

1. Determination of Fourier Series of Discrete Time Periodic Signal from a Matrix Perspective (矩阵视角下确定离散时间周期函数的傅里叶级数),

$$\begin{pmatrix} x[0] \\ x[1] \\ \vdots \\ x[N-1] \end{pmatrix} = \begin{pmatrix} 1 & 1 & \dots & 1 \\ 1 & \omega^{-1} & \dots & \omega^{-(N-1)} \\ \vdots & \vdots & \dots & \vdots \\ 1 & \omega^{-(N-1)} & \dots & \omega^{-(N-1)^2} \end{pmatrix} \begin{pmatrix} a_0 \\ a_1 \\ \vdots \\ a_{N-1} \end{pmatrix}, \omega = e^{j\frac{2\pi}{N}}$$

2. Orthogonality of Discrete Complex Exponentials(离散复指数函数的正交性),

$$\langle \omega^k, \omega^r \rangle = \sum_{n=0}^{N-1} e^{j(k-r)\frac{2\pi}{N}n} = \begin{cases} 0, & k \neq r \\ N, & k = r \end{cases}$$

3. Fourier Series from a Rotating Unit Vector Perspective (傅里叶级数的复空间单位向量旋转视角)。
4. Filtering Concept (滤波器概念)。



Ch 4.

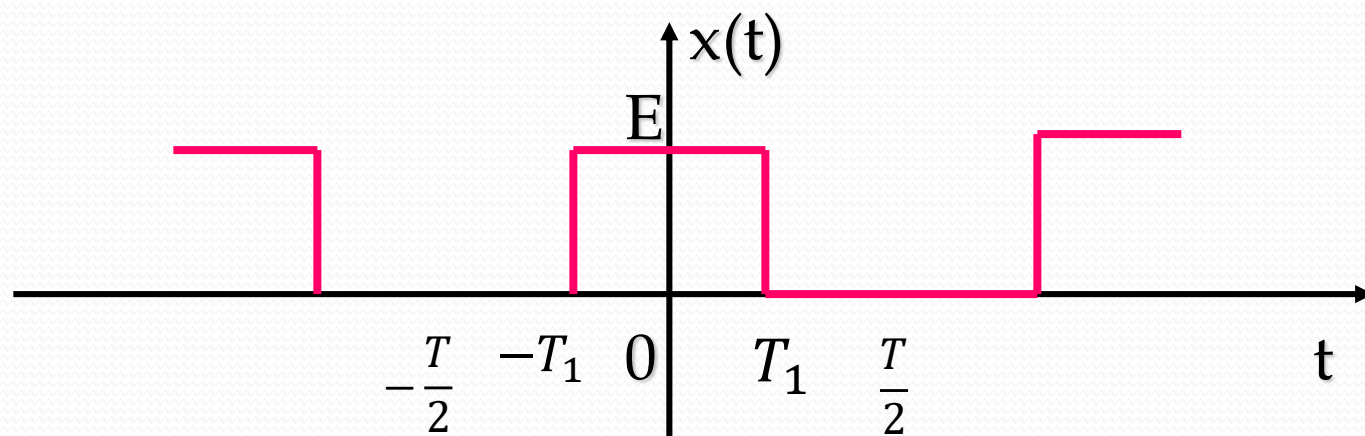
Fourier Transform

- How do we represent aperiodic (非周期的) signals in a similar way?
 - From Fourier Series to Fourier transform
 - Convergence
 - Examples
 - Properties



