

Modeling Durotaxis using Mechanical Bidomain model

November 13, 2018

Analytical treatment - 10/03/2018

The equations formulated below consider a gradient in the extracellular stiffness in a 1-d strand of tissue using the mechanical bidomain model.

Consider u_x and w_x as the intra- and extra-cellular displacements of the bidomain layer. x being the principal direction of the one-dimensional tissue strand, T being the tension, p and q being the intra- and extra-cellular pressure contributions arising mostly from hydrostatic forces.

The standard equations of the bidomain model are as follows [1]:

$$-\frac{\partial p}{\partial x} + \nu(\frac{\partial^2 u_x}{\partial x^2} + \frac{\partial^2 u_x}{\partial y^2}) + \gamma \frac{\partial^2 u_x}{\partial x^2} + \frac{\partial T}{\partial x} = K(u_x - w_x) \quad (1)$$

$$-\frac{\partial q}{\partial x} + \mu(\frac{\partial^2 w_x}{\partial x^2} + \frac{\partial^2 w_x}{\partial y^2}) = -K(u_x - w_x) \quad (2)$$

For this problem we consider:

- $\mu = \mu_0 + gx$
- No contributions along the y axis
- p, q and T are constant and hence their gradients are zero along the x axis
- Slope of μ is a constant
- K is same for intra- and extra-cellular displacements

$$\nu \frac{\partial^2 u_x}{\partial x^2} + \gamma \frac{\partial^2 u_x}{\partial x^2} = K(u_x - w_x) \quad (3)$$

$$\mu_0 \frac{\partial^2 w_x}{\partial x^2} + gx \frac{\partial^2 w_x}{\partial x^2} = -K(u_x - w_x) \quad (4)$$

Simplifying the above expression, replacing u_x in equation (2) from (1):

$$u_x = w_x + \frac{\nu}{K} \frac{\partial^2 u_x}{\partial x^2} + \frac{\gamma}{K} \frac{\partial^2 u_x}{\partial x^2} + \gamma \frac{\partial^2 u_x}{\partial x^2} \quad (5)$$

$$\therefore \frac{\mu_0}{K} \frac{\partial^2 w_x}{\partial x^2} + x \frac{g}{K} \frac{\partial^2 w_x}{\partial x^2} = -(\frac{\nu}{K} \frac{\partial^2 u_x}{\partial x^2} + \frac{\gamma}{K} \frac{\partial^2 u_x}{\partial x^2}) \quad (6)$$

Final form the equation can be written as:

$$\frac{\partial^2 w_x}{\partial x^2} (\frac{\mu_0}{K} + \frac{gx}{K}) + \frac{\partial^2 u_x}{\partial x^2} (\frac{\nu}{K} + \frac{\gamma}{K}) = 0 \quad (7)$$

Analytical treatment - 10/04/2018

The equations formulated below consider a gradient in the extracellular stiffness in a 1-d strand of tissue using the mechanical bidomain model.

Consider u_x and w_x as the intra- and extra-cellular displacements of the bidomain layer. x being the principal direction of the one-dimensional tissue strand, T being the tension, p and q being the intra- and extra-cellular pressure contributions arising mostly from hydrostatic forces.

The intra and extra-cellular stresses arising in the 1-d strand of tissue can be expressed as [2]:

$$\tau_{ix} = -p + 2\nu\epsilon_{ix} + T \quad (8)$$

$$\tau_{ex} = -q + 2\mu(x)\epsilon_{ex} \quad (9)$$

Relationship between the strains and displacement can be written as:

$$\frac{\partial \tau_{ix}}{\partial x} = K(u_x - w_x) \quad (10)$$

$$\frac{\partial \tau_{ex}}{\partial x} = -K(u_x - w_x) \quad (11)$$

Using equation (8) and (9) in (10) and (11) the resulting intra- and extra-cellular equations are:

$$\frac{\partial}{\partial x}(-p + 2\nu\frac{\partial u_x}{\partial x}) = K(u_x - w_x) \quad (12)$$

$$\frac{\partial}{\partial x}(-q + 2\mu(x)\frac{\partial w_x}{\partial x}) = -K(u_x - w_x) \quad (13)$$

