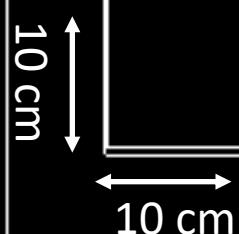
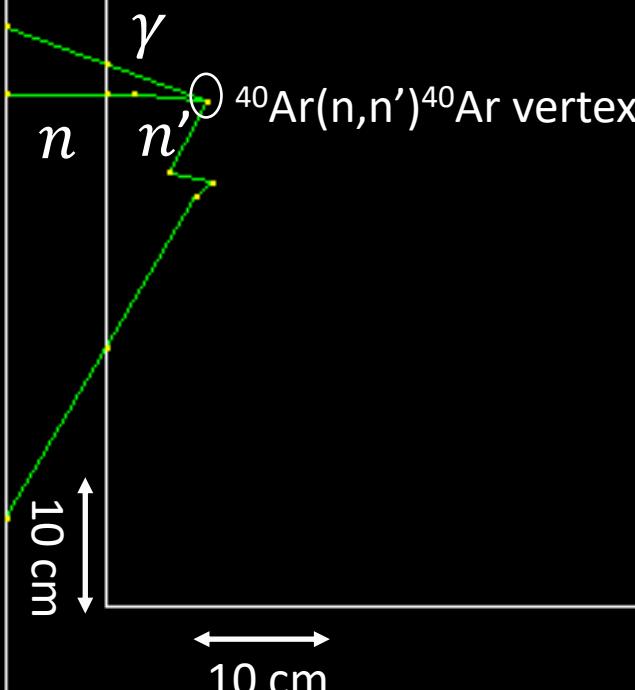


# Examples of 2.5 MeV neutron interaction sequence in LArTPC

$n$  enters the detector, scatters elastically multiple times and exits the detector.

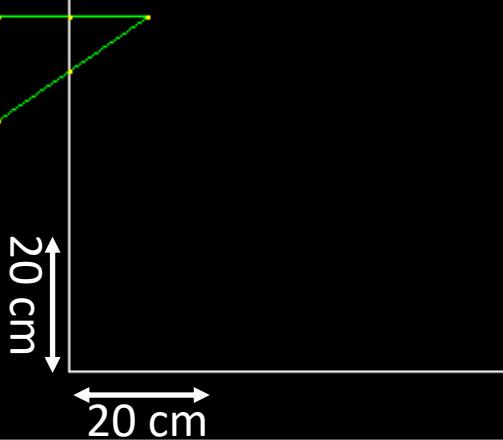


$n$  captured on Ar and performs  $^{40}\text{Ar}(n,n')^{40}\text{Ar}$  process.

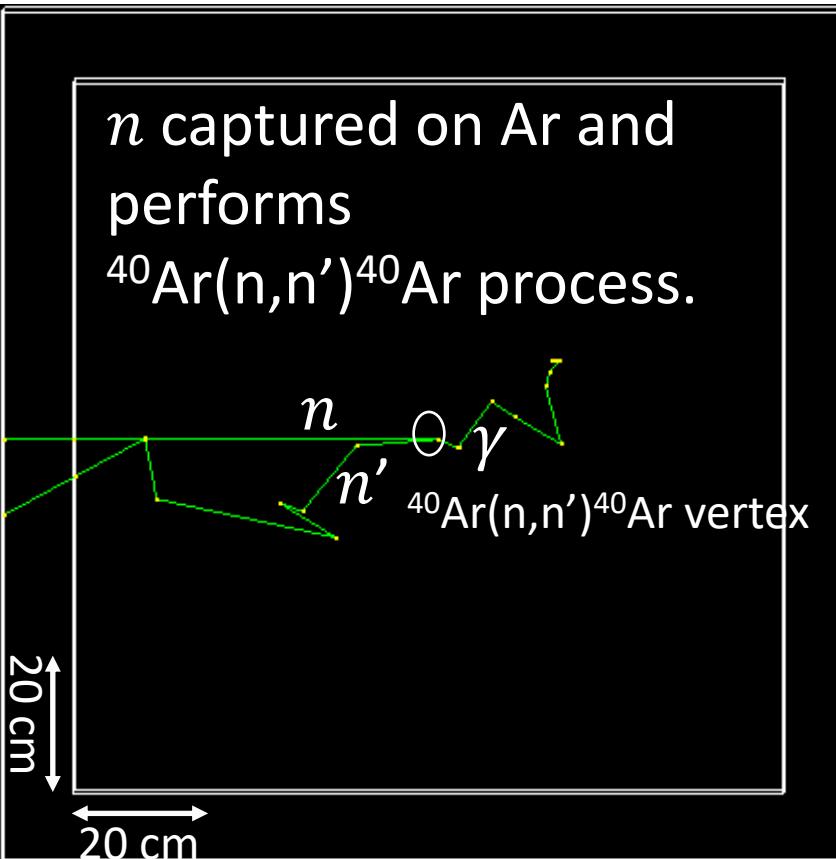


# Examples of 2.5 MeV neutron interaction sequence in LArTPC

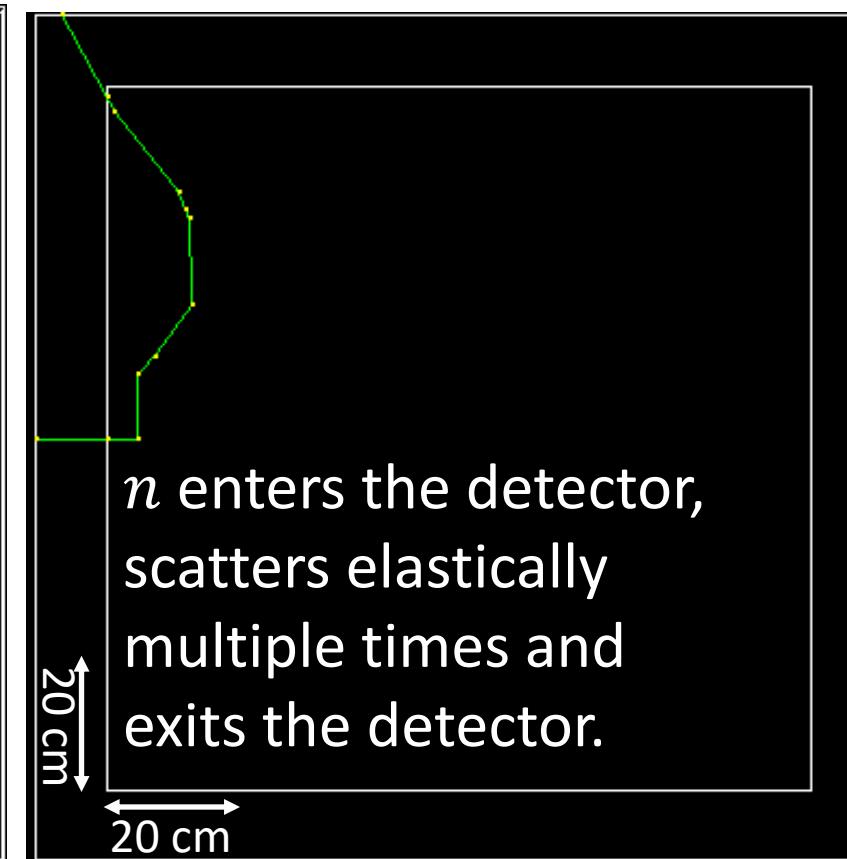
$n$  enters the detector,  
scatters elastically and  
exits the detector.



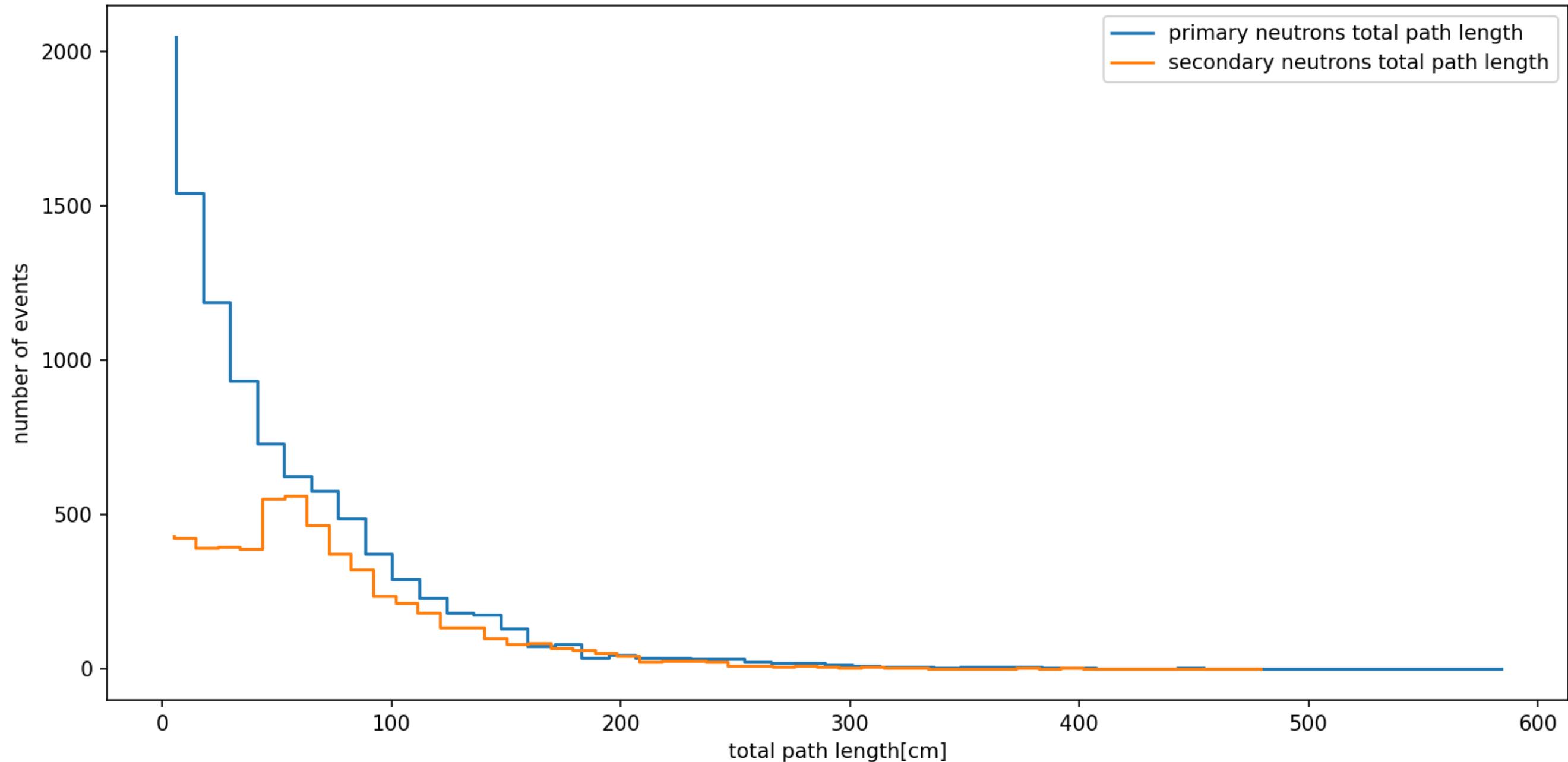
$n$  captured on Ar and  
performs  
 $^{40}\text{Ar}(n,n')^{40}\text{Ar}$  process.



$n$  enters the detector,  
scatters elastically  
multiple times and  
exits the detector.

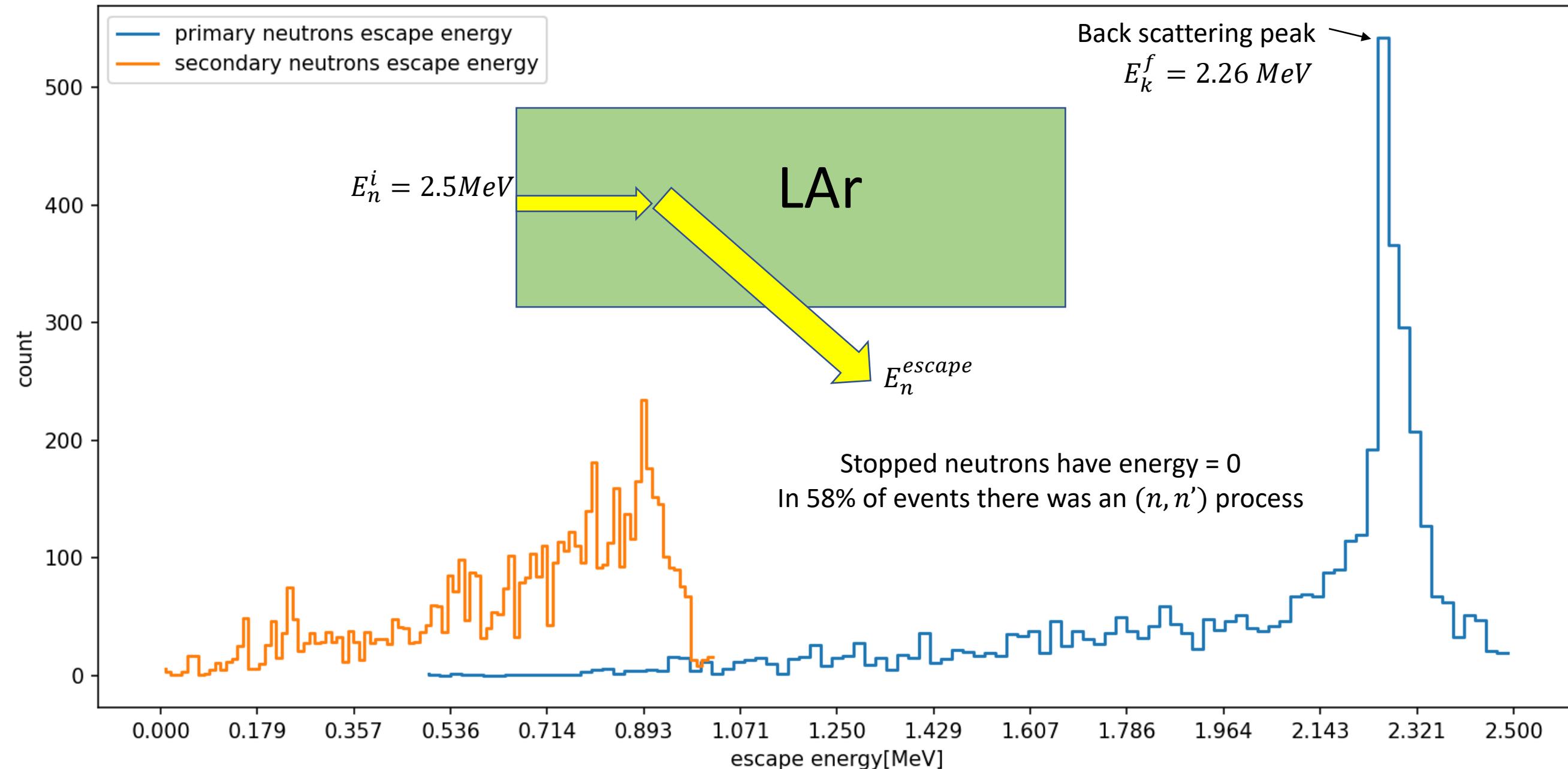


### neutrons total path length

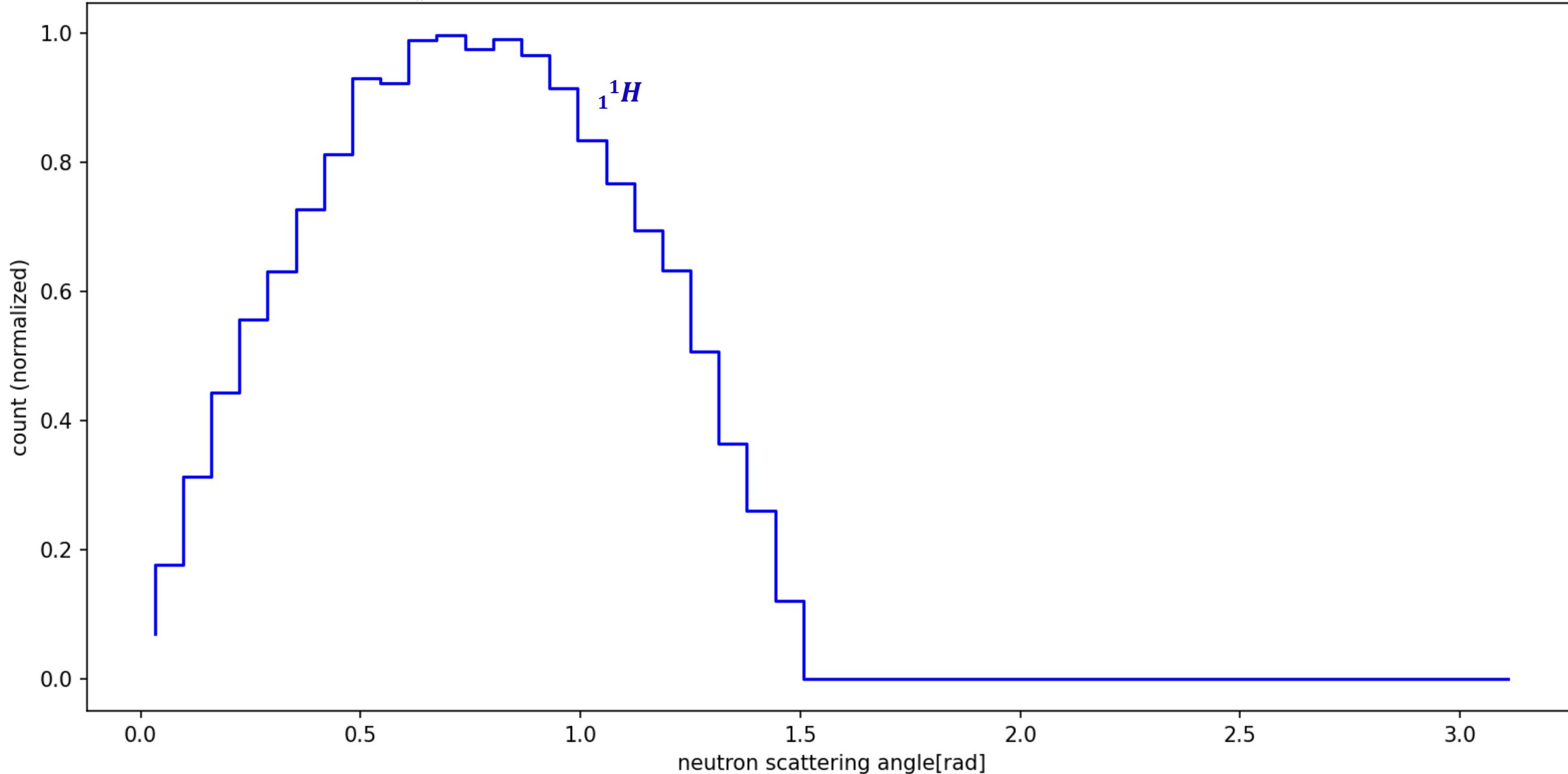


## neutron escape energy distribution

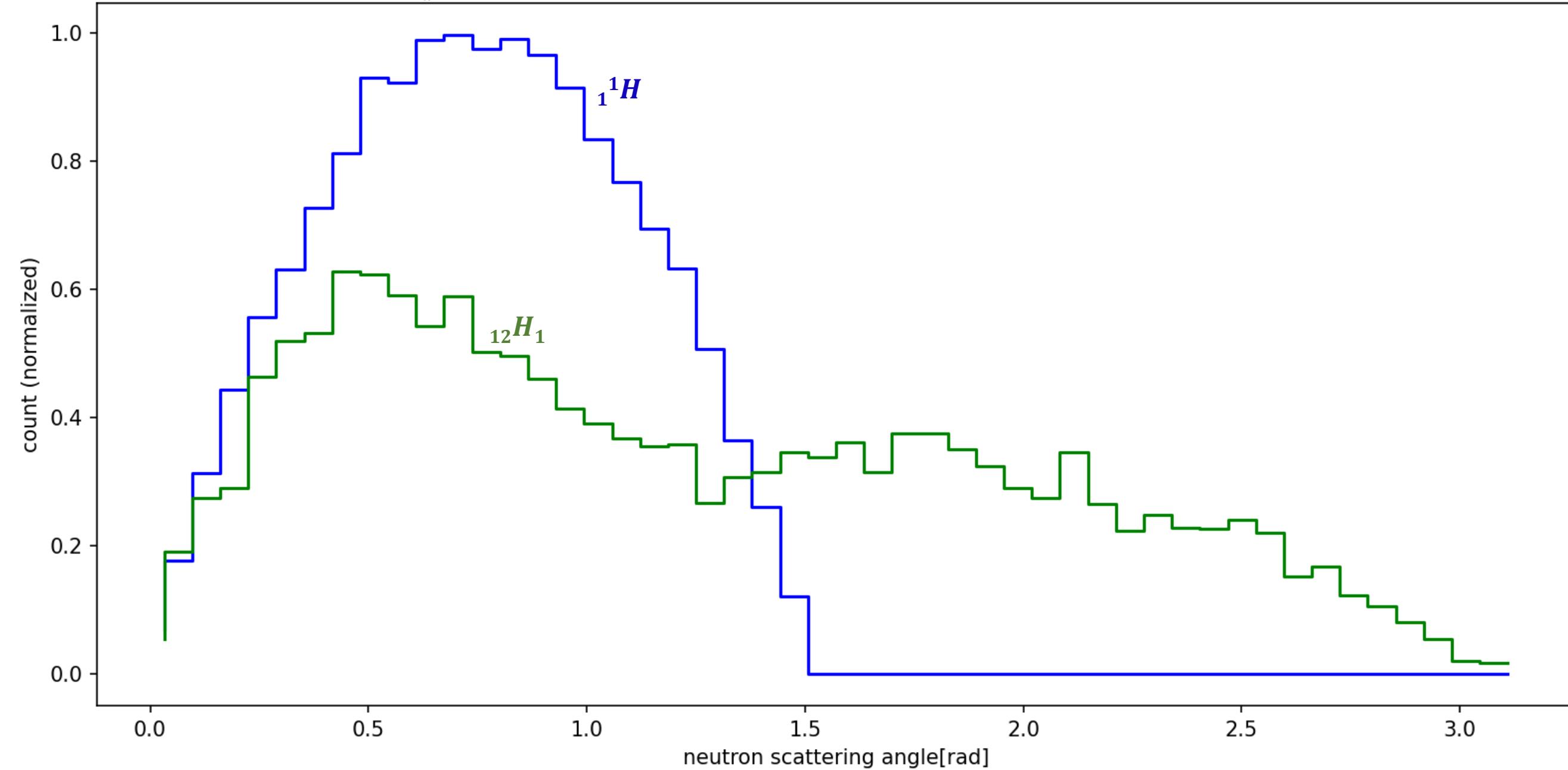
Detector size 1x1x1 meter



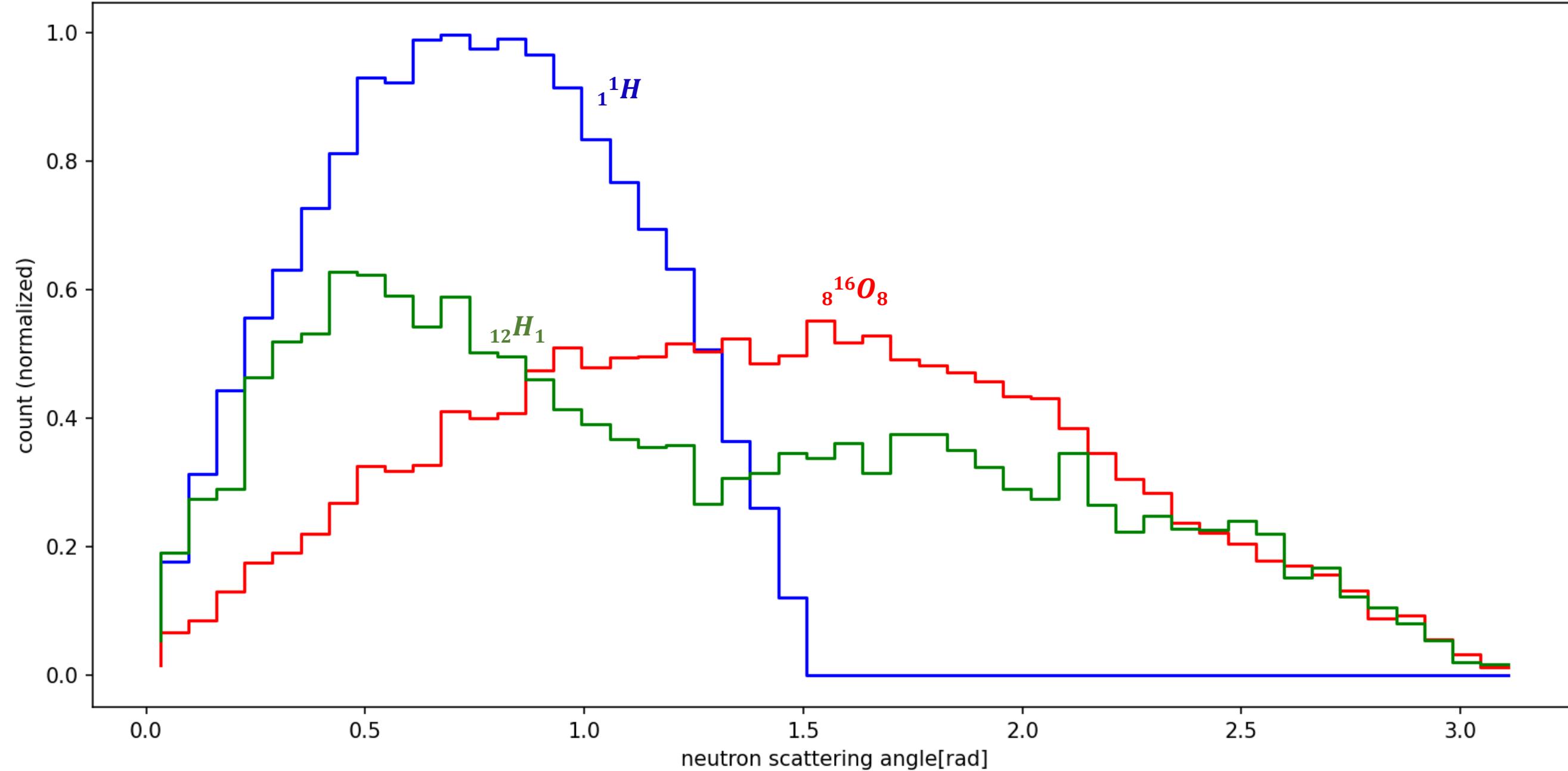
$$\Delta E = E_i \cdot \frac{4Mm_n}{(M+m_n)^2} \cdot \cos^2(\beta) - \text{neutron elastic scattering angle distribution for different materials}$$



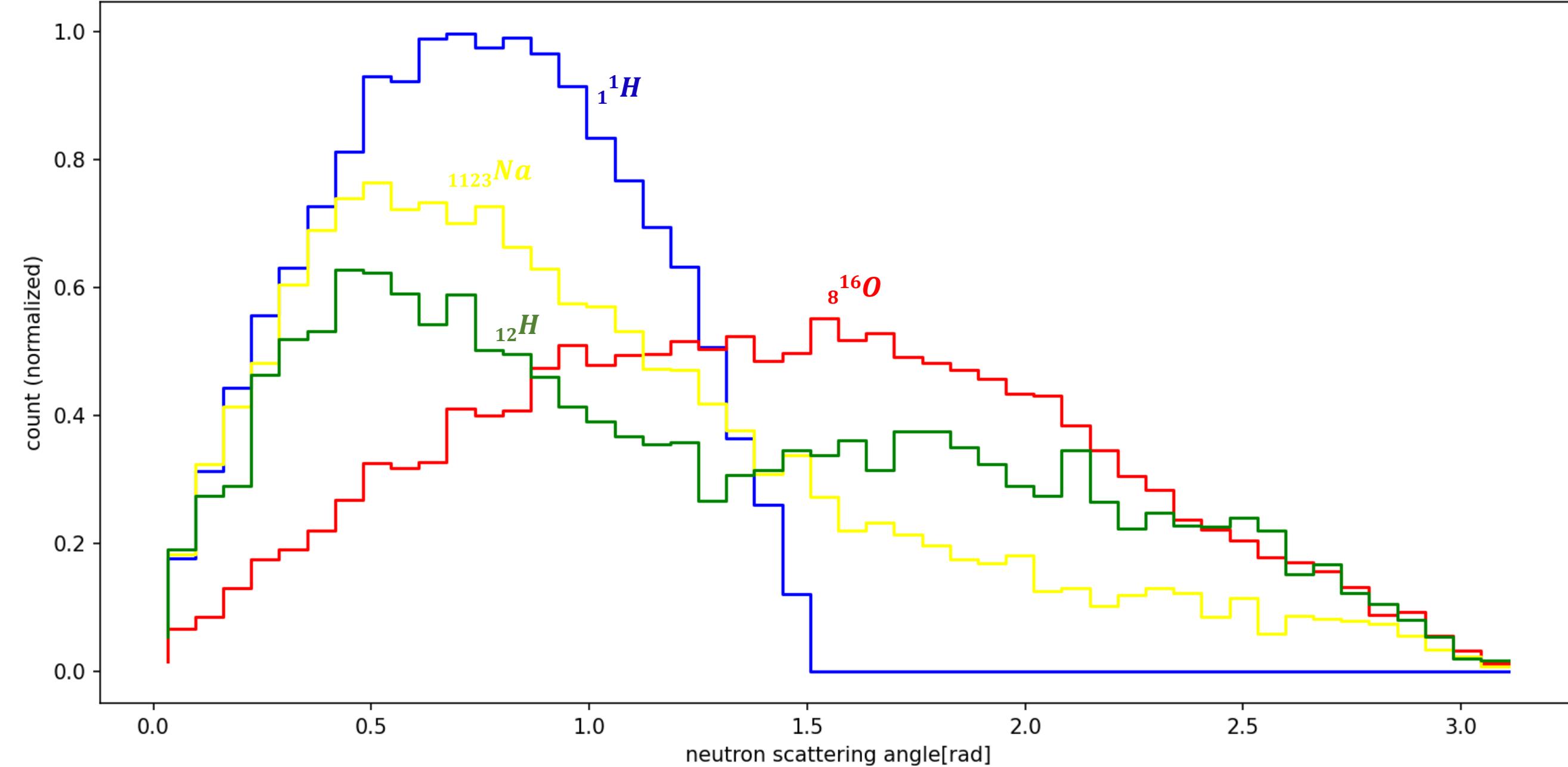
$$\Delta E = E_i \cdot \frac{4Mm_n}{(M+m_n)^2} \cdot \cos^2(\beta) - \text{neutron elastic scattering angle distribution for different materials}$$



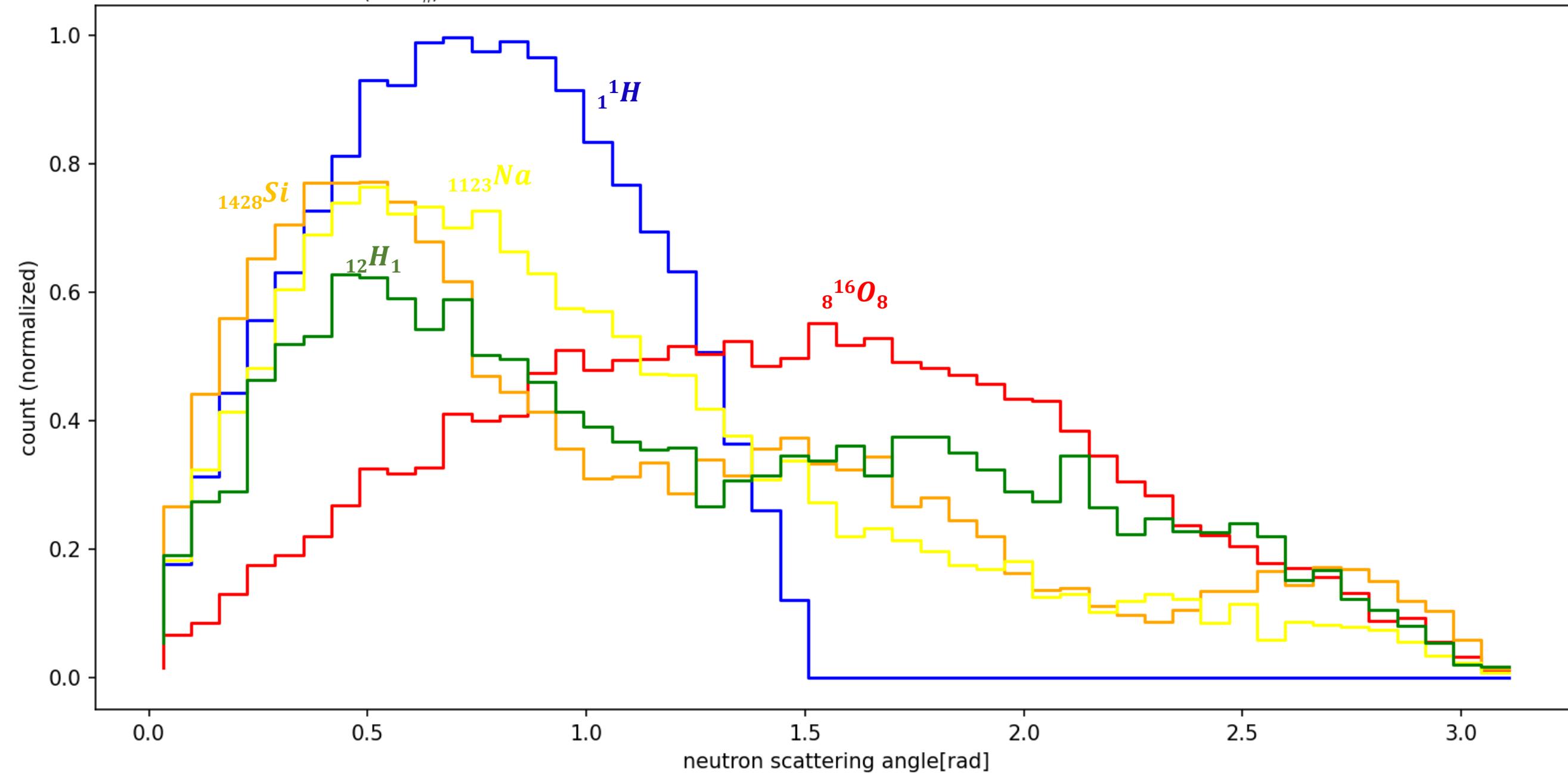
$$\Delta E = E_i \cdot \frac{4Mm_n}{(M + m_n)^2} \cdot \cos^2(\beta) - \text{neutron elastic scattering angle distribution for different materials}$$



