

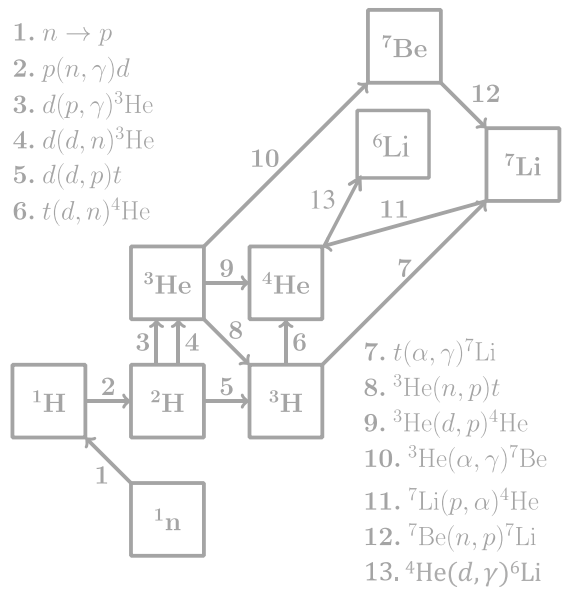
Ab initio description of nuclear reactions with applications to astrophysics

Sofia Quaglioni

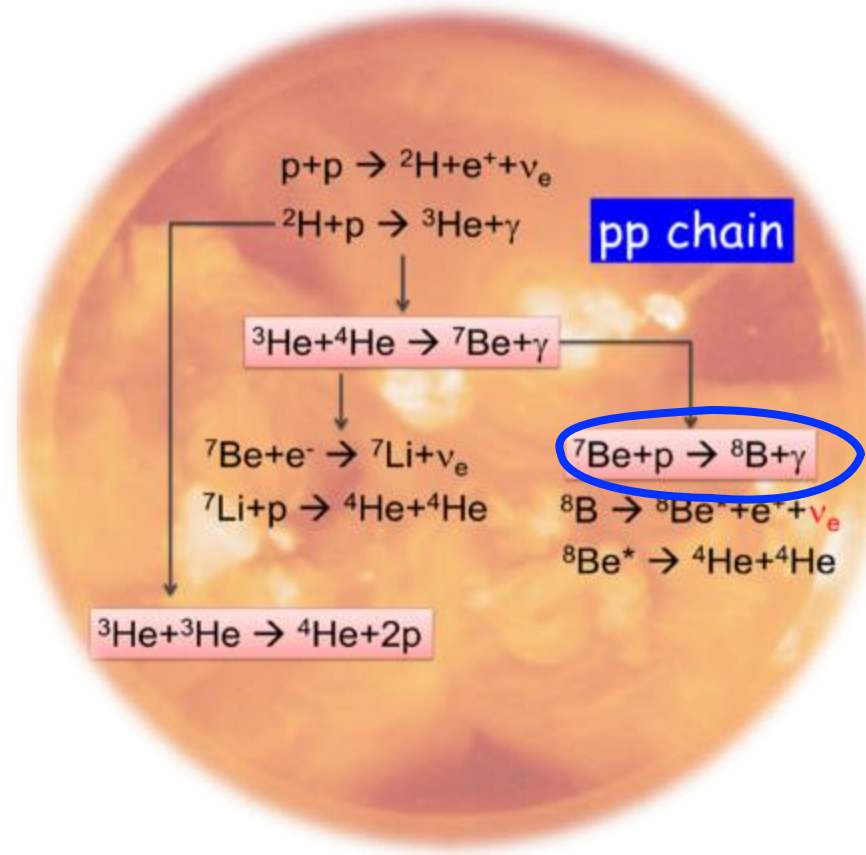


Reactions 'R' Us

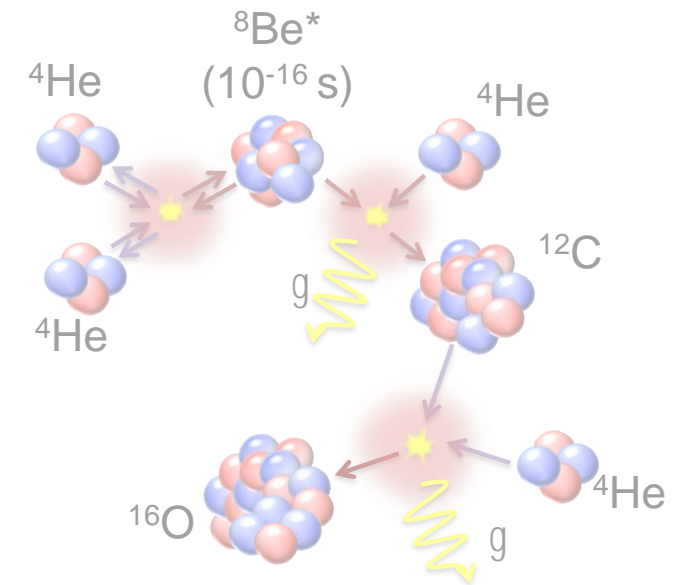
Big Bang



Solar Fusion



Helium Burning



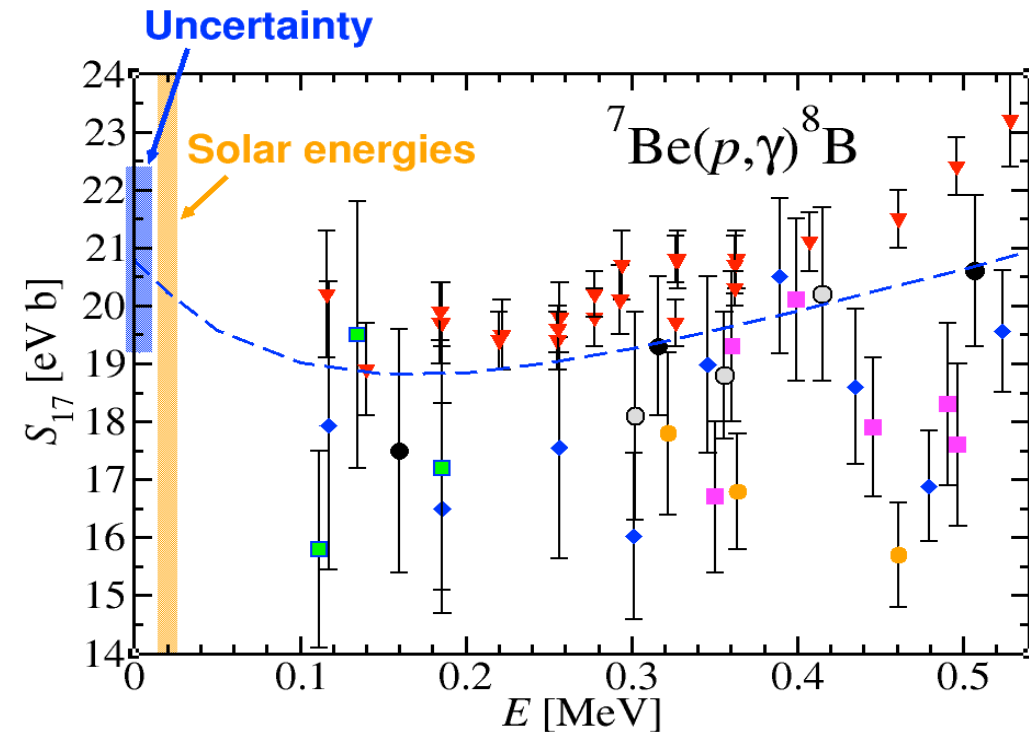
We need reliable theory to accurately evaluate S-factors at stellar energies

Astrophysical S-factor:
nuclear contribution

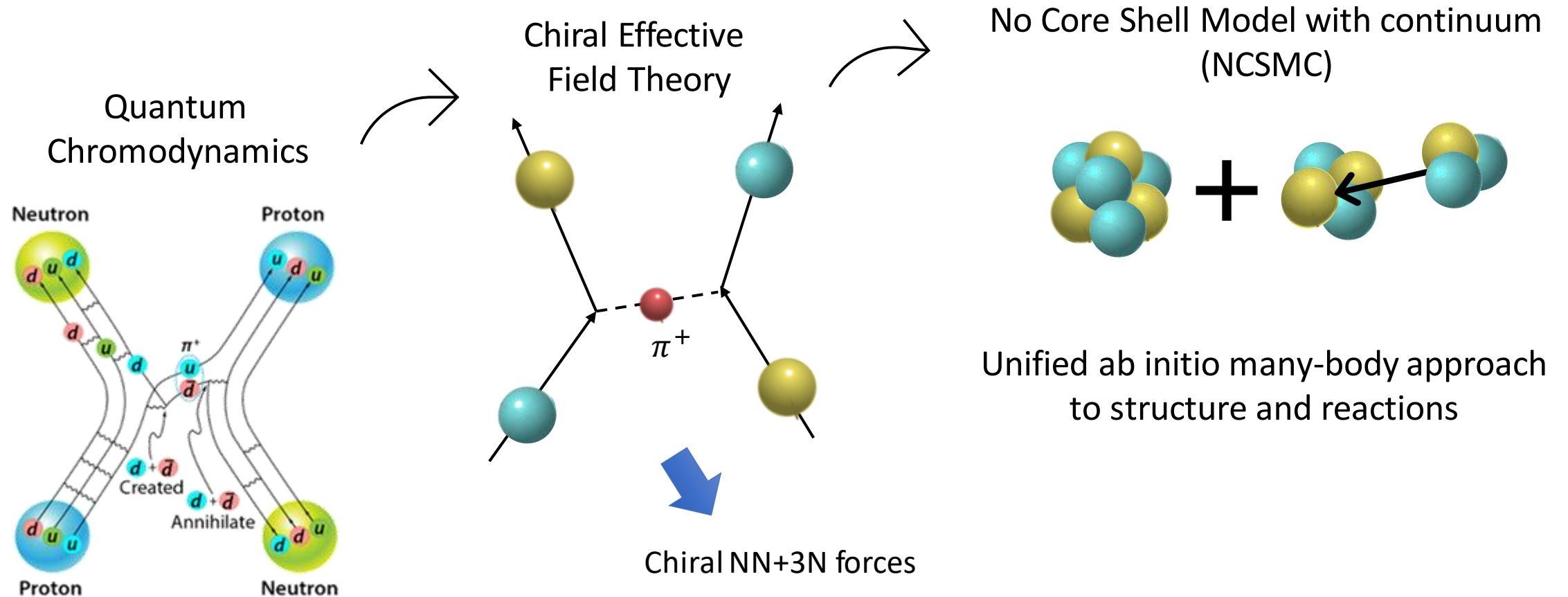
$$\sigma(E) = \frac{S(E)}{E} \exp\left(-\frac{2\pi Z_1 Z_2 e^2}{\hbar \sqrt{2E/m}}\right)$$

