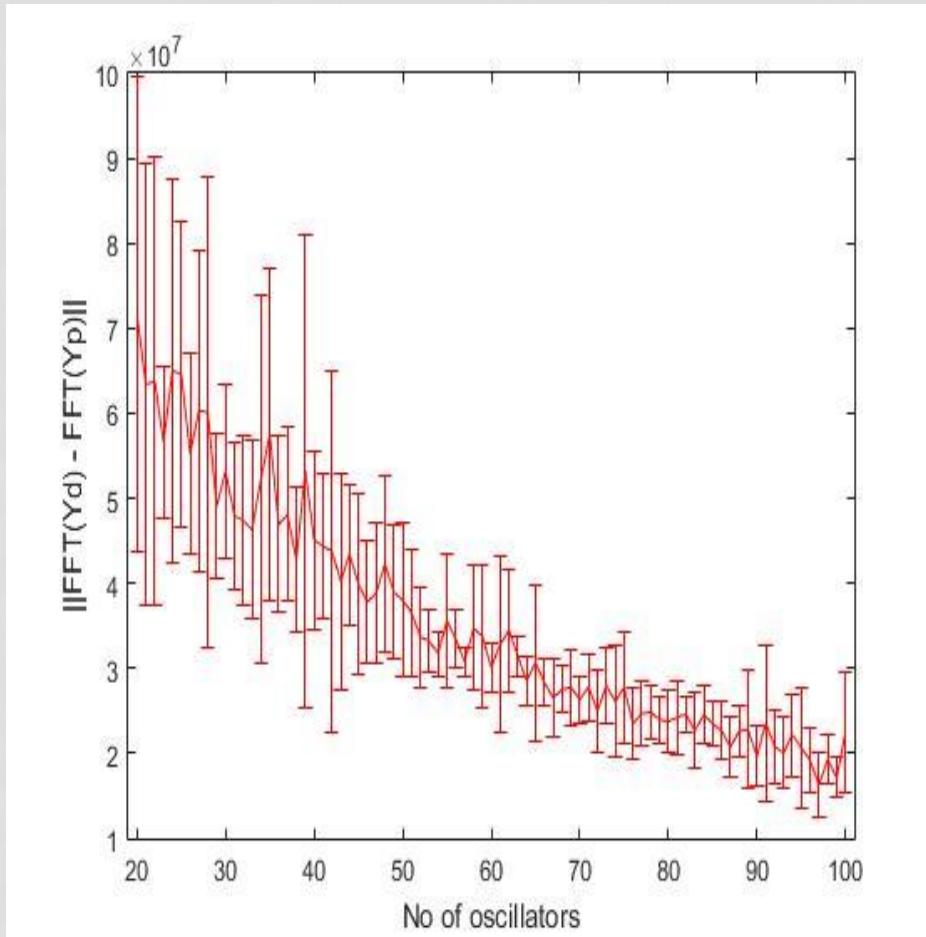
(Biswas, Pallikkulath & Chakravarthy, 2021)



PHASE #1
NETWORK
LEARNS THE
FREQUENCY
SPECTRUM OF A
EEG SIGNAL
COLLECTED
FROM ONE
ELECTRODE AND
IN THE



RECONSTRUCTION ACCURACY INCREASES WITH #OSCILLATORS IN THE NETWORK



FUTURE STEPS

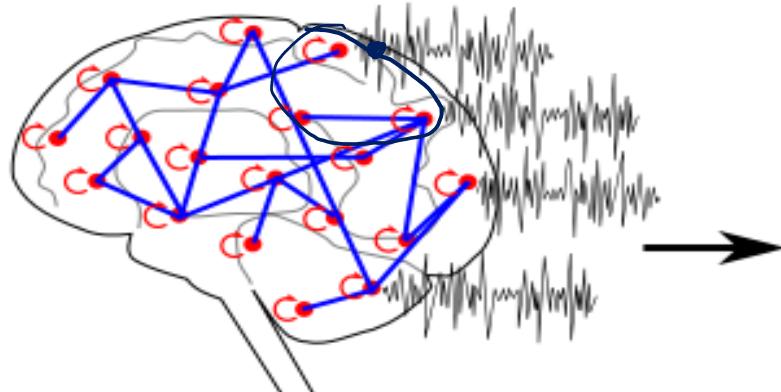
- Large scale Brain models using Oscillatory Deep Neural Networks
- Modeling the Auditory system using Oscillatory Neural Networks
- Generalized Deep Oscillatory Neural Networks

1. LARGE SCALE BRAIN MODELS

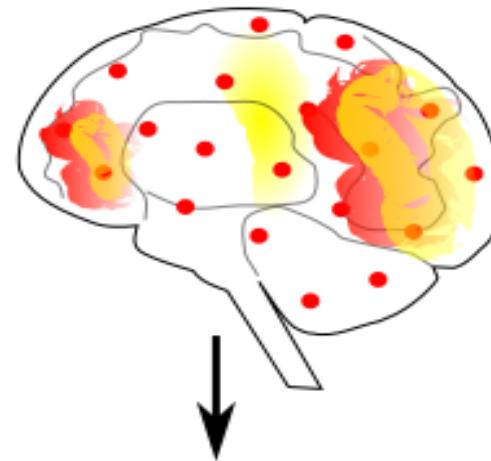
The Virtual Brain (TVB)

- Network of Wilson-Cowan Oscillators
- Connections computed by General optimization

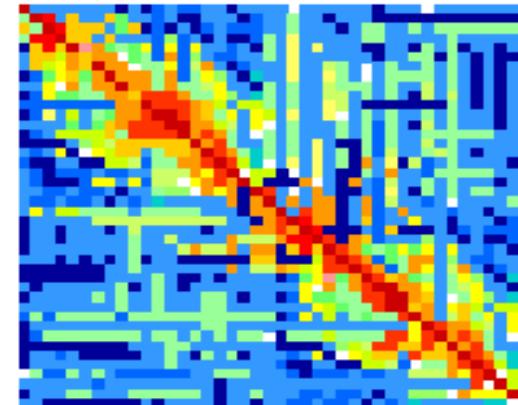
Whole brain simulation
and BOLD signal generation



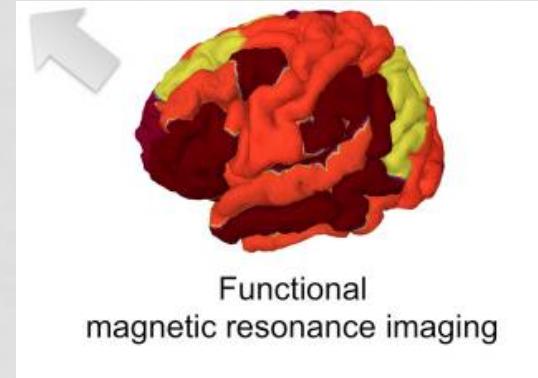
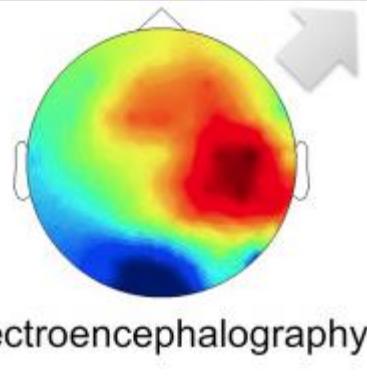
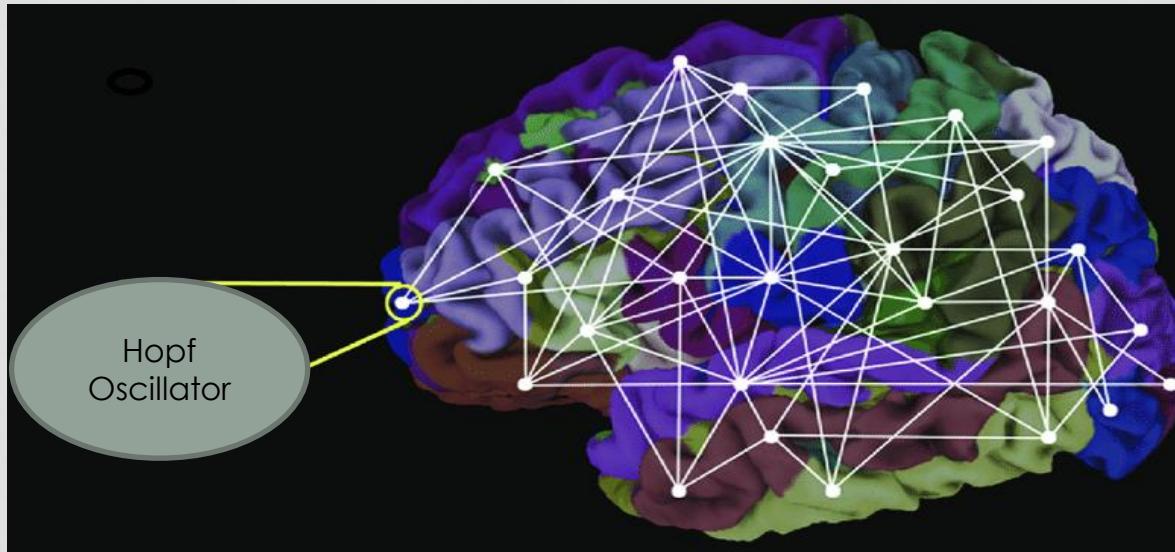
Experimental fMRI data



