

## Block (2/3)

# The Finite Element Method for the Diffusion Equation in 1D

**EE4375 - FEM For EE Applications**

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## 1D FEM: (1/11) One-Dim Problem to Solve (1/2)

Same as Before: Boundary Value Problem  
Second Order Differential Equations

- **given**  $x \in \Omega = (0, 1)$  with left and right boundary point
- **given**:  $f(x)$  given function and  $\alpha$  given number
- **find**:  $u(x)$  such that

$-u''(x) = f(x)$  for  $0 < x < 1$  (differential equation on  $\Omega$ )

$u(x = 0) = 0$  (Dirichlet boundary condition in  $x = 0$ )

$\frac{du}{dx}(x = 1) = \alpha$  (Neumann boundary condition in  $x = 1$ )

- differential equation **not** valid in  $x = 0$  or  $x = 1$

boundary conditions **are** valid in  $x = 0$  or  $x = 1$

# 1D FEM: (1/11) One-Dim Problem to Solve (2/2)

## Residual function $r(x)$

- **definition:**  $r(x) = u''(x) + f(x)$
- **quality of the solution:**  
small (large) residual is indication of good (poor) approximation
- **solve**  $-u''(x) = f(x) + \text{b.c.}$  equivalent to
- **find**  $u(x)$  such that  $r(x) = 0$  and  $u(x)$  satisfies b.c.

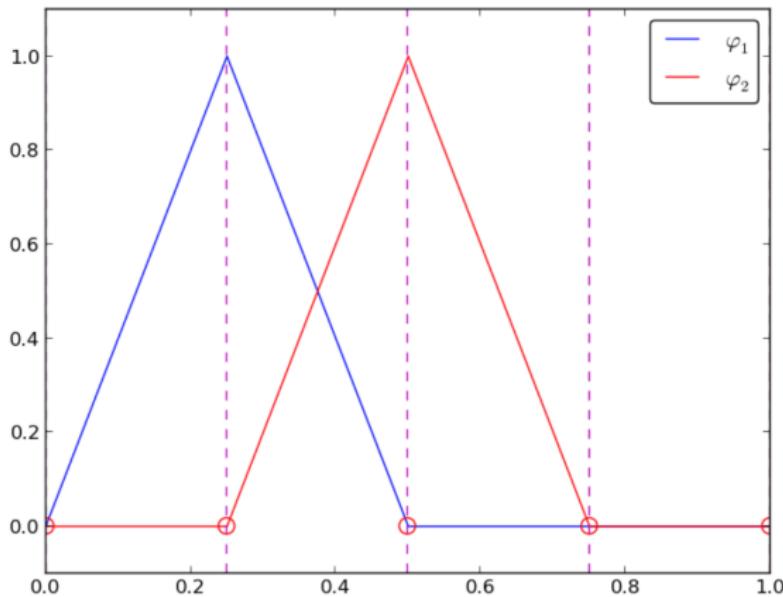
# 1D FEM: (2/11) Mesh and Shape Functions (1/9)

$$\text{Mesh } \Omega^h = \{x_i \mid 1 \leq i \leq N + 1\}$$

- create mesh on  $\Omega = (0, 1)$
- $(N + 1)$  points or **nodes**:  $x_i$  for  $1 \leq i \leq N + 1$
- $N$  intervals or **elements**:  $e_i = [x_i, x_{i+1}]$  for  $1 \leq i \leq N$
- element size  $h_i = x_{i+1} - x_i$
- mesh uniform in case that  $h_i$  is independent of  $i$   
non-uniform otherwise
- Let op (Dutch) - Dhyaan (Hindi) - Zhuyi li (Chinese) - Alaintibah (Arabic):  
 $N$  elements  $e_i$ -  $(N+1)$  nodes  $x_i$  at which 2 on the boundary

# 1D FEM: (2/11) Mesh and Shape Functions (2/9)

## Shape Functions



# 1D FEM: (2/11) Mesh and Shape Functions (3/9)

Definition of shape function  $\phi_i(x)$

- linear Lagrange interpolation function
- for  $x_{i-1} \leq x \leq x_i$  (left of  $x_i$ ):  $\phi_i(x) = \frac{x - x_{i-1}}{x_i - x_{i-1}}$  (positive slope)
- for  $x_i \leq x \leq x_{i+1}$  (right of  $x_i$ ):  $\phi_i(x) = \frac{x_{i+1} - x}{x_{i+1} - x_i}$  (negative slope)
- for  $x \leq x_{i-1}$  or  $x \geq x_{i+1}$ :  $\phi_i(x) = 0$
- each node  $x_i$  (including 2 boundary nodes) has its  $\phi_i(x)$
- mesh defines in total  $N+1$  functions  $\phi_i(x)$
- $\phi_i(x_j) = \delta_{ij}$  (see figure)

# 1D FEM: (2/11) Mesh and Shape Functions (4/9)

## Mesh Terminology

1D Domains  $\Omega$

- two nodes  $x_i$
- interval  $e_i = [x_i, x_{i+1}]$
- interval length  $h_i = x_{i+1} - x_i$