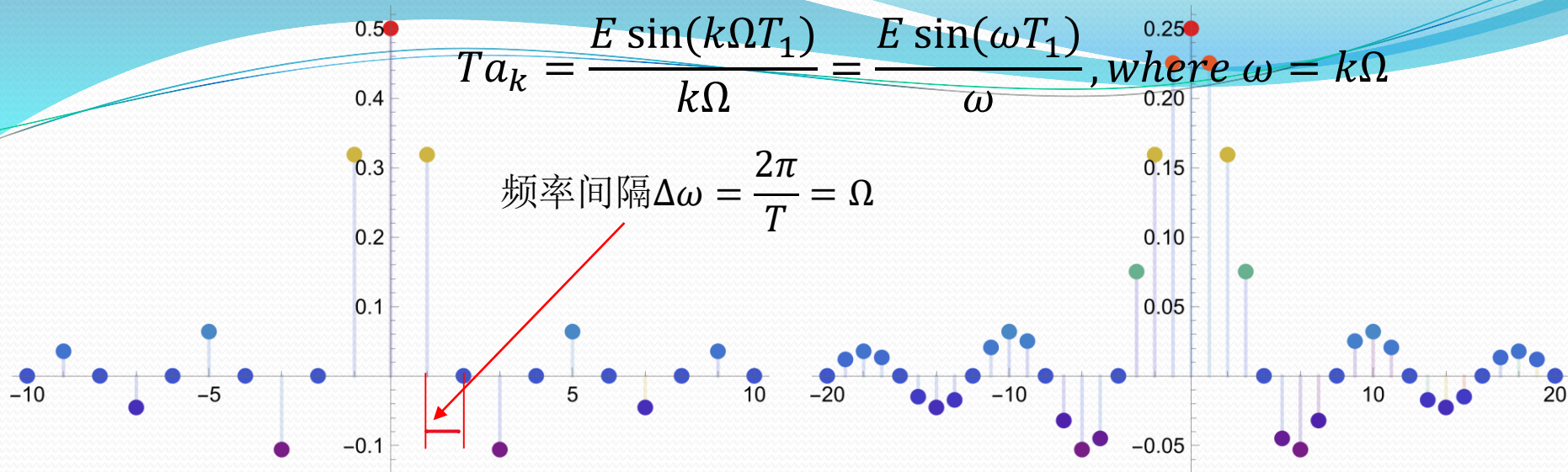
4.1 Representation of Aperiodic Signals: The Continuous-Time Fourier Transform