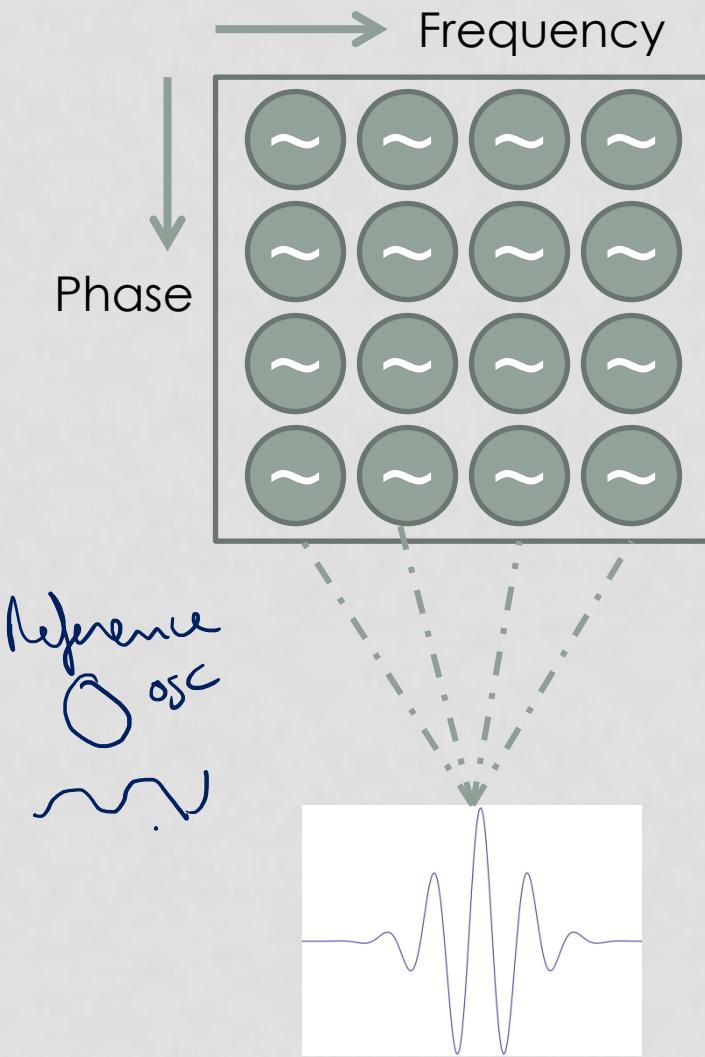
Comparison between simulated
and experimental BOLD signals



LARGE SCALE BRAIN MODELING USING OSCILLATORY NEURAL NETWORK



2. TONOTOPIC MAP MAP OF SOUNDS OR FREQUENCIES

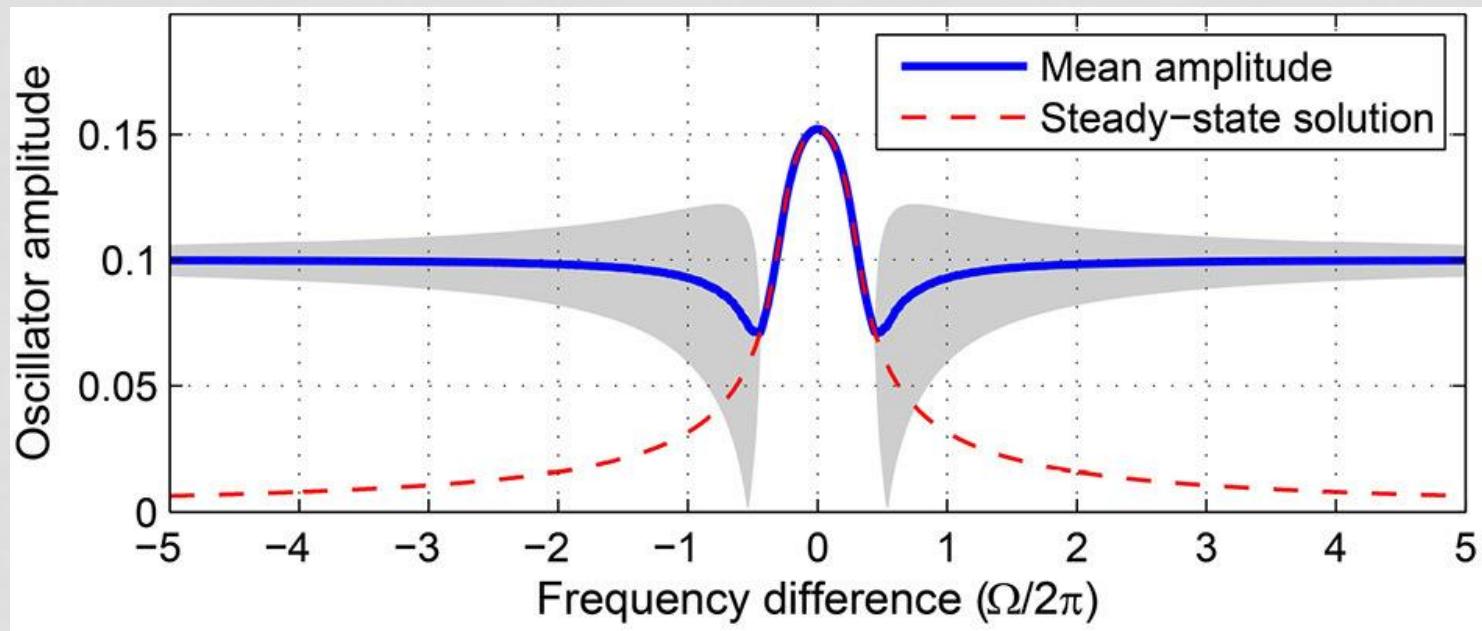


Fourier Series

$$g(t) = a_0 + \sum_{m=1}^{\infty} a_m \cos\left(\frac{2\pi mt}{T}\right) + \sum_{n=1}^{\infty} b_n \sin\left(\frac{2\pi nt}{T}\right)$$

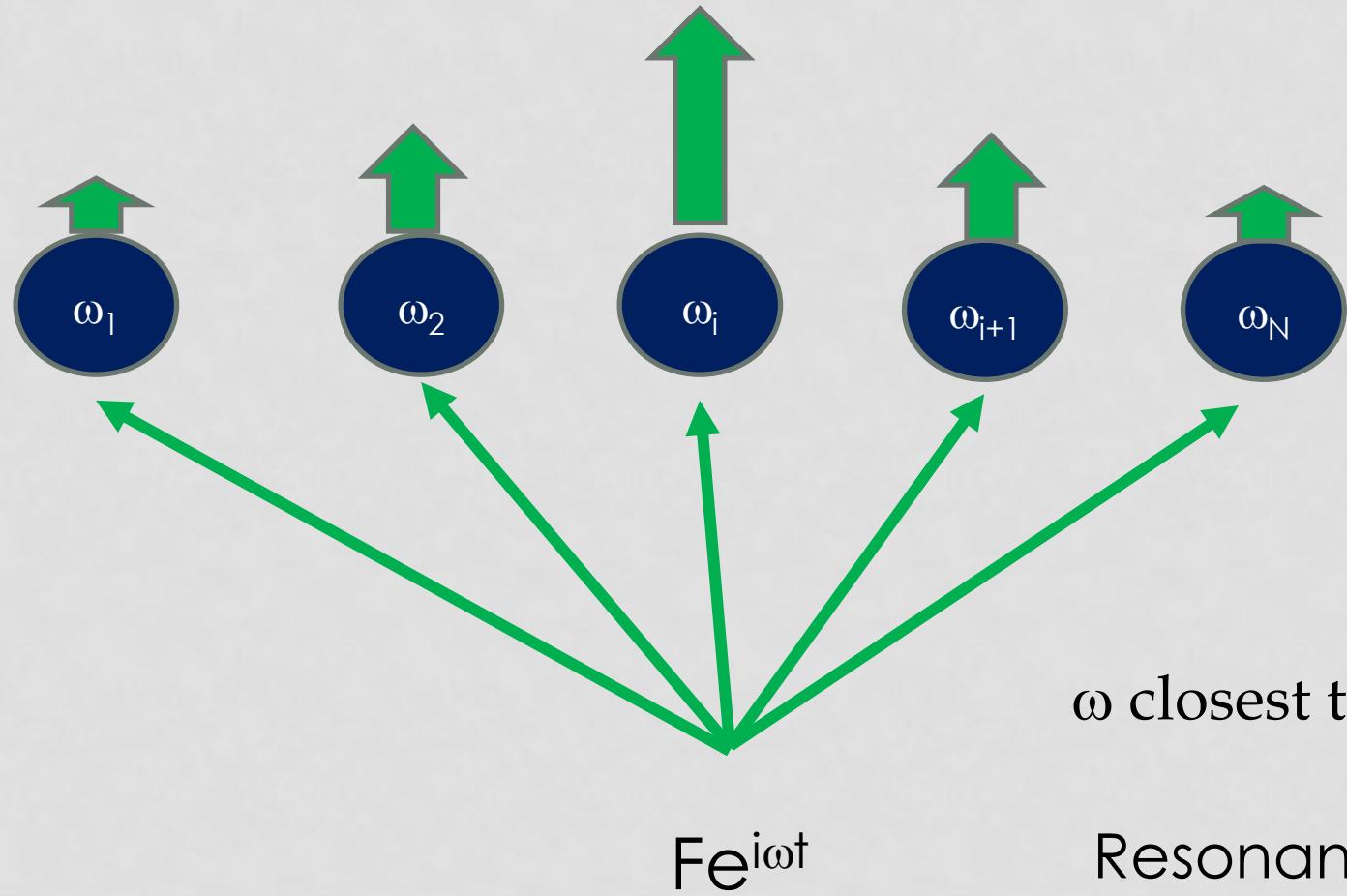
$$= \sum_{m=0}^{\infty} a_m \cos\left(\frac{2\pi mt}{T}\right) + \sum_{n=1}^{\infty} b_n \sin\left(\frac{2\pi nt}{T}\right)$$

