CCD equations (projection method)

Derivation in KM tensor notation

The electronic Hamiltonian can be expressed as follows.

$$H_e = E_{HF} + H_c \qquad E_{HF} = \langle \Phi | H_e | \Phi \rangle \qquad H_c = f_p^q \tilde{a}_q^p + \frac{1}{4} \overline{g}_{pq}^{rs} \tilde{a}_{rs}^{pq} \qquad (1)$$

The CCD approximation parametrizes the wavefunction as $\Psi \approx e^{\hat{T}_2} \Phi$ where $\hat{T}_2 = \frac{1}{4} t^{ij}_{ab} a^{ab}_{ij}$. Substituting this ansatz into the Schrödinger equation and projecting by Φ and Φ^{ab}_{ij} gives the following system of equations.

$$E_{c} = \langle \Phi | H_{c} e^{\hat{T}_{2}} \Phi | \Phi \rangle \qquad \Longrightarrow \qquad E_{c} = \frac{1}{4} \langle \Phi | H_{c} | \Phi_{kl}^{cd} \rangle t_{cd}^{kl}$$

$$E_{c} t_{ab}^{ij} = \langle \Phi_{ij}^{ab} | H_{c} e^{\hat{T}_{2}} \Phi | \Phi \rangle \qquad \Longrightarrow \qquad E_{c} t_{ab}^{ij} = \langle \Phi_{ij}^{ab} | H_{c} | \Phi \rangle + \frac{1}{4} \langle \Phi_{ij}^{ab} | H_{c} | \Phi_{kl}^{cd} \rangle t_{cd}^{kl} + \frac{1}{2} \left(\frac{1}{4} \right)^{2} \langle \Phi_{ij}^{ab} | H_{c} | \Phi_{klmn}^{cdef} \rangle t_{cd}^{kl} t_{ef}^{mn}$$

$$(3)$$

The non-trivial terms to be evaluated are

$$\frac{1}{4} \langle \Phi_{ij}^{ab} | H_c | \Phi_{kl}^{cd} \rangle t_{cd}^{kl} = \frac{1}{4} f_p^q t_{cd}^{kl} (\tilde{a}_{ab}^{ij} \tilde{a}_p^p \tilde{a}_{kl}^{cd})_{\text{f.c.}} + \left(\frac{1}{4}\right)^2 \overline{g}_{pq}^{rs} t_{cd}^{kl} (\tilde{a}_{ab}^{ij} \tilde{a}_{rs}^{pq} \tilde{a}_{kl}^{cd})_{\text{f.c.}}$$
(4)

$$\frac{1}{2} \left(\frac{1}{4}\right)^2 \langle \Phi_{ij}^{ab} | H_c | \Phi_{klmn}^{cdef} \rangle t_{cd}^{kl} t_{ef}^{mn} = \frac{1}{2} \left(\frac{1}{4}\right)^3 \overline{g}_{pq}^{rs} t_{cd}^{kl} t_{ef}^{mn} (\tilde{a}_{ab}^{ij} \tilde{a}_{rs}^{pq} \tilde{a}_{kl}^{cd} \tilde{a}_{mn}^{ef})_{\text{f.c.}}$$

$$(5)$$

where $(\tilde{a}_{ab}^{ij}\tilde{a}_{q}^{p}\tilde{a}_{kl}^{cd})_{\text{f.c.}}$ and $(\tilde{a}_{ab}^{ij}\tilde{a}_{rs}^{pq}\tilde{a}_{kl}^{cd})_{\text{f.c.}}$ can be determined using Wick's theorem

