

Tutorial on Nonlinear Reduced Order Modelling

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Advanced Structural Dynamics

Why ROMs?



High fidelity Models:
High number of DOFs

The dynamics live in a much
lower dimensional space

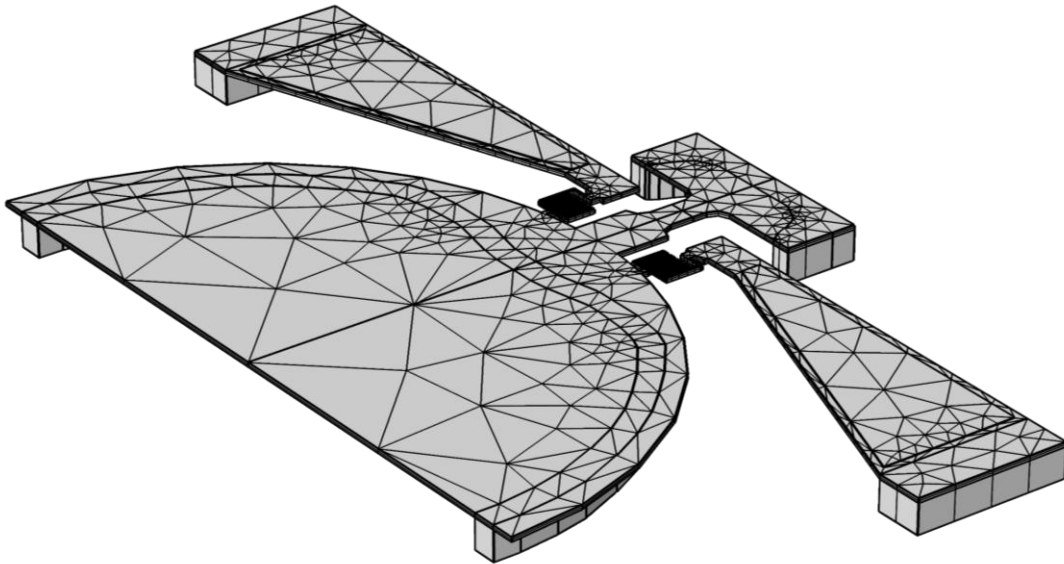
Run **full** dynamic simulations:
Hours-days

With **ROM**:
Seconds-minutes

Hard to accelerate computationally

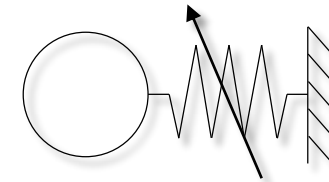
Reduced Order Modelling (ROM)

Large scale 3D finite elements model



Nonlinear Dynamical system
with up to **Millions** of degrees of freedom

Capture the **characteristic** dynamics
with **low dimensional** models



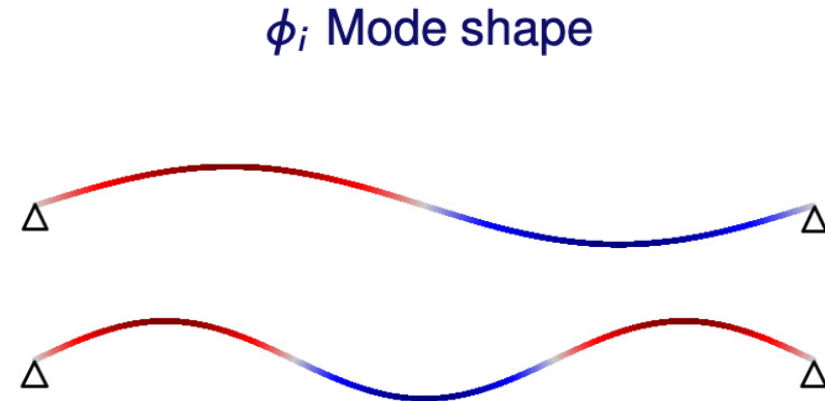
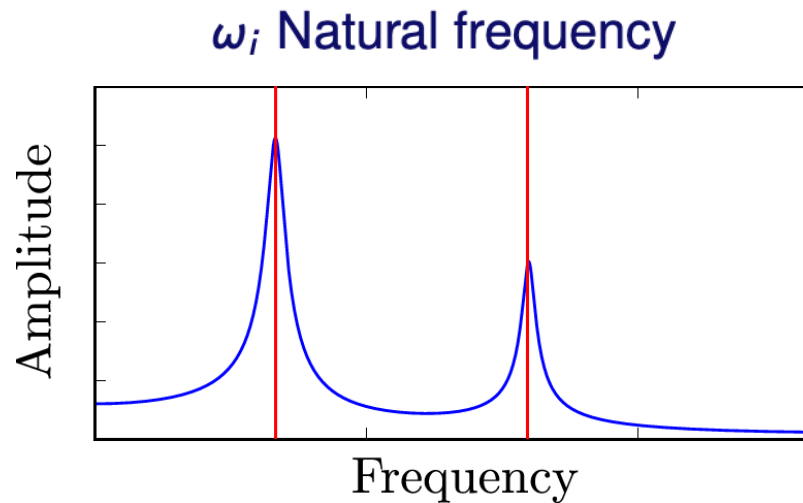
Desired Properties:

- **Small** number of parameters
- Physically **interpretable**
- From FE model
- **Fast** yet **accurate**

The simplest ROM

Large scale 3D finite elements model with up to millions of degrees of freedom:

$$\mathbf{M}\ddot{\mathbf{U}} + \mathbf{K}\mathbf{U} = \mathbf{0}$$



Change of Coordinates
Reduced Dynamics

$$\mathbf{U} = \Phi \mathbf{u}$$

$$\ddot{\mathbf{u}} + \mathbf{\Omega}^2 \mathbf{u} = \mathbf{0}$$

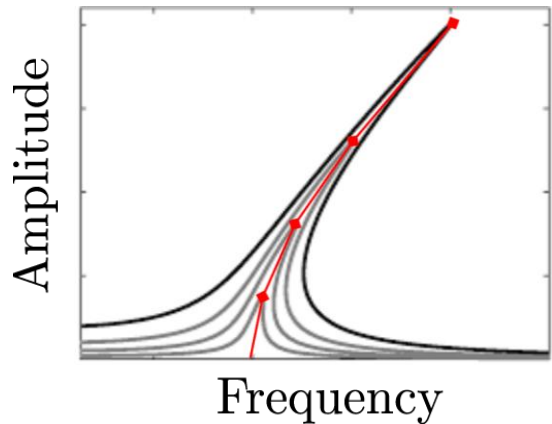
Linear Modal Analysis is the
most trivial of ROMs

What happens when the system is Nonlinear?

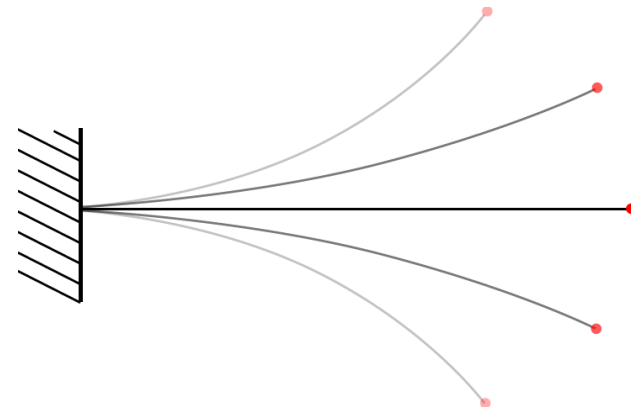
Large scale 3D finite elements model with up to **millions** of degrees of freedom:

$$\mathbf{M}\ddot{\mathbf{U}} + \mathbf{K}\mathbf{U} + \mathbf{G}\mathbf{U}\mathbf{U} + \mathbf{H}\mathbf{U}\mathbf{U}\mathbf{U} = \mathbf{0}$$

Natural frequencies $\omega_{nl} = \omega_{nl}(a)$



Mode shapes $\phi_{nl} = \phi_{nl}(a)$



Change of Coordinates

$$\mathbf{U} = \mathbf{U}(\mathbf{u}, \dot{\mathbf{u}})$$

Reduced Dynamics

$$\ddot{\mathbf{u}} + \omega_n^2 \mathbf{u} + \mathbf{f}(\mathbf{u}, \dot{\mathbf{u}}) = \mathbf{0}$$

The two main ingredients
of every Nonlinear ROM

Naïve approach

Large scale FE model:

$$\mathbf{M}\ddot{\mathbf{U}} + \mathbf{K}\mathbf{U} + \mathbf{G}\mathbf{U}\mathbf{U} + \mathbf{H}\mathbf{U}\mathbf{U}\mathbf{U} = \mathbf{0}$$

$$\mathbf{U} = \boldsymbol{\Phi}\mathbf{u} \rightarrow \ddot{\mathbf{u}} + \boldsymbol{\Omega}^2\mathbf{u} + \mathbf{g}\mathbf{u}\mathbf{u} + \mathbf{h}\mathbf{u}\mathbf{u}\mathbf{u} = \mathbf{0}$$

