

**Esolution**

Place student sticker here

**Note:**

- During the attendance check a sticker containing a unique code will be put on this exam.
- This code contains a unique number that associates this exam with your registration number.
- This number is printed both next to the code and to the signature field in the attendance check list.

## Tensor Networks

**Exam:** IN2388 / Final Exam

**Date:** Monday 2<sup>nd</sup> August, 2021

**Examiner:** Christian Mendl

**Time:** 11:30 – 13:00

### Working instructions

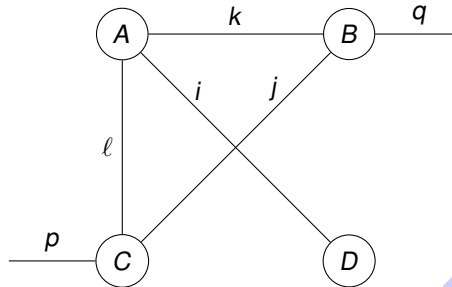
- This exam consists of **10 pages** with a total of **3 problems**.  
Please make sure now that you received a complete copy of the exam.
- The total amount of achievable credits in this exam is 60 credits.
- Detaching pages from the exam is prohibited.
- Allowed resources: open book
- Subproblems marked by \* can be solved without results of previous subproblems.
- **Answers are only accepted if the solution approach is documented.** Give a reason for each answer unless explicitly stated otherwise in the respective subproblem.
- Do not write with red or green colors nor use pencils.

## Problem 1 (20 credits)

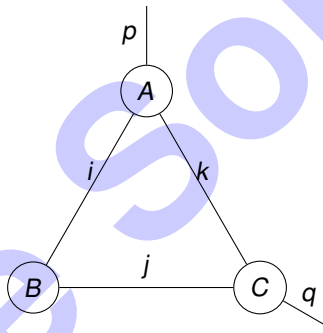
a) Represent the following contraction operation as graphical tensor diagram, labeling each tensor and tensor leg:

$$e_{pq} = \sum_{i,j,k,\ell} a_{\ell ik} b_{kjq} c_{j\ell p} d_i.$$

The legs should be drawn counter-clockwise with respect to the ordering of the tensor indices.



b)\* We consider the following to-be contracted tensor network:



Find an optimal contraction order that minimizes contraction complexity, i.e., the overall computational cost, assuming that the dimensions of the tensor legs obey

$$\dim(p) = \dim(q) =: \ell \ll \dim(j) = \dim(k) =: m \ll \dim(i) =: n,$$

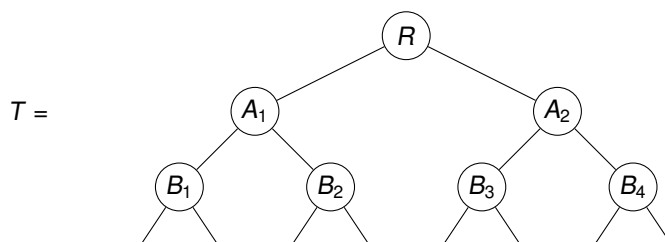
and that at each contraction step, two tensors are contracted together. (An optimality proof is not required here.)

Since  $i$  has the largest dimension, one should contract this bond first. The optimal order is thus  $i$  (i.e., tensor  $A$  with  $B$ ), followed by  $(j, k)$  (i.e.,  $C$  with the output of the first contraction).

c) What is the asymptotic computational cost (in  $\mathcal{O}$ -notation) of the contractions you found in (b)? Assume a literal implementation of contractions based on the summation formulation.

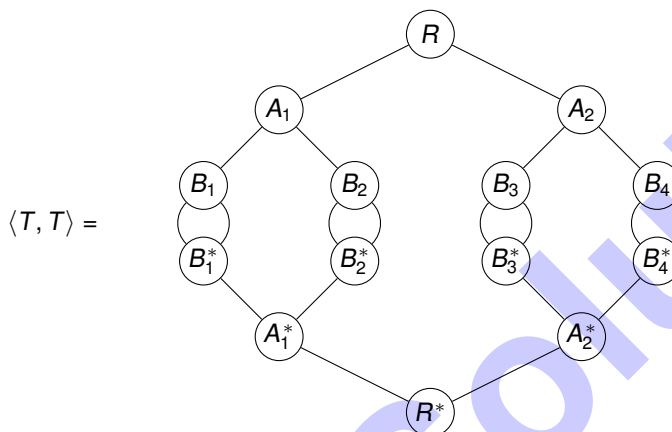
The contraction complexity of the  $A$ - $B$  contraction is  $\mathcal{O}(\ell m^2 n)$ , and of the second contraction  $\mathcal{O}(\ell^2 m^2)$ .