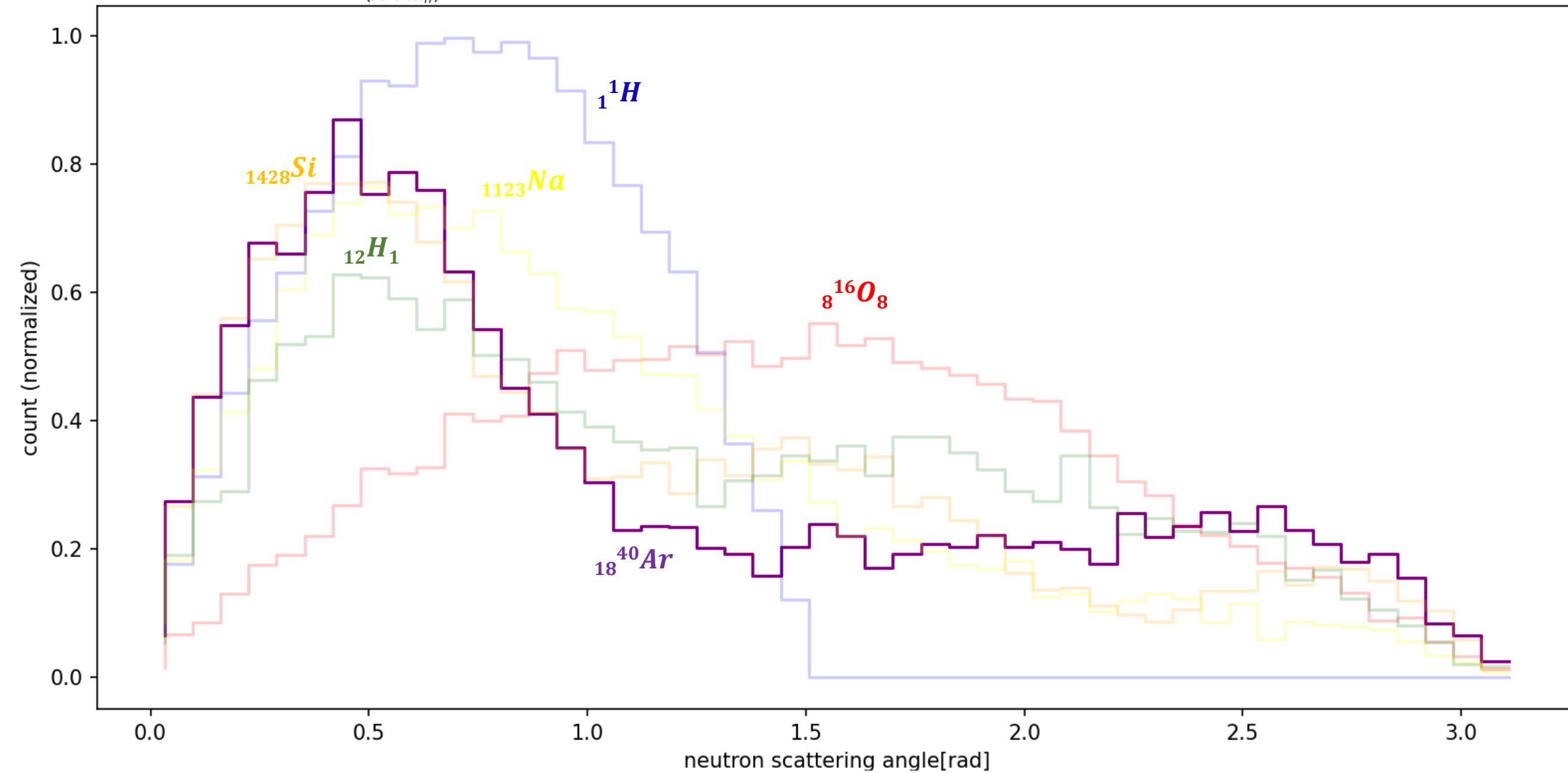
$$\Delta E = E_i \cdot \frac{4Mm_n}{(M+m_n)^2} \cdot \cos^2(\beta) - \text{neutron elastic scattering angle distribution for different materials}$$



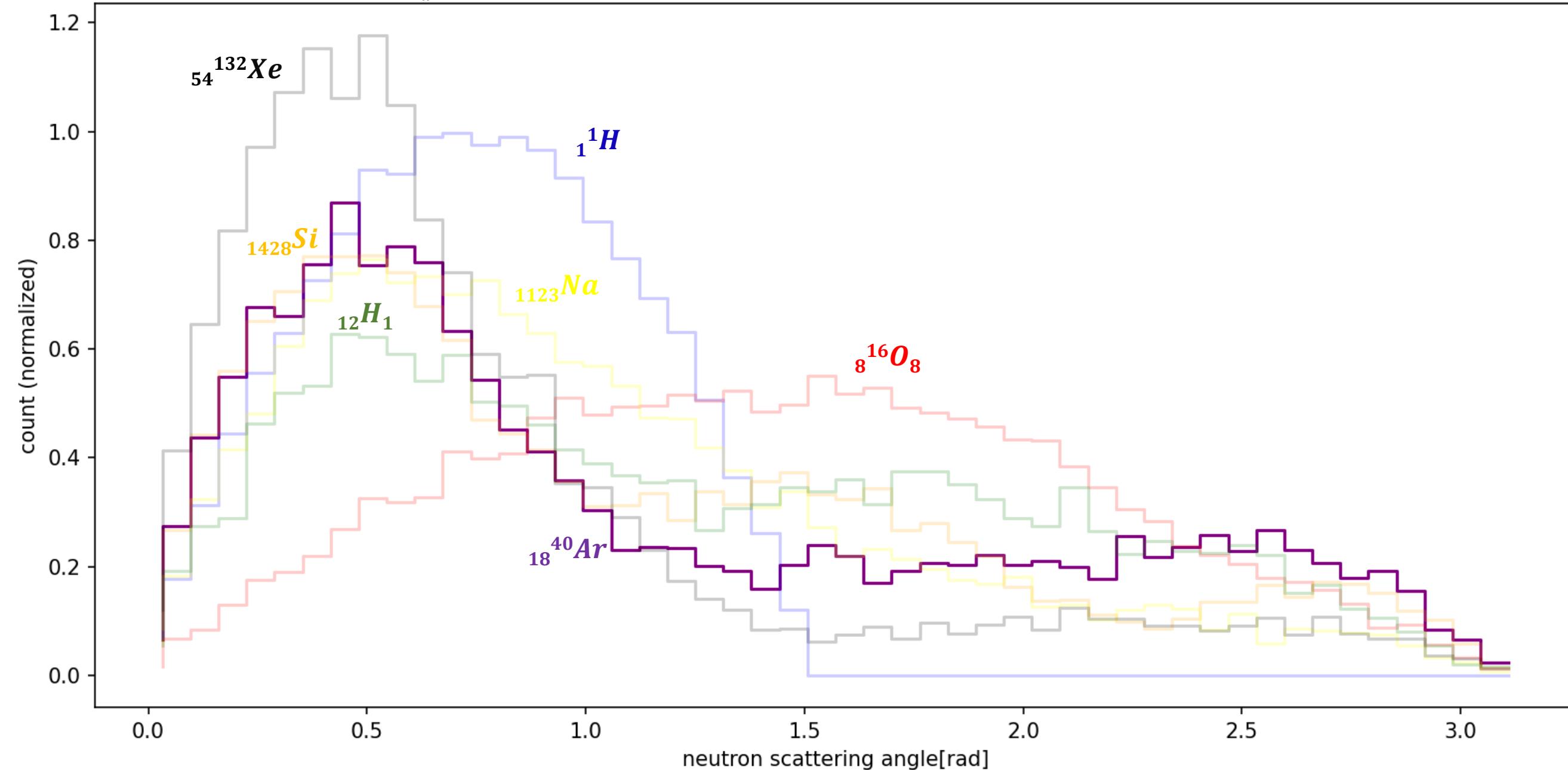
$$\Delta E = E_i \cdot \frac{4Mm_n}{(M+m_n)^2} \cdot \cos^2(\beta) - \text{neutron elastic scattering angle distribution for different materials}$$



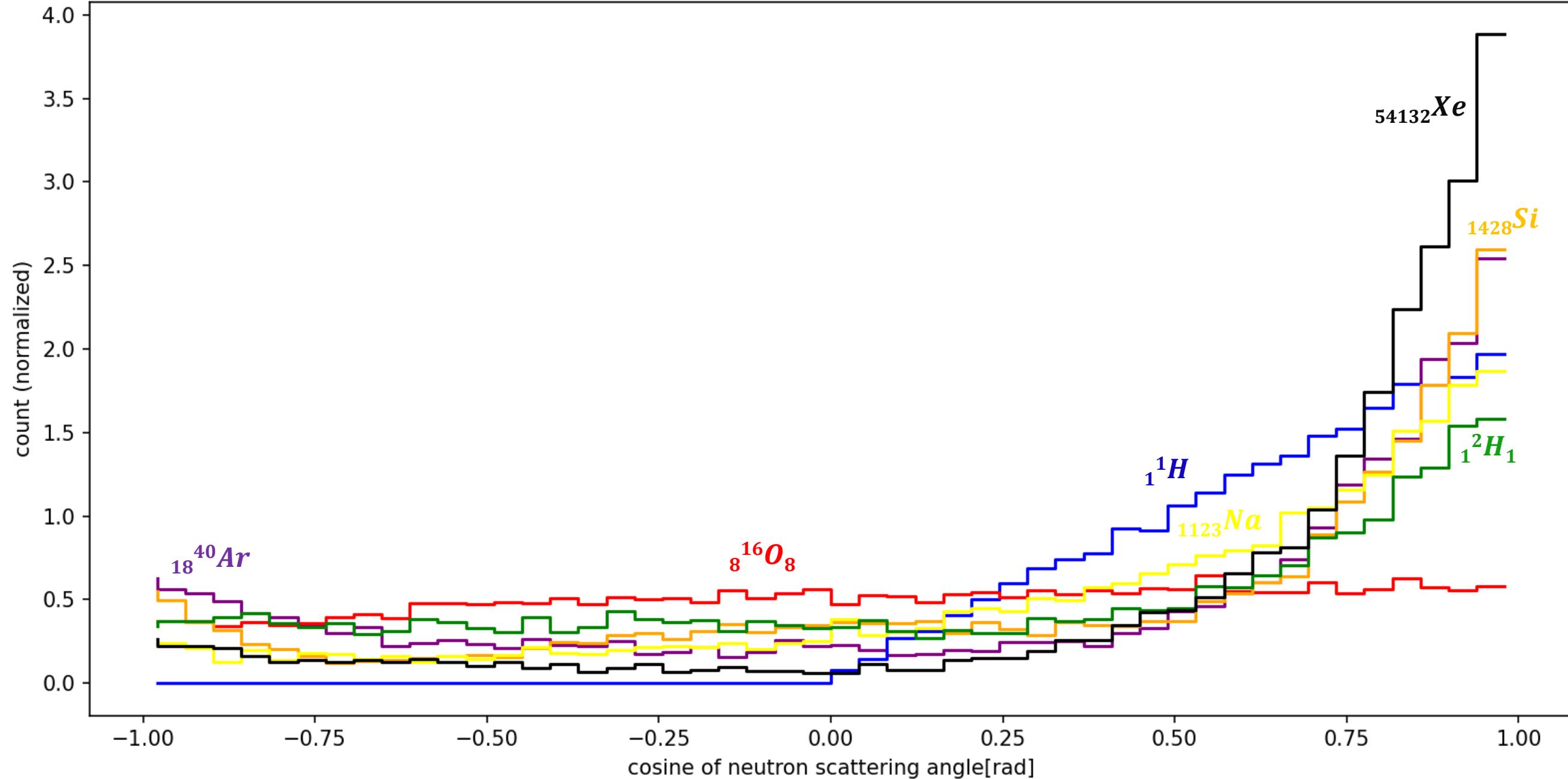
$$\Delta E = E_i \cdot \frac{4Mm_n}{(M+m_n)^2} \cdot \cos^2(\beta) - \text{neutron elastic scattering angle distribution for different materials}$$



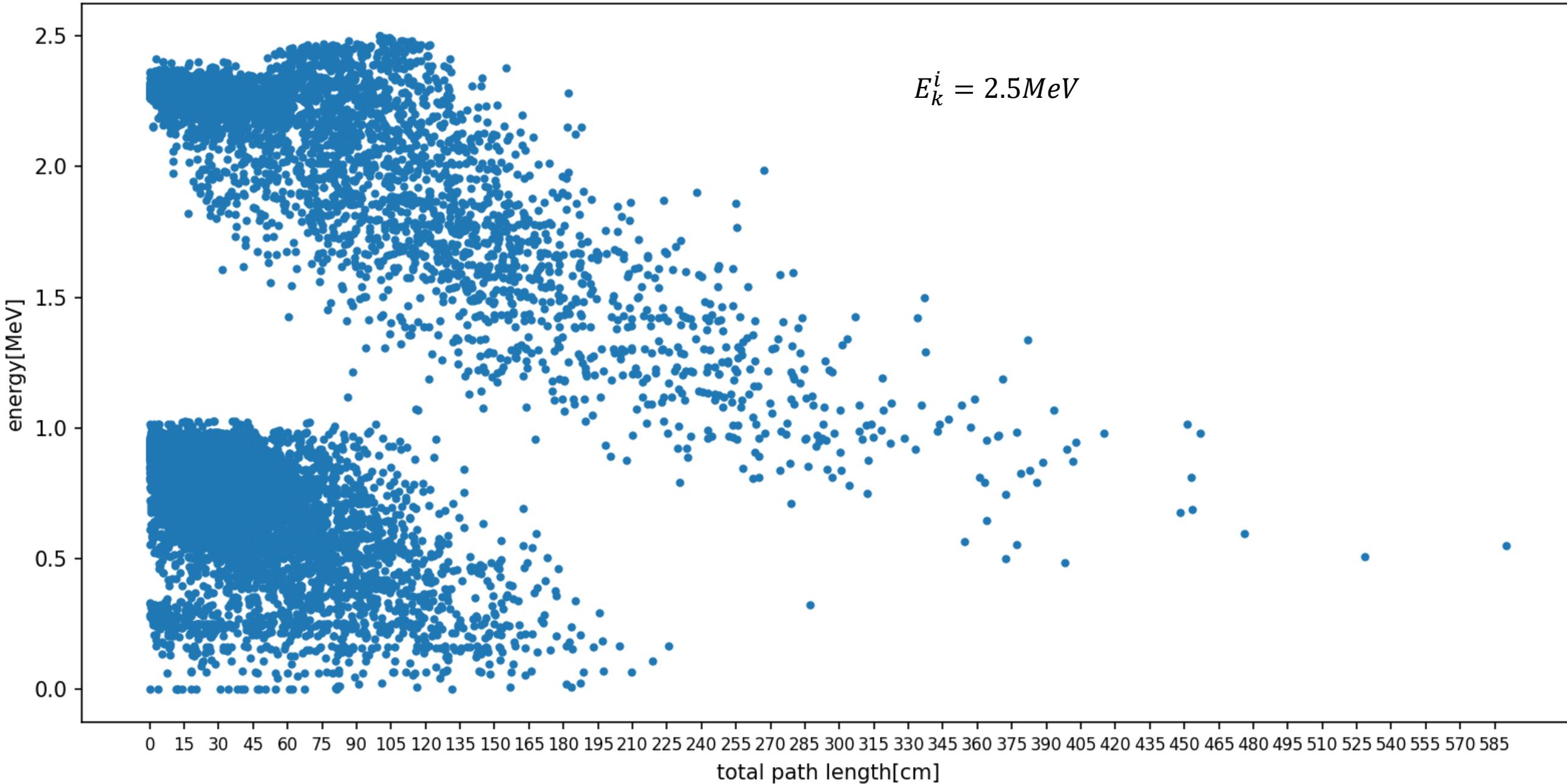
$$\Delta E = E_i \cdot \frac{4Mm_n}{(M+m_n)^2} \cdot \cos^2(\beta) - \text{neutron elastic scattering angle distribution for different materials}$$



$$\frac{4Mm_n}{(M+m_n)^2} \cdot \cos^2(\beta) - \text{neutron elastic scattering angle distribution for different materials}$$

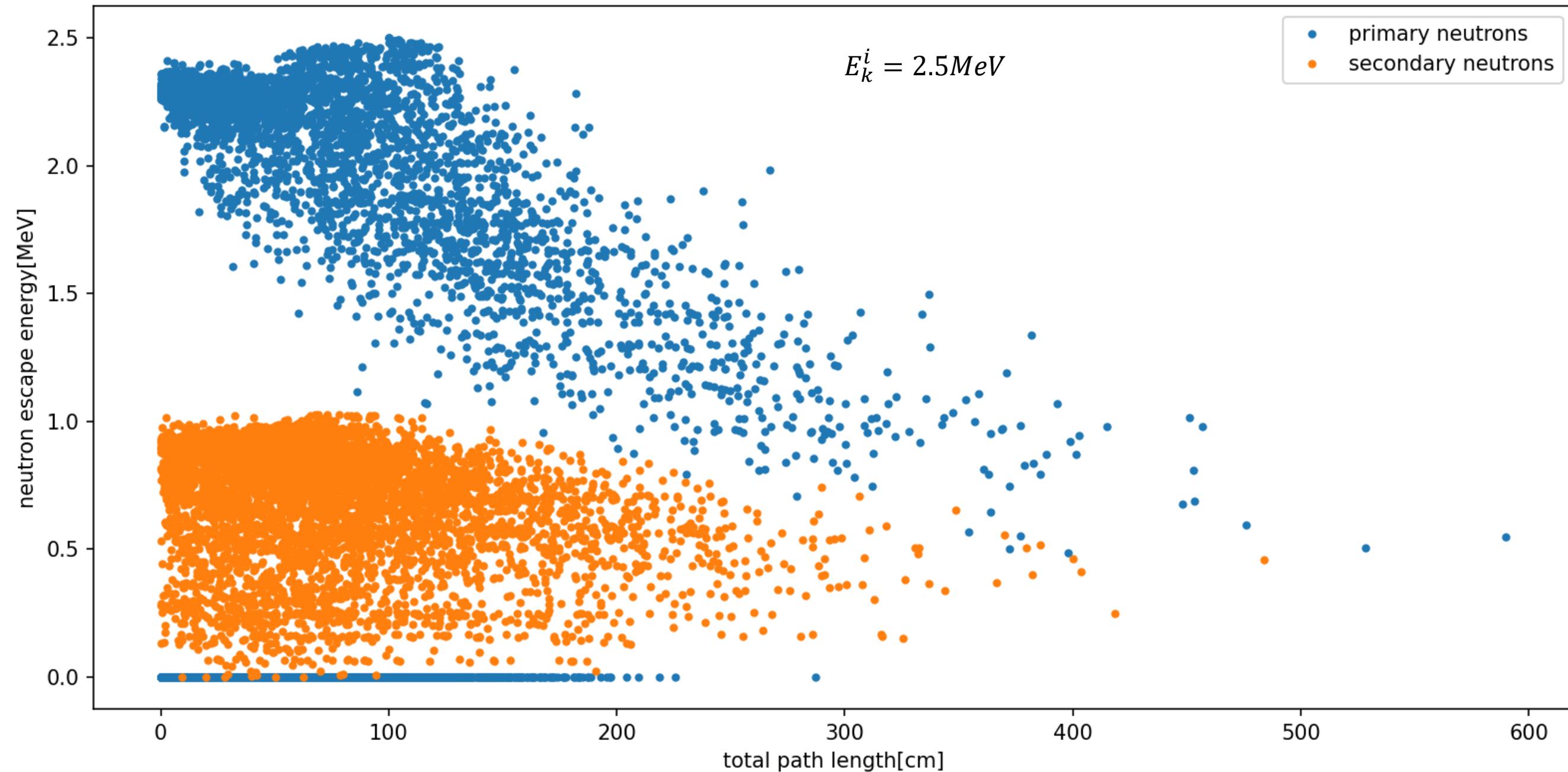


neutron escape energy vs total path length      Detector size 1x1x1 meter



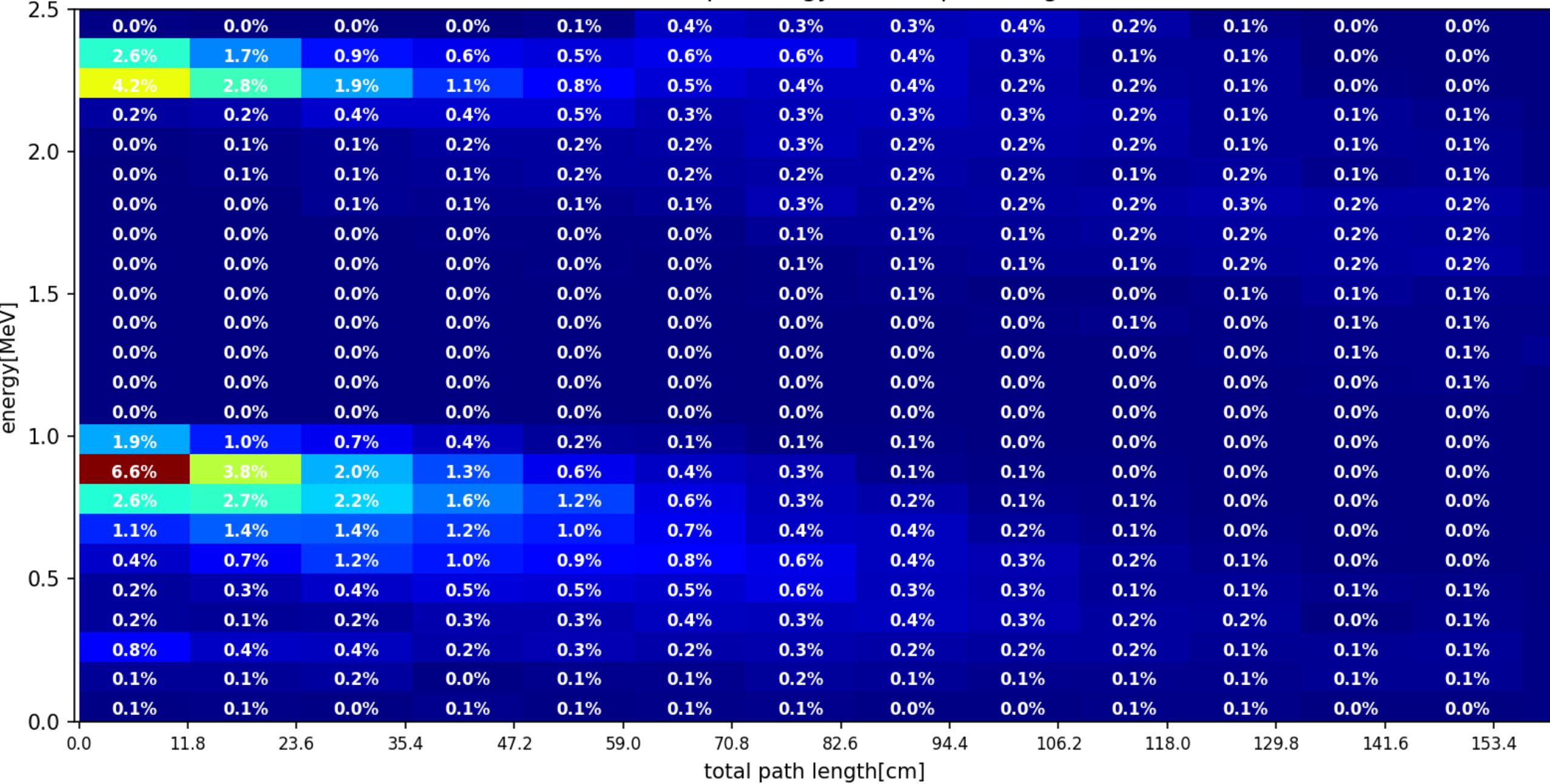
## neutron escape energy vs total path length

Detector size 1x1x1 meter



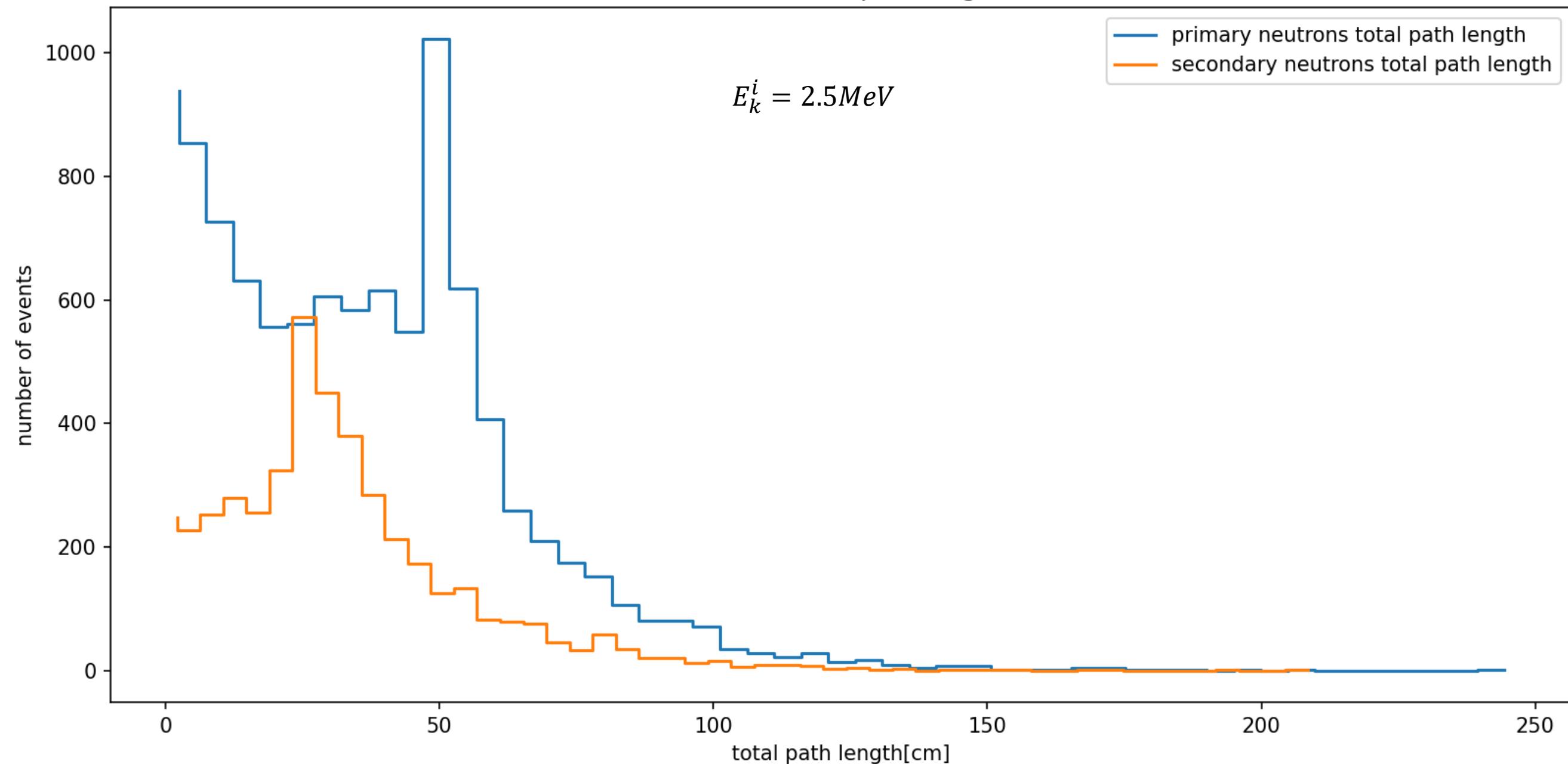
### neutron escape energy vs total path length