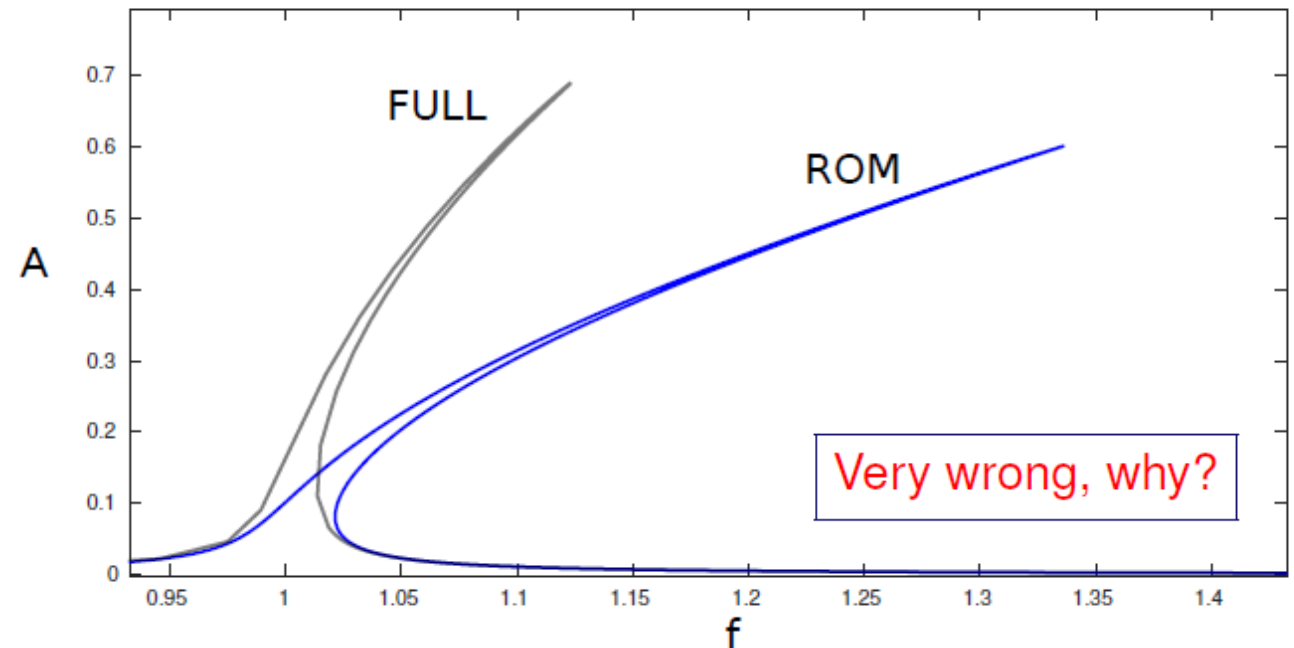
Reduce on the first mode:

all modes zero, $u_1 \neq 0$

Reduced Dynamics

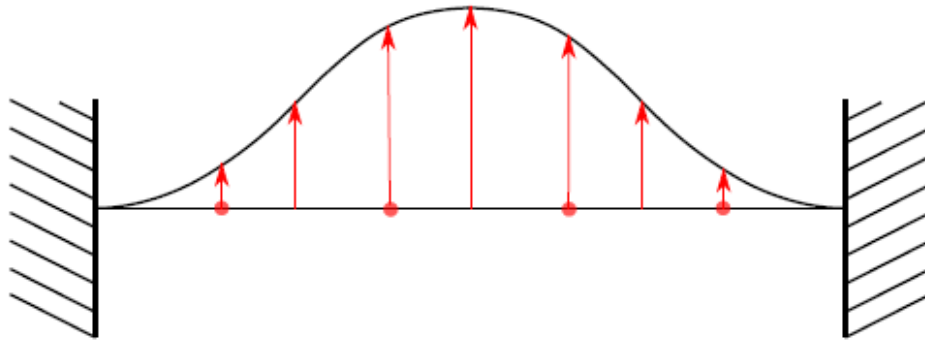
$$\ddot{u}_1 + \omega_1^2 u_1 + g_{11}^1 u_1^2 + h_{111}^1 u_1^3 = 0$$

Classic testcase:
Clamped-Clamped beam

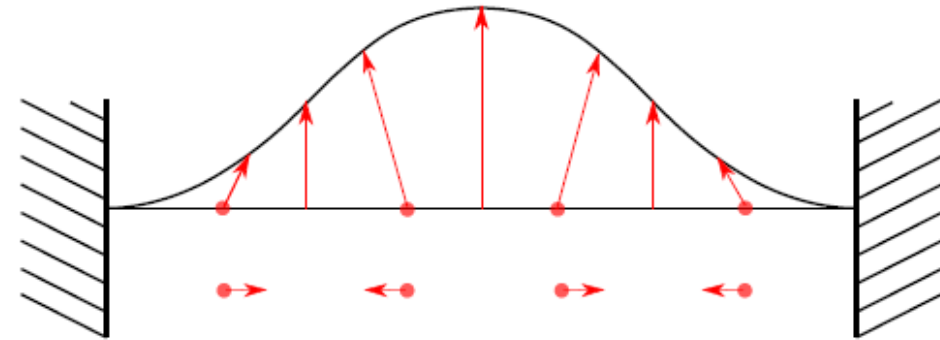


Membrane bending couplings

Linear modeshape



Nonlinear modeshape



Constrains the nonlinear structure on its linear modeshape

Using the Linear mode as a basis:
cannot account for flexural-axial couplings!

4th axial mode has to be included

$$\ddot{u}_1 + \omega_1^2 u_1 + g_{14}^1 u_1 p_4 + h_{111}^1 u_1^3 = 0$$

$$\ddot{u}_4 + \omega_4^2 u_4 + g_{11}^4 u_1^2 = 0$$

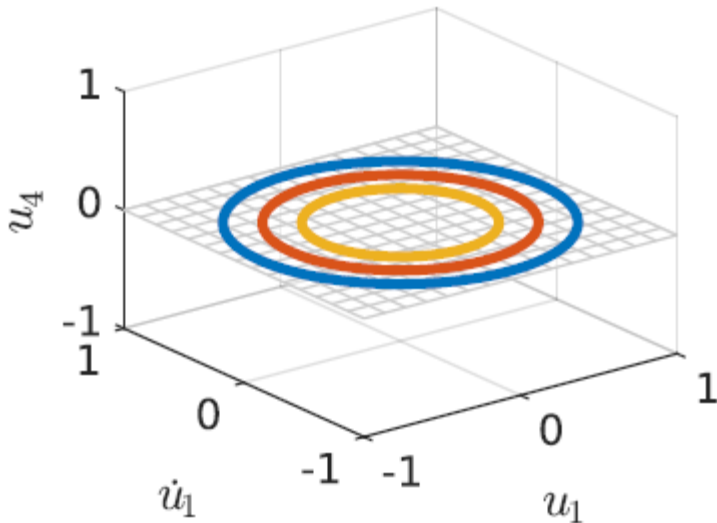
Invariant breaking terms

$$\ddot{u}_4 + \omega_4^2 u_4 + g_{11}^4 u_1^2 = 0$$

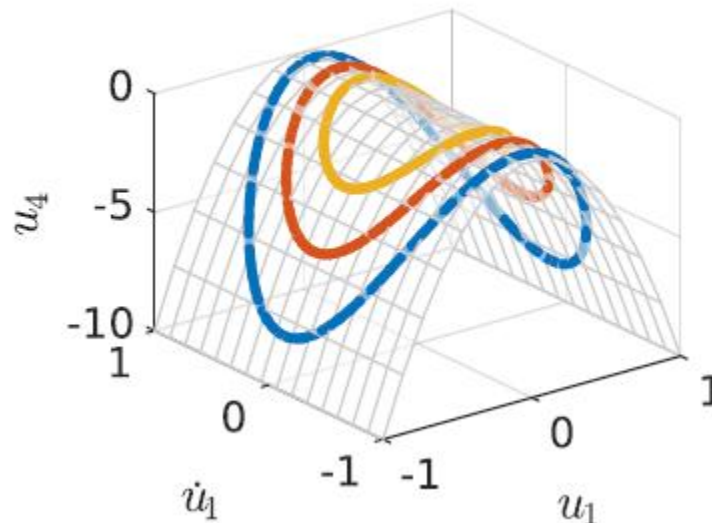
g_{11}^4 is the quadratic invariant breaking term

First modal amplitude nonzero \Rightarrow the coupled mode: forcing like term in u_1^2