d)\* We define  $T$  as the following tree tensor network:

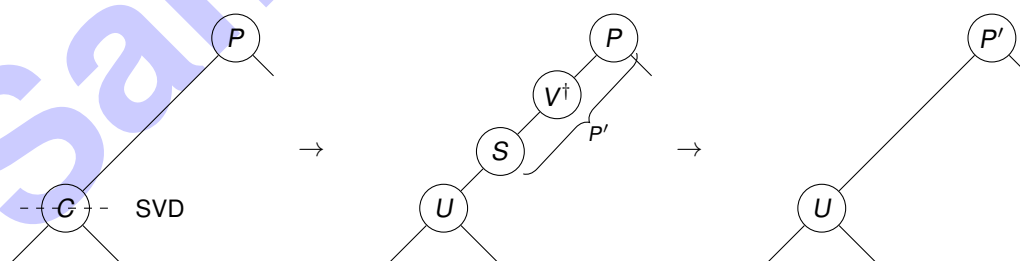


	0
	1
	2
	3

Draw the tensor diagram for evaluating  $\langle T, T \rangle$ .



e) In the context of part (d), we consider a local SVD-splitting operation performed on a node  $C$  of the tree (parent node called  $P$ ), which updates both the child and parent node:



	0
	1
	2
	3
	4
	5
	6

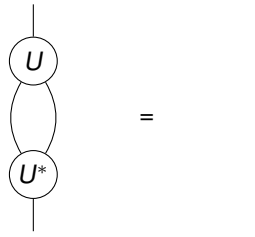
This local operation is applied to each node in the tree (except the root) given a particular ordering. We distinguish between two realizations:

(i) The operation is first performed on the  $A$  tensors in the middle layer before proceeding to the  $B$  tensors.

(ii) The operation is first performed on the  $B$  tensors in the bottom layer before proceeding to the  $A$  tensors.

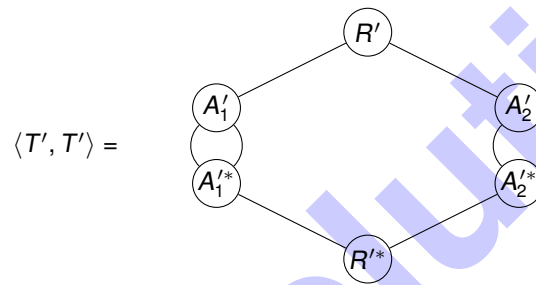
We denote the resulting tree tensor network by  $T'$ . For each of the two cases, simply the tensor diagram for evaluating  $\langle T', T' \rangle$  as far as possible. Also provide a short explanation of your simplifications.

After each local operation, the updated child node  $U$  is an isometry, such that (interpreted as matrix)  $U^\dagger U = I$ . Graphically:

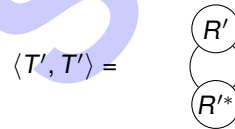


However, the parent node is (in general) no longer an isometry after performing the operation.

- (i) In the first case, performing the local operations on the  $B$  tensors causes the parent  $A$  tensors to lose their isometric property. Hence, the inner product simplifies to the following contraction:



- (ii) In the second case, the parent nodes are always updated after their respective children. This ensures that all nodes are isometries except for the root node. Hence, the inner product  $\langle T', T' \rangle$  is equal to:



## Problem 2 (20 credits)

We consider the following Hamiltonian on a one-dimensional lattice with  $L$  sites and open boundary conditions, where  $J$  is a real parameter:

$$H = -J \sum_{j=2}^{L-1} X_{j-1} Y_j Z_{j+1}.$$

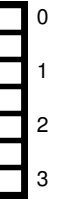
a) What is the matrix dimension of  $H$ ?