$$\begin{array}{l} (\tilde{a}_{ab}^{ij}\tilde{a}_{q}^{p}\tilde{a}_{kl}^{cd})_{\mathrm{f.c.}} = \hat{P}_{(a/b)|k/l}^{(c/d)}; \tilde{a}_{a^{-1}b^{-3}}^{i^{-1}j^{-2}}\tilde{a}_{q^{-2}}^{p^{-2}}\tilde{a}_{k^{-1}l^{-2}}^{c^{-2}d^{-3}} + \hat{P}_{(k/l)}^{(i/j)|c/d})_{i} \tilde{a}_{a^{-1}b^{-3}}^{i^{-1}j^{-3}}\tilde{a}_{q^{-1}}^{p^{-2}}\tilde{a}_{k^{-2}l^{-3}}^{c^{-1}d^{-2}}; \\ = \hat{P}_{(a/b)|k/l}^{(c/d)}, p_{a}^{p}\eta_{c}^{p}\eta_{b}^{h}\chi_{l}^{h}l^{-1} - \hat{P}_{(k/l)}^{(i/j)|c/d}^{(i/j)|c/d}, q_{l}^{p}\eta_{l}^{p}\eta_{c}^{a}\eta_{b}^{b} \\ (\tilde{a}_{ab}^{ij}\tilde{a}_{rs}^{pq}\tilde{a}_{kl}^{cd})_{\mathrm{f.c.}} = \hat{P}_{(a/b)|k/l}^{(c/d)}; \tilde{a}_{a^{-1}b^{-2}}^{i^{-1}j^{-2}}\tilde{a}_{r^{-3}s^{-2}}^{p^{-3}q^{-2}}\tilde{a}_{k^{-1}l^{-2}}^{c^{-3}l^{-2}} + \hat{P}_{(k/l)}^{(i/j)|c/d}; \tilde{a}_{a^{-1}b^{-2}}^{i^{-1}j^{-2}}\tilde{a}_{r^{-1}s^{-2}}^{p^{-3}q^{-2}}\tilde{a}_{k^{-2}l^{-3}}^{c^{-3}l^{-2}}; \\ + \hat{P}_{(r/s|k/l)a/b}^{(r/s)|k/l^{-2}}, \tilde{a}_{a^{-1}b^{-2}}^{i^{-1}j^{-2}}\tilde{a}_{r^{-3}s^{-2}}^{p^{-3}q^{-2}}\tilde{a}_{k^{-2}l^{-3}}^{c^{-2}l^{-3}}; \\ + \hat{P}_{(r/s|k/l)a/b}^{(r/s)|k/l^{-2}}, \tilde{a}_{a^{-1}b^{-2}}^{i^{-1}j^{-2}}\tilde{a}_{r^{-3}s^{-2}}^{p^{-3}q^{-2}}\tilde{a}_{k^{-2}l^{-3}}^{c^{-2}l^{-3}}; \\ \hat{a}_{a^{0}b}^{pq}\tilde{a}_{rs}^{cd}\tilde{a}_{k^{0}}\tilde{a}_{l}\tilde{a}_{mb}^{-1}, \tilde{a}_{a^{-1}b^{-2}}\tilde{a}_{r^{-3}s^{-2}}^{p^{-3}q^{-2}}\tilde{a}_{k^{-2}l^{-3}}^{c^{-2}l^{-2}}; \\ \hat{a}_{a^{0}b}^{pq}\tilde{a}_{rs}^{cd}\tilde{a}_{k^{0}}\tilde{a}_{l}\tilde{a}_{mb}^{-1}, \tilde{a}_{a^{-1}b^{-2}}^{p^{-3}q^{-2}}\tilde{a}_{r^{-3}s^{-2}}^{p^{-3}q^{-2}}\tilde{a}_{k^{-2}l^{-2}}^{p^{-3}q^{-2}}\tilde{a}_{r^{-3}s^{-2}}^{p^{-3}q^{-4}}; \\ \hat{a}_{a^{0}b}^{pq}\tilde{a}_{rs}^{cd}\tilde{a}_{k^{0}}\tilde{a}_{l}\tilde{a}_{mb}^{-1}, \tilde{a}_{a^{0}b^{-2}}^{p^{-3}q^{-2}}\tilde{a}_{a^{0}b^{-2}}^{p^{-3}q^{-3}}\tilde{a}_{k^{-1}l^{-2}}^{p^{-3}q^{-3}}\tilde{a}_{r^{-3}s^{-3}}^{p^{-3}q^{-2}} \\ \hat{a}_{a^{0}b}^{pq}\tilde{a}_{rs}^{pq}\tilde{a}_{l}\tilde{a}_{k^{-1}l^{-2}}^{p^{-3}q^{-3}}\tilde{a}_{r^{-3}s^{-3}}^{p^{-3}q^{-3}}\tilde{a}_{r^{-3}s^{-3}}^{p^{-3}q^{-3}} \\ \hat{a}_{a^{0}b}^{pq}\tilde{a}_{rs}^{pq}\tilde{a}_{l}\tilde{a}_$$

giving the following.