2D Domains  $\Omega$

- three nodes  $x_i$
- triangle formed by  $x_i$
- area of triangle

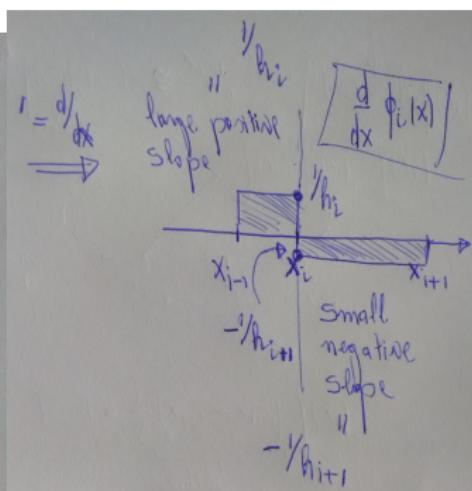
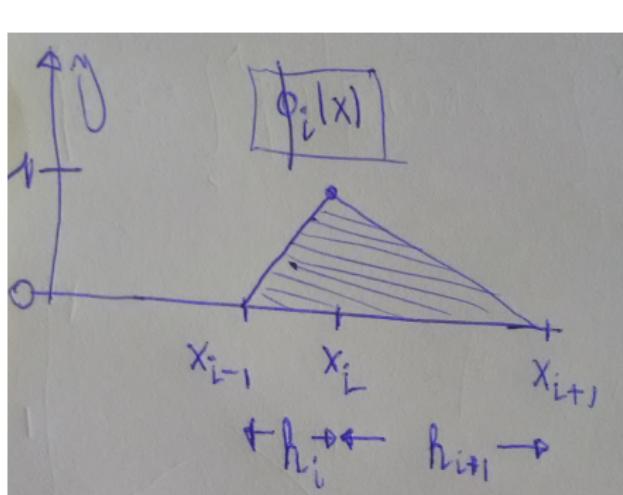
# 1D FEM: (2/11) Mesh and Shape Functions (5/9)

## Shape functions

- the mesh **determines** the set of shape functions
- construction carries over to **non-uniform** meshes
- **one** element  $e_i = [x_i, x_{i+1}]$  "sees" **two** basis functions
  - left node  $x_i$  has basis function  $\phi_i(x)$
  - right node  $x_{i+1}$  has basis function  $\phi_{i+1}(x)$
  - all other basis function  $\phi_j(x)$  with  $j \neq i$  or  $j \neq i + 1$  are zero on  $e_i$

# 1D FEM: (2/11) Mesh and Shape Functions (6/9)

Shape function  $\phi_i(x)$  (left) - derivative  $\phi'_i(x) = \frac{d}{dx}\phi_i(x)$  (right)



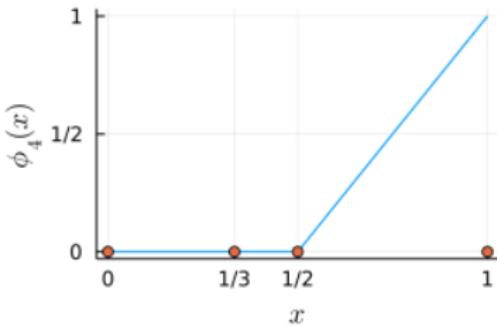
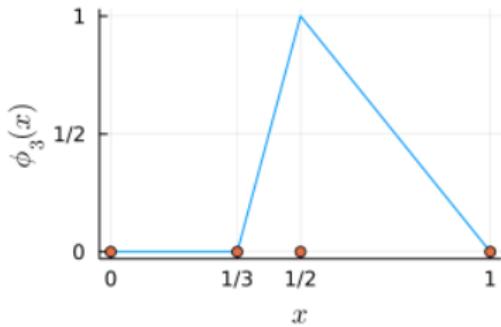
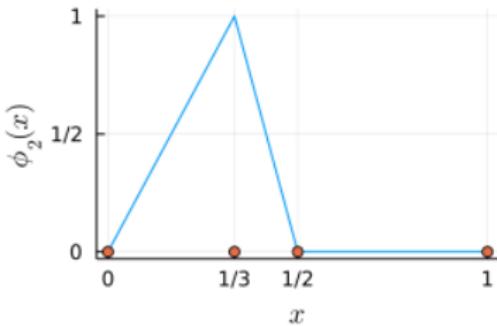
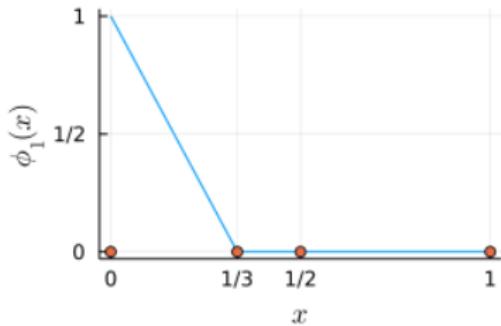
# Mesh and Shape Functions (7/9)

## Exercise

- assume a mesh with the four nodes  
 $x_1 = 0$ ,  $x_2 = 1/3$ ,  $x_3 = 1/2$  and  $x_4 = 1$ ;
- plot the mesh and label the nodes  $x_1$ ,  $x_2$ ,  $x_3$  and  $x_4$ ;
- plot the shape functions  $\phi_1(x)$ ,  $\phi_2(x)$ ,  $\phi_3(x)$  and  $\phi_4(x)$ ;
- plot the first derivative of the shape functions  $\frac{d\phi_1(x)}{dx}$ ,  $\frac{d\phi_2(x)}{dx}$ ,  
 $\frac{d\phi_3(x)}{dx}$  and  $\frac{d\phi_4(x)}{dx}$ ;

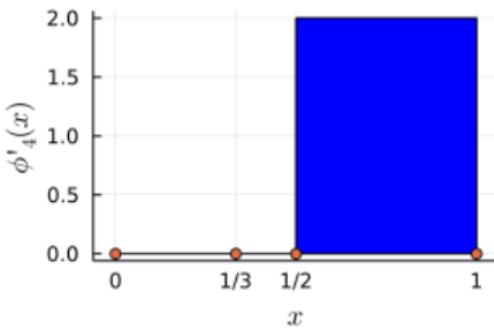
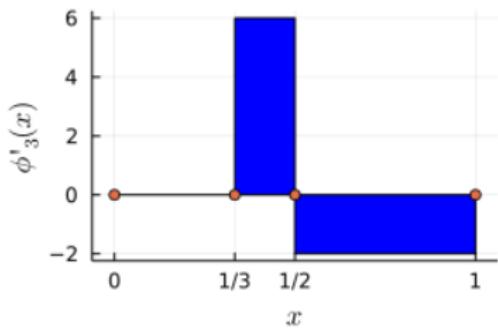
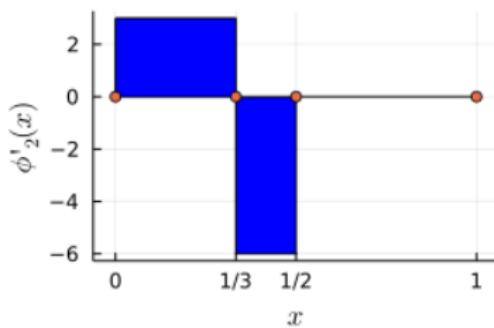
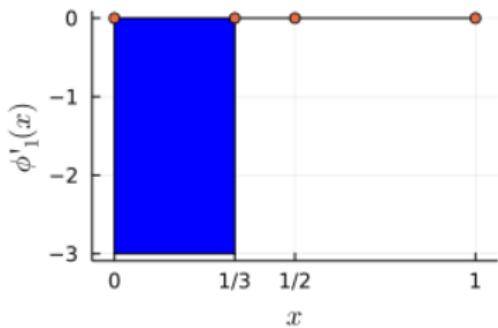
# Mesh and Shape Functions (8/9)