Detector size 1x1x1 meter



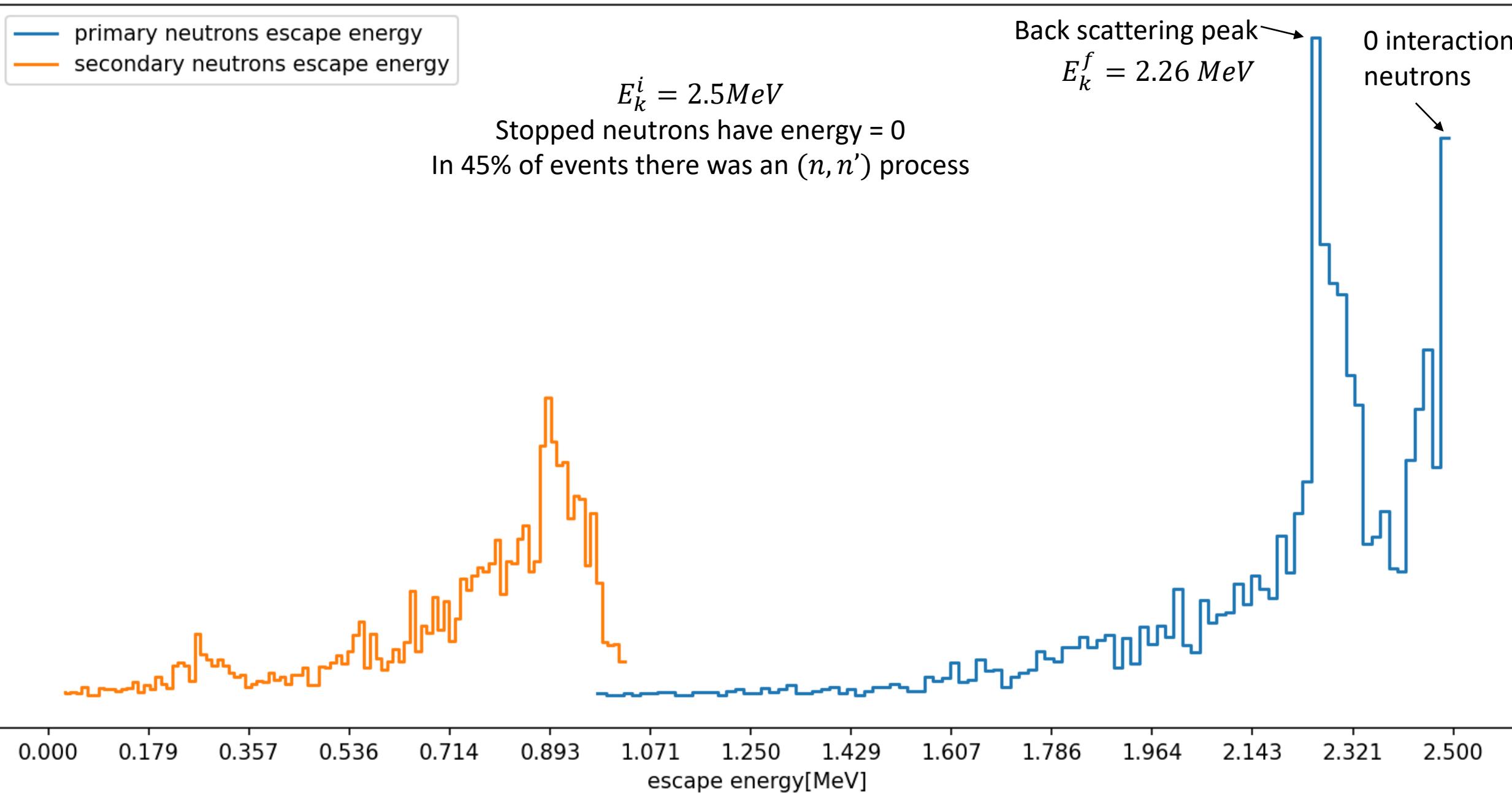
## neutrons total path length

Detector size 0.5x0.5x0.5 meter



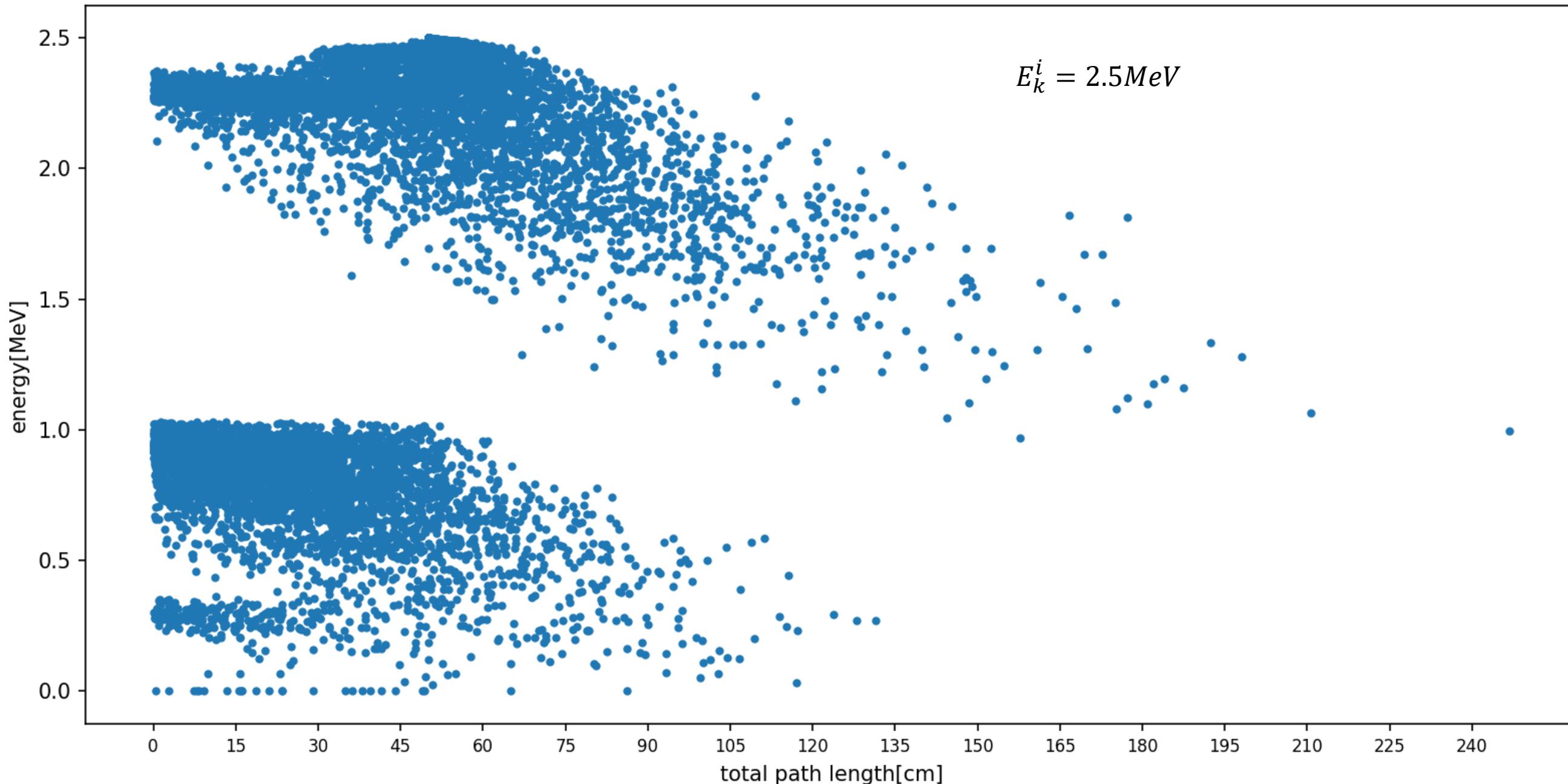
## neutron escape energy distribution

Detector size 0.5x0.5x0.5 meter



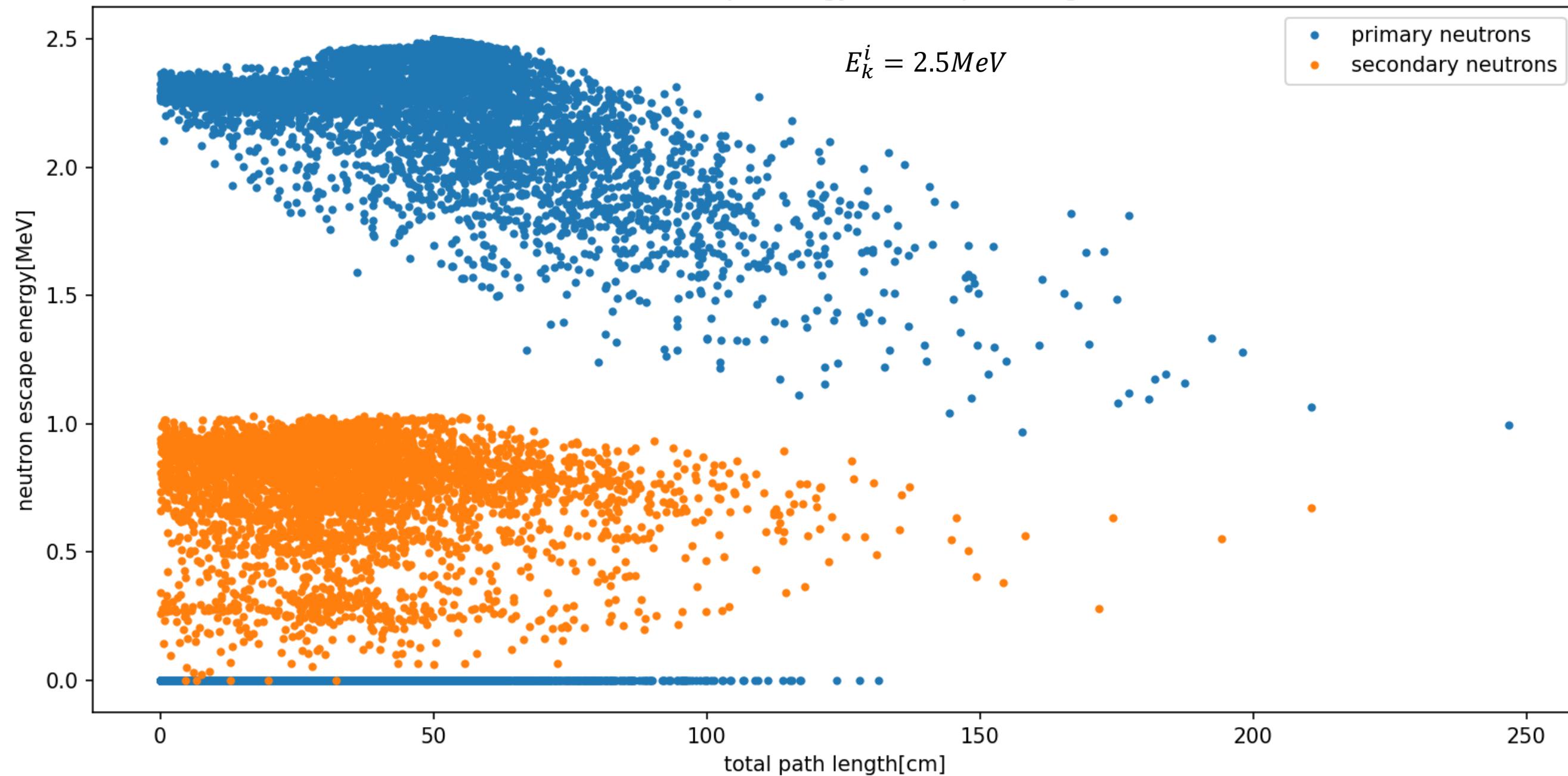
## neutron escape energy vs total path length

Detector size 0.5x0.5x0.5 meter

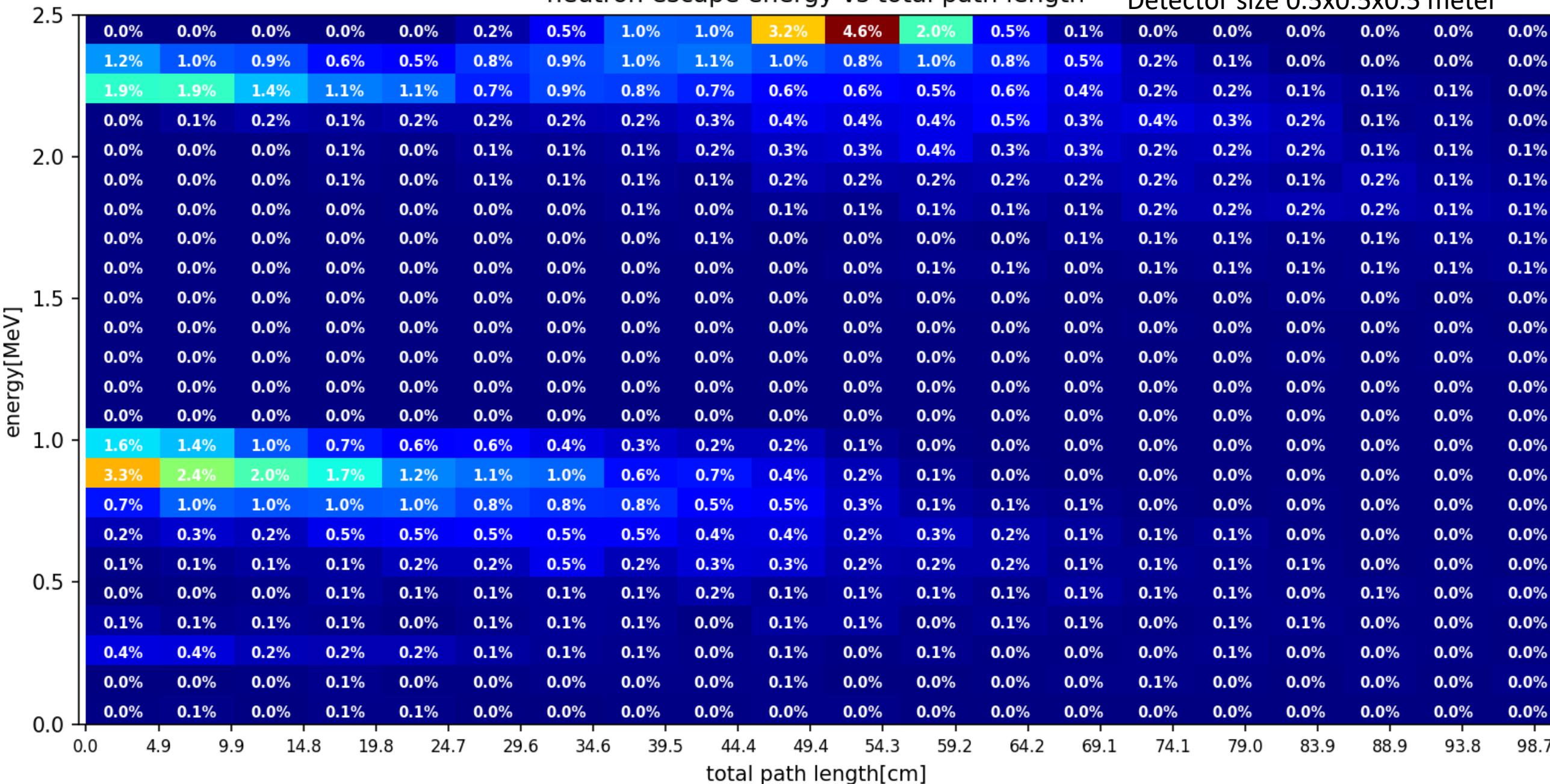


## neutron escape energy vs total path length

Detector size 0.5x0.5x0.5 meter

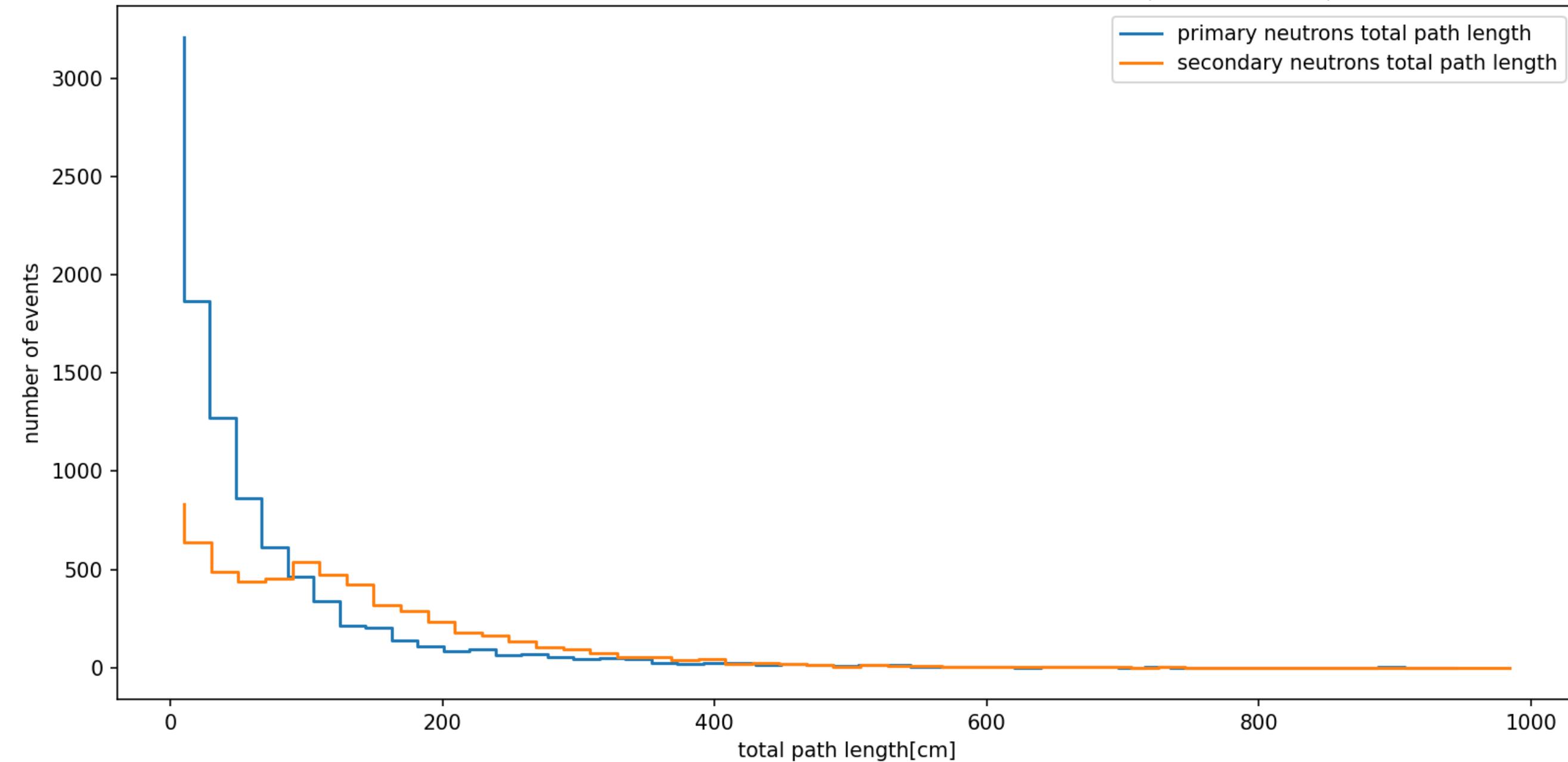


## neutron escape energy vs total path length      Detector size 0.5x0.5x0.5 meter



Detector size 1x2x2 meter  
(thickness = 1m)

neutrons total path length



### neutron escape energy distribution

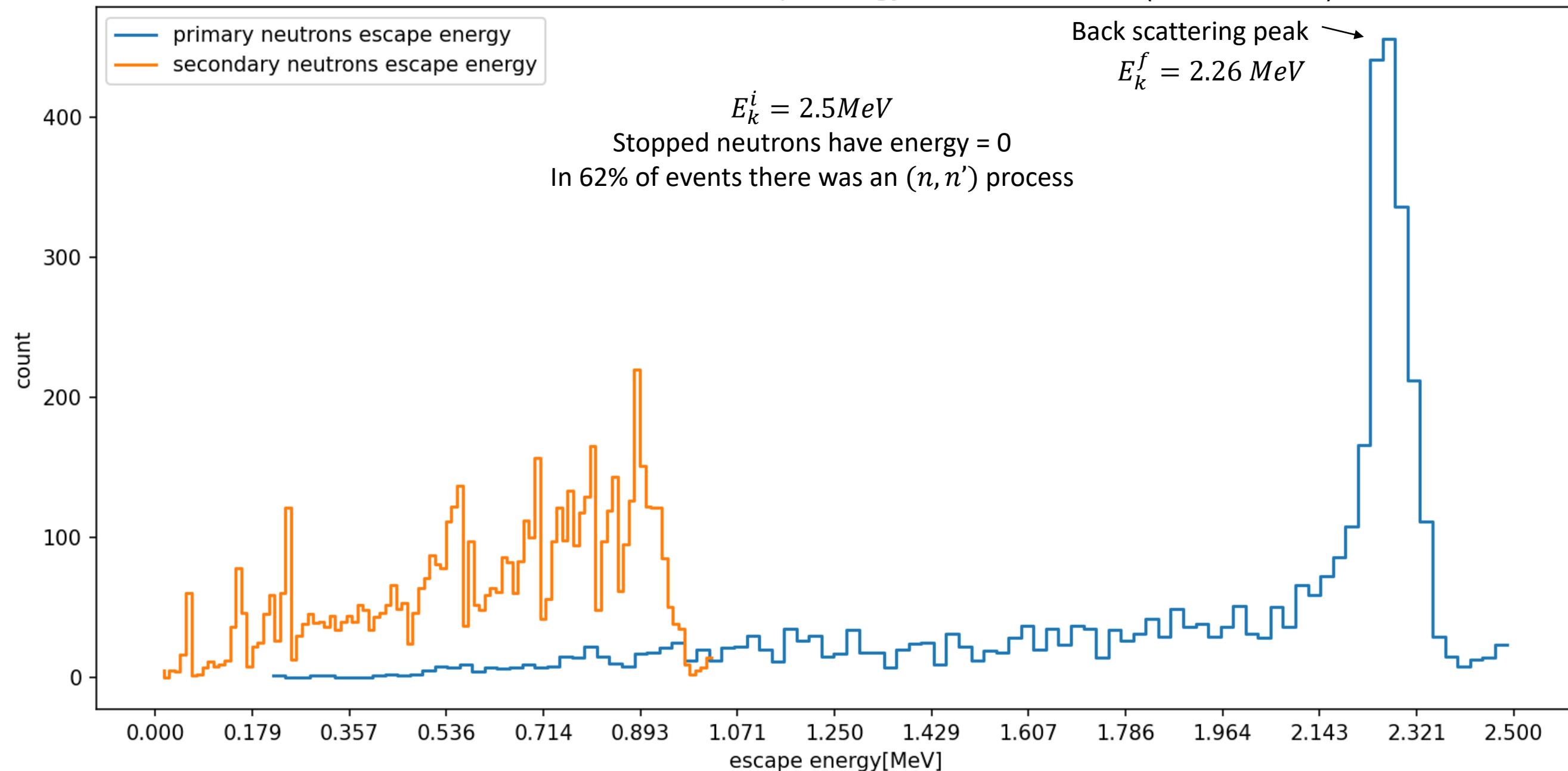
Detector size 1x2x2 meter  
(thickness = 1m)

Back scattering peak →  
 $E_k^f = 2.26 \text{ MeV}$

$$E_k^i = 2.5 \text{ MeV}$$

Stopped neutrons have energy = 0

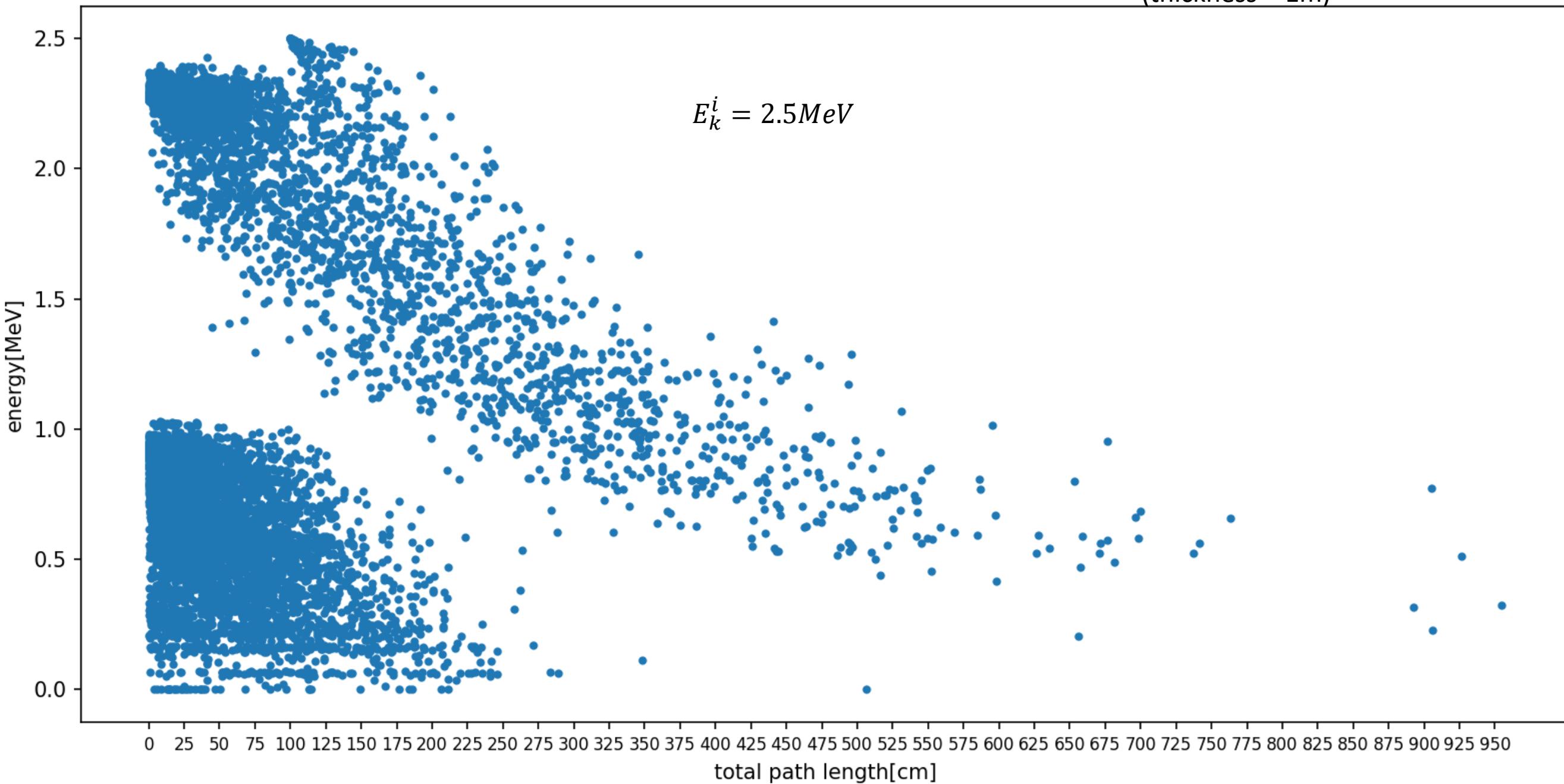
In 62% of events there was an  $(n, n')$  process



Detector size 1x2x2 meter  
(thickness = 1m)

neutron escape energy vs total path length

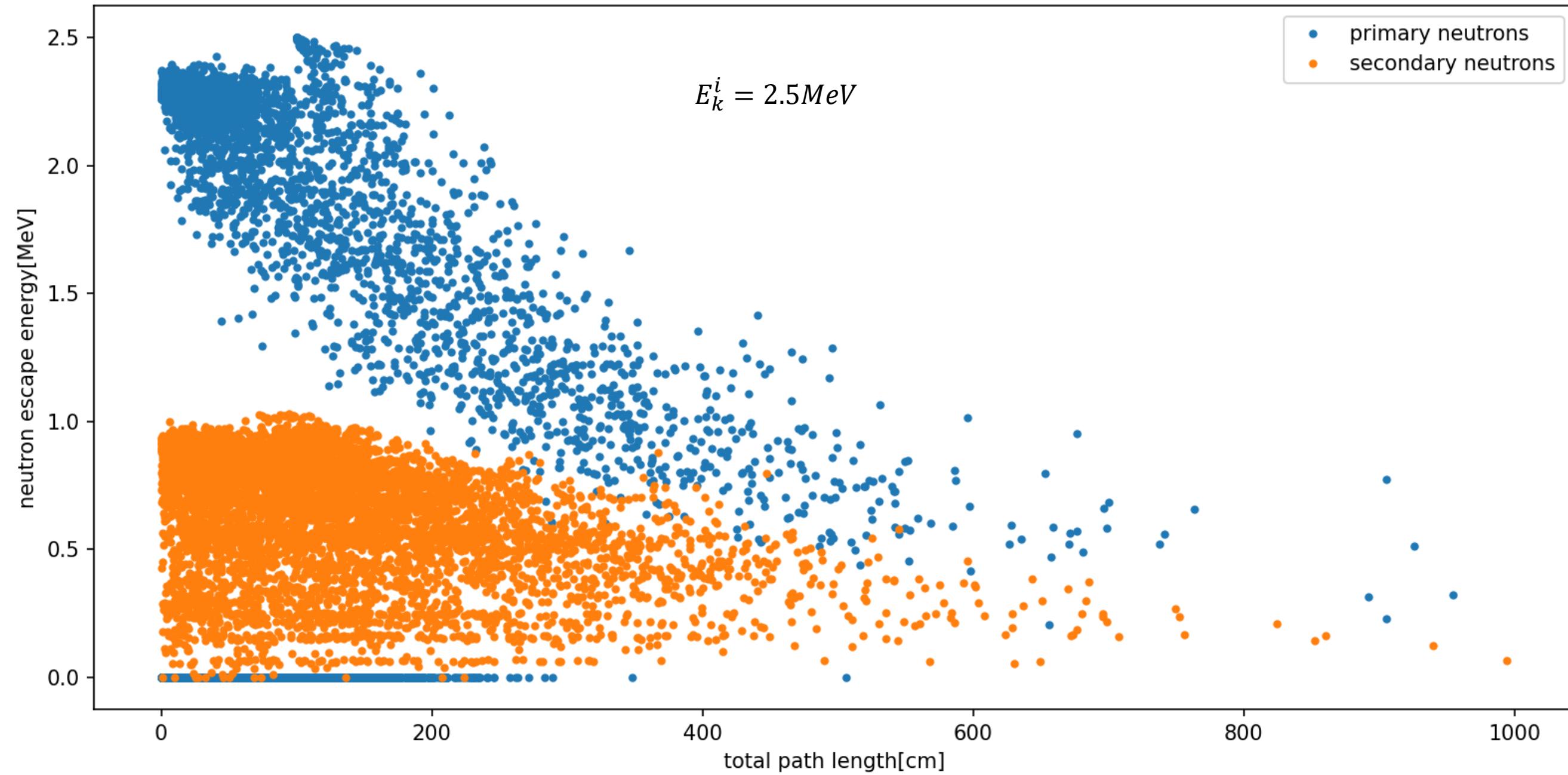
$$E_k^i = 2.5 \text{ MeV}$$



Detector size 1x2x2 meter  
(thickness = 1m)

neutron escape energy vs total path length

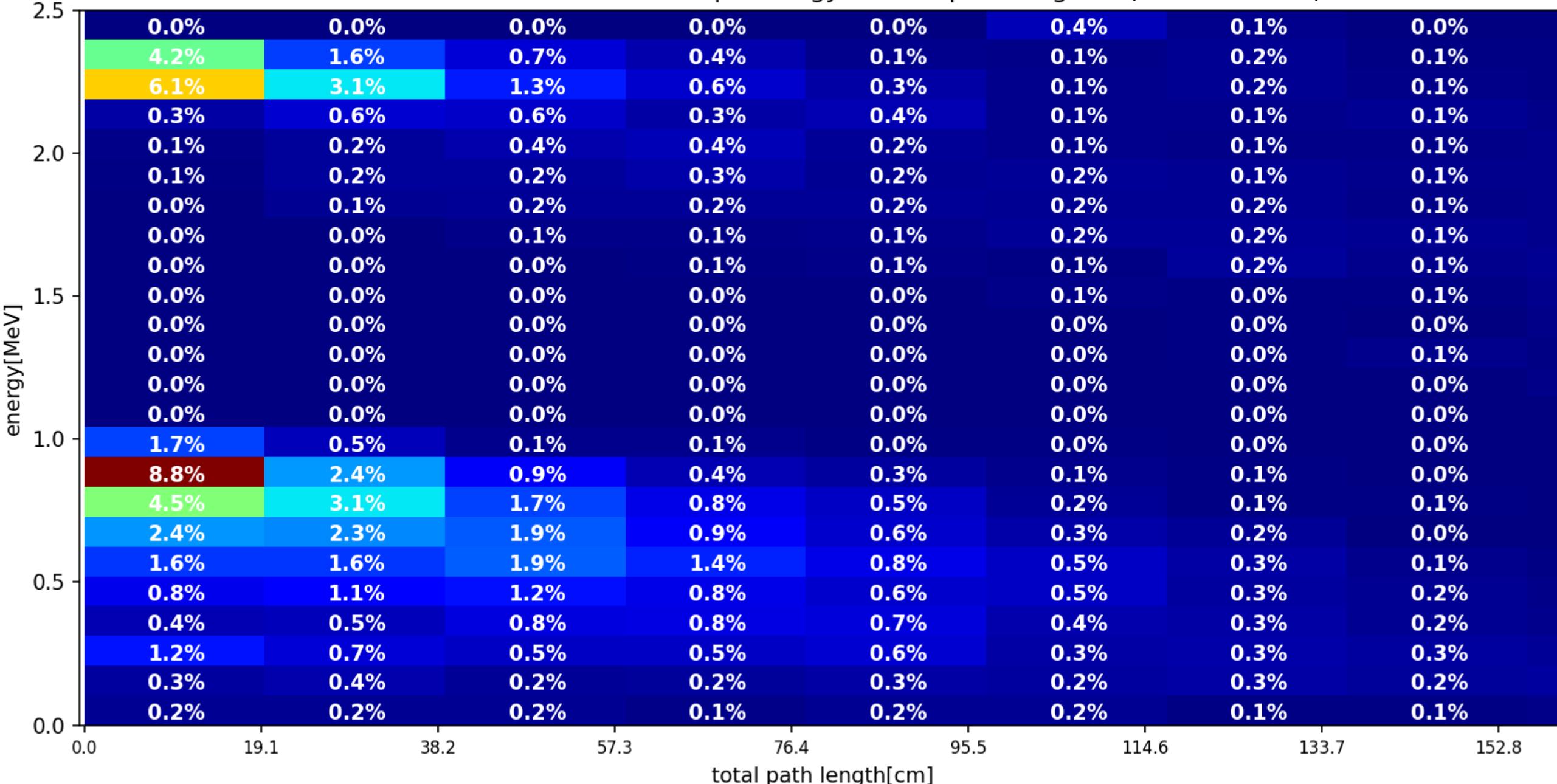
$$E_k^i = 2.5 \text{ MeV}$$



Only for primary neutrons (didn't perform (n,n'))