‘Coulomb’ contribution
(tunneling)

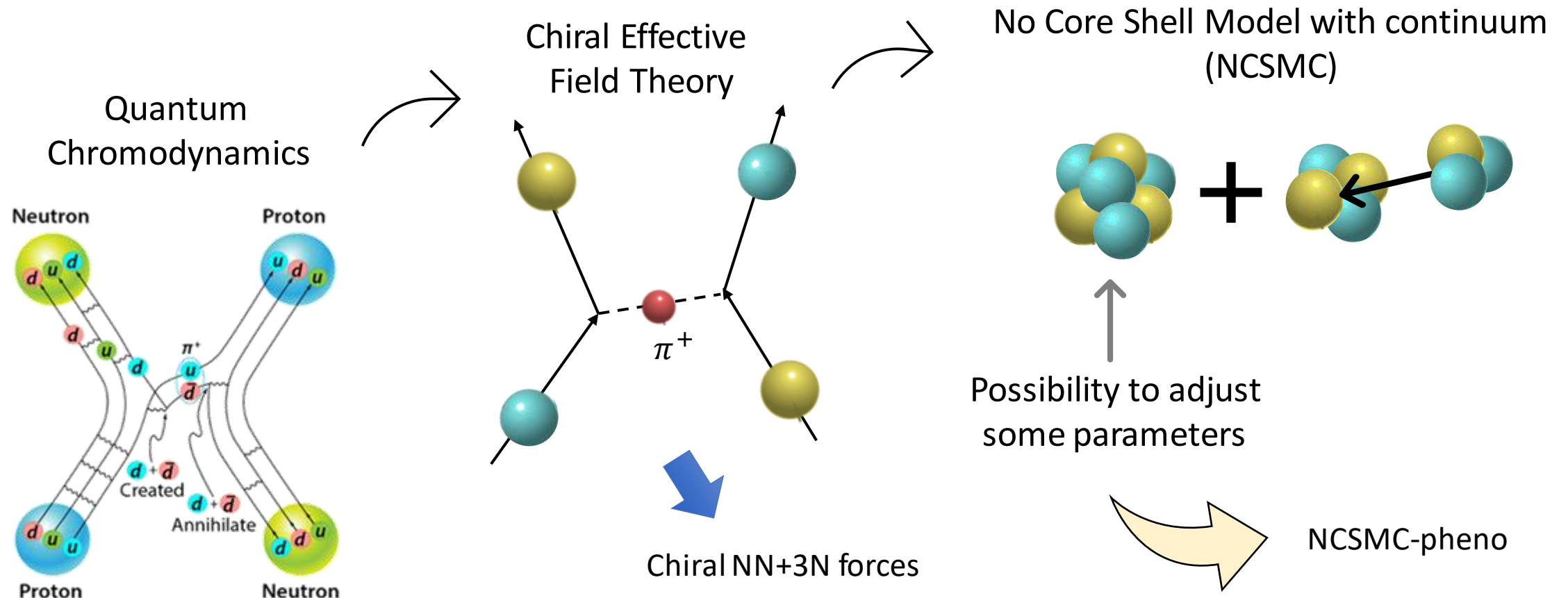
$$S_{17}(0) = 20.8 \pm (0.7)_{\text{exp}} \pm (1.4)_{\text{th}} \text{ eV}\cdot\text{b}$$



We combine nuclear forces derived within chiral effective field theory with ab initio methods

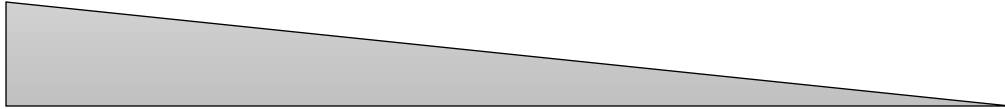


We combine nuclear forces derived within chiral effective field theory with ab initio methods



Can use multiple chiral models, truncation errors to quantify uncertainties

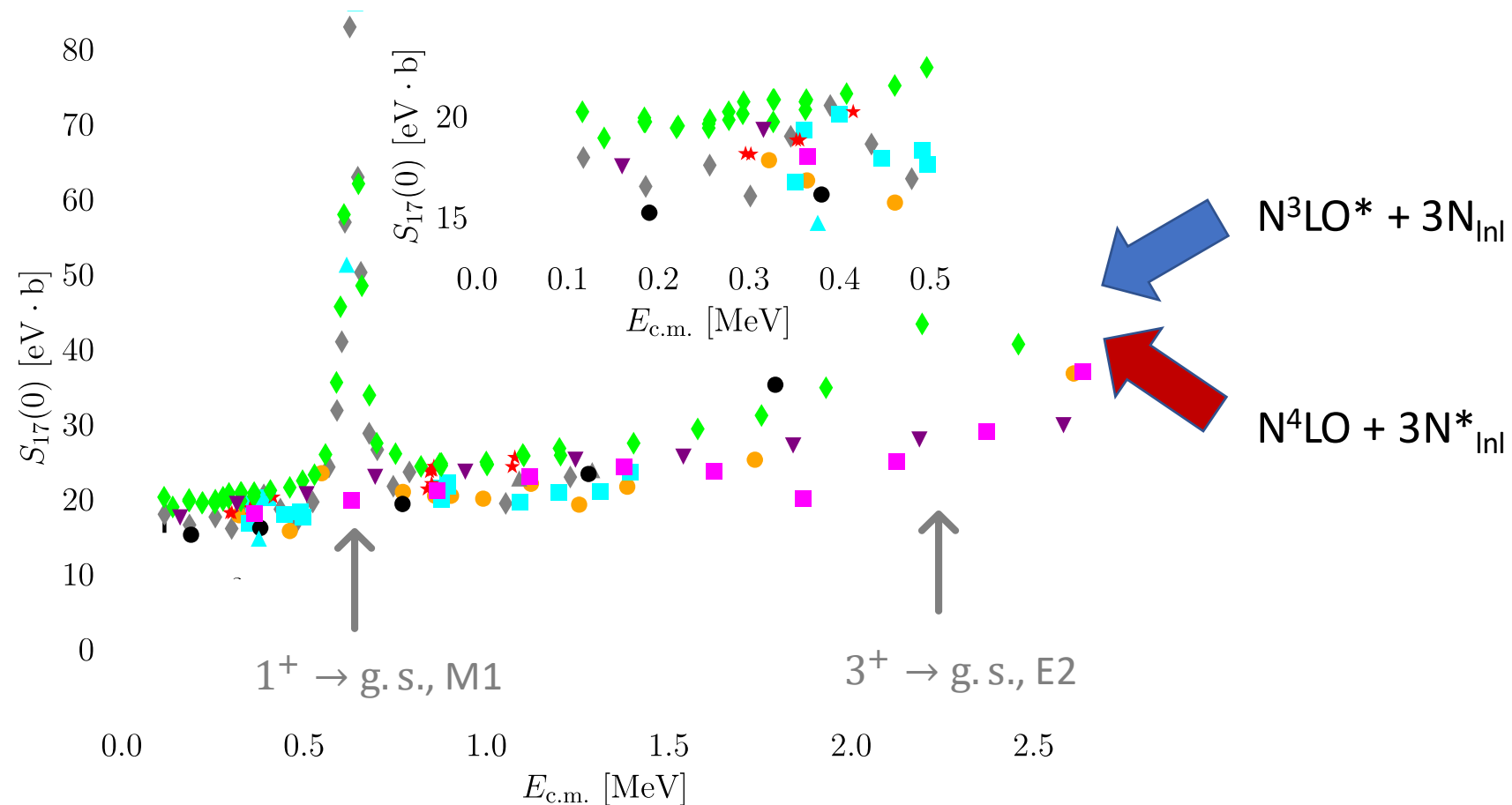
	EMN (2017)				EM (2003)	
	$N^2\text{LO} + 3N_{\text{Inl}}$	$N^3\text{LO} + 3N_{\text{Inl}}$	$N^4\text{LO} + 3N_{\text{Inl}}$	$N^4\text{LO} + 3N_{\text{Inl}}^*$	$N^3\text{LO}^* + 3N_{\text{Inl}}$	$N^3\text{LO}^* + 3N_{\text{loc}}$
NN	$N^2\text{LO}$	$N^3\text{LO}$	$N^4\text{LO}$	$N^4\text{LO}$	$N^3\text{LO}$	$N^3\text{LO}$
3N	$N^2\text{LO}$	$N^2\text{LO}$	$N^2\text{LO}$	$N^2\text{LO} + O(N^3\text{LO})$	$N^2\text{LO}$	$N^2\text{LO (local)}$



