

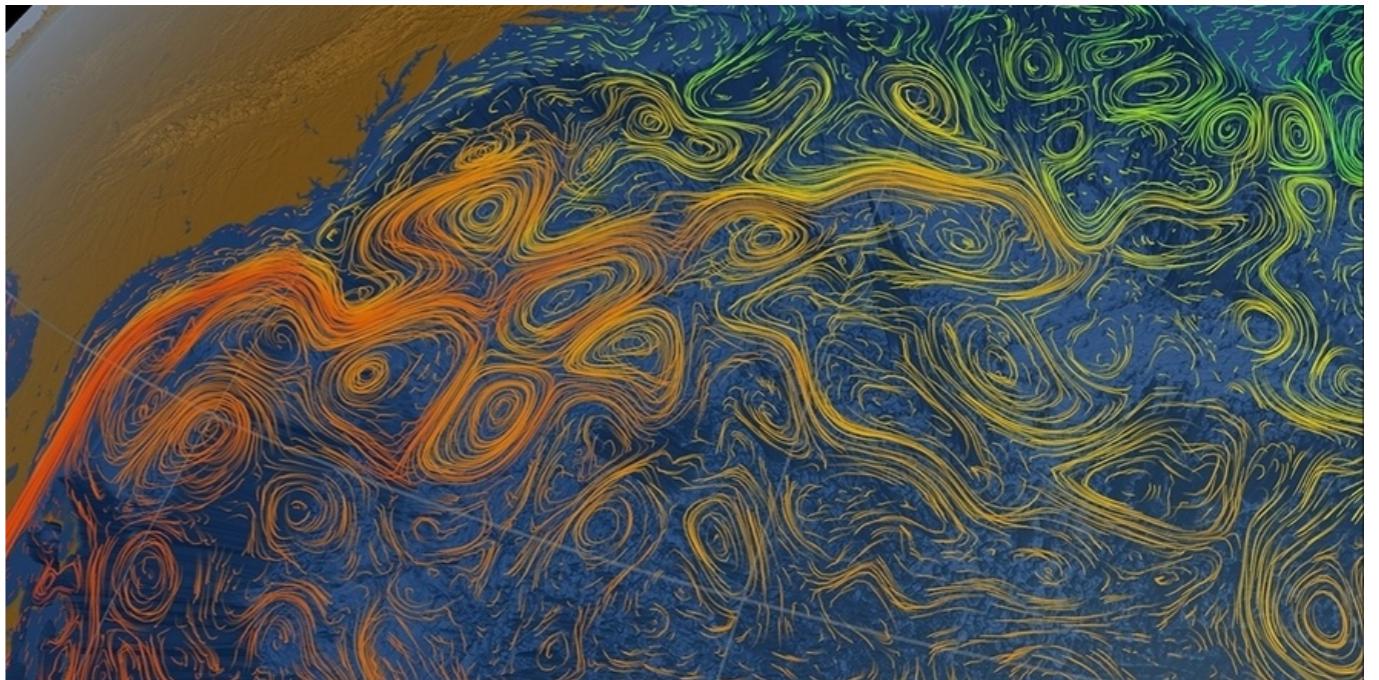
The Phase Space of the Three-Vortex Problem and its Application to Vortex-Dipole Scattering

Atul Anurag, Roy H. Goodman

SIAM Conference on Nonlinear Waves and
Coherent Structures
June 24, 2024

Point Vortices

- A point vortex is a solution to the 2D Euler equations with vorticity concentrated at a single point.
- It is characterized by two quantities: circulation (strength and orientation) and position.
- This model, derived by Helmholtz, idealizes Euler's equations for the motion of interacting point vortices.
- It is a Hamiltonian system, enabling the use of Hamiltonian theory and tools for its study.



Vortices formation while oceans and atmospheres move heat around on Earth and other planetary bodies.

Kirchhoff's Hamiltonian Representation

Consider N -vortices located at $\mathbf{r}_i = (x_i, y_i)$ with the strength Γ_i , then the system of ODEs describing the N -vortex motion can be described by the Hamiltonian,

$$\mathcal{H}(\mathbf{r}_1, \mathbf{r}_2, \mathbf{r}_3, \dots, \mathbf{r}_N) = -\frac{1}{4\pi} \sum_{1 \leq i < j \leq N} \Gamma_i \Gamma_j \log \|\mathbf{r}_i - \mathbf{r}_j\|^2.$$

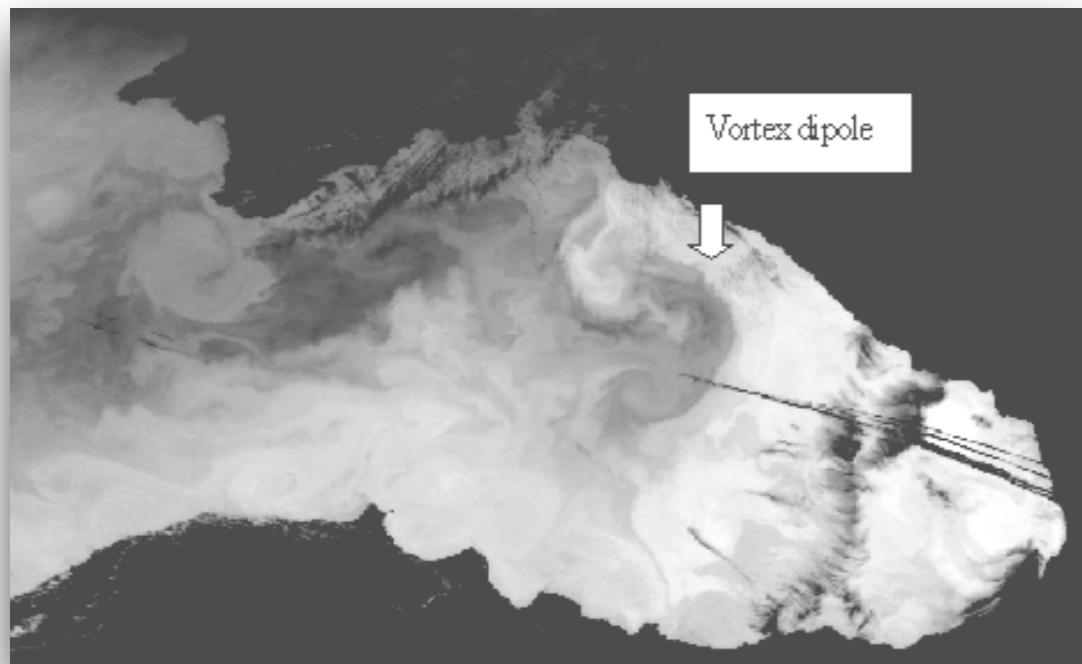
System of $2N$ point vortex equations:

$$\Gamma_i \frac{dx_i}{dt} = \frac{\partial \mathcal{H}}{\partial y_i}, \quad \Gamma_i \frac{dy_i}{dt} = -\frac{\partial \mathcal{H}}{\partial x_i}.$$

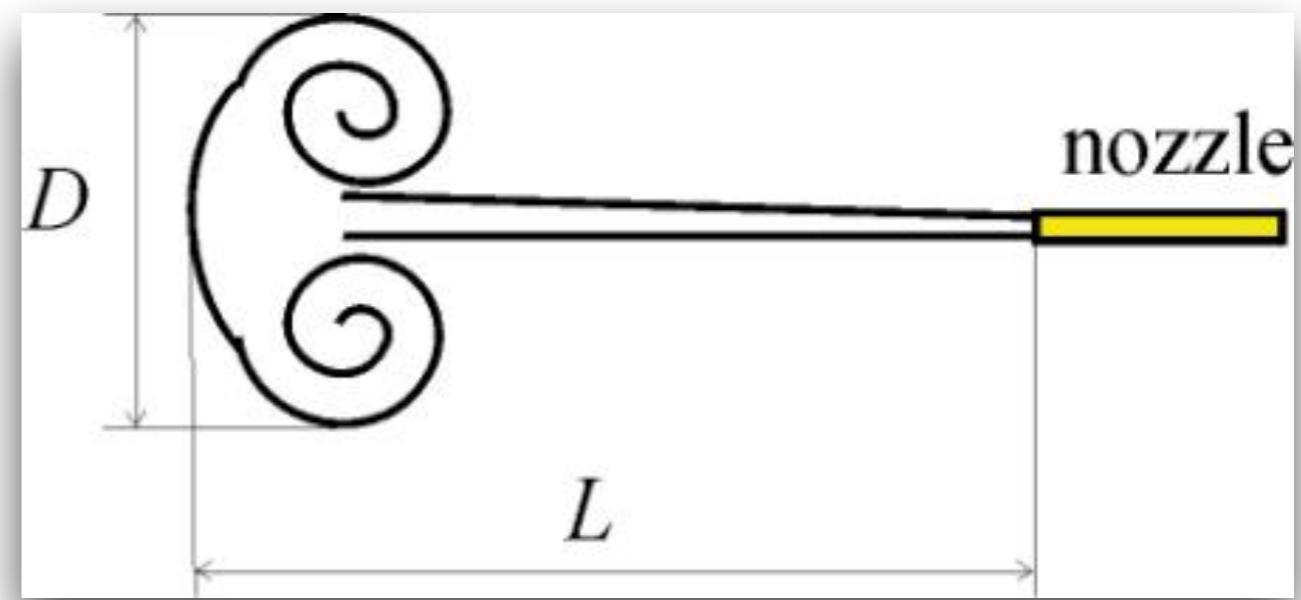
Vortex Dipole

A point vortex dipole consists of two vortices with equal strength but opposite circulation placed close together, forming a dipole.

- Vortices induce motion in each other, making the dipole move as a pair.
- The dipole travels in a straight line perpendicular to the line connecting the two vortices.



Vortex dipole formed by cold water (dark) in the central part of the sea.



Sketch of vortex dipole

Gröbli's Ancient Coordinate System

Specielle Probleme

über die

Bewegung geradliniger paralleler Wirbelfäden.

Inaugural-Dissertation

zur

Erlangung der Doctorwürde

der hohen philosophischen Facultät

der

Georg-August-Universität zu Göttingen

vorgelegt von

WALTER GRÖBLI,

von Oberuzwil, Canton St. Gallen.

The evolution equations derived
by Gröbli

$$\frac{d}{dt} \begin{pmatrix} \ell_{23}^2 \\ \ell_{31}^2 \\ \ell_{12}^2 \end{pmatrix} = 4\sigma A \begin{pmatrix} \Gamma_1 (\ell_{12}^{-2} - \ell_{31}^{-2}) \\ \Gamma_2 (\ell_{23}^{-2} - \ell_{12}^{-2}) \\ \Gamma_3 (\ell_{31}^{-2} - \ell_{23}^{-2}) \end{pmatrix}.$$

ℓ_{ij} be the distance between vortices i and j

A be the area of the triangle formed by the vortices

$\sigma = \pm 1$ indicate the orientation of the triangle spanned
by the three vortices.

Zürich,

Druck von Zürcher und Furrer.

1877.

(6)

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- **Syng (1949)** Trilinear coordinates.
- **Aref (1979)** Rederivation of Gröbli's solution, and use of trilinear coordinates to interpret the results.

Trilinear Coordinate System

Trilinear Coordinate System

For $L \neq 0$, A_{ref} defines new variables

$$b_1 = \frac{l_{23}^2}{\Gamma_1 L}, \quad b_2 = \frac{l_{31}^2}{\Gamma_2 L}, \quad b_3 = \frac{l_{12}^2}{\Gamma_3 L},$$

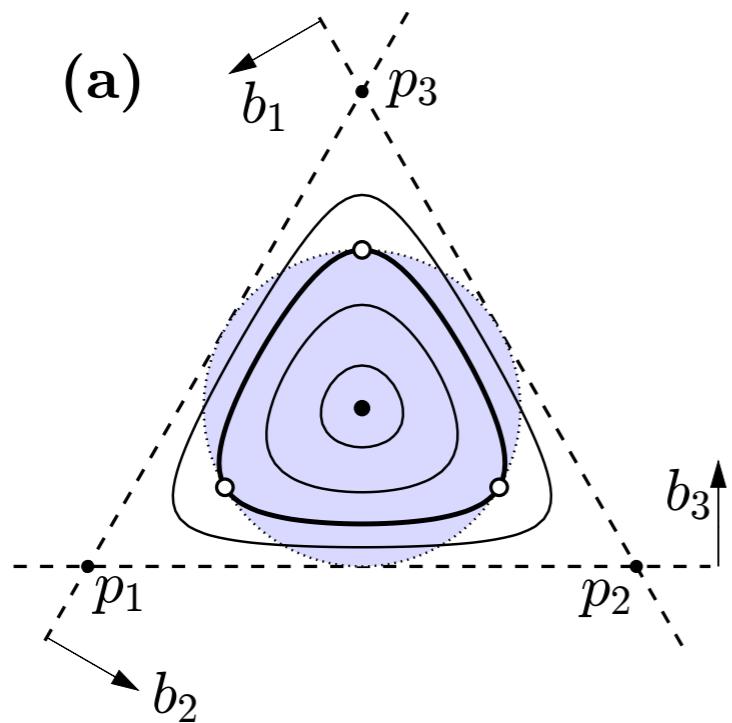
such that $b_1 + b_2 + b_3 = 3$.

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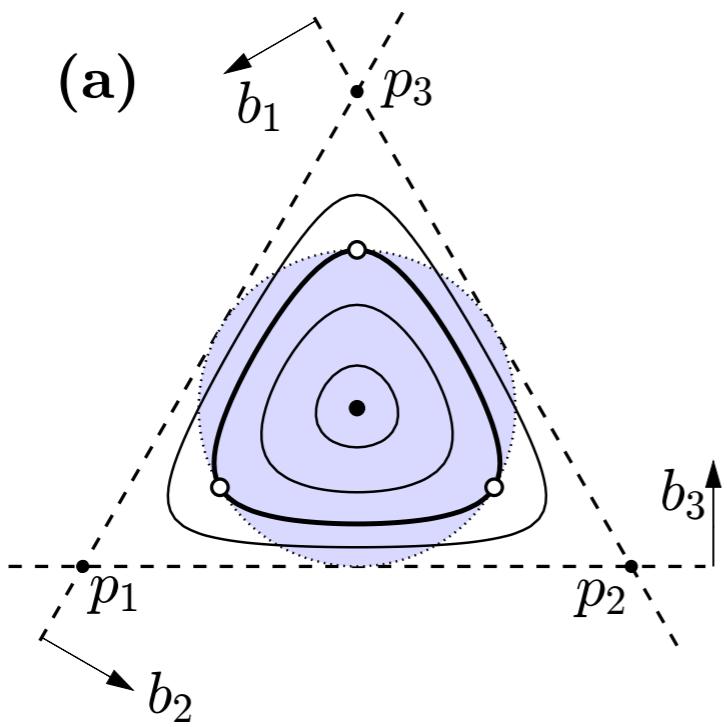
Vortices of circulation $(1, 1, 1)$

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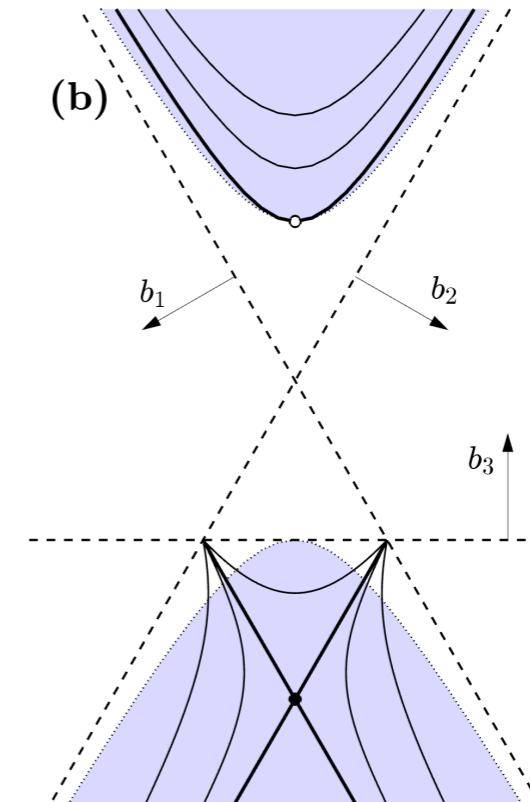
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Vortices of circulation $(1, 1, 1)$



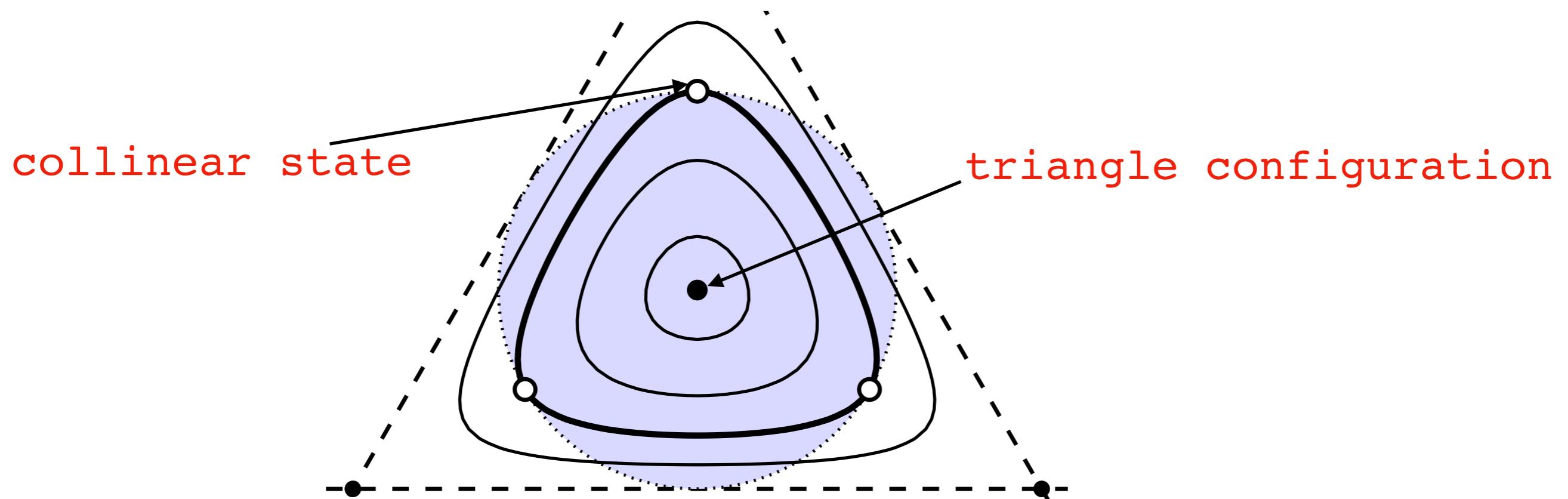
Vortices of circulation $(1, 1, -1)$

Limitations of previous coordinates system

- Not all triples satisfies the **triangular inequality constraint**.
- Introduces **coordinate singularities** at all collinear configuration
- The **linearization fails** at the boundaries due to collinear singularities.

