Two Experiments

Aharonov-Bohm interferometry and Wilson lines

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Outline

Aharonov-Bohm interferometry
Experimental setup and theoretical preparation
The experiment

Wilson lines

Experimental setup and theoretical preparation Measuring Wilson lines
Reconstructing band eigenstates
Determining Wilson line eigenvaluse
Accessing the dispersion relation

Outline

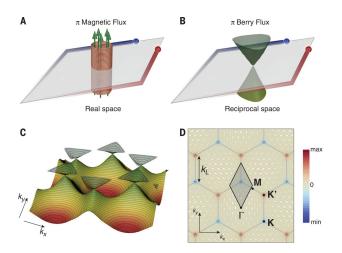
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An Aharonov-Bohm interferometer for determining Bloch band topology

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Bloch state

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$$\psi_{\mathbf{k}}^n(\mathbf{r}) = e^{i\mathbf{k}\mathbf{r}}u_{\mathbf{k}}^n(\mathbf{r})$$

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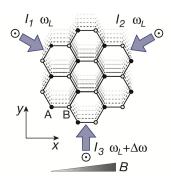
Berry phase

$$\phi_{\mathsf{Berry}} = \oint_C \mathbf{A}_n(\mathbf{k}) d\mathbf{k} = \int_{S_{\mathsf{CDA}}} \Omega_n(\mathbf{k}) d^2\mathbf{k}$$

Hexagonal lattice in real space

Hexagonal lattice in real space

Magnetic field $B = B_0 + \mathbf{r} \cdot \nabla B$ combined with an orthogonal acceleration $\mathbf{a} \perp \nabla B$ of the lattice:



The Hamiltonian is

$$H(t) = \frac{\mathbf{p}^2}{2m} + V[\mathbf{r} - \mathbf{R}]$$
$$-\mu \mathbf{r} \cdot \nabla B - \mu B_0$$

in co-moving frame

$$H(t) = \frac{\mathbf{p}^2}{2m} + V(\mathbf{r}) - \mathcal{F}_{\mu} \cdot \mathbf{r} + \varepsilon_{\mu}(t)$$

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$$\mathcal{F}_{\mu}$$
: $\mathbf{k} \rightarrow \mathbf{k} + \mathcal{F}_{\mu}t$

Hamiltonian

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Ansatz

$$ilde{\Psi}(t) = e^{\mathrm{i}\eta(t)} \psi_{k_0 + \mathcal{F}_\mu t}^n \ \eta = \phi_\mathsf{dyn} + \phi_\mathsf{Berry}$$

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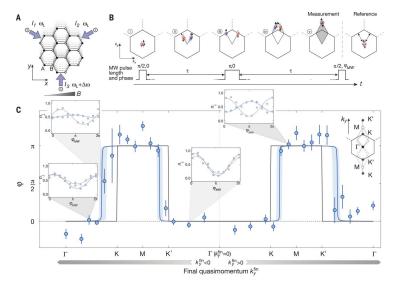
Phase

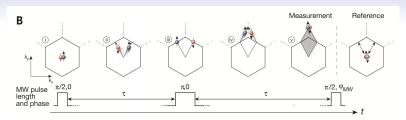
$$\phi_{ ext{dyn}} = \int_0^T [E_1(\mathbf{k} + \mathcal{F}_\mu t) + \varepsilon_\mu t] dt$$

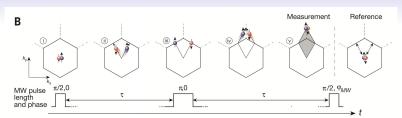
$$\phi_{ ext{Berry}} = \mathrm{i} \int_C \langle u_\mathbf{k}^1 | \nabla_\mathbf{k} | u_\mathbf{k}^1
angle d\mathbf{k}$$

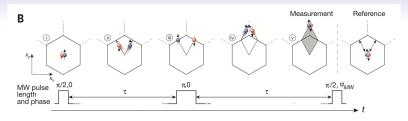
Experiment procedure

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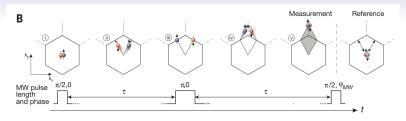