HOPF OSCILLATORS EXHIBITS RESONANCE

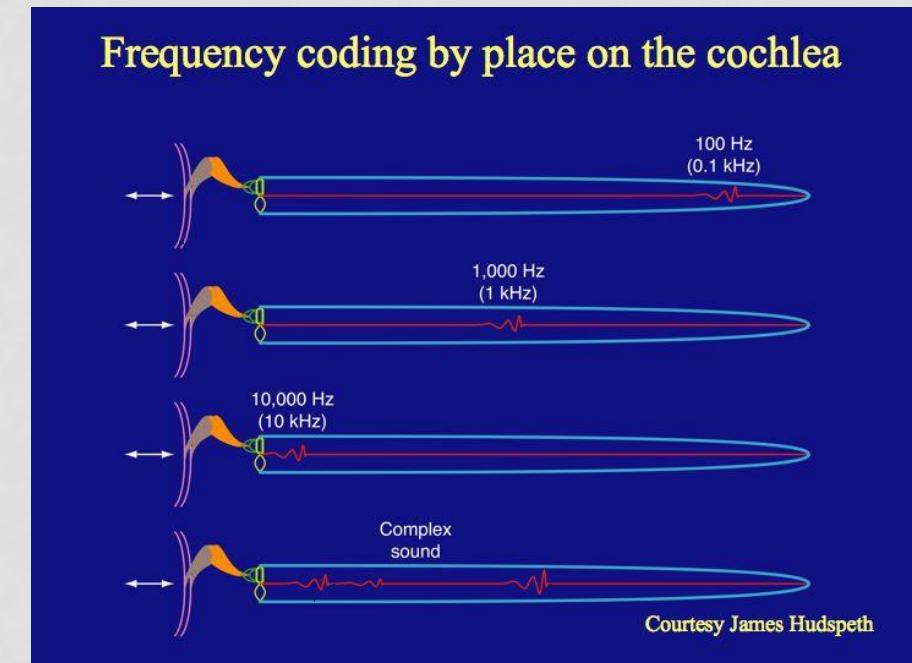
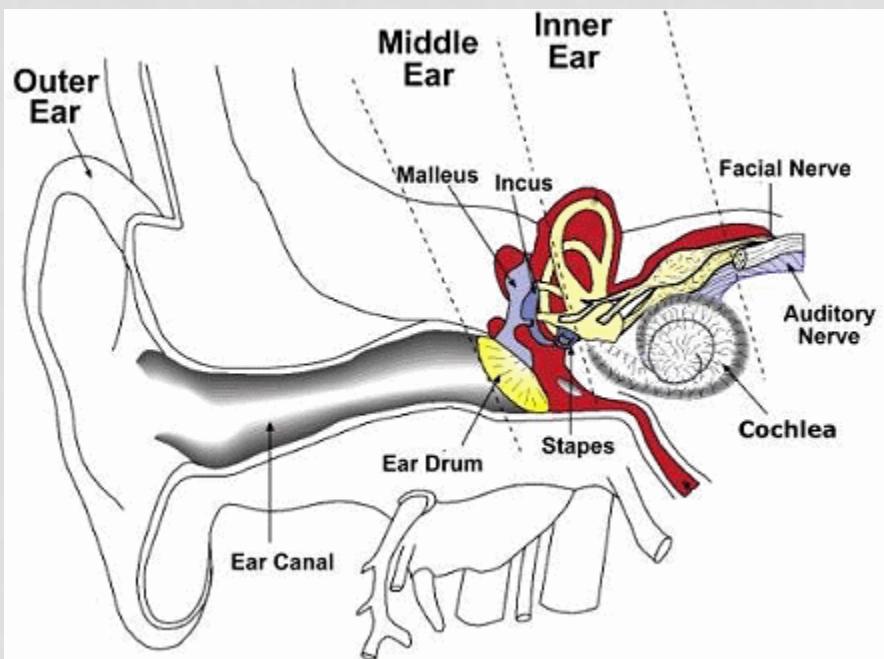


$$\Omega = \omega - \omega_0$$

(Kim & Large,)

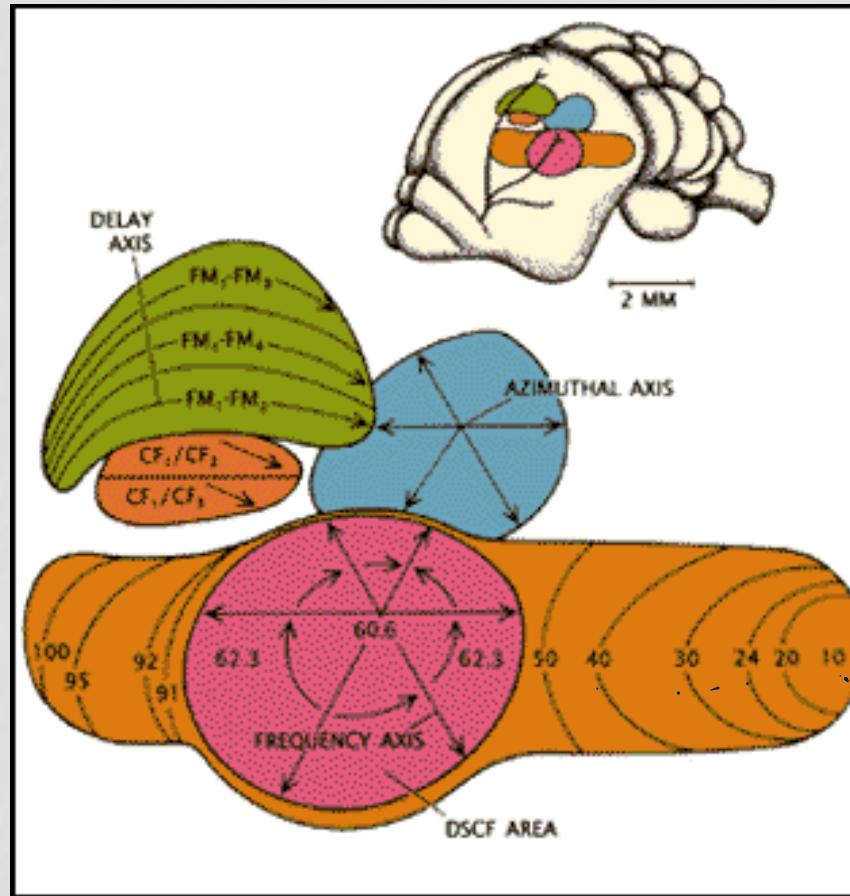
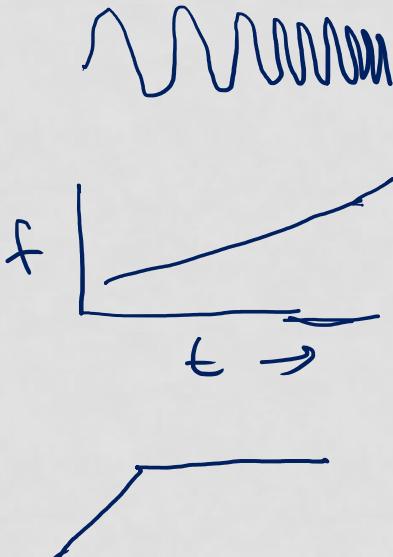