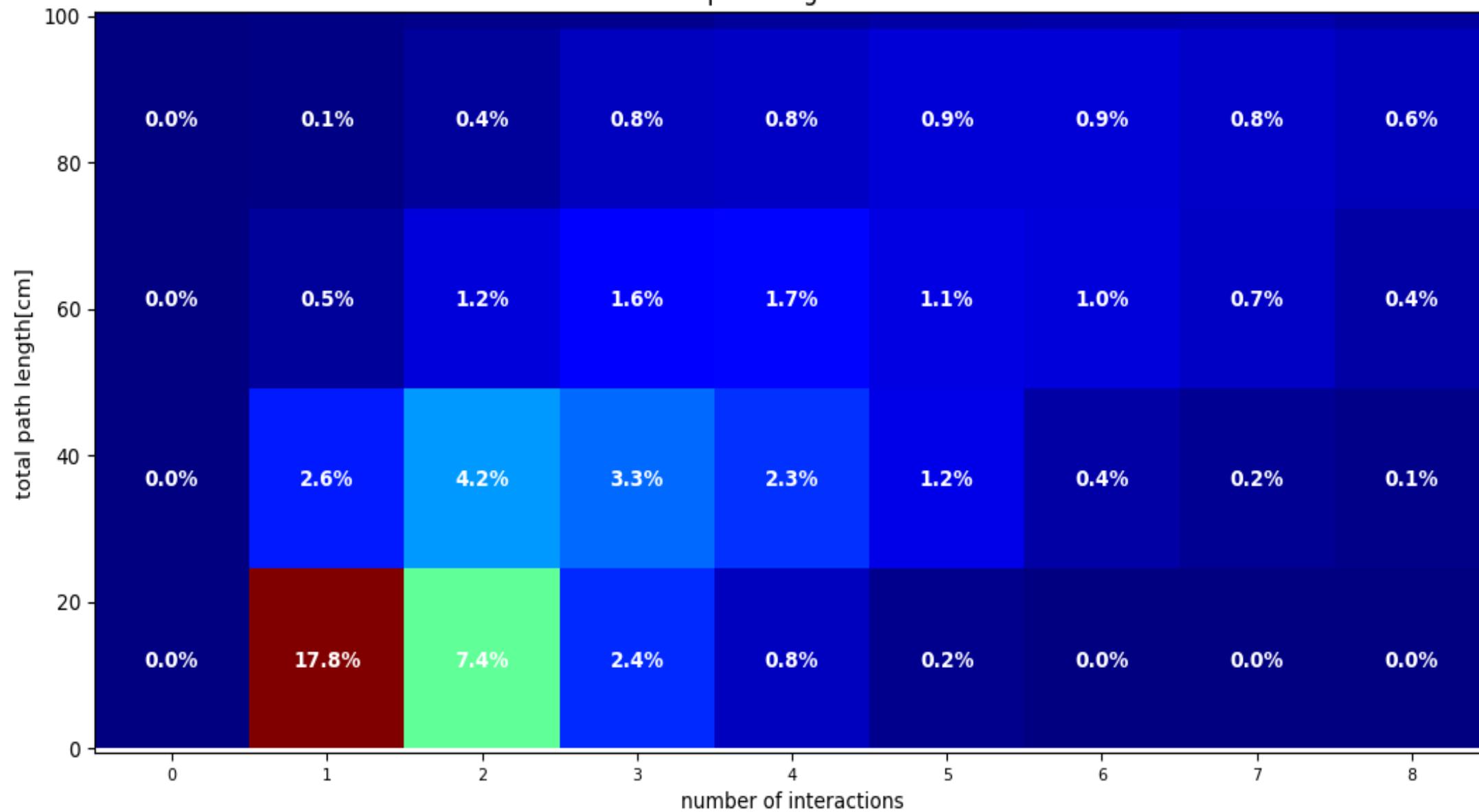
Detector size 1x2x2 meter  
(thickness = 1m)

neutron escape energy vs total path length

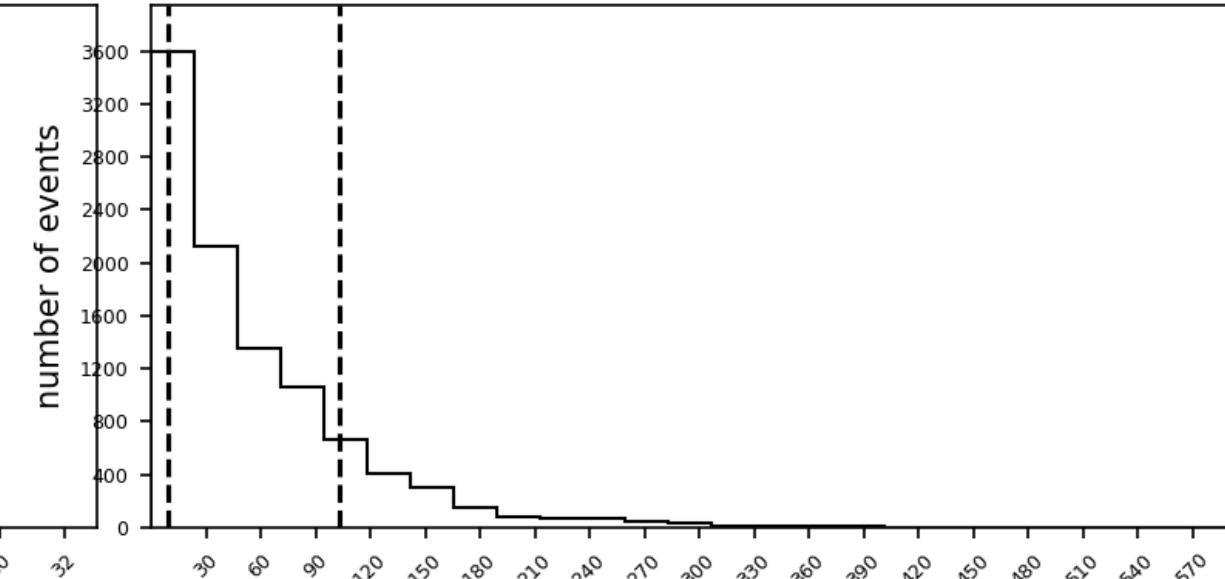
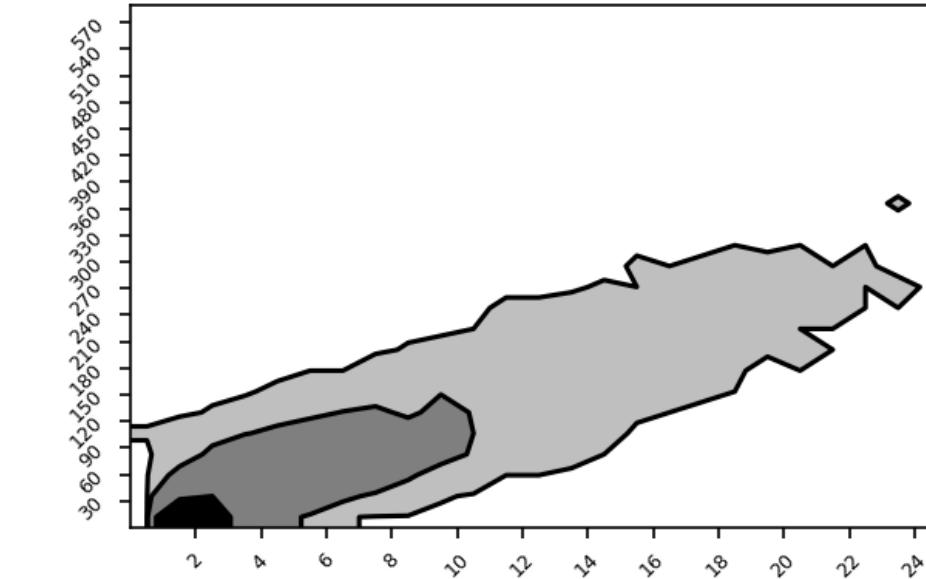
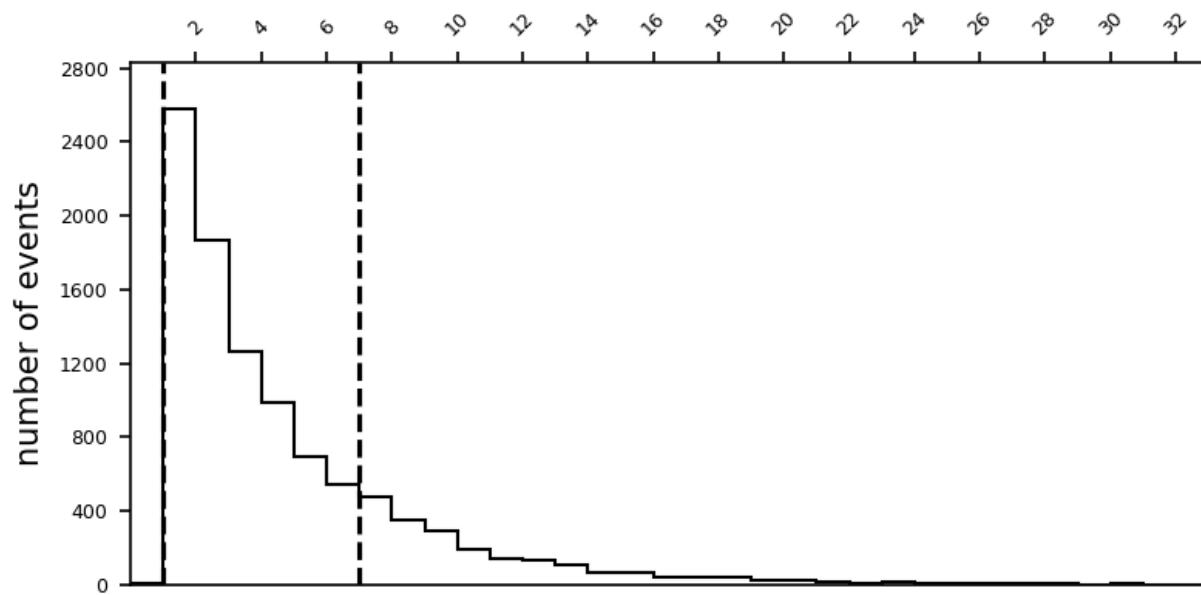


Only for primary neutrons (didn't perform (n,n'))

## 2.5 MeV neutrons total path length vs number of interactions



Only for primary neutrons (didn't perform (n,n'))

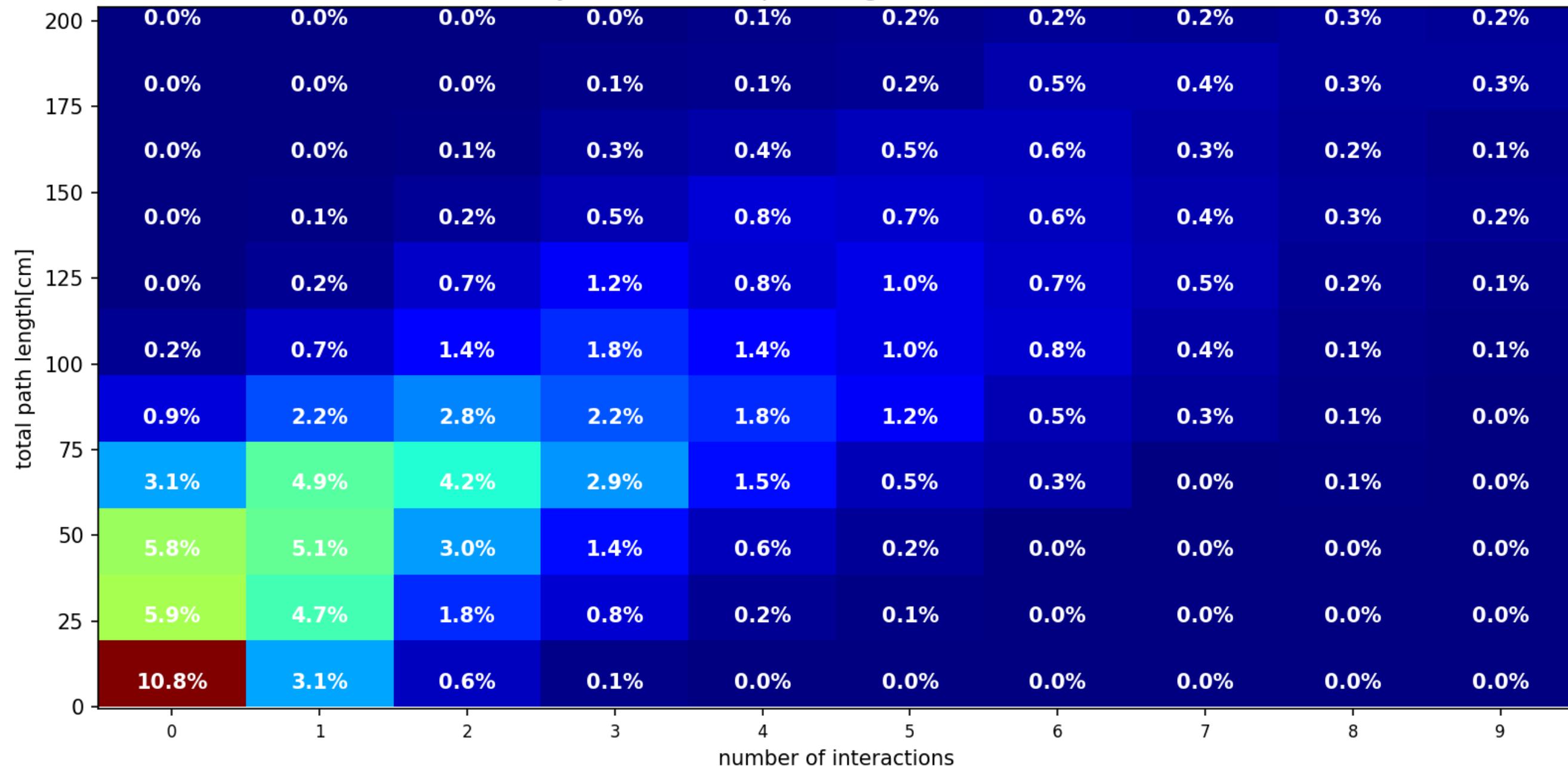


numberOfInteractions

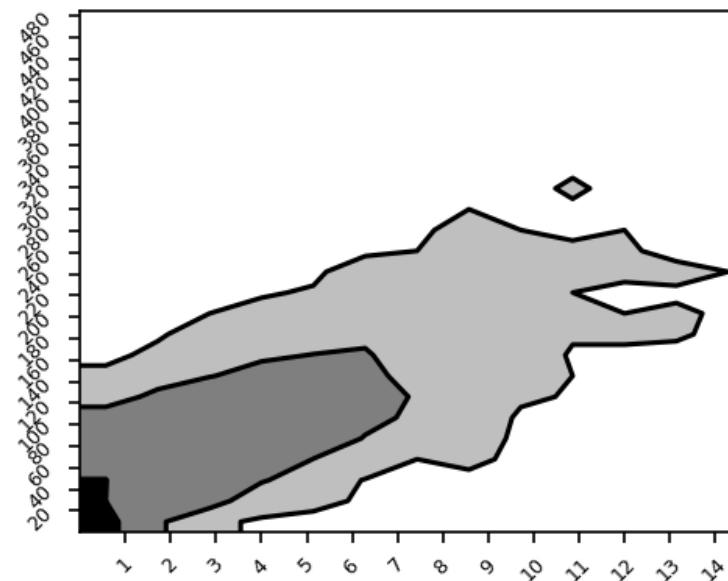
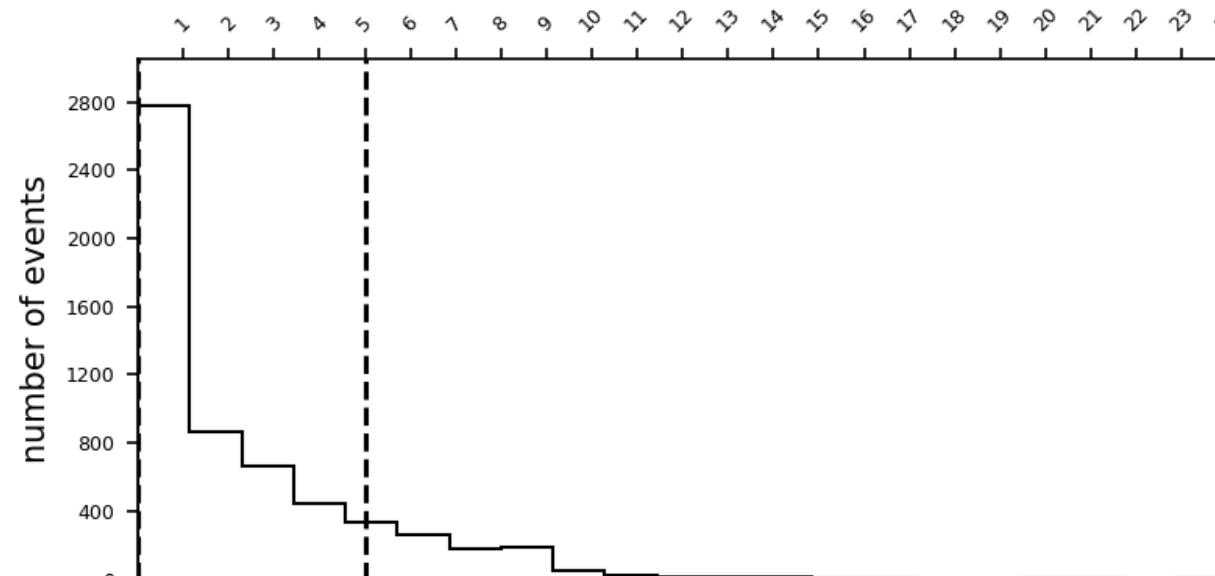
total path length[cm]

TotalPathLength

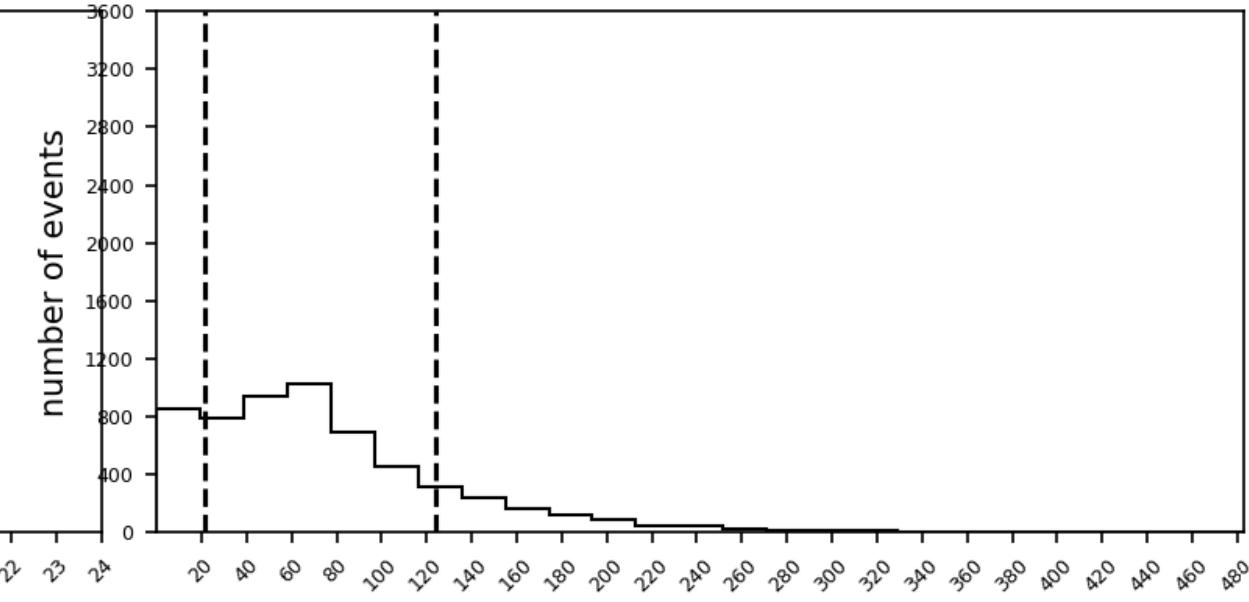
### Secondary neutrons total path length vs number of interactions



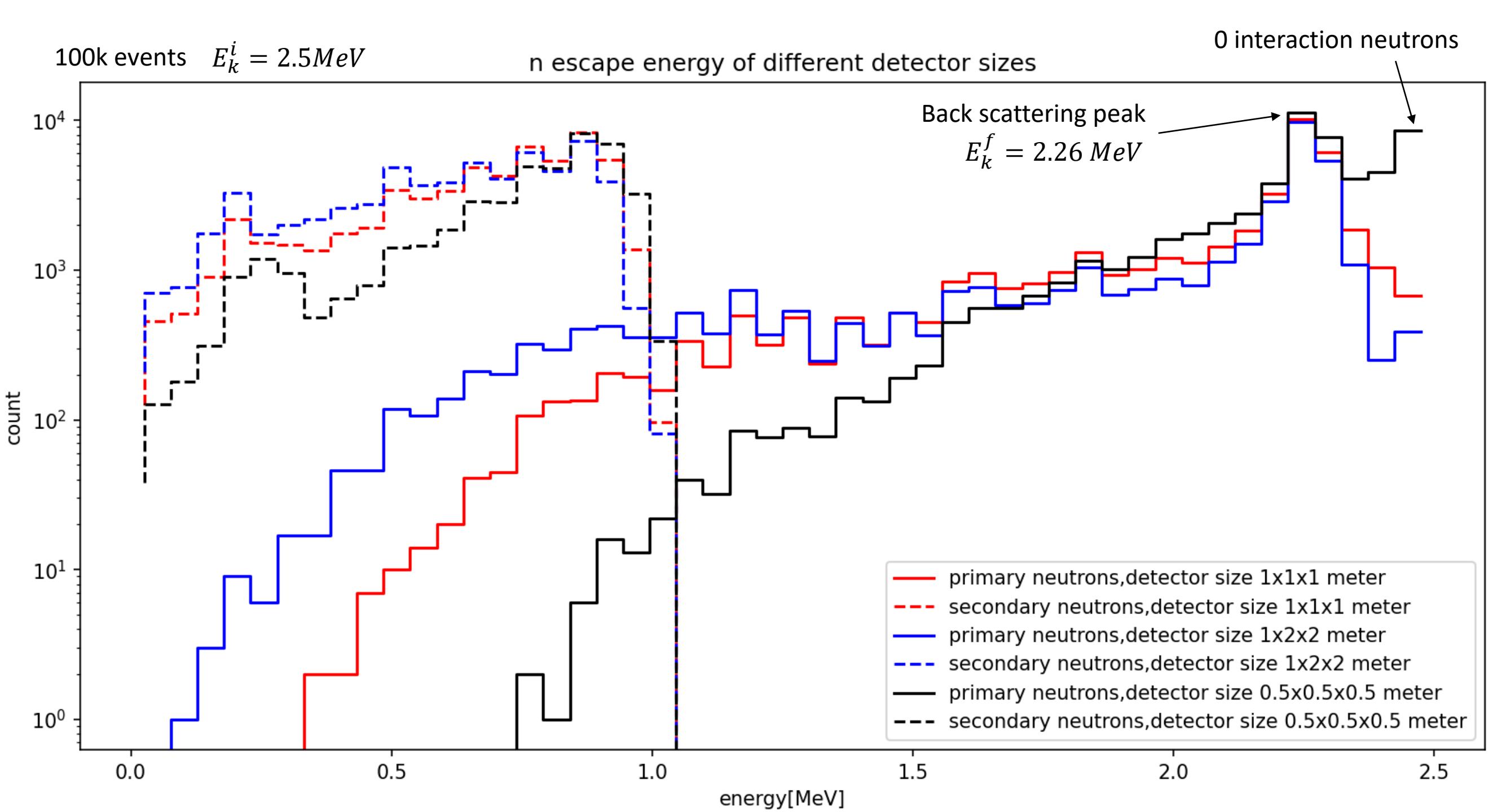
TotalPathLength



numberOfInteractions

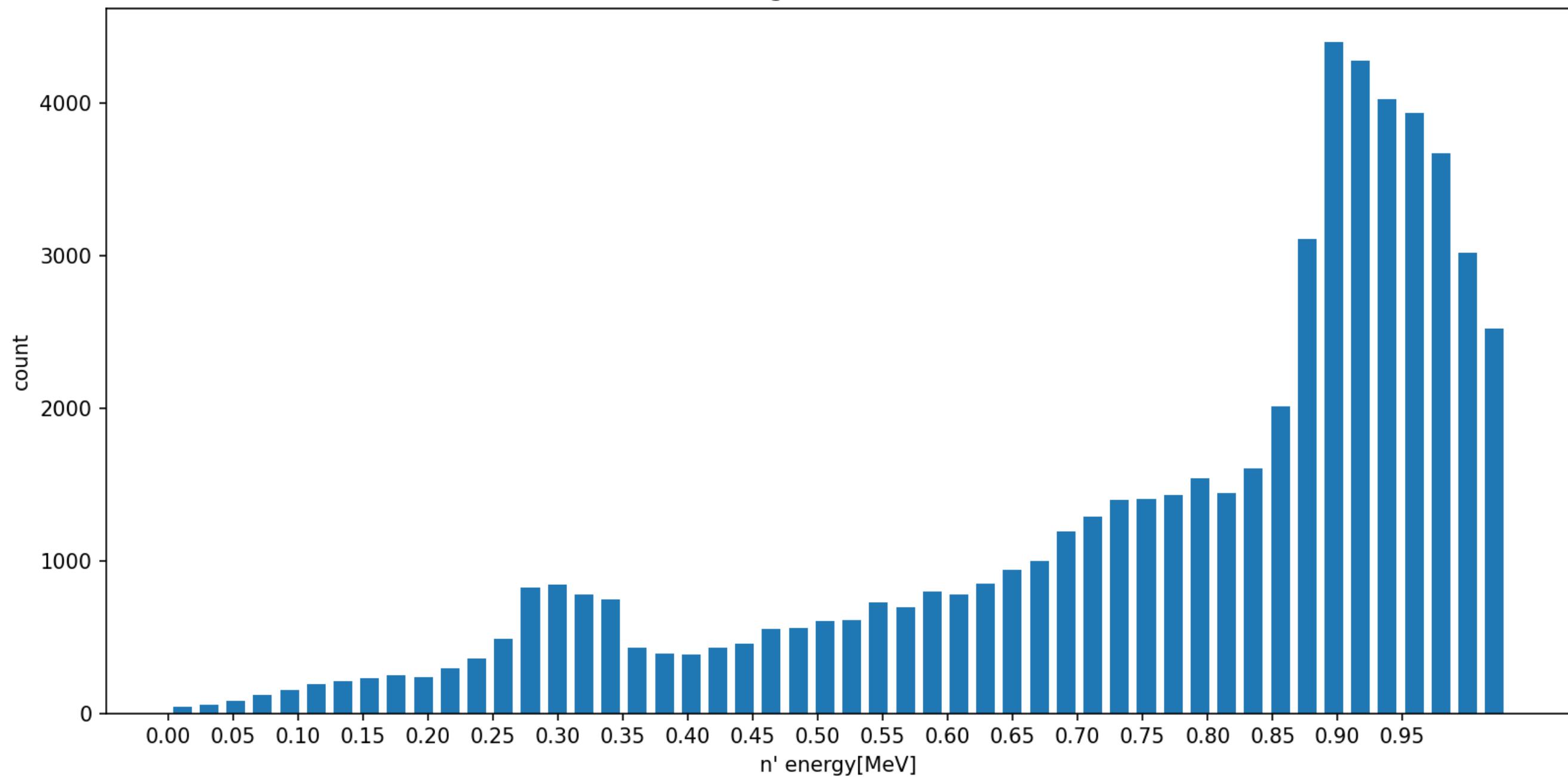


total path length[cm]

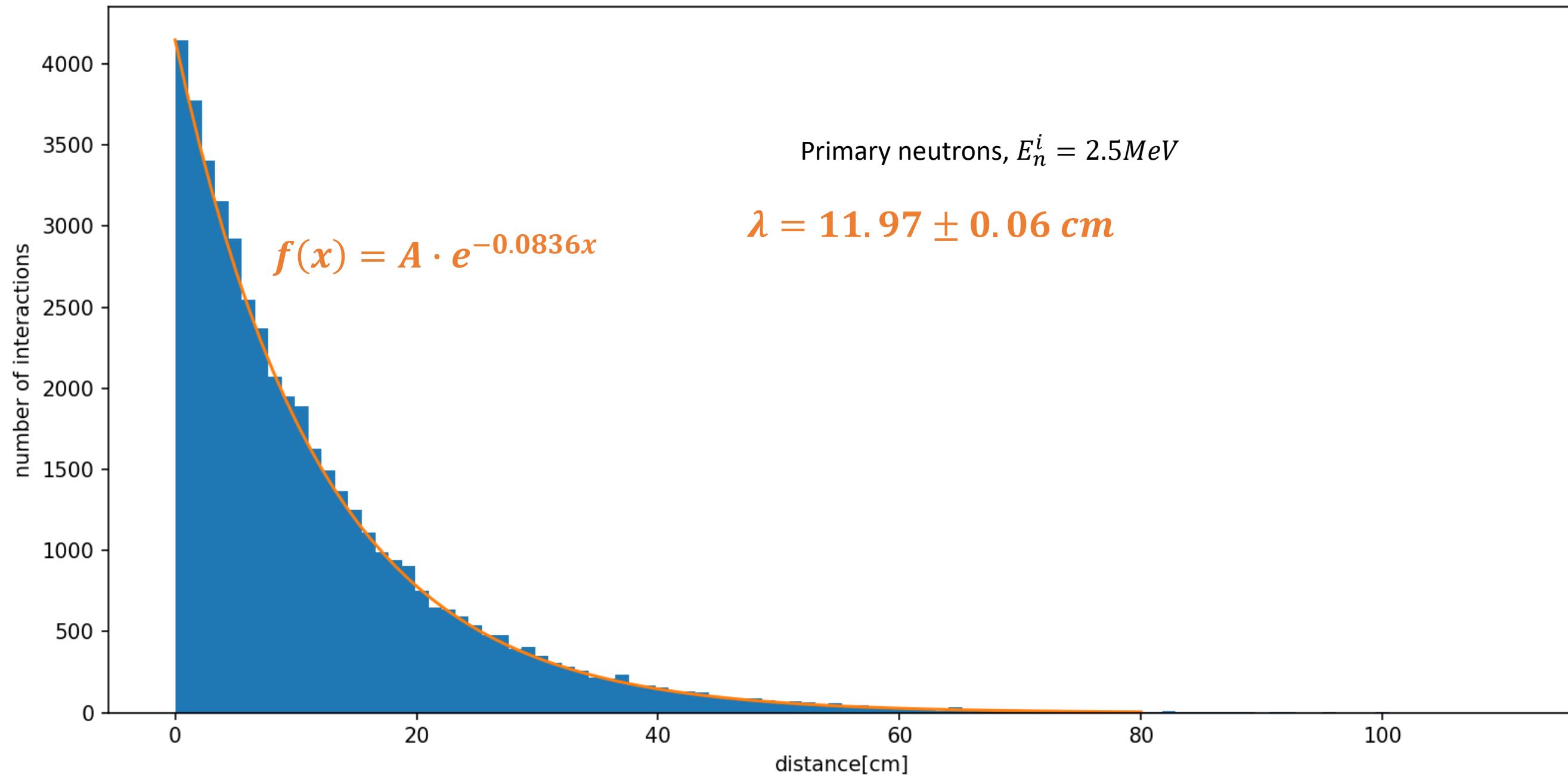


$n'$  energies after creation

100k events

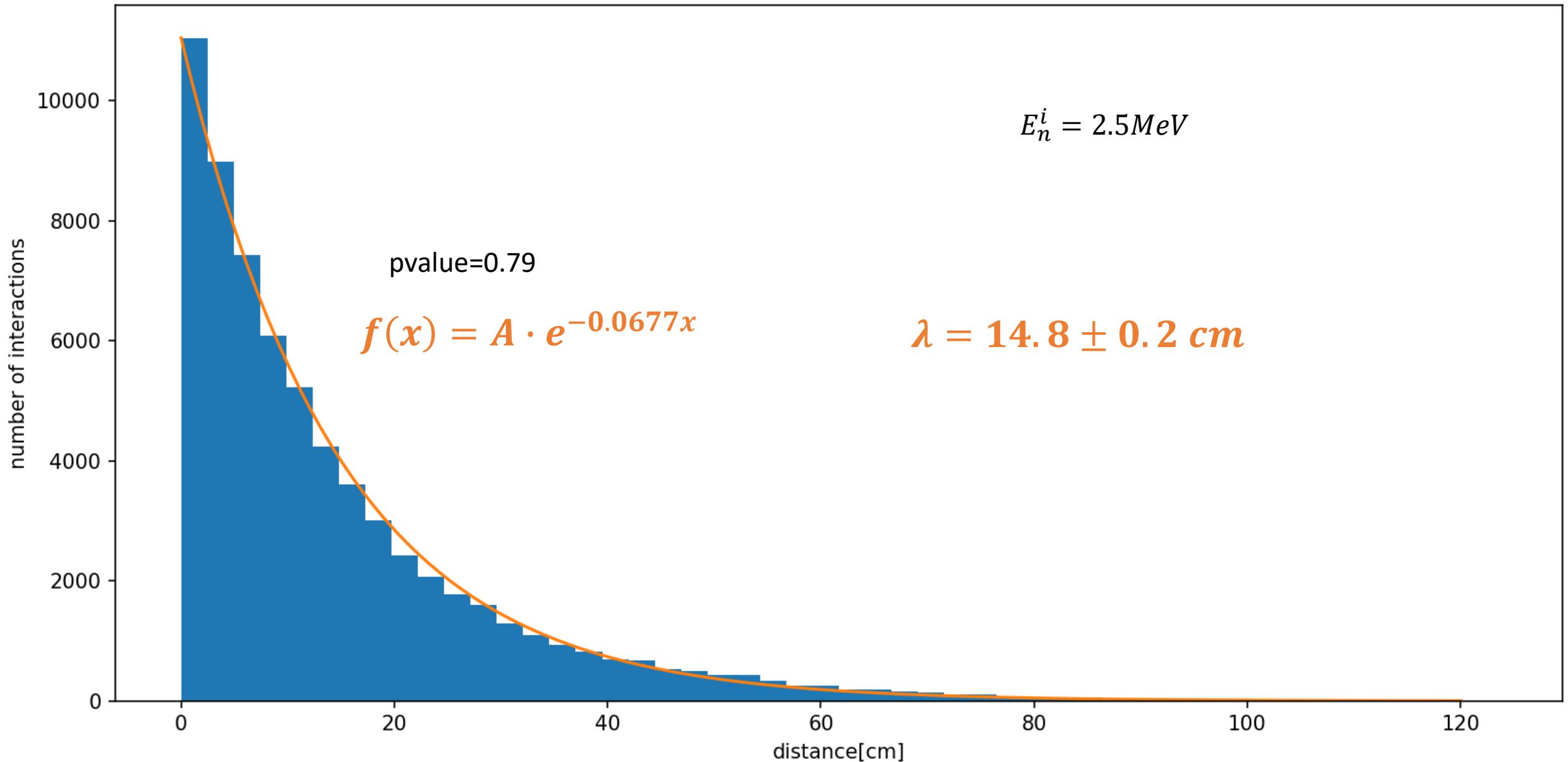


Distribution does not include secondary neutrons  
distance between interactions histogram



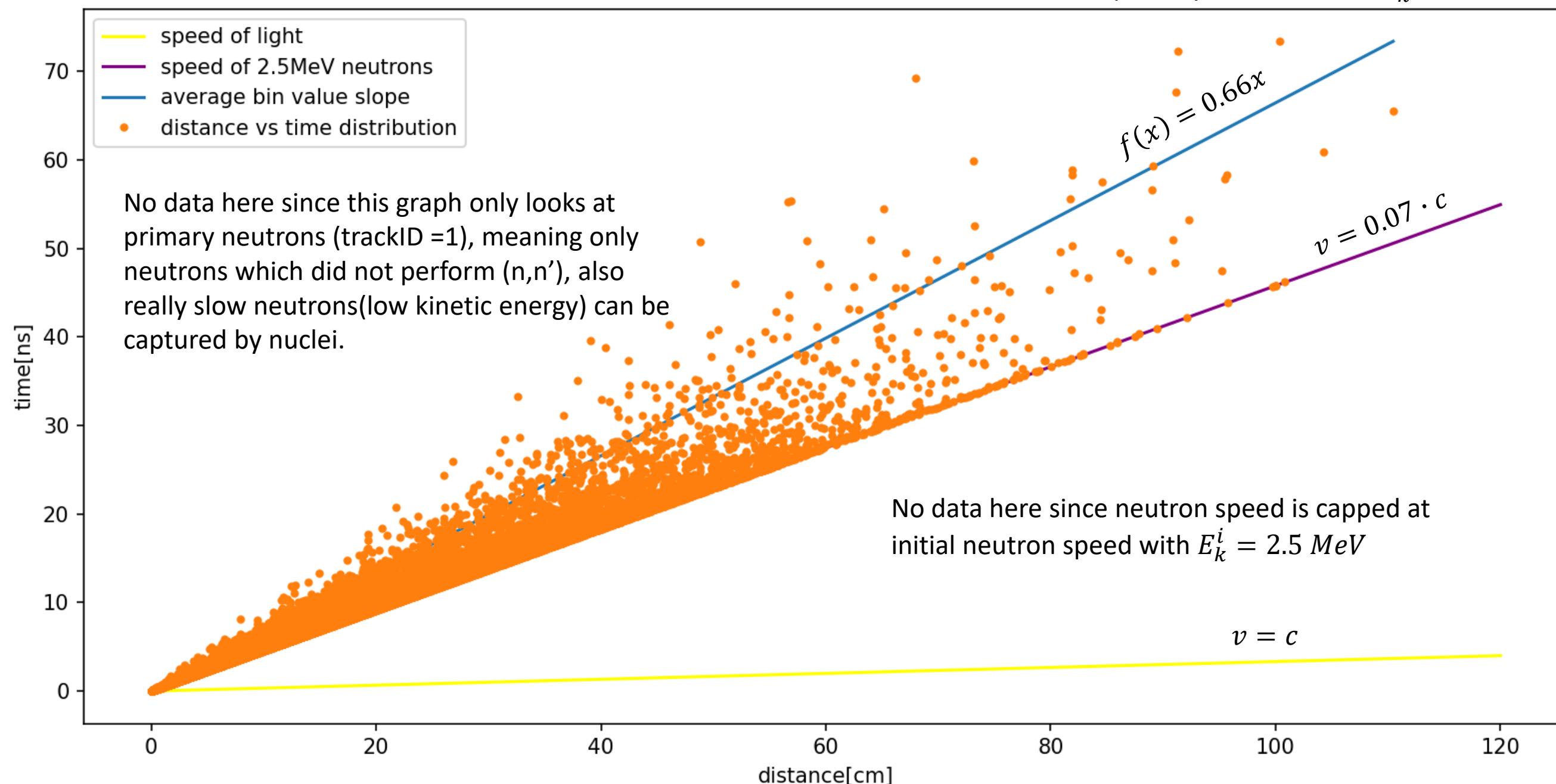
Distribution also includes secondary neutrons

distance between interactions histogram



Only for primary neutrons (didn't perform (n,n'))

