# DENSITY CUMULANT THEORY FOR GROUND AND EXCITED ELECTRONIC STATES

by

#### Andreas Victor Copan

(Under the Direction of Henry F. Schaefer III)

#### Abstract

Here is my abstract. Here is my abstract.

INDEX WORDS: Cool stuff, Deep learning, Big data, Crowdsourcing, Computer vision, Computer science

# DENSITY CUMULANT THEORY FOR GROUND AND EXCITED ELECTRONIC STATES

by

Andreas Victor Copan

B.A., Bethel University, 2013

A Dissertation Submitted to the Graduate Faculty of The University of Georgia in Partial Fulfillment of the

Requirements for the Degree

DOCTOR OF PHILOSOPHY

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In	memory	of	Valery	And reiye vich	and	Valtraut	Kirsch	Copan

### Acknowledgments

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## Chapter 1

## Introduction

Here is the introduction. My dissertation introduces widgets. If I wish to cite someone, I could refer to them in text as in Pedersen, Fernández, and Koch<sup>1</sup>, or parenthetically<sup>1</sup>. I can also cite just the year 2001 or mention the author, Pedersen, Fernández, and Koch.

This dissertation template is handy. I like it and use it every time I write a dissertation.

— Abraham Lincoln 2001

#### 1.1 Stuff

Here is a section on stuff.

#### 1.1.1 Artificial Neural Networks

See Figure 1.1 for a neural network.

Cool.

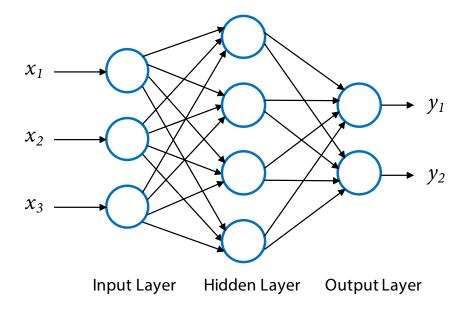


Figure 1.1: An example neural network with two final outputs. Notice how each neuron in one layer connects to each neuron in the following layer. This is called *fully connected*.

### Chapter 2

Ground-State Density Cumulant Theory:

Thermochemical and Kinetic Benchmark Calculations\*

#### 2.1 Abstract

We present an extensive benchmark study of density cumulant functional theory (DCFT) for thermochemistry and kinetics of closed- and open-shell molecules. The performance of DCFT methods (DC-06, DC-12, ODC-06, and ODC-12) is compared to that of coupled-electron pair methods (CEPA<sub>0</sub> and OCEPA<sub>0</sub>) and coupled-cluster theory (CCSD and CCSD(T)) for the description of noncovalent interactions (A24 database), barrier heights of hydrogen-transfer reactions (HTBH38), radical stabilization energies (RSE30), adiabatic ionization energies (AIE), and covalent bond stretching in diatomic molecules. Our results indicate that out of four DCFT methods the ODC-12 method is the most reliable and accurate DCFT formulation to date. Compared to CCSD, ODC-12 shows superior results for all benchmark tests employed in our study. With respect to coupled-pair

<sup>\*</sup>A. V. Copan, A. Yu. Sokolov, and H. F. Schaefer., J. Chem. Theory Comput. **10**, 2389 (2014). Adapted with permission of the American Chemical Society.

theories, ODC-12 outperforms CEPA<sub>0</sub>, and shows similar accuracy to the orbital-optimized CEPA<sub>0</sub> variant (OCEPA<sub>0</sub>) for systems at equilibrium geometries. For covalent bond stretching, ODC-12 is found to be more reliable than OCEPA<sub>0</sub>. For the RSE30 and AIE datasets, ODC-12 shows competitive performance with CCSD(T). In addition to benchmark results, we report new reference values for the RSE30 dataset computed using coupled cluster theory with up to perturbative quadruple excitations.