$$\begin{split} &\frac{1}{4} f_p^a t_{cd}^{kl} (\tilde{a}_{ab}^{ij} \tilde{a}_p^a \tilde{a}_{kl}^{cd})_{\text{f.c.}} = f_p^a t_{cd}^{kl} \hat{P}_{(a/b)} \eta_p^a \eta_q^c \eta_b^d \gamma_i^b \gamma_j^l - f_p^a t_{cd}^{kl} \hat{P}^{(i/j)} \gamma_i^a \gamma_k^p \gamma_l^l \eta_a^c \eta_b^d \\ &= \hat{P}_{(a/b)} f_a^c t_b^{ij} - \hat{P}^{(i/j)} f_k^i t_a^{kj} \\ &(\frac{1}{4})^2 \, \overline{g}_{pq}^{rs} t_{cd}^{kl} (\tilde{a}_{ab}^{ij} \tilde{a}_{rs}^{pq} \tilde{a}_{kl}^{cd})_{\text{f.c.}} = \frac{1}{4} \overline{g}_{pq}^{rs} t_{cd}^{kl} \hat{P}_{(a/b)} \eta_p^a \eta_p^a \eta_p^c \eta_s^d \gamma_k^i \gamma_l^j + \frac{1}{4} \overline{g}_{pq}^{rs} t_{cd}^{kl} \hat{P}^{(i/j)} \gamma_r^i \gamma_s^j \gamma_k^p \gamma_l^q \eta_a^c \eta_b^d \\ &- \overline{g}_{pq}^{rs} t_{cd}^{kl} \hat{P}_{(a/b)} \eta_r^i \gamma_r^k \gamma_l^j \eta_a^q \eta_s^c \eta_b^d \\ &= \frac{1}{4} \hat{P}_{(a/b)} \overline{g}_{ad}^{cd} t_{cd}^i + \frac{1}{4} \hat{P}^{(i/j)} \overline{g}_{kl}^{il} t_{ab}^{kl} - \hat{P}_{(a/b)}^{(i/j)} \overline{g}_{ka}^{ic} t_{cb}^{kj} \\ &= \frac{1}{2} \overline{g}_{ab}^{cd} t_{cd}^{ij} + \frac{1}{2} \overline{g}_{kl}^{ij} t_{ab}^k + \hat{P}_{(a/b)}^{(i/j)} \overline{g}_{ak}^{ic} t_{bc}^{jk} \\ &= \frac{1}{2} \overline{g}_{ab}^{rs} t_{cd}^{kl} + \hat{P}_{(a/b)}^{ij} \overline{g}_{ab}^{ic} t_{bc}^{kl} + \hat{P}_{(a/b)}^{ij} \overline{g}_{ab}^{ic} t_{bc}^{jk} \\ &= \frac{1}{2} \overline{g}_{pq}^{rs} t_{cd}^{kl} t_{ef}^{rm} \gamma_k^i \gamma_l^j \gamma_p^m \gamma_n^q \alpha_n^c \eta_b^a \eta_r^c \eta_b^f \\ &= \frac{1}{2} \overline{g}_{pq}^{rs} t_{cd}^{kl} t_{ef}^{rm} \hat{\gamma}_k^i \gamma_l^j \gamma_p^m \gamma_n^q \alpha_n^c \eta_b^a \eta_r^c \eta_b^f \\ &- \frac{1}{2} \overline{g}_{pq}^{rs} t_{cd}^{kl} t_{ef}^{rm} \hat{P}_{(a/b)}^i \gamma_k^i \gamma_l^j \gamma_p^m \gamma_n^q \alpha_n^c \eta_b^a \eta_r^c \eta_b^f \\ &- \frac{1}{2} \overline{g}_{pq}^{rs} t_{cd}^{kl} t_{ef}^{rm} \hat{P}_{(i/j)}^i \gamma_k^i \gamma_l^j \gamma_p^m \gamma_n^q \alpha_n^c \eta_b^a \eta_r^c \eta_b^f \\ &+ \frac{1}{4} \overline{g}_{pq}^{rs} t_{cd}^{kl} t_{ef}^r \hat{P}_{(a/b)}^i \gamma_k^i \gamma_l^i \gamma_p^m \gamma_n^a \eta_a^c \eta_b^a \eta_r^c \eta_b^f \\ &= \frac{1}{2} \overline{g}_{pq}^{rs} t_{cd}^{kl} t_{ef}^{rm} \hat{P}_{(a/b)}^i \gamma_k^i \gamma_l^i \gamma_p^m \gamma_n^a \eta_b^c \eta_b^c \eta_b^c \\ &= \frac{1}{4} \overline{g}_{ef}^{rs} t_{cd}^{kl} t_{ef}^{rm} \hat{P}_{(a/b)}^i \gamma_k^i \gamma_l^i \gamma_l^m \gamma_n^a \eta_b^c \eta_b^c \eta_b^c \\ &= \frac{1}{4} \overline{g}_{ef}^{rs} t_{cd}^{kl} t_{ef}^i \hat{P}_{(a/b)}^i \gamma_l^i \gamma_l^i$$