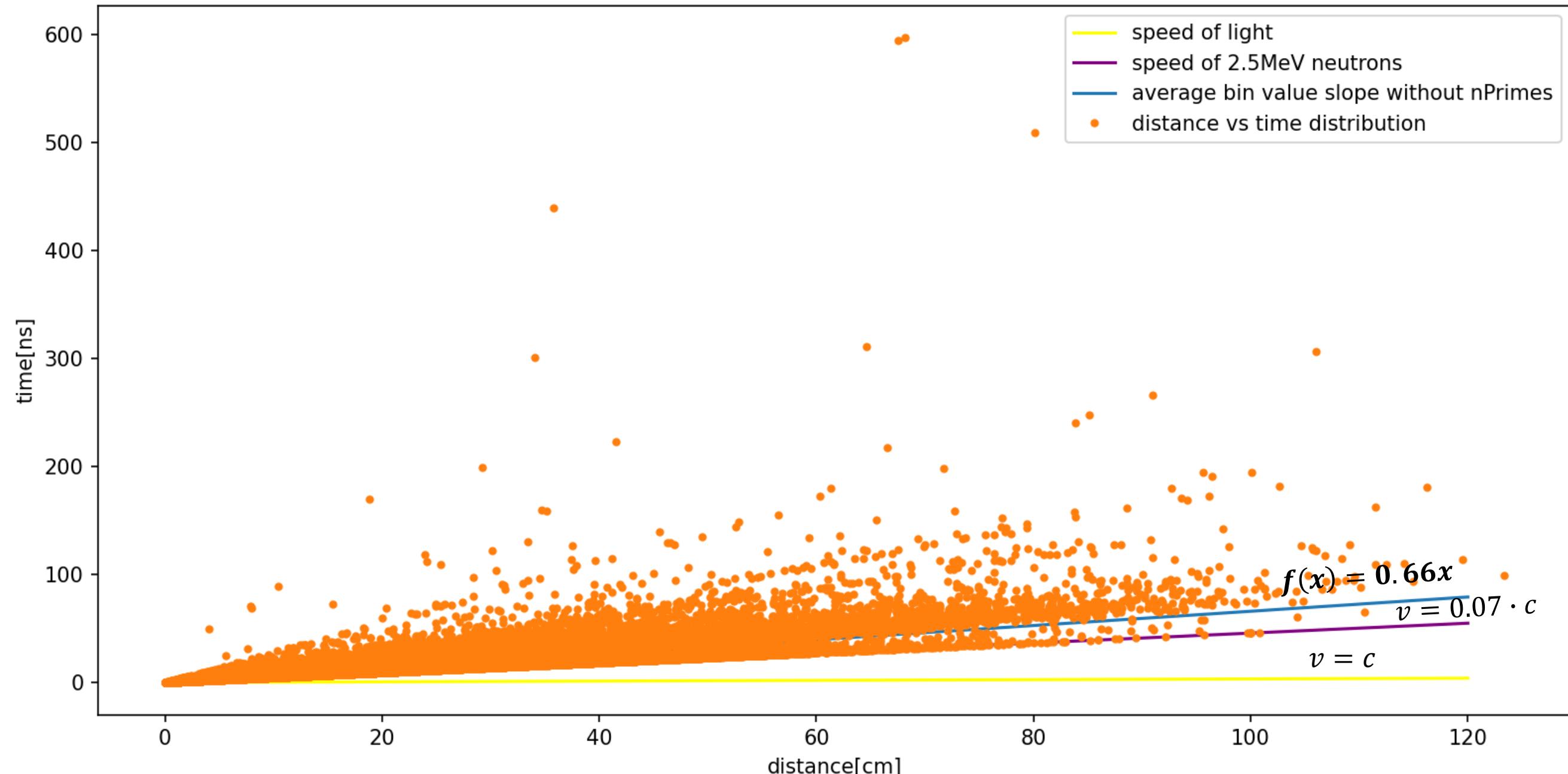
distance vs time between interactions Of primary neutrons with  $E_k^i = 2.5 \text{ MeV}$



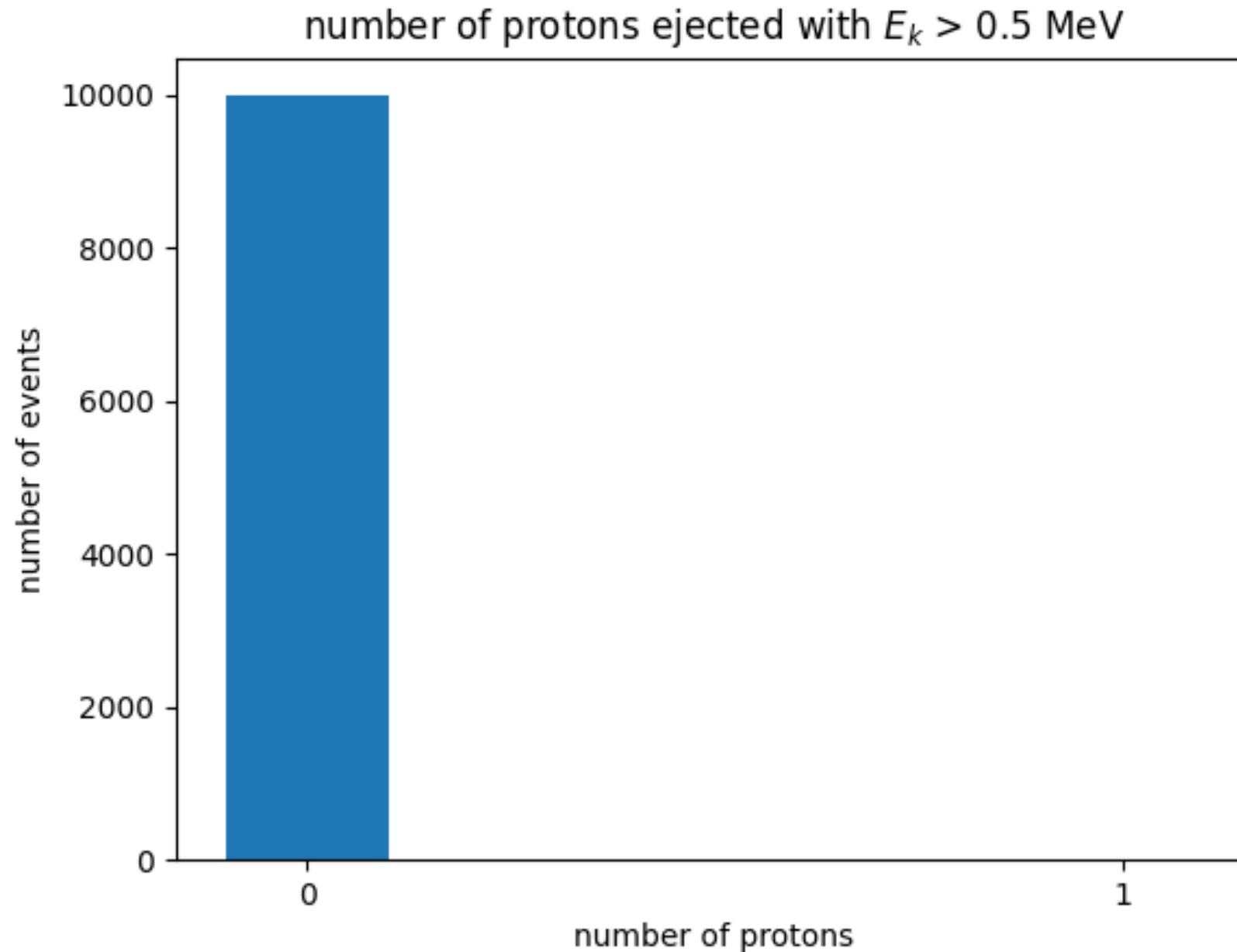
$E_k^i = 2.5\text{MeV}$

distance vs time between interactions

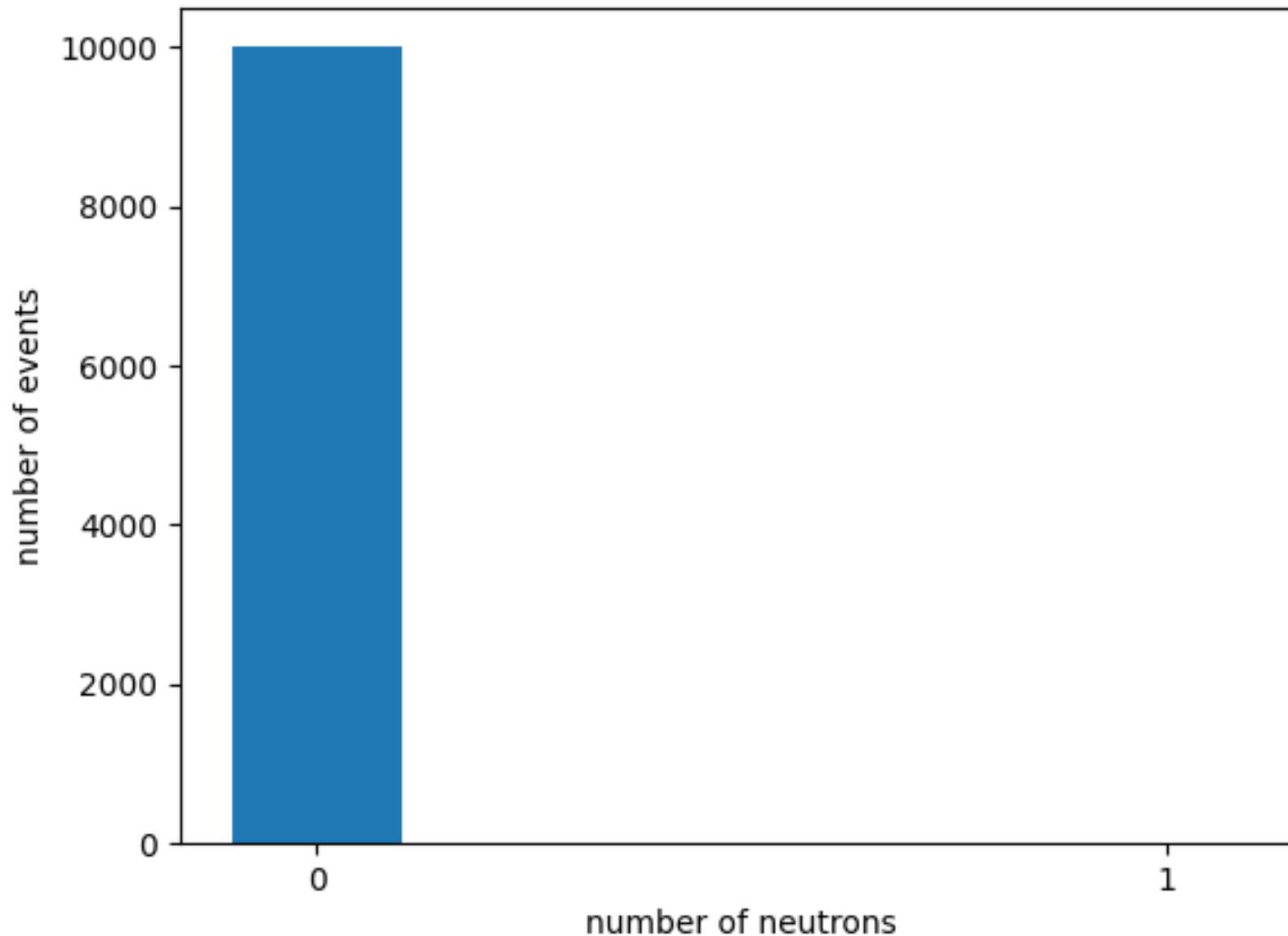
Also includes secondary neutrons ( $n'$ )



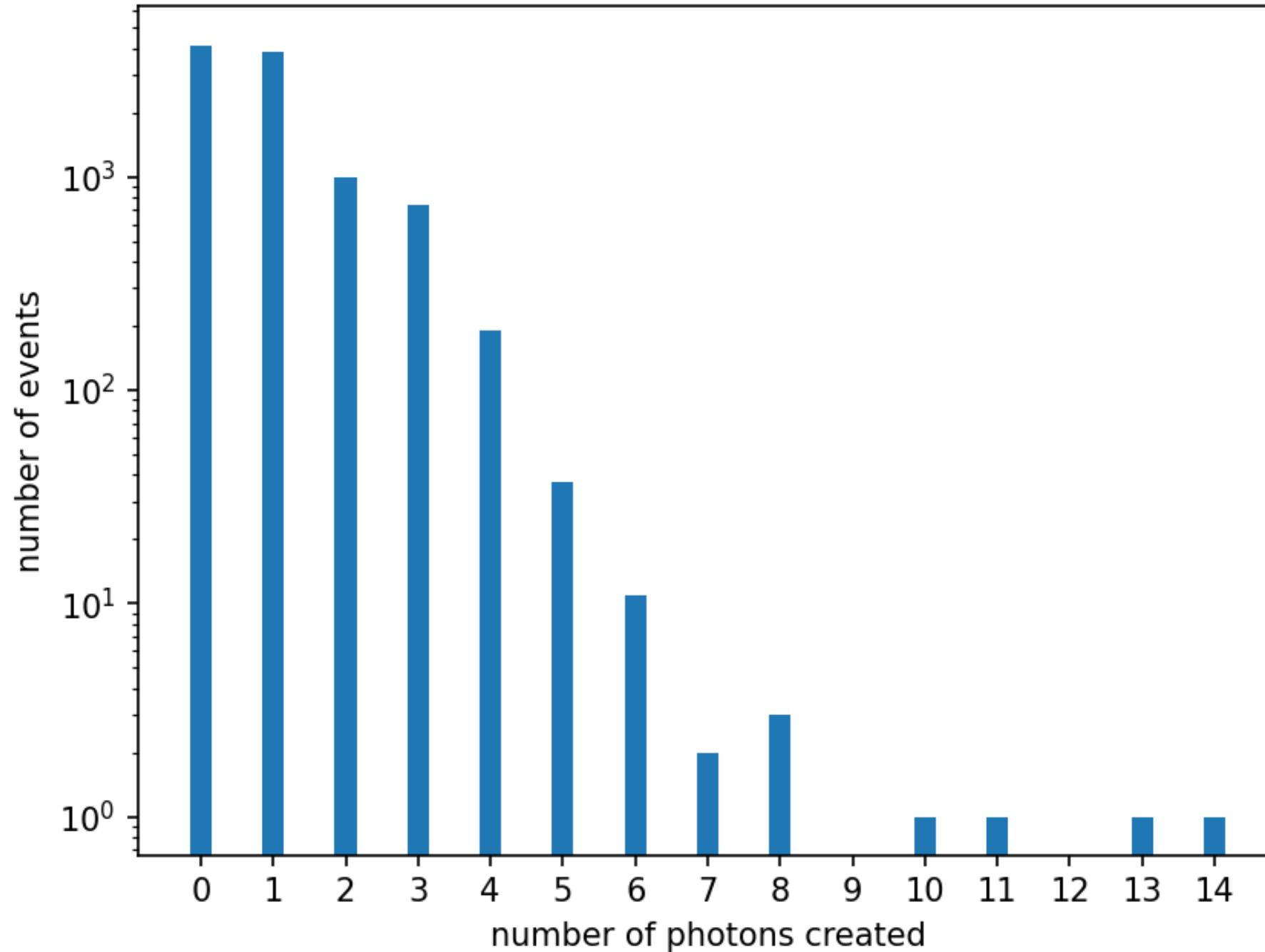
No protons exit the detector, though some are created in (n,p) process with Argon 36

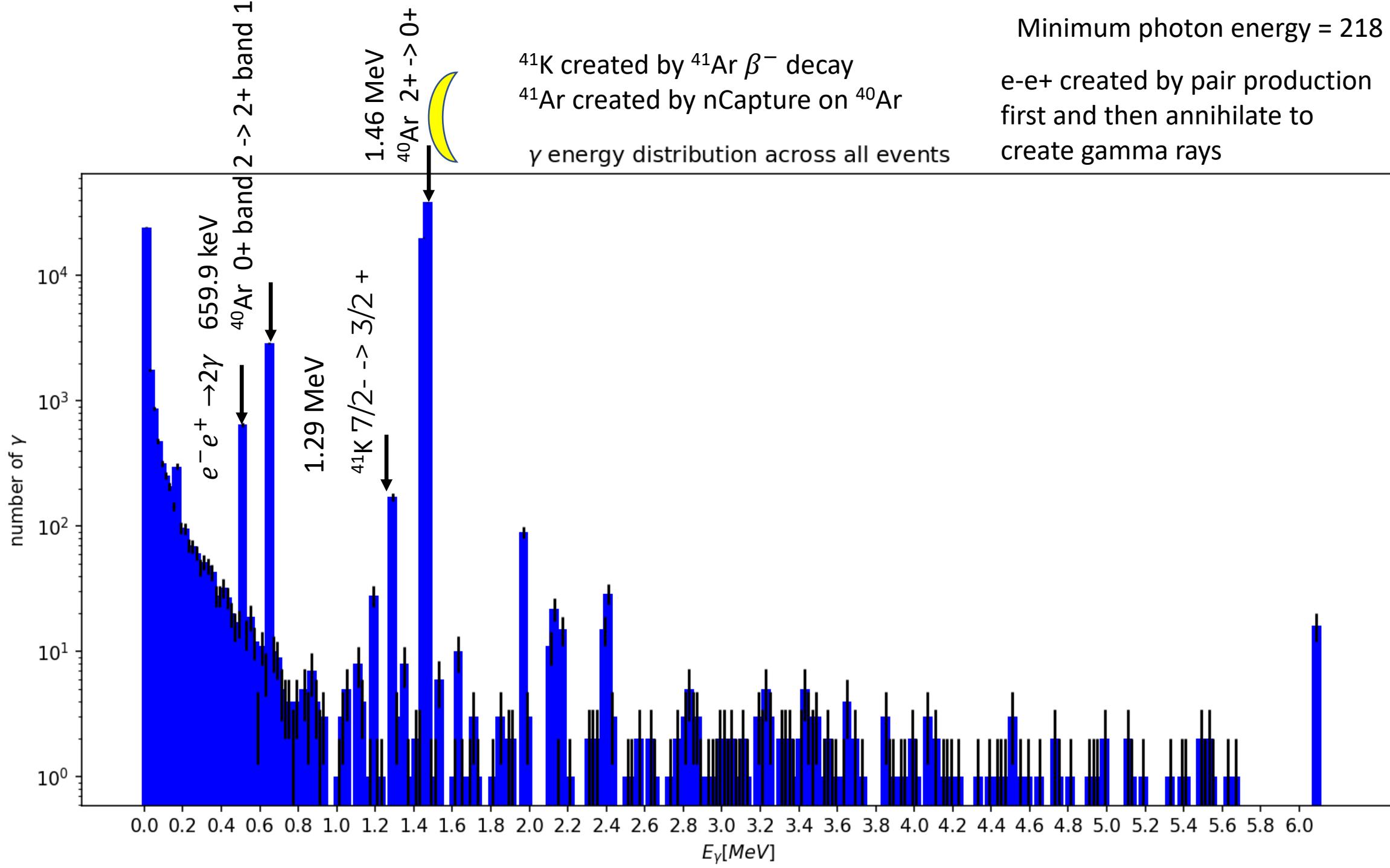


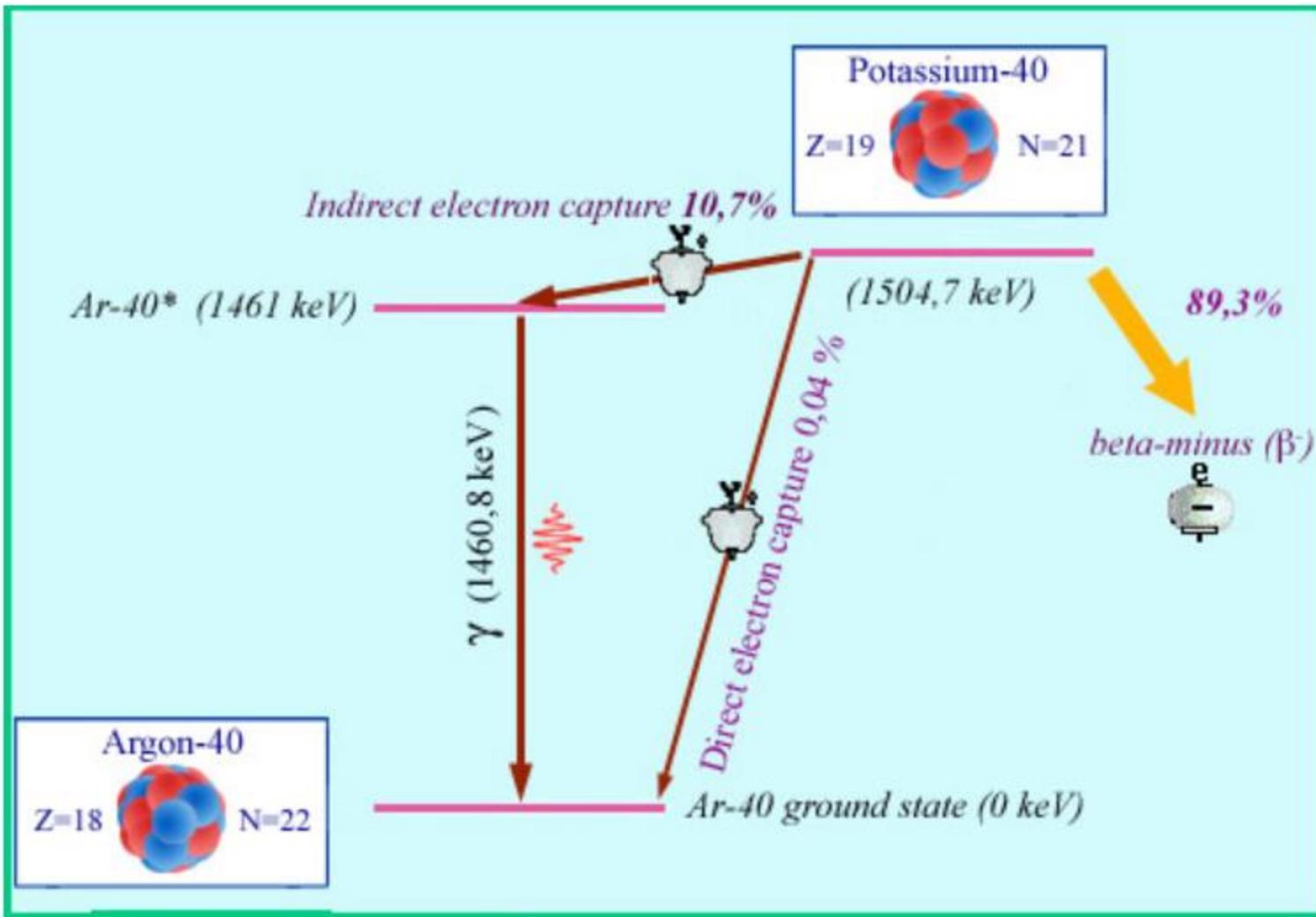
number of new neutrons ejected in an event



number of photons created per event



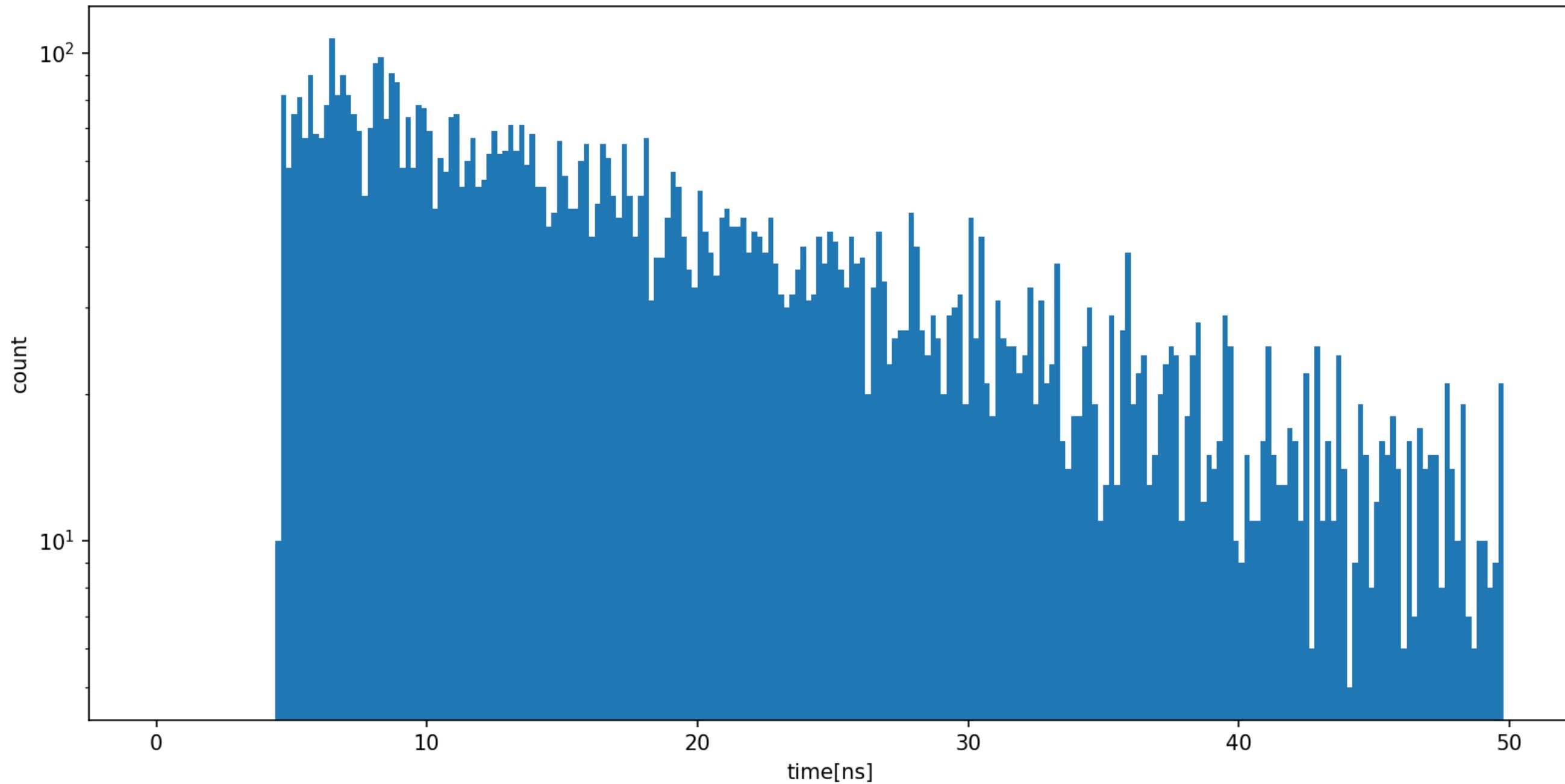




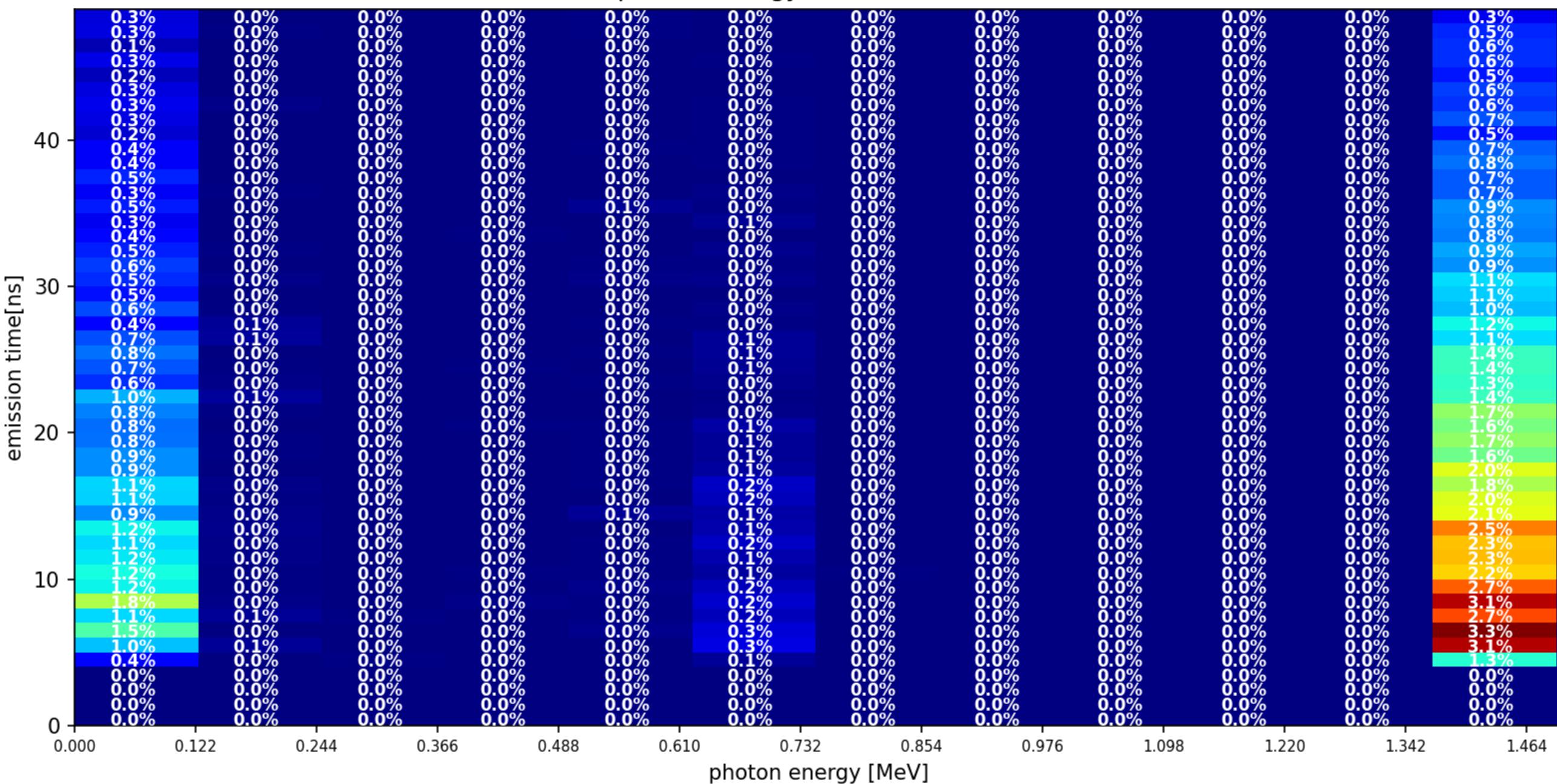
### From potassium 40 to argon 40

The electron capture which causes potassium 40 to transform into argon 40 in its ground state takes place in only 0.04% of cases. Far more frequently (10.68% of the time), an indirect capture leads to an excited argon atom which needs to return to its ground state by emitting a gamma ray at an energy of 1.46 MeV. Without this characteristic gamma ray,