Recall the spectrum of periodic signals

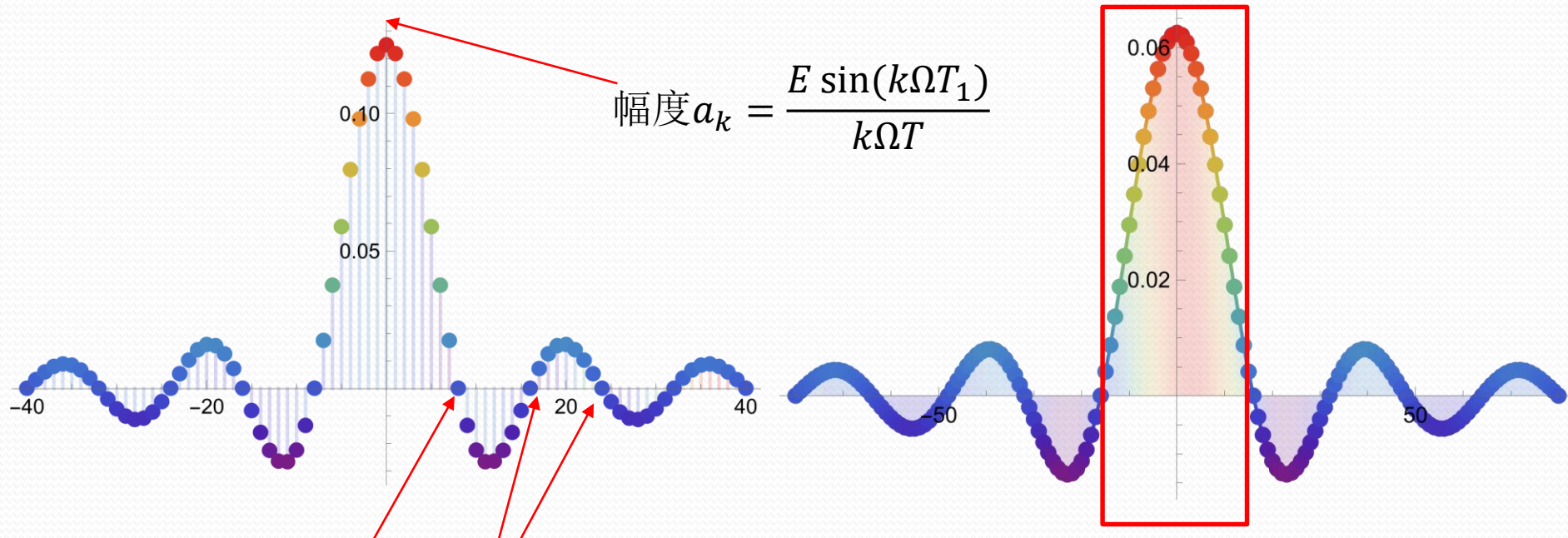


$$Ta_k = \frac{E \sin(k\Omega T_1)}{k\Omega} = \frac{E \sin(\omega T_1)}{\omega}, \text{ where } \omega = k\Omega$$

$$\text{频率间隔 } \Delta\omega = \frac{2\pi}{T} = \Omega$$

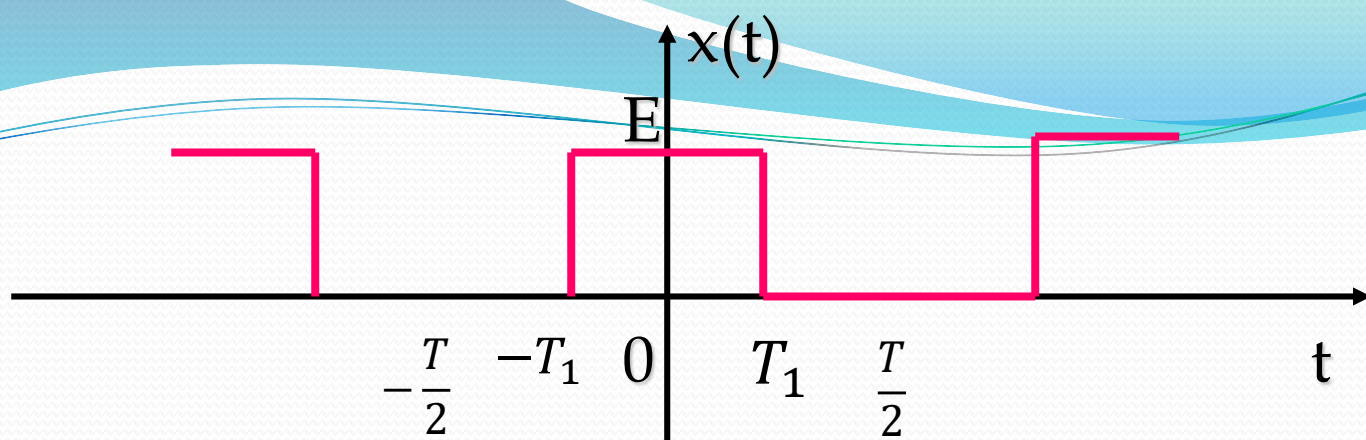


$$\text{幅度 } a_k = \frac{E \sin(k\Omega T_1)}{k\Omega T}$$



$$\text{过零点位置 } \omega = \frac{m\pi}{T_1}$$

$$\text{带宽 } B_\omega = \frac{2\pi}{T_1}$$



From this example, we can see:

1. When the period, T , tends to infinite, periodic signal turns to be an aperiodic impulse signal.

$$T \rightarrow \infty$$

2. Meanwhile, the frequency turns to be a continuous independent variable.

$$\Omega = \frac{2\pi}{T} \rightarrow 0 \rightarrow d\omega \quad n\Omega \rightarrow \omega$$