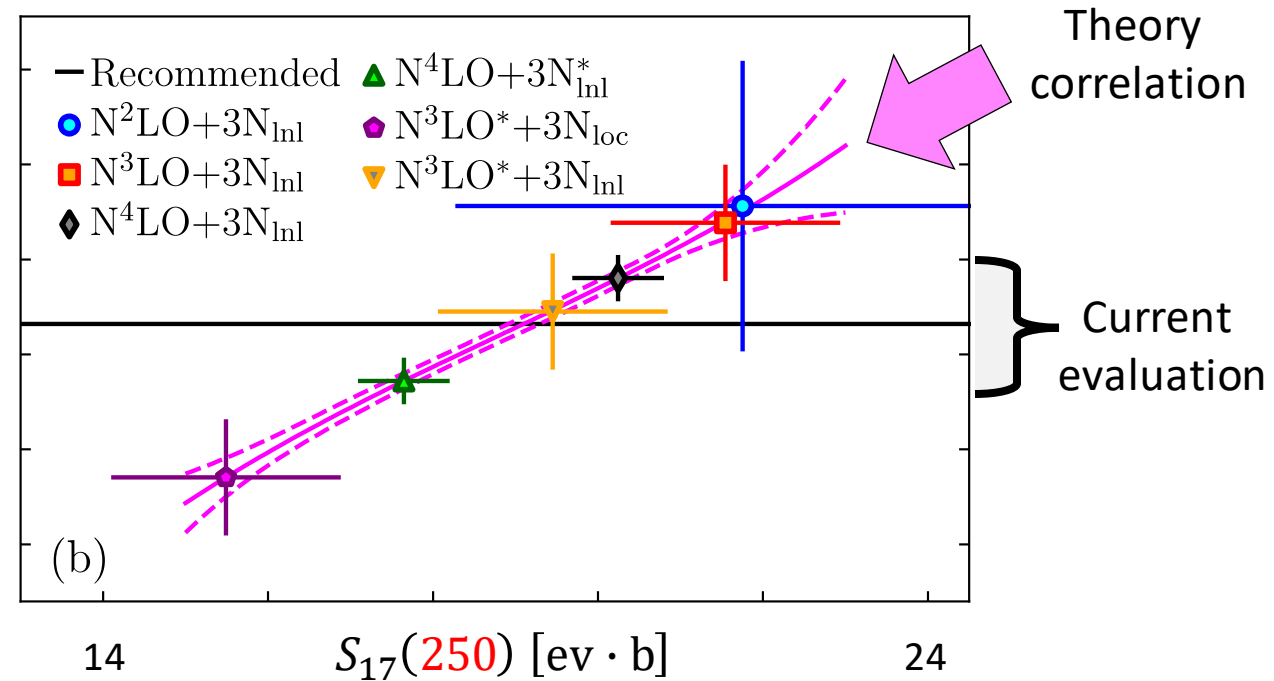
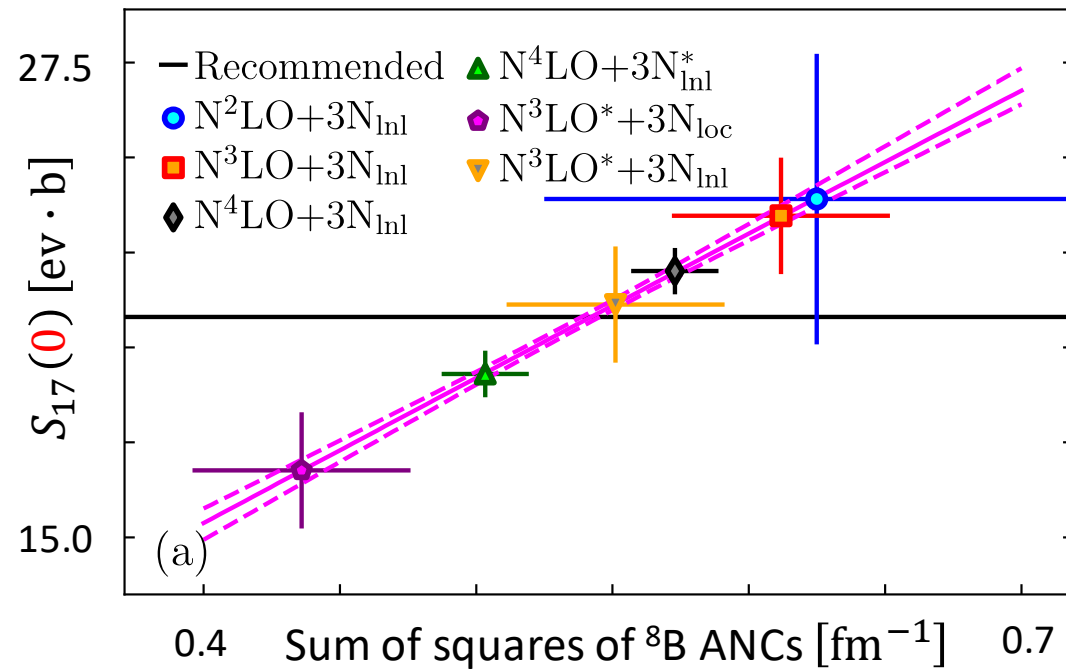
Chiral truncation uncertainty

$\left(\frac{Q}{\Lambda_\chi}\right)^{n+1}$

Chiral NN+3N forces describe expt. ${}^7\text{Be}(p,\gamma){}^8\text{B}$ S-factor with varying success for energies up to 2.5 MeV

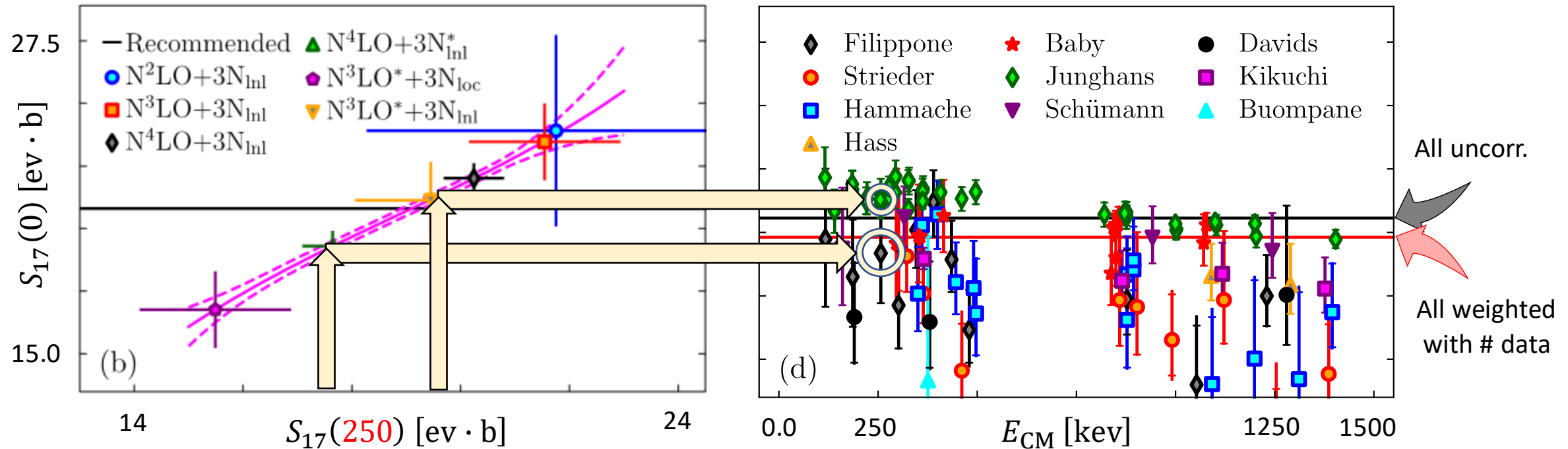


Can extract universal correlation functions
leveraging calculations with different interactions ...



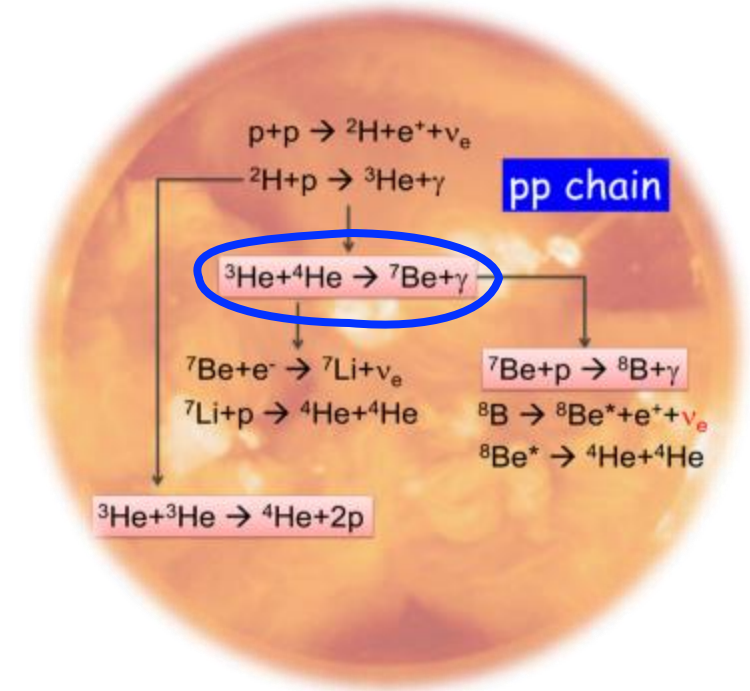
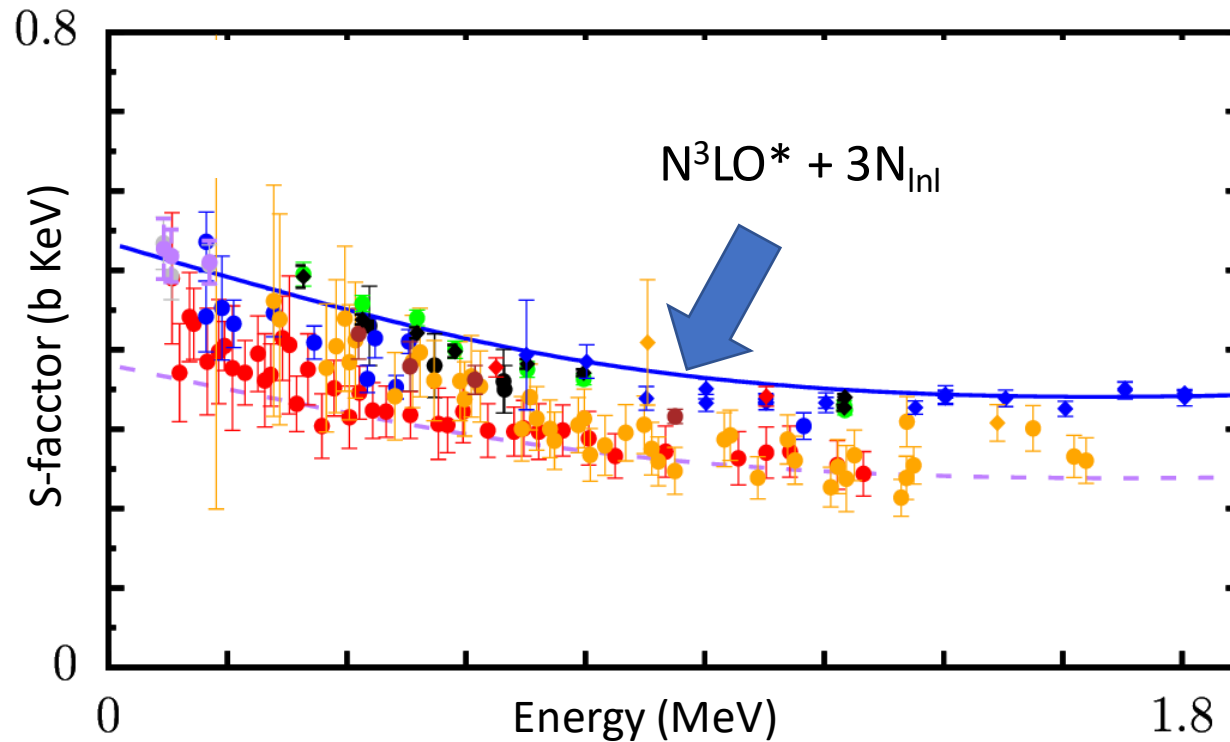
... combine them with experimental data
to arrive at an improved evaluation of $S_{17}(0)$