## Solution (1/2)



# Mesh and Shape Functions (9/9)

## Solution (1/2)



# 1D FEM: (3/11) Linear Combination of Shape Functions (1/5)

## Linear Combination of Shape Functions

- set of functions  $\{\phi_i(x) | 1 \leq i \leq N + 1\}$
- linear combinations of these functions  $\phi_i(x)$  can be made
- $V_0^h(\Omega)$ : function space defined by all linear combinations

$$V_0^h(\Omega) = \text{span}\{\phi_1(x), \dots, \phi_{N+1}(x)\}$$

- $u^h(x) \in V_0^h(\Omega)$ : there exists coordinates  $c_1, \dots, c_{N+1}$  such that

$$u^h(x) = c_1 \phi_1(x) + \dots + c_{N+1} \phi_{N+1}(x)$$

# 1D FEM: (3/11) Linear Combination of Shape Functions (2/5)

## Analogy with Geometry

- geometry in 2D plane:  $\mathbf{x} = x \mathbf{i} + y \mathbf{j}$
- $x$  and  $y$  coordinates -  $\mathbf{i} = (1, 0)$  and  $\mathbf{j} = (0, 1)$  basis vectors
- geometry in 3D space:  $\mathbf{x} = x \mathbf{i} + y \mathbf{j} + z \mathbf{k}$
- $x, y$  and  $z$ : coordinates
- $\mathbf{i} = (1, 0, 0), \mathbf{j} = (0, 1, 0)$  and  $\mathbf{k} = (0, 0, 1)$ : basis vectors

# 1D FEM: (3/11) Linear Combination of Shape Functions (3/5)

Analogy with Fourier Analysis

- $f(x)$ : given function on domain  $x \in \Omega = (0, 1)$
- $f^N(x)$ : approximation of  $f(x)$  with all harmonics up to order  $N$

$$f^N(x) = a_0 + \sum_{k=1}^N a_k \cos(k \pi x) + b_k \sin(k \pi x)$$

- $a_0, a_k$  and  $b_k$  for  $1 \leq k \leq N$ : **coordinates**
- $1, \cos(k \pi x)$  and  $\sin(k \pi x)$ : Fourier modes or **basis functions**

# 1D FEM: (3/11) Linear Combination of Shape Functions (4/5)

## Application to Finite Elements

- $u(x)$ : exact solution of the boundary value problem
- $u^h(x)$ : finite element approximation to  $u(x)$  computed on  $\Omega^h$
- $u^h(x) \in V_0^h(\Omega) = \text{span}\{\phi_1(x), \dots, \phi_{N+1}(x)\}$
- expansion of  $u^h(x)$  as linear combination of shape function

$$u^h(x) = c_1 \phi_1(x) + \dots + c_{N+1} \phi_{N+1}(x)$$

- $c_1, \dots, c_{N+1}$  coordinates -  $\phi_1(x), \dots, \phi_{N+1}(x)$  basis functions
- non-uniform meshes can be accommodated

# 1D FEM: (3/11) Linear Combination of Shape Functions (5/5)

## Exercises

- assume a non-uniform mesh on  $\Omega$  with 4 nodes  $x = 0$ ,  $x = 1/3$ ,  $x = 1/2$  and  $x = 1$
- plot  $\phi_1(x)$ ,  $\phi_2(x)$ ,  $\phi_3(x)$  and  $\phi_4(x)$
- plot linear combinations of  $\phi_1(x)$ ,  $\phi_2(x)$ ,  $\phi_3(x)$  and  $\phi_4(x)$
- assume a non-uniform mesh on  $\Omega$  with  $(N+1)$  nodes
- plot  $\phi_i(x) + \phi_{i+1}(x)$  on  $e_i$

# 1D FEM: (4/11) Strong vs. Weak Equal to Zero (1/8)

Strong Equal to Zero (works as expected)

- assume  $g(x)$  scalar function in  $x$  with domain  $\Omega = (0, 1)$  (as  $u(x)$  and  $f(x)$  before)
- analytically:**

$g(x) = 0$  if and only if for all  $x \in \Omega$  we have that  $g(x) = 0$

$g(x) = 0$  in continuous strong form  $\Leftrightarrow \forall x \in \Omega : g(x) = 0$

- numerically:** we only have mesh points  $\Omega^h$  at disposal - enforce above in discrete nodes only

$g(x) = 0$  in discrete strong form  $\Leftrightarrow \forall x_i \in \Omega^h : g(x_i) = 0$

# 1D FEM: (4/11) Strong vs. Weak Equal to Zero (2/8)

## Discussion of Strong Equal to Zero

- **analytically:**  $g(x) = 0 \quad \forall x \in \Omega = (0, 1)$   
no space to choose  $g(x)$  for input  $x$
- **numerically:**  $g(x_i) = 0 \quad \forall x_i \in \Omega^h = \{x_j | 1 \leq j \leq N + 1\}$   
lots of space to choose  $g(x)$  for  $x \neq x_i$   
e.g.  $g(\frac{x_i+x_{i+1}}{2}) = \text{whatever-you-please}$