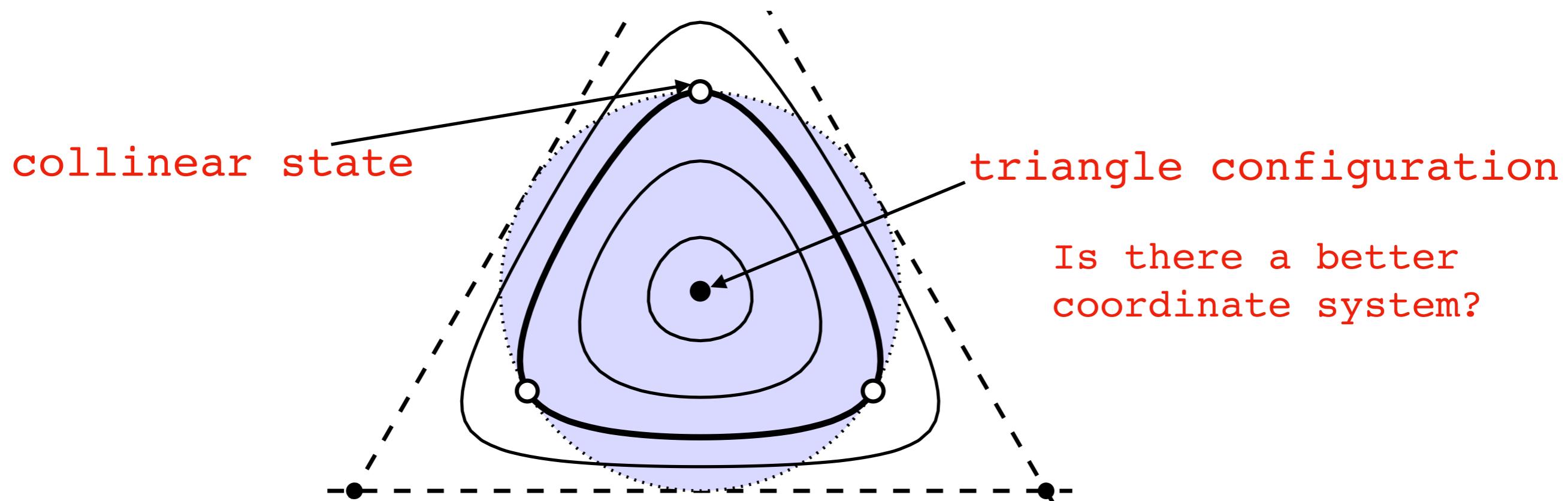
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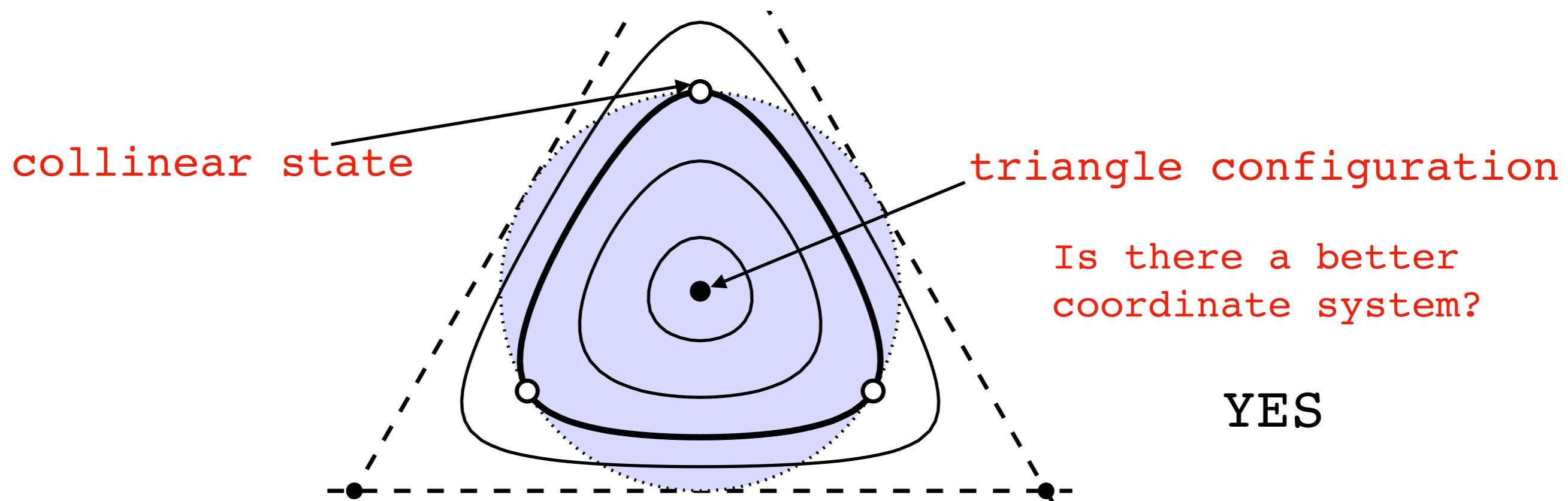
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Jacobi coordinates

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Jacobi Coordinates (classical)

The first transformation

$$\tilde{\mathbf{r}}_1 = \mathbf{r}_1 - \mathbf{r}_2; \quad \tilde{\Gamma}_1 = \frac{\Gamma_1 \Gamma_2}{\Gamma_1 + \Gamma_2};$$

$$\tilde{\mathbf{r}}_2 = \frac{\Gamma_1 \mathbf{r}_1 + \Gamma_2 \mathbf{r}_2}{\Gamma_1 + \Gamma_2}; \quad \tilde{\Gamma}_2 = \Gamma_1 + \Gamma_2;$$

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The second transformation

$$\begin{aligned}R_1 &= \tilde{\mathbf{r}}_1; & \kappa_1 &= \tilde{\Gamma}_1; \\ R_2 &= \tilde{\mathbf{r}}_2 - \tilde{\mathbf{r}}_3; & \kappa_2 &= \frac{\tilde{\Gamma}_2 \tilde{\Gamma}_3}{\tilde{\Gamma}_2 + \tilde{\Gamma}_3}; \\ R_3 &= \frac{\tilde{\Gamma}_2 \tilde{\mathbf{r}}_2 + \tilde{\Gamma}_3 \tilde{\mathbf{r}}_3}{\tilde{\Gamma}_2 + \tilde{\Gamma}_3}; & \kappa_3 &= \tilde{\Gamma}_2 + \tilde{\Gamma}_3.\end{aligned}$$

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Transformed coordinates in terms
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where $\Gamma_1 + \Gamma_2 + \Gamma_3 \neq 0$

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Jacobi
transformation is
applied recursively.

Nambu Bracket

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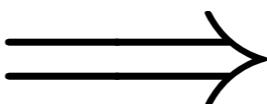
- Generalization of Poisson Brackets

Nambu Bracket

- Generalization of Poisson Brackets
- Considers a canonical triplet

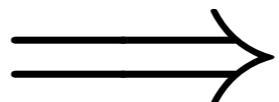
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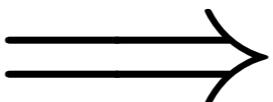
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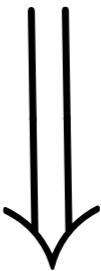
$(R_1, R_2) \rightarrow (X, Y, Z, \Theta)$
 $(\Theta \text{ conserved})$

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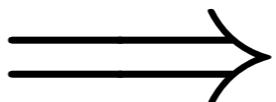


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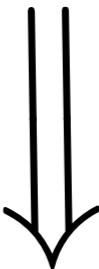


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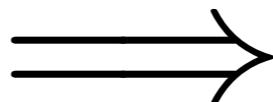
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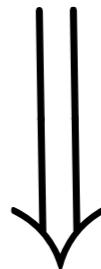
The system with coordinates (X, Y, Z) ,
the evolution equations

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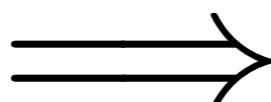


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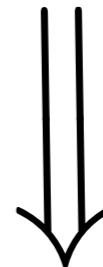
$$\frac{d}{dt} \begin{pmatrix} X \\ Y \\ Z \end{pmatrix} = \nabla C \times \nabla H$$

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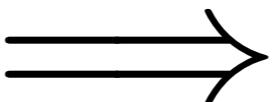
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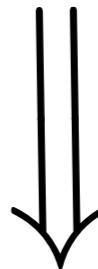
$$\frac{dF(X, Y, Z)}{dt} = \{F, G\}_C$$

Nambu Bracket

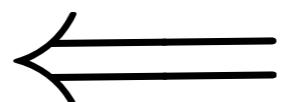
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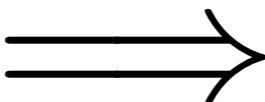


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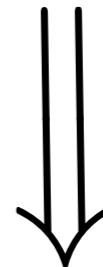
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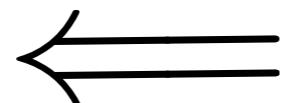
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The system with coordinates (X, Y, Z) ,
the evolution equations

The reduced system
is free from
Coordinate
singularities



$$\frac{d}{dt} \begin{pmatrix} X \\ Y \\ Z \end{pmatrix} = \nabla C \times \nabla H$$

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Application to the Three-Vortex System with Arbitrary Circulation

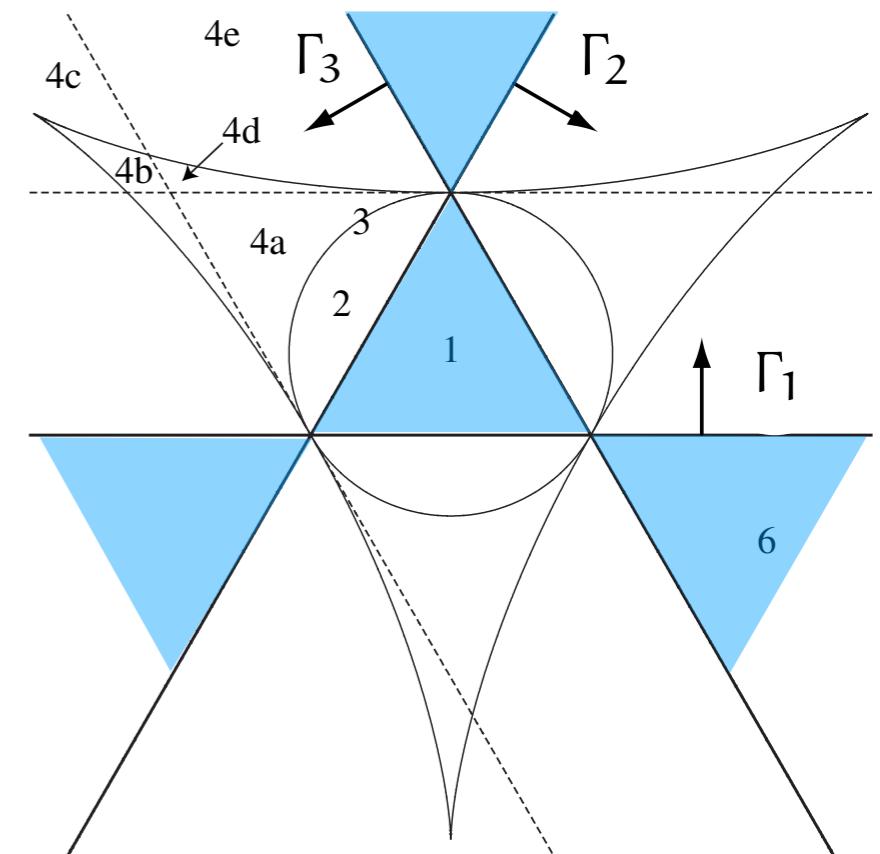
Application to the Three-Vortex System with Arbitrary Circulation

The three-vortex system conserve an energy when $\kappa_2 > 0$

$$\begin{aligned} H(X, Y, Z, \Theta) = & -\frac{\Gamma_1 \Gamma_2}{2} \log \left(\frac{Z + \Theta}{2\kappa_1} \right) \\ & -\frac{\Gamma_2 \Gamma_3}{2} \log \left(\frac{\Theta - Z}{2\kappa_2} + \frac{\kappa_1(Z + \Theta)}{2\Gamma_2^2} - \frac{kX}{\Gamma_2} \right) \\ & -\frac{\Gamma_1 \Gamma_3}{2} \log \left(\frac{\Theta - Z}{2\kappa_2} + \frac{\kappa_1(Z + \Theta)}{2\Gamma_1^2} + \frac{kX}{\Gamma_1} \right) \end{aligned}$$

where $k^2 = \frac{\kappa_1}{\kappa_2}$.

$$\Theta^2 = Z^2 + X^2 + Y^2$$



Barycentric coordinates for circulation space when $\Gamma_1 + \Gamma_2 + \Gamma_3 = 1$

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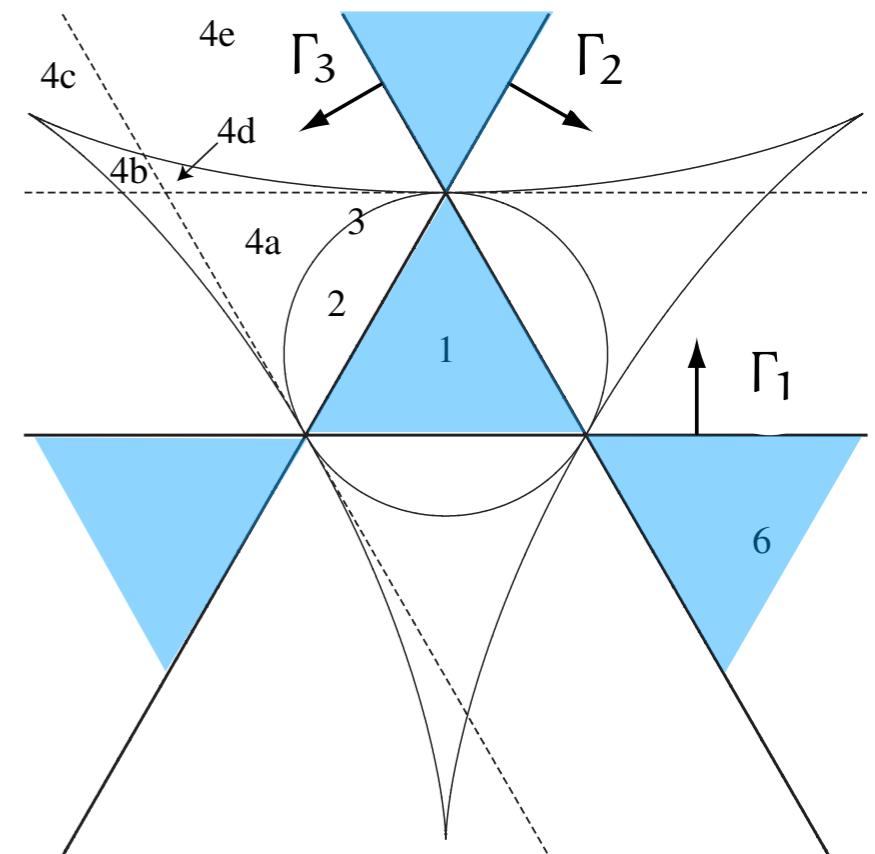
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The three-vortex system conserve an energy when $\kappa_2 < 0$

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$$\text{where } l^2 = \frac{-\kappa_1}{\kappa_2}.$$

$$\Theta^2 = Z^2 - X^2 - Y^2$$



Barycentric coordinates for circulation space when $\Gamma_1 + \Gamma_2 + \Gamma_3 = 1$

Application to the Three-Vortex System with Arbitrary Circulation

The three-vortex system conserve an energy when $\kappa_2 > 0$

$$\begin{aligned} H(X, Y, Z, \Theta) = & -\frac{\Gamma_1 \Gamma_2}{2} \log \left(\frac{Z + \Theta}{2\kappa_1} \right) \\ & -\frac{\Gamma_2 \Gamma_3}{2} \log \left(\frac{\Theta - Z}{2\kappa_2} + \frac{\kappa_1(Z + \Theta)}{2\Gamma_2^2} - \frac{kX}{\Gamma_2} \right) \\ & -\frac{\Gamma_1 \Gamma_3}{2} \log \left(\frac{\Theta - Z}{2\kappa_2} + \frac{\kappa_1(Z + \Theta)}{2\Gamma_1^2} + \frac{kX}{\Gamma_1} \right) \end{aligned}$$

$$\text{where } k^2 = \frac{\kappa_1}{\kappa_2}.$$

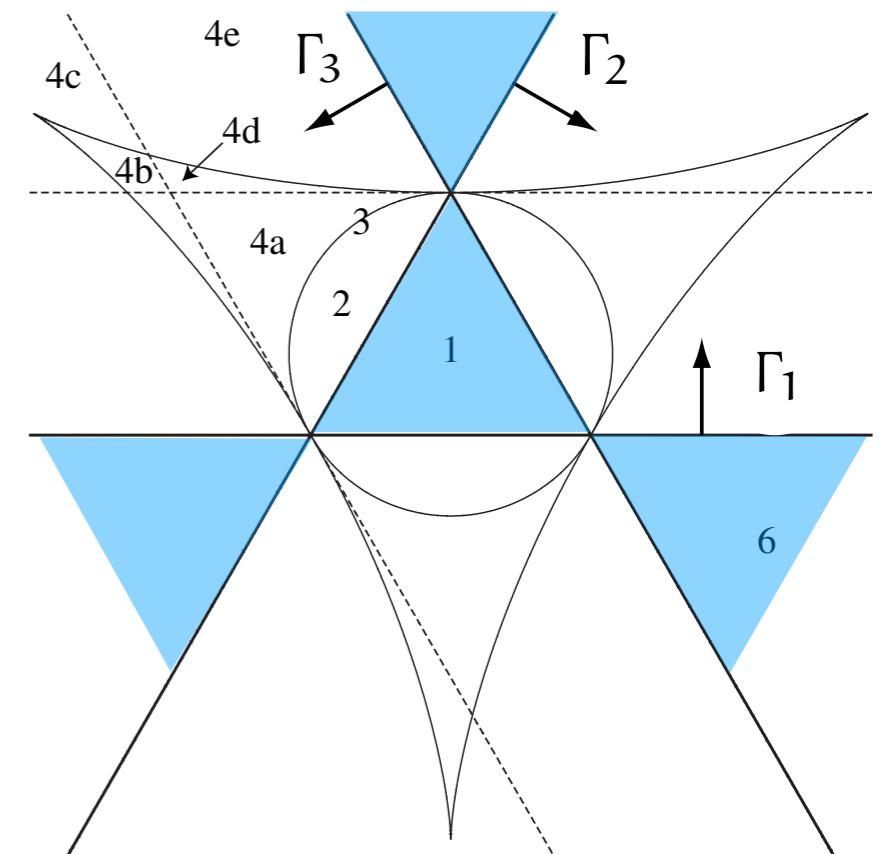
$$\Theta^2 = Z^2 + X^2 + Y^2$$

The three-vortex system conserve an energy when $\kappa_2 < 0$

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$$\Theta^2 = Z^2 - X^2 - Y^2$$



Barycentric coordinates for circulation space when $\Gamma_1 + \Gamma_2 + \Gamma_3 = 1$

Dynamics based of the sign of κ_2 :

Application to the Three-Vortex System with Arbitrary Circulation

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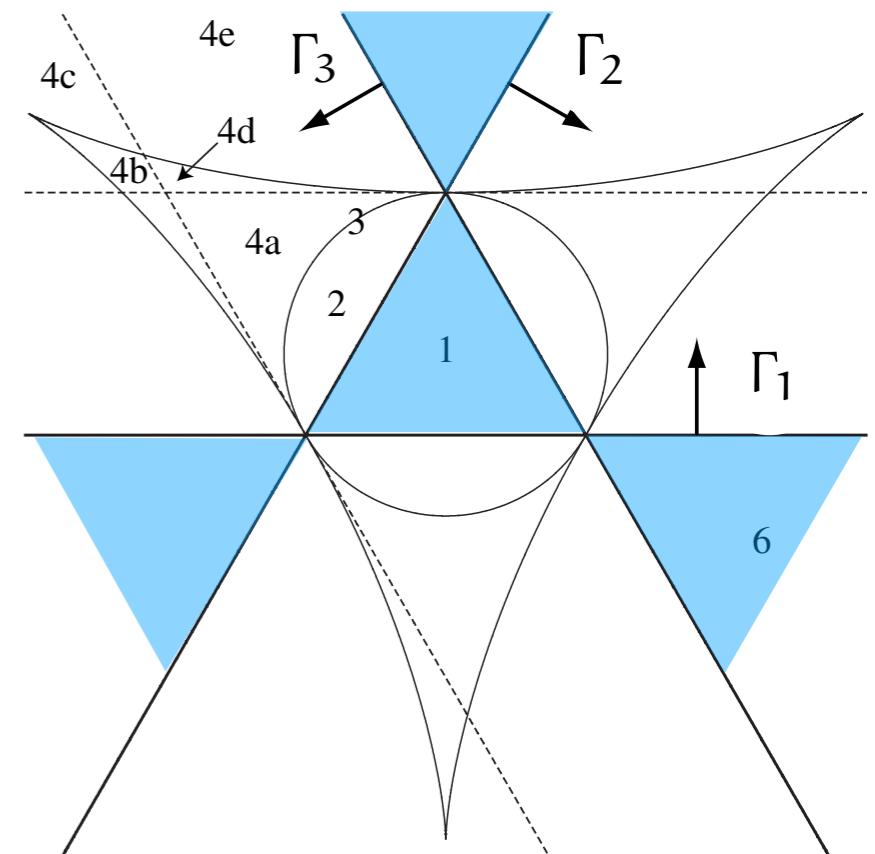
$$\Theta^2 = Z^2 + X^2 + Y^2$$