Linear Invariance



Nonlinear Invariant manifold



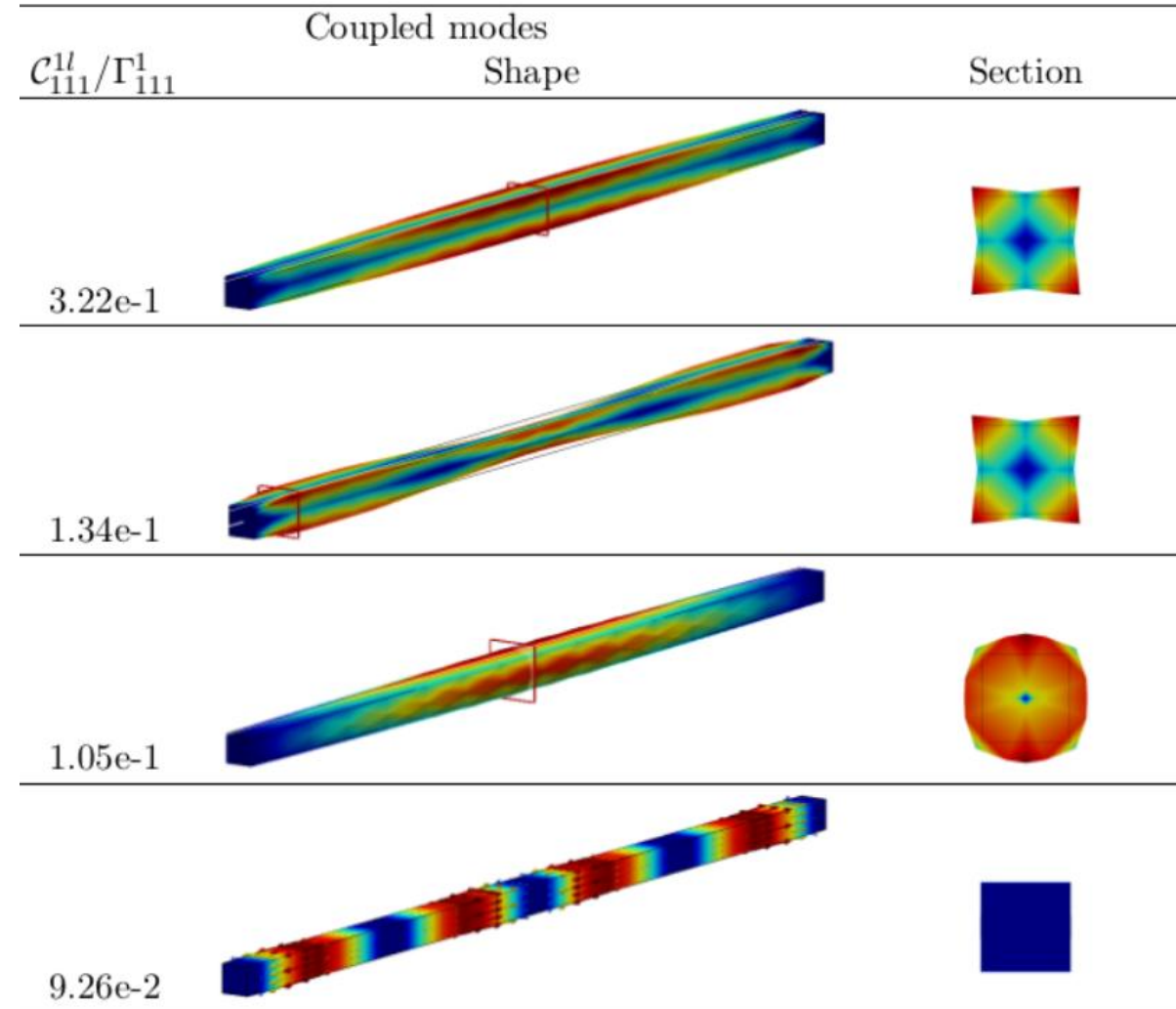
4th axial mode has to be included

So what's the problem?

- Many modes coupled to a single master mode
- Distributed all over the full spectrum

(Q1) How can we find them?

(Q2) How can we 'enslave' them?



Outline

- ROM by condensation
- Linear VS Nonlinear ROMs
- ROM by projection
- ROM by invariance

(A2) Static condensation (1D beam elements)

$$\ddot{u}_1 + \omega_1^2 u_1 + g_{14}^1 u_1 p_4 + h_{111}^1 u_1^3 = 0$$

$$\ddot{u}_4 + \omega_4^2 u_4 + g_{11}^4 u_1^2 = 0$$

Now, perfect!

Slow-Fast Assumption:

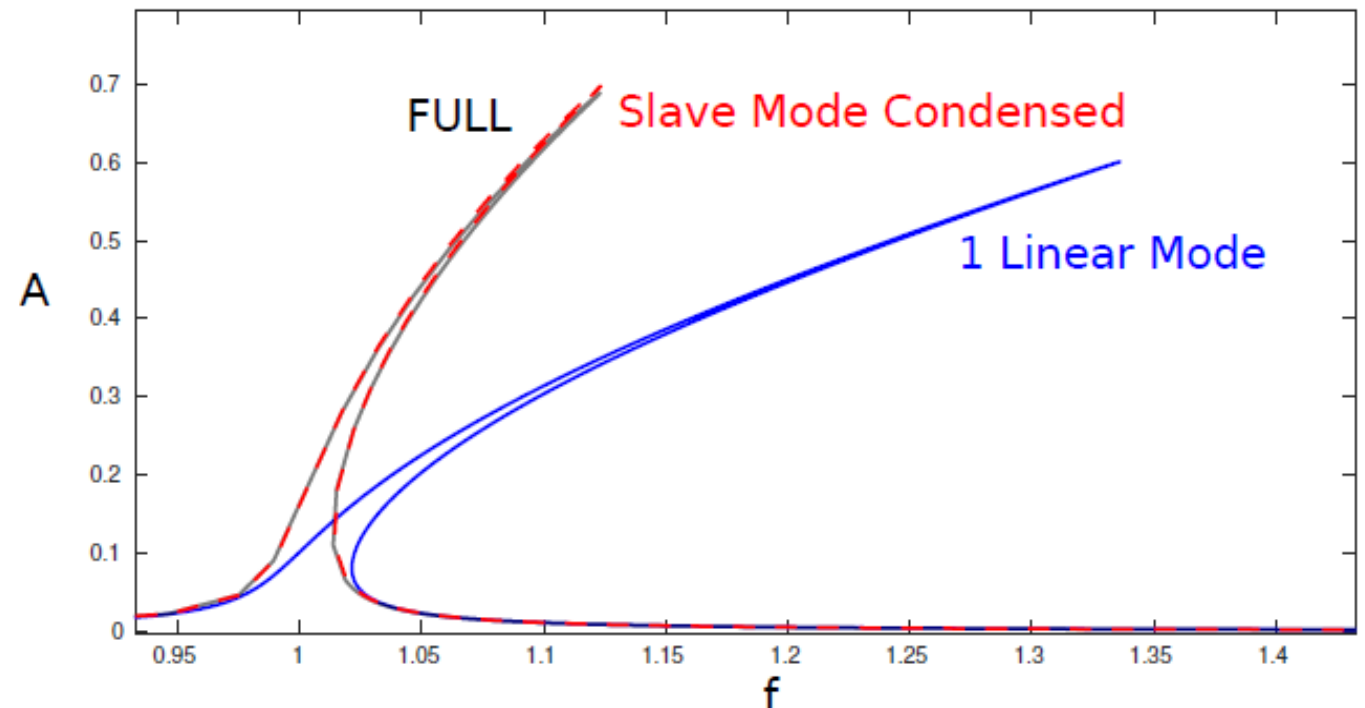
$$\omega_4 \gg \omega_1$$

$$u_4 = -\frac{g_{11}^4}{\omega_4^2} u_1^2$$

$$\ddot{u}_1 + \omega_1^2 u_1 + \left(h_{111}^1 - \frac{g_{14}^1 g_{11}^4}{\omega_4^2} \right) u_1^3 = 0$$

Correction Factor:

$$\frac{g_{14}^1 g_{11}^4}{\omega_4^2}$$



(A2) Static condensation (3D beam elements)

$$\ddot{u}_1 + \omega_1^2 u_1 + g_{1s}^1 u_1 u_s + h_{111}^1 u_1^3 = 0$$

$$\ddot{u}_s + \omega_s^2 u_s + g_{11}^s u_1^2 = 0$$

Slow-Fast Assumption:

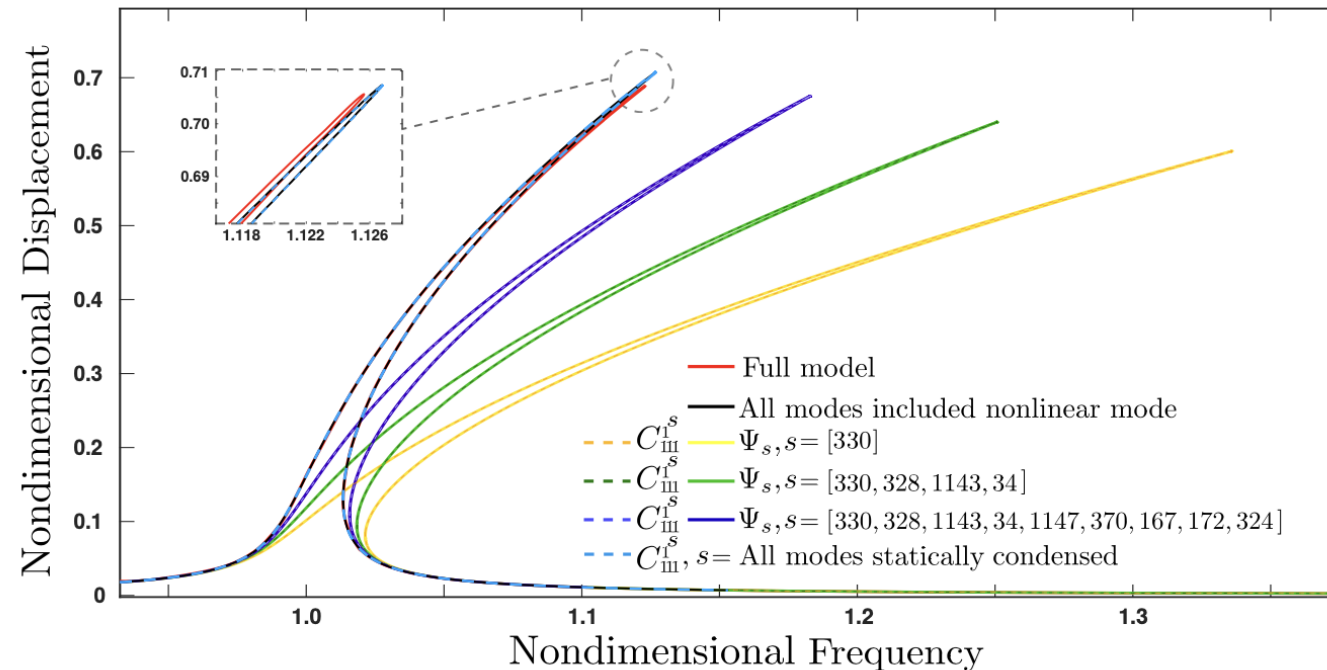
$$\omega_s \gg \omega_1$$

$$u_s = -\frac{g_{11}^s}{\omega_s^2} u_1^2$$

$$\ddot{u}_1 + \omega_1^2 u_1 + \left(h_{111}^1 - \sum \frac{g_{1s}^1 g_{11}^s}{\omega_s^2} \right) u_1^3 = 0$$

Correction Factor: $\sum \frac{g_{1s}^1 g_{11}^s}{\omega_s^2}$

(Q1) How can we find them?



(A1) EITHER.. Implicit condensation (ICE)

$$\mathbf{K}\mathbf{U} + \mathbf{G}\mathbf{U}\mathbf{U} + \mathbf{H}\mathbf{U}\mathbf{U}\mathbf{U} = \mathbf{M}\phi_1\alpha$$

$$\omega_1^2 u_1 + g_{1s}^1 u_1 u_s + h_{111}^1 u_1^3 = \alpha$$

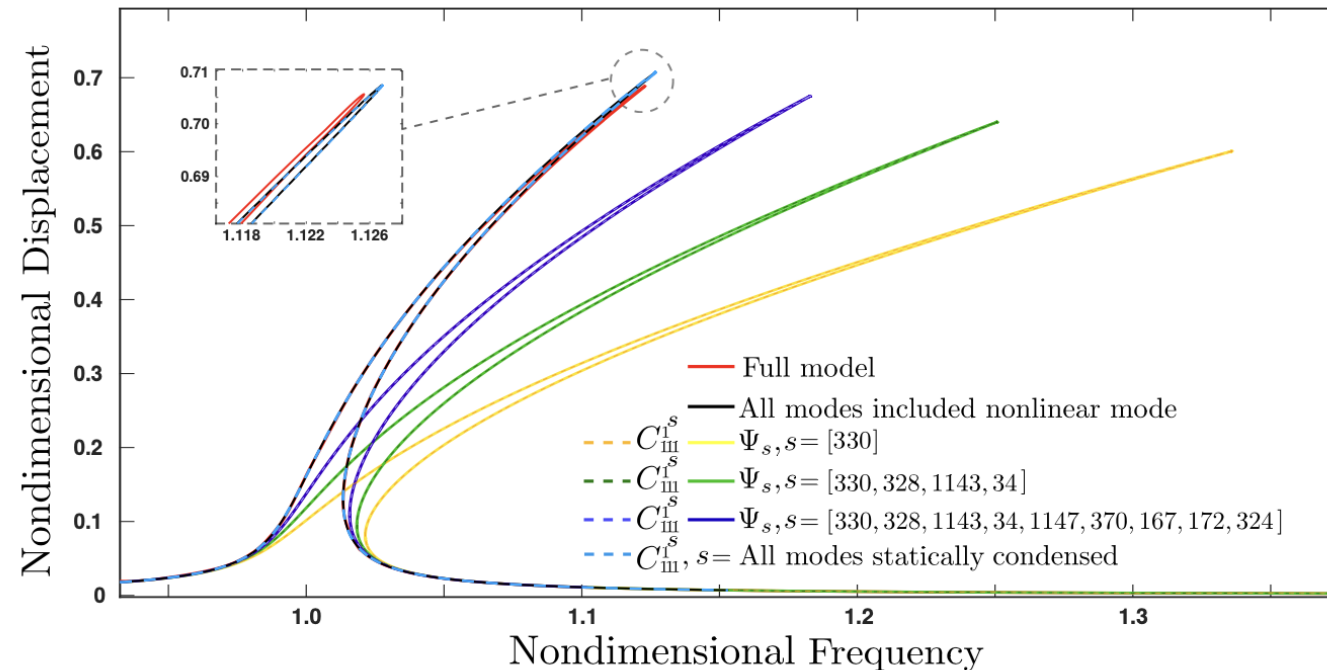
$$\omega_s^2 u_s + g_{11}^s u_1^2 = 0$$

$$u_s = -\frac{g_{11}^s}{\omega_s^2} u_1^2$$

$$\omega_1^2 u_1 + \left(h_{111}^1 - \sum \frac{g_{1s}^1 g_{11}^s}{\omega_s^2} \right) u_1^3 = \alpha$$

$$\alpha = \alpha(u_1)$$

$$\ddot{u}_1 + \alpha(u_1) = 0$$



(A1) OR.. Static Modal Derivative (SMD)

Perturbed eigenproblem:

$$\left(-\tilde{\omega}_1^2(u_1)\mathbf{M} + \tilde{\mathbf{K}}(u_1) \right) \tilde{\phi}_1(u_1) = \mathbf{0}$$

Neglecting inertial contribution:

$$\mathbf{K} \frac{\partial \tilde{\phi}_1(u_1)}{\partial u_1} + \frac{\partial \tilde{\mathbf{K}}(u_1)}{\partial u_1} \phi_1 = \mathbf{0}$$

$$\frac{\partial \tilde{\mathbf{K}}}{\partial u_1} \phi_1 = \mathbf{G} \phi_1 \phi_1^2$$

$$\theta_{11} = \frac{\partial \tilde{\phi}_1(u_1)}{\partial u_1}$$

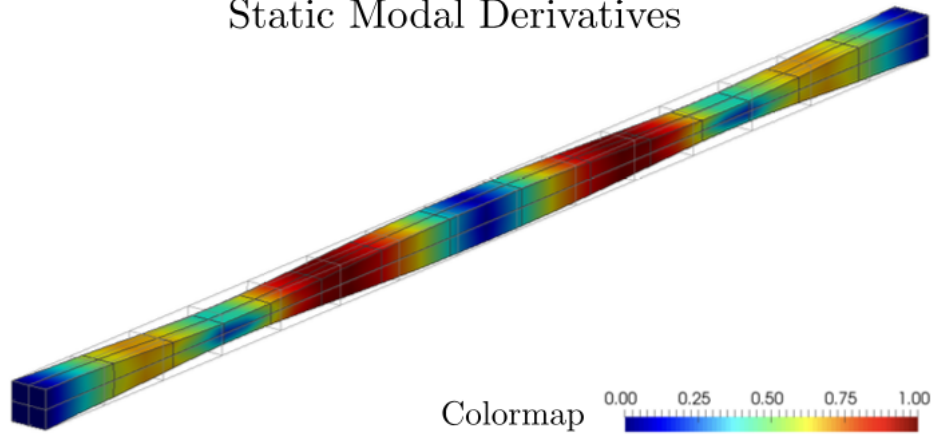
Static modal derivative:

$$\theta_{11} = -\mathbf{K}^{-1} \mathbf{G} \phi_1 \phi_1^2$$

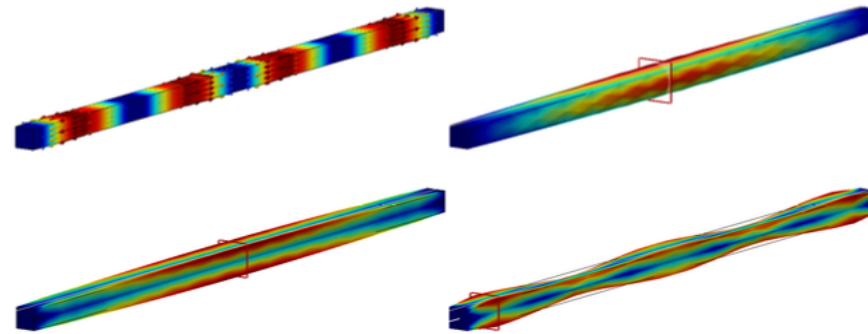
(A1) OR.. Static Modal Derivative (SMD)

$$\theta_{11} = -\mathbf{K}^{-1} \mathbf{G} \phi_1 \phi_1$$

Static Modal Derivatives



Non-bending modes



Inclusion of all slave modes under the slow fast assumption:

$$\theta_{11} = - \sum 2\phi_s \frac{g_{11}^s}{\omega_s^2}$$

Recap so far

- ❖ Finding the coupling (A1)
 - either by nonlinear *static* analysis: **ICE/POD** (*inverse* problem)
 - or by *direct* force computation: **SMD** (direct problem + *linear solve*)
- ❖ Including the found couplings (A2)
 - either keeping the found modes in as an additional DOF (linear ROM)
 - or *keeping* the *aggregated* vector in as an additional DOF: **dual modes** (linear ROM)
 - Or *condensing* the couplings and only keeping master: (nonlinear ROM)
- ❖ Coming up other nonlinear ROMs (A2)
 - ❖ ROM by *projection*: **QM/IC-ICE**
 - ❖ ROM by *invariance*: **SP/NF/SSM/DNF/DPIM**

(A2) Intuition on quadratic manifold (QM)

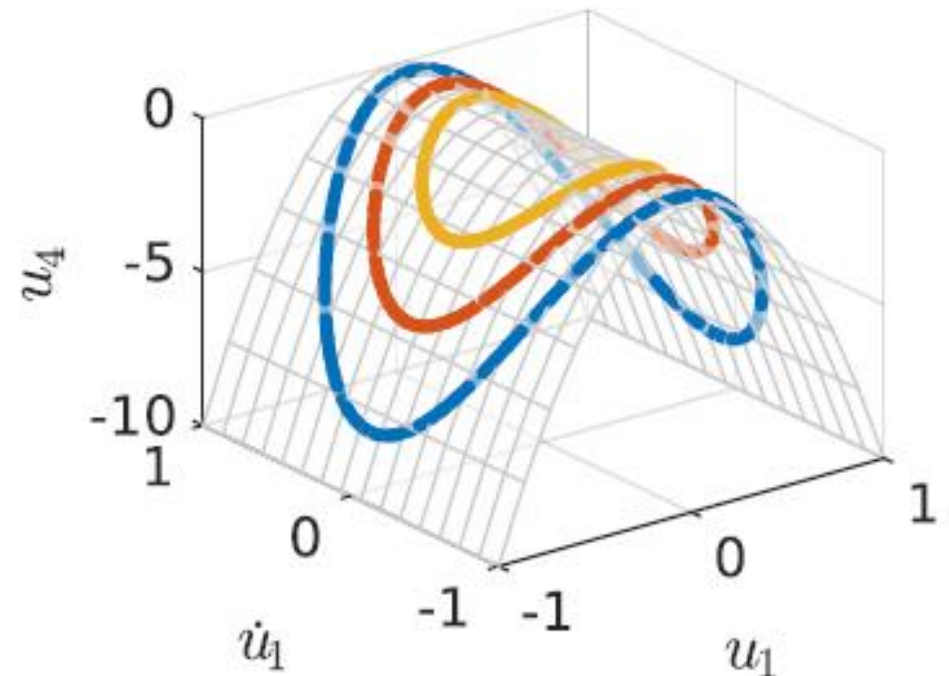
$$\theta_{11} = \frac{\partial \tilde{\phi}_1(u_1)}{\partial u_1}$$

$$\theta_{11} = - \sum 2\phi_s \frac{g_{11}^s}{\omega_s^2}$$

$$\frac{1}{2} \theta_{11} u_1^2 = - \sum \phi_s \frac{g_{11}^s}{\omega_s^2} u_1^2$$

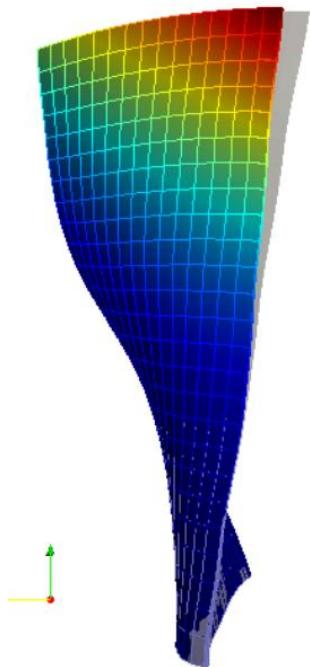
Nonlinear Change of coordinates:

$$\mathbf{U}(u_1) = \phi u_1 + \frac{1}{2} \theta u_1^2$$

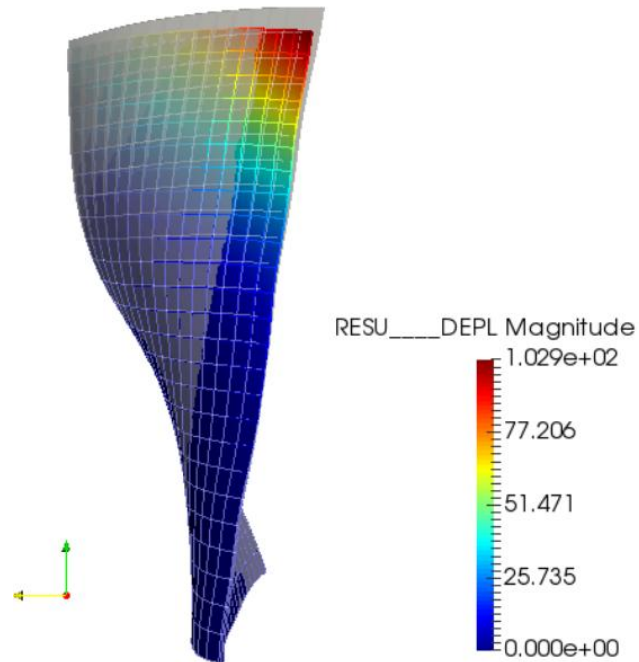


(A2) Intuition on quadratic manifold (QM)

Mode: ϕ

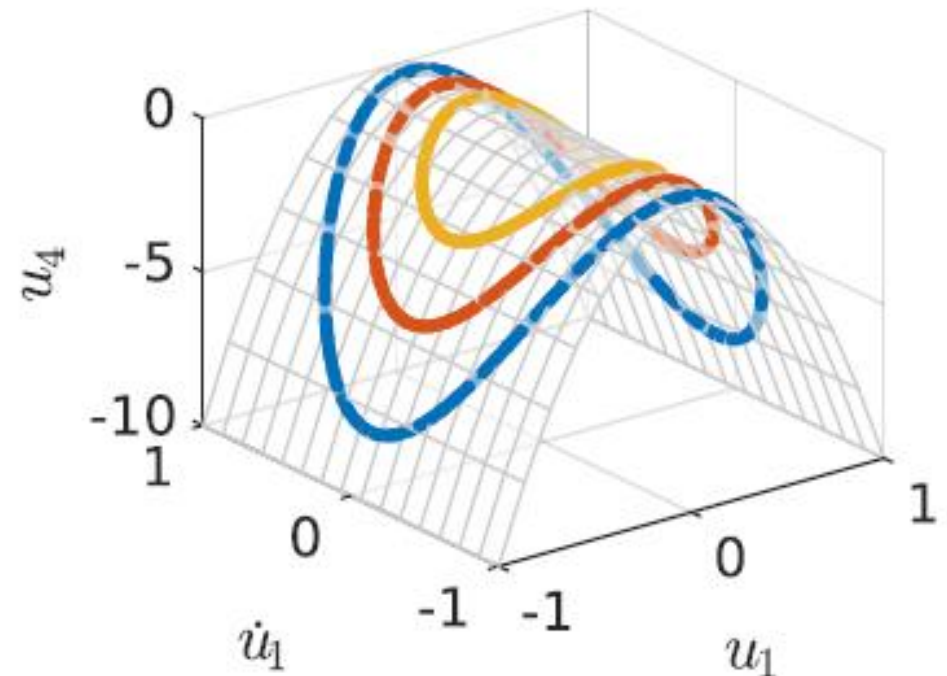


Modal derivative: θ



Nonlinear Change of coordinates:

$$\mathbf{U}(u_1) = \phi u_1 + \frac{1}{2} \theta u_1^2$$



ROM by projection (QM / IC-ICE)

Nonlinear change of coordinates

$$U(u_1) = \phi u_1 + \frac{1}{2}\theta u_1^2$$

$$\ddot{U}(u_1) = \phi \ddot{u}_1 + \frac{1}{2}\theta (\ddot{u}_1 u_1 + \dot{u}_1^2)$$

Projection on tangent space!

