## 1. Diagram notation

**Notation 1.1.** *Diagram notation.* In diagram notation, particle-hole operators are written as oriented lines extending from a vertex. Particle annihilation operators enter the vertex from below, particle creation operators leave the vertex at the top, and single-excitation operators have both creation and annihilation lines. Contractions are represented by joining particle-hole lines with compatible position and orientation.

$$\stackrel{\downarrow}{\uparrow} \equiv a_p$$
  $\stackrel{\downarrow}{\downarrow} \equiv a_p^{\dagger}$   $\stackrel{\downarrow}{\downarrow} \equiv a_p^{\dagger} a_q = a_q^p$   $\stackrel{\downarrow}{\downarrow} \equiv \overline{a_p} \overline{a_q} = a_q^{p^{ullet}}$ 

Quasiparticle operators with respect to  $\Phi$  are distinguished by the use of closed-circle vertices, with particle lines pointing upward and with hole lines pointing downward. Single-excitation operators split into four cases (vv, vo, ov, and oo) representing the virtual and occupied blocks of  $a_q^p$ . Internal contractions of single-excitation operators (bubble contractions) are implicitly taken to be hole contractions with respect to  $\Phi$ .

Higher excitation operators are depicted by joining single-excitation operators with a solid line. Contracted operators are implicitly normal ordered together. Normal-ordered products of uncontracted operators are joined with a dotted line.

$$\underbrace{\underbrace{-\cdots }_{m \text{ times}}}_{\text{m times}} \equiv :a_{q_1}^{p_1} \cdots a_{q_m}^{p_m} := a_{q_1 \cdots q_m}^{p_1 \cdots p_m} \underbrace{-\cdots }_{n \text{ times}} \equiv :a_{q_1 \cdots q_m}^{p_1 \cdots p_m} a_{s_1 \cdots s_n}^{r_1 \cdots r_n} :\underbrace{-\cdots }_{m \text{ times}} \equiv :a_{q_1 \cdots q_m}^{p_1 \cdots p_m} a_{s_1 \cdots s_n}^{r_1 \cdots r_n} :\underbrace{-\cdots }_{n \text{ times}} \equiv :a_{q_1 \cdots q_m}^{p_1 \cdots p_m} a_{s_1 \cdots s_n}^{r_1 \cdots r_n} :\underbrace{-\cdots }_{n \text{ times}} = :a_{q_1 \cdots q_m}^{p_1 \cdots p_m} a_{s_1 \cdots s_n}^{r_1 \cdots r_n} :\underbrace{-\cdots }_{n \text{ times}} = :a_{q_1 \cdots q_m}^{p_1 \cdots p_m} a_{s_1 \cdots s_n}^{r_1 \cdots r_n} :\underbrace{-\cdots }_{n \text{ times}} = :a_{q_1 \cdots q_m}^{p_1 \cdots p_m} a_{s_1 \cdots s_n}^{r_1 \cdots r_n} :\underbrace{-\cdots }_{n \text{ times}} = :a_{q_1 \cdots q_m}^{p_1 \cdots p_m} a_{s_1 \cdots s_n}^{r_1 \cdots r_n} :\underbrace{-\cdots }_{n \text{ times}} = :a_{q_1 \cdots q_m}^{p_1 \cdots p_m} a_{s_1 \cdots s_n}^{r_1 \cdots r_n} :\underbrace{-\cdots }_{n \text{ times}} = :a_{q_1 \cdots q_m}^{p_1 \cdots p_m} a_{s_1 \cdots s_n}^{r_1 \cdots r_n} :\underbrace{-\cdots }_{n \text{ times}} = :a_{q_1 \cdots q_m}^{p_1 \cdots p_m} a_{s_1 \cdots s_n}^{r_1 \cdots r_n} :\underbrace{-\cdots }_{n \text{ times}} = :a_{q_1 \cdots q_m}^{p_1 \cdots p_m} a_{s_1 \cdots s_n}^{r_1 \cdots r_n} :\underbrace{-\cdots }_{n \text{ times}} = :a_{q_1 \cdots q_m}^{p_1 \cdots p_m} a_{s_1 \cdots s_n}^{r_1 \cdots r_n} :\underbrace{-\cdots }_{n \text{ times}} = :a_{q_1 \cdots q_m}^{p_1 \cdots p_m} a_{s_1 \cdots s_n}^{r_1 \cdots r_n} :\underbrace{-\cdots }_{n \text{ times}} = :a_{q_1 \cdots q_m}^{p_1 \cdots p_m} a_{s_1 \cdots s_n}^{r_1 \cdots r_n} :\underbrace{-\cdots }_{n \text{ times}} = :a_{q_1 \cdots q_m}^{p_1 \cdots p_m} a_{s_1 \cdots s_n}^{r_1 \cdots r_n} :\underbrace{-\cdots }_{n \text{ times}} = :a_{q_1 \cdots q_m}^{p_1 \cdots p_m} a_{s_1 \cdots s_n}^{r_1 \cdots r_n} :\underbrace{-\cdots }_{n \text{ times}} = :a_{q_1 \cdots q_m}^{p_1 \cdots p_m} a_{s_1 \cdots s_n}^{r_1 \cdots r_n} :\underbrace{-\cdots }_{n \text{ times}} = :a_{q_1 \cdots q_m}^{p_1 \cdots p_m} a_{s_1 \cdots s_n}^{r_1 \cdots r_n} :\underbrace{-\cdots }_{n \text{ times}} = :a_{q_1 \cdots q_m}^{p_1 \cdots p_m} a_{s_1 \cdots s_n}^{r_1 \cdots r_n} :\underbrace{-\cdots }_{n \text{ times}} = :a_{q_1 \cdots q_n}^{p_1 \cdots p_m} a_{s_1 \cdots s_n}^{r_1 \cdots r_n} :\underbrace{-\cdots }_{n \text{ times}} = :a_{q_1 \cdots q_n}^{p_1 \cdots p_n} a_{s_1 \cdots s_n}^{r_1 \cdots r_n} :\underbrace{-\cdots }_{n \text{ times}} = :a_{q_1 \cdots q_n}^{p_1 \cdots p_n} a_{s_1 \cdots s_n}^{r_1 \cdots r_n} :\underbrace{-\cdots }_{n \text{ times}} = :a_{q_1 \cdots q_n}^{p_1 \cdots p_n} a_{s_1 \cdots s_n}^{r_1 \cdots r_n} :\underbrace{-\cdots }_{n \text{ times}} = :a_{q_1 \cdots q_n}^{p_1 \cdots p_n} a_{s_1 \cdots s_n}^{r_1 \cdots r_n} :\underbrace{-\cdots }_{n \text{ times}} = :a_{q_1 \cdots q_n}^{p_1 \cdots p_n} a_{s_1 \cdots s_n}^{r_1 \cdots r_n} a_{s_1 \cdots s_n}^{r_1 \cdots$$