$H$  is a  $2^L \times 2^L$  matrix.



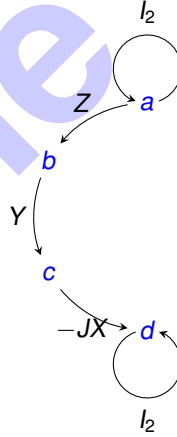
b)\* For a given index  $j \in \{2, \dots, L-1\}$ , represent  $X_{j-1} Y_j Z_{j+1}$  in terms of Kronecker products of identity matrices and the Pauli matrices  $X$ ,  $Y$ ,  $Z$ .

$$X_{j-1} Y_j Z_{j+1} = \underbrace{I_2 \otimes \dots \otimes I_2}_{j-2 \text{ terms}} \otimes X \otimes Y \otimes Z \otimes \underbrace{I_2 \otimes \dots \otimes I_2}_{L-j-1 \text{ terms}}$$



c)\* Construct a finite state automaton and corresponding MPO tensors for representing  $H$  as matrix product operator. You should separately specify the MPO tensors  $A^j$  for  $j = 2, \dots, L-1$  and the boundary tensors  $A^1, A^L$ . Also provide the dimensions of these tensors.

An automaton with four states, denoted  $a, b, c, d$ , and the following transitions generates the Hamiltonian:

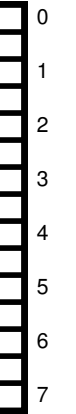


(The prefactor  $-J$  can equivalently be attached to  $Y$  or  $Z$  as well.)  
The corresponding MPO tensors are then

$$A^1 = \begin{pmatrix} 0 & 0 & -JX & I_2 \end{pmatrix}, \quad A^j = \begin{pmatrix} I_2 & 0 & 0 & 0 \\ Z & 0 & 0 & 0 \\ 0 & Y & 0 & 0 \\ 0 & 0 & -JX & I_2 \end{pmatrix} \text{ for } j = 2, \dots, L-1, \quad A^L = \begin{pmatrix} I_2 \\ Z \\ 0 \\ 0 \end{pmatrix}.$$

$A^1$  is equal to the last row of  $A^j$  (corresponding to  $d$  as final state), and  $A^L$  equal to the first column of  $A^j$  (corresponding to  $a$  as initial state).

With the convention that the leading two dimensions are the physical dimensions,  $A^1 \in \mathbb{C}^{2 \times 2 \times 1 \times 4}$ ,  $A^j \in \mathbb{C}^{2 \times 2 \times 4 \times 4}$  for  $j = 2, \dots, L-1$ , and  $A^L \in \mathbb{C}^{2 \times 2 \times 4 \times 1}$ .



0 ☐ d)\* Specify an operator of the form  $P_k = A_{k-1} B_k C_{k+1}$  with each  $A, B, C$  a Pauli matrix and  $k \in \{2, \dots, L-1\}$ , such that  $P_k$  commutes with  $H$ . (A proof of the commuting property is not required.)

1 ☐

2 ☐  $P_k = Z_{k-1} Y_k X_{k+1}$  has the desired property, since it commutes with each summand  $X_{j-1} Y_j Z_{j+1}$  of  $H$ . This is a consequence of the anti-commutation relation of the Pauli matrices:

3 ☐  $XY = -YX, \quad YZ = -ZY, \quad ZX = -XZ.$

4 ☐

0 ☐ e)\* Consider the partitioning of the Hamiltonian as  $H = H_a + H_b$  with

$$H_a = -J \sum_{j=2, \text{mod}(j,4)=2}^{L-2} (X_{j-1} Y_j Z_{j+1} + X_j Y_{j+1} Z_{j+2}) \quad \text{and} \quad H_b = -J \sum_{j=4, \text{mod}(j,4)=0}^{L-2} (X_{j-1} Y_j Z_{j+1} + X_j Y_{j+1} Z_{j+2}).$$

1 ☐

2 ☐

3 ☐ Is it possible to represent the matrix exponentials  $e^{-iH_a t}$  or  $e^{-iH_b t}$  (with  $t \in \mathbb{R}$ ) exactly in quantum circuit form, assuming that you can use arbitrary one-, two- and three-qubit gates? Briefly justify your answer.