# 1D FEM: (4/11) Strong vs. Weak Equal to Zero (3/8)

## Example of Strong Equal to Zero

- choose:  $g(x) = r(x) = u''(x) + f(x)$  (residual function)
- analytically:

$$g(x) = 0 + \text{bc} \Leftrightarrow -u''(x) = f(x) + \text{bc} \Leftrightarrow \text{problem to solve}$$

- numerically:

$$g(x_i) = 0 + \text{bc} \Leftrightarrow -u''(x_i) = f(x_i) + \text{bc} \Leftrightarrow \text{finite diff. method}$$

# 1D FEM: (4/11) Strong vs. Weak Equal to Zero (4/8)

## Weak Equal to Zero (using averaging)

- first idea

$$g(x) = 0 \text{ in weak form} \Leftrightarrow \int_{\Omega} g(x) dx = \int_0^1 g(x) dx = 0$$

- strong equal zero implies weak equal zero

$$g(x) = 0 \Rightarrow \int_0^1 g(x) dx = 0$$

- converse is false (try e.g.  $g(x) = x - 1/2$ )
- hence the terminology

# 1D FEM: (4/11) Strong vs. Weak Equal to Zero (5/8)

Weak Equal to Zero (using weighted avering)

- towards a more proper definition

$$g(x) = 0 \text{ in weak form} \Leftrightarrow \int_{\Omega} g(x) v(x) dx = \int_0^1 g(x) v(x) dx = 0$$

- in previous slide  $v(x) = 1$  was chosen
- What is  $v(x)$ ? How many are there? How to choose them?

# 1D FEM: (4/11) Strong vs. Weak Equal to Zero (6/8)

Inspiration from Linear Algebra

- vector  $\mathbf{v} = \begin{pmatrix} v_1 \\ v_2 \\ \vdots \\ v_n \end{pmatrix}$

- unit-vector  $\mathbf{i}$   $i$ -th component equal to one, all other components equal to zero
- $\mathbf{v} = \mathbf{0} \Leftrightarrow v_i = 0 \quad 1 \leq i \leq n$   
again  $N + 1$  equations indexed by  $i$
- $\mathbf{v} = \mathbf{0} \Leftrightarrow \mathbf{v}^T \mathbf{i} = \langle \mathbf{v}, \mathbf{i} \rangle = 0 \quad 1 \leq i \leq n$
- how to bring this to finite elements?

# 1D FEM: (4/11) Strong vs. Weak Equal to Zero (7/8)

## Applied to Finite Elements

- $N + 1$  basis functions  $\phi_i(x)$  defined by the mesh  $\Omega^h$
- inner product:  $\langle g(x), \phi_i(x) \rangle = \int_0^1 g(x) \phi_i(x) dx$
- $\langle g(x), \phi_i(x) \rangle$  coordinate of  $g(x)$  along the basis function  $\phi_i(x)$
- **numerically**: essential in remainder of the course

$$g(x) = 0 \text{ in discrete weak form} \Leftrightarrow \forall x_i \in \Omega^h : \langle g(x), \phi_i(x) \rangle = 0$$

again  $N + 1$  equations indexed by  $i$

- **analytically**: see later

# 1D FEM: (4/11) Strong vs. Weak Equal to Zero (8/8)

## Applied to Finite Elements

- choose  $g(x) = r(x) = u''(x) + f(x)$
- enforce  $g(x) = 0$  plus boundary conditions in discrete weak form

$$\langle g(x), \phi_i(x) \rangle = 0 \quad \text{for } 1 \leq i \leq N+1$$

$$\Leftrightarrow \int_0^1 g(x) \phi_i(x) dx = 0 \quad \text{for } 1 \leq i \leq N+1$$

$$\Leftrightarrow \int_0^1 -u''(x) \phi_i(x) dx = \int_0^1 f(x) \phi_i(x) dx \quad \text{for } 1 \leq i \leq N+1$$

# 1D FEM: (5/11) Computation of the RHS Vector (1/4)

## Terminology

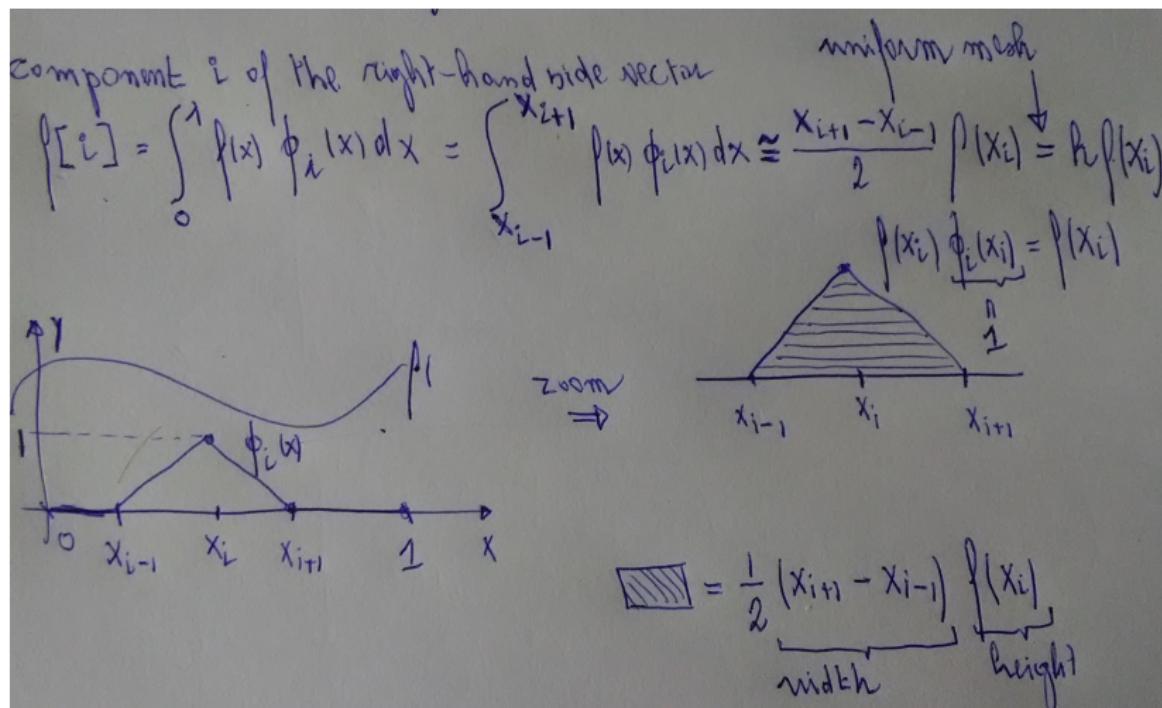
- $f(x)$ : source term or right-hand side in the differential equation
- $\alpha$ : value imposed on the derivative
- $\phi_i(x)$ : basis function centered on  $x_i \in \Omega^h$
- $f_i$ :  $i$ -th component of the right-hand side vector  $\mathbf{f}$
- relationship (see previous slide)

