# Finite-Dimensional Irreducible Representations of the Weyl Algebra in Characteristic p>0

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7 August, 2025

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## What is Weyl Algebra?

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Then we quotient out  $k\langle x,y\rangle$  by the ideal generated by the element yx-xy-1. This is the first Weyl algebra  $A_1$ :

$$A_1 = k \langle x, y \rangle / \langle yx - xy - 1 \rangle$$
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#### Theorem 1

A basis for  $A_1$  is  $\{x^iy^j \mid i,j \geq 0\}$ .

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So Weyl Algebra is the algebra generated by position and momentum operators!

We can similarly define the n-th Weyl Algebra  $A_n$  by taking the "canonnical commutation relations" of n position and n momentum operators.

$$A_n = \frac{k \langle x_1, x_2, \dots, x_n, y_1 y_2, \dots, y_n \rangle}{\langle y_j x_i - x_i y_j - \delta_{ij} \mid 1 \leq i, j \leq n \rangle}.$$

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But in this talk, we'll mainly focus on  $A_1$ .

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## Representation

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#### Definition 1

A **representation** of an (associative) algebra A is a vector space V along with a homomorphism  $\rho: A \to \operatorname{End}(V)$ .

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#### Definition 1

A **representation** of an (associative) algebra A is a vector space V along with a homomorphism  $\rho: A \to \operatorname{End}(V)$ .

**Abuse of notation:** for  $a \in A$  and  $\mathbf{v} \in V$ ,  $\rho(a) : V \to V$  is a linear map, so that  $\rho(a)\mathbf{v} \in V$ . We shall often write  $a\mathbf{v}$  instead of  $\rho(a)\mathbf{v}$ .

#### Definition 2

A **subrepresentation** of a representation V is a subspace  $U \subseteq V$  such that  $\rho(a)$   $U \subseteq U$  for every  $a \in A$ .

In such a case, we call U invariant under the action of A.

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#### Definition 3

We call a representation V **irreducible** if the only subrepresentations are 0 and V itself.

#### Definition 4

Let  $\rho_1:A\to \operatorname{End}(V_1)$  and  $\rho_2:A\to \operatorname{End}(V_2)$  be two representations of A. Then a **homomorphism of representations** is a linear map  $\phi:V_1\to V_2$  such that

$$\phi \circ \rho_1(a) = \rho_2(a) \circ \phi,$$

for every  $a \in A$ . In other words, the following diagram commutes for every  $a \in A$ :

$$egin{array}{ccc} V_1 & \stackrel{\phi}{\longrightarrow} & V_2 \ 
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An isomorphism of representations is an invertible homomorphism of representations.

## Proposition 2

Let A be an algebra over a field k, and  $\phi: V_1 \to V_2$  is a nonzero homomorphism of representations.

- **1** If  $V_1$  is irreducible,  $\phi$  is injective.
- ② If  $V_2$  is irreducible,  $\phi$  is surjective.
- $\odot$  If both are irreducible,  $\phi$  is an isomorphism.

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#### Proof.

Just look at Ker  $\phi$  (or im  $\phi$ ), and they are subrepresentations. So irreducibility forces that Ker  $\phi = 0$  (or im  $\phi = V_2$ ).

## Corollary 3

Let A be an algebra over an **algebraically closed** field k, and V is a finite dimensional irrep. If  $\phi:V\to V$  is a homomorphism of representations, then  $\phi=c\,\mathbbm{1}_V$  for some  $c\in k$ .

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#### Proof.

If  $\lambda$  is an eigenvalue of  $\phi$ , then  $\phi - \lambda \mathbb{1}_V$  is a homomorphism of algebras, which is not an isomorphism. So it has to be 0.

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#### Corollary 4

Let A be an algebra over an algebraically closed field k, and V is a finite dimensional irrep. If  $a \in Z(A)$ , then  $\rho(a) = c \mathbb{1}_V$  for some  $c \in k$ .

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In QM, we work with the nicest field ever,  $\mathbb{C}$ , which is algebraically closed and has characteristic 0.

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By Stone-Von Neumann theorem, every irreducible representation of the position and momentum operators are **unitarily equivalent** to the usual ones: position operator is multiplication by x, momentum operator is  $\frac{\partial}{\partial x}$ .

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By Stone-Von Neumann theorem, every irreducible representation of the position and momentum operators are **unitarily equivalent** to the usual ones: position operator is multiplication by x, momentum operator is  $\frac{\partial}{\partial x}$ .