From Fourier Series to Fourier Transform

$$F_n = \frac{1}{T} \int_{-\frac{T}{2}}^{\frac{T}{2}} f(t) \cdot e^{-jn\Omega t} \cdot dt$$

$$f(t) = \sum_{n=-\infty}^{\infty} F_n \cdot e^{jn\Omega t}$$

As $T \rightarrow \infty$

$$F_n T = F_n \cdot \frac{2\pi}{\Omega} = \int_{-\infty}^{\infty} f(t) \cdot e^{-jn\Omega t} \cdot dt$$

$$f(t) = \sum_{n=-\infty}^{\infty} F_n T \cdot e^{jn\Omega t} \cdot \frac{1}{T}$$

Denote:
$$F(j\omega) = \lim_{T \rightarrow \infty} \frac{F_n}{1/T} = \lim_{T \rightarrow \infty} F_n T$$

$$T \rightarrow \infty \quad \Omega = \frac{2\pi}{T} \rightarrow \rho \quad n\Omega \rightarrow \omega \quad \Delta(n\Omega) = \Omega \rightarrow d\omega$$

$$\sum_{n=-\infty}^{\infty} \rightarrow \int_{-\infty}^{\infty}, \quad \frac{1}{T} = \frac{\Omega}{2\pi} \rightarrow \frac{1}{2\pi} d\omega$$

$$F_n T = \int_{-\frac{T}{2}}^{\frac{T}{2}} f(t) \cdot e^{-jn\Omega t} \cdot dt$$

$$\Rightarrow F(j\omega) = \int_{-\infty}^{\infty} f(t) \cdot e^{-j\omega t} \cdot dt$$

**Fourier
Transform**

**Inverse
Fourier Transform**

$$f(t) = \sum_{n=-\infty}^{\infty} F_n T \cdot e^{jn\Omega t} \cdot \frac{1}{T} \Rightarrow \frac{1}{2\pi} \int_{-\infty}^{\infty} F(j\omega) e^{j\omega t} d\omega$$

$$F(j\omega) = \int_{-\infty}^{\infty} f(t)e^{-j\omega t} dt \quad f(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} F(j\omega)e^{j\omega t} d\omega$$

Fourier Transform and Inverse Fourier Transform also can be represented as the following manner

$$F(f) = \int_{-\infty}^{\infty} f(t)e^{-j2\pi ft} dt$$

$$f(t) = \int_{-\infty}^{\infty} F(f) \cdot e^{j2\pi ft} df$$

$$f = \frac{1}{T}, \omega = 2\pi f$$

❖ $F(j\omega)$ is different from F_n

➤ $F(j\omega)$ is defined as a density function;

➤ $F(j\omega)$ is a continuous function of ω (for aperiodic);

➤ $F(j\omega)$ contains component with frequency ranging from 0 to infinite;

➤ No harmonic relations among these components.

$$F(j\omega) = \int_{-\infty}^{\infty} f(t) e^{-j\omega t} dt = |F(j\omega)| e^{j\phi(\omega)}$$

$$f(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} F(j\omega) e^{j\omega t} d\omega$$

$$= \frac{1}{2\pi} \int_{-\infty}^{\infty} |F(j\omega)| e^{j\omega t + j\phi(\omega)} d\omega$$

If $f(t)$ is real, $|F(j\omega)|$ is an even function of ω , and the phase of $F(j\omega)$ is an odd function of ω .

WHY?

The sufficient condition for the existence of Fourier Transform

$$\int_{-\infty}^{\infty} |f(t)| dt < \infty$$

- Some singular functions, which do not satisfy the sufficient condition, have corresponding Fourier transforms if unit impulse function is applicable.
- We will interpret it in the chapter of Laplace transform.