The three-vortex system conserve an energy when $\kappa_2 < 0$

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Barycentric coordinates for circulation space when $\Gamma_1 + \Gamma_2 + \Gamma_3 = 1$

Dynamics based of the sign of κ_2 :

$$\kappa_2 > 0$$

(X, Y, Z, Θ) represents a sphere

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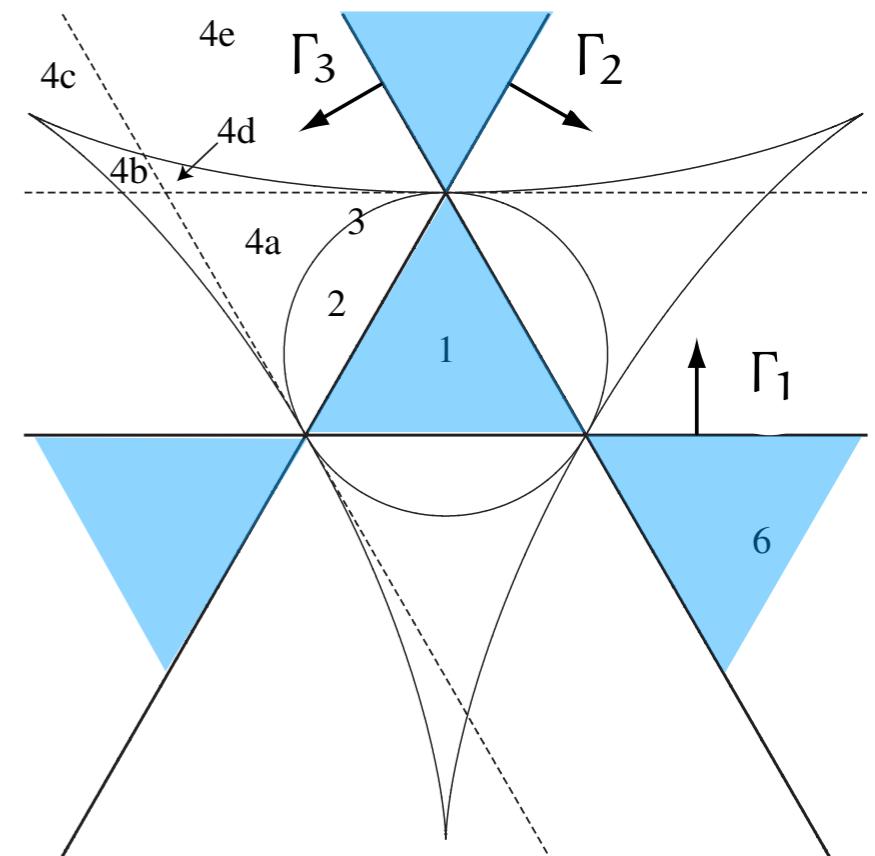
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Barycentric coordinates for circulation space when $\Gamma_1 + \Gamma_2 + \Gamma_3 = 1$

Dynamics based of the sign of κ_2 :

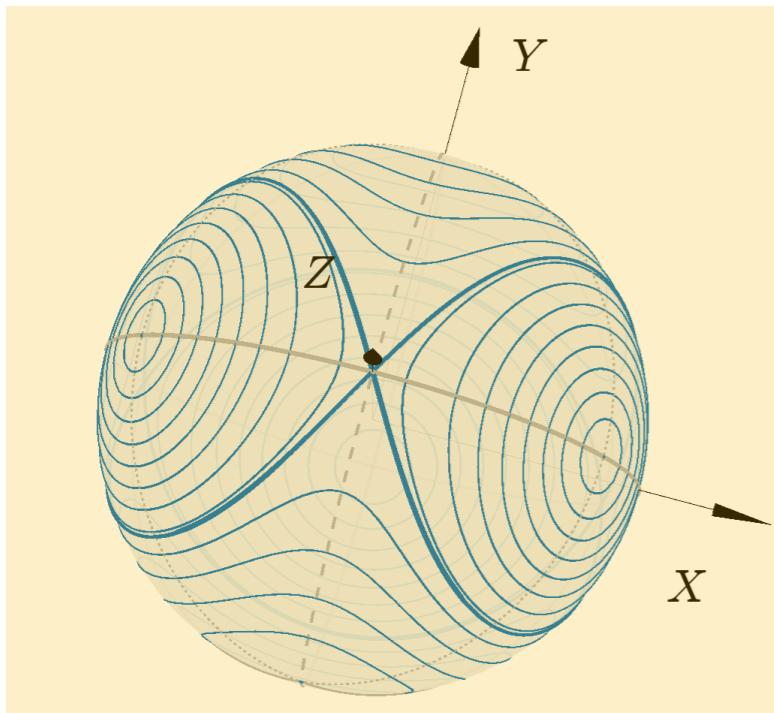
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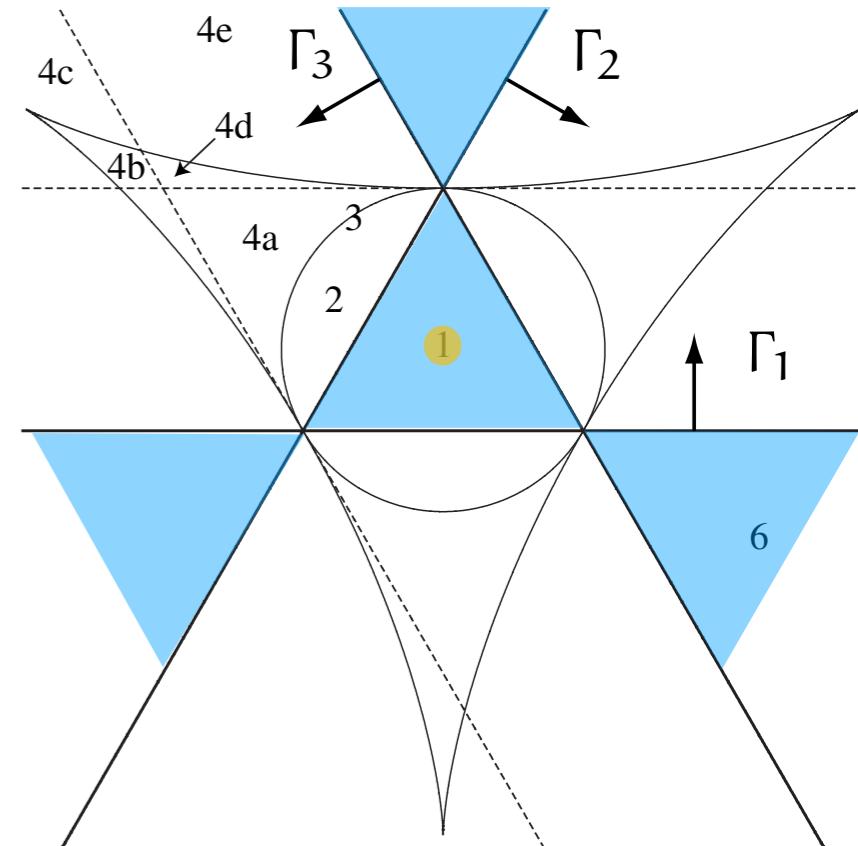
$$\kappa_2 < 0$$

(X, Y, Z, Θ) represents a upper sheet of a two-sheeted hyperboloid

$\kappa_2 > 0$ (Sphere)



$$(\Gamma_1, \Gamma_2, \Gamma_3) = \left(\frac{1}{3}, \frac{1}{3}, \frac{1}{3} \right)$$



Barycentric coordinates for circulation space when $\Gamma_1 + \Gamma_2 + \Gamma_3 = 1$

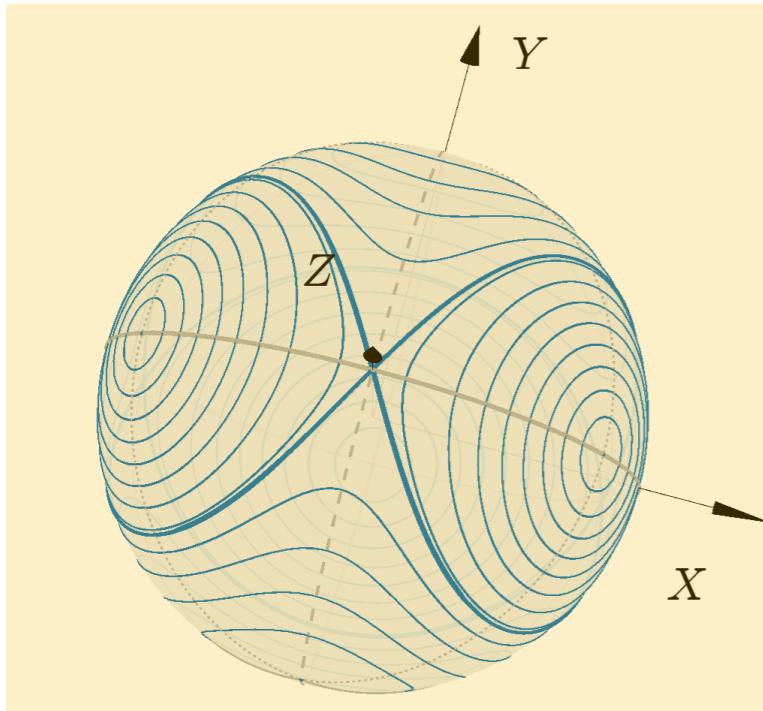
Fixed Point Analysis (Jacobian)

$$J(0, 1, 0) = \begin{pmatrix} 0 & 0 & 3 \\ 0 & 0 & 0 \\ -3 & 0 & 0 \end{pmatrix}, \quad \text{and} \quad J(1, 0, 0) = \begin{pmatrix} 0 & 3 & 0 \\ 9 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}.$$

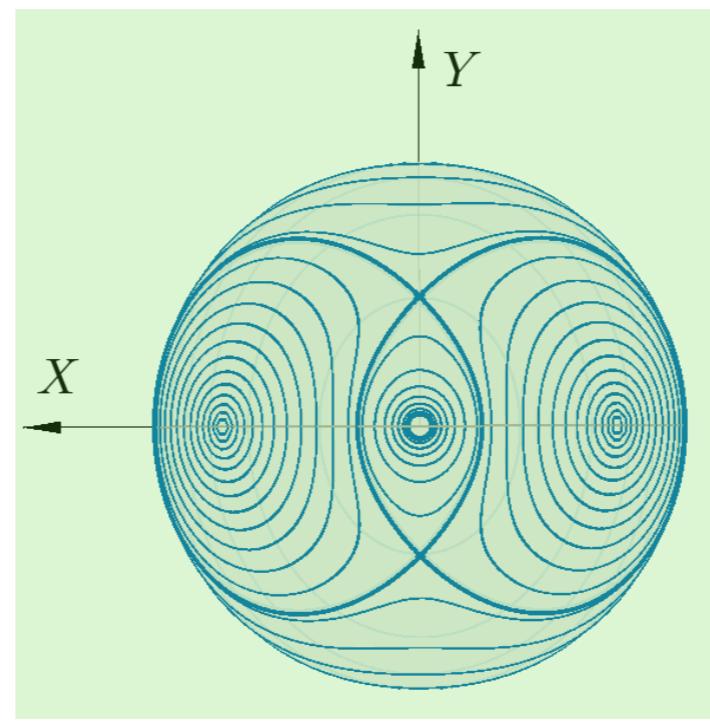
Triangular
Equilibrium
(stable)

Collinear
Equilibrium
(saddle)

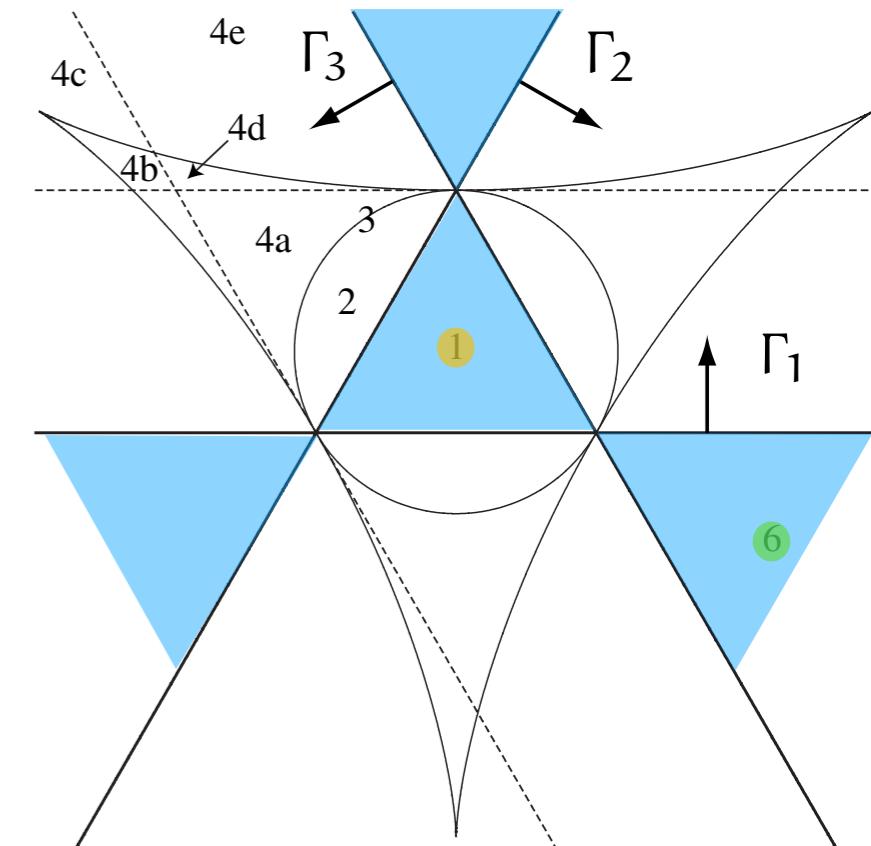
$\kappa_2 > 0$ (Sphere)



$$(\Gamma_1, \Gamma_2, \Gamma_3) = \left(\frac{1}{3}, \frac{1}{3}, \frac{1}{3}\right)$$



$$(\Gamma_1, \Gamma_2, \Gamma_3) = (-2, -1, 4)$$



Barycentric coordinates for circulation space when $\Gamma_1 + \Gamma_2 + \Gamma_3 = 1$

Fixed Point Analysis (Jacobian)

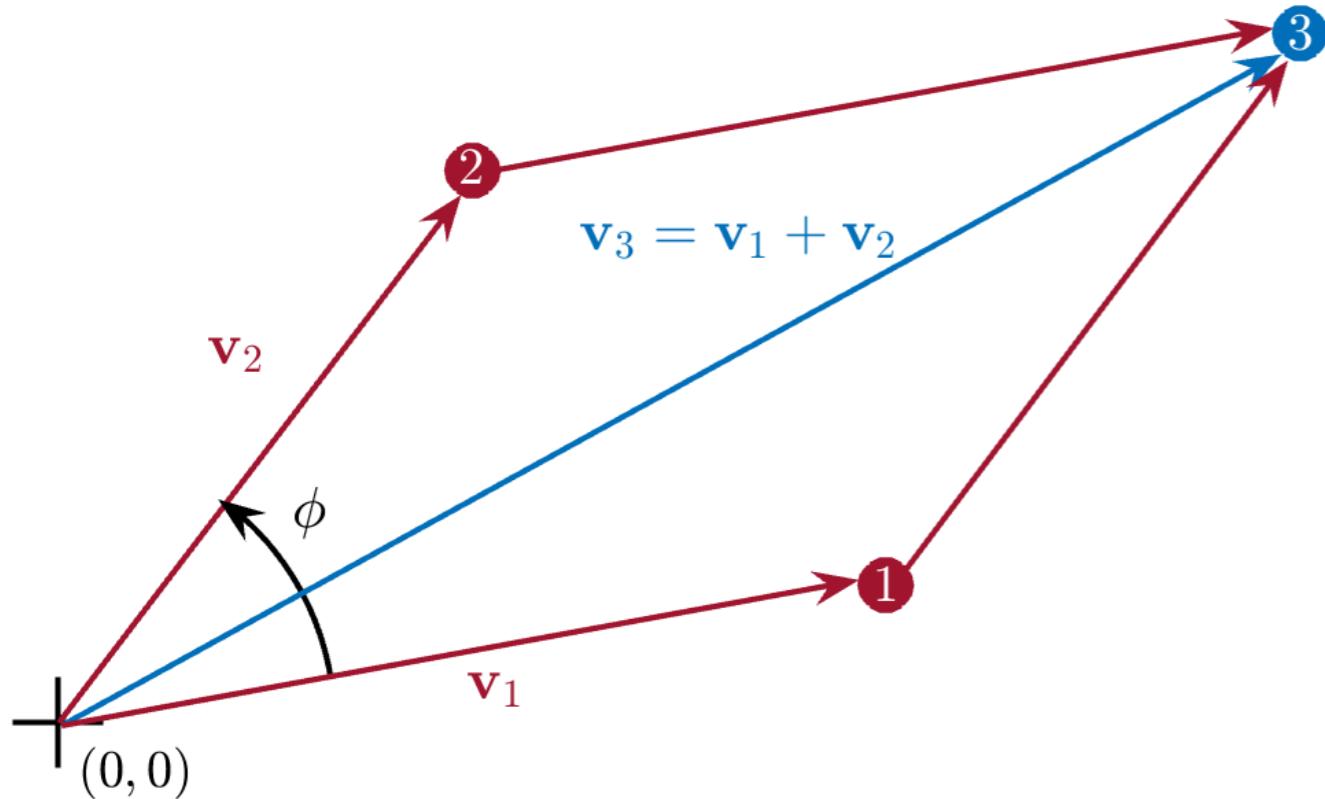
$(0, 0, 1)$ is a center (stable).

$$J(0, 1, 0) = \begin{pmatrix} 0 & 0 & 3 \\ 0 & 0 & 0 \\ -3 & 0 & 0 \end{pmatrix}, \quad \text{and} \quad J(1, 0, 0) = \begin{pmatrix} 0 & 3 & 0 \\ 9 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}. \quad (0, \pm \frac{3}{5}, \pm \frac{4}{5}) \text{ is a saddle.}$$

Triangular
Equilibrium
(stable)

Collinear
Equilibrium
(saddle)

$$\Gamma = (1, 1, -1) : \kappa_2 < 0 \text{ (Hyperboloid)}$$



New Coordinates

$$X = -\|\mathbf{v}_1\|^2 + \|\mathbf{v}_2\|^2;$$

$$Y = 2 \|\mathbf{v}_1\| \|\mathbf{v}_2\| \sin \phi;$$

$$Z = \|\mathbf{v}_1\|^2 + \|\mathbf{v}_2\|^2;$$

$$\Theta = -2 \|\mathbf{v}_1\| \|\mathbf{v}_2\| \cos \phi.$$

which satisfy:

$$\frac{dX}{dt} = \frac{-2Y}{Z + \Theta} + \frac{4ZY}{Z^2 - X^2};$$

$$\frac{dY}{dt} = \frac{2X}{Z + \Theta};$$

$$\frac{dZ}{dt} = \frac{4XY}{Z^2 - X^2},$$

$$\Theta^2 = Z^2 - X^2 - Y^2$$

(conserved quantity)