$$f_i = \int_0^1 f(x) \phi_i(x) dx$$

## 1D FEM: (5/11) Computation of the RHS Vector (2/4)

$$\begin{aligned}f_i &= \int_0^1 f(x) \phi_i(x) dx = \int_{x_{i-1}}^{x_{i+1}} f(x) \phi_i(x) dx \quad (\text{restrict to } e_{i-1} \cup e_i) \\&= \int_{x_{i-1}}^{x_i} f(x) \phi_i(x) dx + \int_{x_i}^{x_{i+1}} f(x) \phi_i(x) dx \quad (\text{separate } e_{i-1} \text{ and } e_i) \\&\approx \frac{h_{i-1}}{2} [f(x_{i-1}) \phi_i(x_{i-1}) + f(x_i) \phi_i(x_i)] + \frac{h_i}{2} [f(x_i) \phi_i(x_i) + f(x_{i+1}) \phi_i(x_{i+1})] \\&\quad (\text{use trapezoidal rule twice}) \\&= \frac{h_{i-1}}{2} [f(x_{i-1}) 0 + f(x_i) 1] + \frac{h_i}{2} [f(x_i) 1 + f(x_{i+1}) 0] \\&\quad (\text{use properties of } \phi_i(x)) \\&= \frac{h_{i-1} + h_i}{2} f(x_i) \\&= h f(x_i) \quad (\text{in case of uniform mesh})\end{aligned}$$

# 1D FEM: (5/11) Computation of the RHS Vector (3/4)



# 1D FEM: (5/11) Computation of the RHS Vector (4/4)

Treatment of the Dirichlet and Neumann boundary conditions

- **Dirichlet** boundary condition

fix  $u(x)$  - grounding:

same as for Finite Difference method

- **Neumann** boundary condition

fix  $\frac{d}{dx} u(x)$  - imposed symmetry or flux:

different from Finite Difference method - details later

# 1D FEM: (6/11) Mathematical Preliminaries (1/4)

Wish - Desire - Aim - Ambition?

- manipulate –  $\int_0^1 u''(x) \phi_i(x) dx$
- challenge:  $u(x)$  is unknown
- observe:  $u''(x)$  has double prime -  $\phi_i(x)$  has no prime

# 1D FEM: (6/11) Mathematical Preliminaries (2/4)

Derivative: Integration by Parts in one variable (1/2)

- fundamental theorem of calculus:

assume  $0 < x < 1$  or  $x \in \Omega = (0, 1)$  and  $F'(x) = \frac{dF(x)}{dx}$

$$\int_{\Omega} F'(x) dx = \int_0^1 F'(x) dx = [F(x)]_0^1 = F(1) - F(0)$$

- choose  $F(x) = \frac{du}{dx} v(x)$  (motivation on next slide)

observe  $u(x)$  has prime and  $v(x)$  has no prime

- then  $F'(x) = \frac{d}{dx} \left( \frac{du}{dx} v(x) \right) = u''(x) v(x) + u'(x) v'(x)$

# 1D FEM: (6/11) Mathematical Preliminaries (3/4)

## Derivative: Integration by Parts in one variable (2/2)

- from  $\int_0^1 F'(x) dx = [F(x)]_0^1$  we obtain that

$$\int_0^1 [u''(x) v(x) + u'(x) v'(x)] dx = [u'(x) v(x)]_0^1$$

- after rearranging terms

$$-\int_0^1 u''(x) v(x) dx = \int_0^1 u'(x) v'(x) dx - [u'(x) v(x)]_0^1$$

LHS:  $u(x)$  has double prime and  $v(x)$  has no prime

RHS: both  $u(x)$  and  $v(x)$  have one prime - additional fudge term  
(used later)

# 1D FEM: (6/11) Mathematical Preliminaries (4/4)

## Integration: Quadrature by Trapezoidal Rule

- trapezoidal rule:  $0 \leq a < b \leq 1$

$$\int_a^b g(x) dx \approx \frac{b-a}{2} [g(a) + g(b)]$$

- Simpson rule:  $0 \leq a < b \leq 1$

$$\int_a^b g(x) dx \approx \frac{b-a}{6} \left[ g(a) + 4 g\left(\frac{a+b}{2}\right) + g(b) \right]$$

- typically applied to  $a = x_i$  and  $b = x_{i+1}$

# 1D FEM: (7/11) Discrete Weak Form (1/4)

Apply Integration by Parts on Weak Formulation

- earlier we set the residual  $r(x) = u''(x) + f(x)$  to zero in weak form and arrived at

$$\int_0^1 -u''(x) \phi_i(x) dx = \int_0^1 f(x) \phi_i(x) dx \quad \text{for } 1 \leq i \leq N+1$$