4 ☐

Yes, this is possible since the individual terms in  $H_a$  pairwise commute, such that  $e^{-iH_a t}$  becomes a product of individual matrix exponentials  $e^{JX_{j-1} Y_j Z_{j+1} t}$ , which can be realized as three-qubit gates. The same argument applies to  $e^{-iH_b t}$ .

We define a linear map  $\mathcal{E} : \mathbb{C}^{n \times n} \rightarrow \mathbb{C}^{n \times n}$  ( $n \in \mathbb{N}$ ) as

$$\mathcal{E}(\rho) = \sum_{j=1}^s K_j \rho K_j^\dagger$$

with given matrices  $K_j \in \mathbb{C}^{n \times n}$  satisfying  $\sum_{j=1}^s K_j^\dagger K_j = I$ .

- (i) Hermitian matrices to Hermitian matrices, and
- (ii) positive semi-definite matrices to positive semi-definite matrices.

(i) Let  $\rho \in \mathbb{C}^{n \times n}$  be a Hermitian matrix,  $\rho^\dagger = \rho$ , then

$$\mathcal{E}(\rho)^\dagger = \sum_{j=1}^s \left( K_{j\rho} K_j^\dagger \right)^\dagger = \sum_{j=1}^s (K_j^\dagger)^\dagger \rho^\dagger K_j^\dagger = \sum_{j=1}^s K_{j\rho} K_j^\dagger = \mathcal{E}(\rho),$$

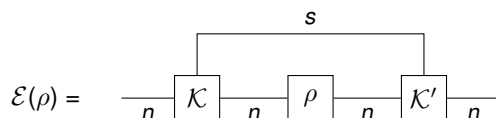
that is,  $\mathcal{E}(\rho)$  is Hermitian as well. Here we have used that  $(AB)^\dagger = B^\dagger A^\dagger$  for any two matrices  $A, B$  of compatible dimensions, and that  $(A^\dagger)^\dagger = A$ .

(ii) Let  $\rho \in \mathbb{C}^{n \times n}$  be a positive semi-definite matrix, i.e.,  $\langle v, \rho v \rangle \geq 0 \ \forall v \in \mathbb{C}^n$ . We have to show that  $\langle v, \mathcal{E}(\rho)v \rangle \geq 0 \ \forall v \in \mathbb{C}^n$ :

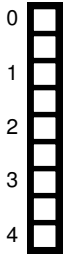
$$\langle v, \mathcal{E}(\rho)v \rangle = \sum_{j=1}^s \left\langle v, K_{j\rho} K_j^\dagger v \right\rangle = \sum_{j=1}^s \left\langle K_j^\dagger v, \rho K_j^\dagger v \right\rangle = \sum_{j=1}^s \langle \varphi_j, \rho \varphi_j \rangle \geq 0,$$

where we have defined  $\varphi_j = K_j^\dagger v$  for  $j = 1, \dots, s$ .

b)\* For the following, we define two new tensors  $\mathcal{K}$  and  $\mathcal{K}'$  of degree 3 by  $\mathcal{K}_{j,\dots} = K_j$  and  $\mathcal{K}'_{j,\dots} = K_j^\dagger$  for all  $j = 1, \dots, s$ . Draw the tensor network representing the application of  $\mathcal{E}$  to  $\rho$  in terms of  $\rho$ ,  $\mathcal{K}$  and  $\mathcal{K}'$ , clearly indicating the dimension of each leg.



The dimension of each horizontal leg is  $n$ , and the bond connecting  $\mathcal{K}$  with  $\mathcal{K}'$  has dimension  $s$ .



c) Given a unitary matrix  $V = (v_{ij}) \in \mathbb{C}^{s \times s}$ , we introduce the new linear map  $\mathcal{F} : \mathbb{C}^{n \times n} \rightarrow \mathbb{C}^{n \times n}$  by