$$S_{17}(0) = 19.8 \pm 0.3 \text{ eV}\cdot\text{b}$$

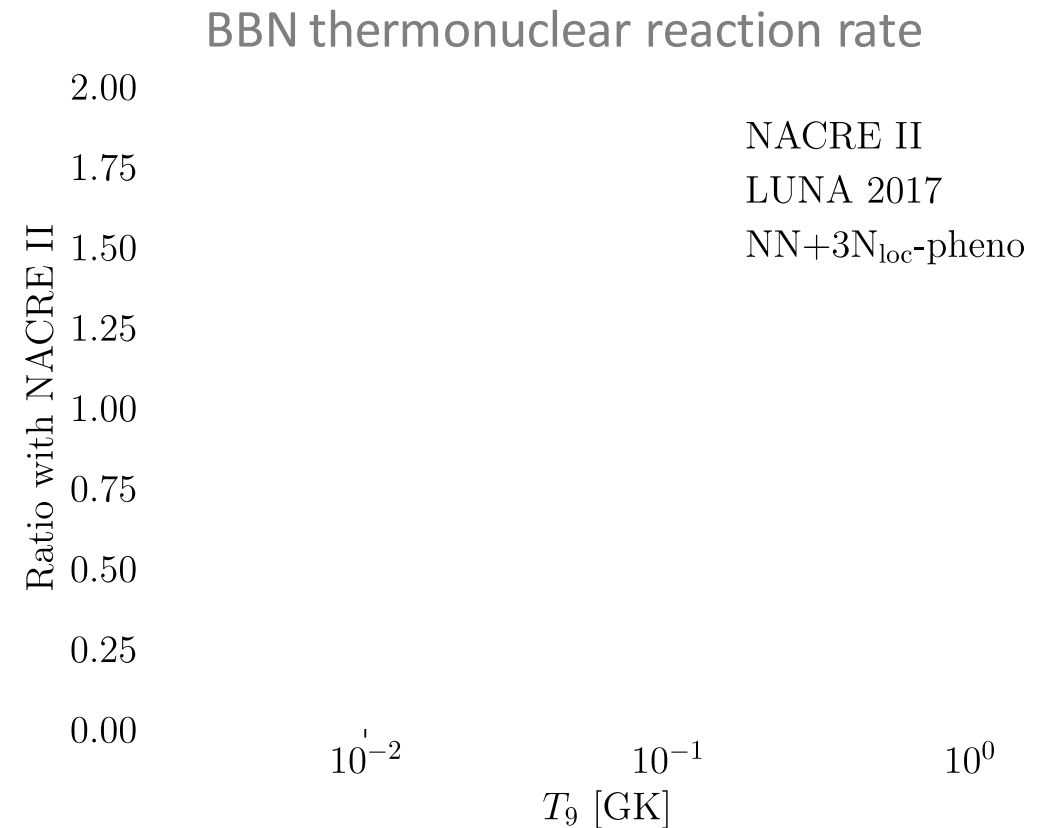
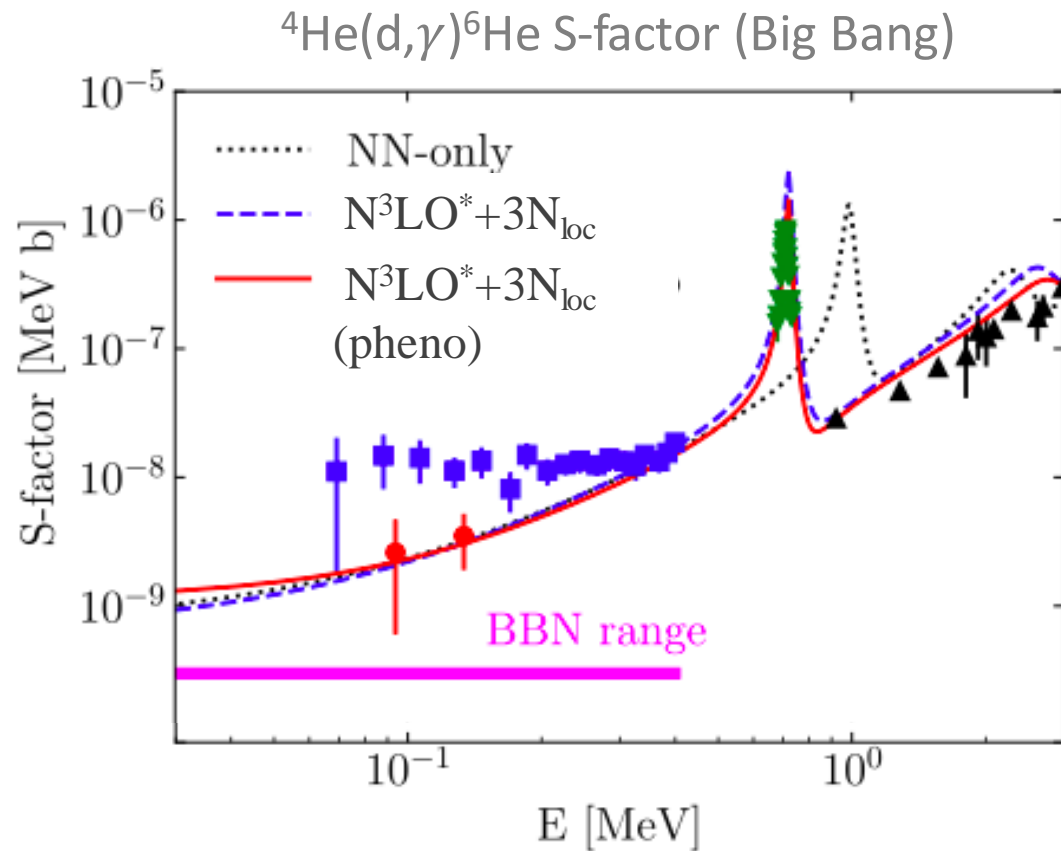