FOURIER DECOMPOSITION IN COCHLEA



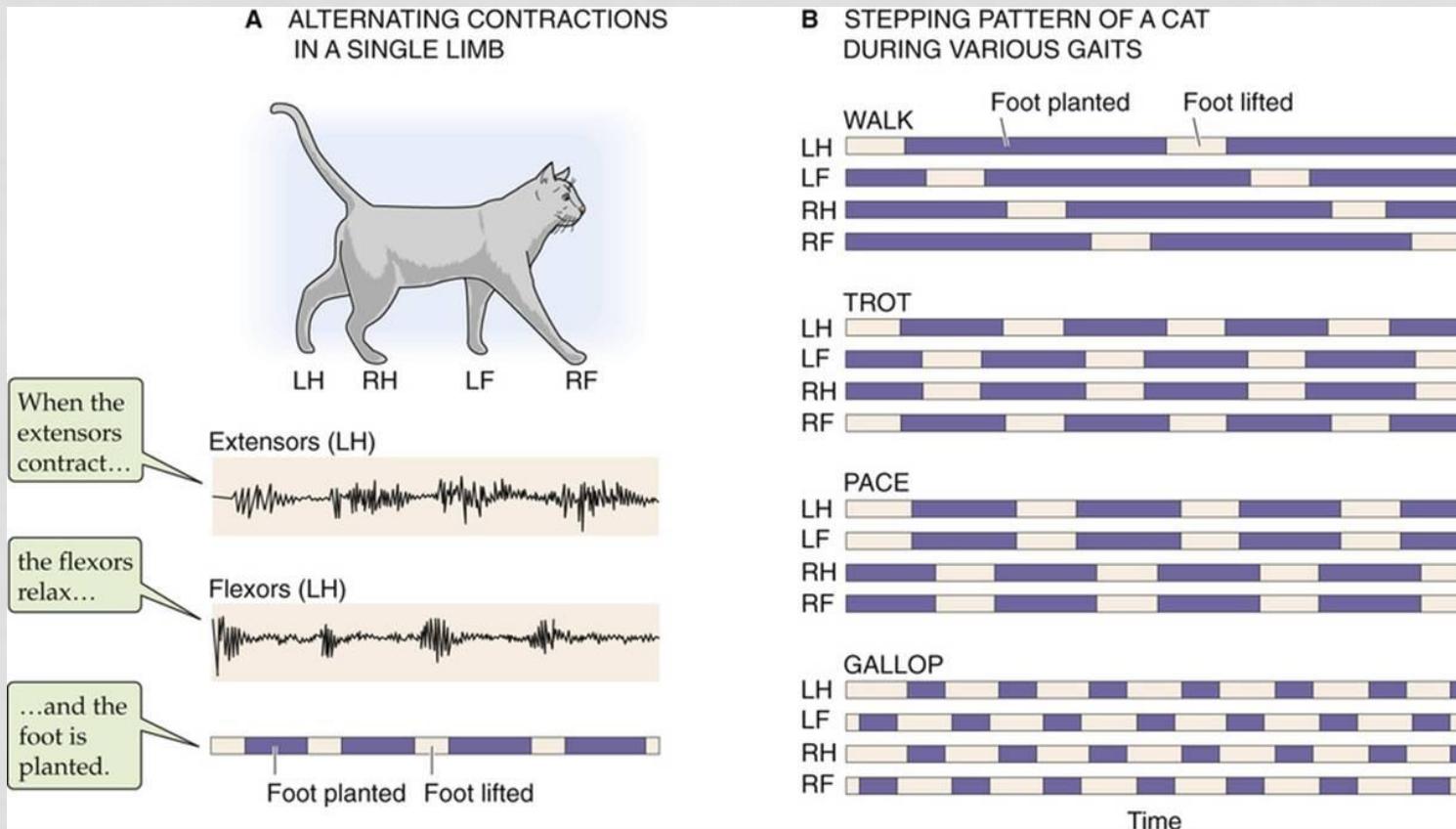
AUDITORY CORTEX OF BATS

chirp signals.