 $\Phi$ -normal-ordering is indicated by the use of double-circle vertices,  $\otimes$  and  $\otimes$  instead of  $\circ$  and  $\bullet$ .

**Definition 1.1.** m-electron operators in Diagram notation. The primary building blocks of a graph are m-electron operators, which can be represented in two equivalent ways. The Goldstone representation depicts an operator as a label attached to the corresponding excitation operator, whereas the Hugenholtz representation depicts the operator as a single vertex with m outgoing and incoming lines. Note that  $\left(\frac{1}{m!}\right)^2 \sum_{\text{Einstein}}$  is baked into the definition (see def 1.2 and ax 1.1).

$$\begin{array}{c}
\boxed{v} & \stackrel{p_1}{\longrightarrow} \cdots \stackrel{p_m}{\longrightarrow} = \overline{v}_{p_1 \cdots p_m}^{q_1 \cdots q_m} a_{q_1 \cdots q_m}^{p_1 \cdots p_m} a_{q_1 \cdots q_m}^{p_1 \cdots p_m} = \overbrace{v}^{p_1 \cdots p_m} a_{q_1 \cdots q_m}^{p_1 \cdots p_m} a_{q_1 \cdots q_m}^{p_1 \cdots p_m} a_{q_1 \cdots q_m}^{p_1 \cdots p_m} = \overline{v}_{p_1 \cdots p_m}^{q_2 \cdots q_m} a_{q_1 \cdots q_m}^{p_2 \cdots p_m} = \overline{v}_{p_{\pi(1)} \cdots p_{\pi(m)}}^{p_2 \cdots p_m} a_{q_{\sigma(1)} \cdots q_{\sigma(m)}}^{p_{\pi(1)} \cdots p_{\pi(m)}} a_{q_{\sigma(1)} \cdots q_{\sigma(m)}}^{p_{\sigma(1)} \cdots p_{\sigma(m)}} a_{q_{\sigma(1)} \cdots q_{\sigma(m)}}^{p_{\sigma(1)} \cdots p_{\sigma(m)}}^{p_{\sigma(1)} \cdots p_{\sigma(m)}}^{p_{\sigma(1)} \cdots p_{\sigma(m)}}^{p_{\sigma(1)} \cdots p_{\sigma(m)}} a_{q_{\sigma(1)} \cdots q_{\sigma(m)}}^{p_{\sigma(1)} \cdots p_{\sigma(m)}}^{p_{\sigma(1)} \cdots p_{\sigma(m)}}^{p_{\sigma(1)} \cdots p_{\sigma(m)}}^{p_{\sigma(1)} \cdots$$

The labeled diagrams on the right represent just the summand of the operator, which highlights the difference between representations. Both summands correspond to an excitation operator weighted by its antisymmetrized interaction tensor, but whereas the Goldstone summand specifies an ordering for the indices of its corresponding algebraic term, the Hugenholtz summand does not. Since the phases of  $\overline{v}_{p_1...p_m}^{q_1...q_m}$  and  $a_{q_1...q_m}^{p_1...p_m}$  cancel under index permutation, the two labeled diagrams are actually equal – a Hugenholtz summand can be expanded into a Goldstone summand by simply choosing an arbitrary ordering for the indices. In practice, the symmetry of the Hugenholtz operator simplifies the enumeration of Wick expansions whereas the Goldstone operator makes it easier to evaluate a graph's overall phase.

<sup>&</sup>lt;sup>1</sup>In the original paper [J. Goldstone, *P. Roy. Soc. A* **239**, (1957)], Goldstone's diagrams were actually defined in terms of non-antisymmetrized integrals. The *antisymmetrized Goldstone diagrams* used here are sometimes called *Brandow diagrams*.

**Definition 1.2.** Graph. A graph<sup>2</sup> is a 4-tuple G = (O, L, h, t) where O is a set of electron operators, L is a set of lines, and h and t are maps from L to O that return the head  $h(l) \in O$  and tail end  $t(l) \in O$  of every line  $l \in L$ . Here, we allow for external lines in which either h(l) or t(l) equals e, the free end, which is formally considered a member of O. Lines no free end are termed internal. If l and l' share the same tail and the same head, they are termed equivalent lines and we write  $l \sim l'$ . If o and o' can be exchanged, attaching the ends of one to the other and vice versa, without altering G then they are termed interchangeable operators. Repeated copies of the same operator are formally distinguished as elements of O and termed identical operators, denoted  $o \sim o'$ . If identical operators are interchangeable, they are termed equivalent, denoted  $o \sim o'$ . See ex A.2. The rules of interpretation for translating G into an algebraic expression are given in ax 1.1.

**Definition 1.3.** Summand graph. A summand graph<sup>3</sup>  $\Sigma(G)$  of G is a 3-tuple  $\Sigma(G) = (G, S, s)$  where S is a set of symbols and  $s: L \to S$  is a label map assigning one symbol  $s(l) \in S$  to each line  $l \in L$ . Pictorially, this corresponds to labeling each line in G with an index, p, q, r, s, etc.  $\Sigma(G)$  translates directly into an algebraic summand according to def 1.1 and nt 1.1, with top-to-bottom ordering in the graph corresponding to left-to-right ordering in the expression.

**Definition 1.4.** Degeneracy. The line degeneracy or simply degeneracy of G is the number of permutational symmetries in  $\Sigma(G)$ , a positive integer denoted dg(G). Formally, if  $S = \{s_1, \ldots, s_n\}$  with  $s(l_i) = s_i$ , then dg(G) is the number of label permutations  $\Sigma_{\pi}(G) = (G, S, s_{\pi})$ , where  $\pi \in S_n$  and  $s_{\pi}(l_i) = s_{\pi(i)}$ , that don't change the summand graph,  $\Sigma_{\pi}(G) = \Sigma(G)$ . If G has no identical operators, then  $dg(G) = |L_1|! \cdots |L_m|!$  where  $L = L_1 \cup \cdots \cup L_m$  partitions L into equivalent lines and  $|L_i|$  denotes the number of elements in the set  $L_i$ .