Repeat with computed correlations
for energies up to 1.5 MeV

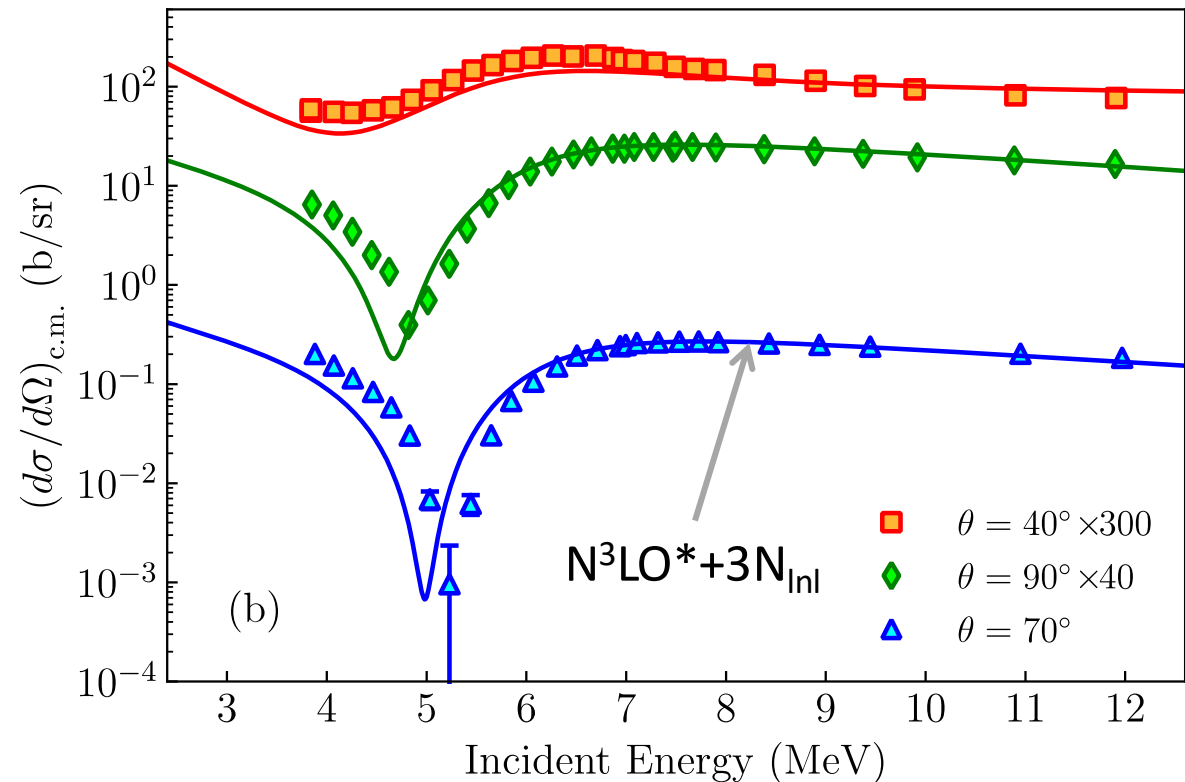
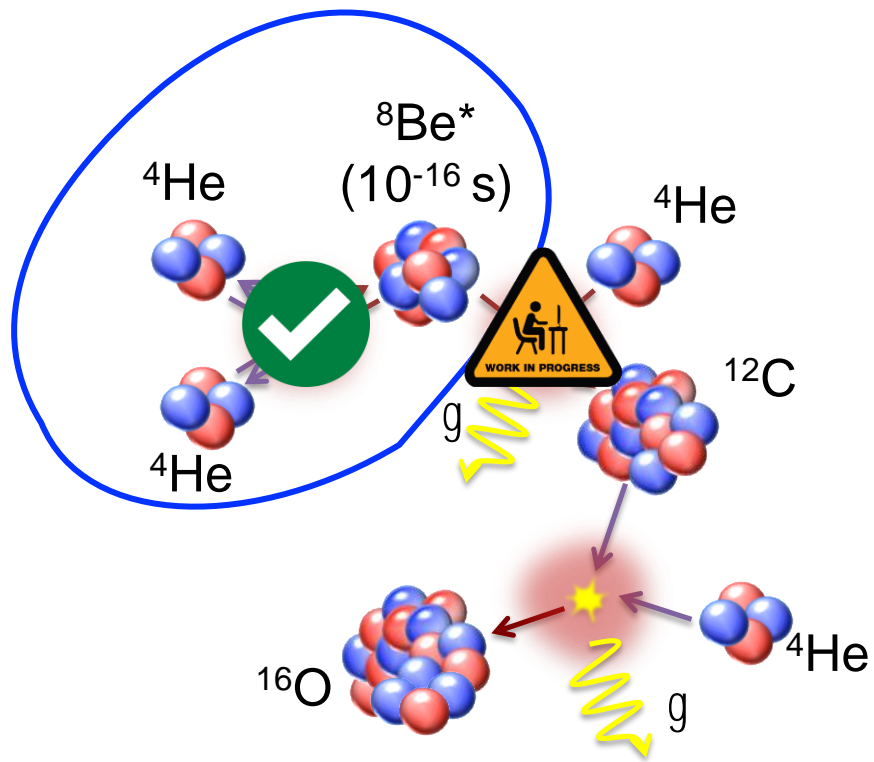
Now employing NN+3N forces to compute ${}^3\text{He}(\alpha,\gamma){}^7\text{Be}$,
perform analogous ab initio informed evaluation



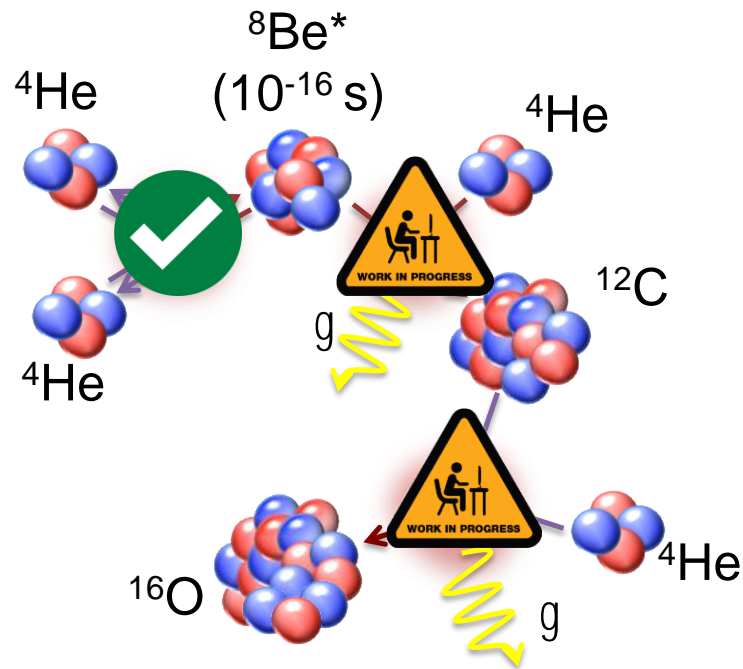
The NCSMC also successfully applied to BBN reactions
(yielding again significantly reduced uncertainties)



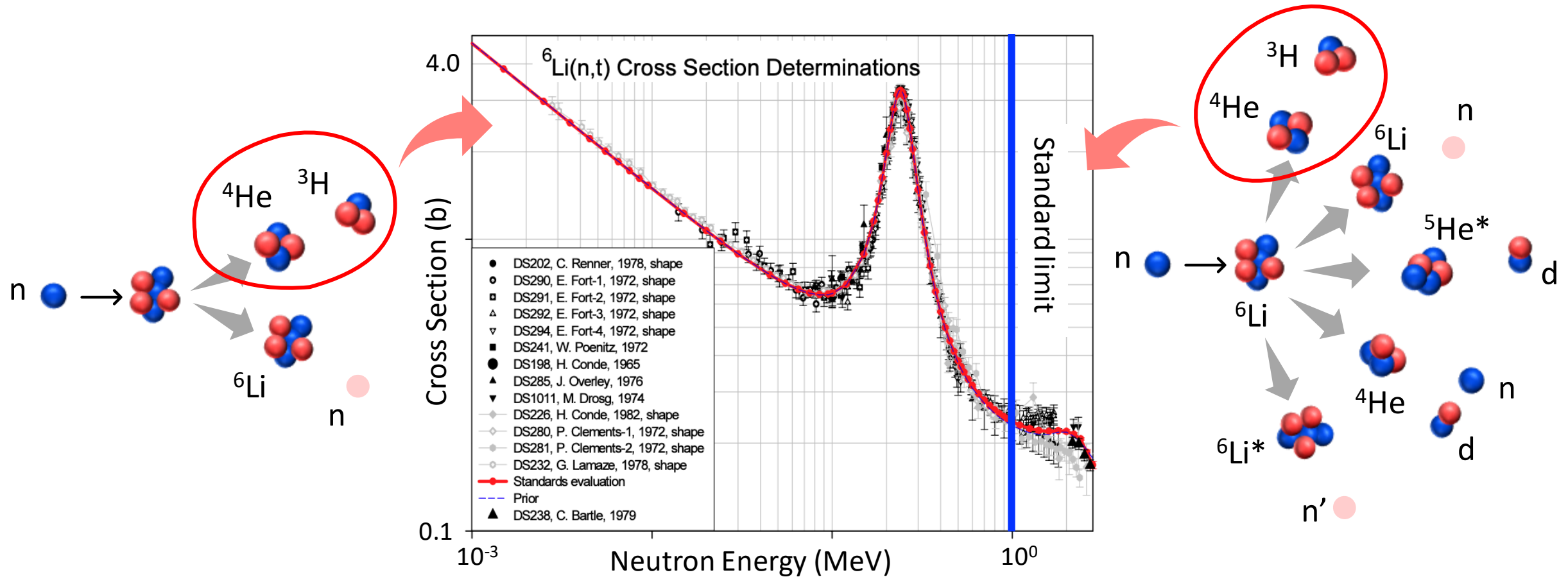
Smart formalism, GPUs also enabled description of $^4\text{He} + ^4\text{He}$ scattering, first stage of helium burning