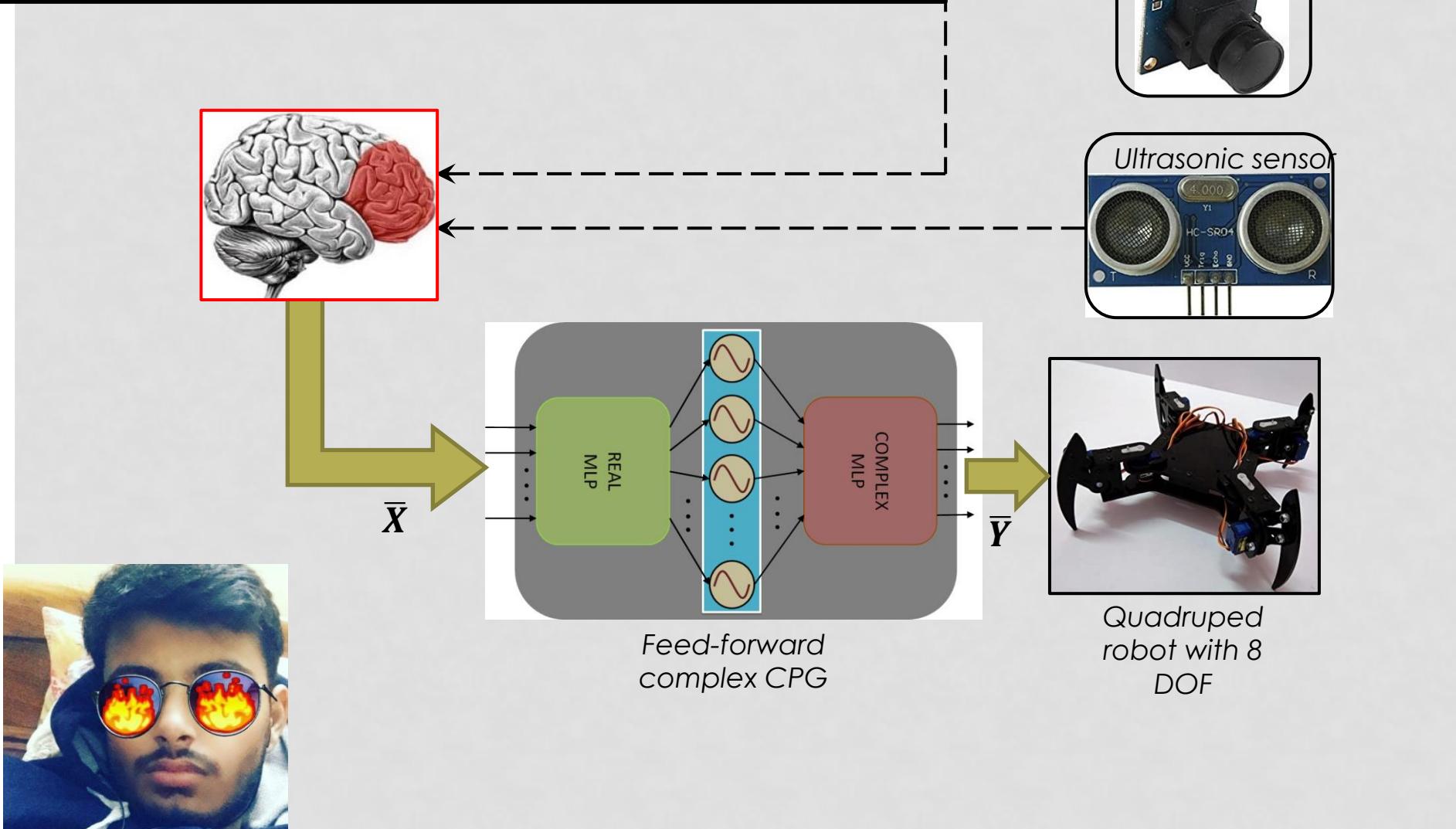


Nobuo
Suga

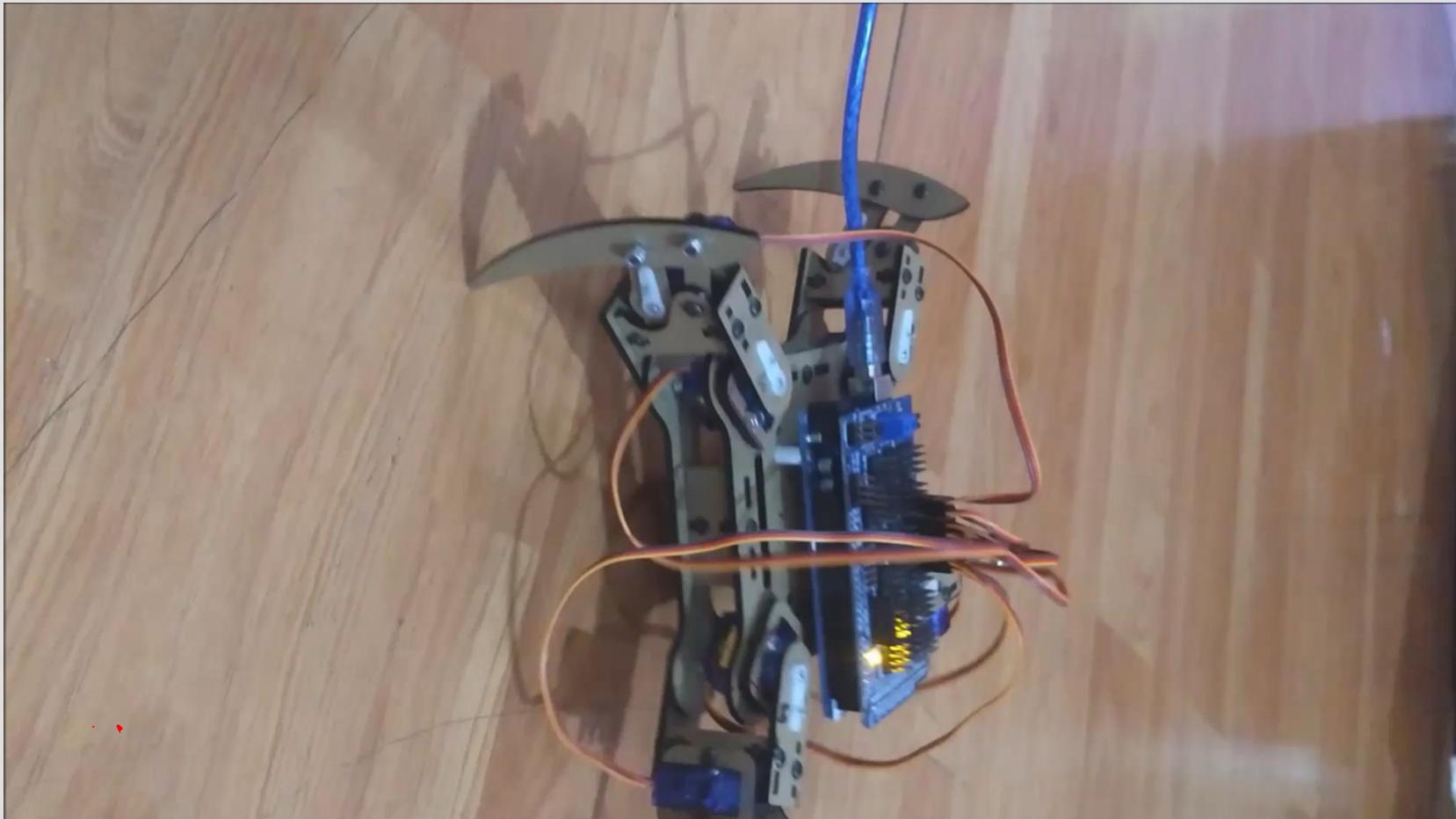
3. MODELING LOCOMOTOR RHYTHMS



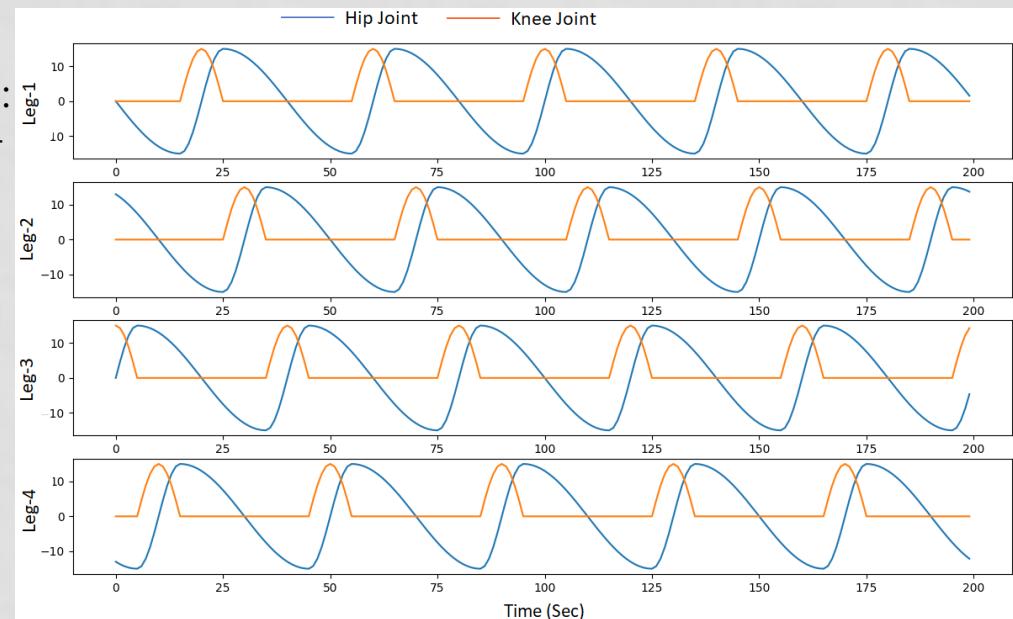
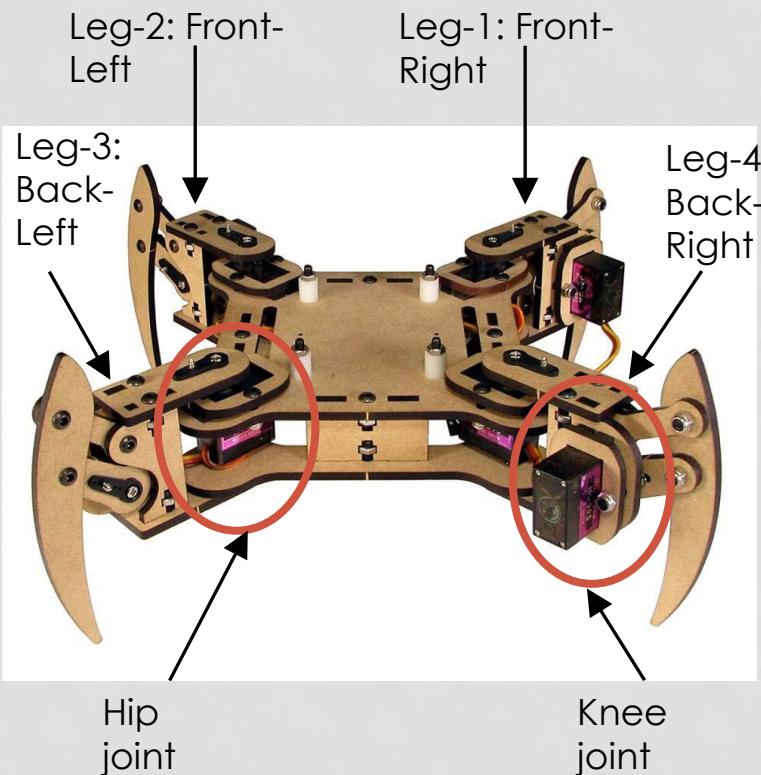
A feed-forward supervised oscillatory CPG model to control robotic locomotion



CURRENT IMPLEMENTATION

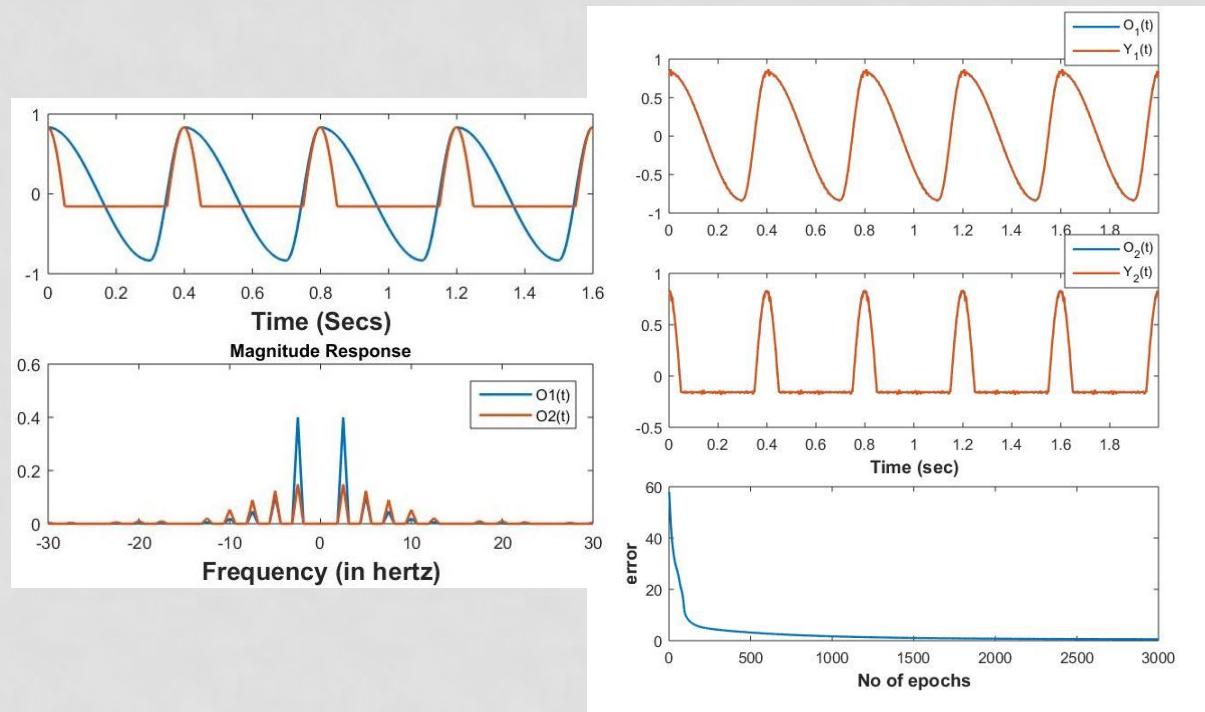


MEPED V2 QUADRUPED ROBOT:

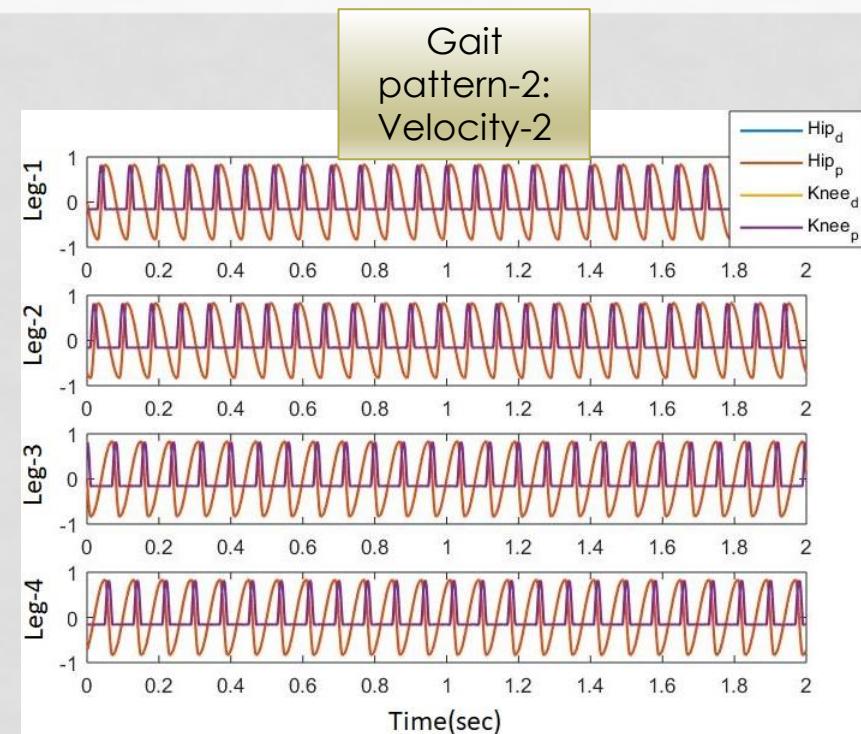
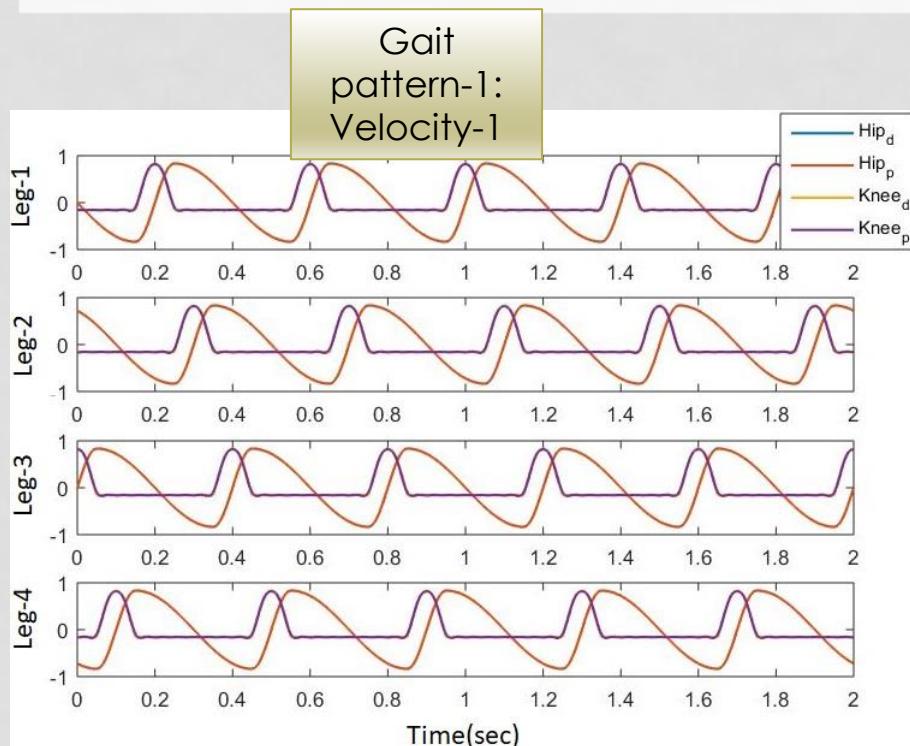


SUPERVISED CPG NETWORK TO PRODUCE GAIT RHYTHMS:

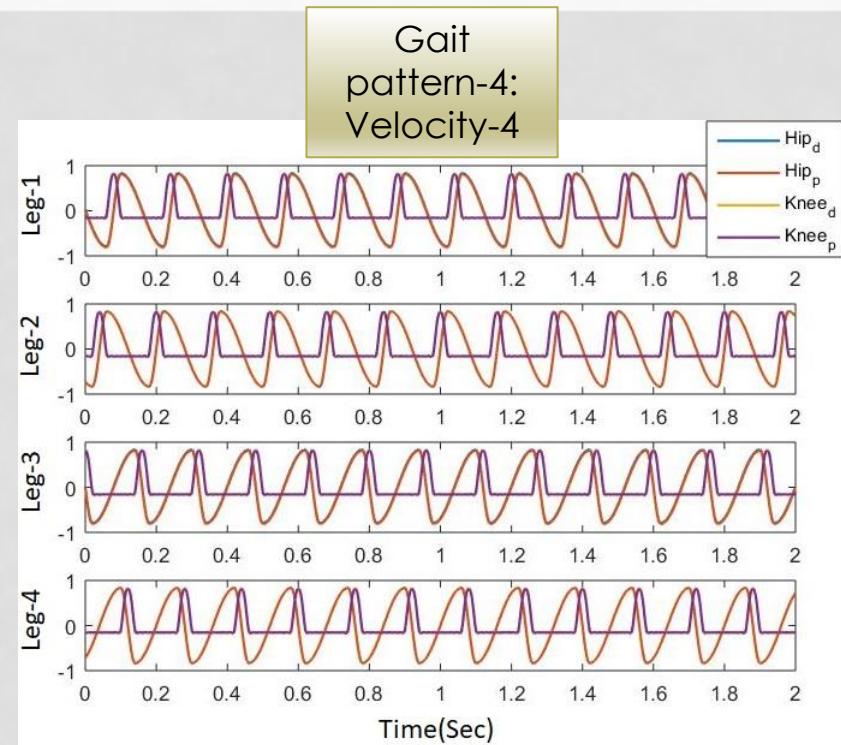
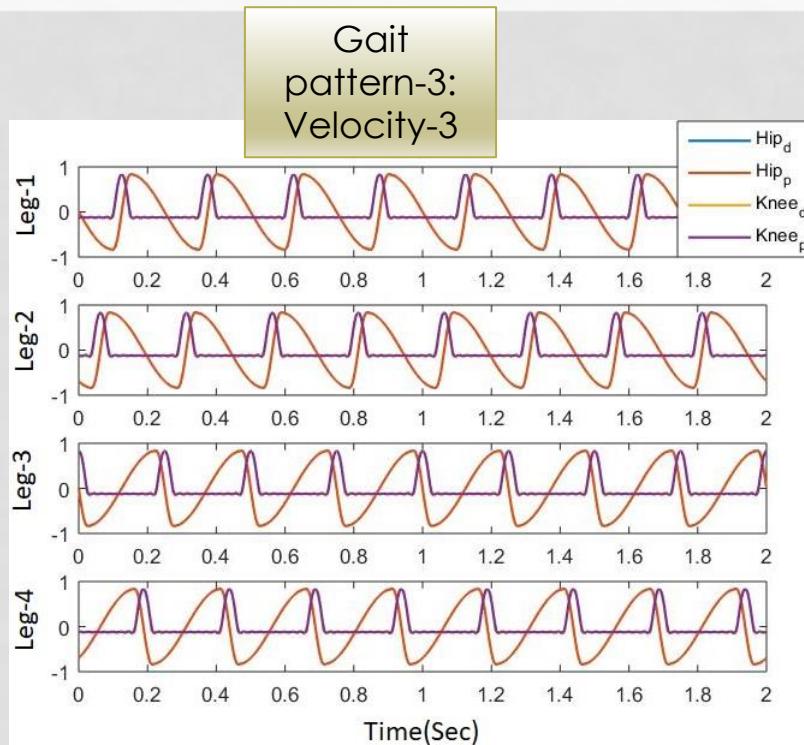
- As the gait rhythms are periodic in nature the constitutive frequency components are harmonics of the fundamental frequency (f_d) of the gait rhythm.
- A perceptron with complex weights projecting the activity of reservoir of oscillators with natural frequencies as harmonics of f_d can learn to produce desired gait pattern with reasonable accuracy.



THE FOUR DIFFERENT GAIT PATTERNS LEARNT BY THE NETWORK:



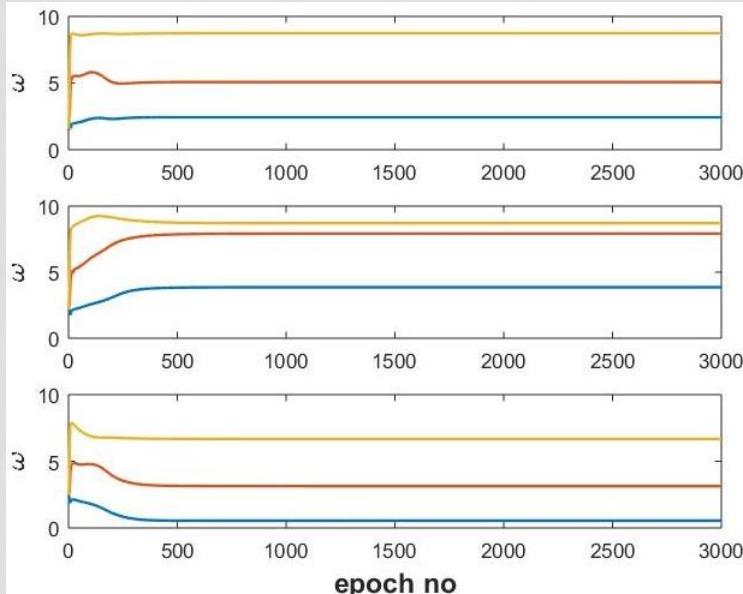
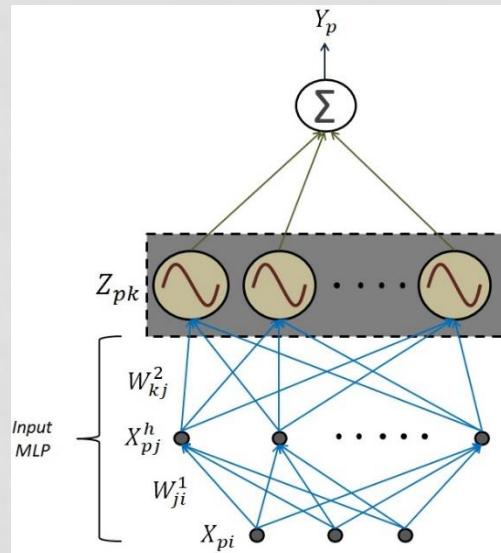
DESIRED AND THE NETWORK LEARNT GAIT PATTERNS:



A FEED-FORWARD HYBRID NETWORK WITH HOPF OSCILLATORS TO LEARN FOURIER DECOMPOSITION OF MULTIPLE SIGNALS

Network-1:

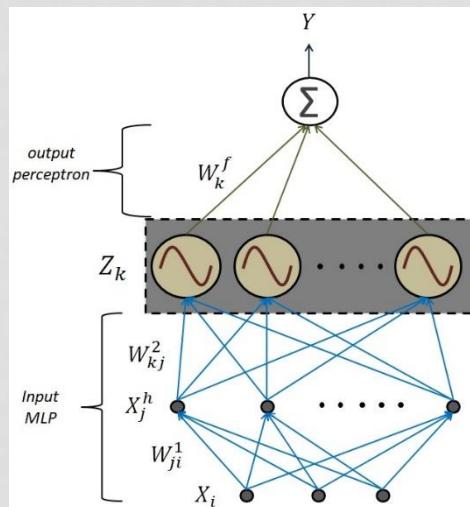
- The output of the “input MLP” is the natural frequencies of the oscillators.
- The parameters of the MLP are supposed to learn the frequency components in the desired output signal.



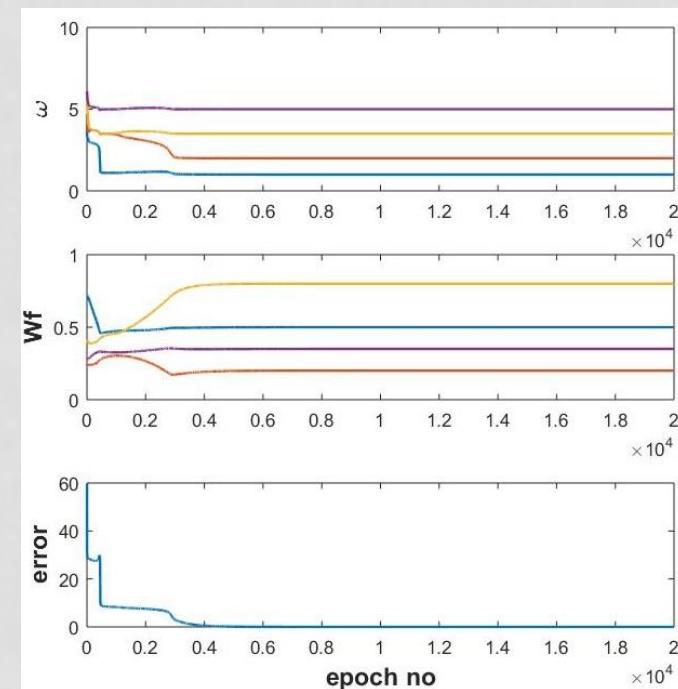
	Desired input			Desired output
1	0.7544	0.8892	0.7104	$\cos(8.7403t) + \cos(5.0725t) + \cos(2.4314t)$
2	0.2465	0.4514	0.9312	$\cos(3.8588t) + \cos(8.7079t) + \cos(7.9213t)$
3	0.6157	0.9312	0.6311	$\cos(6.6784t) + \cos(0.5457t) + \cos(3.1394t)$

NETWORK-2: CAN LEARN FREQUENCY AS WELL AS THE MAGNITUDE OF A OUTPUT SIGNAL

- The real weights of the “output” perceptron learn the magnitude of the various frequency components learnt by the oscillators.

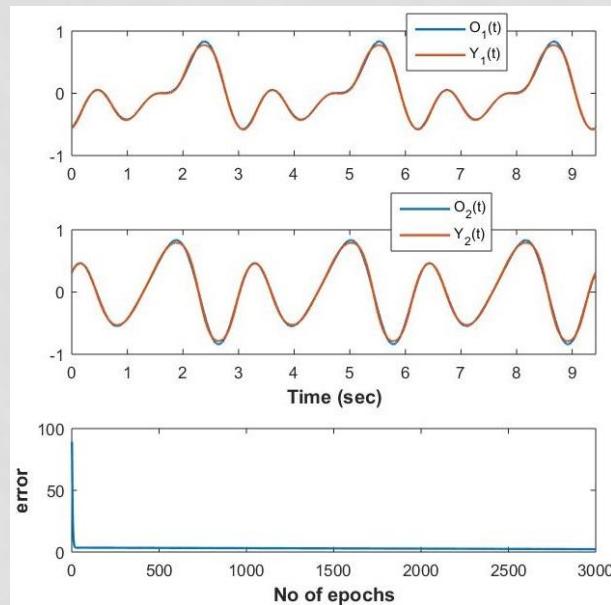
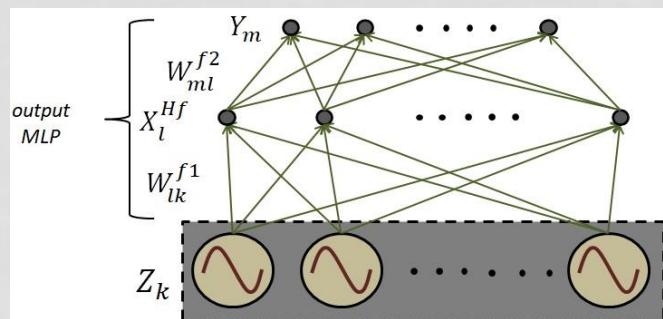


	Desired input			Desired output
1	0.4108	0.7703	0.3400	$0.2\cos(2t) + 0.8\cos(3.5t) + 0.5\cos(t) + 0.5\cos(5t)$



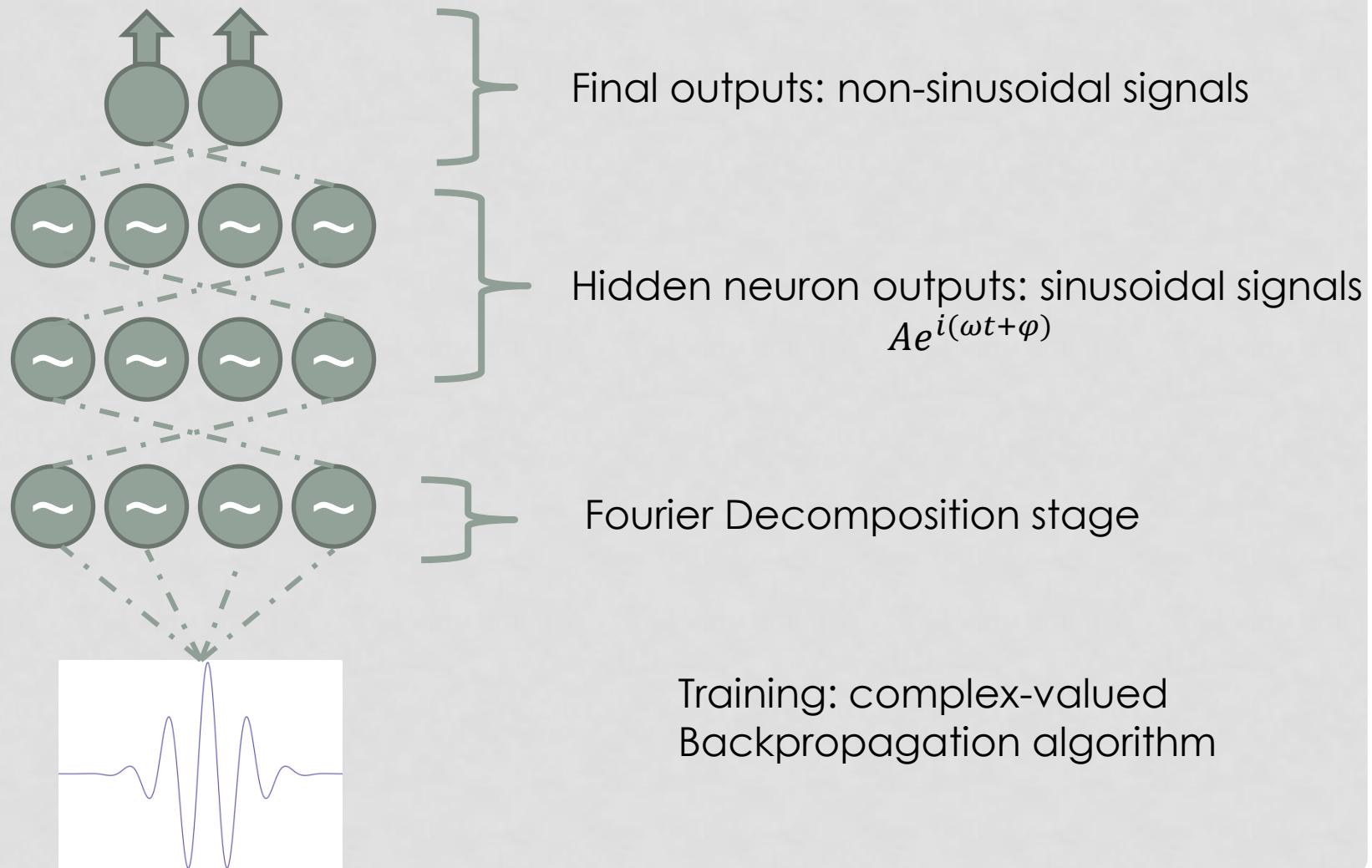
NETWORK-3: CAN LEARN THE MAGNITUDE AND THE PHASE OFFSET OF THE DESIRED OUTPUT SIGNAL

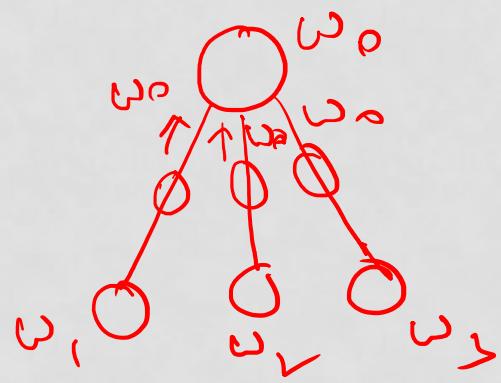
- Given that the constitutive frequency components of the desired output signals at the output nodes are same as the natural frequencies of the Hopf oscillators the network can learn to predict the output by tuning the parameters of the feed-forward complex MLP.