#### 2.2 Introduction

Recent developments in *ab initio* quantum chemistry have resulted in a variety of computational models for studying molecules. Apart from concerns about efficiency and accuracy, several concepts have evolved as criteria for judging the merits of a particular method. Energy-based criteria typically define an "ideal" approximation as one yielding correlation energies that are size-consistent, extensive<sup>2</sup>, well-defined (giving continuous, unique potential surfaces), and variational.<sup>3</sup> While it has been argued that the practical benefits of variationality are rather limited,<sup>4</sup> the efficiency of gradient computations, at least, is improved by formulating a theory in terms of a Hermitian and stationary energy functional.<sup>5</sup> With respect to scope and stability, methods that show consistent performance for open-shell systems, strongly correlated states, and non-equilibrium geometries are particularly valuable.<sup>4</sup>

The incorrect scaling of truncated configuration interaction (CI) energies with system size has inspired the development of size-extensive alternatives. Among the earliest formulations, the coupled electron pair approximations (CEPAs)<sup>6-10</sup> attracted much attention in 1970s, 11-15 offering rigorous extensivity and sizeconsistency while retaining much of the linearity <sup>16</sup> of CI in their equations. CEPA methods, however, have been shown to rapidly deteriorate as the molecular geometry deviates from equilibrium 16 and yield energies that vary under the rotation of the occupied orbitals. Partly in light of such defects, CEPA has been largely displaced by coupled-cluster (CC) theory. 4;17-23 In addition to size-extensivity, CC offers orbital invariance and improved stability for non-equilibrium structures <sup>16</sup>, but has a non-Hermitian energy functional and non-linear equations which are not readily amenable to parallel implementation. Although neither class of methods is strictly variational, VCEPA (variational CEPA) has been shown to be effectively equivalent to its non-variational counterpart. 24 Various other modifications to resolve the deficiencies of traditional CEPA have been explored, including self-consistent size-consistent CI,  $^{25;26}$  orbital-invariant CEPA,  $^{27;28}$  and orbital-invariant CEPA, optimized CEPA formulations. <sup>29–32</sup> Recently, the CEPA methods have been revived by Neese and co-workers <sup>24;33;34</sup> who developed the local pair-natural-orbital CEPA (LPNO-CEPA) methods and have implemented them for massively parallel computer architectures.

It has recently been demonstrated <sup>35–38</sup> that CEPA methods naturally arise in the context of theories that obtain the molecular energies from density cumulants, the connected and extensive components of the reduced density matrices (RDMs). <sup>39–44</sup> The advantage of cumulant-based theories is that, unlike their RDM-based counterparts, <sup>45–47</sup> they are naturally size-extensive and size-consistent. <sup>42;48</sup>

We have recently achieved the first implementation <sup>49;50</sup> of density cumulant functional theory (DCFT), proposed by Kutzelnigg in 2006. <sup>35</sup> In DCFT, the molecular energy is obtained in terms of a mean-field one-particle RDM and the two-particle density cumulant, constrained to be at least approximately N-representable (i.e. to correspond to a physical N-electron wavefunction). Like traditional CC theory, DCFT is size-extensive and orbital-invariant, but it has the additional advantage of a stationary and Hermitian energy functional, which simplifies the computation of molecular properties. In the original DCFT formulation (DC-06) <sup>35;49;50</sup> N-representability conditions derived from second-order Møller-Plesset perturbation theory (MPPT) were used, <sup>51</sup> yielding equations similar to those of the simplest CEPA model (CEPA<sub>0</sub>), <sup>8;10</sup> but including higher-order terms in the description of one-particle correlation effects. Using the same set of conditions, we have developed new formulations of DCFT that take advantage of an improved description of the one-particle density matrix (DC-12) <sup>52</sup> and full orbital optimization (ODC-06 and ODC-12 methods). <sup>53</sup>

Our previous studies <sup>49;50;52;53</sup> demonstrated for a limited set of systems that the DC-06, DC-12, ODC-06 and ODC-12 methods generally yield molecular energies and properties competitive with those obtained by CCSD and CCSD(T), but may exhibit unstable performance due to imbalances in the description of electron correlation. Herein, we present an extensive benchmark of the DCFT methods with respect to thermochemical and kinetic molecular properties, including noncovalent interactions, barrier heights in hydrogen-transfer reactions, radical stabilization energies, and adiabatic ionization energies for challenging electron-

dense systems. We conclude our benchmark study by testing the performance of DCFT for covalent bond stretching in diatomic molecules.