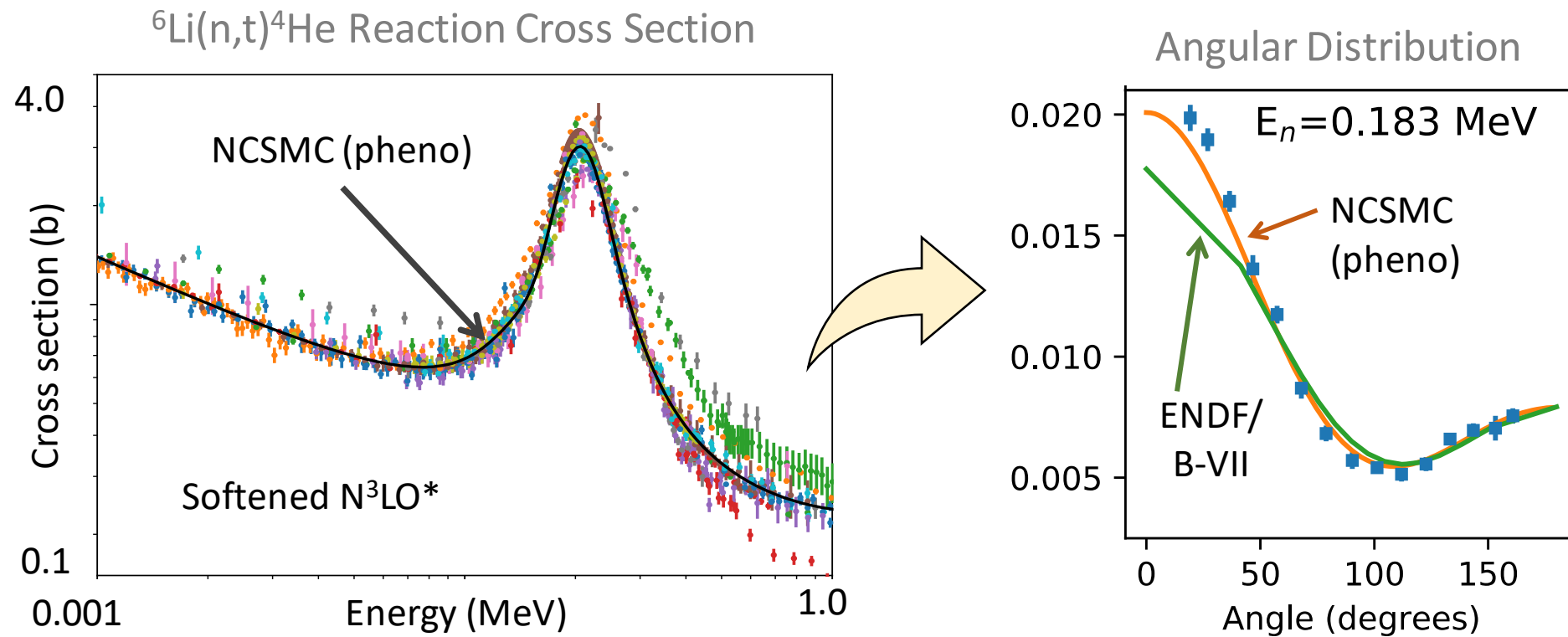
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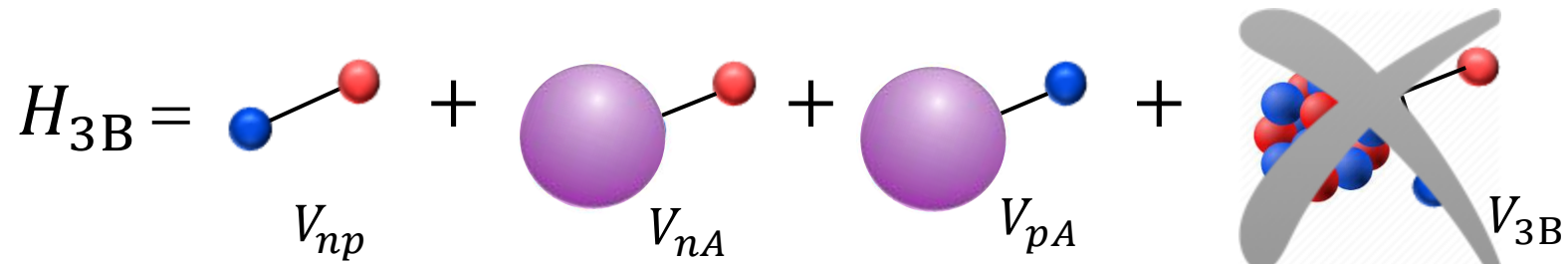
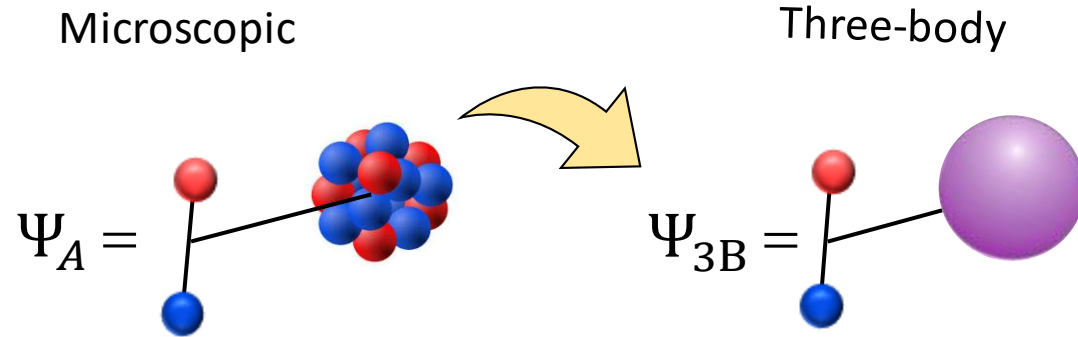
Ab initio reaction theory can also aid in improving evaluation of neutron standard cross sections



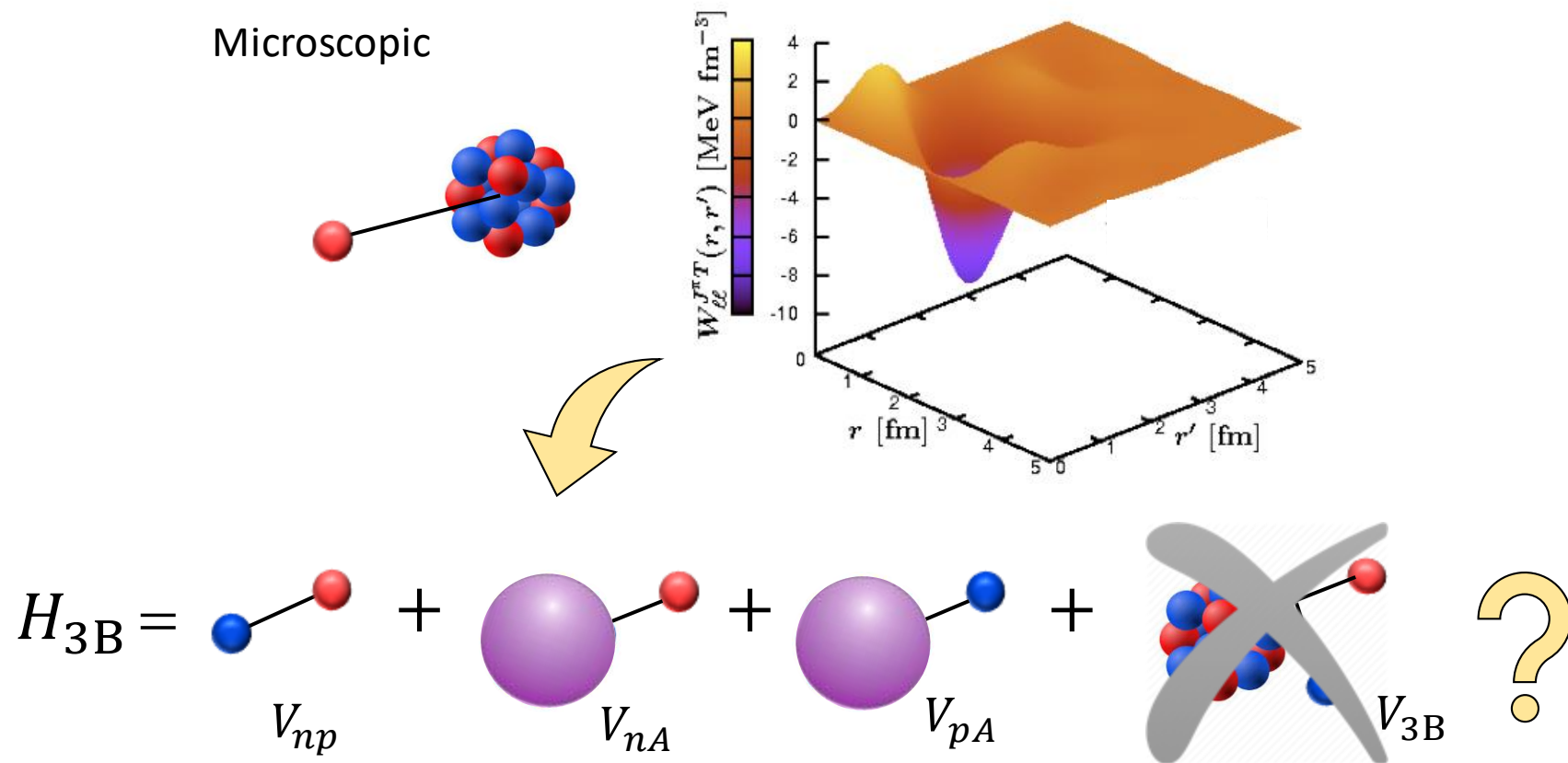
Calculation adjusted to reproduce reaction cross section
yields improved predictions for angular distributions