- apply integration by part to the LHS

$$\int_0^1 u'(x) \phi'_i(x) dx = \int_0^1 f(x) \phi_i(x) dx + [u'(x) \phi_i(x)]_0^1 \quad \text{for } 1 \leq i \leq N+1$$

observe: minus sign in LHS disappeared

- first order derivative on both  $u(x)$  and  $\phi_i(x)$

## 1D FEM: (7/11) Discrete Weak Form (2/4)

- boundary conditions:  $u(x = 0) = 0$  and  $u'(x = 1) = \alpha$
- boundary term in the RHS of the weak form

$$\begin{aligned}[u'(x) \phi_i(x)]_0^1 &= u'(x = 1) \phi_i(x = 1) - u'(x = 0) \phi_i(x = 0) \\ &= \alpha \phi_i(x = 1) - (\text{no information given}) \phi_i(x = 0)\end{aligned}$$

- Dirichlet and Neumann boundary conditions treated differently
- for now save assume that

$$[u'(x) \phi_i(x)]_0^1 = \alpha \phi_i(x = 1) - 0 = \alpha \phi_i(x = 1)$$

# 1D FEM: (7/11) Discrete Weak Form (3/4)

Discrete Weak Form Becomes

$$\int_0^1 u'(x) \phi_i'(x) dx = \int_0^1 f(x) \phi_i(x) dx + \alpha \phi_i(x=1) \quad \text{for } 1 \leq i \leq N+1$$

- assume  $u(x)$  approximate by  $u^h(x)$  where  $u^h(x) = \sum_{j=1}^{N+1} c_j \phi_j(x)$
- thus  $u'(x)$  approximate by  $[u^h(x)]' = \sum_{j=1}^{N+1} c_j \phi'_j(x)$
- then for  $1 \leq i \leq N+1$

$$\sum_{j=1}^{N+1} \int_0^1 \phi'_j(x) \phi'_i(x) dx c_j = \int_0^1 f(x) \phi_i(x) dx + \alpha \phi_i(x=1)$$

# 1D FEM: (7/11) Discrete Weak Form (4/4)

Discrete Weak Form Becomes

- for  $1 \leq i \leq N + 1$

$$\sum_{j=1}^{N+1} \int_0^1 \phi'_i(x) \phi'_j(x) dx c_j = \int_0^1 f(x) \phi_i(x) dx + \alpha \phi_i(x=1)$$

- can be written in the form: for  $1 \leq i \leq N + 1$

$$\sum_{j=1}^{N+1} A_{ij} c_j = f_i$$

- and thus as a  $N + 1$  by  $N + 1$  linear system

$$A \mathbf{c} = \mathbf{f}$$

- index  $i$  counts equations - index  $j$  counts unknowns

# 1D FEM: (8/11) Linear System Formulation (1/8)

## Expression for Matrix and Vector Elements

- Matrix elements:

$$A_{ij} = \int_0^1 \phi'_i(x) \phi'_j(x) dx \text{ for } 1 \leq i, j \leq N + 1$$

- Vector elements:

$$f_i = \int_0^1 f(x) \phi_i(x) dx + \alpha \phi_i(x = 1) \text{ for } 1 \leq i \leq N + 1$$

# 1D FEM: (8/11) Linear System Formulation (2/8)

## Properties of Matrix $A$

- diagonal of  $A$  is **positive** -  $A_{ii} > 0$

$$A_{ii} = \int_0^1 \underbrace{\phi'_i(x) \phi'_i(x)}_{>0} dx > 0$$

- $A$  is **symmetric** -  $A = A^T$

$$A_{ij} = \int_0^1 \phi'_i(x) \phi'_j(x) dx = \int_0^1 \phi'_j(x) \phi'_i(x) dx = A_{ji}$$

- off-diagonal** of  $A$  is negative

$$A_{i,i+1} = \int_0^1 \phi'_i(x) \phi'_{i+1}(x) dx = \int_{x_i}^{x_{i+1}} \underbrace{\phi'_i(x)}_{<0} \underbrace{\phi'_{i+1}(x)}_{>0} dx < 0$$

- $A$  is **tri-diagonal** and thus sparse
- $A$  many other cool properties  $\Rightarrow$  **fast solvers** for  $A\mathbf{c} = \mathbf{f}$  exist

# 1D FEM: (8/11) Linear System Formulation (3/8)

Remember that

- length of element  $e_i = [x_i, x_{i+1}]$  is denoted by  $h_i = x_{i+1} - x_i$
- length of element  $e_{i-1} = [x_{i-1}, x_i]$  is denoted by  $h_{i-1} = x_i - x_{i-1}$
- slope of  $\phi_i(x)$  on  $e_i$  is equal to  $\frac{-1}{h_i}$
- slope of  $\phi_{i+1}(x)$  on  $e_i$  is equal to  $\frac{1}{h_i}$

# 1D FEM: (8/11) Linear System Formulation (4/8)

## Diagonal Matrix Elements $A_{ii}$

$$\begin{aligned} A_{ii} &= \int_0^1 \phi'_i(x) \phi'_i(x) dx = \int_{x_{i-1}}^{x_{i+1}} \phi'_i(x) \phi'_i(x) dx \quad (\text{restrict to } e_{i-1} \cup e_i) \\ &= \int_{x_{i-1}}^{x_i} \phi'_i(x) \phi'_i(x) dx + \int_{x_i}^{x_{i+1}} \phi'_i(x) \phi'_i(x) dx \quad (\text{separate } e_{i-1} \text{ and } e_i) \\ &= \int_{x_{i-1}}^{x_i} \left(\frac{1}{h_{i-1}}\right)^2 dx + \int_{x_i}^{x_{i+1}} \left(\frac{-1}{h_i}\right)^2 dx \quad (\text{use information on derivatives}) \\ &= h_{i-1} \left(\frac{1}{h_{i-1}}\right)^2 + h_i \left(\frac{1}{h_i}\right)^2 \quad (\text{integration of constant integrands}) \\ &= \frac{1}{h_{i-1}} + \frac{1}{h_i} \quad (\text{in case of non-uniform mesh}) \\ &= \frac{2}{h} \quad (\text{in case of uniform mesh for which all } h_i \text{ are equal}) \end{aligned}$$