For the problem we have assumed $\mu = \mu_0 + gx$, where g is a constant. Rewriting (12) and (13) we have:

$$-\frac{\partial p}{\partial x} + 2\nu\frac{\partial^2 u_x}{\partial x^2} = K(u_x - w_x) \quad (14)$$

$$-\frac{\partial q}{\partial x} + 2(g\frac{\partial w_x}{\partial x} + \frac{\partial^2 w_x}{\partial x^2}\mu(x)) = -K(u_x - w_x) \quad (15)$$

Analytical treatment - 11/13/2018

In order to obtain an analytical solution for the problem initially we consider $g = 0$. The boundary limits of the 1-dimensional problem assume the length of the domain spanning from $x=-L$ to $x=+L$. The stresses are taken to be zero at each of the boundaries. We begin by implementing a trial solution for u_x and w_x .

$$u_x = Ax + B \sinh\left(\frac{x}{\sigma}\right) \quad (16)$$

$$w_x = Cx + D \sinh\left(\frac{x}{\sigma}\right) \quad (17)$$

For estimating the values of A,B,C,D we assume additionally $p=q=0$. Rewriting (12):

$$2\nu \frac{\partial^2}{\partial x^2} \left(Ax + B \sinh\left(\frac{x}{\sigma}\right) \right) = K \left(Ax + B \sinh\left(\frac{x}{\sigma}\right) - Cx - D \sinh\left(\frac{x}{\sigma}\right) \right) \quad (18)$$

$$2\nu \frac{\partial}{\partial x} \left(A + \frac{B}{\sigma} \cosh\left(\frac{x}{\sigma}\right) \right) = K \left((A - C)x + (B - D) \sinh\left(\frac{x}{\sigma}\right) \right) \quad (19)$$

$$2\nu \frac{B}{\sigma^2} \sinh\left(\frac{x}{\sigma}\right) = K \left((A - C)x + (B - D) \sinh\left(\frac{x}{\sigma}\right) \right) \quad (20)$$

and (13)

$$2\mu(x) \frac{\partial^2}{\partial x^2} \left(Cx + D \sinh\left(\frac{x}{\sigma}\right) \right) = -K \left(Ax + B \sinh\left(\frac{x}{\sigma}\right) - Cx - D \sinh\left(\frac{x}{\sigma}\right) \right) \quad (21)$$

$$2\mu(x) \frac{\partial}{\partial x} \left(C + \frac{D}{\sigma} \cosh\left(\frac{x}{\sigma}\right) \right) = -K \left((A - C)x + (B - D) \sinh\left(\frac{x}{\sigma}\right) \right) \quad (22)$$

$$2\mu(x) \frac{D}{\sigma^2} \sinh\left(\frac{x}{\sigma}\right) = -K \left((A - C)x + (B - D) \sinh\left(\frac{x}{\sigma}\right) \right) \quad (23)$$

From (20) and (23) we have:

$$D = -\frac{\nu}{\mu(x)} B \quad (24)$$

Comparing coefficients of hyperbolic sine terms we see that (20) and (23) is satisfied only when the coefficient of the linear term is zero. Therefore,

$$A = C \quad (25)$$

Hence the length constant¹ σ has the value

$$\sigma = \sqrt{\frac{2\nu\mu}{K(\nu + \mu)}} \quad (26)$$

¹This might also be a variable since μ is a function of x for the problem

Bibliography

- [1] Bradley J Roth *The Mechanical Bidomain Model: A Review* ISRN Tissue Eng. 2013
- [2] Sharma, Kharananda; Al-Asuoad, Nofe; Shillor, Meir; Roth, Bradley J. *Intracellular, extracellular, and membrane forces in remodeling and mechanotransduction: The mechanical bidomain model* Journal of Coupled Systems and Multiscale Dynamics, Volume 3, Number 3, September 2015, pp. 200-207(8)