So the only irrep of  $A_1 = k \langle x, y \rangle / \langle yx - xy - 1 \rangle$  looks like

$$\rho(y) = \frac{\mathrm{d}}{\mathrm{d}t}, \quad \rho(x) = t \tag{1}$$

acting on an infinite dimensional space, maybe k[t].

Is there any finite dimensional irreps?

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Let V be a finite dimensional irrep. Since yx - xy = 1 in  $A_1$ ,

$$\rho(y)\rho(x) - \rho(x)\rho(y) = \mathbb{1}_V.$$
 (2)

Taking trace on both sides, we get  $0 = \dim V$ .

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So, the finite dimensional irreps are 0-dimensional.

Unless, a positive number is **EQUAL TO** 0.

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As before,  $0 = \dim V$  indicates that  $p \mid \dim V$ , i.e.  $\dim V \ge p$ .

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Let's now look at the Center of  $A_1$ .

## Proposition 5

$$[x^{i}y^{j}, x] = jx^{i}y^{j-1} [x^{i}y^{j}, y] = -ix^{i-1}y^{j}$$
 (3)

#### Proof.

Induction on i and j.



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### Proposition 5

$$\begin{aligned}
[x^i y^j, x] &= j x^i y^{j-1} \\
[x^i y^j, y] &= -i x^{i-1} y^j
\end{aligned} \tag{3}$$

#### Proof.

Induction on i and j.

So  $k[x^p, y^p] \in Z(A_1)$ . In fact,  $Z(A_1) = k[x^p, y^p]$ .

#### Definition 5

Let  $V \neq 0$  be a representation of A. We say that a vector  $\mathbf{v} \in V$  is **cyclic** if it generates V, i.e.  $A\mathbf{v} = \{\rho(a)\mathbf{v} \mid a \in A\} = V$ .

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#### Theorem 6

A representation V of A is irreducible if and only if every nonzero vector  $\mathbf{v} \in V$  is cyclic.

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A representation V of A is irreducible if and only if every nonzero vector  $\mathbf{v} \in V$  is cyclic.

#### Proof.

If **v** is not cyclic, then A**v** := { $\rho$  (a) **v** |  $a \in A$ } is a proper subrepresentation.

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#### Theorem 6

A representation V of A is irreducible if and only if every nonzero vector  $\mathbf{v} \in V$  is cyclic.

#### Proof.

If  $\mathbf{v}$  is not cyclic, then  $A\mathbf{v}:=\{\rho\left(a\right)\mathbf{v}\mid a\in A\}$  is a proper subrepresentation. Conversely, if  $W\subseteq V$  is a subrepresentation, then for  $\mathbf{w}\in W,\ V=A\mathbf{w}\subseteq W$ , proving that W=V.

 $\rho(y^p) = c \, \mathbb{1}_V$  since  $y^p \in Z(A_1)$ . Let **v** be an eigenvector of  $\rho(y)$  with the eigenvalue  $\lambda$  (where  $\lambda^p = c$ ). Then **v** is cyclic!

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$$\operatorname{span}\left\{\rho\left(x^{i}y^{j}\right)\mathbf{v}\mid i,j\geq0\right\}=V$$

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Therefore, dim V = p.

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$$\mathbf{v}_2 = \rho\left(P(x)\right)\mathbf{v}_1,\tag{4}$$

for a polynomial P(x) of degree at most p-1.

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for a polynomial P(x) of degree at most p-1. Then

$$[y, P(x)] = P'(x). \tag{5}$$

$$\rho\left(P'(x)\right)\mathbf{v}_{1} = \rho(y)\rho\left(P(x)\right)\mathbf{v}_{1} - \rho\left(P(x)\right)\rho(y)\mathbf{v}_{1}$$

$$\rho(P'(x)) \mathbf{v}_1 = \rho(y) \rho(P(x)) \mathbf{v}_1 - \rho(P(x)) \rho(y) \mathbf{v}_1$$
  
= \rho(y) \mathbf{v}\_2 - \rho(P(x)) \lambda \mathbf{v}\_1

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$$= \mathbf{0}.$$

Therefore,  $\rho(P'(x)) = 0$ , so that  $\mathbf{v}_2$  is a constant multiple of  $\mathbf{v}_1$ .

Therefore, not only does  $\rho(y)$  have only one eigenvalue, it has only one linearly independent eigenvector.