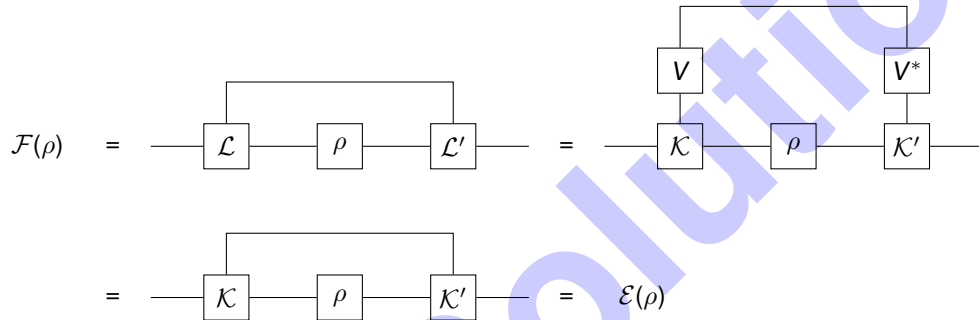
$$\mathcal{F}(\rho) = \sum_{i=1}^s L_i \rho L_i^\dagger \quad \text{with } L_i = \sum_{j=1}^s v_{ij} K_j \text{ for } i = 1, \dots, s.$$

Show that  $\mathcal{E}$  and  $\mathcal{F}$  represent the same map.

Similar to  $\mathcal{E}$ , we can represent  $\mathcal{F}$  as tensor network using analogous degree-3 tensors  $\mathcal{L}_{i,:,:} = L_i$  and  $\mathcal{L}'_{i,:,:} = L_i^\dagger$ . By definition,  $\mathcal{L}$  is related to  $\mathcal{K}$  and  $\mathcal{L}'$  to  $\mathcal{K}'$  via



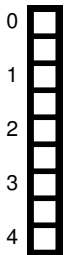
Thus, for any  $\rho \in \mathbb{C}^{n \times n}$ ,



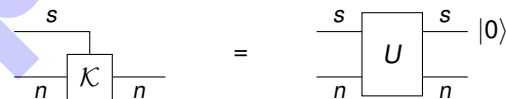
For the third equal sign we have used that  $V$  is unitary.

Alternative solution:

$$\mathcal{F}(\rho) = \sum_{i=1}^s L_i \rho L_i^\dagger = \sum_{i,j,j'=1}^s v_{ij} v_{ij'}^* K_j \rho K_{j'}^\dagger = \sum_{j,j'=1}^s \left( \sum_{i=1}^s v_{ij} v_{ij'}^* \right) K_j \rho K_{j'}^\dagger = \sum_{j=1}^s K_j \rho K_j^\dagger = \mathcal{E}(\rho).$$

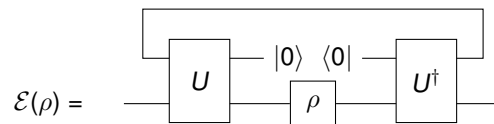


d)  $\mathcal{K}$  interpreted as  $(sn) \times n$  matrix is an isometry by definition; we can thus extend it to a unitary matrix  $U \in \mathbb{C}^{(sn) \times (sn)}$  such that (with labels at the legs denoting dimensions, and  $|0\rangle$  the first unit vector in  $\mathbb{C}^s$ ):



Draw a tensor diagram that represents  $\mathcal{E}(\rho)$  in terms of  $\rho$ ,  $U$ ,  $U^\dagger$  and  $|0\rangle$ .

The tensor  $\mathcal{K}'$ , or equivalently  $K_j^\dagger$  for  $j = 1, \dots, s$ , can be expressed by taking the adjoint of  $U$ . One arrives at:



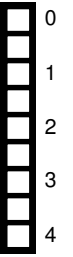


e) Based on your solution of part (d), express  $\mathcal{E}(\rho)$  using a partial trace operation.

The diagram in part (d) can be interpreted as “tracing out” the top register with dimension  $s$ . As formula:

$$\mathcal{E}(\rho) = \text{tr}_1 [U(|0\rangle\langle 0| \otimes \rho) U^\dagger]$$

Remark: This is a derivation of Stinespring’s dilation theorem in finite dimensions.



Sample Solution

Additional space for solutions—clearly mark the (sub)problem your answers are related to and strike out invalid solutions.

A large grid of graph paper for solutions, with a diagonal watermark reading "Sample Solution".