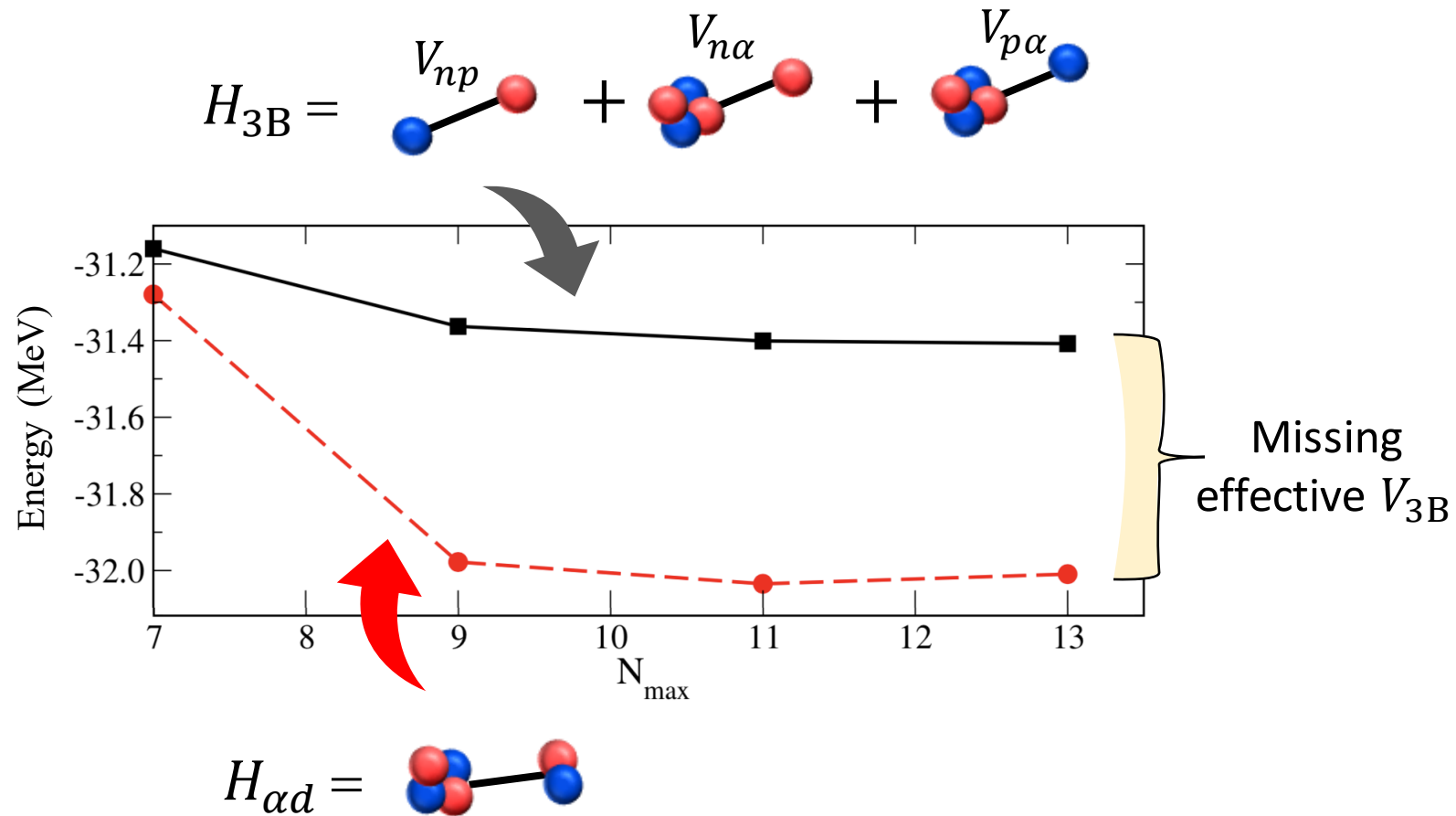
For heavier systems,
few-body reaction models are more effective



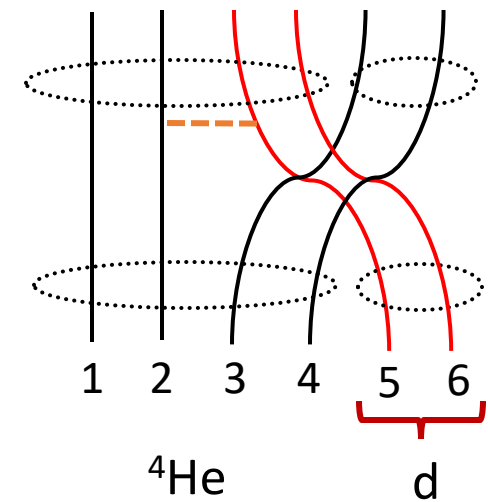
Can we bridge the gap between ab initio calculations and few-body models of nuclear reactions?



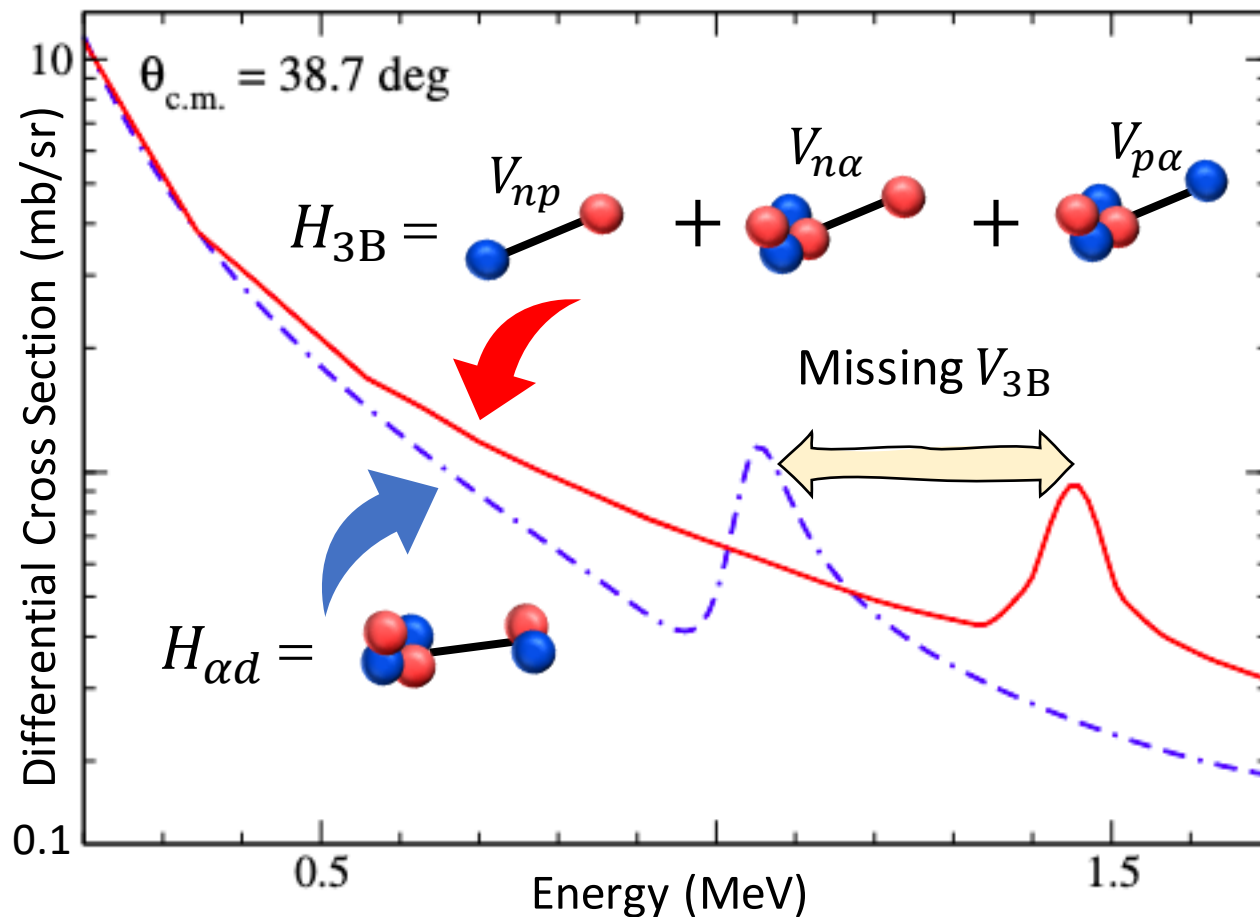
In test ground ${}^4\text{He}(\text{g.s.})+d$ system, omission of 3-body force causes ~ 600 keV underbinding for the ${}^6\text{Li}$ ground state ...



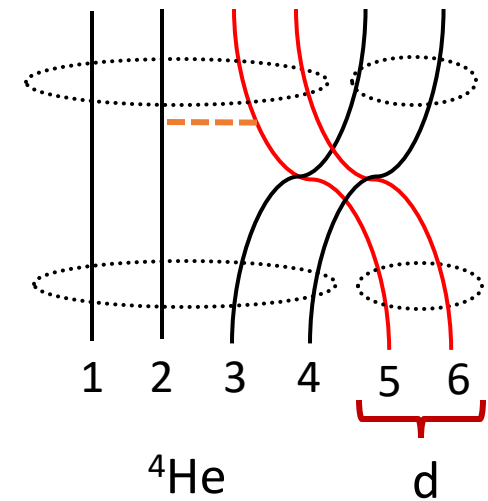
Simultaneous interaction
& two-nucleon exchange



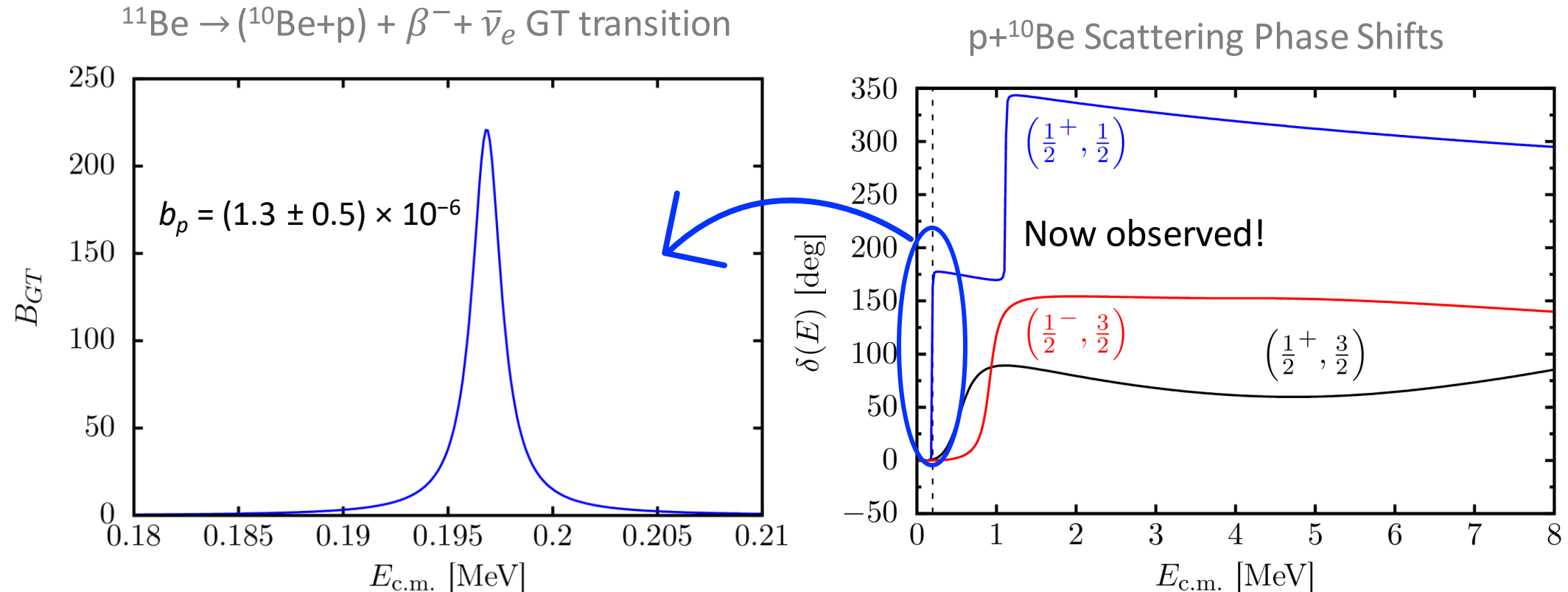
... ~ 400 keV shift to higher energy of 3^+ ^4He -d resonance