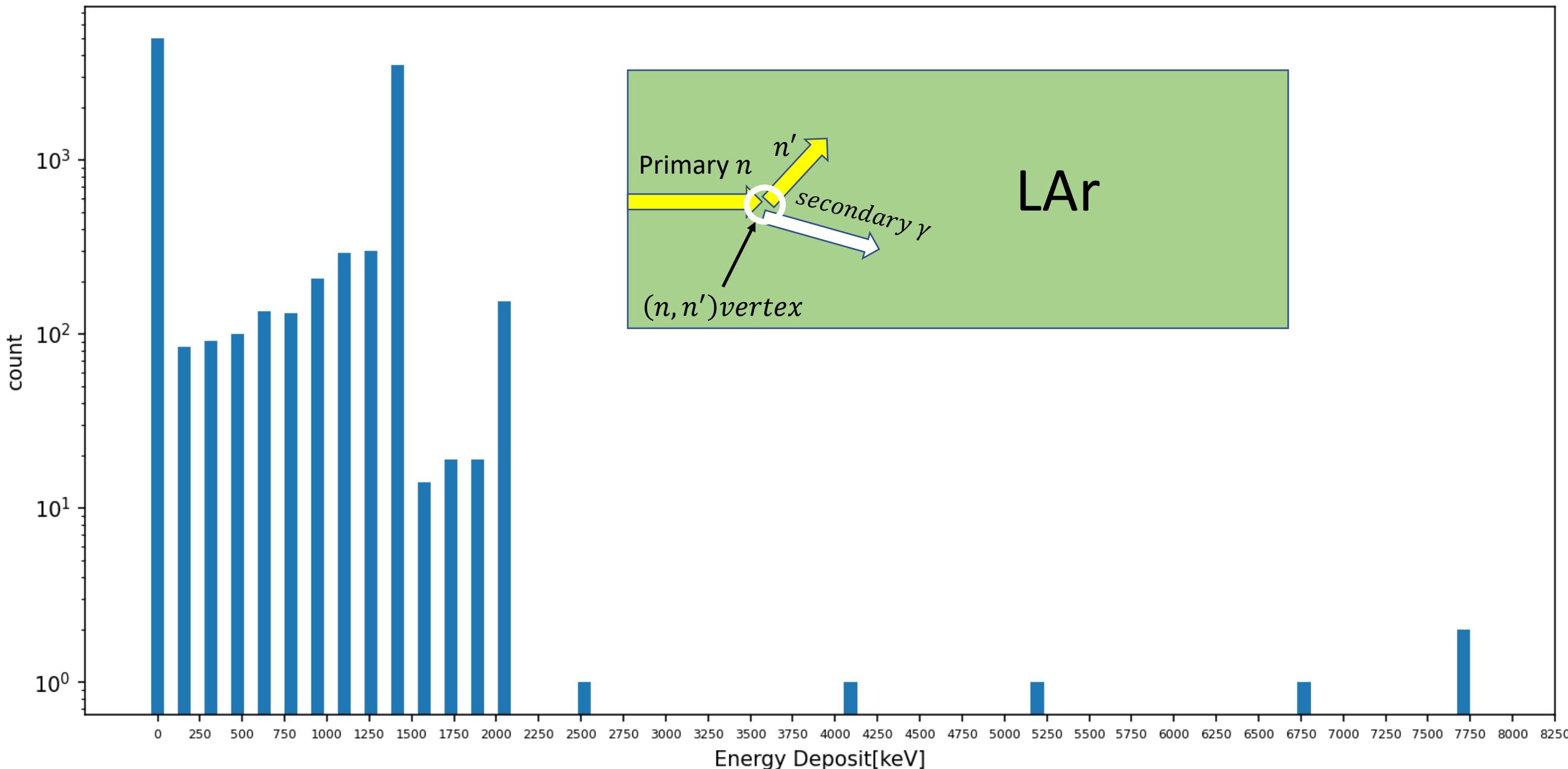
### photon emission time



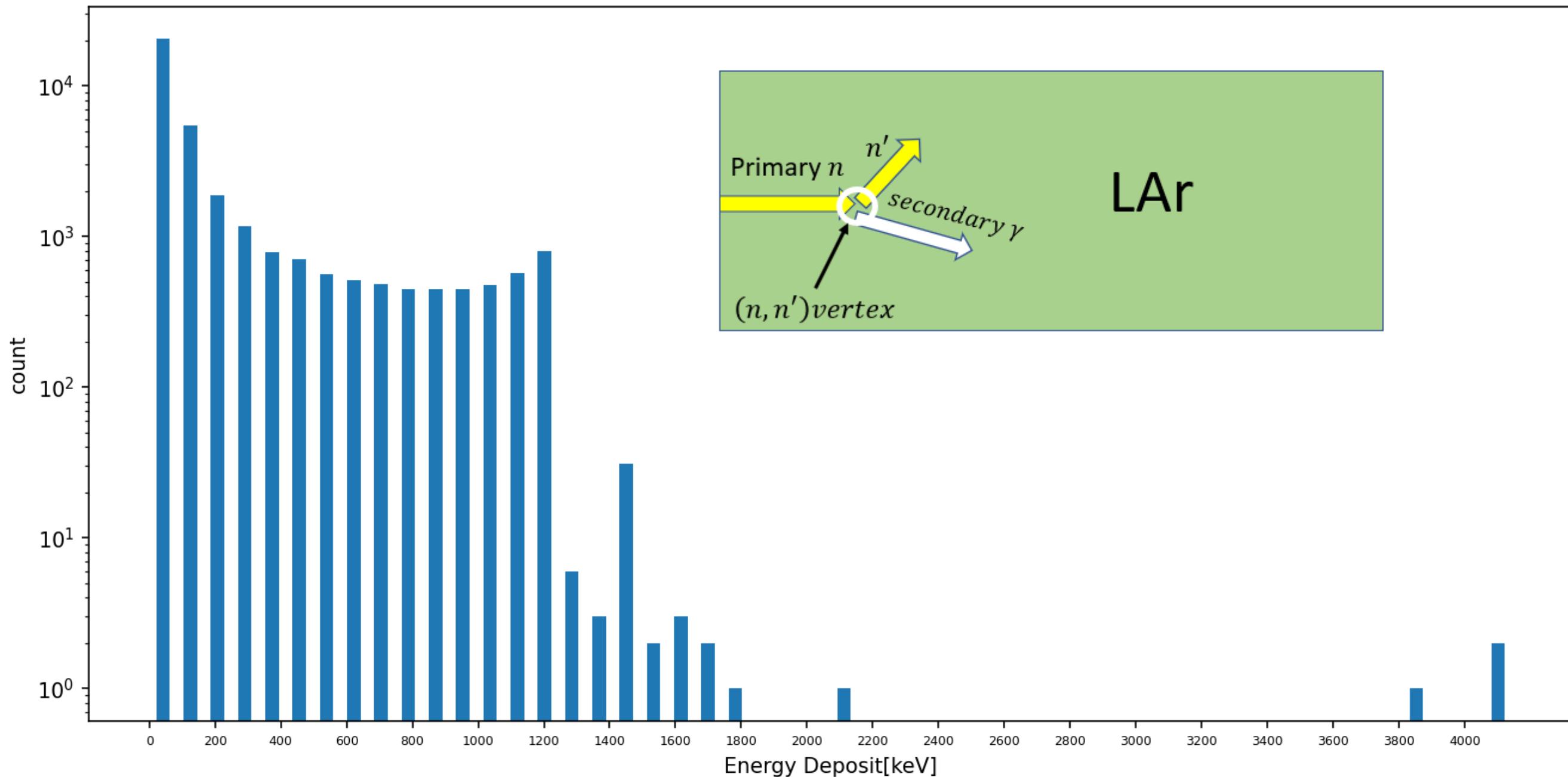
photon energy vs emission time



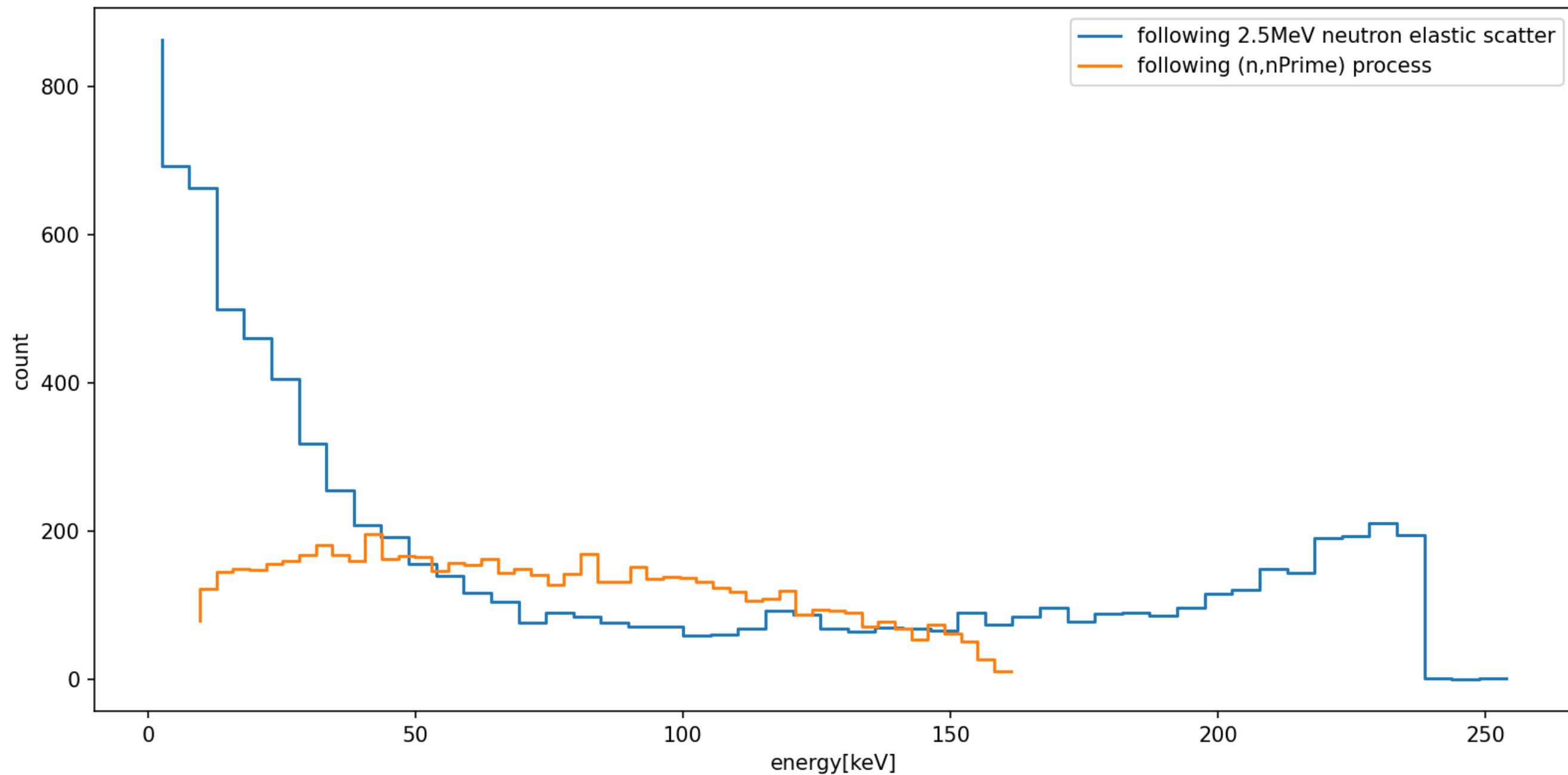
# secondary $\gamma$ energy deposit per event

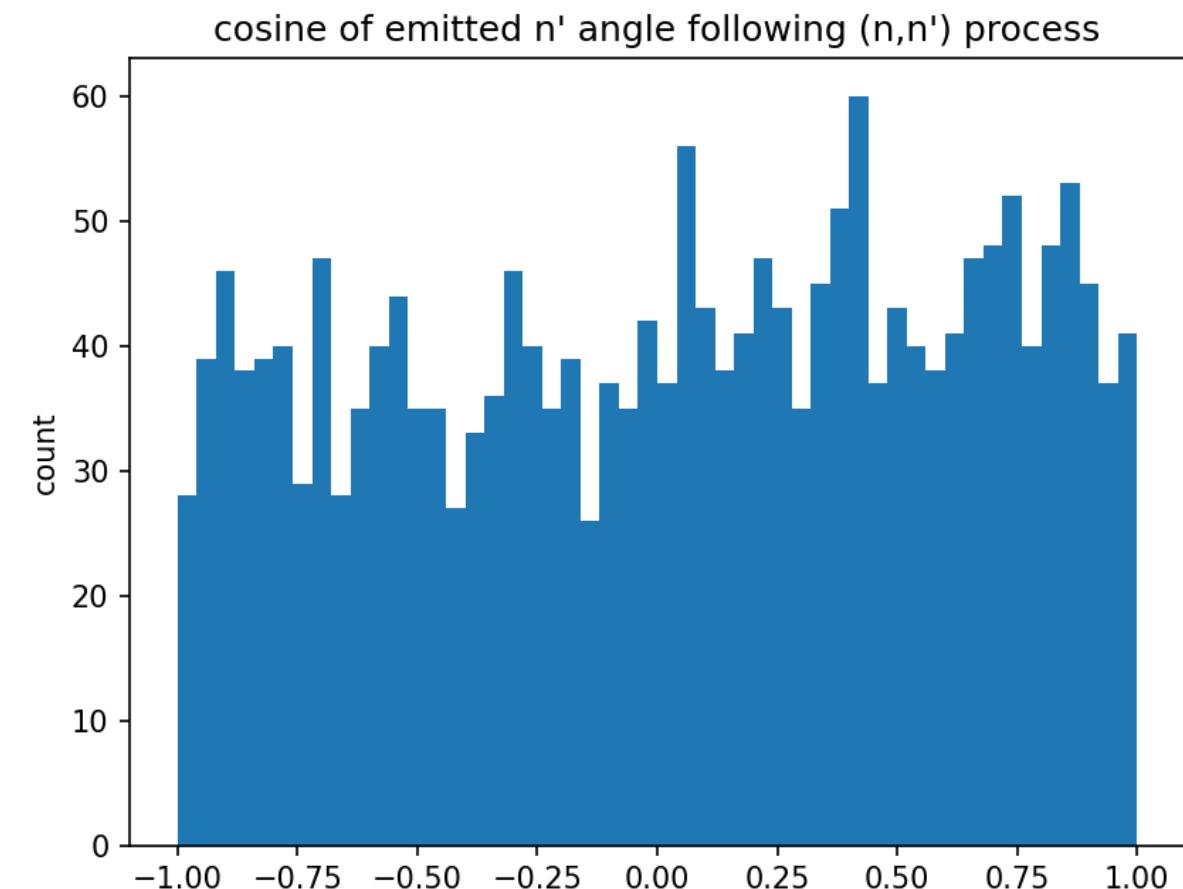
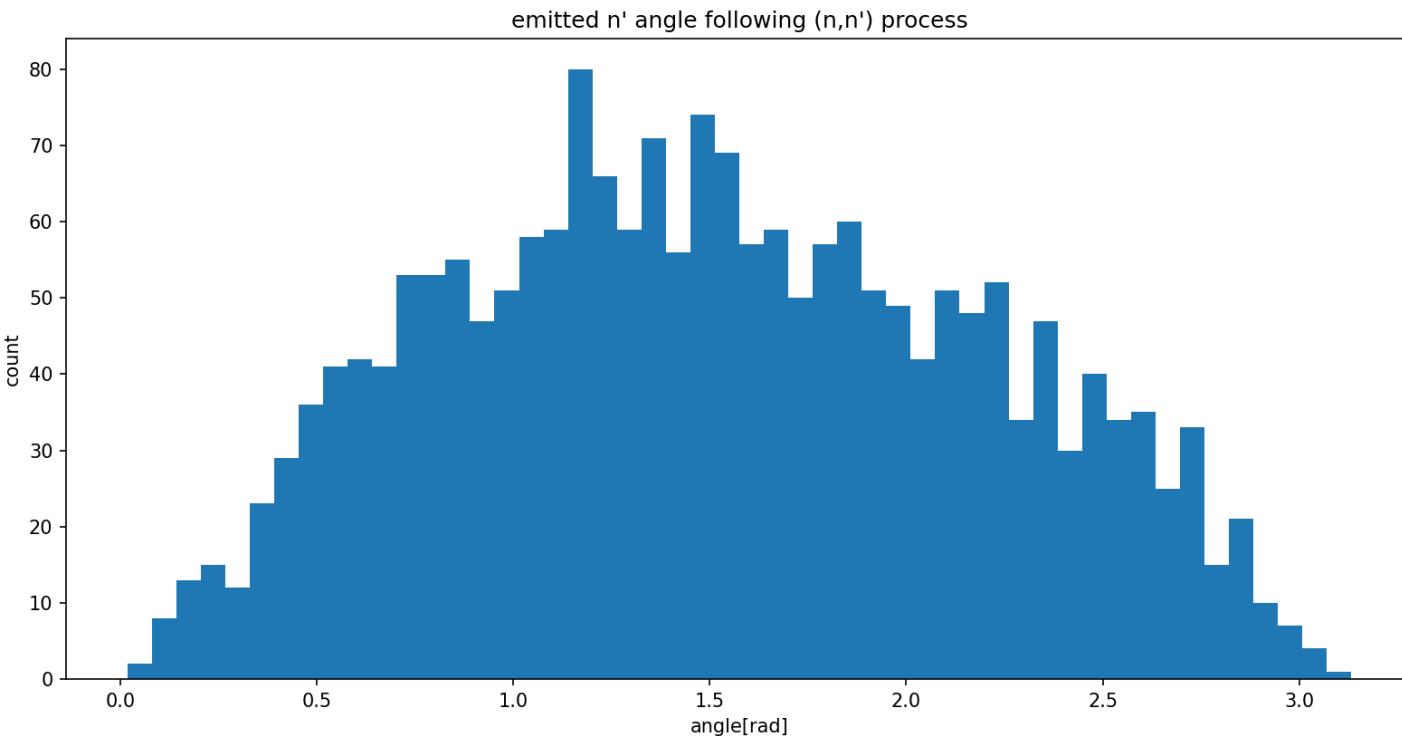


# secondary $\gamma$ energy deposit per $\gamma$

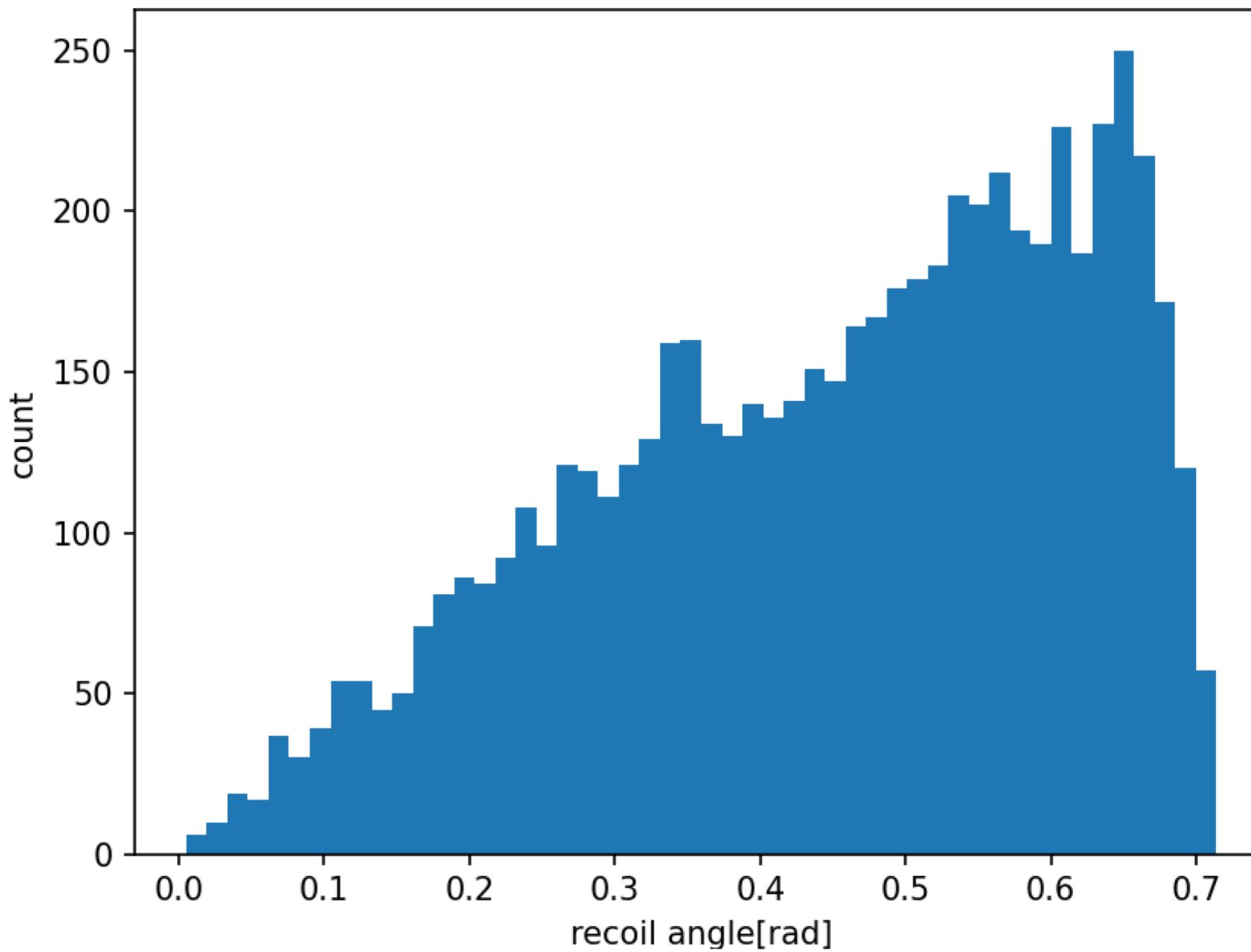


# $^{40}\text{Ar}$ nucleus recoil energies

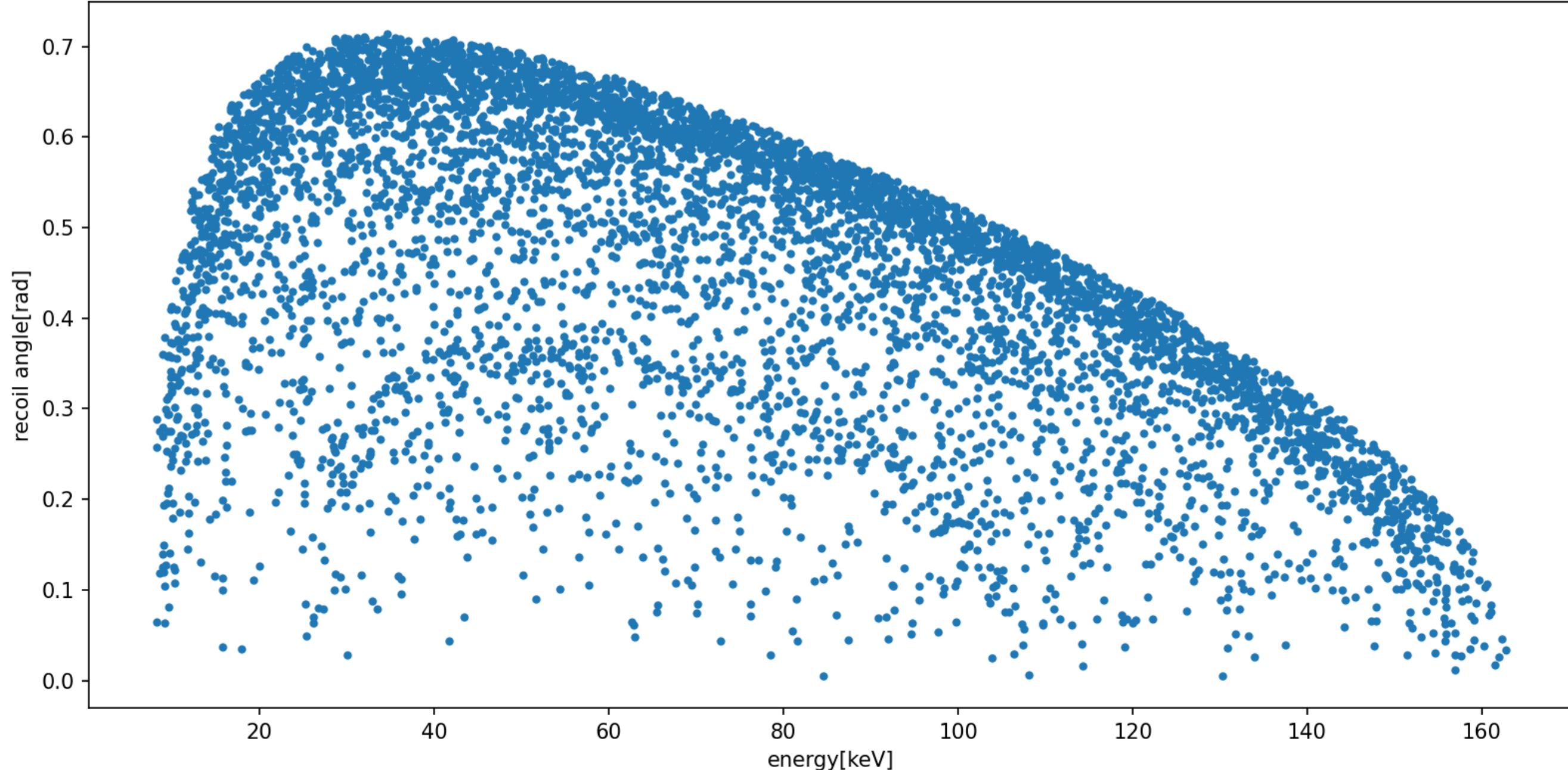




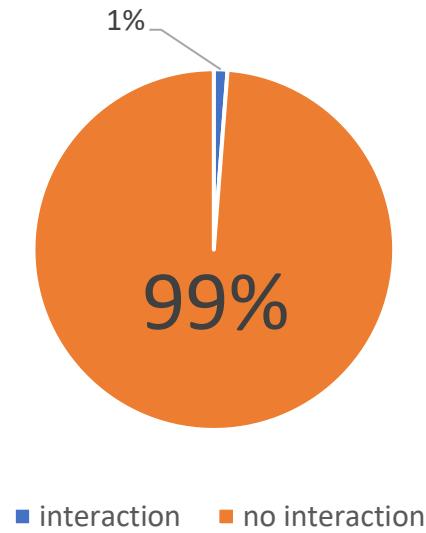
$^{40}\text{Ar}$  nucleus recoil angles following (n,n') process



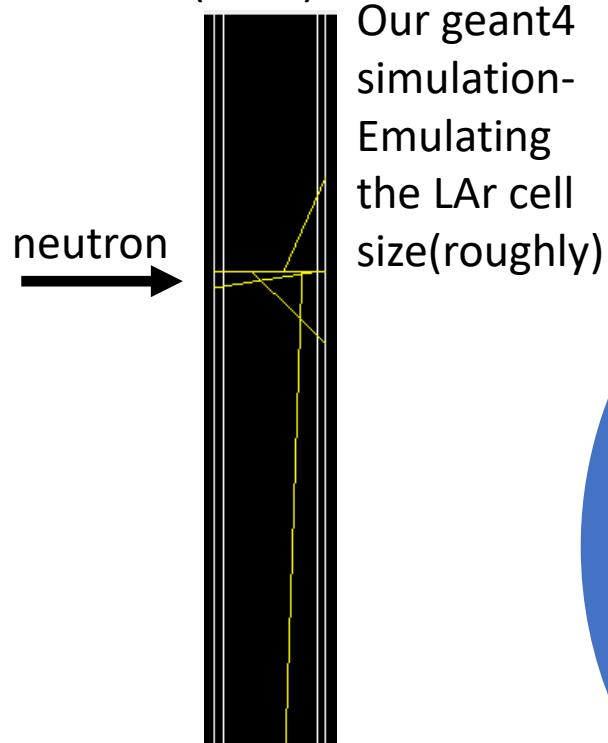
$^{40}\text{Ar}$  nucleus recoil energies vs recoil angle following (n,n') process



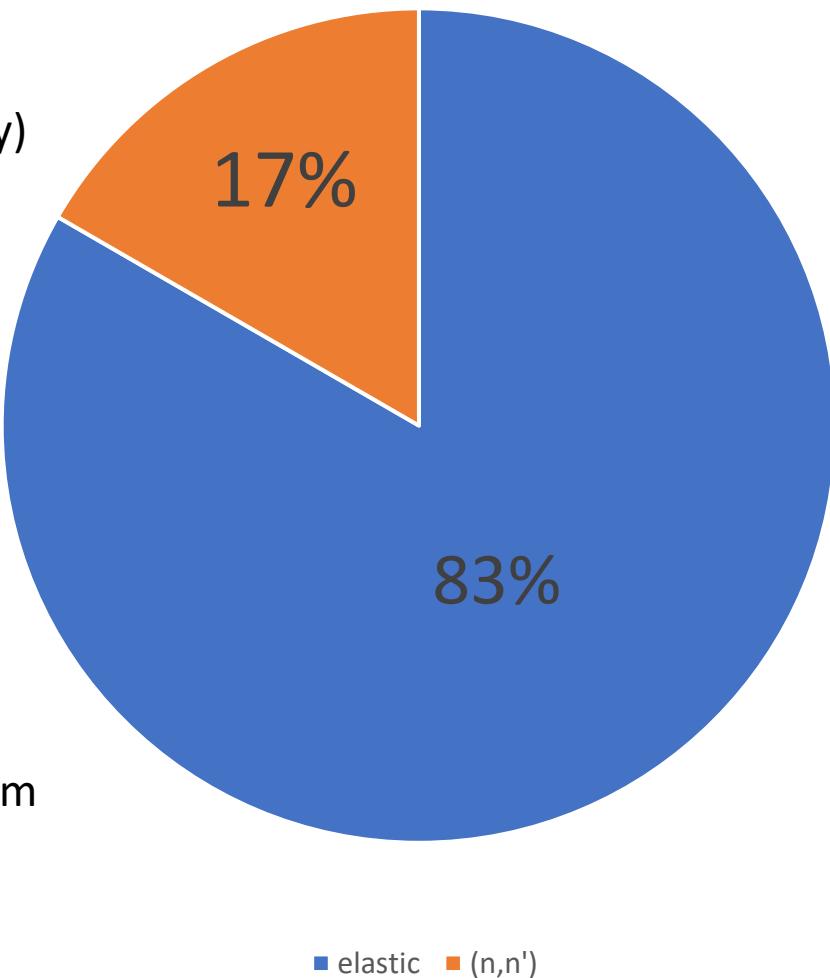
## Interaction statistics



■ interaction ■ no interaction



## Interaction statistics



■ elastic ■  $(n,n')$

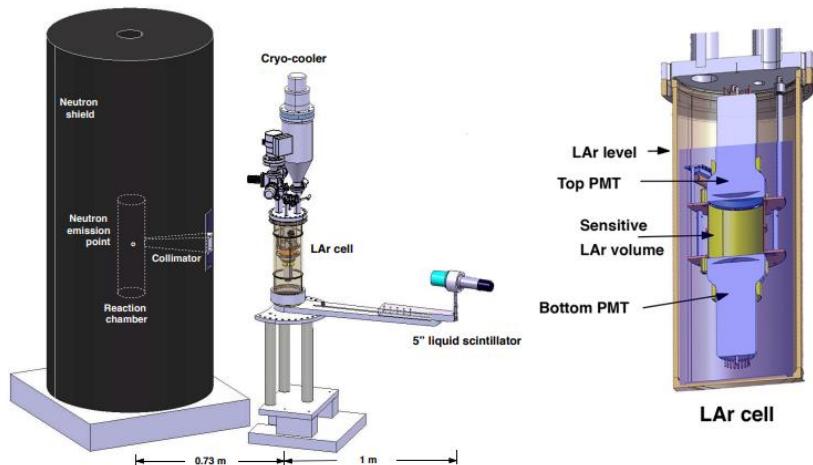


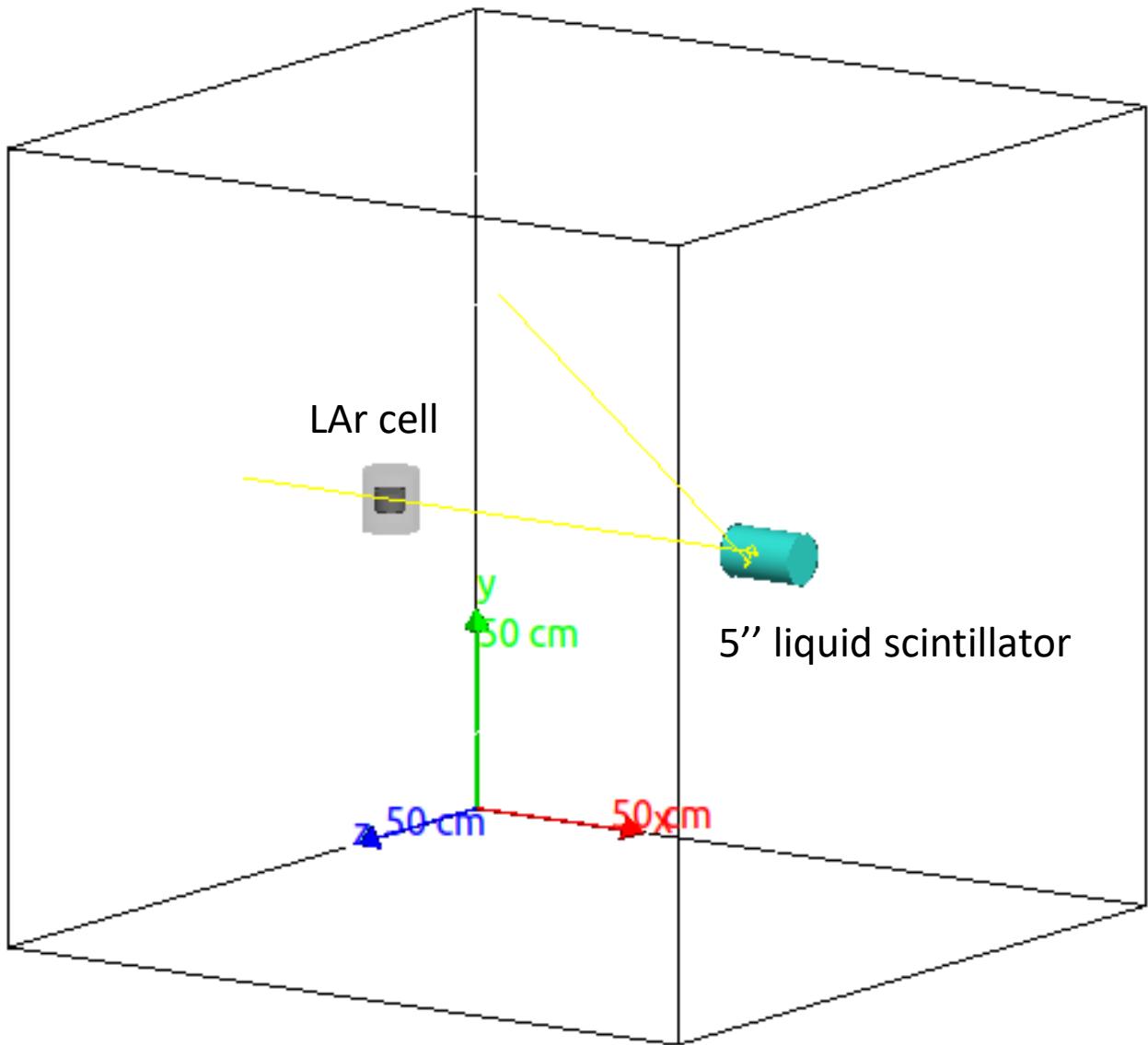
Figure 3. Left: neutron gun, argon cell and liquid scintillation counter used to measure  $L_{eff}$  in LAr. Right: sketch of the LAr cell with its vacuum chamber (see text).

clear recoils, relative to electrons, between 11 keV and 120 keV in liquid argon. Single elastic neutron-argon scattering dominates thanks to the small active volume of our argon cell and the well collimated neutron beam, while the contamination from neutron inelastic scattering is negligible. A  $\chi^2$  minimization is performed leaving  $L_{eff}$  and the energy resolution as free parameters. The

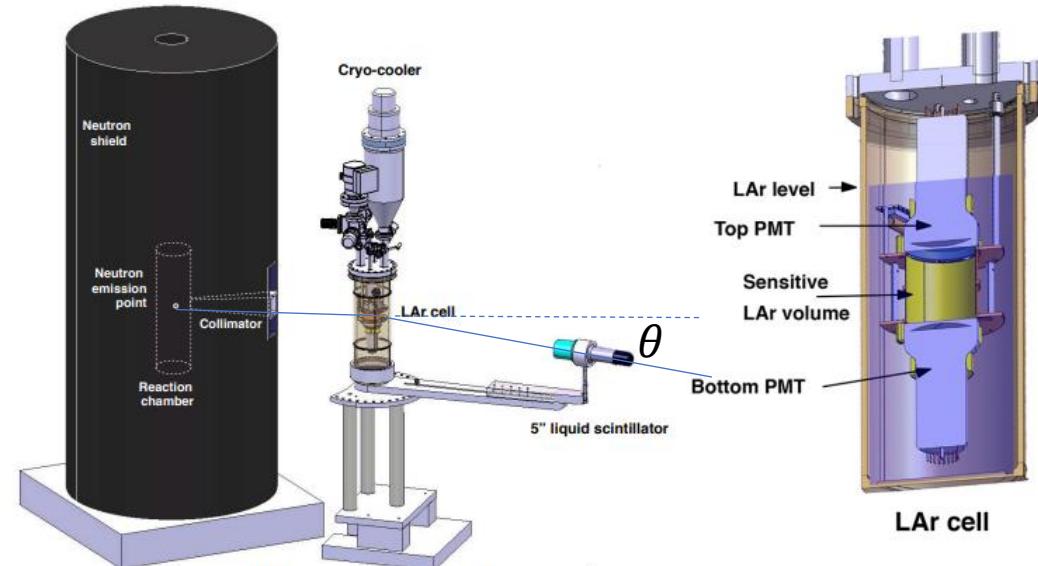
?