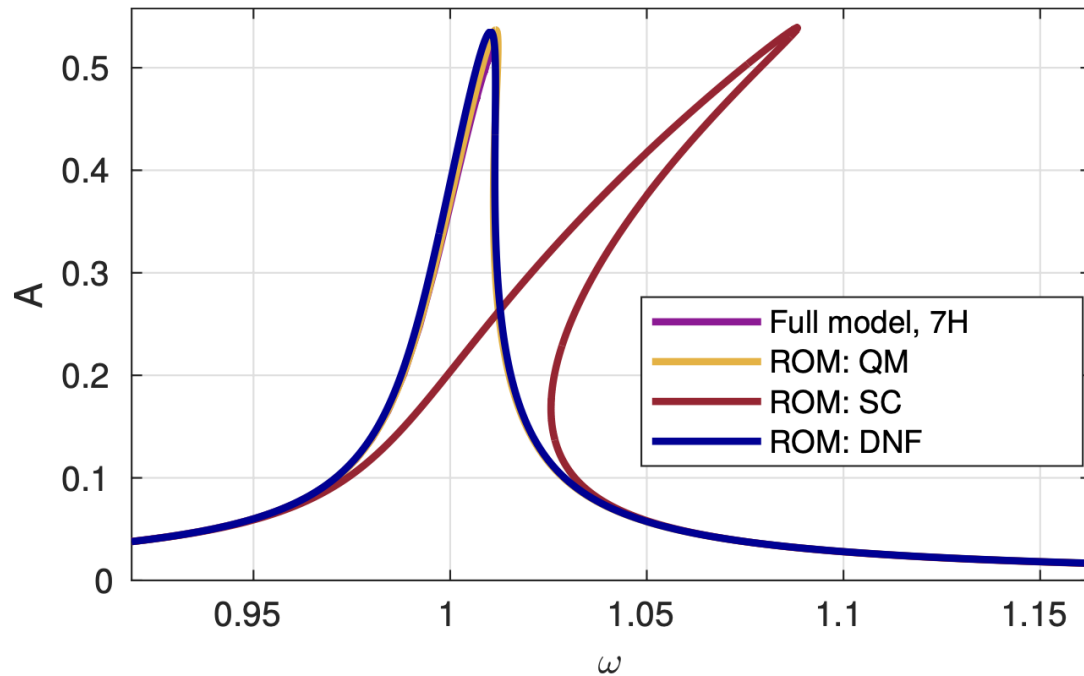
$$\nabla U(u_1) = \phi + \theta u_1$$

Reduced dynamics equation:

$$(\nabla U(u_1))^T \left(\mathbf{M} \ddot{U}(u_1) + \mathbf{K} U(u_1) + \mathbf{G} U(u_1) U(u_1) + \mathbf{H} U(u_1) U(u_1) U(u_1) \right) = \mathbf{0}$$

(A2) Inertial nonlinearity in QM

Cantilever beam



Nonlinear Change of coordinates:

$$\mathbf{U}(u_1) = \phi u_1 + \frac{1}{2} \theta u_1^2$$

nonlinear inertia component

$$\ddot{\mathbf{U}}(u_1) = \phi \ddot{u}_1 + \frac{1}{2} \theta (\ddot{u}_1 u_1 + \dot{u}_1^2)$$

SC: $\ddot{u}_1 + \omega_1^2 u_1 + (h_{111}^1 - \sum \frac{g_{1s}^1 g_{11}^s}{\omega_s^2}) u_1^3 = 0$

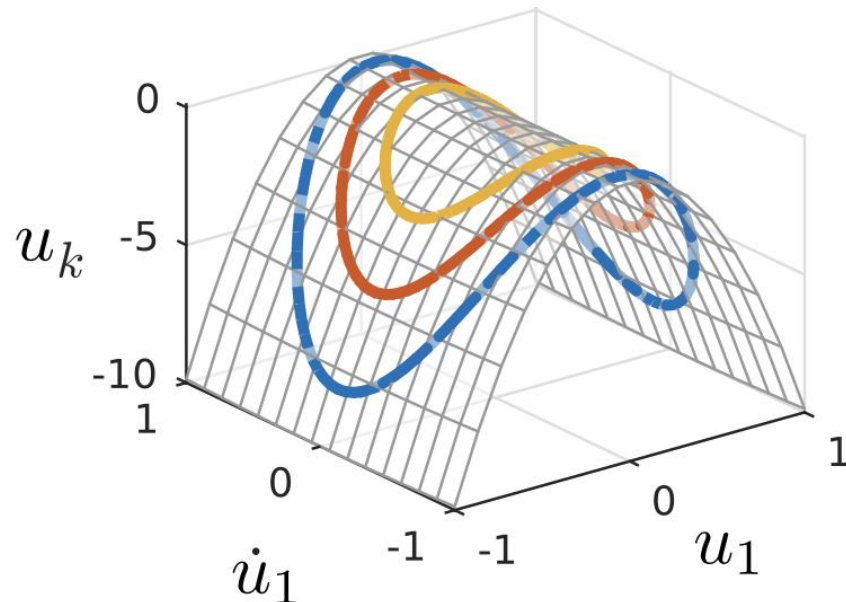
QM: $\ddot{u}_1 + \omega_1^2 u_1 + (h_{111}^1 - \sum \frac{g_{1s}^1 g_{11}^s}{\omega_s^2}) u_1^3 + \sum \frac{2g_{1s}^1 g_{11}^s}{\omega_s^4} (\ddot{u}_1 u_1^2 + u_1 \dot{u}_1^2) = 0$

Displacement only change of coordinates

Nonlinear change of coordinates

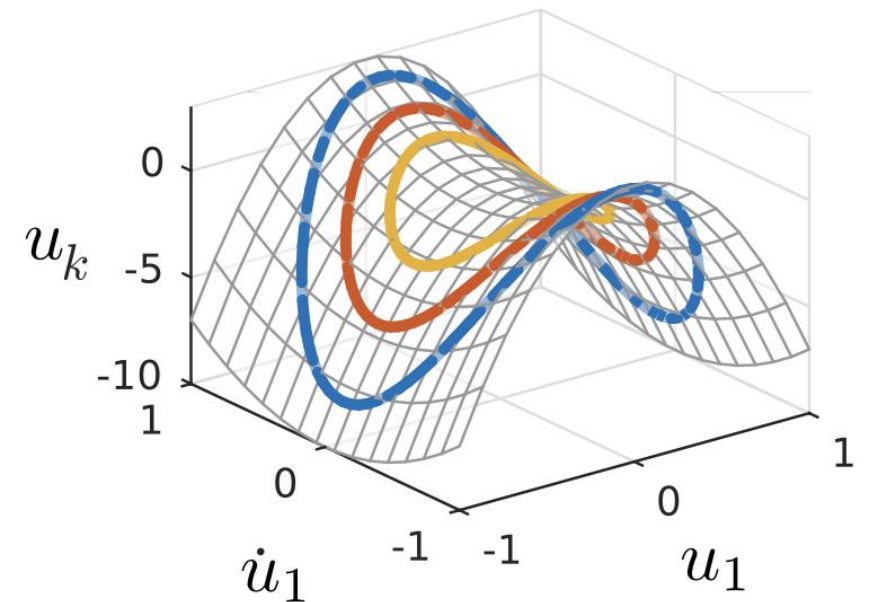
$$U(u_1) = \phi u_1 + \frac{1}{2}\theta u_1^2$$

Velocity
independent
manifold



Velocity
dependent
manifold

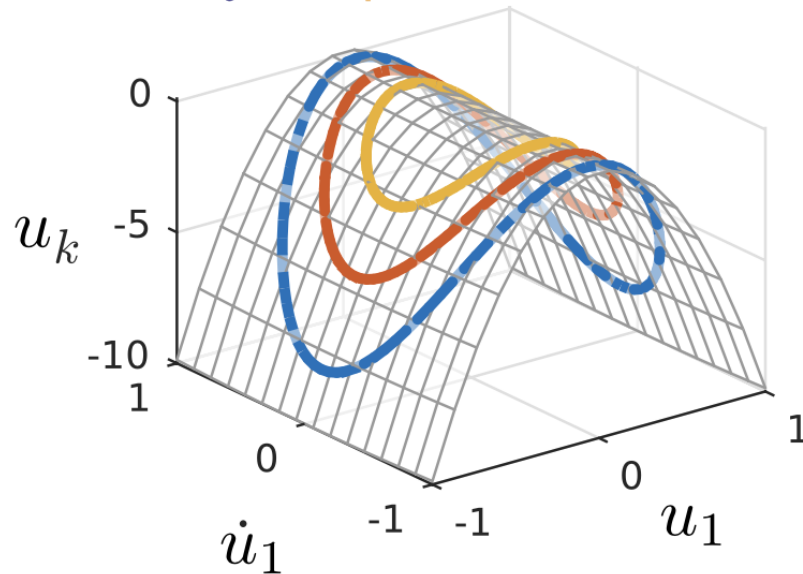
$$\mathbf{U} = \mathbf{U}(u_1, \dot{u}_1)$$