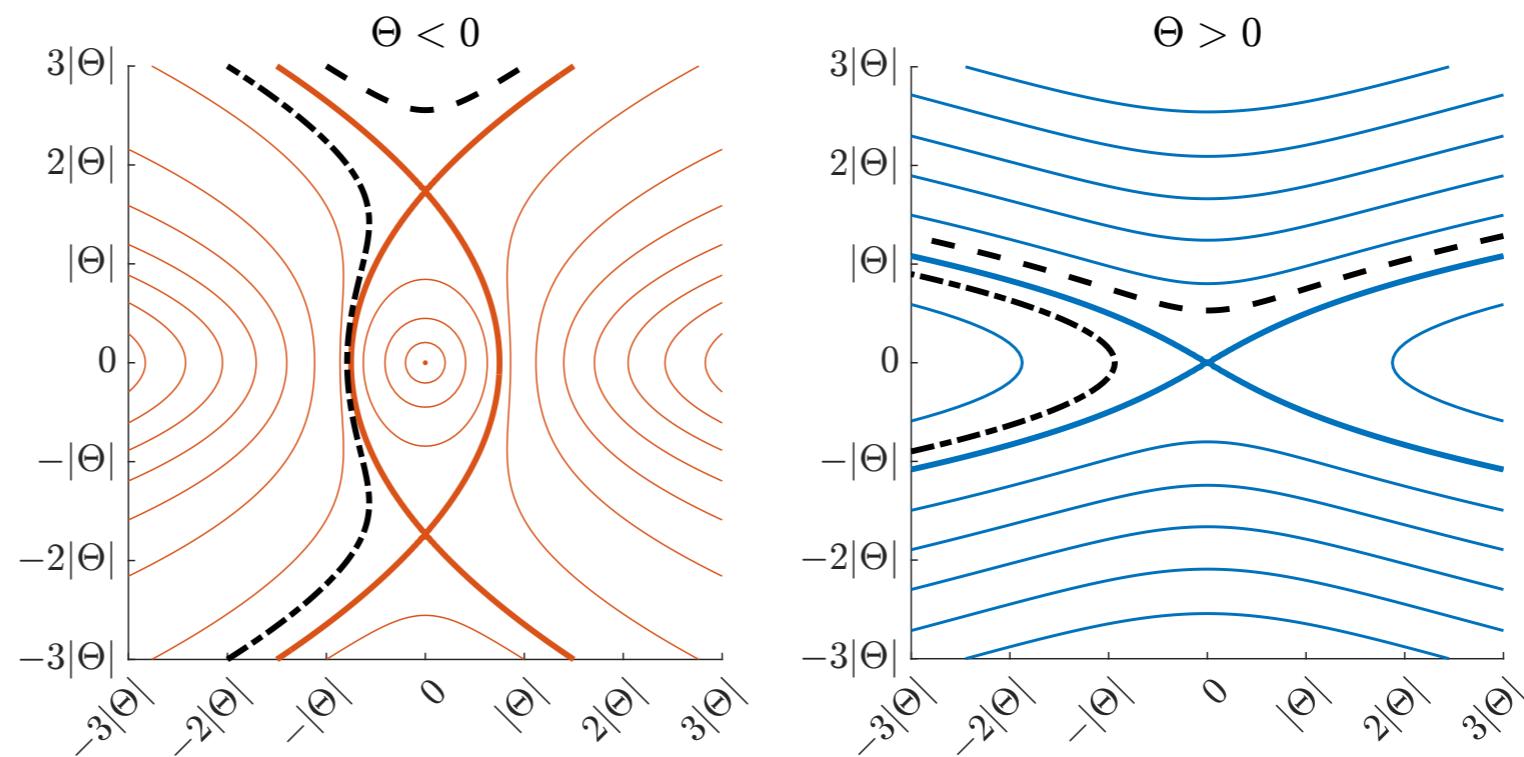
Phase Space when $(\Gamma_1, \Gamma_2, \Gamma_3) = (1, 1, -1)$

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The space is the upper sheet of a two-sheeted hyperbola.

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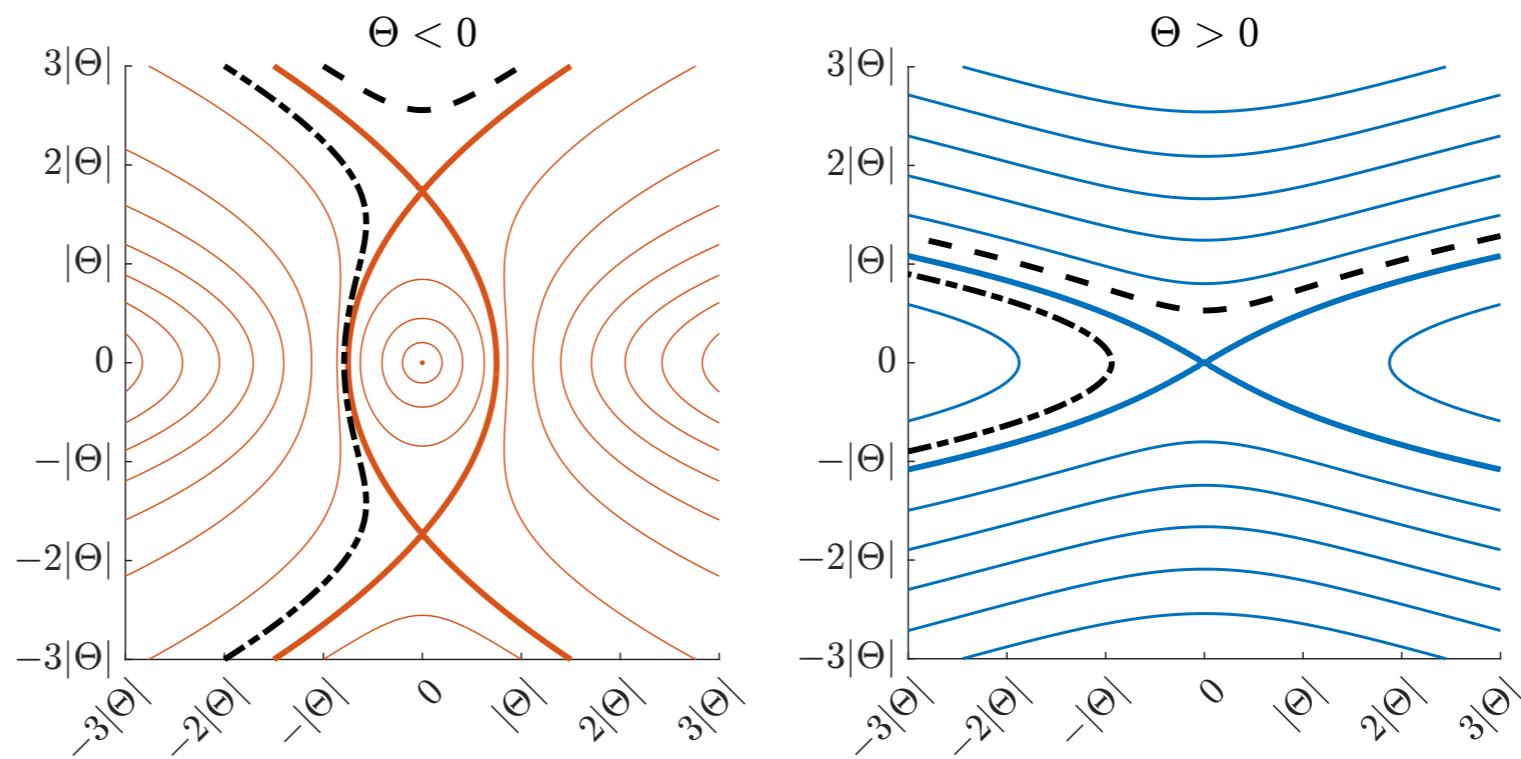
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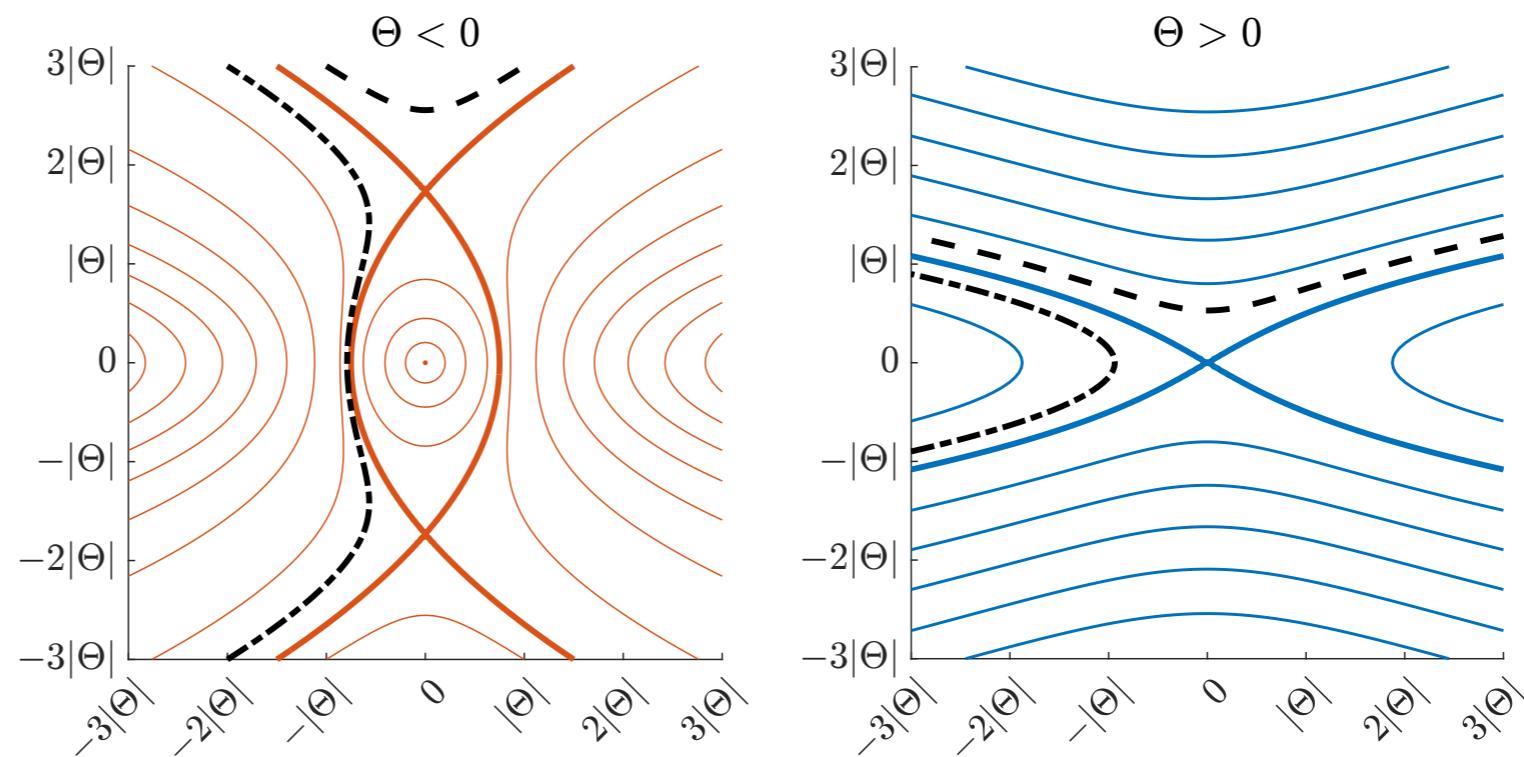
The level sets of the Hamiltonian are plotted in the phase planes



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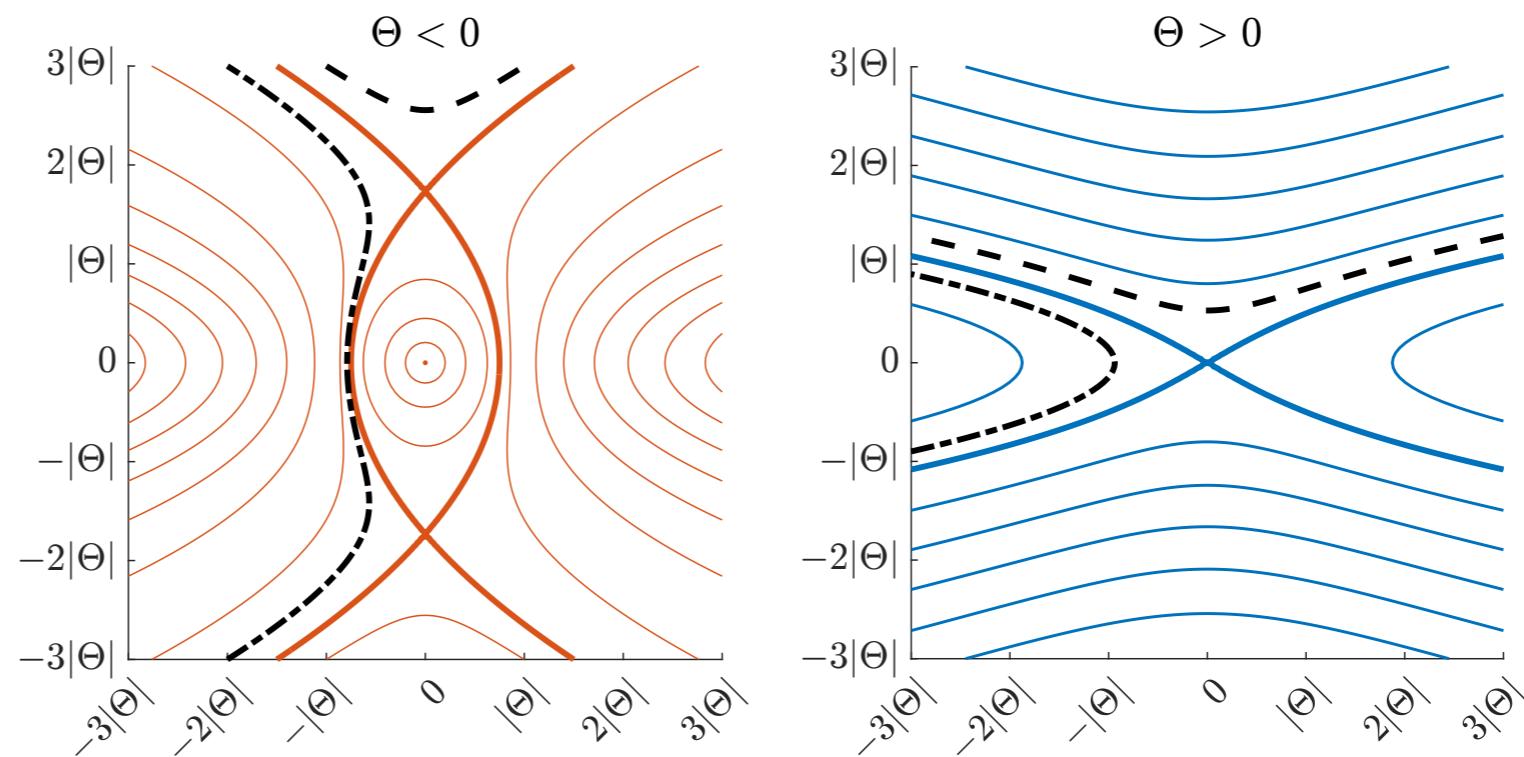


What do they represent in Aref's Phase Plane diagram?

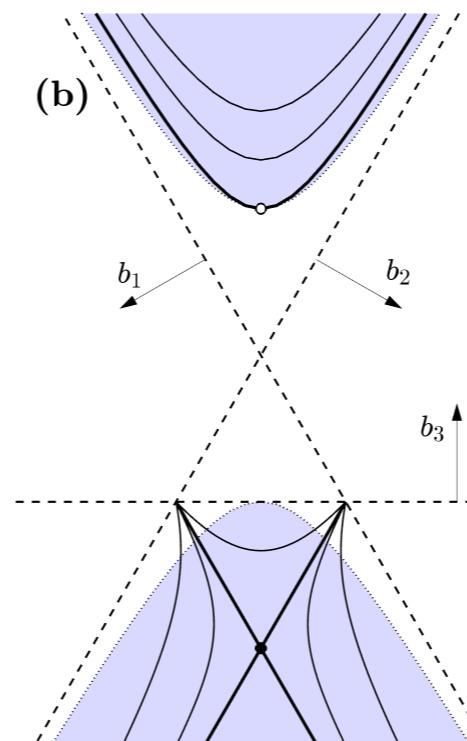
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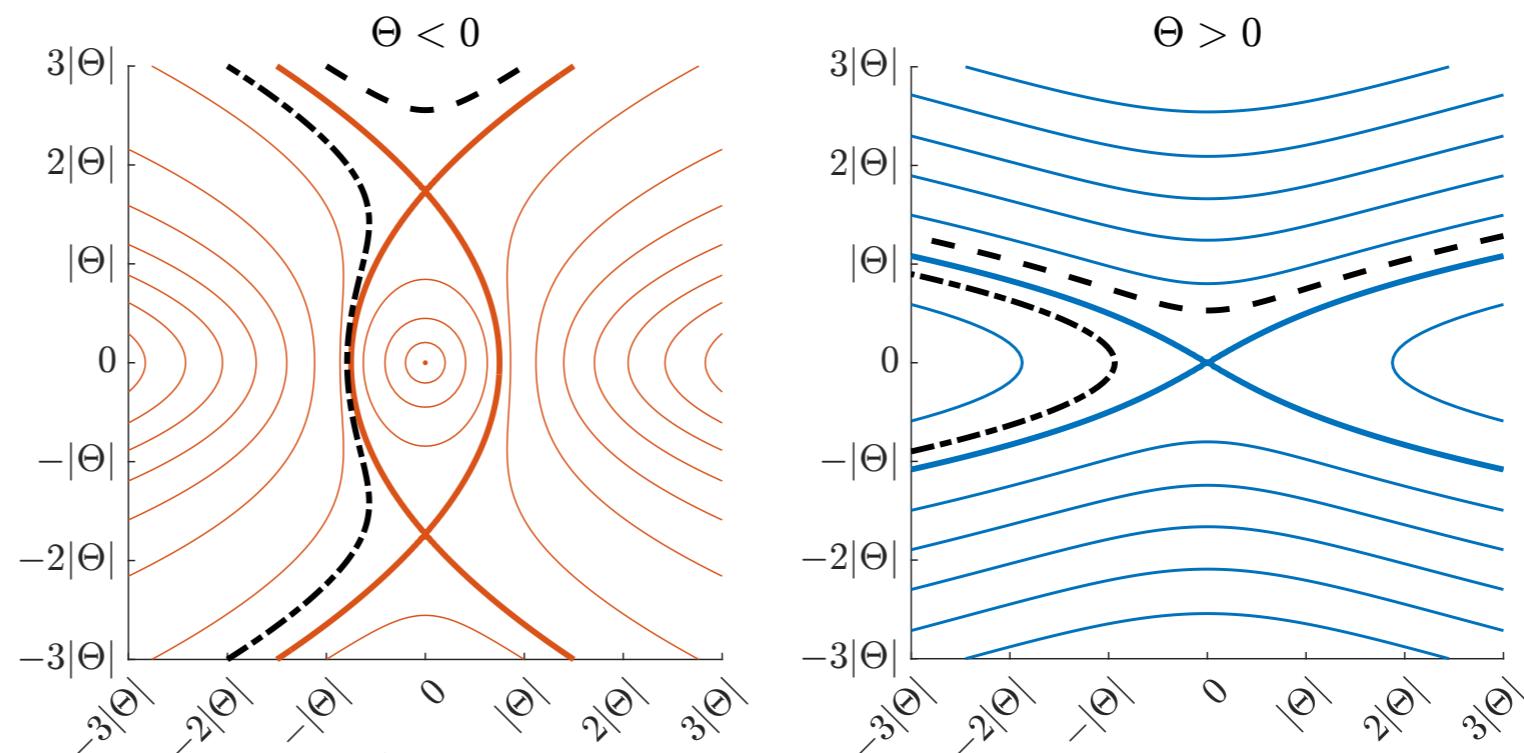