Consequence -  
In the paper they neglect inelastic scatters because they use elastic kinematics in order to calculate nucleus recoil, and thus in 17% of events they have an error in their calculations.

## My geant4 simulator



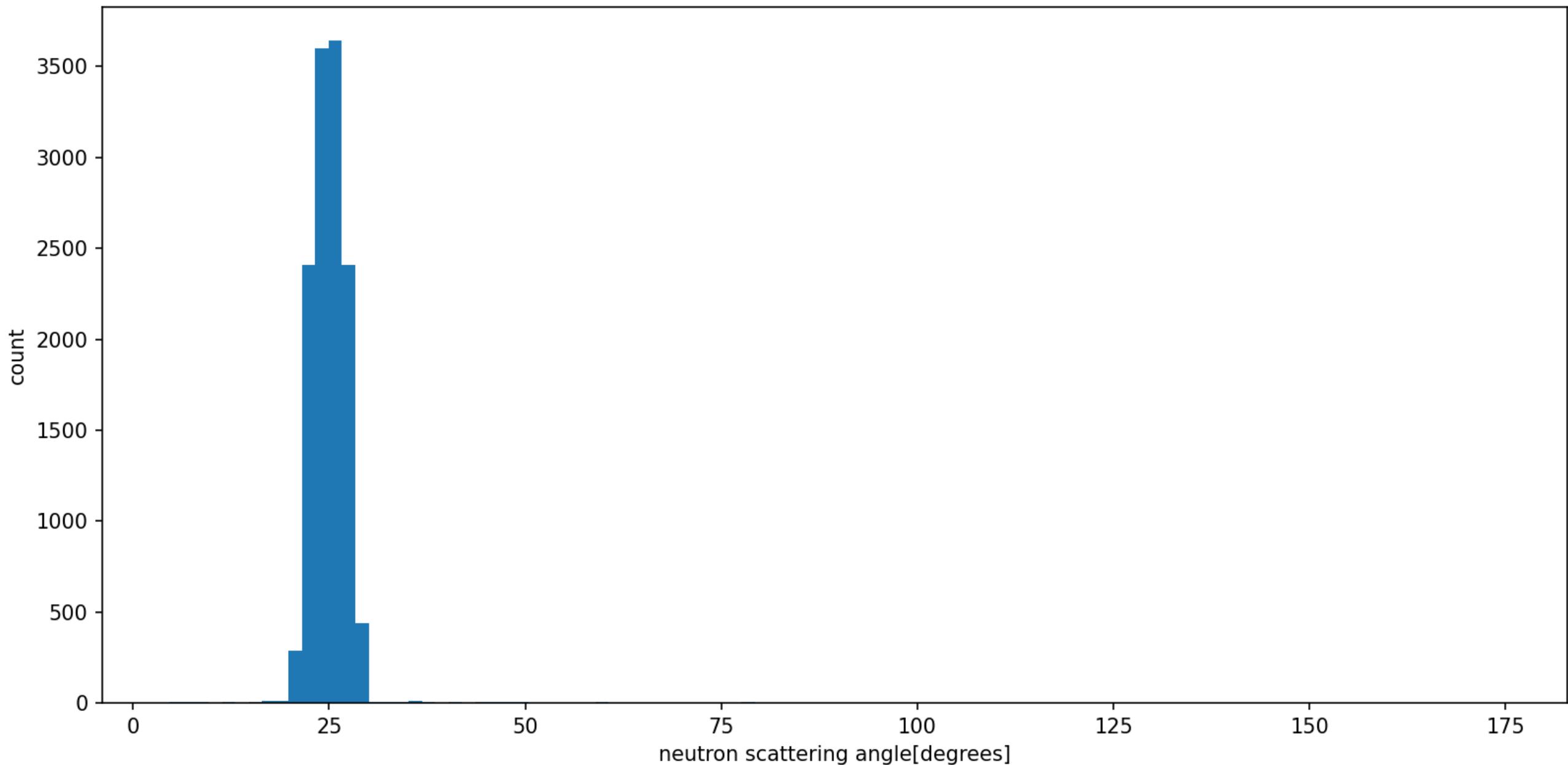
## Article experimental setup



**Figure 3.** Left: neutron gun, argon cell and liquid scintillation counter used to measure  $L_{eff}$  in LAr. Right: sketch of the LAr cell with its vacuum chamber (see text).

in the experiment there is also a 1.5mm aluminium layer around the sensitive volume of LAr, but the effect on the neutrons is minimal(sub percent)

Also, in reality the beam is collimated into 0.2% 4pi steradians which gives an effective surface area of 0.013 [ $cm^2$ ] of neutrons hitting the LAr, and in my simulator I use a pencil beam (0 width)



## Article experimental setup

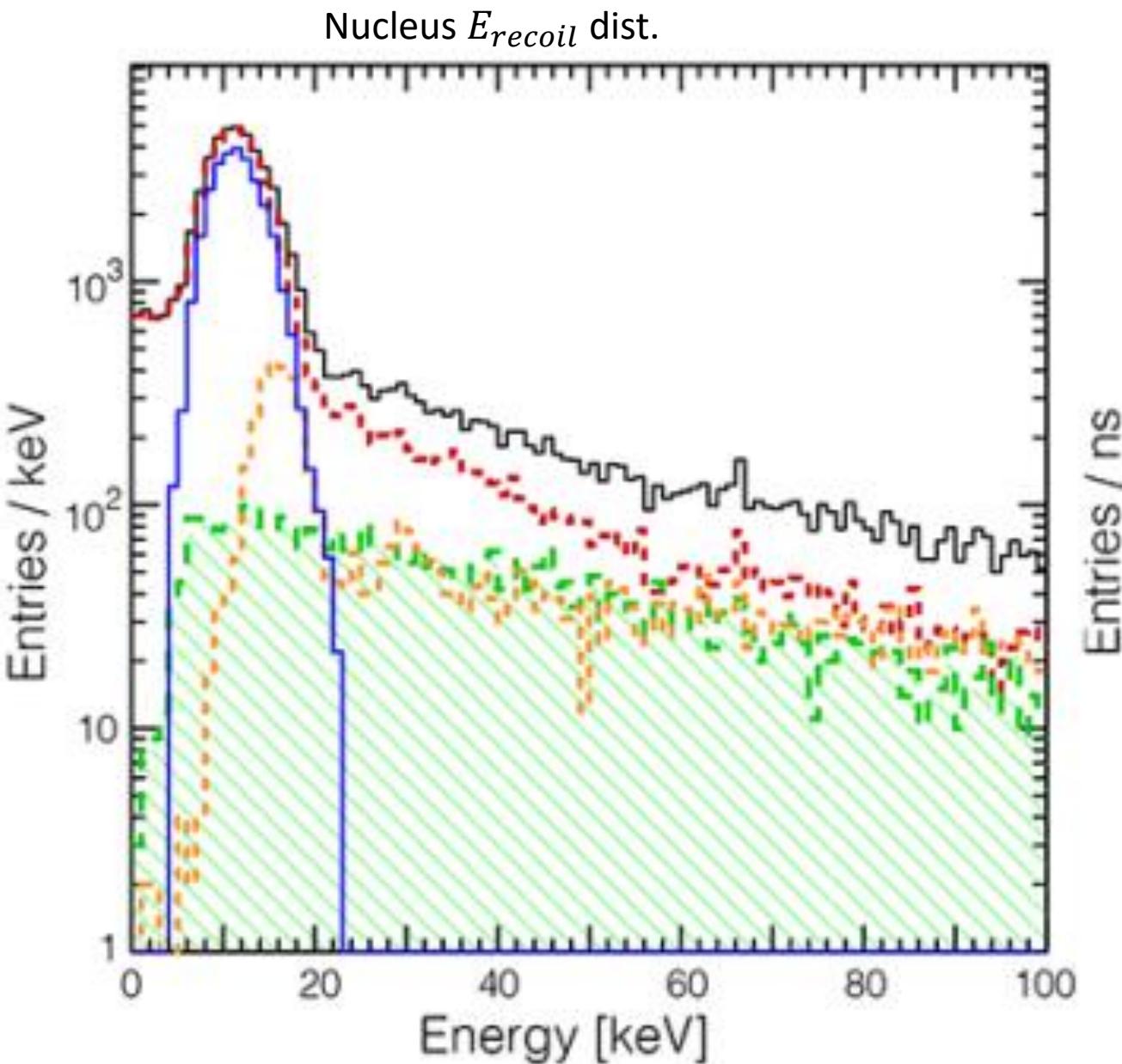
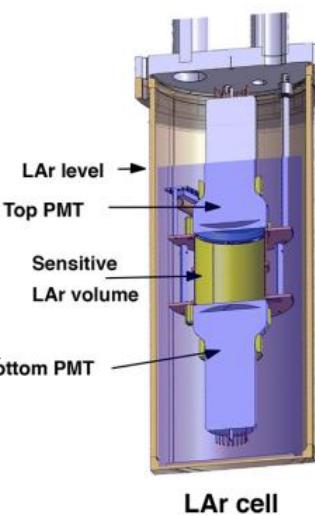
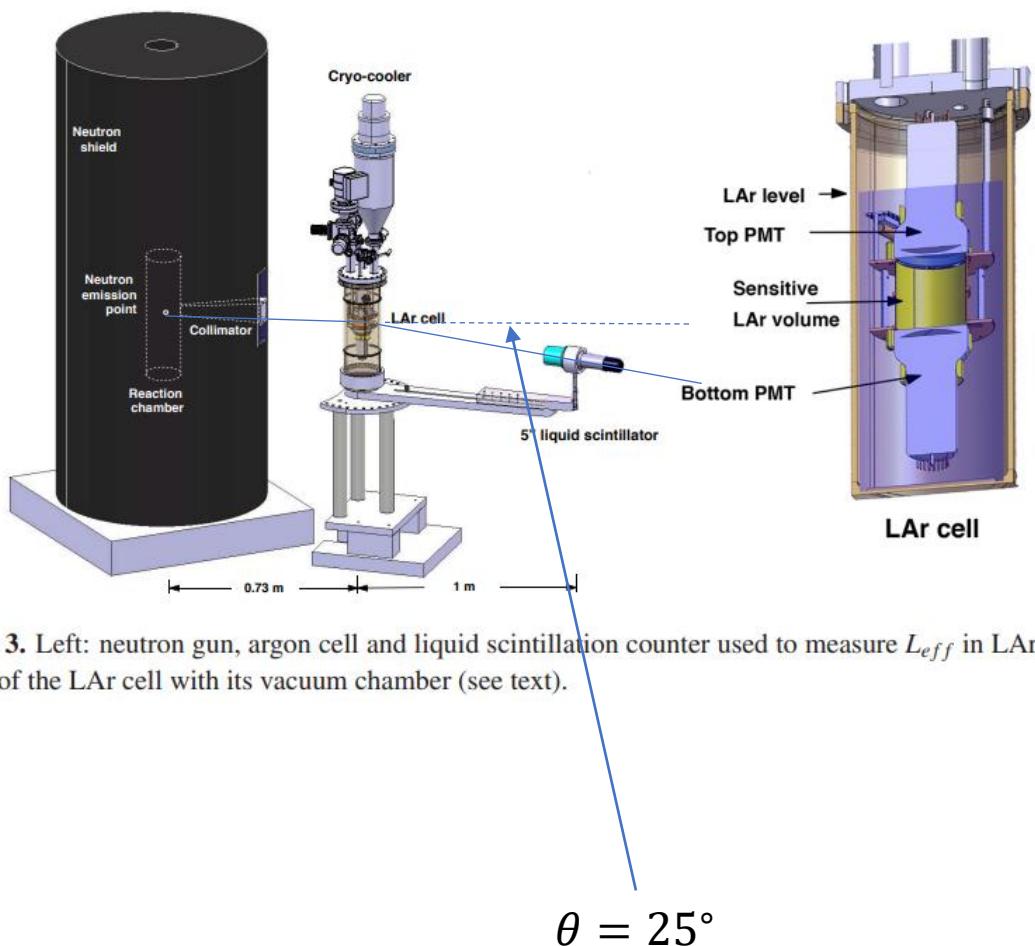
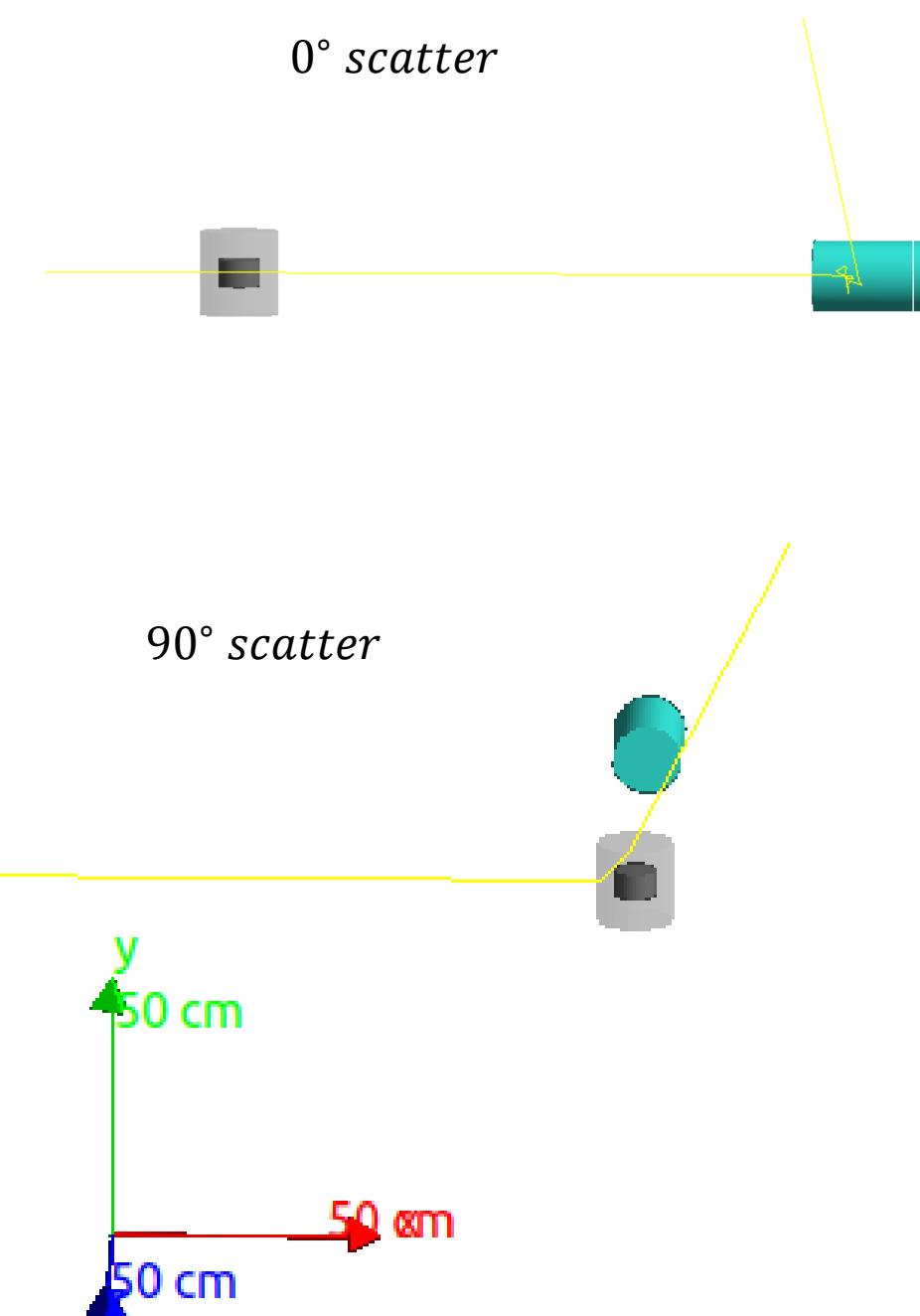
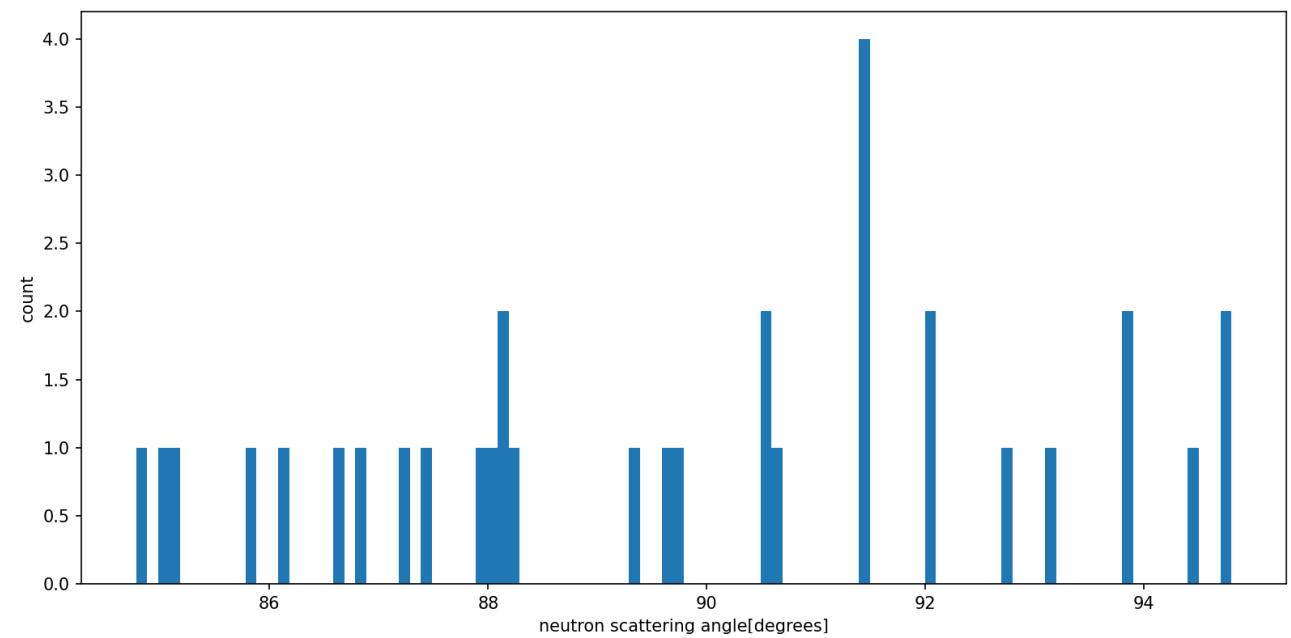
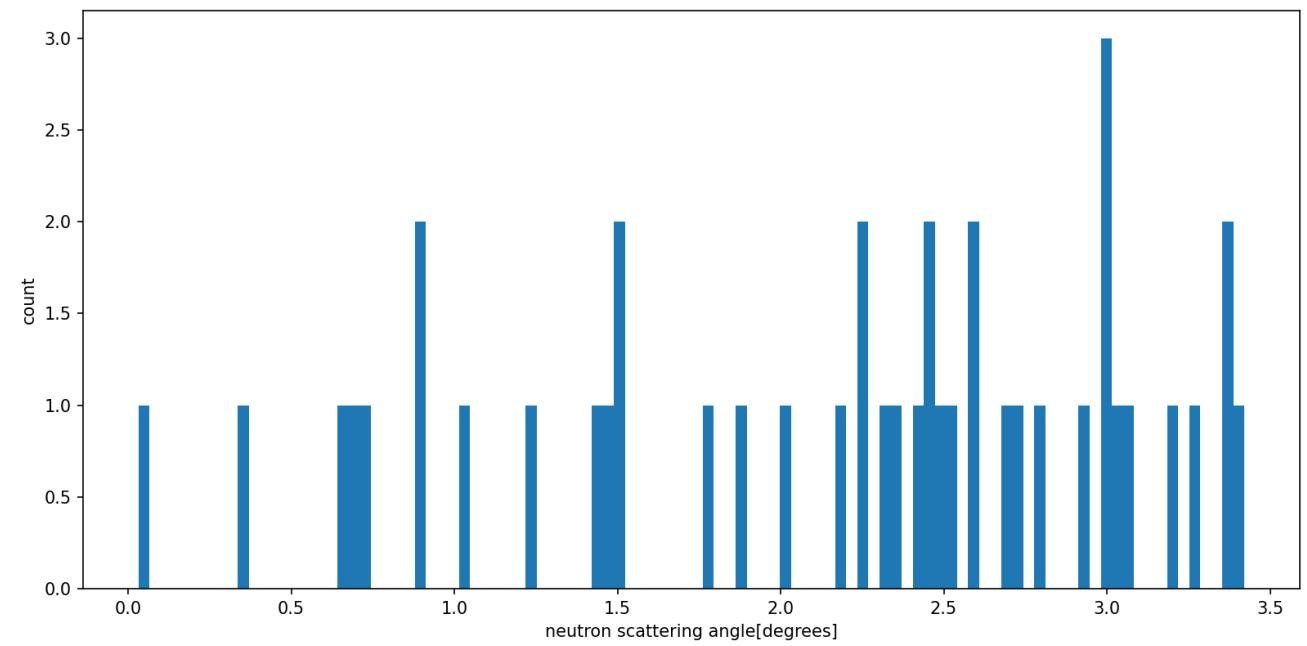
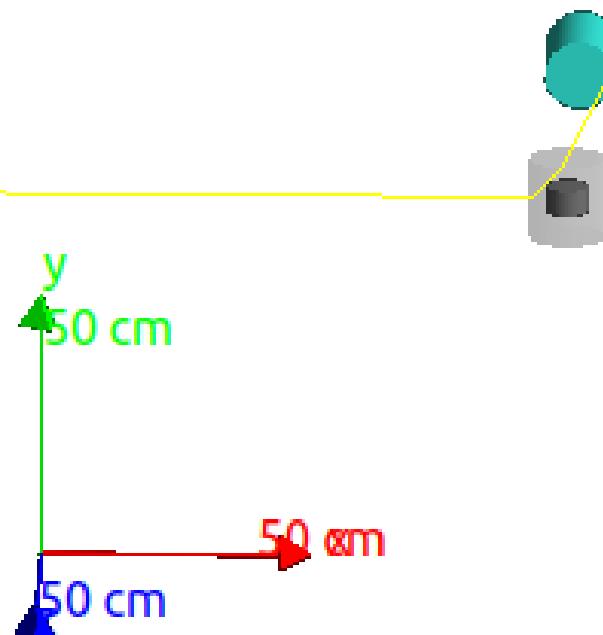


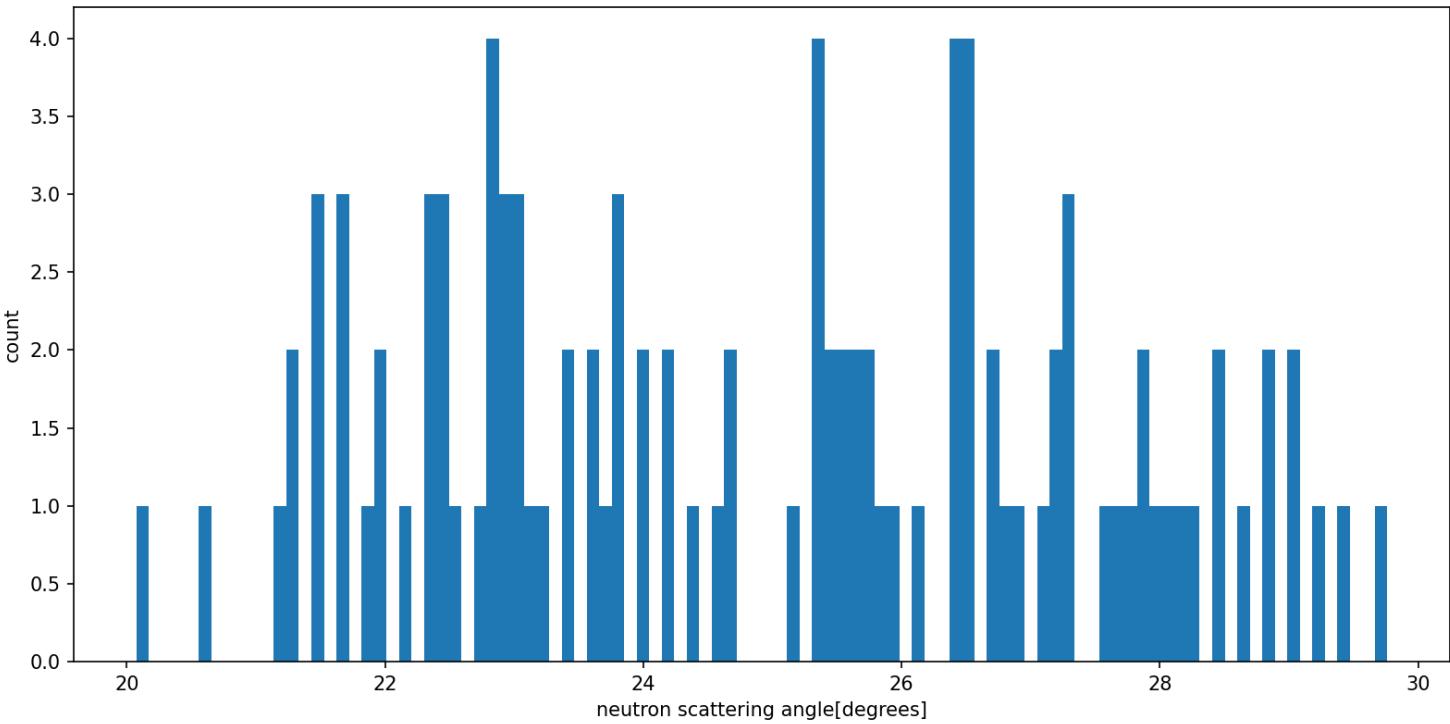
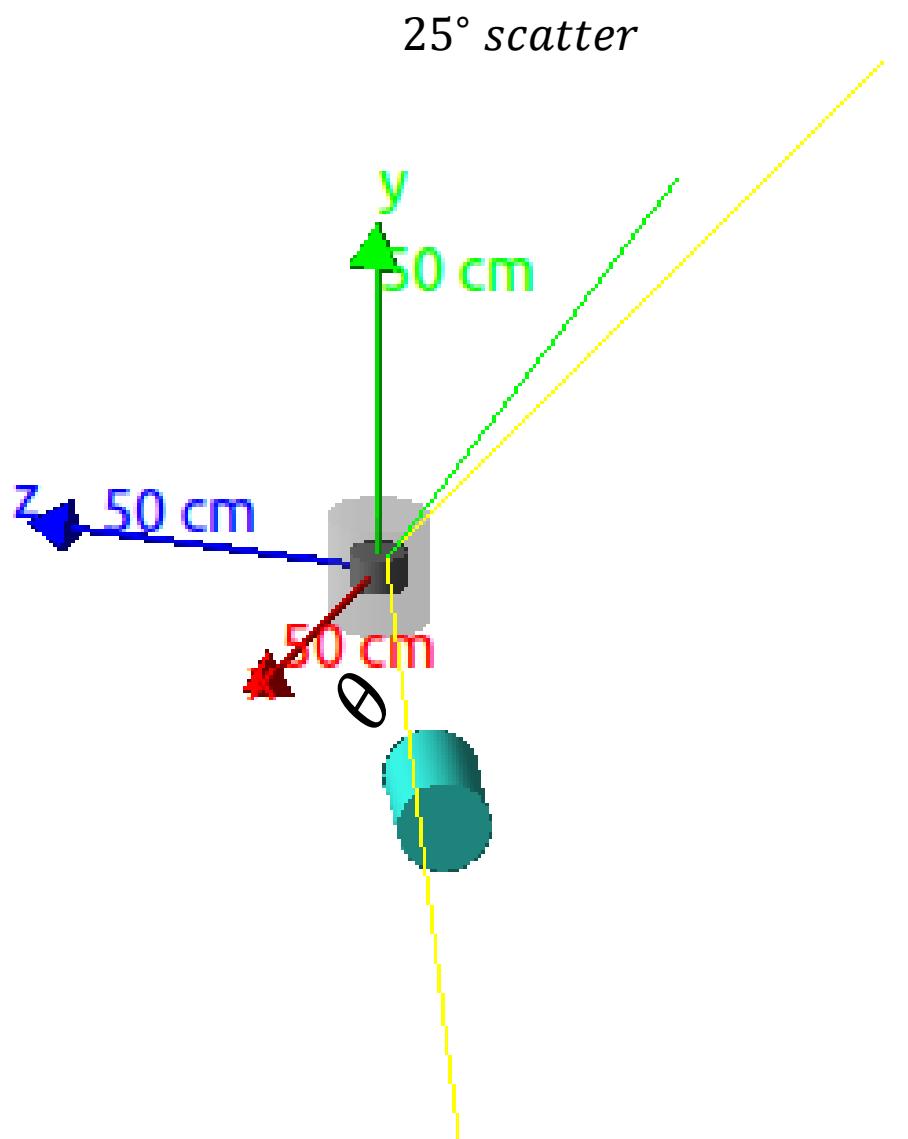
Figure 3. Left: neutron gun, argon cell and liquid scintillation counter used to measure  $L_{eff}$  in LAr. Right: sketch of the LAr cell with its vacuum chamber (see text).