**Axiom 1.1.** Rules of interpretation. The algebraic interpretation of G is obtained from  $\Sigma(G)$  as follows. See ex A.1.

- 1. Multiply  $\Sigma(G)$  by  $dg(G)^{-1}$ , termed the degeneracy factor.
- 2. Sum each index in  $\Sigma(G)$  over its range.

**Example 1.1.** The one- and two-electron components of  $H_e$  expand into occupied/virtual blocks as follows.

$$\begin{split} h_p^q a_q^p &= h_a^b a_b^a + h_a^i a_i^a + h_i^a a_a^i + h_i^j a_j^i \\ \frac{1}{4} \overline{g}_{pq}^{rs} a_{rs}^{pq} &= \frac{1}{4} \overline{g}_{ab}^{cd} a_{cd}^{ab} + \frac{1}{2} \overline{g}_{ab}^{ci} a_{cd}^{ab} + \frac{1}{2} \overline{g}_{ab}^{ci} a_{bc}^{ai} + \frac{1}{4} \overline{g}_{ab}^{ij} a_{ij}^{ab} + \overline{g}_{ij}^{bj} a_{ab}^{ij} + \frac{1}{4} \overline{g}_{ij}^{ab} a_{ab}^{ij} + \frac{1}{2} \overline{g}_{ia}^{jk} a_{jk}^{ia} + \frac{1}{2} \overline{g}_{ij}^{ka} a_{ij}^{ia} + \frac{1}{4} \overline{g}_{ij}^{kl} a_{kl}^{ij} \\ \end{split}$$

Defining  $\boxtimes \xrightarrow{\uparrow} \equiv h_p^q a_q^p$  and  $\updownarrow \longrightarrow \updownarrow \equiv \frac{1}{4} \overline{g}_{pq}^{rs} a_{rs}^{pq}$ , these equations are expressed in terms of Goldstone diagrams as follows.

The degeneracy factors fall into three cases:  $\{l_1, l_2\} \cup \{l_3, l_4\} \implies \operatorname{dg}(G)^{-1} = \frac{1}{2 \cdot 2}; \{l_1, l_2\} \cup \{l_3\} \cup \{l_4\} \implies \operatorname{dg}(G)^{-1} = \frac{1}{2 \cdot 1 \cdot 1};$  and  $\{l_1\} \cup \{l_2\} \cup \{l_3\} \cup \{l_4\} \implies \operatorname{dg}(G)^{-1} = \frac{1}{1 \cdot 1 \cdot 1 \cdot 1}.$  In terms of Hugenholtz diagrams, these are written as follows.

Example 1.2. The  $\Phi$ -normal Wick expansion of the one- and two-electron components of  $H_e$  are as follows.

$$\begin{split} h_{p}^{q} a_{q}^{p} &= h_{p}^{q} \left( \tilde{a}_{q}^{p} + \tilde{a}_{q^{\circ}}^{p^{\circ}} \right) = h_{p}^{q} \tilde{a}_{q}^{p} + h_{p}^{q} \gamma_{q}^{p} \\ \frac{1}{4} \overline{g}_{pq}^{rs} a_{rs}^{pq} &= \frac{1}{4} \overline{g}_{pq}^{rs} \left( \tilde{a}_{rs}^{pq} + \hat{P}_{(r/s)}^{(p/q)} \tilde{a}_{r^{\circ}s}^{p^{\circ}q} + \hat{P}_{(r/s)} \tilde{a}_{r^{\circ}s^{\circ\circ}}^{p^{\circ}q^{\circ\circ}} \right) = \frac{1}{4} \overline{g}_{pq}^{rs} \tilde{a}_{rs}^{pq} + \overline{g}_{pq}^{rs} \gamma_{r}^{p} \tilde{a}_{s}^{q} + \frac{1}{2} \overline{g}_{pq}^{rs} \gamma_{r}^{p} \gamma_{s}^{q} \end{split}$$

In terms of Goldstone diagrams, these equations are written as

and, in terms of Hugenholtz diagrams, they are written as follows.