Therefore, not only does  $\rho(y)$  have only one eigenvalue, it has only one linearly independent eigenvector. So in a suitable basis, it looks like a Jordan block:

$$\rho(y) = \begin{bmatrix} \lambda & 1 & 0 & \cdots & 0 \\ 0 & \lambda & 1 & \cdots & 0 \\ 0 & 0 & \lambda & \cdots & 0 \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & 0 & \cdots & \lambda \end{bmatrix}$$
 (6)

So all that is left is to compute is  $\rho(x)$ . But we don't quite know what the basis is yet. So we're gonna stick to the previous basis that we showed earlier:

$$B = \left\{ \rho\left(x^{i}\right)\mathbf{v} \mid 0 \leq i \leq p-1 \right\} = \left\{\mathbf{v}, \rho(x)\mathbf{v}, \rho\left(x^{2}\right)\mathbf{v}, \dots, \rho\left(x^{p-1}\right)\mathbf{v} \right\}.$$

How  $\rho(x)$  acts on these basis vectors is pretty simple. Since  $\rho(x^p)$  is a scalar, say  $\rho(x^p) = \mu \mathbb{1}_V$ ,

$$\rho(x) \left[\rho(x^{i}) \mathbf{v}\right] = \begin{cases} \rho(x^{i+1}) \mathbf{v} & \text{if } i \neq p-1\\ \mu \mathbf{v} & \text{if } i = p-1 \end{cases}$$
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So  $\rho(x)$  shifts all the basis vectors once to the next, and for the final one, it scales by  $\mu$ . So the matrix representation of  $\rho(x)$  in this (ordered) basis  $B = \{\mathbf{v}, \rho(x)\mathbf{v}, \rho(x^2)\mathbf{v}, \dots, \rho(x^{p-1})\mathbf{v}\}$  is

$$\begin{bmatrix} 0 & 0 & 0 & \cdots & 0 & \mu \\ 1 & 0 & 0 & \cdots & 0 & 0 \\ 0 & 1 & 0 & \cdots & 0 & 0 \\ 0 & 0 & 1 & \cdots & 0 & 0 \\ \vdots & \vdots & \vdots & \ddots & \vdots & \vdots \\ 0 & 0 & 0 & \cdots & 1 & 0 \end{bmatrix}$$

(8)

Finally, we just need to compute  $\rho(y)$ 's action on these basis vectors. Since  $[y,x^i]=ix^{i-1}$ ,

$$\rho(y) \left[ \rho(x^{i}) \mathbf{v} \right] = \rho\left( \left[ y, x^{i} \right] \right) \mathbf{v} + \rho(x^{i}) \rho(y) \mathbf{v}$$
$$= i\rho(x^{i-1}) \mathbf{v} + \lambda \rho(x^{i}) \mathbf{v}. \tag{9}$$

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$$\begin{bmatrix}
\lambda & 1 & 0 & \cdots & 0 & 0 \\
0 & \lambda & 2 & \cdots & 0 & 0 \\
0 & 0 & \lambda & \cdots & 0 & 0 \\
\vdots & \vdots & \vdots & \ddots & \vdots & \vdots \\
0 & 0 & 0 & \cdots & \lambda & p-1 \\
0 & 0 & 0 & \cdots & 0 & \lambda
\end{bmatrix}$$
(10)

So in characteristic p, the irreducible representations of  $A_1$  are p-dimensional, and they look like

$$\rho(x) = \begin{bmatrix} 0 & 0 & 0 & \cdots & 0 & \mu \\ 1 & 0 & 0 & \cdots & 0 & 0 \\ 0 & 1 & 0 & \cdots & 0 & 0 \\ 0 & 0 & 1 & \cdots & 0 & 0 \\ \vdots & \vdots & \vdots & \ddots & \vdots & \vdots \\ 0 & 0 & 0 & \cdots & 1 & 0 \end{bmatrix}, \quad \rho(y) = \lambda \, \mathbb{1} + \begin{bmatrix} 0 & 1 & 0 & \cdots & 0 & 0 \\ 0 & 0 & 2 & \cdots & 0 & 0 \\ 0 & 0 & 0 & \cdots & 0 & 0 \\ \vdots & \vdots & \vdots & \ddots & \vdots & \vdots \\ 0 & 0 & 0 & \cdots & 0 & p-1 \\ 0 & 0 & 0 & \cdots & 0 & 0 \end{bmatrix}$$

Looks familiar?