4.2 Spectrums of Some Typical Aperiodic Signals

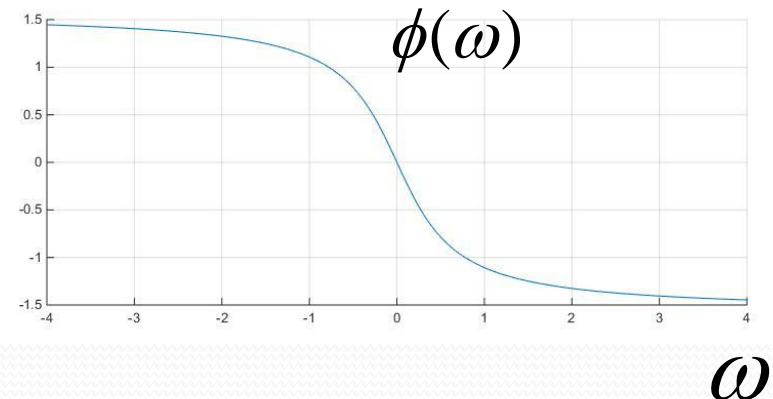
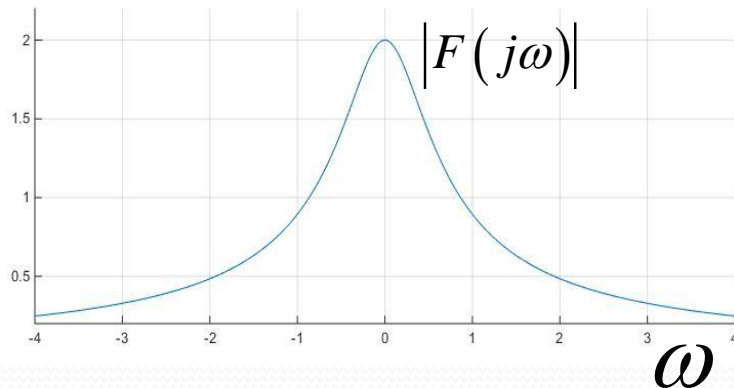
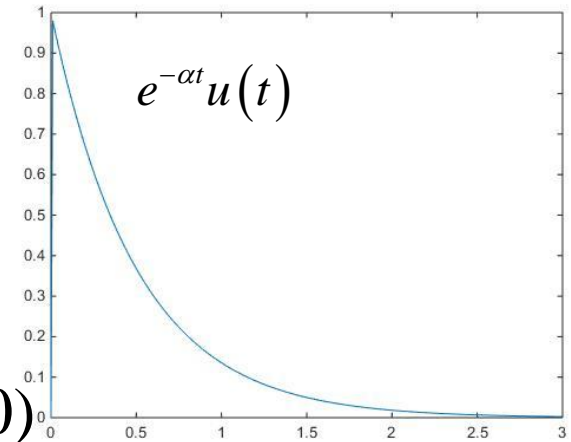
- One-side exponential signal
- Dual-side exponential signal
- Rectangular pulse signal
- Unit impulse function
- Direct current signal

● One-Side Exponential Signal

$$f(t) = \begin{cases} e^{-\alpha t} & (t \geq 0) \\ 0 & (t < 0) \end{cases}$$

$$F(j\omega) = \int_{-\infty}^{\infty} f(t)e^{-j\omega t} dt = \frac{1}{\alpha + j\omega} \quad (\alpha > 0)$$

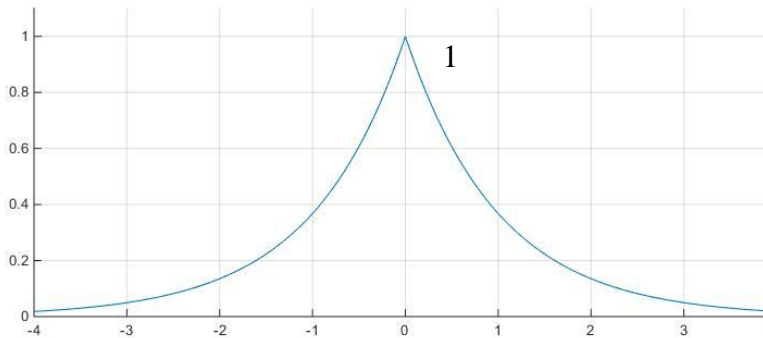
$$|F(j\omega)| = \frac{1}{\sqrt{\alpha^2 + \omega^2}} \quad \phi(\omega) = -\arctan\left(\frac{\omega}{\alpha}\right)$$