Simultaneous interaction
& two-nucleon exchange



NCSMC extended to describe exotic ^{11}Be β p emission,
 supports large branching ratio due to narrow $\frac{1}{2}^+$ resonance



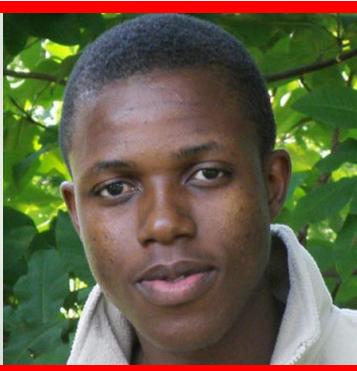
The ab initio structure and reactions team



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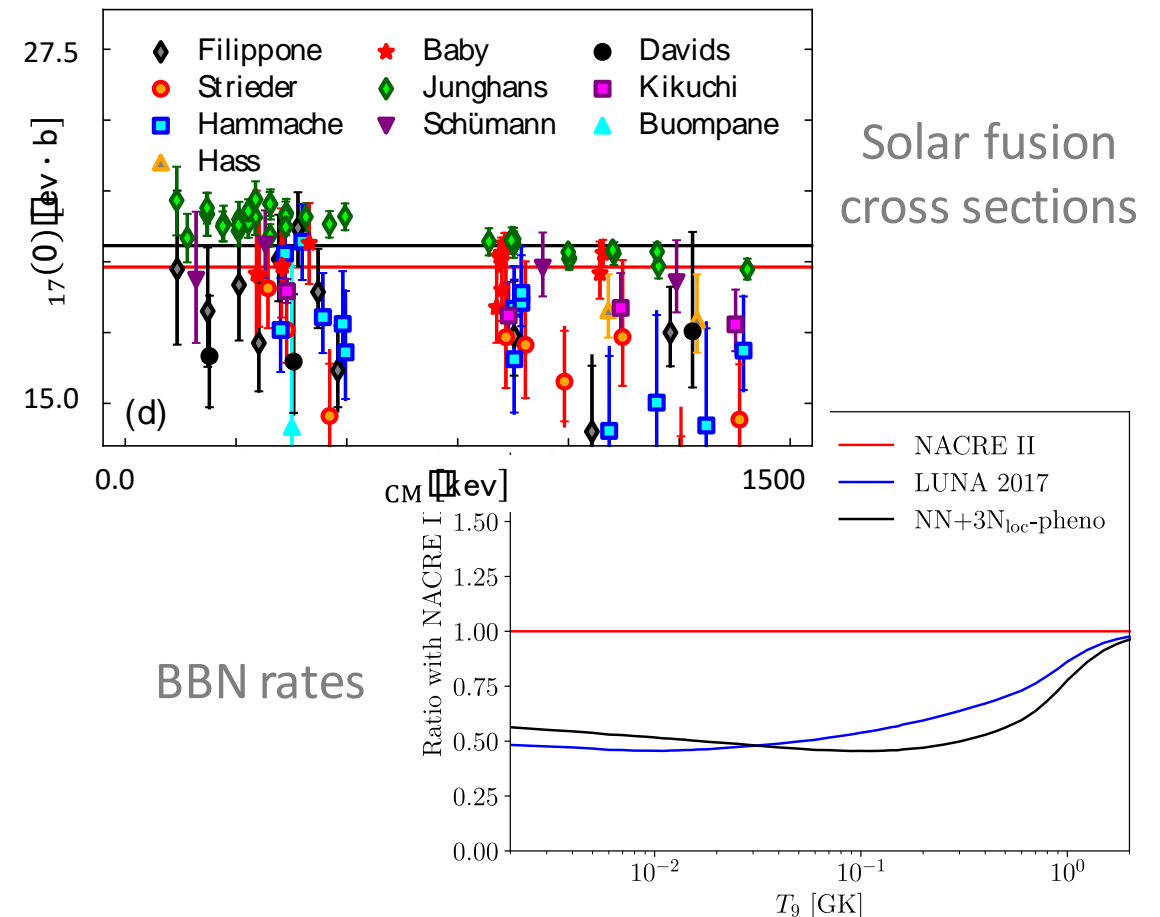
Predictive ab initio calculations are enabling substantially reduced uncertainties for astrophysical rates

New evaluation protocol combines:

- Ab initio calculations with chiral NN+3N forces
- Expt. data both at low and higher energies

Progress also on other fronts:

- Predictions of neutron standard cross sections
- Ab initio informed few-body reaction models
- Predictions of beta-delayed particle emission





Structure, scattering and reactions obtained with unified treatment of bound and unbound states

$$\Psi = \sum_{\lambda} c_{\lambda} \left| \begin{array}{c} \text{Potential well diagram} \end{array} \right\rangle + \sum_{\nu} \int dr u_{\nu}(r) \left| \begin{array}{c} \text{Scattering state diagram} \end{array} \right\rangle$$

No Core Shell Model with continuum
(NCSMC)

Structure, scattering and reactions obtained with unified treatment of bound and unbound states

$$\Psi = \underbrace{\sum_{\lambda} c_{\lambda} \left| \begin{array}{c} \text{Diagram of nucleons in a potential well} \end{array} \right\rangle}_{\text{Bound states}} + \sum_{\nu} \int dr u_{\nu}(r) \left| \begin{array}{c} \text{Diagram of two nucleons, one in a red cloud and one in a blue cloud} \end{array} \right\rangle$$

Static solutions for aggregate system,
describe all nucleons close together

Structure, scattering and reactions obtained with unified treatment of bound and unbound states

$$\Psi = \sum_{\lambda} c_{\lambda} \left| \begin{array}{c} \text{Potential well with} \\ \text{bound states and} \\ \text{clusters} \end{array} \right\rangle + \underbrace{\sum_{\nu} \int dr u_{\nu}(r) \left| \begin{array}{c} \text{Separated} \\ \text{projectile and} \\ \text{target} \end{array} \right\rangle}_{\text{Continuous microscopic cluster states, describe separated projectiles \& targets}}$$

Continuous microscopic cluster states,
describe separated projectiles & targets

Phenomenological correction obtained by treating NCSM eigenenergy as an adjustable parameter

The diagram illustrates the derivation of a phenomenological correction to the NCSM eigenenergy. It shows the relationship between the NCSM Hamiltonian matrix, the correction term h , and the RGM Hamiltonian matrix.

On the left, a blue box contains the expression $E_\lambda \delta_{\lambda\lambda'}$, which is circled in red. A blue arrow points from this box to the top-left element of the matrix H_{NCSM} in the equation below.

At the top center, a green box contains the expression $\langle (A) \left| H \hat{A}_n^\dagger \right| \begin{smallmatrix} (a) \\ (A-a) \end{smallmatrix} \rangle$, with a diagram of two nuclei (one with A nucleons, one with $A-a$ nucleons) and a vector r . A green arrow points from this box to the top-right element h of the matrix H_{NCSM} .

At the bottom left, a red box contains the expression $\langle \begin{smallmatrix} (A-a) \\ (a) \end{smallmatrix} \left| \hat{A}_n^\dagger H \hat{A}_n^\dagger \right| \begin{smallmatrix} (a) \\ (A-a) \end{smallmatrix} \rangle$, with a diagram of two nuclei and vectors r and r' . A red arrow points from this box to the bottom-right element H_{RGM} of the matrix.

The central equation is:

$$\begin{pmatrix} H_{NCSM} & h \\ h & H_{RGM} \end{pmatrix} \begin{pmatrix} c \\ u \end{pmatrix} = E \begin{pmatrix} 1_{NCSM} & g \\ g & N_{RGM} \end{pmatrix} \begin{pmatrix} c \\ u \end{pmatrix}$$

On the right, a blue box contains the expression $\delta_{\lambda\lambda'}$. A blue arrow points from this box to the top-left element 1_{NCSM} of the matrix in the equation.

At the top right, a green box contains the expression $\langle (A) \left| \hat{A}_n^\dagger \right| \begin{smallmatrix} (a) \\ (A-a) \end{smallmatrix} \rangle$, with a diagram of two nuclei and a vector r . A green arrow points from this box to the top-right element g of the matrix.

At the bottom right, a red box contains the expression $\langle \begin{smallmatrix} (A-a) \\ (a) \end{smallmatrix} \left| \hat{A}_n^\dagger \hat{A}_n^\dagger \right| \begin{smallmatrix} (a) \\ (A-a) \end{smallmatrix} \rangle$, with a diagram of two nuclei and vectors r and r' . A red arrow points from this box to the bottom-right element N_{RGM} of the matrix.

All other characteristics of the S-matrix still predicted from ab initio theory.