#### 2.3 Overview of DCFT

In this section a short overview of DCFT is presented. For details on the theory the reader is referred to our earlier publications. <sup>49;52;53</sup> In the RDM methods <sup>54</sup> the exact molecular energy is expressed as a functional of the one- and two-particle reduced density matrices,  $\gamma_1$  and  $\gamma_2$  (1-RDM and 2-RDM):

$$E = h_p^q \gamma_q^p + \frac{1}{2} g_{pq}^{rs} \gamma_{rs}^{pq} , \qquad [\boldsymbol{\gamma}_1]_q^p \equiv \gamma_q^p , \qquad [\boldsymbol{\gamma}_2]_{rs}^{pq} \equiv \gamma_{rs}^{pq} . \tag{2.1}$$

In eq. (2.1),  $h_p^q$  and  $g_{pq}^{rs}$  are the usual one- and two-electron integrals in the orthonormal spin-orbital basis  $\{\psi_p\}$  and summation over the repeated indices is implied. Expressing  $\gamma_1$  through  $\gamma_2$  via the partial trace relation  $\sum_r \gamma_{qr}^{pr} = (N-1)\gamma_q^p$ , the energy functional (2.1) can be minimized by varying  $\gamma_2$  subject to N-representability constraints. This is the essence of the variational 2-RDM approach. <sup>54</sup>

In DCFT, some of the challenges of the 2-RDM approach are circumvented by expanding  $\gamma_2$  in terms of its irreducible components – the 1-RDM and the two-particle cumulant (denoted by  $\lambda_2$ ):

$$\gamma_{rs}^{pq} = \gamma_r^p \gamma_s^q - \gamma_r^q \gamma_s^p + \lambda_{rs}^{pq}. \tag{2.2}$$

In eq. (2.2),  $\lambda_2$  describes the correlated part of  $\gamma_2$  that cannot be expressed via  $\gamma_1$ . The cumulant also determines the correlation contribution to  $\gamma_1$ , allowing

the 1-RDM to be decomposed as the sum of an idempotent 1-RDM ( $\kappa$ ) and a correlation correction ( $\tau$ ):

$$\gamma_1 = \kappa + \tau \,. \tag{2.3}$$

The correlation component  $\tau$  is fully specified by  $\lambda_2$ , whereas  $\kappa$  is independent of  $\lambda_2$ . eqs. (2.2) and (2.3) allow us to write an equivalent energy expression with  $\kappa$  and  $\lambda_2$  as independent functional parameters:

$$E[\kappa, \lambda_{2}] = \frac{1}{2} (h_{p}^{q} + f_{p}^{q}) (\kappa_{q}^{p} + \tau_{q}^{p}) + \frac{1}{4} \overline{g}_{pq}^{rs} \lambda_{pq}^{rs},$$

$$f_{p}^{q} = h_{p}^{q} + \overline{g}_{pr}^{qs} (\kappa_{s}^{r} + \tau_{s}^{r}), \quad \overline{g}_{rs}^{pq} = g_{rs}^{pq} - g_{rs}^{qp}.$$
(2.4)

## Chapter 3

Linear-Response Density Cumulant Theory for Excited States: First Implementation and Benchmark Calculations\*

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<sup>\*</sup>A. V. Copan and A. Yu. Sokolov (to be submitted in J. Chem. Theory Comput).

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## Chapter 4

## Linear-Response Density Cumulant Theory for Excited States: Better Algorithms, Bigger Systems

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Count	$x  ext{ Error}$	y Error
4	2.30	2.30
5	2.10	1.97
9	1.92	1.72
13	1.82	1.64

Table 4.1: Effect of count on error using our model. The error reduces as the count increases.

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## Chapter 5

#### Conclusion

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## Appendix A

#### **Extended Results**

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