Derivation of the Schwarz Inequality (Cauchy-Schwarz Inequality)

Let $|\psi\rangle$ and $|\phi\rangle$ be two vectors in a Hilbert space. The Schwarz inequality states:

$$|\langle \psi | \phi \rangle|^2 \le \langle \psi | \psi \rangle \cdot \langle \phi | \phi \rangle$$

Step 1: Define a Vector

Define a new vector $|\chi\rangle$ as:

$$|\chi\rangle = |\phi\rangle - \lambda|\psi\rangle$$

for any complex number λ . Then:

$$\langle \chi | \chi \rangle = \langle \phi - \lambda \psi | \phi - \lambda \psi \rangle \ge 0$$

Step 2: Expand the Inner Product

$$\langle \chi | \chi \rangle = \langle \phi | \phi \rangle - \lambda^* \langle \phi | \psi \rangle - \lambda \langle \psi | \phi \rangle + |\lambda|^2 \langle \psi | \psi \rangle$$

Step 3: Choose Optimal λ

To minimize $\langle \chi | \chi \rangle$, choose:

$$\lambda = \frac{\langle \psi | \phi \rangle}{\langle \psi | \psi \rangle}$$

Step 4: Substitute λ Back In

Substituting into the expression for $\langle \chi | \chi \rangle$:

$$\langle \chi | \chi \rangle = \langle \phi | \phi \rangle - \frac{|\langle \psi | \phi \rangle|^2}{\langle \psi | \psi \rangle}$$

Since $\langle \chi | \chi \rangle \geq 0$, we get:

$$\langle \phi | \phi \rangle - \frac{|\langle \psi | \phi \rangle|^2}{\langle \psi | \psi \rangle} \ge 0$$

Multiplying both sides by $\langle \psi | \psi \rangle$, we obtain:

$$\langle \psi | \psi \rangle \langle \phi | \phi \rangle \ge |\langle \psi | \phi \rangle|^2$$

Final Result

$$\left| |\langle \psi | \phi \rangle|^2 \le \langle \psi | \psi \rangle \cdot \langle \phi | \phi \rangle \right|$$

Equality Condition

Equality holds if and only if:

$$|\phi\rangle = \lambda |\psi\rangle$$

i.e., the vectors $|\phi\rangle$ and $|\psi\rangle$ are linearly dependent.

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Objective

We aim to derive the generalized uncertainty relation for two Hermitian operators A and B, acting on a quantum state $|\psi\rangle$:

$$\sigma_A \sigma_B \ge \frac{1}{2} \left| \langle [A, B] \rangle \right|.$$
 (1)

This relation sets a lower bound on the product of standard deviations (uncertainties) of measurements of the observables A and B.

Step 1: Define Deviation Operators

Let the expectation values be:

$$\langle A \rangle = \langle \psi | A | \psi \rangle, \qquad \langle B \rangle = \langle \psi | B | \psi \rangle.$$

We define the deviation operators (also known as fluctuation operators):

$$\Delta A = A - \langle A \rangle, \qquad \Delta B = B - \langle B \rangle.$$

These represent how the operators differ from their mean values when acting on the state $|\psi\rangle$.

Step 2: Define Vectors in Hilbert Space

Define the following state vectors:

$$|\phi\rangle = \Delta A |\psi\rangle, \qquad |\chi\rangle = \Delta B |\psi\rangle.$$

These vectors represent how the state $|\psi\rangle$ is modified by the action of the deviation operators.

Step 3: Apply the Schwarz Inequality

The Schwarz inequality states:

$$|\langle \phi | \chi \rangle|^2 \le \langle \phi | \phi \rangle \cdot \langle \chi | \chi \rangle.$$

Substitute the vectors:

$$|\langle \psi | \Delta A \Delta B | \psi \rangle|^2 \le \langle \psi | (\Delta A)^2 | \psi \rangle \cdot \langle \psi | (\Delta B)^2 | \psi \rangle.$$

So we get:

$$|\langle \Delta A \Delta B \rangle|^2 \le \sigma_A^2 \sigma_B^2$$
,

where $\sigma_A^2 = \langle (\Delta A)^2 \rangle$, and similarly for σ_B .

Step 4: Expand the Expectation Value

We now expand the complex quantity $\langle \Delta A \Delta B \rangle$:

$$\Delta A \Delta B = \frac{1}{2} [\Delta A, \Delta B] + \frac{1}{2} \{\Delta A, \Delta B\}.$$

This identity splits the operator product into a commutator (antisymmetric) and an anticommutator (symmetric) part.

Take the expectation value:

$$\langle \Delta A \Delta B \rangle = \frac{1}{2} \langle [\Delta A, \Delta B] \rangle + \frac{1}{2} \langle \{\Delta A, \Delta B\} \rangle.$$

Now, note that:

$$[\Delta A, \Delta B] = [A, B],$$

because constants commute and have zero expectation value.

So:

$$\langle \Delta A \Delta B \rangle = \frac{1}{2} \langle [A,B] \rangle + \frac{1}{2} \langle \{\Delta A, \Delta B\} \rangle.$$

Step 5: Separate Real and Imaginary Parts

We now write:

$$|\langle \Delta A \Delta B \rangle|^2 = (\text{Re}\langle \Delta A \Delta B \rangle)^2 + (\text{Im}\langle \Delta A \Delta B \rangle)^2.$$

Using the earlier expansion:

$$\operatorname{Re}\langle \Delta A \Delta B \rangle = \frac{1}{2} \langle \{ \Delta A, \Delta B \} \rangle,$$
$$\operatorname{Im}\langle \Delta A \Delta B \rangle = \frac{1}{2i} \langle [A, B] \rangle.$$

So:

$$|\langle \Delta A \Delta B \rangle|^2 = \left(\frac{1}{2} \langle \{\Delta A, \Delta B\} \rangle \right)^2 + \left(\frac{1}{2i} \langle [A, B] \rangle \right)^2.$$

Step 6: Final Inequality

Combining with the Schwarz inequality:

$$\sigma_A^2\sigma_B^2 \geq \left(\frac{1}{2}\langle\{\Delta A,\Delta B\}\rangle\right)^2 + \left(\frac{1}{2i}\langle[A,B]\rangle\right)^2.$$

This is the Robertson–Schrödinger uncertainty relation.

Special Case: Heisenberg Uncertainty Principle

If the anticommutator term is zero or ignored, we obtain:

$$\sigma_A\sigma_B\geq \frac{1}{2}\left|\langle [A,B]\rangle\right|.$$

This is the standard uncertainty principle.