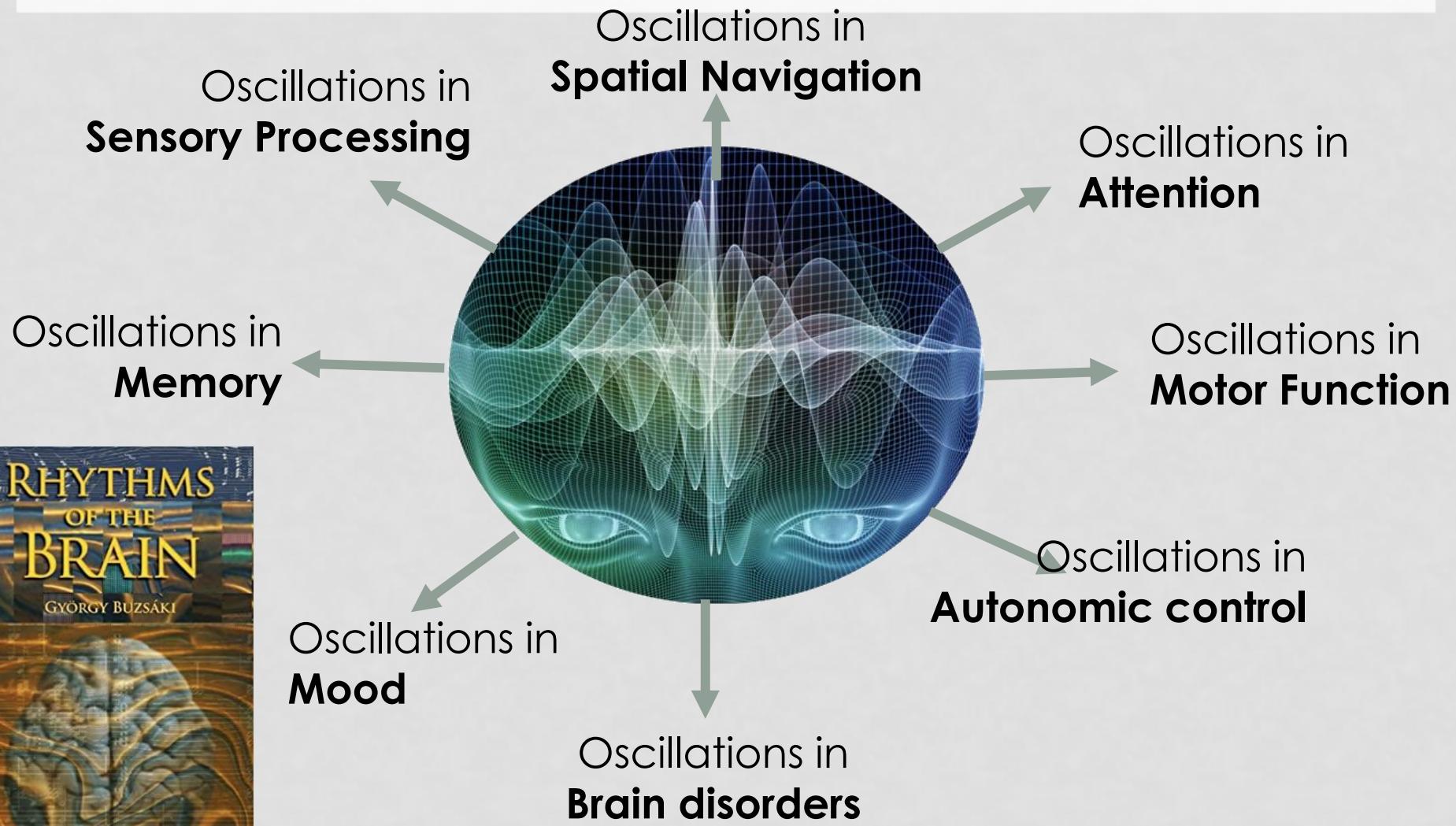
Desired input	Desired output 1	Desired output 2
$Z(t)$ $(\omega = [2; 4; 6])$	$0.9963 \cos(2t + 4.3585)$ + $0.7657 \cos(4t + 5.9176)$ + $0.6521 \cos(6t + 3.9679)$	$0.5754 \cos(2t + 3.8775)$ + $1.2634 \cos(4t + 4.0434)$ + $0.3763 \cos(6t + 0.9521)$

4. GENERALIZED OSCILLATORY DEEP NEURAL NETWORKS





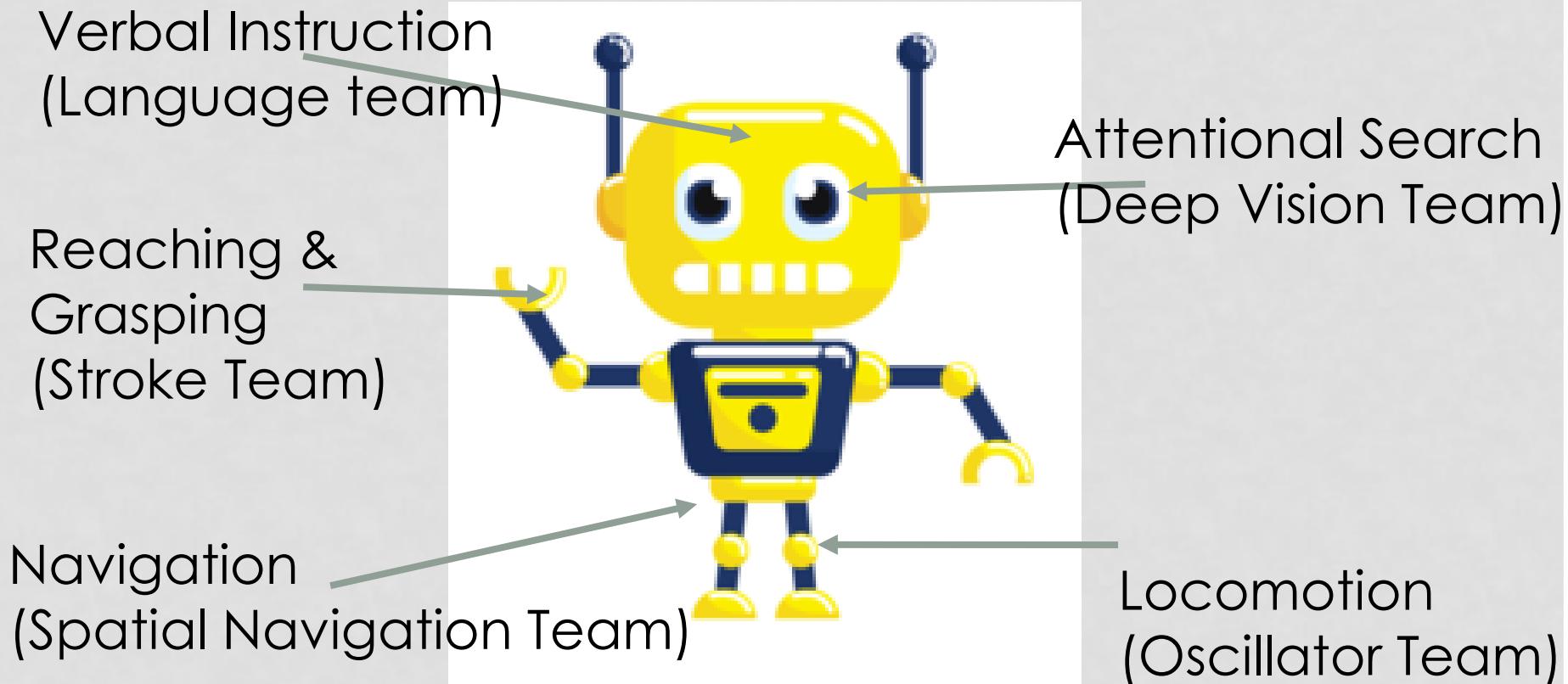
AN OSCILLATOR THEORY OF BRAIN FUNCTION



APPLICATION

HANU: HIERARCHICAL ADAPTIVE NAVIGATING UNIT

A SEARCH-AND-LOCATE ROBOT WITH
A BRAIN-INSPIRED NAVIGATIONAL SYSTEM



ACKNOWLEDGEMENTS

- Dipayan Biswas
- Aasit
- Surya Kiran
- Shreyas
- Ankit

thank
you