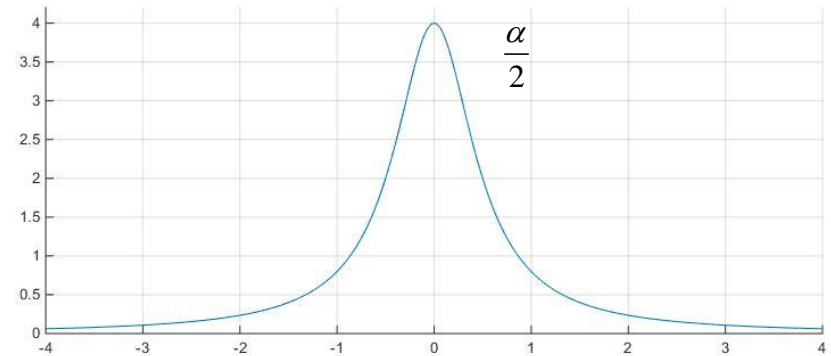
● Dual-Side Exponential Signal

$$f(t) = e^{-\alpha|t|} \quad (-\infty < t < +\infty)$$

$$F(j\omega) = \frac{2\alpha}{\alpha^2 + \omega^2} \quad \varphi(\omega) = 0$$



$$f(t) = e^{-\alpha|t|} \quad (\alpha > 0)$$



$$F(j\omega) = \frac{2\alpha}{\alpha^2 + \omega^2}$$

● Rectangular pulse signal

$$f(t) = \begin{cases} E & (|t| \leq \frac{\tau}{2}) \\ 0 & (|t| > \frac{\tau}{2}) \end{cases}$$

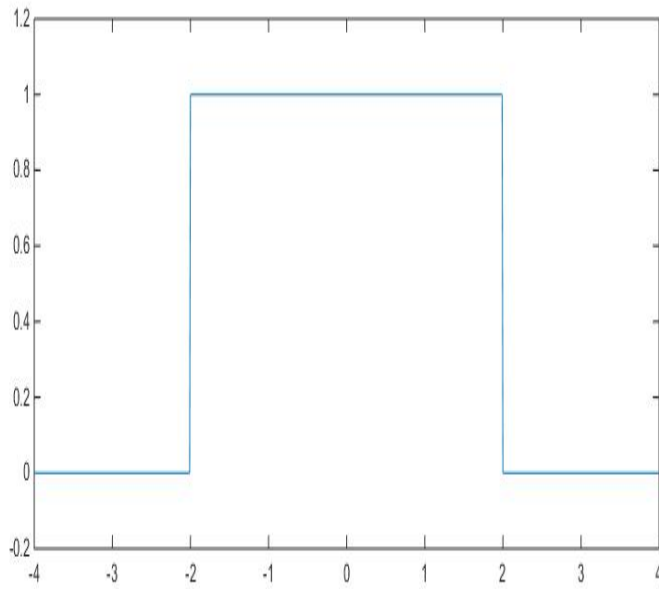
$$F(j\omega) = \int_{-\tau/2}^{\tau/2} E e^{-j\omega t} dt = \frac{2E}{\omega} \sin\left(\frac{\omega\tau}{2}\right)$$

$$= E\tau \left(\frac{\sin(\frac{\omega\tau}{2})}{\frac{\omega\tau}{2}} \right) = E\tau \operatorname{Sa}\left(\frac{\omega\tau}{2}\right)$$

$$|F(j\omega)| = E\tau \left| \operatorname{Sa}\left(\frac{\omega\tau}{2}\right) \right| \quad \varphi(\omega) = \begin{cases} 0 & \left(\frac{4n\pi}{\tau} < |\omega| < \frac{2(2n+1)\pi}{\tau} \right) \\ \pi & \left(\frac{2(2n+1)\pi}{\tau} < |\omega| < \frac{4(n+1)\pi}{\tau} \right) \end{cases}$$

Very
Important!