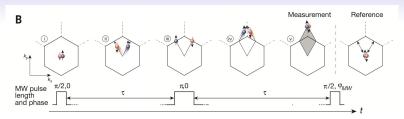




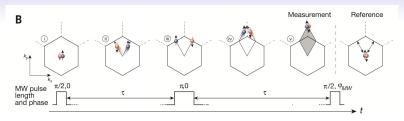
1. ^{87}Rb BEC initial state $|\uparrow\rangle = |F = 2, m_F = 1\rangle$; $\pi/2$ -microwave pulse;



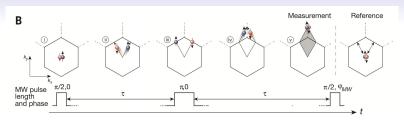
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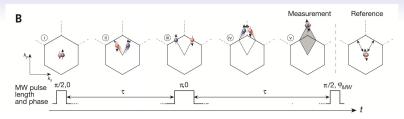
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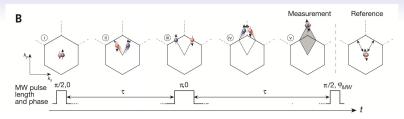


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zero-area reference: V-shape path.

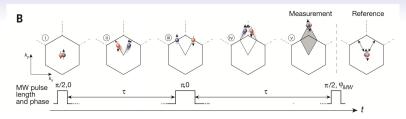


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Reversing the lattice acceleration after π pulse.



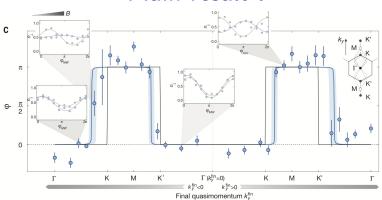
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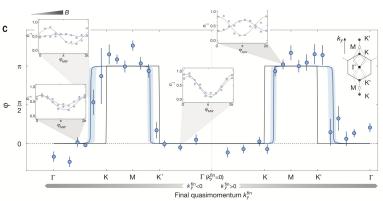
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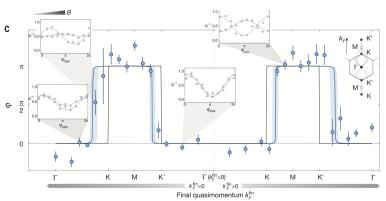
Reversing the lattice acceleration after π pulse.

$$n_{\uparrow,\downarrow} \propto 1 \pm \cos(\varphi + \varphi_{MW})$$





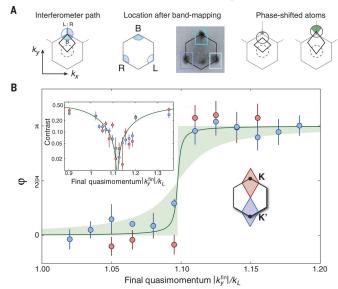
Broadening of the edges — caused by momentum spread.

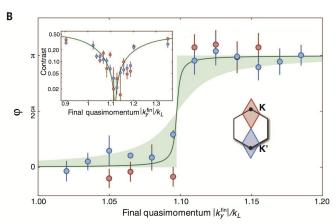


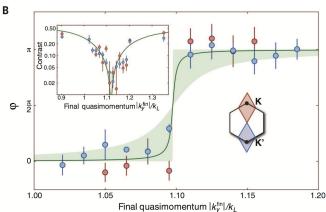
- Broadening of the edges caused by momentum spread.
- Systematic errors.

Self-referenced interferometry at Dirac point

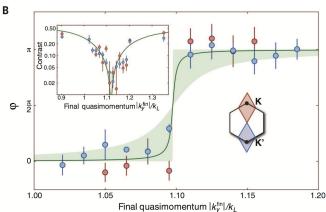
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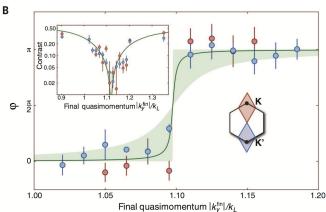




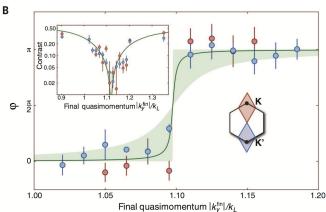
Contrast: $(n_{\downarrow}^{\mathsf{max}} - n_{\downarrow}^{\mathsf{min}}) / (n_{\downarrow}^{\mathsf{max}} + n_{\downarrow}^{\mathsf{min}})$



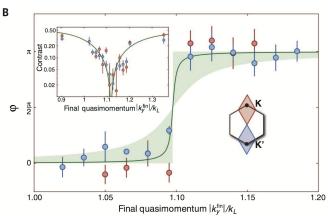
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- $\varphi = (\varphi_L + \varphi_R)/2 \varphi_B$