Vortices of circulation $(1, 1, -1)$

Phase Space when $(\Gamma_1, \Gamma_2, \Gamma_3) = (1, 1, -1)$

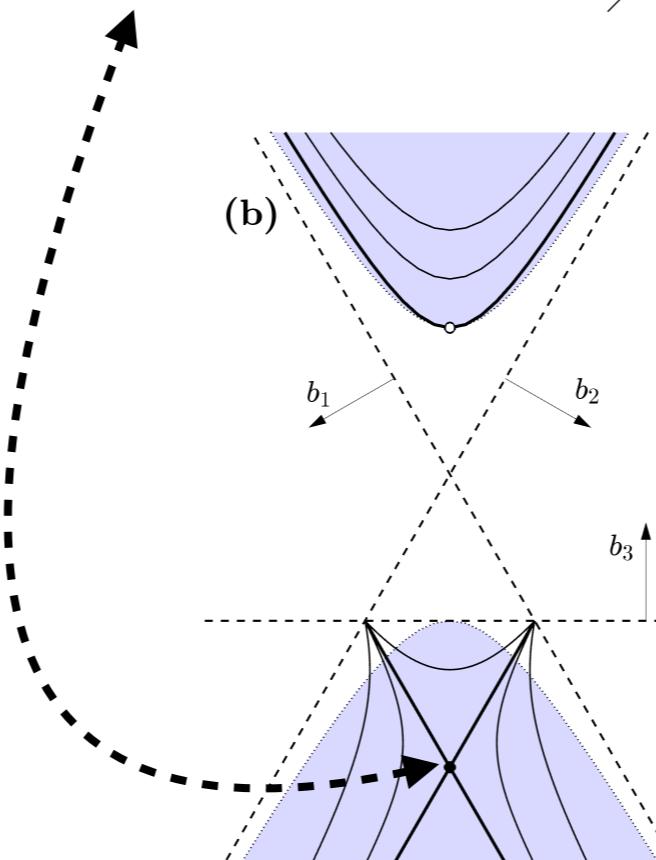
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Triangular configuration where at the separatrix

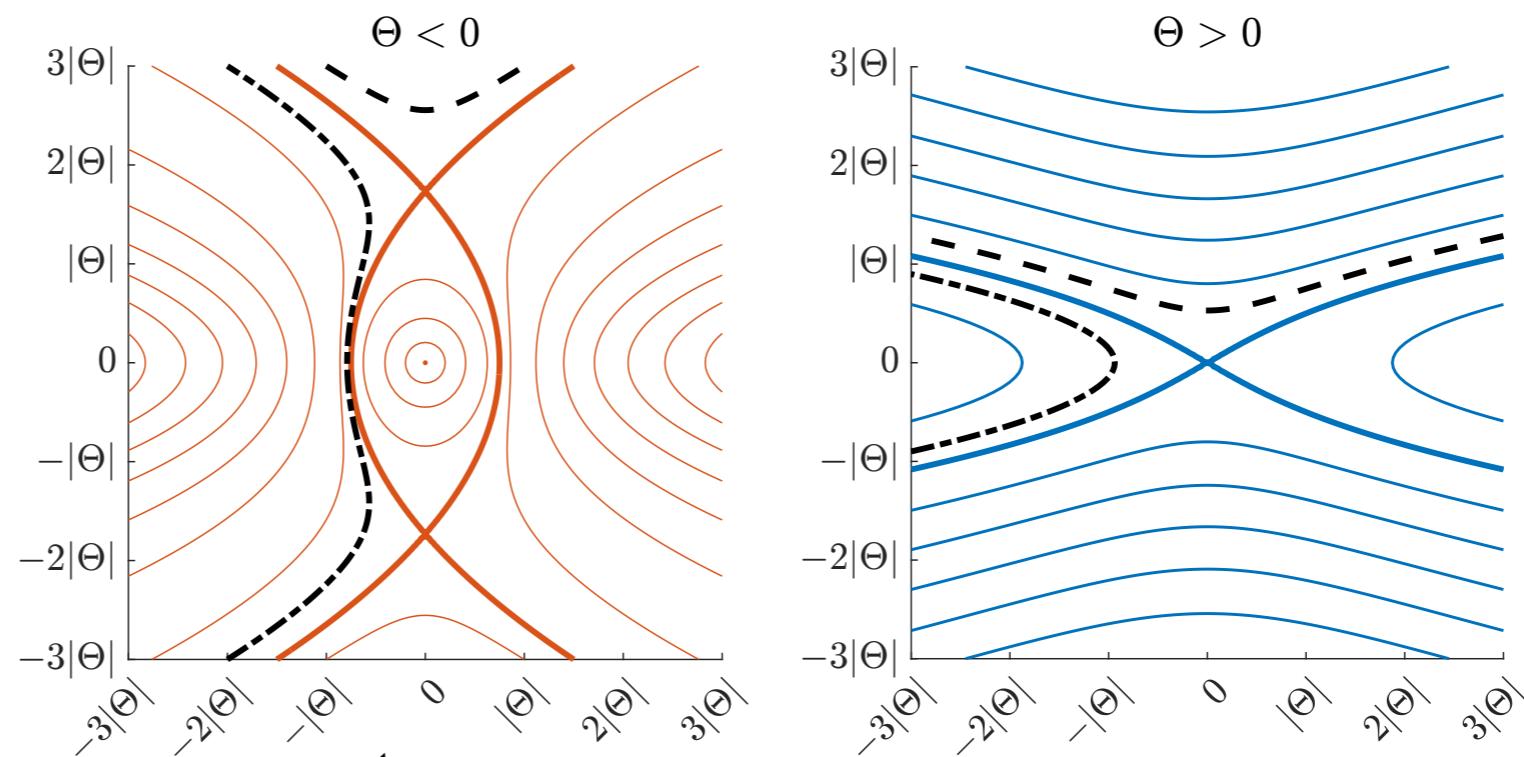


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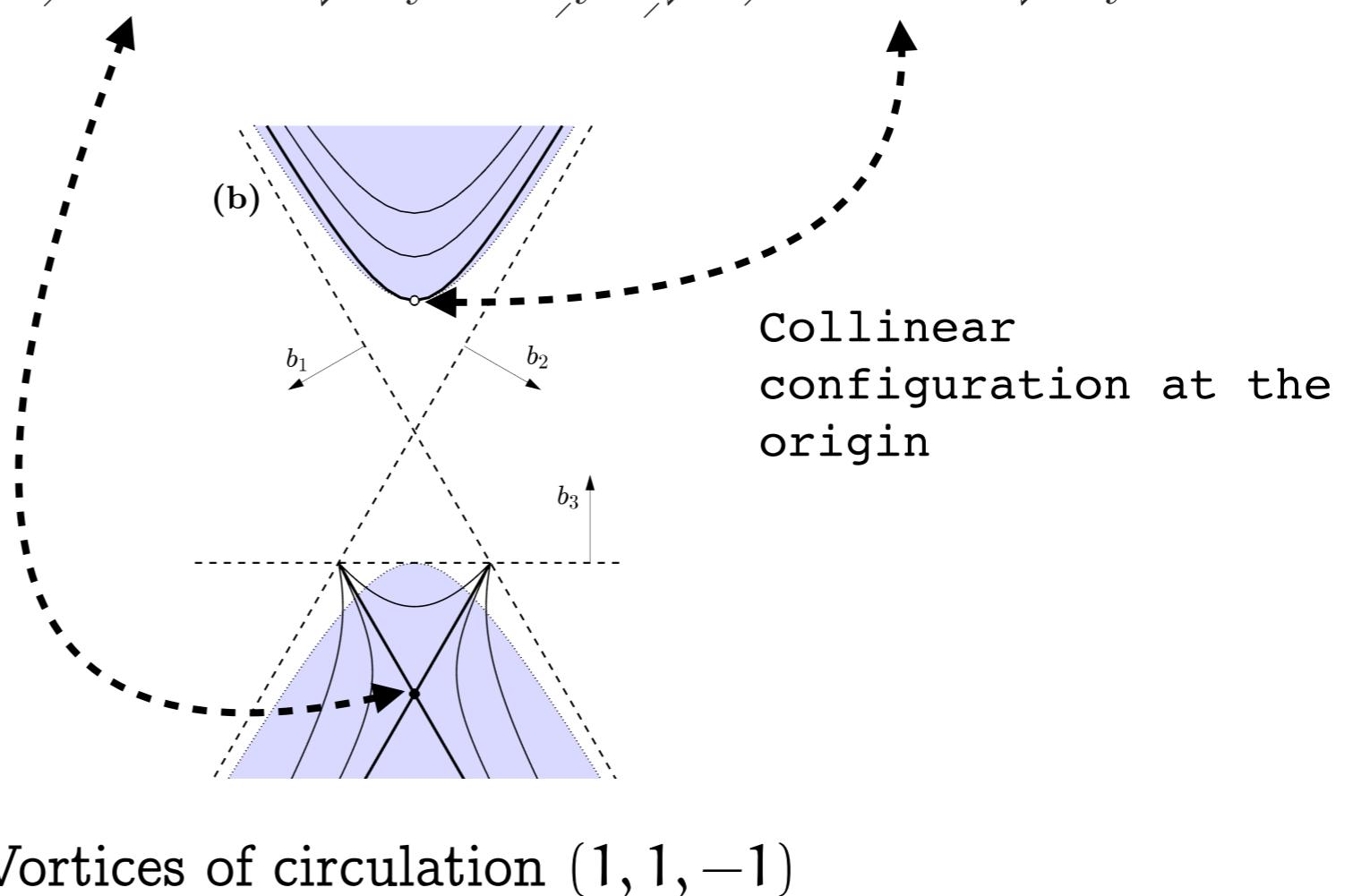
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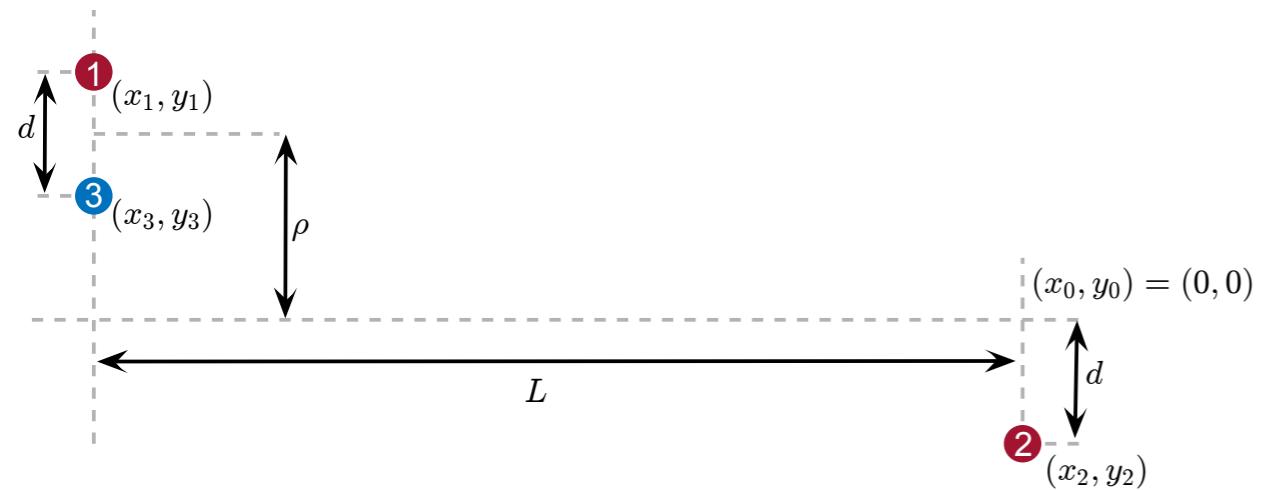
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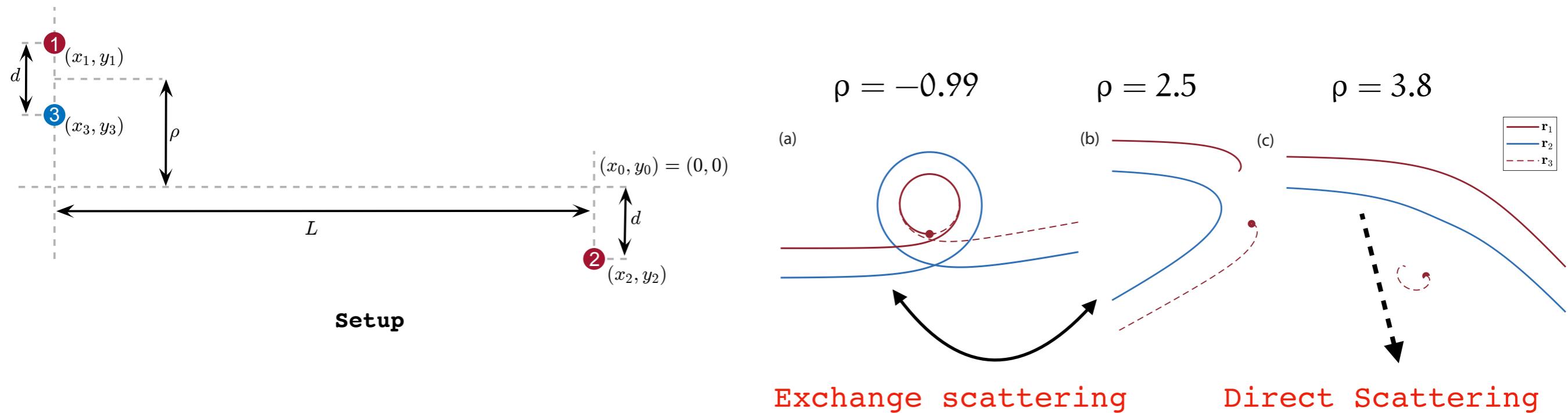
Application to vortex dipole scattering for $\Gamma = (1, 1, -1)$

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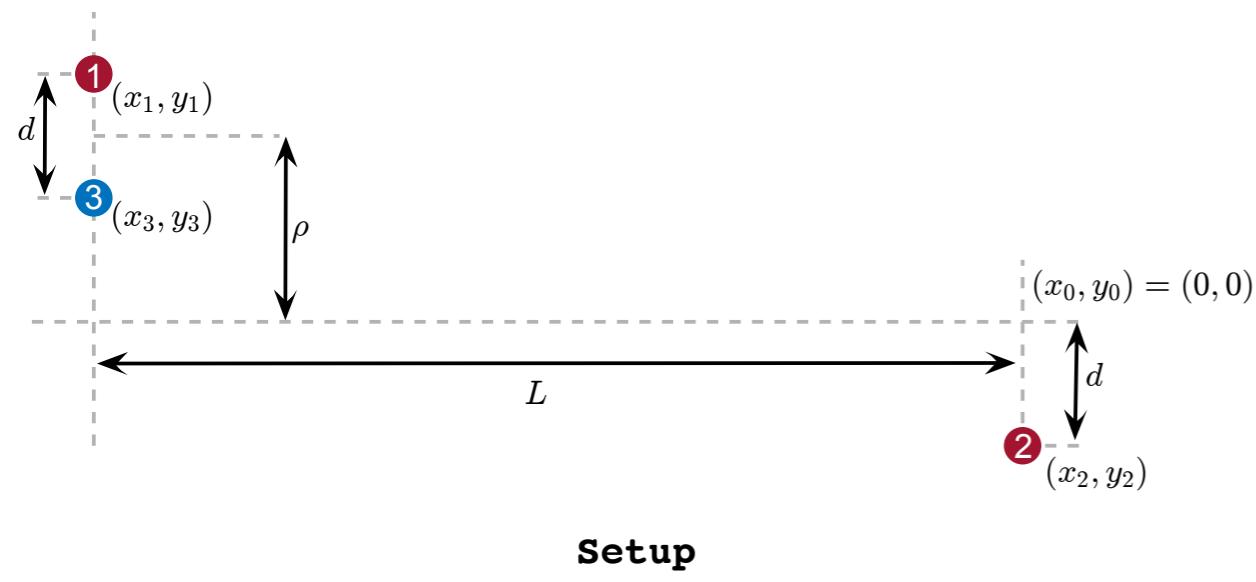


Setup

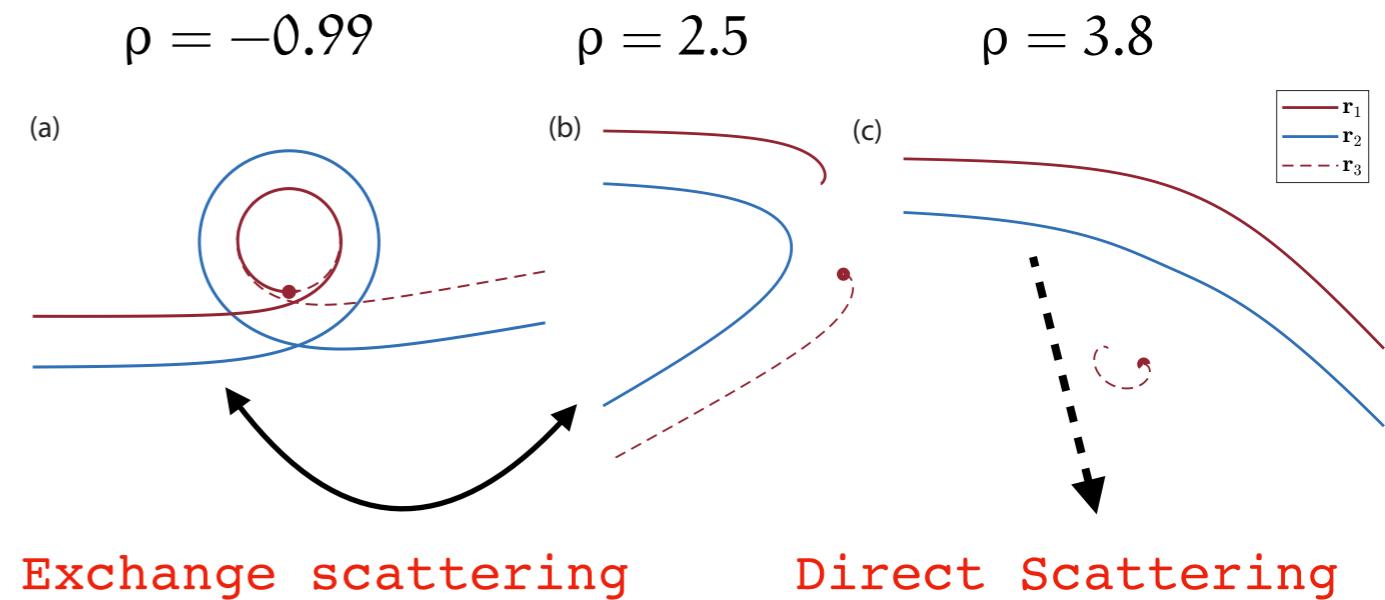
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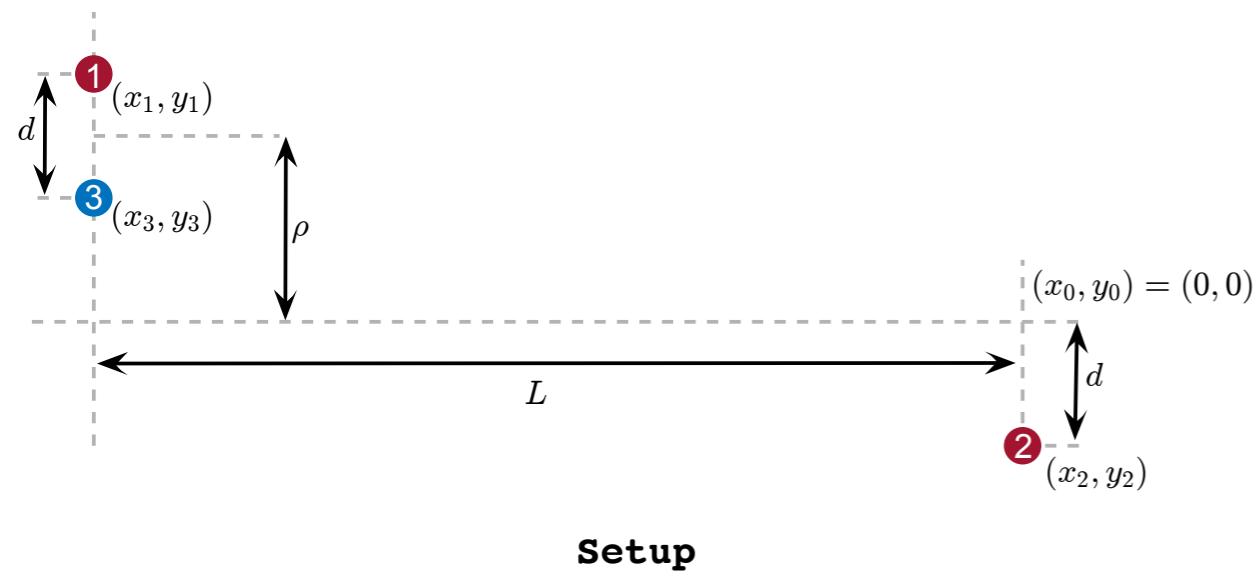
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H and Θ are determined by the parameter ρ