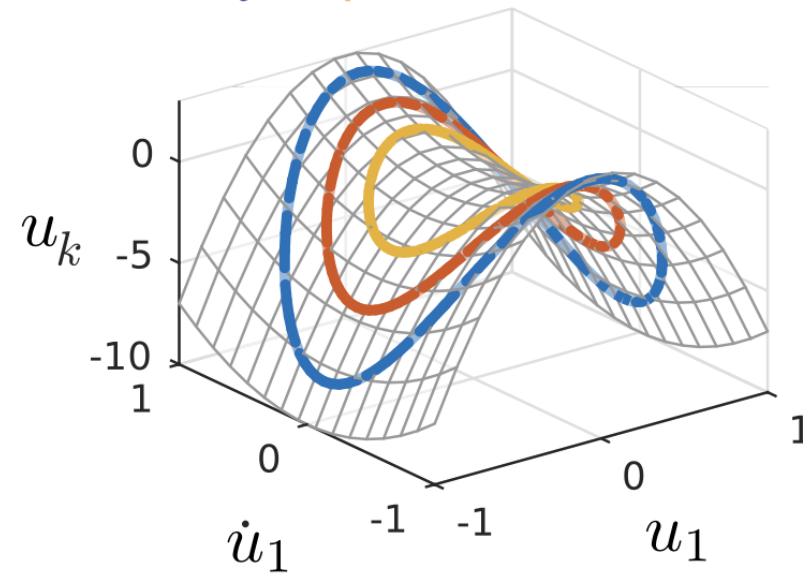
When is it needed?

When is the velocity dependence needed?

Velocity Independent Manifold



Velocity Dependent Manifold

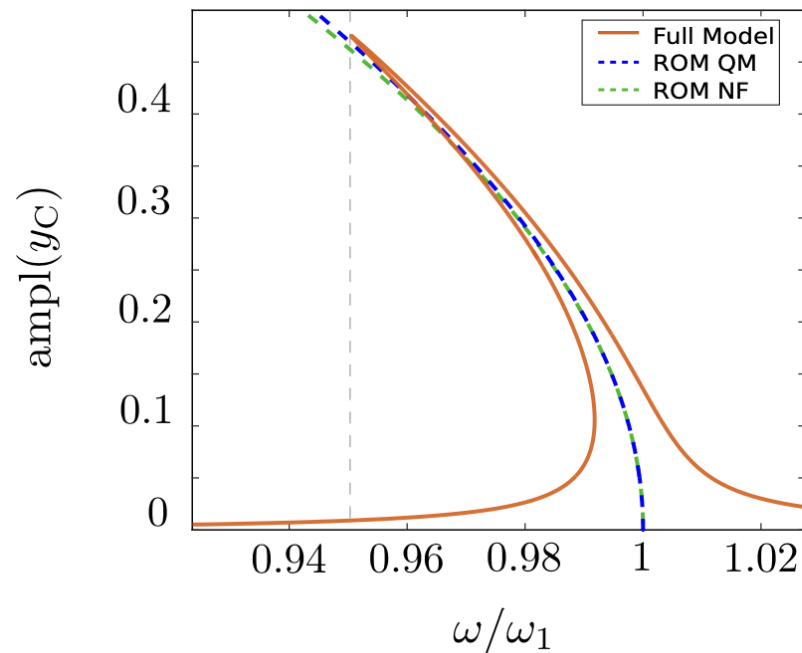
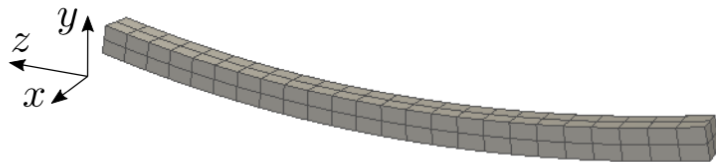


Velocity Independence: Frequency Coupled Mode \gg Frequency Master Mode

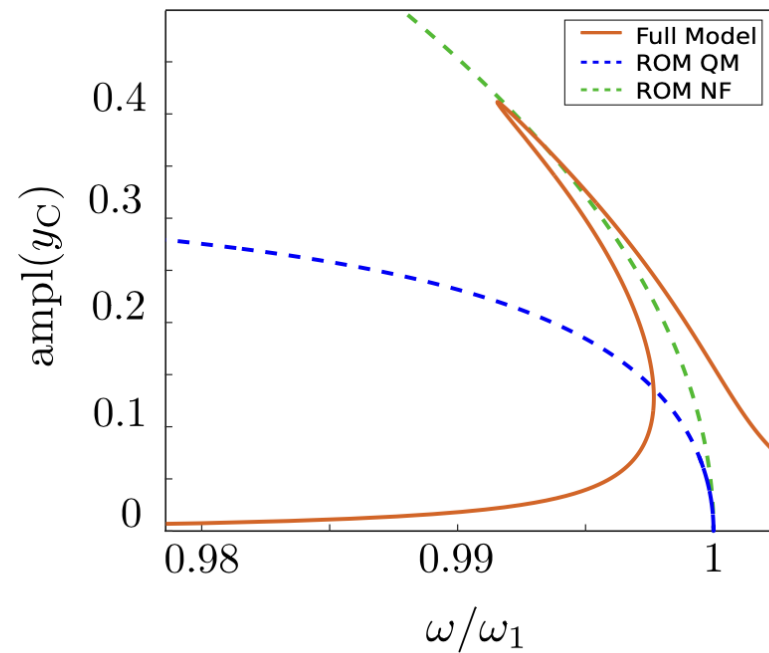
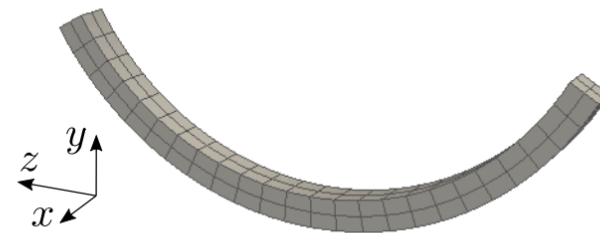
$$\omega_C > 3\omega_M$$

Example shallow VS non-shallow arch

Shallow: $\omega_C \approx 3.4\omega_M$



Non-shallow: $\omega_C \approx 1.7\omega_M$



$$\mathbf{U} = \mathbf{U}(u_1, \dot{u}_1)$$

(A2) ROM by invariance (Shaw-Pierre SP)

Start by rewriting the system in first order:

$$\dot{\mathbf{U}} = \mathbf{V}$$

$$\mathbf{M}\dot{\mathbf{V}} + \mathbf{K}\mathbf{U} + \mathbf{G}\mathbf{U}\mathbf{U} + \mathbf{H}\mathbf{U}\mathbf{U}\mathbf{U} = \mathbf{0}$$

Change of coordinates in displacement and velocity:

$$\mathbf{U} = \mathbf{U}(u_1, \mathbf{v}_1)$$

$$\mathbf{V} = \mathbf{V}(u_1, \mathbf{v}_1)$$

Ansatz for reduced model (no projection needed):

$$\dot{u}_1 = \mathbf{v}_1$$

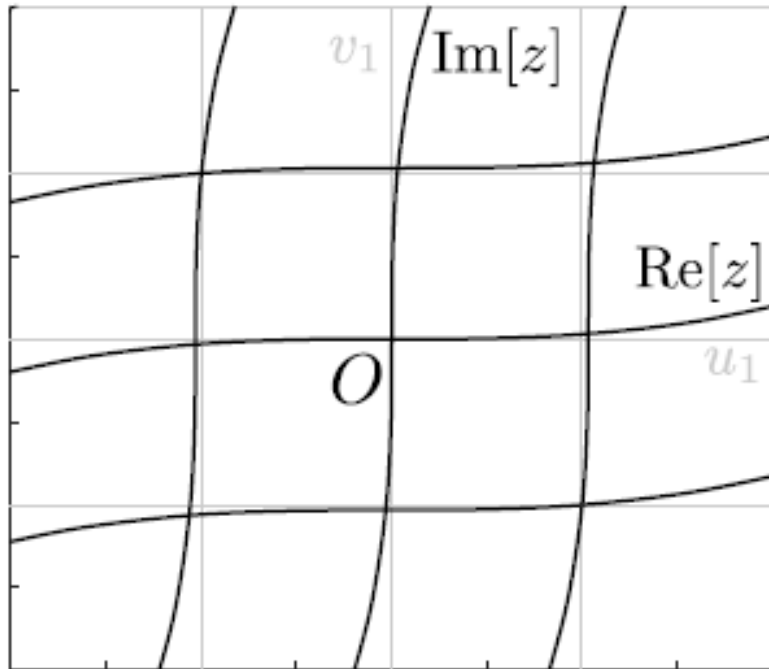
$$\dot{\mathbf{v}}_1 + f(u_1, \mathbf{v}_1) = 0$$