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- $\varphi = (\varphi_L + \varphi_R)/2 \varphi_B = 0.95(10)\pi$



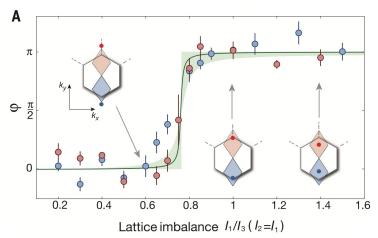
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- $\varphi = (\varphi_L + \varphi_R)/2 \varphi_B = 0.95(10)\pi$
- Berry curvature localization $\delta k_w \simeq 10^{-4} k_L (\Delta \simeq h \times 3 \text{Hz})$

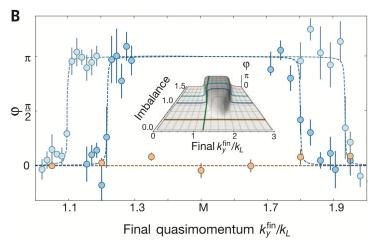
Imbalance lattice mapping

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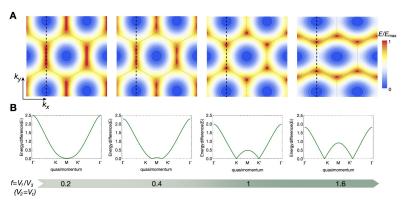
Imbalance lattice mapping Self-referenced phase

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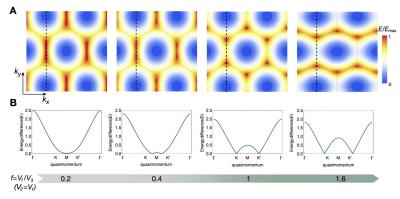


ab initio calculation of imbalanced lattice

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ab initio calculation of imbalanced lattice



Seeing Dirac points annihilating clear.

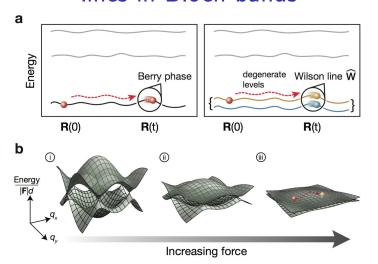
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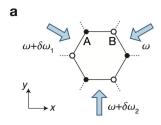
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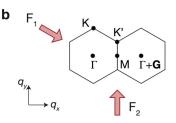
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Experimental reconstruction of Wilson lines in Bloch bands



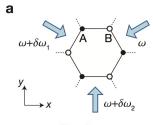


Real space



Reciprocal space

b

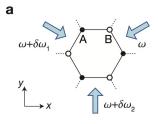


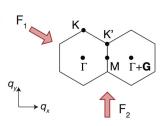
 $\begin{array}{c|c} F_1 & K_0 & K' \\ \hline & \Gamma & M & \Gamma + \mathbf{G} \\ \hline & q_y & & & \Gamma \\ \end{array}$

Real space

Reciprocal space

• force **F** :
$$q(t) = q(0) + Ft$$





Real space

Reciprocal space

- force **F** : q(t) = q(0) + Ft
- unitary time-evolution operator (Wilson line matrix)

$$W_{\mathbf{q}(0)\to\mathbf{q}(t)} = \mathcal{P}\exp(i\int_C A_{\mathbf{q}}d\mathbf{q})$$

 A_q : Wilczek-Zee connection \mathcal{P} : Path-ordering (non-Abelian)

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Wilson line

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- $n \neq n'$: inter-band Berry connections induce inter-band transition

total Hamiltonian:

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constant force

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initial state

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constant force F

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- lattice

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constant force F

- initial state $|\psi(0)\rangle = \sum_n \alpha^n(0) |\Phi_{\mathbf{q}_0}^n\rangle$ $|\alpha^n(0)|^2$ gives the population in the n^{th} band at t=0
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Wilson lines and Wilczek-Zee connections Example (two-band system)

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where

$$\xi_{\mathbf{q}(t)}^{n,n'} = \mathbf{A}_{\mathbf{q}(t)}^{n,n'} \cdot \mathbf{F} = \mathbf{i} \langle u_{\mathbf{q}(t)}^n | \partial_t | u_{\mathbf{q}(t)}^{n'} \rangle$$