<sup>&</sup>lt;sup>2</sup>In group theory jargon this is essentially a *directed multigraph*, except that the vertical ordering of operators matters.

<sup>&</sup>lt;sup>3</sup>In group theory jargon this is an edge-labeled directed multigraph.

Example 1.3. The  $\Phi$ -normal Wick expansion of  $H_e$  in terms of Goldstone diagrams is

$$H_e = \boxtimes - \downarrow + \downarrow \\ - \downarrow$$

where  $\otimes$ — $\stackrel{\bullet}{\Diamond}$  represents the Φ-normal-ordered Fock operator,  $f_p^q \tilde{a}_q^p$ .

**Definition 1.5.** Contraction. A graph contraction is a map  $G \mapsto c(G)$  joining one or more compatible external lines in G. For example, c might replace  $l_1$  and  $l_2$ , which have ends  $(t(l_1), h(l_1)) = (o_1, e)$  and  $(t(l_2), h(l_2)) = (e, o_2)$ , with  $l_{12}$ , which has ends  $(t(l_{12}), h(l_{12})) = (o_1, o_2)$ . Contractions c and c' of G are equivalent if c(G) = c'(G). The number of c' equivalent to c is called the pattern degeneracy of c, denoted pat(c). The complete set of inequivalent graph contractions of G is here denoted Ctr(G).

## Definition 1.6. Connected graph.

**Definition 1.7.** Equivalent subgraph. If O' and L' are subsets of O and L then G' = (O', L', h, t) is a subgraph of G, denoted  $G' \subseteq G$ . Given an operator subset  $O' \subseteq O$ , the subgraph induced by O' is  $G[O'] \equiv (O', L', h, t)$  where L' contains all lines in L except for those incident on operators not in O'. If  $O_1$  and  $O_2$  are disjoint subsets of O such that  $G[O_1]$  and  $G[O_2]$  can be exchanged

If  $O_1$  and  $O_2$  are disjoint subsets of O inducing identical subgraphs  $G[O_1] \simeq G[O_2]$  that can be interchanged without altering G, we say that they are *equivalent subgraphs* and write  $G[O_1] \sim G[O_2]$ . See example A.3 for an example. If  $O_1 = \{o_1\}$  and  $O_2 = \{o_2\}$ , then  $o_1$  and  $o_2$  are equivalent operators. A pair of equivalent operators  $o \sim o'$  such that  $G[O_1 \cup \{o\}]$  and  $G[O_2 \cup \{o'\}]$  are equivalent subgraphs is called an *extension* of the equivalent subgraphs  $G[O_1] \sim G[O_2]$ .

**Lemma 1.1.** The pattern degeneracy of a contraction c of the graph G is given by  $\operatorname{pat}(c) = \frac{\operatorname{dg}(G)}{\operatorname{dg}(c(G))}$ 

Proof: First, consider c' and G' which equal c and G except that identical operators are treated as distinct. If  $L = L_1 \cup \cdots \cup L_m$  partitions L into equivalent lines, let  $L_{i,j}$  denote the subset of lines in  $L_i$  contracted with lines from  $L_j$  under c. Let  $L_{i,\text{ext}}$  denote the (possibly empty) subset of  $L_i$  which is external in c(G). The number of equivalent ways of partitioning  $L_i$  for contraction is then " $|L_i|$  choose  $|L_{i,\text{ext}}|$ ,  $|L_{i,1}|$ ,  $\cdots$ ,  $|L_{i,m}|$ ". The number of equivalent ways of contracting  $L_{i,j}$  with  $L_{j,i}$  is  $|L_{i,j}|! = |L_{j,i}|!$ . Therefore, the total pattern degeneracy of c' is

$$\operatorname{pat}(c') = \prod_{i=1}^{m} \binom{|L_i|}{|L_{i,\text{ext}}|, |L_{i,1}|, \dots, |L_{i,m}|} \prod_{i=1}^{m} \prod_{j=i+1}^{m} |L_{i,j}|! = \frac{|L_1|! \cdots |L_m|!}{\prod_{i=1}^{m} |L_{i,\text{ext}}|! |L_{i,1}|! \cdots |L_{i,i-1}|!} = \frac{\operatorname{dg}(G')}{\operatorname{dg}(c'(G'))}.$$

Theorem 1.1. Wick's theorem for diagrams.