# 1D FEM: (8/11) Linear System Formulation (5/8)

## Upper-Diagonal Matrix Elements

$$\begin{aligned} A_{i,i+1} &= \int_0^1 \phi'_i(x) \phi'_{i+1}(x) dx = \int_{x_i}^{x_{i+1}} \phi'_i(x) \phi'_{i+1}(x) dx \quad (\text{restrict to } e_i) \\ &= \int_{x_i}^{x_{i+1}} \frac{-1}{h_i} \frac{1}{h_i} dx \quad (\text{use information on derivatives}) \\ &= h_i \frac{-1}{h_i} \frac{1}{h_i} \quad (\text{integration of constant integrands}) \\ &= \frac{-1}{h_i} \quad (\text{in case of non-uniform mesh}) \\ &= \frac{-1}{h} \quad (\text{in case of uniform mesh for which all } h_i \text{ are equal}) \end{aligned}$$

# 1D FEM: (8/11) Linear System Formulation (6/8)

Summary: Matrix  $A$  and Right-Hand Vector  $\mathbf{f}$

Internal mesh points  $x_i$  for  $2 \leq i \leq N$

- Matrix

$$A[i, :] = \begin{pmatrix} 0 & \dots & 0 & \underbrace{-1/h_{i-1}}_{i-1} & \underbrace{1/h_{i-1} + 1/h_i}_i & \underbrace{-1/h_i}_{i+1} & 0 & \dots & 0 \end{pmatrix}$$

- Vector

$$f[i] = \frac{h_{i-1} + h_i}{2} f(x_i)$$

# 1D FEM: (8/11) Linear System Formulation (7/8)

Summary: Matrix  $A$  and Right-Hand Vector  $\mathbf{f}$

Treatment of the Boundary Conditions

- Dirichlet boundary conditions:  $u(x = 0) = 0$ :

modify first equation in linear system

see finite difference method

- Neumann boundary conditions:  $u'(x = 1) = \alpha$

add term  $\alpha \phi_i(x = 1)$  to vector  $\mathbf{f}$

see lab sessions for details

## Exercise

- assume again a non-uniform mesh on  $\Omega$  with 4 nodes  $x = 0$ ,  $x = 1/3$ ,  $x = 1/2$  and  $x = 1$
- compute matrix  $A$  and vector  $f$

# 1D FEM: (9/11) Element-by-Element Construction of the Vector (1/7)

Remember that

- $i$ -th component of the right-hand side vector  $\mathbf{f}$  is given by

$$f_i = \int_{x_{i-1}}^{x_{i+1}} f(x) \phi_i(x) dx = \int_{e_{i-1}}^{e_i} f(x) \phi_i(x) dx + \int_{e_i}^{x_{i+1}} f(x) \phi_i(x) dx$$

# 1D FEM: (9/11) Element-by-Element Construction of the Vector (2/7)

Element  $e_i$  contributes to  $f_i$  and  $f_{i+1}$  only

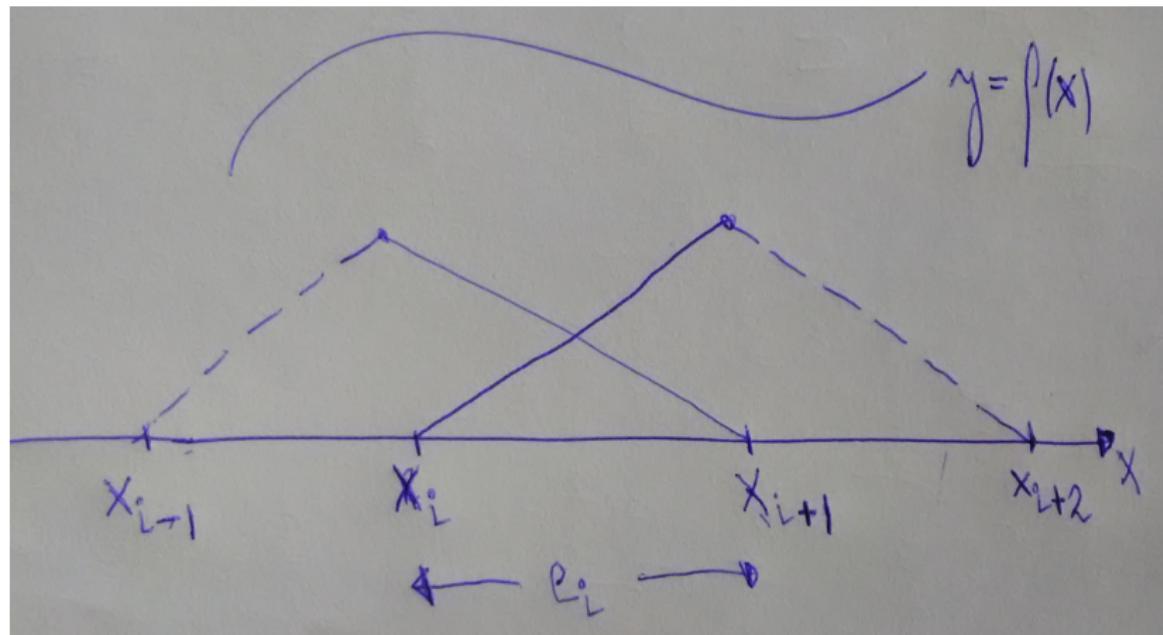
- $i$ -th component of the right-hand side vector  $\mathbf{f}$

$$f_i = \int_{e_{i-1}} f(x) \phi_i(x) dx + \int_{e_i} f(x) \phi_i(x) dx$$

- $i+1$ -th component of the right-hand side vector  $\mathbf{f}$

$$f_{i+1} = \int_{e_i} f(x) \phi_{i+1}(x) dx + \int_{e_{i+1}} f(x) \phi_{i+1}(x) dx$$

# 1D FEM: (9/11) Element-by-Element Construction of the Vector (3/7)



# 1D FEM: (9/11) Element-by-Element Construction of the Vector (4/7)

How does element  $e_i$  contribute to the vector  $\mathbf{f}$ ?