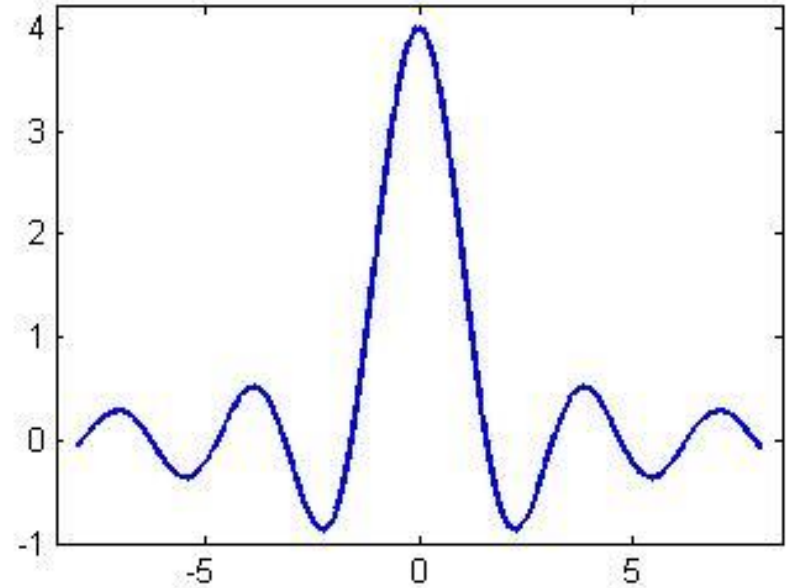
$$\operatorname{Sa}(x) = \frac{\sin x}{x}$$



$$g(t)$$

$$\tau = 4$$

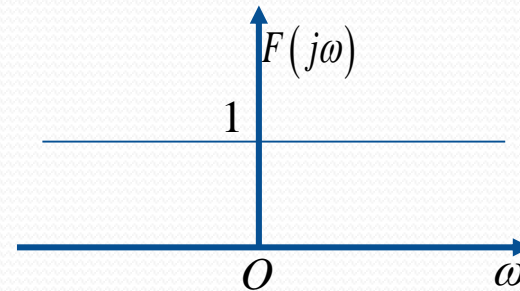
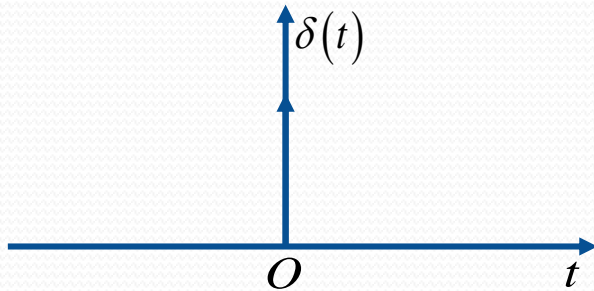
$$E = 1$$



$$G(j\omega) = E\tau \text{Sa}\left(\frac{\omega\tau}{2}\right) = 4\text{Sa}(2\omega)$$

● FT of Unit Impulse Function

$$F(j\omega) = \int_{-\infty}^{\infty} \delta(t) e^{-j\omega t} dt = 1$$



Simple,
but important.