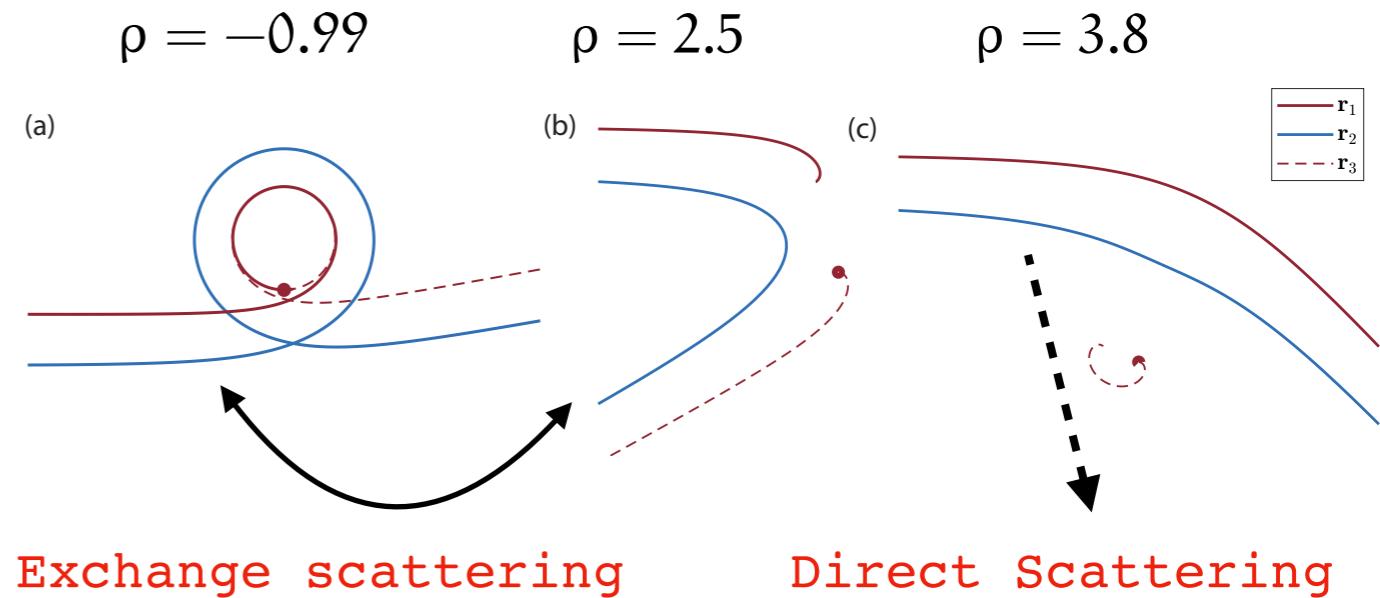
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Setup

H and Θ are determined by the parameter ρ

At the critical values of ρ , the nature of the scattering changes

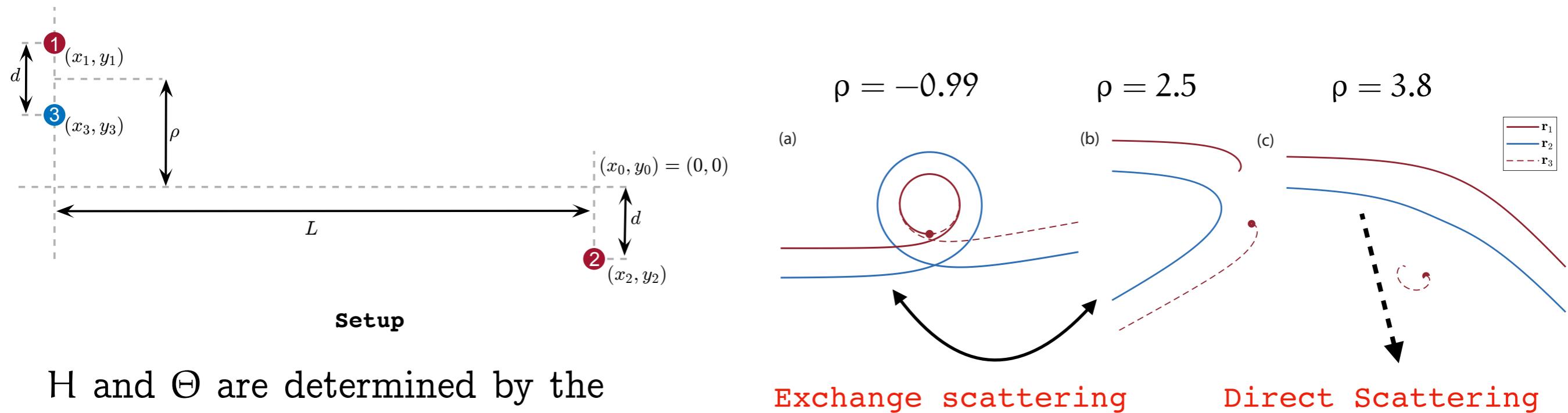


Exchange scattering

Direct Scattering

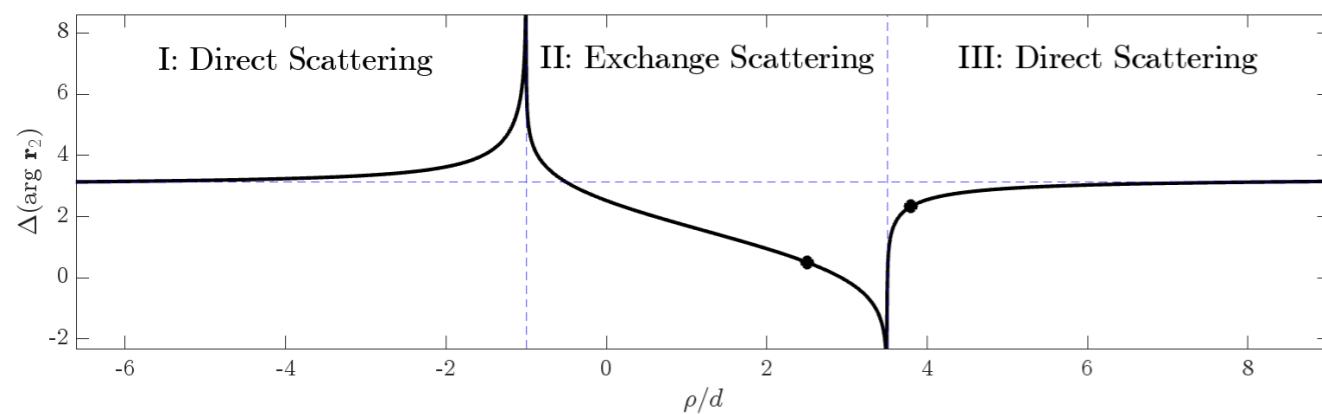
$\begin{cases} \text{---} & r_1 \\ \text{---} & r_2 \\ \text{---} & r_3 \end{cases}$

Application to vortex dipole scattering for $\Gamma = (1, 1, -1)$

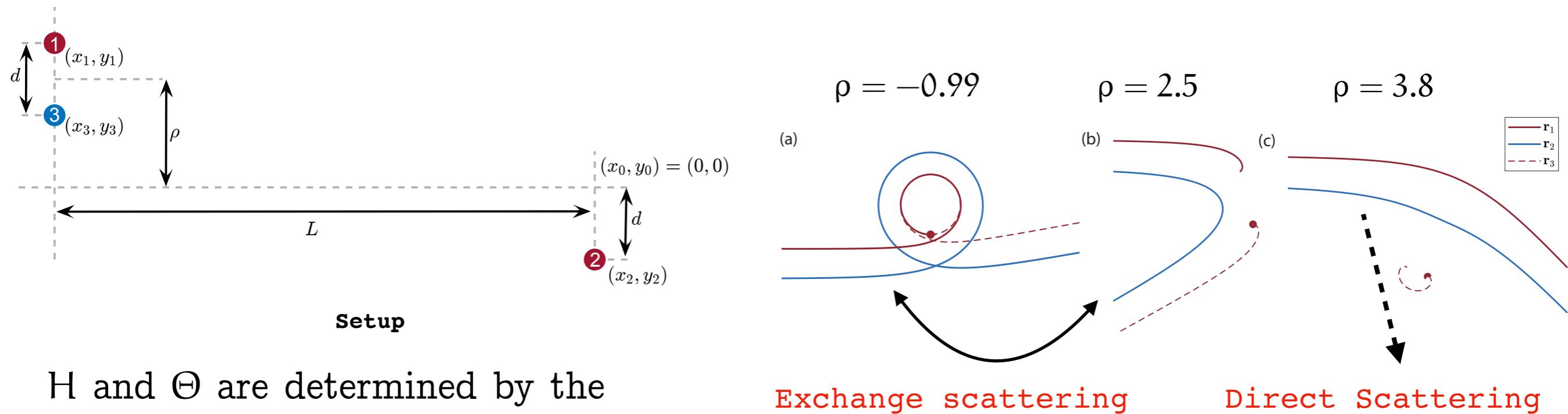


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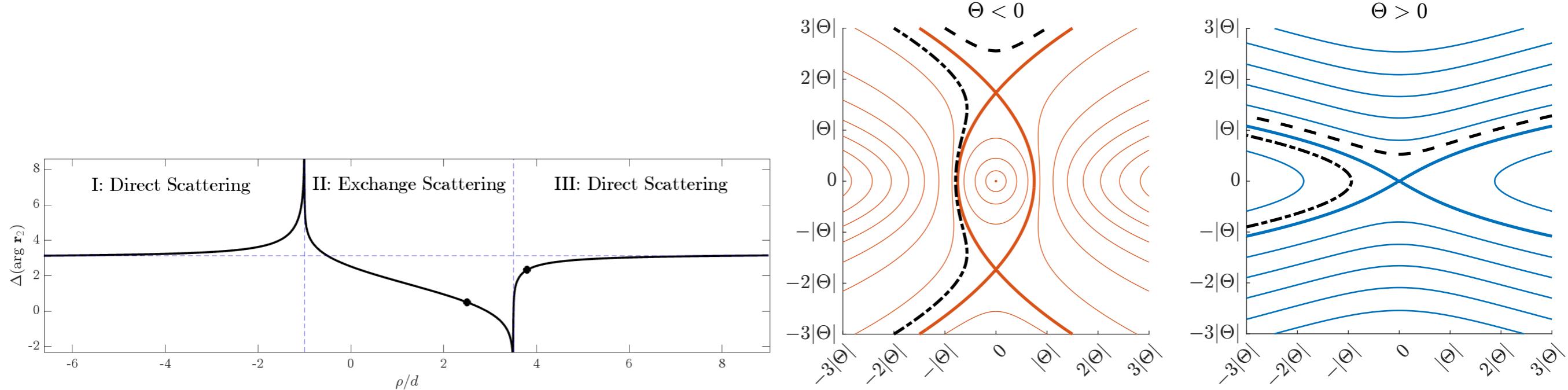


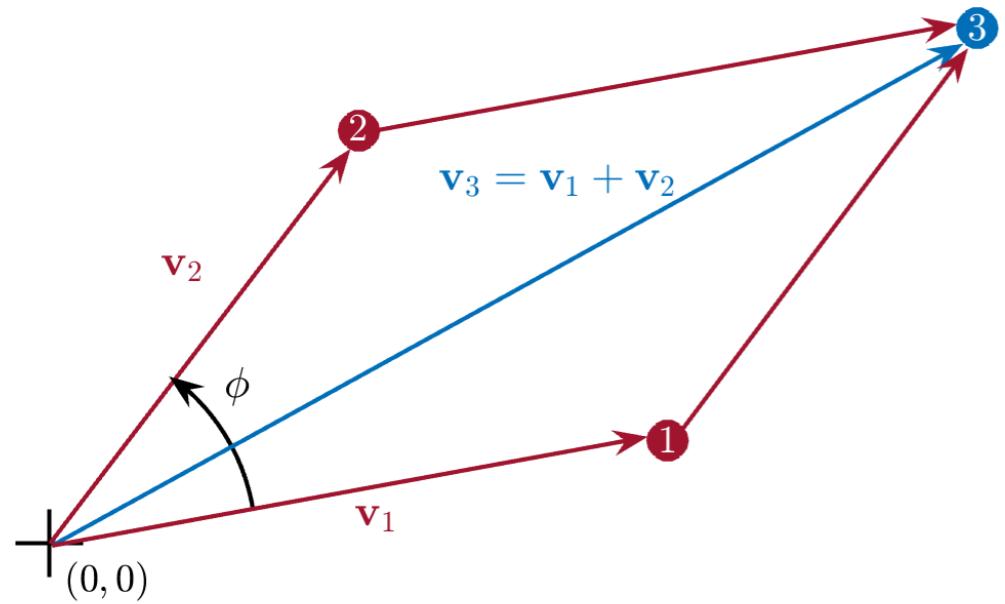
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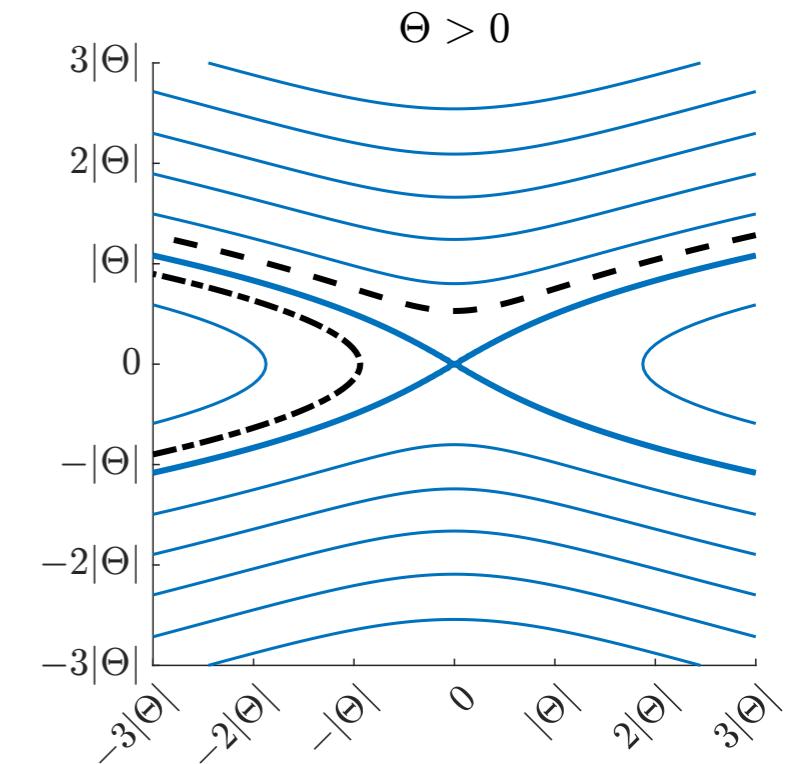
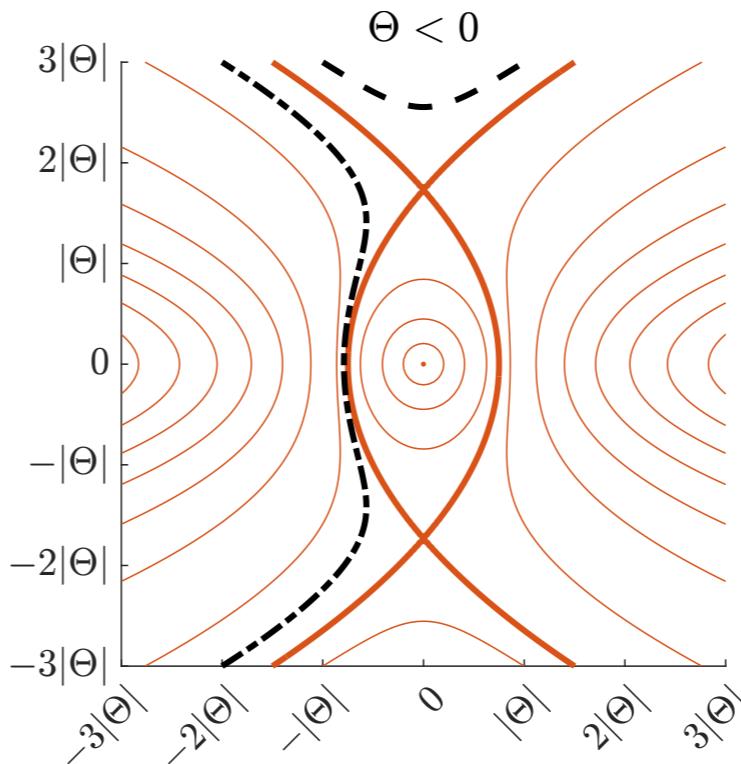


$$X = -\|v_1\|^2 + \|v_2\|^2;$$

$$Y = 2 \|v_1\| \|v_2\| \sin \phi.$$

X vanishes when the triangle of vortices is isosceles.

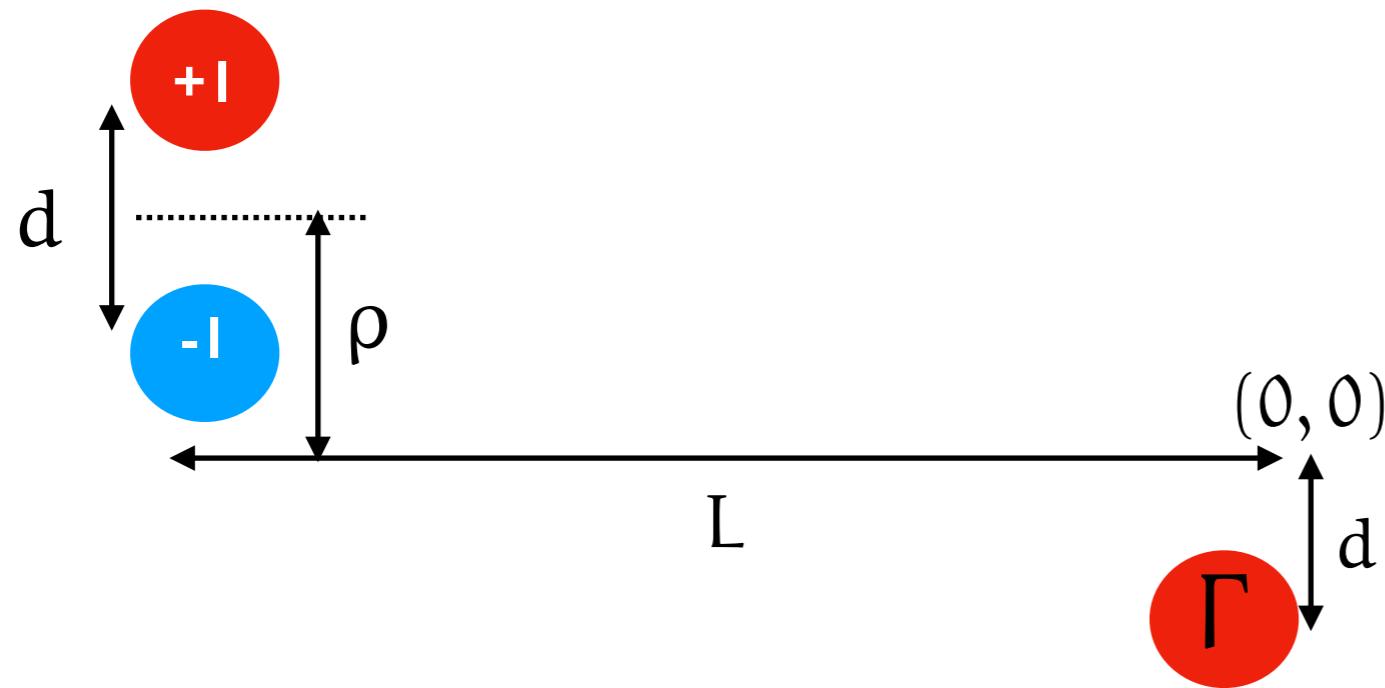
Y vanishes when the three vortices are collinear, which is not a singularity of the coordinate system.