Substituting these terms, along with the second Slater rule $(\langle \Phi_{ij}^{ab}|H_c|\Phi\rangle = \overline{g}_{ab}^{ij})$ and $\langle \Phi|H_c|\Phi_{kl}^{cd}\rangle = \overline{g}_{kl}^{cd}$, into equations 2 and 3 gives the following.

$$E_{c} = \frac{1}{4} \overline{g}_{kl}^{cd} t_{cd}^{kl}$$

$$E_{c} t_{ab}^{ij} = \overline{g}_{ab}^{ij} + \hat{P}_{(a/b)} f_{a}^{c} t_{cb}^{ij} - \hat{P}^{(i/j)} f_{k}^{i} t_{ab}^{kj} + \frac{1}{2} \overline{g}_{ab}^{cd} t_{cd}^{ij} + \frac{1}{2} \overline{g}_{kl}^{ij} t_{ab}^{kl} + \hat{P}_{(a/b)}^{(i/j)} \overline{g}_{ak}^{ic} t_{bc}^{jk}$$

$$+ \left(\frac{1}{4} \overline{g}_{kl}^{cd} t_{cd}^{kl} \right) t_{ab}^{ij} - \frac{1}{2} \hat{P}_{(a/b)} \overline{g}_{kl}^{cd} t_{ac}^{ij} t_{bd}^{kl} - \frac{1}{2} \hat{P}^{(i/j)} \overline{g}_{kl}^{cd} t_{ab}^{ik} t_{cd}^{jl} + \frac{1}{4} \overline{g}_{kl}^{cd} t_{cd}^{ij} t_{ab}^{kl} + \hat{P}^{(i/j)} \overline{g}_{kl}^{cd} t_{ac}^{ik} t_{bd}^{jl}$$

$$(7)$$

Finally, using the energy expression 6 allows us to cancel the left-hand side of equation 7 with the seventh term on the right, leaving the traditional CCD \hat{T}_2 amplitude equations.

$$0 = \overline{g}_{ab}^{ij} + \hat{P}_{(a/b)} f_a^c t_{cb}^{ij} - \hat{P}^{(i/j)} f_k^i t_{ab}^{kj} + \frac{1}{2} \overline{g}_{ab}^{cd} t_{cd}^{ij} + \frac{1}{2} \overline{g}_{kl}^{ij} t_{ab}^{kl} + \hat{P}_{(a/b)}^{(i/j)} \overline{g}_{ak}^{ic} t_{bc}^{jk} - \frac{1}{2} \hat{P}_{(a/b)} \overline{g}_{kl}^{cd} t_{ac}^{ij} t_{bd}^{kl} - \frac{1}{2} \hat{P}_{(a/b)}^{ij} \overline{g}_{kl}^{cd} t_{ab}^{ij} + \frac{1}{4} \overline{g}_{kl}^{cd} t_{cd}^{ij} t_{ab}^{kl} + \hat{P}^{(i/j)} \overline{g}_{kl}^{cd} t_{ac}^{ik} t_{bd}^{jl}$$

$$(8)$$

Assuming a canonical HF reference wavefunction (so that $f_a^c = \varepsilon_a \delta_a^c$ and $f_k^i = \varepsilon_i \delta_k^i$) and moving the second and third terms to the left-hand side gives the CCD working equations

$$t_{ab}^{ij} = \mathcal{E}_{ab}^{ij} \left(\overline{g}_{ab}^{ij} + \frac{1}{2} \overline{g}_{ab}^{cd} t_{cd}^{ij} + \frac{1}{2} \overline{g}_{kl}^{ij} t_{ab}^{kl} + \hat{P}_{(a/b)}^{(i/j)} \overline{g}_{ak}^{ic} t_{bc}^{jk} - \frac{1}{2} \hat{P}_{(a/b)} \overline{g}_{kl}^{cd} t_{ab}^{ij} t_{bd}^{kl} - \frac{1}{2} \hat{P}^{(i/j)} \overline{g}_{kl}^{cd} t_{ab}^{ik} t_{cd}^{jl} \right)$$

$$+ \frac{1}{4} \overline{g}_{kl}^{cd} t_{cd}^{ij} t_{ab}^{kl} + \hat{P}^{(i/j)} \overline{g}_{kl}^{cd} t_{ac}^{ik} t_{bd}^{jl}$$

$$\text{where } \mathcal{E}_{ab}^{ij} = \frac{1}{\varepsilon_i + \varepsilon_i - \varepsilon_a - \varepsilon_b}.$$

$$(10)$$