● Fourier Transform of Direct Current Signal

For $f(t) = e^{-\alpha|t|}$ $(-\infty < t < +\infty)$

when $\alpha \rightarrow 0$, $f(t) \rightarrow 1$

$$\text{then } \lim_{\alpha \rightarrow 0} F(j\omega) = \lim_{\alpha \rightarrow 0} \frac{2\alpha}{\alpha^2 + \omega^2} = \begin{cases} 0, & \omega \neq 0 \\ \infty, & \omega = 0 \end{cases}$$

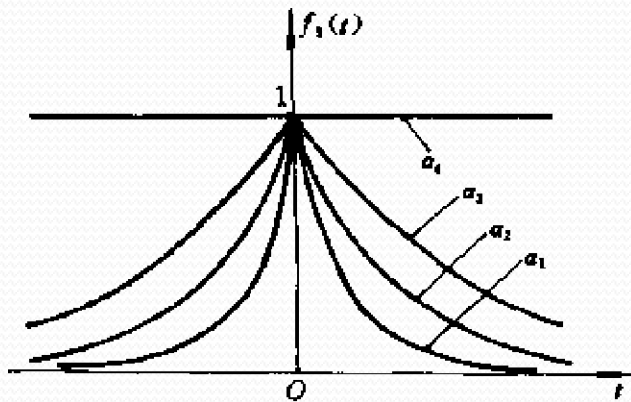
The strength this impulse has is:

$$\lim_{\alpha \rightarrow 0} \int_{-\infty}^{\infty} \frac{2\alpha}{\alpha^2 + \omega^2} d\omega = \lim_{\alpha \rightarrow 0} \int_{-\infty}^{\infty} \frac{2}{1 + \left(\frac{\omega}{\alpha}\right)^2} d\left(\frac{\omega}{\alpha}\right)$$

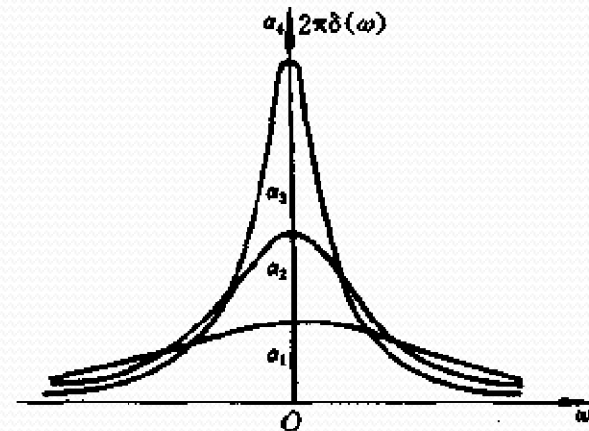
$$= \lim_{\alpha \rightarrow 0} 2 \arctan\left(\frac{\omega}{\alpha}\right) \Bigg|_{-\infty}^{\infty} = 2\pi$$

$$\therefore \lim_{\alpha \rightarrow 0} \frac{2\alpha}{\alpha^2 + \omega^2} = 2\pi\delta(\omega)$$

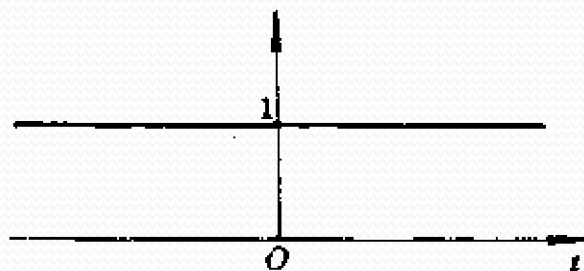
$$\therefore F[1] = 2\pi\delta(\omega)$$



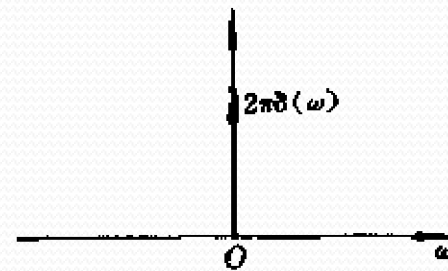
(a) $f_1(t) = e^{-\alpha_1|t|}$



(b) $F_1(j\omega) = \frac{2\alpha_1}{\alpha_1^2 + \omega^2}$



(a)



(b)

Summary

- From Fourier Series to Fourier Transform
- Spectrums of Some Typical Aperiodic Signals

Assignments

- 4.1
- 4.2