Example (two-band system)

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$$\mathbf{A}_{\mathbf{q}(t)}^{n,n'} = \mathbf{i} \langle u_{\mathbf{q}}^{n} | \nabla_{\mathbf{q}} | u_{\mathbf{q}}^{n'} \rangle \Big|_{\mathbf{q} = \mathbf{q}(t)}$$

$$\mathrm{i}\partial_t \begin{pmatrix} \alpha^1(t) \\ \alpha^2(t) \end{pmatrix} = \begin{pmatrix} -\xi_{\mathbf{q}(t)}^{1,1} & -\xi_{\mathbf{q}(t)}^{1,2} \\ -\xi_{\mathbf{q}(t)}^{2,1} & -\xi_{\mathbf{q}(t)}^{2,2} \end{pmatrix} \begin{pmatrix} \alpha^1(t) \\ \alpha^2(t) \end{pmatrix}$$

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$$|\psi(t)\rangle = \mathcal{T} \exp(\mathrm{i} \int \xi_{\mathbf{q}(t)} dt) |\psi(0)\rangle \equiv \mathrm{W} |\psi(0)\rangle$$

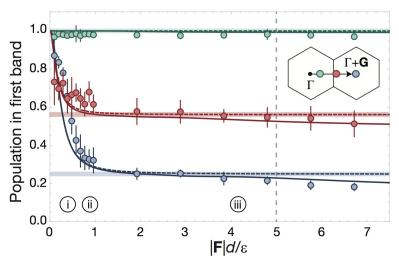
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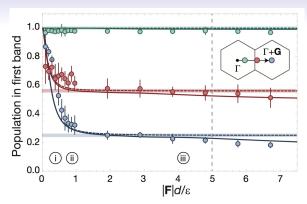
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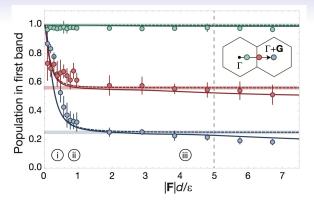
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ight)$

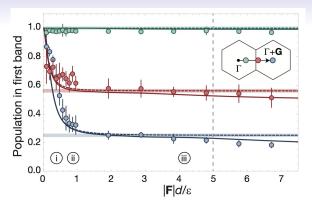
Measuring Wilson lines





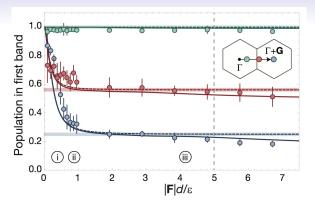


Matrix elements of Wilson line operator



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$$W_{\mathbf{Q} \to \mathbf{q}}^{m,n} = \langle \Phi_{\mathbf{q}}^m | e^{i(\mathbf{q} - \mathbf{Q}) \cdot \hat{\mathbf{r}}} | \Phi_{\mathbf{Q}}^n \rangle = \langle u_{\mathbf{q}}^m | u_{\mathbf{Q}}^n \rangle$$



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• saturation value $W^{11}_{\Gamma \to \mathbf{q}} = \langle u^1_{\mathbf{q}} | u^1_{\Gamma} \rangle$ of population after transport measures overlap between $|u\rangle$

• cell-periodic Bloch state as pseudo-spin

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- \bullet cell-periodic Bloch functions at a fixed reference quasimomentum \boldsymbol{Q}

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$$|1\rangle = |u_{\mathbf{Q}}^1\rangle \qquad |2\rangle = |u_{\mathbf{Q}}^2\rangle$$

Reconstructing band eigenstates

- cell-periodic Bloch state as pseudo-spin
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$$|1\rangle = |u_{\mathbf{Q}}^1\rangle \qquad |2\rangle = |u_{\mathbf{Q}}^2\rangle$$

• such that $|u_{\mathbf{q}}^1\rangle=\cos(\frac{\theta_{\mathbf{q}}}{2})|1\rangle+\sin(\frac{\theta_{\mathbf{q}}}{2})e^{\mathrm{i}\phi_{\mathbf{q}}}|2\rangle$

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Reconstructing band eigenstates