Time derivatives:

$$\dot{\mathbf{U}} = \partial_{u_1} \mathbf{U} \mathbf{v}_1 - \partial_{\mathbf{v}_1} \mathbf{U} f(u_1, \mathbf{v}_1)$$

$$\dot{\mathbf{V}} = \partial_{u_1} \mathbf{V} \mathbf{v}_1 - \partial_{\mathbf{v}_1} \mathbf{V} f(u_1, \mathbf{v}_1)$$

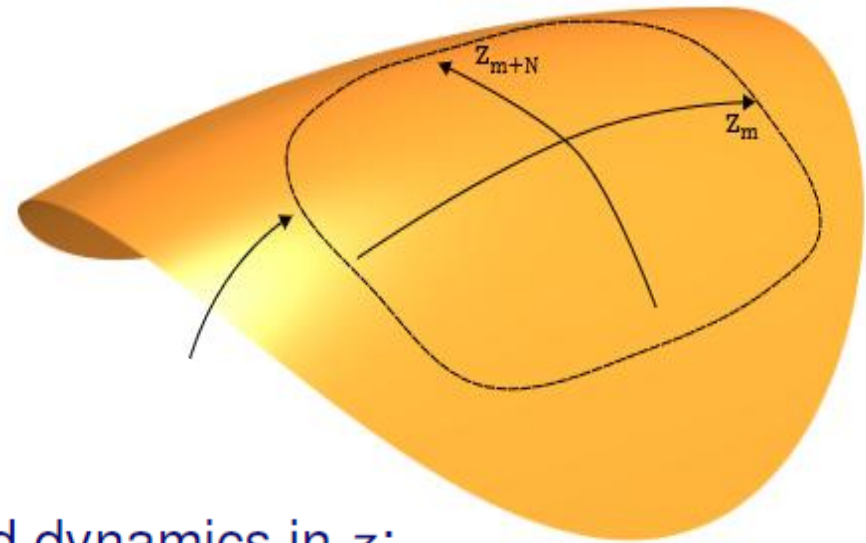
(A2) ROM by invariance (Normal Form Style)



Change of coordinates:

$$\mathbf{U} = \mathbf{U}(z)$$

Parametrisation of manifold as an embedding in the complex normal coordinates z



Reduced dynamics in z :

$$\dot{z} = f(z)$$

(A2) Invariance equation (SP/SSM/DNF/DPIM)

Change of coordinates:

$$\mathbf{U} = \mathbf{U}(z)$$

Reduced dynamics in z :

$$\dot{z} = f(z)$$

Time derivatives:

$$\dot{\mathbf{U}} = \nabla_z \mathbf{U} \cdot f;$$

$$\dot{\mathbf{V}} = \nabla_z \mathbf{V} \cdot f$$

Invariance Equation:

$$\mathbf{M} \nabla_z \mathbf{U} \cdot f = \mathbf{M} \mathbf{V}$$

$$\mathbf{M} \nabla_z \mathbf{V} \cdot f + \mathbf{K} \mathbf{U} + \mathbf{G} \mathbf{U} \mathbf{U} + \mathbf{H} \mathbf{U} \mathbf{U} \mathbf{U} = \mathbf{0}$$

Homological order p :

$$(\sigma^2 \mathbf{M} + \mathbf{K}) \psi^{(p)} = F^{(p)}$$

Homological order 2

$$\theta = (-\mathbf{K})^{-1} \mathbf{G}_{11}$$

$$\chi = ((2\omega_1)^2 \mathbf{M} - \mathbf{K})^{-1} \mathbf{G}_{11}$$

(A2) Velocity dependent manifold

SP:

$$\mathbf{U} = \phi u_1 + \mathbf{a} u_1^2 + \mathbf{b} v_1^2$$

QM:

$$\mathbf{U} = \phi u_1 + \theta u_1^2$$

Non-intrusive FE computation of \mathbf{a} , \mathbf{b} :

$$\theta = (-\mathbf{K})^{-1} \mathbf{G}_{11}$$

Standard Static Analysis

$$\chi = ((2\omega_1)^2 \mathbf{M} - \mathbf{K})^{-1} \mathbf{G}_{11}$$

Non-Standard Analysis

$$\mathbf{a} = \frac{\theta + \chi}{2}$$

$$\mathbf{b} = \frac{\theta - \chi}{2\omega_1^2}$$

Slow-fast: $\chi \approx \theta$, $\mathbf{a} \approx \theta$ $\mathbf{b} \approx \mathbf{0}$

Recap

Starting from the equation of motion:

$$\mathbf{M}\ddot{\mathbf{U}} + \mathbf{K}\mathbf{U} + \mathbf{G}\mathbf{U}\mathbf{U} + \mathbf{H}\mathbf{U}\mathbf{U}\mathbf{U} = \mathbf{0} \quad (4)$$

Change of coordinates

Modal ROM:

$$u_s = \mathbf{a}_s^{(2)} u_1^2 + \dots$$

Direct ROM:

$$\mathbf{U} = \phi_1 u_1 + \boldsymbol{\psi}^{(2)} u_1^2 + \dots$$

Sought **geometry** of the approximated manifold is given by the unknown $\mathbf{a}_s^{(2)}$ or $\boldsymbol{\psi}^{(2)}$

Reduced dynamics

$$\ddot{u}_1 + \omega_1^2 u_1 + \mathbf{f}^{(2)} u_1^2 + \dots = 0$$

Sought **dynamics** on the approximated manifold is given by the unknown $\mathbf{f}^{(2)}$

SC/ICE

Assume polynomial form in u_1 :

$$\mathbf{U} = \phi_1 u_1 + \boldsymbol{\psi}^{(2)} u_1^2 + \dots$$

Neglect inertia and balance high order terms:

$$\mathbf{K}\boldsymbol{\psi}^{(2)} u_1^2 + \mathbf{G}\phi_1 \phi_1 u_1^2 = \mathbf{0}$$

$$\vdots$$

Obtain reduced dynamics by projection of static forces:

$$\begin{aligned} \ddot{u}_1 + \phi_1^T (\mathbf{K}\mathbf{U} + \mathbf{G}\mathbf{U}\mathbf{U} + \mathbf{H}\mathbf{U}\mathbf{U}\mathbf{U}) = \\ \ddot{u}_1 + \omega_1^2 u_1 + \boldsymbol{f}^{(2)} u_1^2 + \boldsymbol{f}^{(3)} u_1^3 + \dots = 0 \end{aligned}$$

ICE method works similarly but in an implicit manner: easy to go to higher orders.

QM/IC-ICE

Assume polynomial form in u_1 :

$$\mathbf{U} = \phi_1 u_1 + \boldsymbol{\psi}^{(2)} u_1^2 + \dots$$

Neglect inertia and balance high order terms:

$$\mathbf{K} \boldsymbol{\psi}^{(2)} u_1^2 + \mathbf{G} \phi_1 \phi_1 u_1^2 = \mathbf{0}$$

$$\vdots$$

Obtain reduced dynamics by **projection of whole e.o.m. on a tangent space**:

$$\begin{aligned} (\phi_1^\top + u_1 \boldsymbol{\psi}^{(2)\top})(\mathbf{M}\ddot{\mathbf{U}} + \mathbf{K}\mathbf{U} + \mathbf{G}\mathbf{U}\mathbf{U} + \mathbf{H}\mathbf{U}\mathbf{U}\mathbf{U}) = \\ \ddot{u}_1 + \omega_1^2 u_1 + \boldsymbol{f}_{NL}(u_1, \dot{u}_1, \ddot{u}_1) = 0 \end{aligned}$$

Appearance of **inertial terms** in acceleration and velocity

SP/SSM/DNF/DPIM

Assume polynomial form in u_1 and \dot{u}_1 :

$$\mathbf{U} = \phi_1 u_1 + \psi^{(2,0)} u_1^2 + \psi^{(1,1)} u_1 \dot{u}_1 + \psi^{(0,2)} \dot{u}_1^2 + \dots$$

Assume reduced dynamics too: polynomial in u_1 and \dot{u}_1 :

$$\ddot{u}_1 + \omega_1^2 u_1 + f^{(2,0)} u_1^2 + f^{(1,1)} u_1 \dot{u}_1 + f^{(0,2)} \dot{u}_1^2 + \dots = 0$$

Balance same monomials in invariance equation:

$$\begin{aligned} \mathbf{M} \nabla_z \mathbf{U} \cdot \mathbf{f} &= \mathbf{M} \mathbf{V} \\ \mathbf{M} \nabla_z \mathbf{V} \cdot \mathbf{f} + \mathbf{K} \mathbf{U} + \mathbf{G} \mathbf{U} \mathbf{U} + \mathbf{H} \mathbf{U} \mathbf{U} \mathbf{U} &= 0 \end{aligned}$$

No need to enforce projection!

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Thank you for your attention