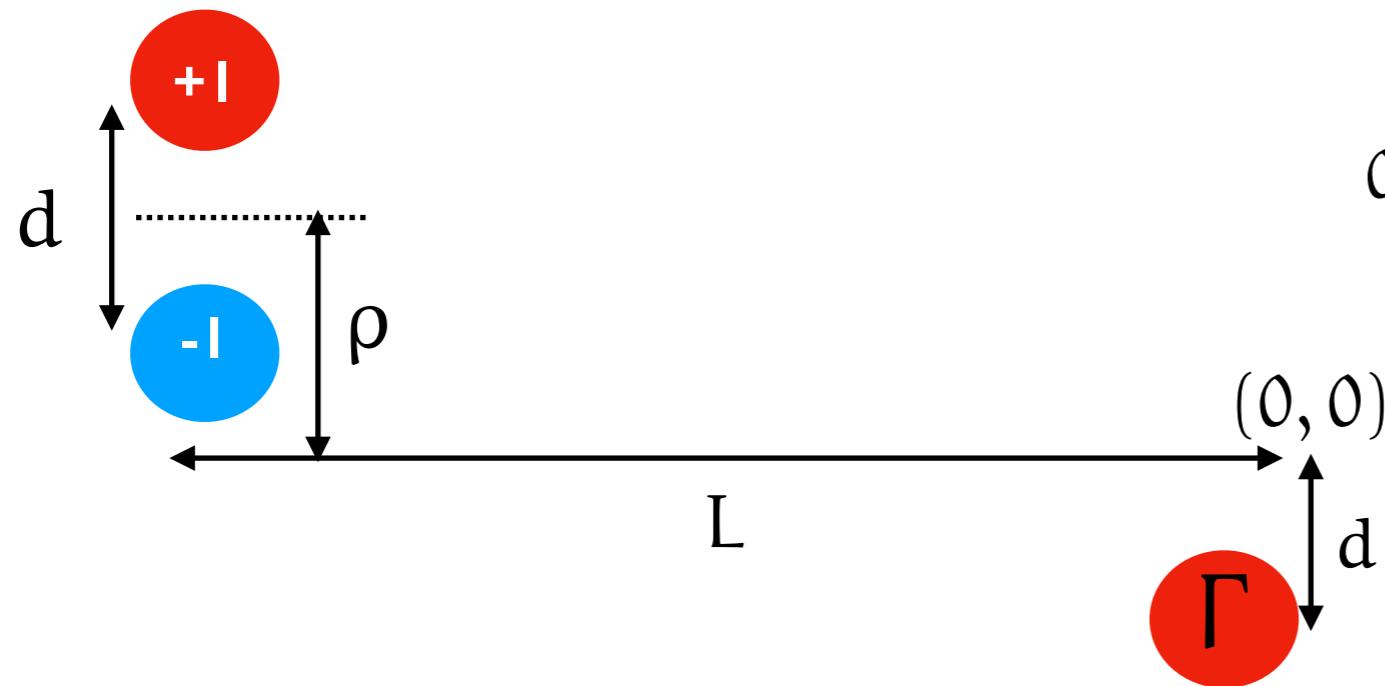
Above the separatrix, $X \rightarrow +\infty$ as $t \rightarrow +\infty$, $\|v_2\|$ diverges, an *exchange scattering*.

Below the separatrix, $X \rightarrow -\infty$ as $t \rightarrow +\infty$, $\|v_1\|$ diverges, a *direct scattering*.

Generalizing Dipole Scattering

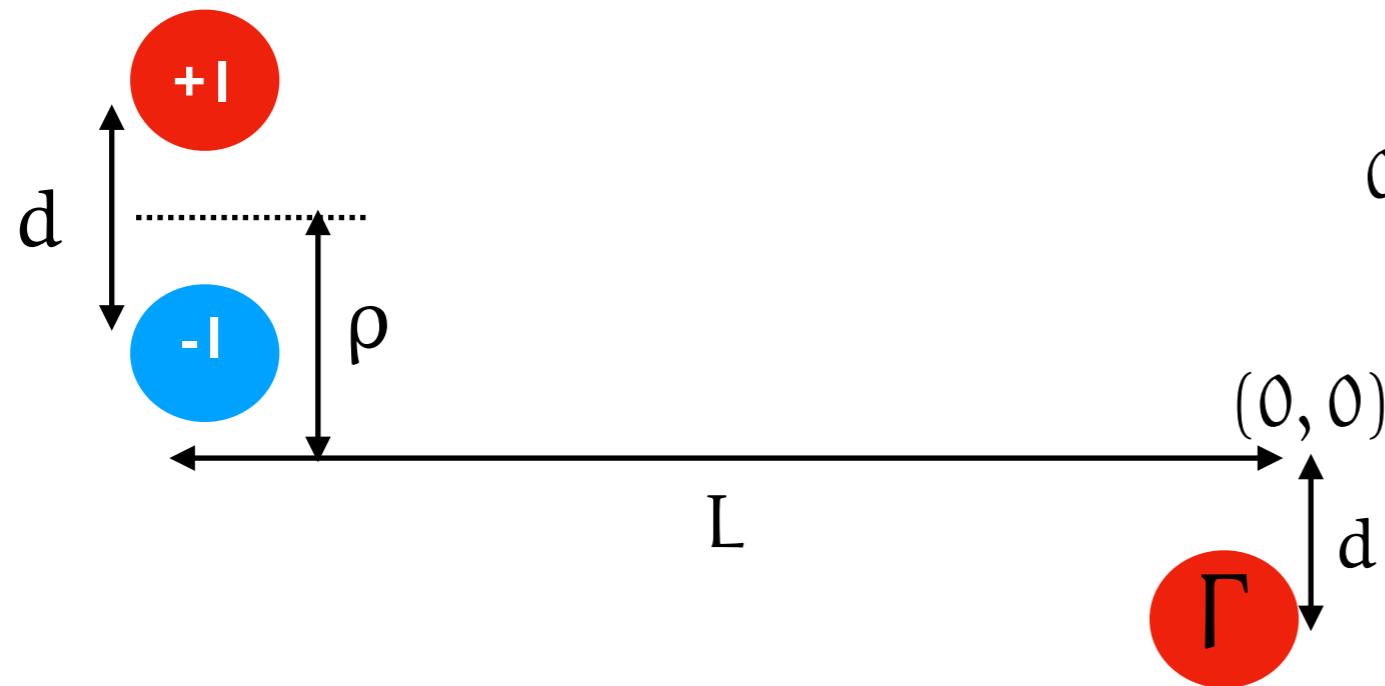


Generalizing Dipole Scattering



$0 < \Gamma \neq 1$: No Exchange scattering

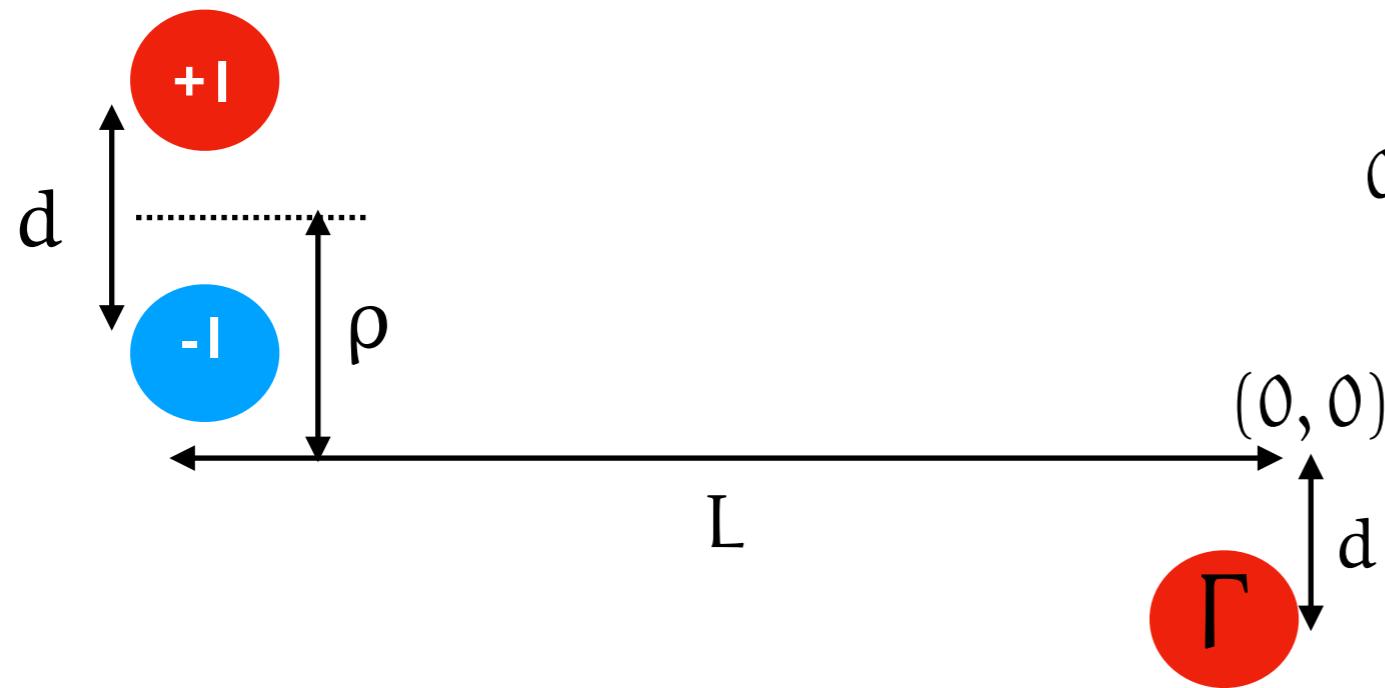
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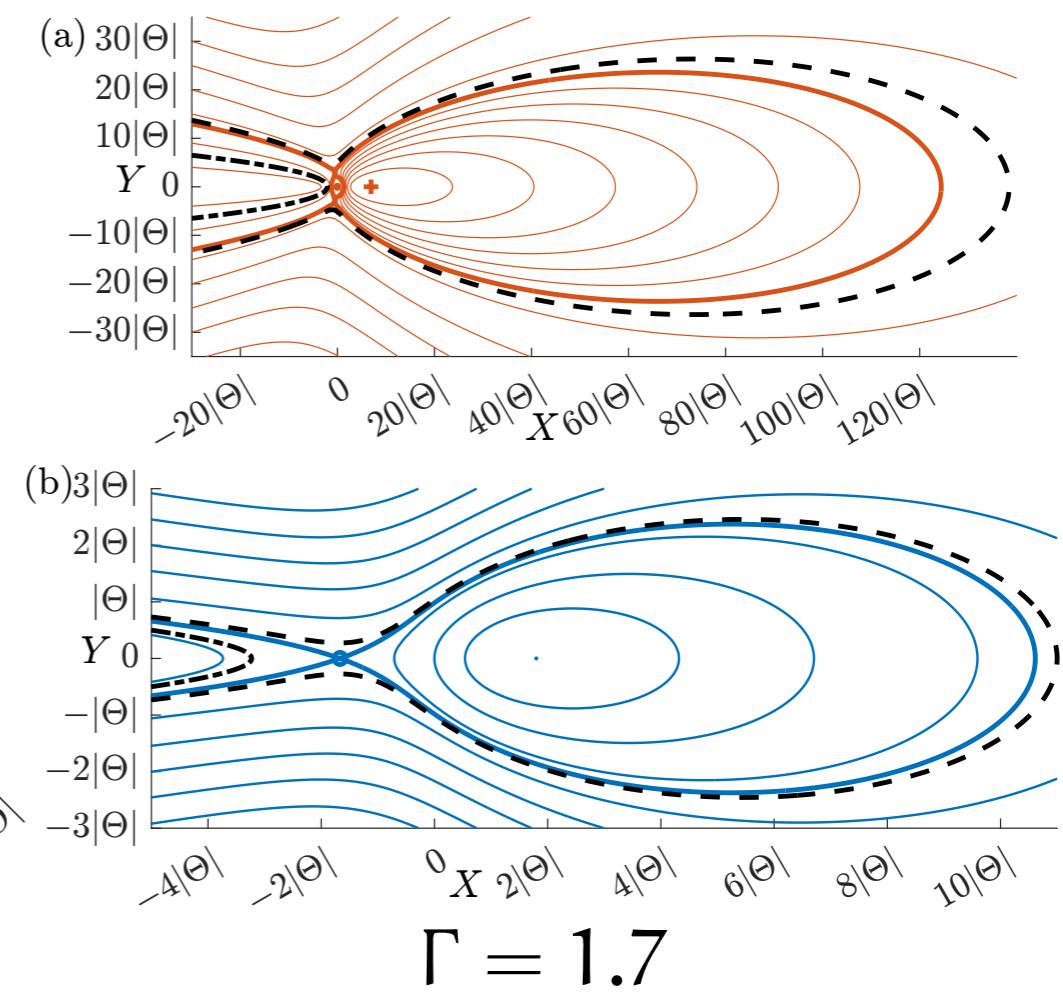
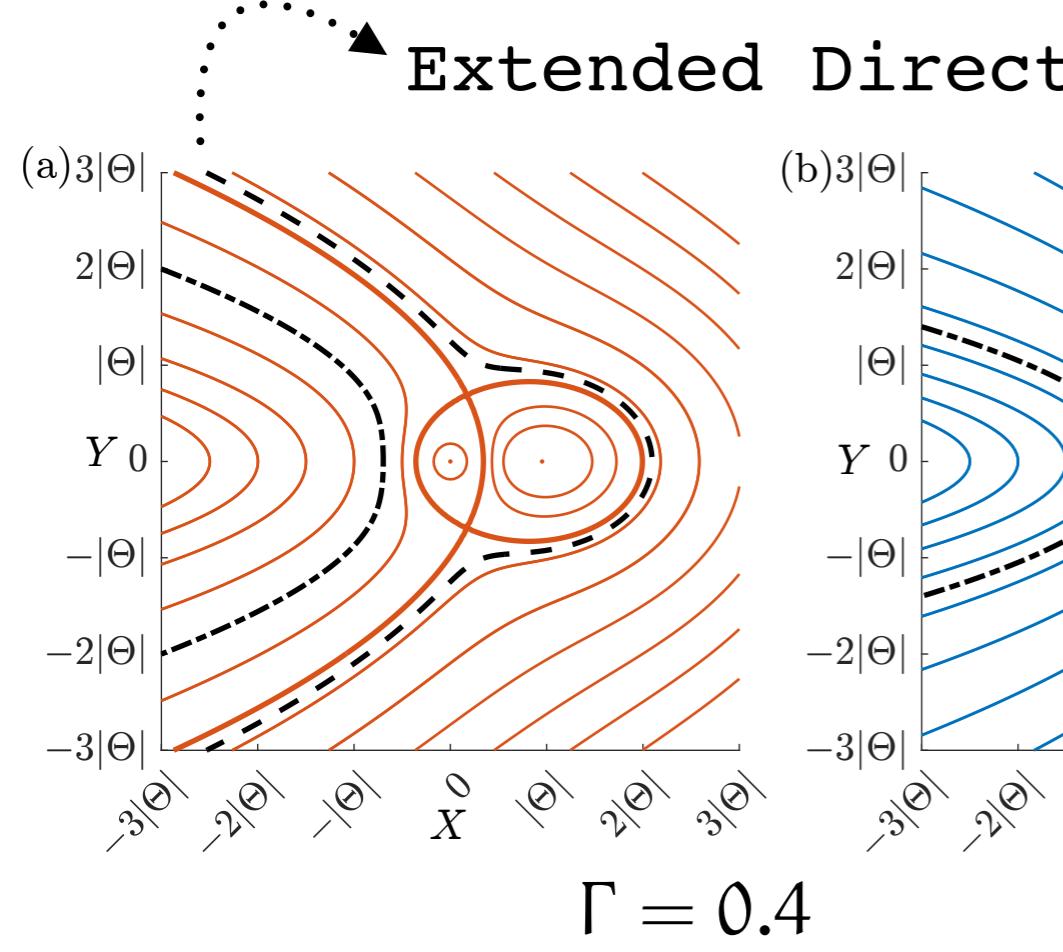
What replaces it?