- we here fix element  $e_i$  (rather than fixing component  $f_i$ )
- element  $e_i = [x_i, x_{i+1}]$  sees shape functions  $\phi_i(x)$  and  $\phi_{i+1}(x)$   
all other basis functions are zero on the element  $e_i$
- the element  $e_i$  therefore contributes to **two** component of  $\mathbf{f}$   
these components are  $f_i$  and  $f_{i+1}$   
(same indices as basis functions)
- this leads to an element-by-element construction of vector  $\mathbf{f}$

# 1D FEM: (9/11) Element-by-Element Construction of the Vector ( $5/7$ )

How does element  $e_i$  contribute to the vector  $\mathbf{f}$ ?

- $f_{e_i} \in \mathbb{R}^2$  contribution of element  $e_i$  to global vector  $\mathbf{f}$

$$f_{e_i} = \begin{pmatrix} \int_{e_i} f(x) \phi_i(x) dx \\ \int_{e_i} f(x) \phi_{i+1}(x) dx \end{pmatrix}$$

- use trapezoidal rule of integration (see earlier)

$$f_{e_i} = \begin{pmatrix} \int_{e_i} f(x) \phi_i(x) dx \\ \int_{e_i} f(x) \phi_{i+1}(x) dx \end{pmatrix} = \frac{h_i}{2} \begin{pmatrix} f(x_i) \\ f(x_{i+1}) \end{pmatrix}$$

- given mesh  $\Omega^h$  and source  $f(x)$ ,  $f_{e_i}$  for each  $e_i$  can be computed

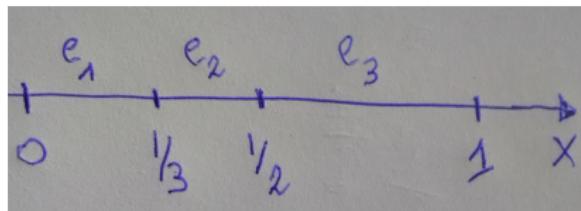
# 1D FEM: (9/11) Element-by-Element Construction of the Vector (6/7)

## Finite Element Assembly of the Vector $\mathbf{f}$

- loop over all of the  $N$  elements  $e_i$  in the mesh  $\Omega^h$
- on  $e_i$  compute the local element vector  $f_{e_i} \in \mathbb{R}^2$   
the local element vector has **two** components
- add local element vector to the global vector  $\mathbf{f} \in \mathbb{R}^{N+1}$   
the global vector  $\mathbf{f}$  has  **$N + 1$**  components
- $\mathbf{f} = \mathbf{f} + f_{e_i}$   
assembly requires taking the mesh connectivity into account  
connectivity here between nodes  $x_i$  and the elements  $e_i$  into account

# 1D FEM: (9/11) Element-by-Element Construction of the Vector (7/7)

Assembly for Three-Element Mesh



$$f = \begin{pmatrix} e_1 \\ e_2 \\ e_3 \end{pmatrix}$$

A hand-drawn diagram of a 3x3 matrix. The matrix has non-zero entries in the main diagonal and the super-diagonal. The matrix is:

$$\begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix}$$

# 1D FEM: (10/11) Element-by-Element Construction of the Matrix (1/4)

How does element  $e_i$  contribute to the vector  $A$ ?

- element  $e_i = [x_i, x_{i+1}]$  sees shape functions  $\phi_i(x)$  and  $\phi_{i+1}(x)$
- the element  $e_i$  therefore contributes to **four** component of  $A$   
these components are  $A_{ii}$ ,  $A_{i,i+1} = A_{i+1,i}$  and  $A_{i+1,i+1}$

# 1D FEM: (10/11) Element-by-Element Construction of the Matrix (2/4)

How does element  $e_i$  contribute to the vector  $A$ ?

- $A_{e_i} \in \mathbb{R}^{2 \times 2}$  contribution of element  $e_i$  to global vector  $A$

$$A_{e_i} = \begin{pmatrix} \int_{e_i} \phi'_i(x) \phi'_i(x) dx & \int_{e_i} \phi'_i(x) \phi'_{i+1}(x) dx \\ \int_{e_i} \phi'_i(x) \phi_{i+1}(x) dx & \int_{e_i} \phi'_{i+1}(x) \phi'_{i+1}(x) dx \end{pmatrix}$$

- use derivative of the shape functions

$$A_{e_i} = \frac{1}{h_i} \begin{pmatrix} 1 & -1 \\ -1 & 1 \end{pmatrix}$$

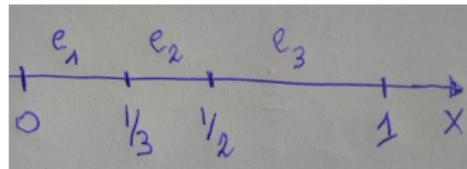
# 1D FEM: (9/11) Element-by-Element Construction of the Matrix (3/4)

## Finite Element Assembly of the Matrix $A$

- loop over all of the  $N$  elements  $e_i$  in the mesh  $\Omega^h$
- on  $e_i$  compute the local element matrix  $A_{e_i} \in \mathbb{R}^{2 \times 2}$   
the local element matrix has **two by two** components
- add local element matrix to the global matrix  $A \in \mathbb{R}^{(N+1) \times (N+1)}$   
the global matrix  $A$  has  **$N + 1$  by  $N + 1$**  components
- $A = A + A_{e_i}$   
assembly requires taking the mesh connectivity into account  
connectivity here between nodes  $x_i$  and the elements  $e_i$  into account

## 1D FEM: (10/11) Element-by-Element Construction of the Matrix (4/4)

## Assembly for Three-Element Mesh



$$A = \left( \begin{array}{ccc} e_1 & & \\ & e_2 & \\ & & e_3 \end{array} \right)$$

# 1D FEM: (11/11) Implementation in Julia

- see Lab Session