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Explicit upwind scheme for $f_{e,i}$:

$$\frac{f_{s,k,l}^{n+1} - f_{s,k,l}^n}{\Delta t} + \left(\frac{v_l}{c_s E_k^n} \right)_+ D_{k,l}^- f_s^n + \left(\frac{v_l}{c_s E_k^n} \right)_- D_{k,l}^+ f_s^n = S_{s,k,l}^n$$

with $S_e = 0$, $S_i = \nu f_e$, and $D_{k,l}^\pm f := \pm \left(\frac{f_{k\pm 1,l} - f_{k,l}}{\Delta x}, \frac{f_{k,l\pm 1} - f_{k,l}}{\Delta v} \right)^t$.

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$$E_k^{n+1} = \frac{1}{\lambda^2} \sum_{\kappa} \frac{1}{\Delta x} (n_{i,\kappa}^n - n_{e,\kappa}^n)$$

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- CFL condition to ensure the stability of the explicit scheme.
- First-order approximation, diffusive.

Semi-Lagrangian scheme

Strang splitting :

$$\frac{\Delta t}{2} \begin{cases} \partial_t f_s + v \partial_x f_s = 0 & \text{Linear advection at constant speed} \\ \lambda^2 \partial_x E = n_i - n_e & \text{Resolution by (numerical) integration} \\ \partial_t f_i = \nu f_e & \text{Pointwise ODE} \end{cases}$$

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- The analytical solution to each step is known.
- Use of the 1D solver to enforce boundary conditions.

Fixed-point algorithm

Equilibrium state :

$$\begin{cases} v\partial_x f_s - c_s\partial_x\varphi_s\partial_v f_s = S_s, & \frac{d}{d\tau}[f_s(x_s(\tau), v_s(\tau))] = S_s(x_s(\tau), v_s(\tau)) \\ -\lambda^2\partial_{xx}^2\varphi = n_i - n_e \end{cases}$$

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Suppose φ^k is known. Iteration of

$$f_e^{k+1}(x, v) \quad := \quad f_{e,b}(x_e(-\tau), v_e(-\tau))$$

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$$\begin{aligned} f_e^{k+1}(x, v) &:= f_{e,b}(x_e(-\tau), v_e(-\tau)) \\ f_i^{k+1}(x, v) &:= f_{i,b}(x_i(-\tau), v_i(-\tau)) + \int_{-\tau}^0 \nu f_e^{k+1}(x_i(s), v_i(s)) ds \end{aligned}$$

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$$-\lambda^2\partial_{xx}^2\varphi^{k+1} := \int_v [f_i^{k+1}(\cdot, v) - f_e^{k+1}(\cdot, v)] dv$$

Merci.