Proof:

$$\tfrac{1}{\operatorname{dg}(G)} \sum_{\operatorname{Einstein}} \Sigma(G) = \tfrac{1}{\operatorname{dg}(G)} \sum_{\operatorname{Einstein}} : \Sigma(G) : + \sum_{c \in \operatorname{Ctr}(G)} \operatorname{pat}(c) \, \tfrac{1}{\operatorname{dg}(G)} \sum_{\operatorname{Einstein}} : \Sigma(c(G)) :$$

and since  $\operatorname{pat}(c)\operatorname{dg}(G)^{-1}=\operatorname{dg}(c(G))^{-1}$ , this is equivalent to  $G=:G:+\sum_{c\in\operatorname{Ctr}(G)}:c(G):$ 

**Proposition 1.1.** pat $(c) = \frac{dg(G)}{dg(c(G))}$ . Let G be graph and c(G) a contraction of G. Then the pattern degeneracy of c is equal to the line degeneracy of the original graph divided by the line degeneracy of the contracted graph. Proof:

$$G = (O, L)$$

$$O = O_1 \cup \dots \cup O_h = \bigcup_{i=1}^h O_{i,1} \cup \dots \cup O_{i,n_i}$$

First, define G' = (O', L) which equals G except that copies of the same operator are treated as distinct.

G = (O, L). First, consider G' = (O', L) which equals G except that,

$$G' = (O', L)$$

$$L = L_1 \cup \dots \cup L_g = \bigcup_{i=1}^g L_{i, \text{ext}} \cup L_{i, 1} \cup \dots \cup L_{i, g}$$

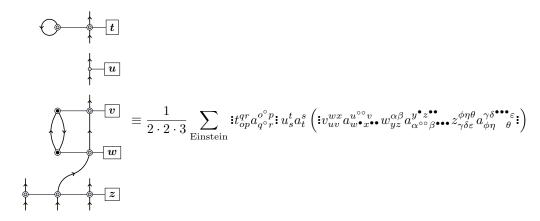
$$\lim(G') = |L_1|! \dots |L_g|!$$

$$\lim(\text{ctr}(G')) = \prod_{i=1}^g (|L_{i, \text{ext}}|! |L_{i, 1}|! \dots |L_{i, i-1}|!) = \prod_{i=1}^g (|L_{i, \text{ext}}|! |L_{i, i+1}|! \dots |L_{i, g}|!)$$

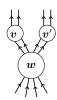
$$\operatorname{pat}(\operatorname{ctr}(G')) = \prod_{i=1}^{g} \left[ \binom{|L_{i}|}{|L_{i,\operatorname{ext}}|, |L_{i,1}|, \cdots, |L_{i,g}|} \cdot \prod_{j=1}^{i-1} |L_{i,j}|! \right] = \prod_{i=1}^{g} \left[ \frac{|L_{i}|!}{|L_{i,\operatorname{ext}}|! |L_{i,1}|! \cdots |L_{i,g}|!} \cdot |L_{i,1}|! \cdots |L_{i,i-1}|! \right] = \prod_{i=1}^{g} \left[ \frac{|L_{i}|!}{|L_{i,\operatorname{ext}}|! |L_{i,i+1}|! \cdots |L_{i,g}|!} \right] = \frac{\operatorname{lin}(G')}{\operatorname{lin}(\operatorname{ctr}(G'))}$$

## A. Parenthetical results

Example A.1. Graph.



**Example A.2.** Equivalent operators. If v and v' in the following graph are identical, then they are equivalent operators,  $v \sim v'$ . Otherwise, v and v' are interchangeable but inequivalent operators.



**Example A.3.** Equivalent subgraphs. If  $v \simeq v'$  and  $x \simeq x'$  in the following graph, then  $G[\{v, x\}]$  and  $G[\{v', x'\}]$  are equivalent subgraphs,  $G[\{v, x\}] \sim G[\{v', x'\}]$ . Otherwise, they are interchangeable but inequivalent subgraphs.