Generalizing Dipole Scattering

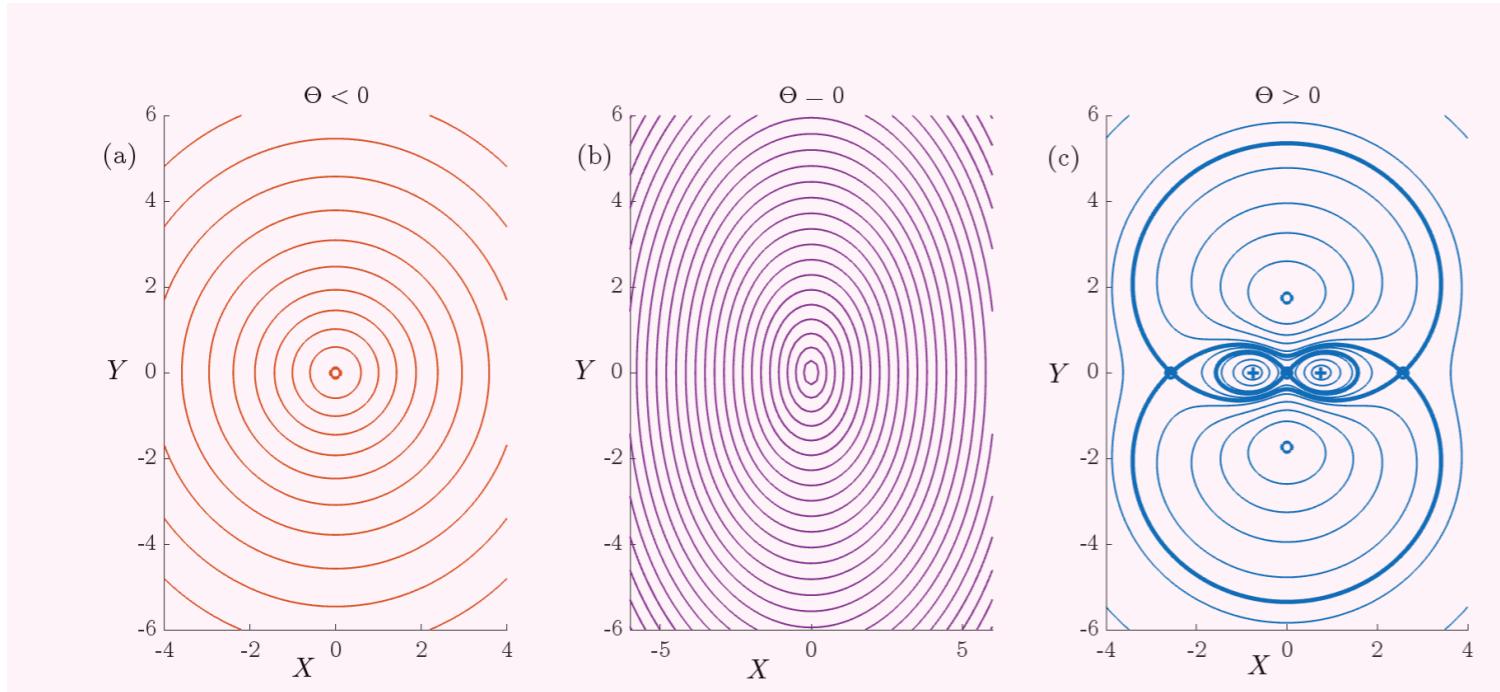


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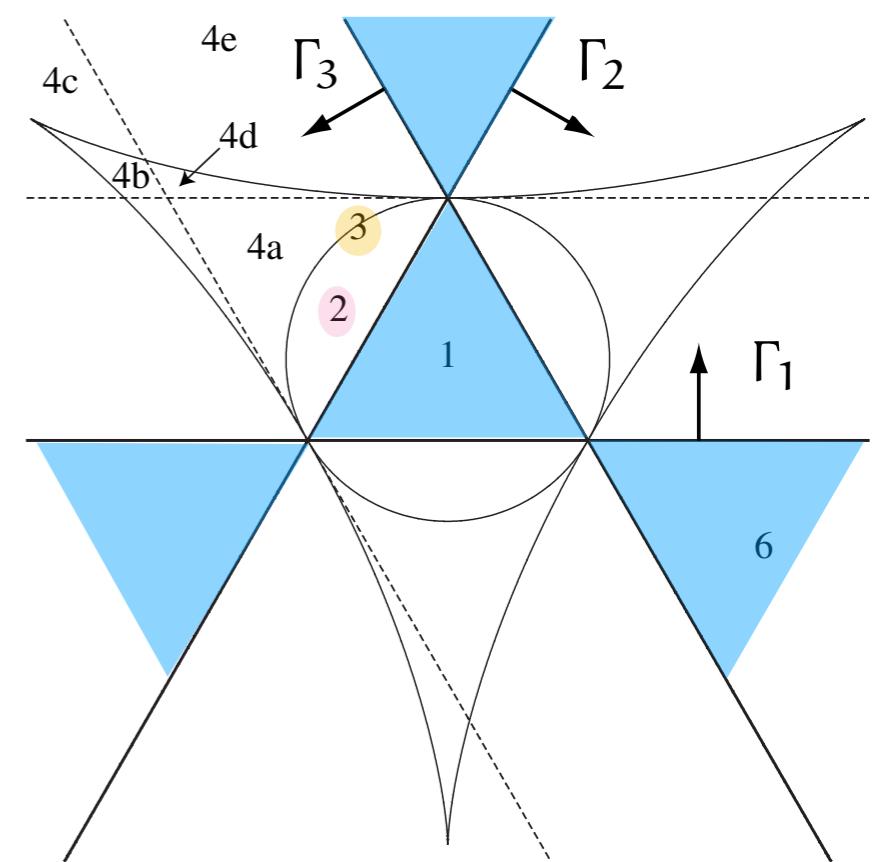
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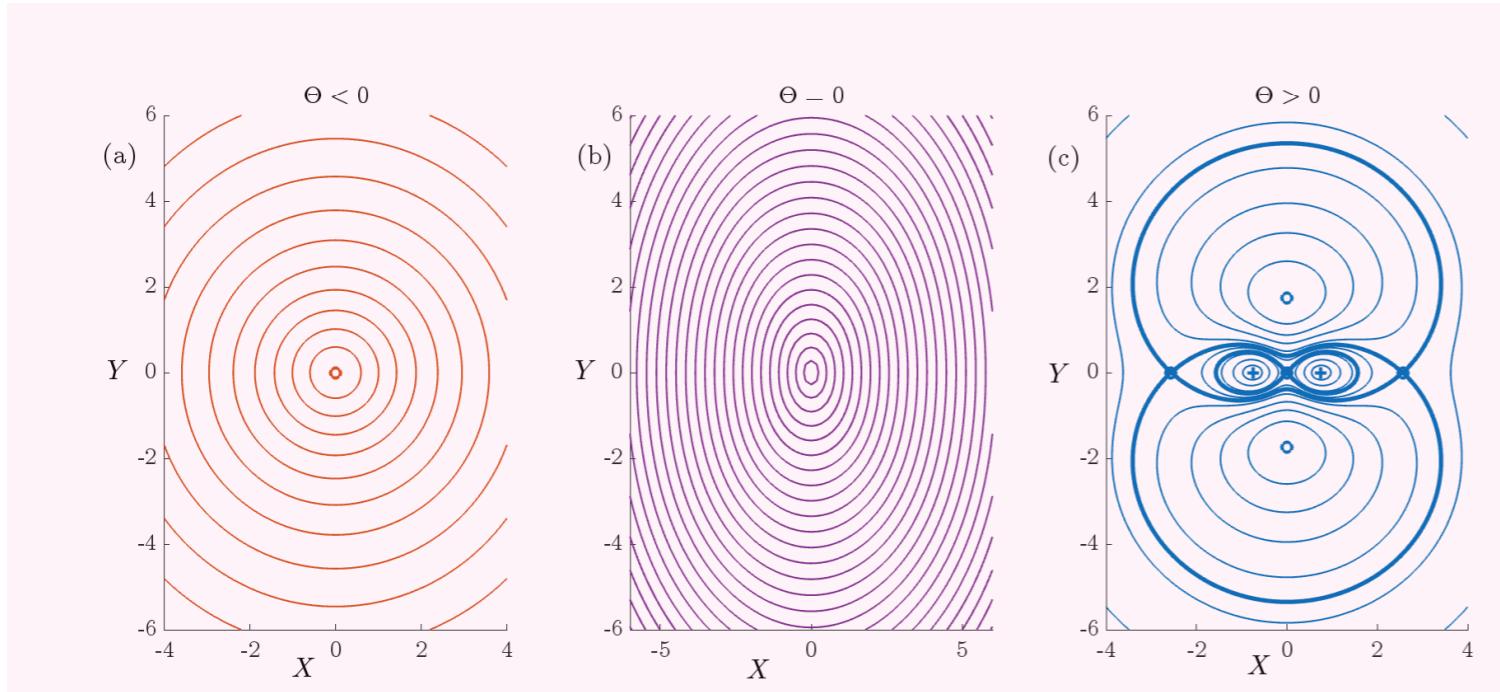
Global Phase Plane $\kappa_2 < 0$



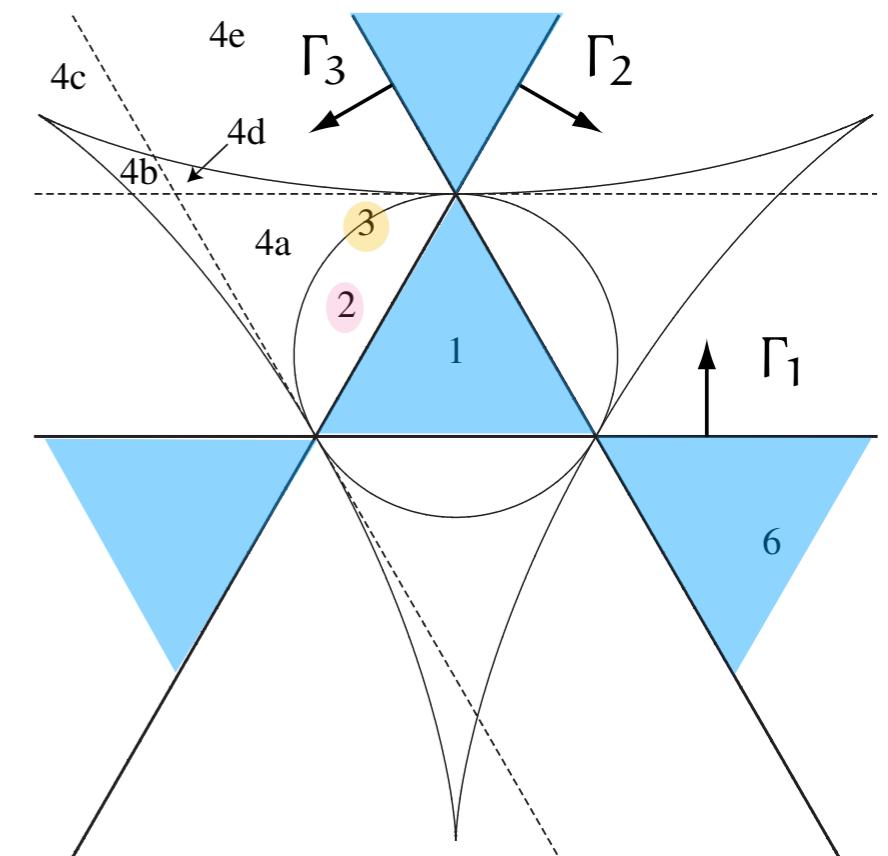
$$\Gamma = \left(\frac{5}{9}, \frac{5}{9}, -\frac{1}{9} \right)$$



Global Phase Plane $\kappa_2 < 0$



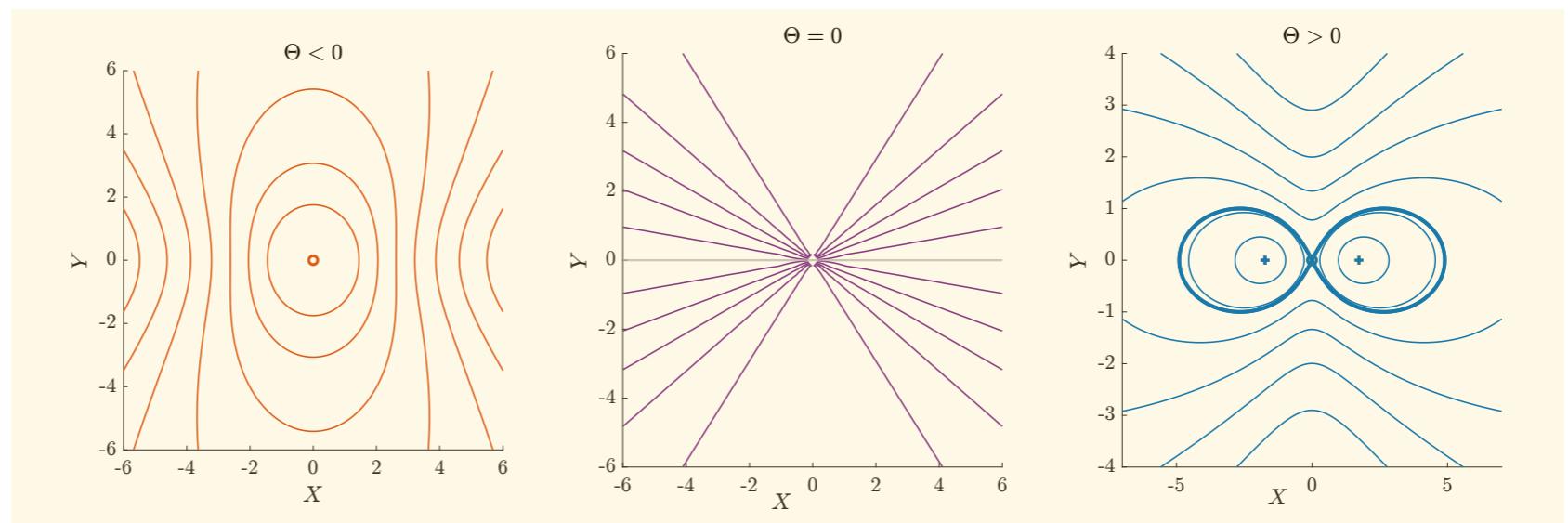
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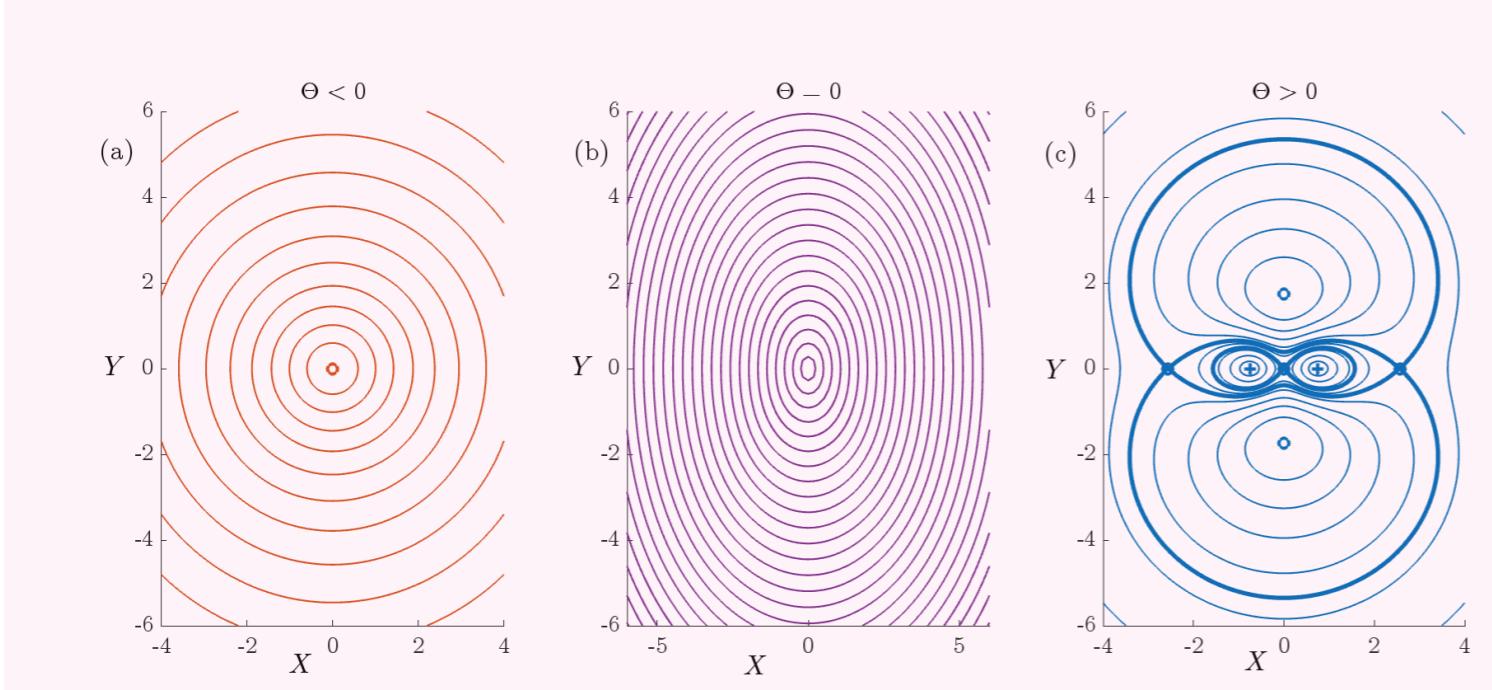
Vortex Collapse

$$\Gamma = \left(\frac{2}{3}, \frac{2}{3}, -\frac{1}{3} \right)$$

3 $\Gamma_1\Gamma_2 + \Gamma_2\Gamma_3 + \Gamma_1\Gamma_3 = 0$

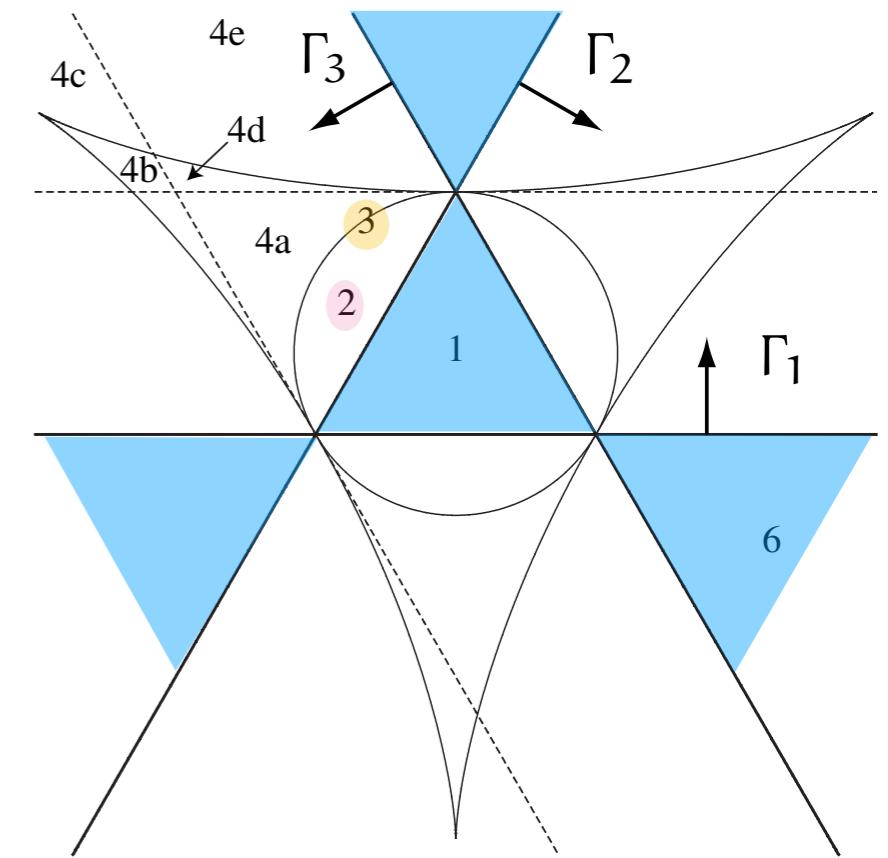


Global Phase Plane $\kappa_2 < 0$



$$\Gamma = \left(\frac{5}{9}, \frac{5}{9}, -\frac{1}{9} \right)$$

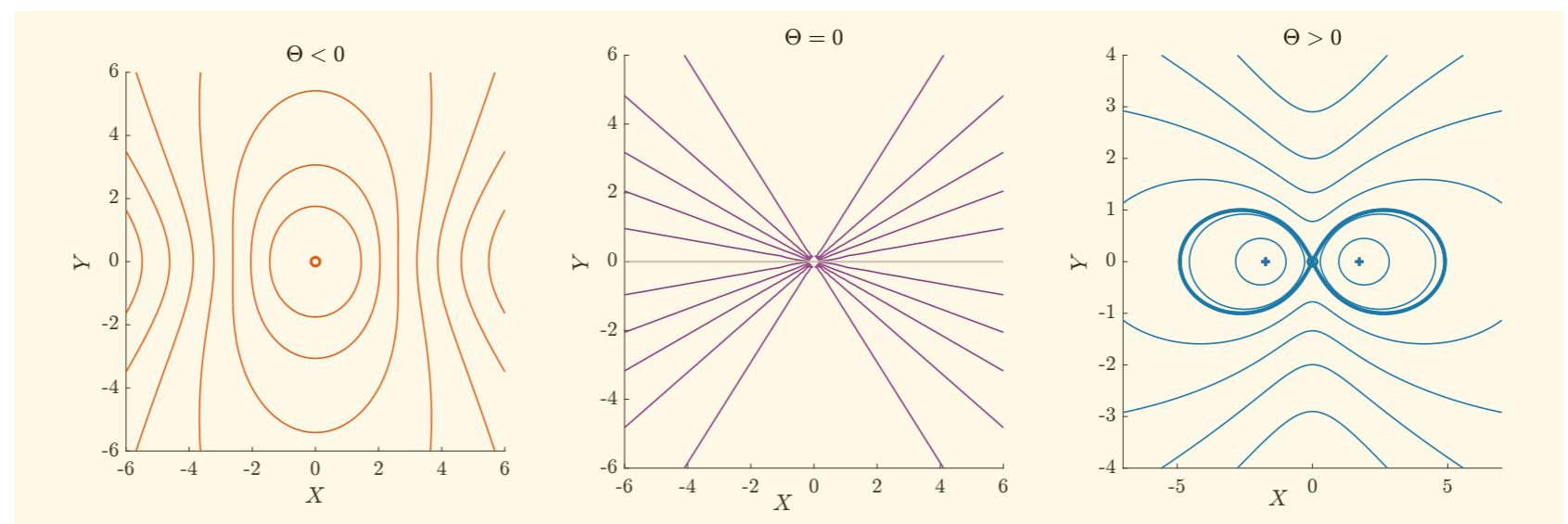
- alters both the topology and the linear stability of all 5 equilibria.
- the linear arrangements become unstable.
- The triangular arrangements in **Region 2** exhibit neutral stability.



Vortex Collapse

$$\Gamma = \left(\frac{2}{3}, \frac{2}{3}, -\frac{1}{3} \right)$$

3) $\Gamma_1\Gamma_2 + \Gamma_2\Gamma_3 + \Gamma_1\Gamma_3 = 0$



Conclusion

- Finish the stability analysis of the three-vortex (arbitrary circulation).
- Complete the case where circulation sums to zero.
- Review more literature to find more applications of the new coordinate system.

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Thank you!
Questions?