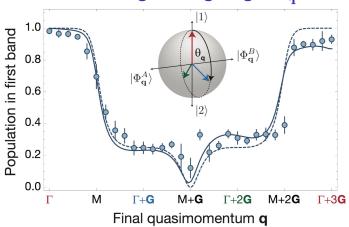
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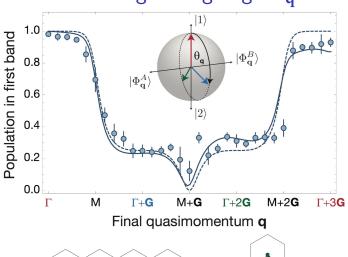
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- Throughout this work basis states are chosen at reference point $\mathbf{Q} = \Gamma$

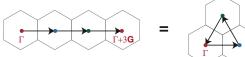
Measuring mixing angle $\theta_{\mathbf{q}}$

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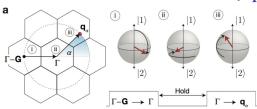




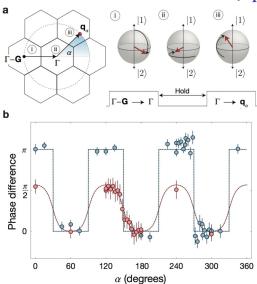


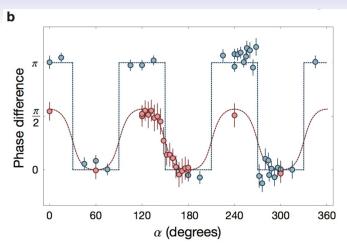
Measuring relative phase $\phi_{\mathbf{q}}$

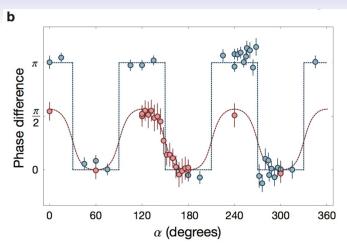
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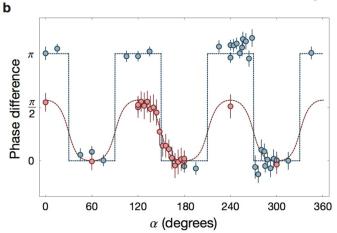


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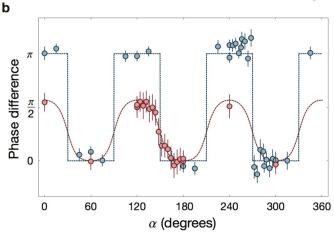




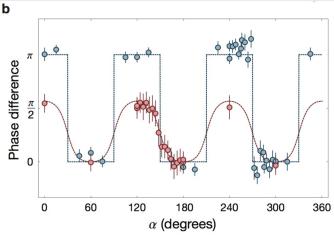




• AB-site degeneracy (blue) — π jump



- AB-site degeneracy (blue) π jump
- AB-site offset (red) by elliptically-polarized lattice beam
 continuous varying phase



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- AB-site offset (red) by elliptically-polarized lattice beam
 continuous varying phase
- 3-fold symmetry of system (no matter whether offset)



Determining Wilson line eigenvaluse

 $W_{\boldsymbol{q} \rightarrow \boldsymbol{q} + \boldsymbol{G}}$

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- the U(2) Wilson line $\Longrightarrow U(1)$ global phase multiplied by a SU(2) matrix.

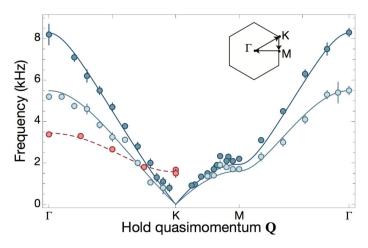
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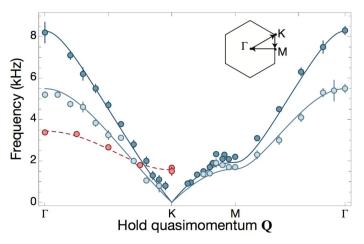
$$W_{q \to q + G}$$

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- experimental data analysis gives: $\xi = 1.03(2)\pi/3$
- theoretical expected value: $\xi = \pi/3$

Mapping dispersion relation



Mapping dispersion relation



by varying the reference quasimomentum ${f Q}$

Reference

- Immanuel Bloch et al., An Aharonov-Bohm interferometer for determining Bloch band topology, Science **347**, 288-292 (2015).
- Immanuel Bloch et al., Experimental reconstruction of Wilson lines in Bloch bands, arXiv:1509.02185v2 [cond-mat.quant-gas].
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- J. Zak, *Berrys phase for energy bands in solids*, Phys. Rev. Lett. **62**, 2747 (1989).

Thank you!