Consider the space  $\mathbb{C}[t]$  of all polynomials over  $\mathbb{C}$ . In this infinite dimensional space, the matrix representation of  $\frac{\mathrm{d}}{\mathrm{d}t}$  in the ordered basis  $\{1,t,t^2,t^3,\ldots\}$  looks like this

```
\begin{bmatrix} 0 & 1 & 0 & 0 & 0 & \cdots \\ 0 & 0 & 2 & 0 & 0 & \cdots \\ 0 & 0 & 0 & 3 & 0 & \cdots \\ 0 & 0 & 0 & 0 & 4 & \cdots \\ \vdots & \vdots & \vdots & \vdots & \vdots & \ddots \end{bmatrix}
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Consider the space  $\mathbb{C}[t]$  of all polynomials over  $\mathbb{C}$ . In this infinite dimensional space, the matrix representation of  $\frac{\mathrm{d}}{\mathrm{d}t}$  in the ordered basis  $\{1,t,t^2,t^3,\ldots\}$  looks like this

and the "matrix" representation of multiplication by t looks like this

```
\begin{bmatrix} 0 & 0 & 0 & 0 & \cdots \\ 1 & 0 & 0 & 0 & \cdots \\ 0 & 1 & 0 & 0 & \cdots \\ 0 & 0 & 1 & 0 & \cdots \\ 0 & 0 & 0 & 1 & \cdots \\ \vdots & \vdots & \vdots & \vdots & \ddots \end{bmatrix}
```

$$\rho(x) = \begin{bmatrix}
0 & 0 & 0 & \cdots & 0 & \mu \\
1 & 0 & 0 & \cdots & 0 & 0 \\
0 & 1 & 0 & \cdots & 0 & 0 \\
0 & 0 & 1 & \cdots & 0 & 0 \\
\vdots & \vdots & \vdots & \ddots & \vdots & \vdots \\
0 & 0 & 0 & \cdots & 1 & 0
\end{bmatrix}$$
(11)

So  $\rho(x)$  looks exactly like multiplication by t on k[t], except for  $t^p$  being identified with the scalar  $\mu$ .

$$\rho(x) = \begin{bmatrix}
0 & 0 & 0 & \cdots & 0 & \mu \\
1 & 0 & 0 & \cdots & 0 & 0 \\
0 & 1 & 0 & \cdots & 0 & 0 \\
0 & 0 & 1 & \cdots & 0 & 0 \\
\vdots & \vdots & \vdots & \ddots & \vdots & \vdots \\
0 & 0 & 0 & \cdots & 1 & 0
\end{bmatrix}$$
(11)

So  $\rho(x)$  looks exactly like multiplication by t on k[t], except for  $t^p$  being identified with the scalar  $\mu$ . So our representation space is gonna be  $k[t]/\langle t^p - \mu \rangle$ .

$$\rho(y) = \lambda \, \mathbb{1} + \begin{bmatrix} 0 & 1 & 0 & \cdots & 0 & 0 \\ 0 & 0 & 2 & \cdots & 0 & 0 \\ 0 & 0 & 0 & \cdots & 0 & 0 \\ \vdots & \vdots & \vdots & \ddots & \vdots & \vdots \\ 0 & 0 & 0 & \cdots & 0 & p-1 \\ 0 & 0 & 0 & \cdots & 0 & 0 \end{bmatrix}$$

$$(12)$$

And  $\rho(y)$  looks exactly like  $\lambda$  plus  $\frac{\mathrm{d}}{\mathrm{d}t}$  on  $k[t]/\langle t^p - \mu \rangle$ .

To summarize, the irreducible representations of  $A_1$  in characteristic p are on the representation space  $k[t]/\langle t^p - \mu \rangle$ , where

$$\rho(x) = \text{multiplication by } t$$

$$\rho(y) = \lambda + \frac{\mathrm{d}}{\mathrm{d}t}$$
(13)

for  $\lambda, \mu \in k$ .

### **Concluding Remarks**

So, in finite dimensions and char k=p>0 as well, the representation of Weyl Algebra looks exactly like those for  $k=\mathbb{C}$ .

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What's next? Defining Quantum Mechanics over an algebraically closed field k with char k = p > 0?

### Concluding Remarks

So, in finite dimensions and char k=p>0 as well, the representation of Weyl Algebra looks exactly like those for  $k=\mathbb{C}$ .

What's next? Defining Quantum Mechanics over an algebraically closed field k with char k = p > 0?

Apparently some people are trying to do that! And it might have some cool applications in Quantum Computing! See references!

### References

- Introduction to Representation Theory, by Etingof et al.
- ② Categorical Quantum Computing with Finite Fields, by Matthew Varughese. https://www.cs.ox.ac.uk/people/bob.coecke/Varughese.pdf
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### Thank you for joining!

The slides are available in my webpage https://atonurc.github.io/assets/weyl\_talk.pdf