$0^\circ$  scatter

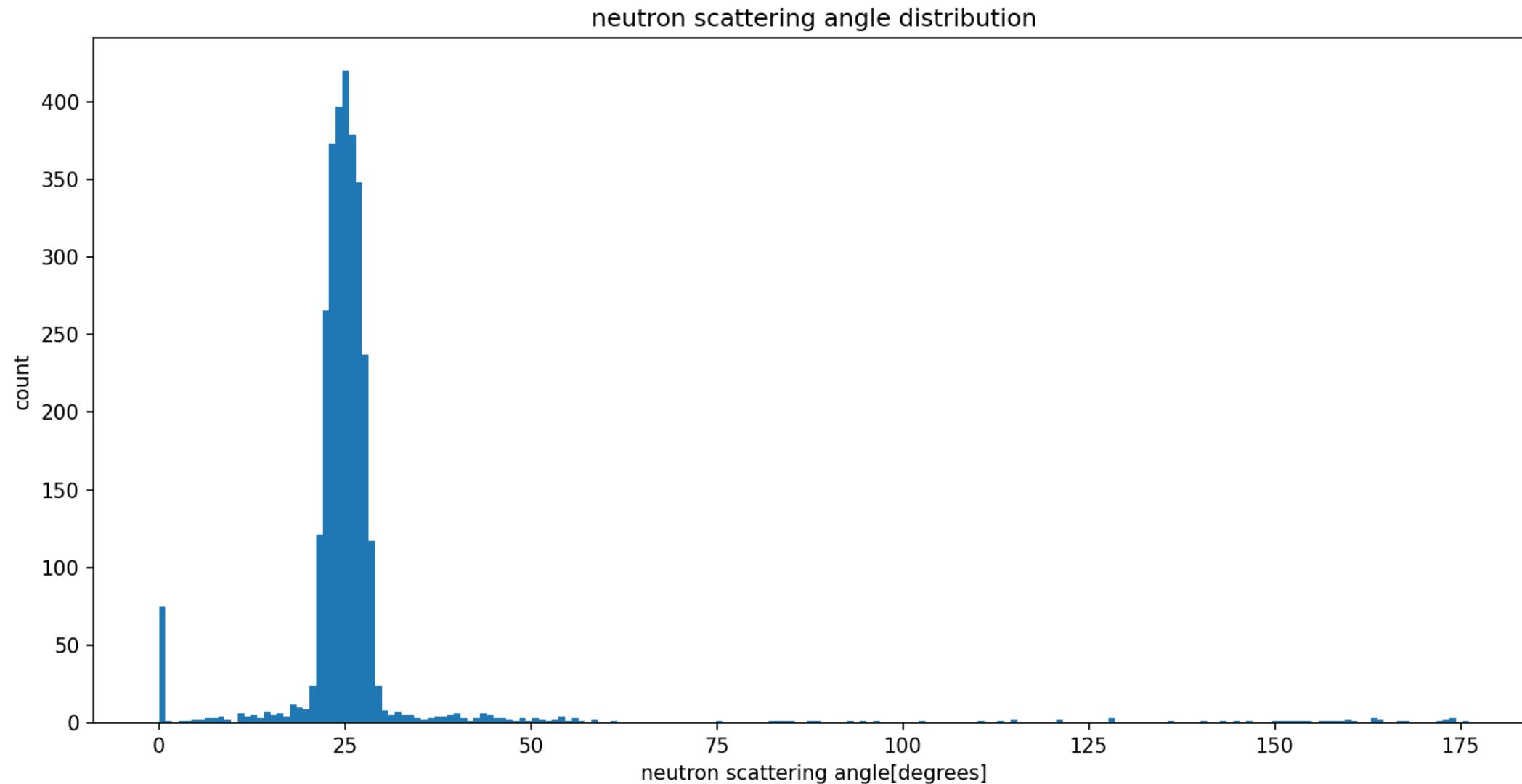


$90^\circ$  scatter





Liquid scintillator placed at 25° angle relative to LAr cell,  
Showing the neutron scattering angle distribution within the active volume of the cell without external scatters

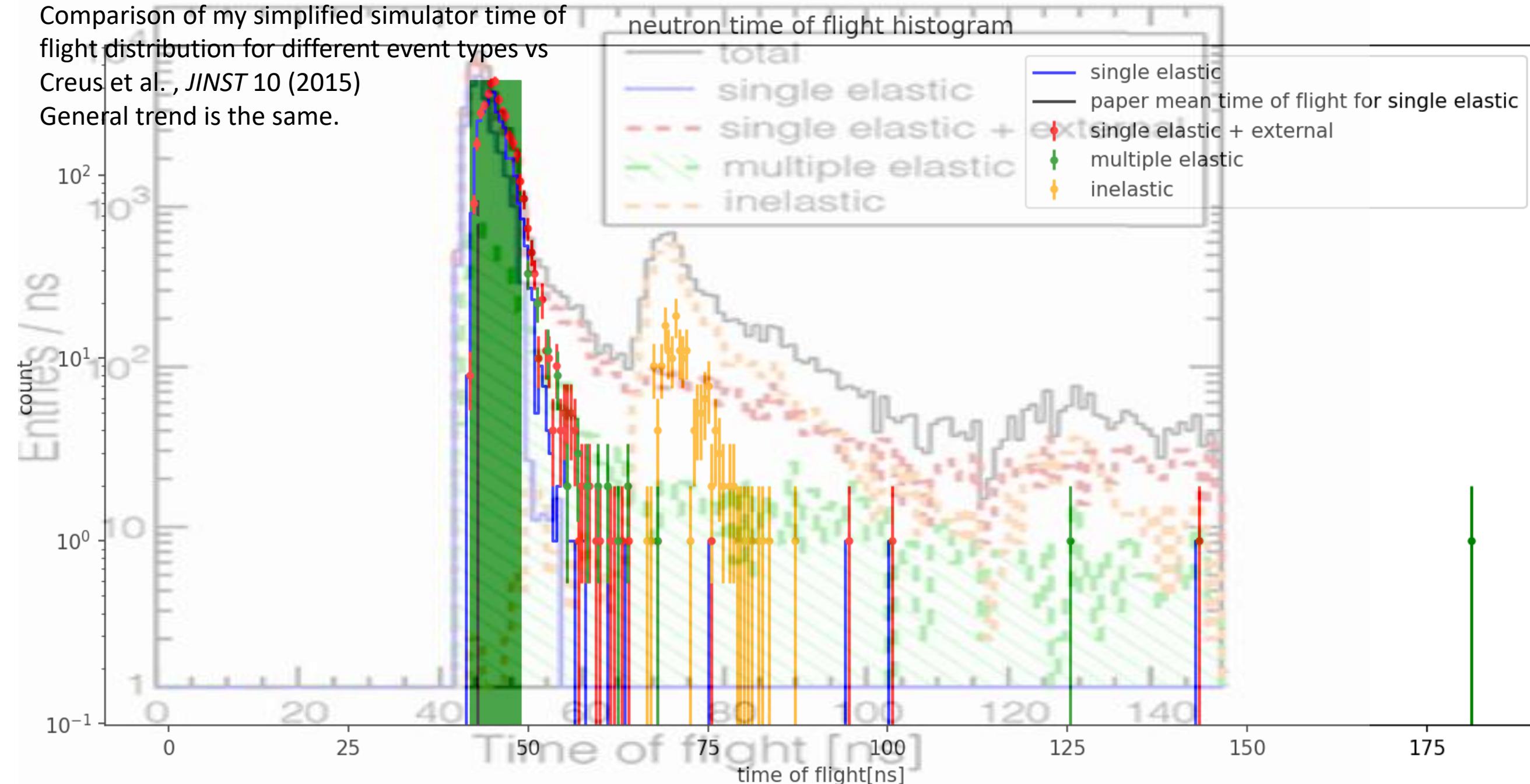


Comparison of my simplified simulator time of flight distribution for different event types vs

Creus et al. , JINST 10 (2015)

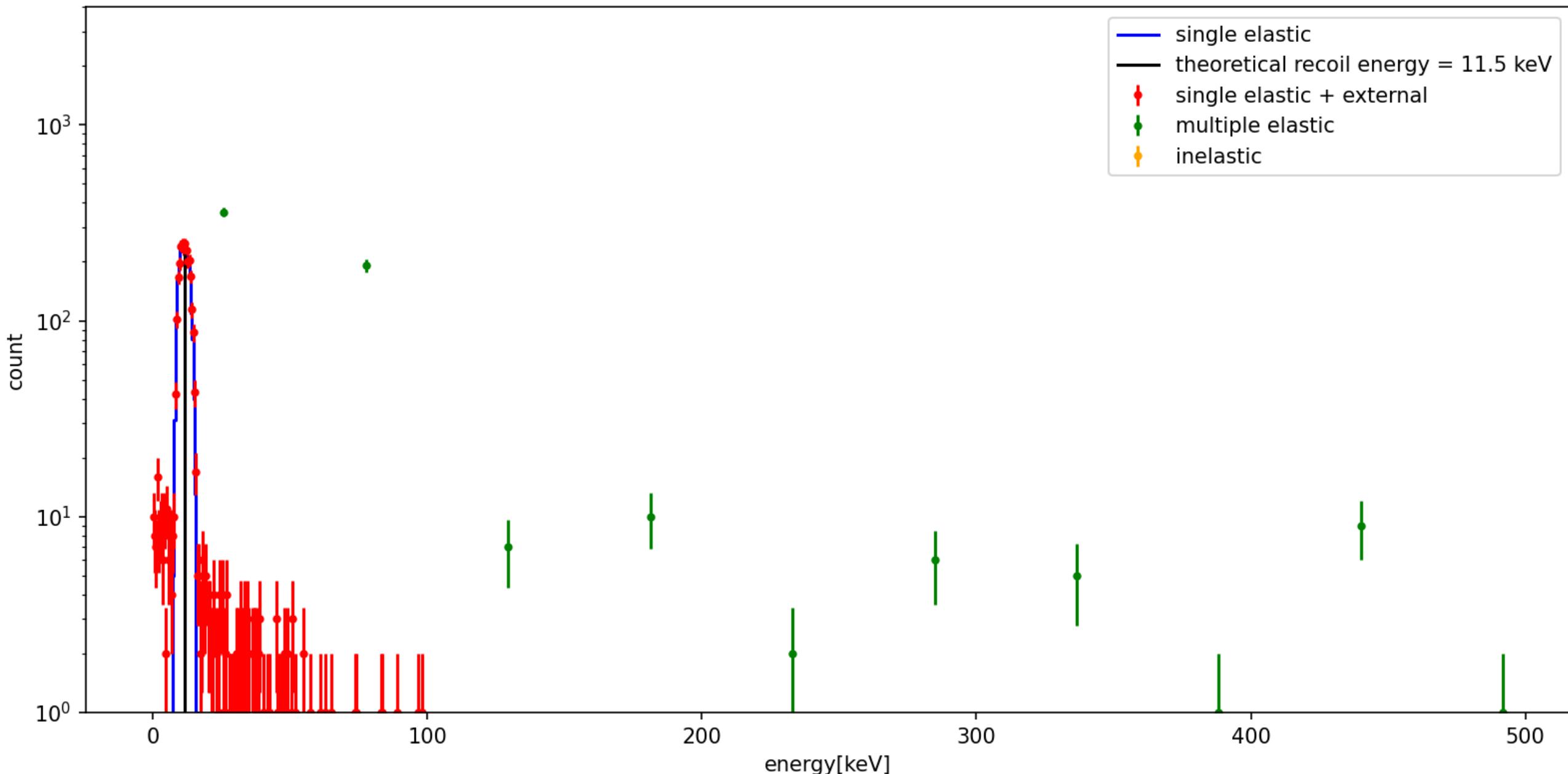
General trend is the same.

neutron time of flight histogram

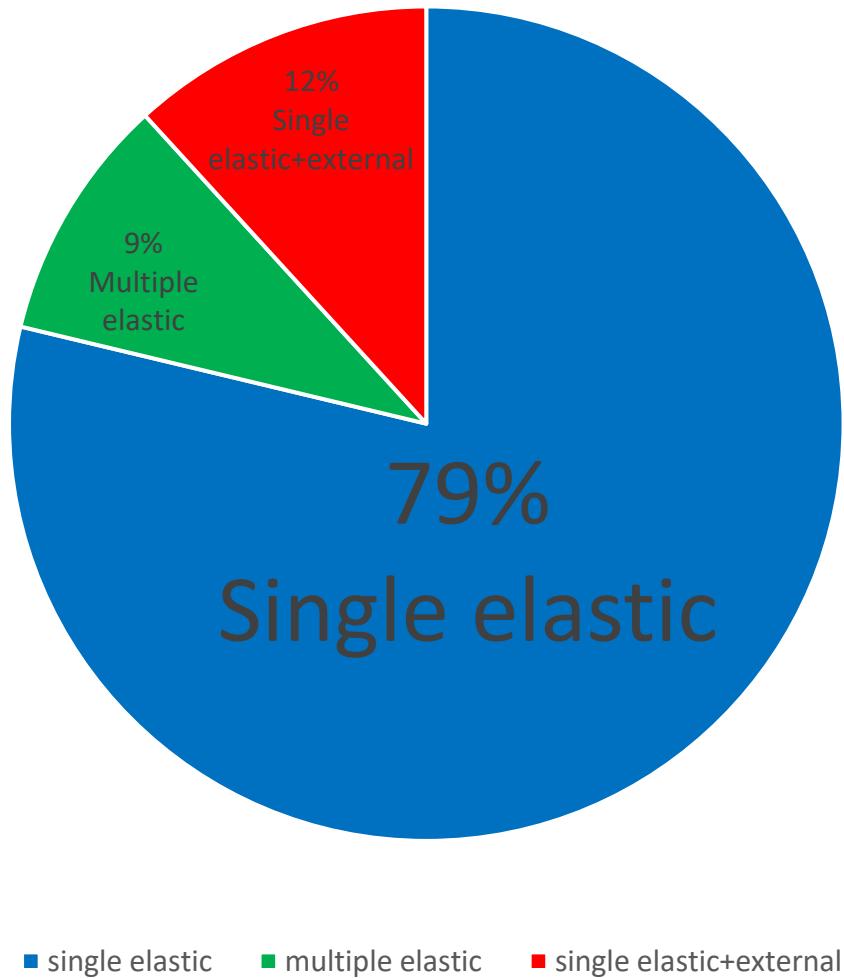


After time of flight cut -  $42 \text{ [ns]} < \text{time} < 49 \text{ [ns]}$

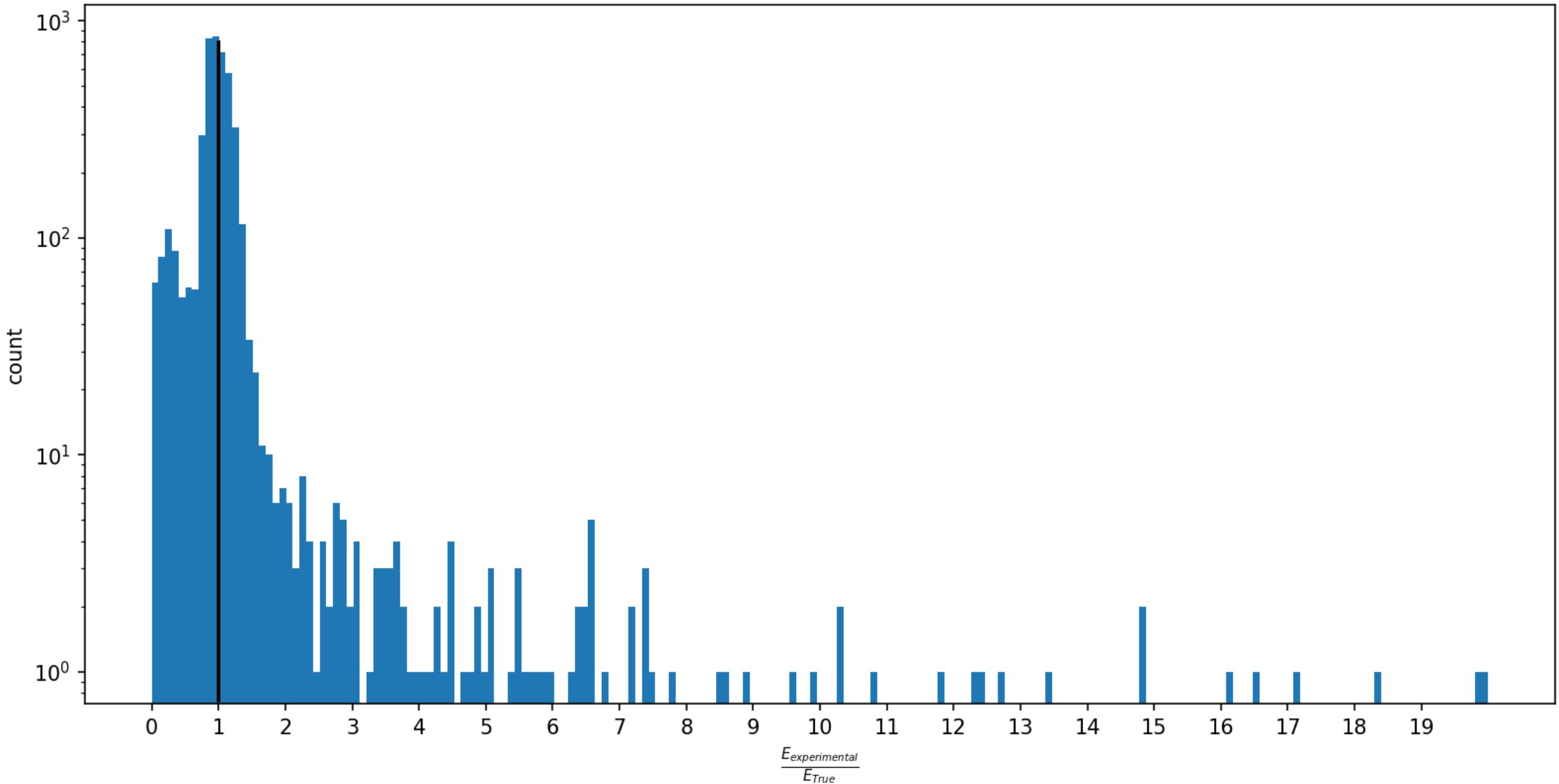
recoil nucleus energy distribution following elastic scatter with  $\theta = 25^\circ$

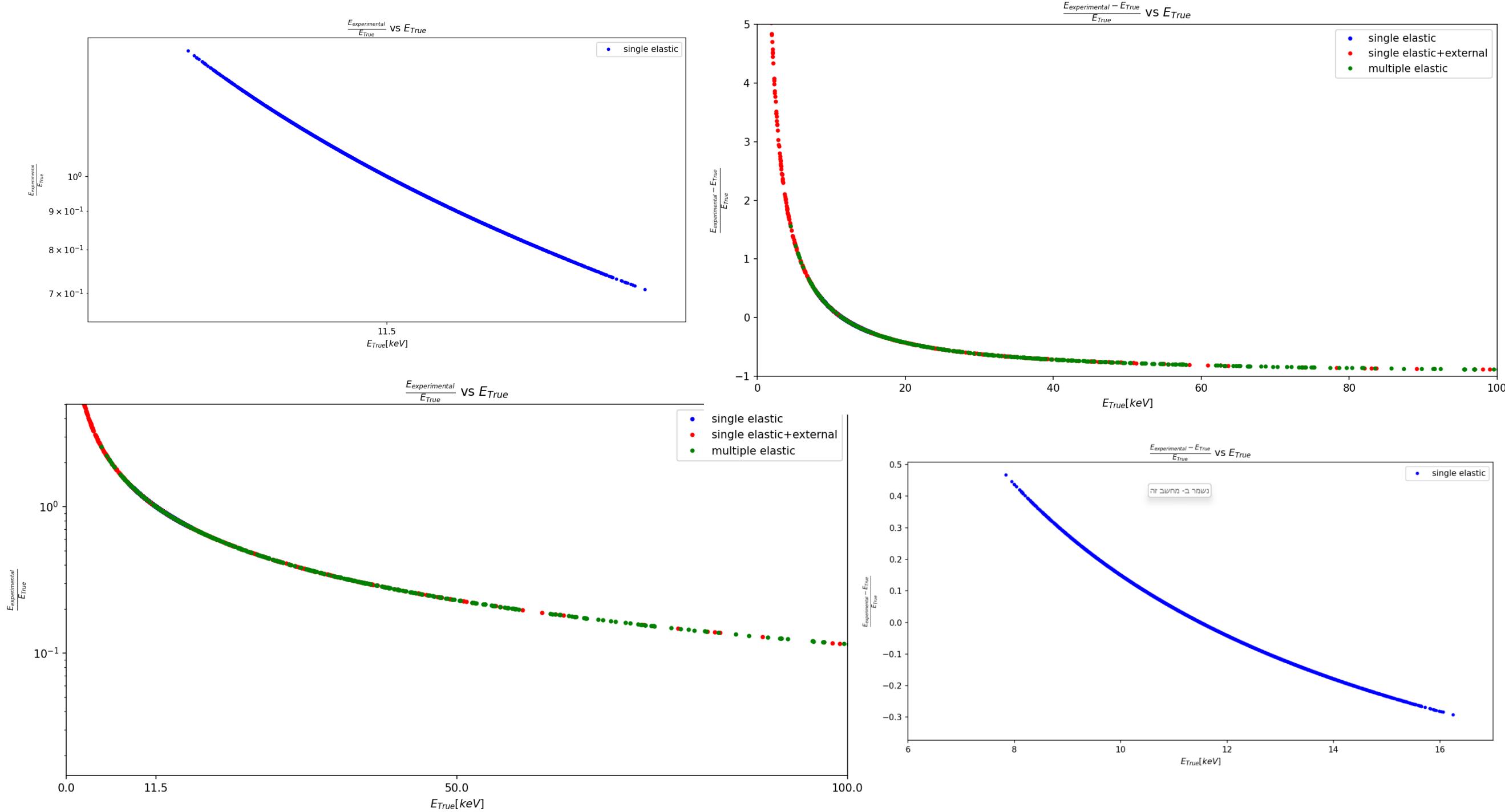


Event Types after time of flight cut – 25 degree scatter

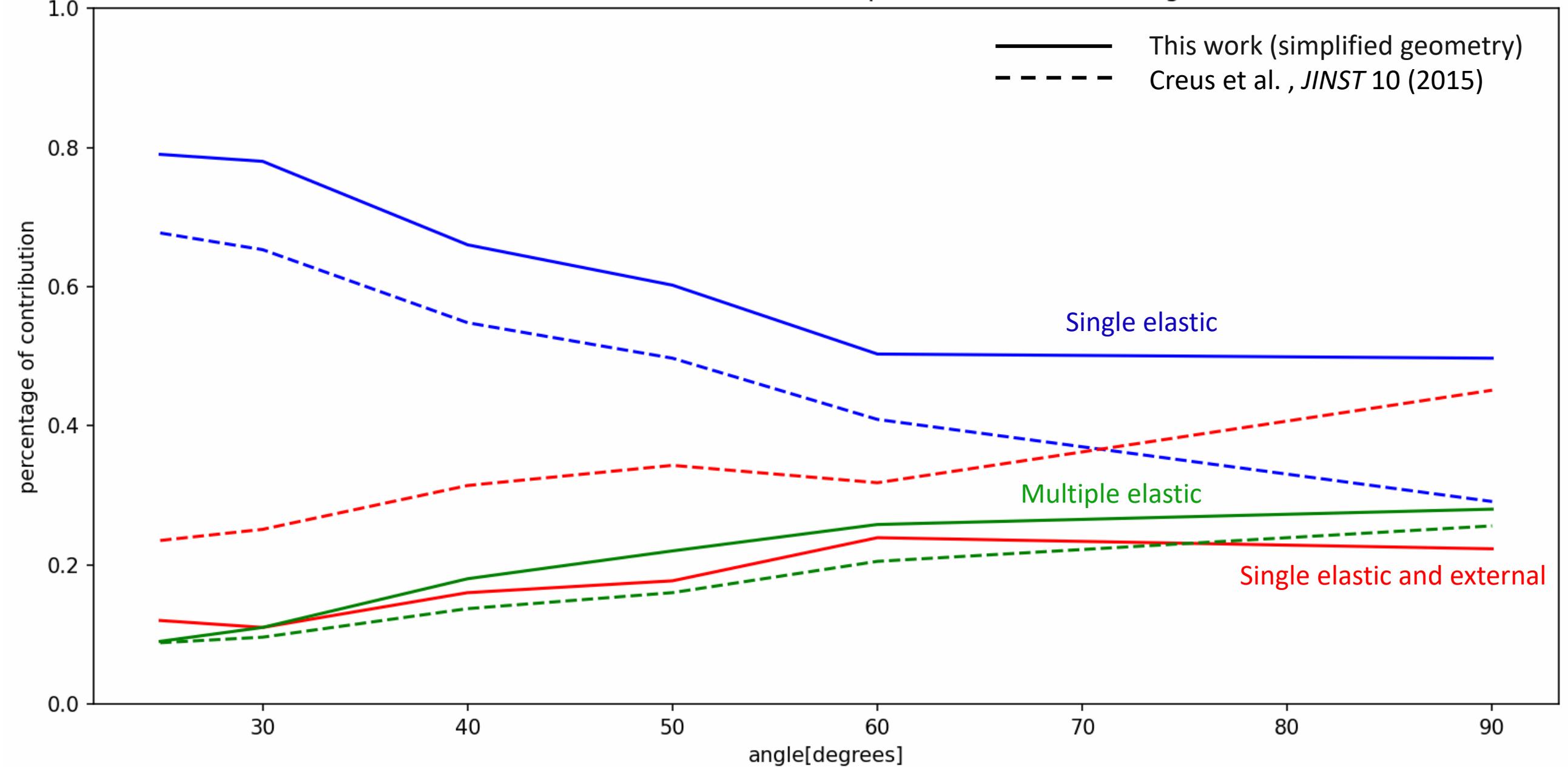


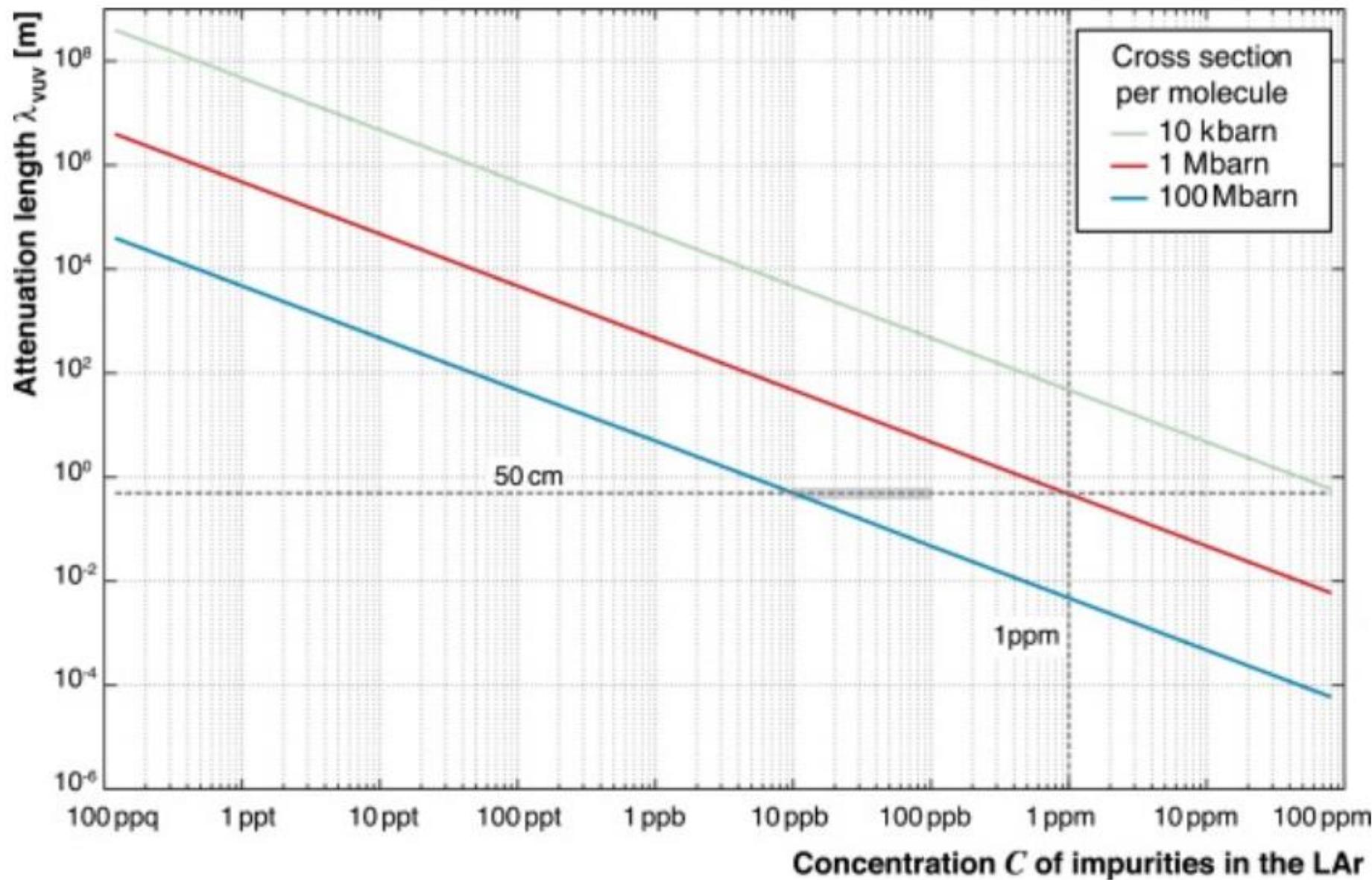
ratio of reconstructed recoil energy and true recoil energy per event





### Fractional contributions to the recoil spectrum after time-of-flight cut





In order to make sure we understand how to properly set the required parameters for scintillation light in geant4, we first want to make some sanity checks.

First and foremost, we want to see that the distribution of the amount of scintillation photons created matches the parameter we set for the SCINTILLATIONYIELD.

Doing this we run into a problem – we don't seem to get the right amount of photons.

The following slides explain why that is

When setting the SCINTILLATIONYIELD of a material in geant4, geant takes into account the birks constant when calculating how many scintillation photons to produce in an event, even for electrons. This means you will see fewer scintillation photons than your input parameter.

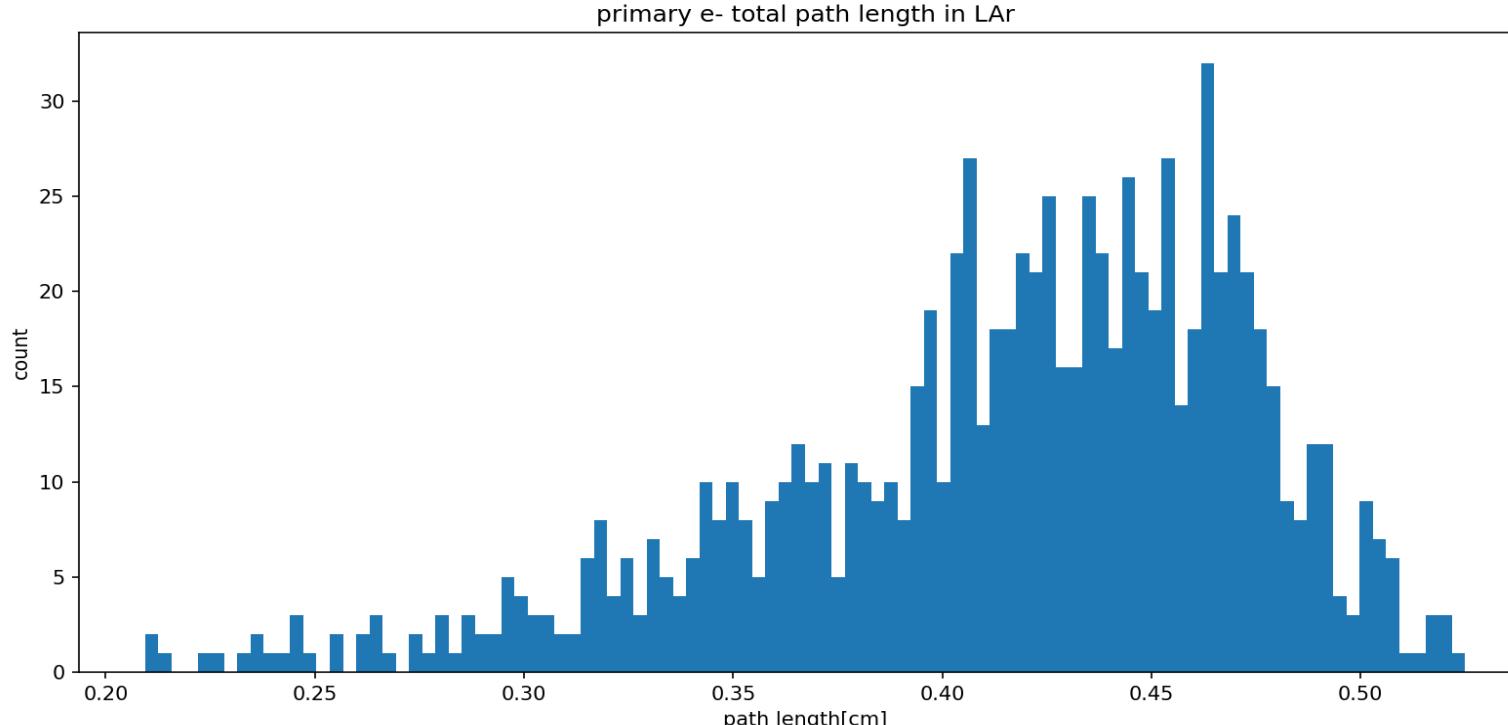
Even if you set birks constant = 0, the physics package overrides and sets the LAr birks constant = 0.032 mm/MeV

This tracks according to the data we gathered – using the average path length of the primary electron in the detector (with  $E_{k_i} = 1\text{ MeV}$ ), we can calculate the light yield attenuation assuming birks constant = 0.032 using:

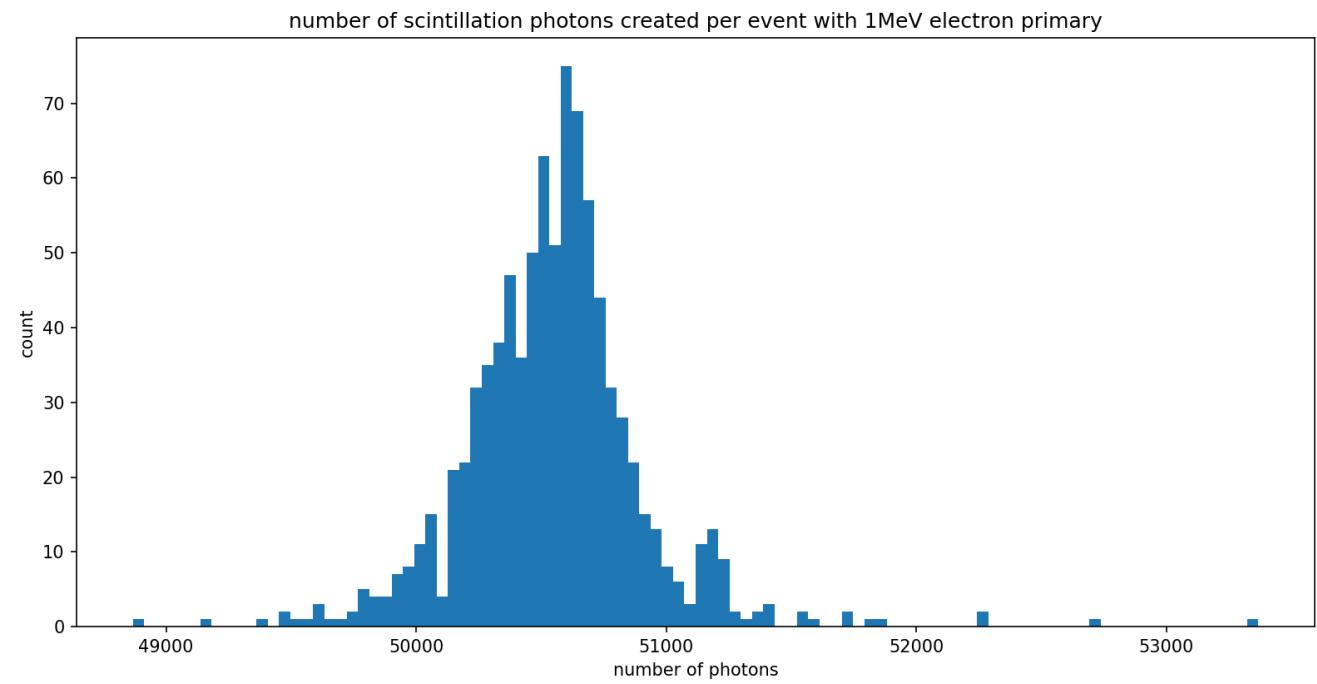
$$\text{attenuation} = \frac{1}{1 + 0.032 \cdot \text{avg. path length}} = 1.007 - 1.009$$

if you can the amount of scintillation photons per event, you can deduce the attention of the light yield by  $\frac{51000}{\text{num of scint photons}}$

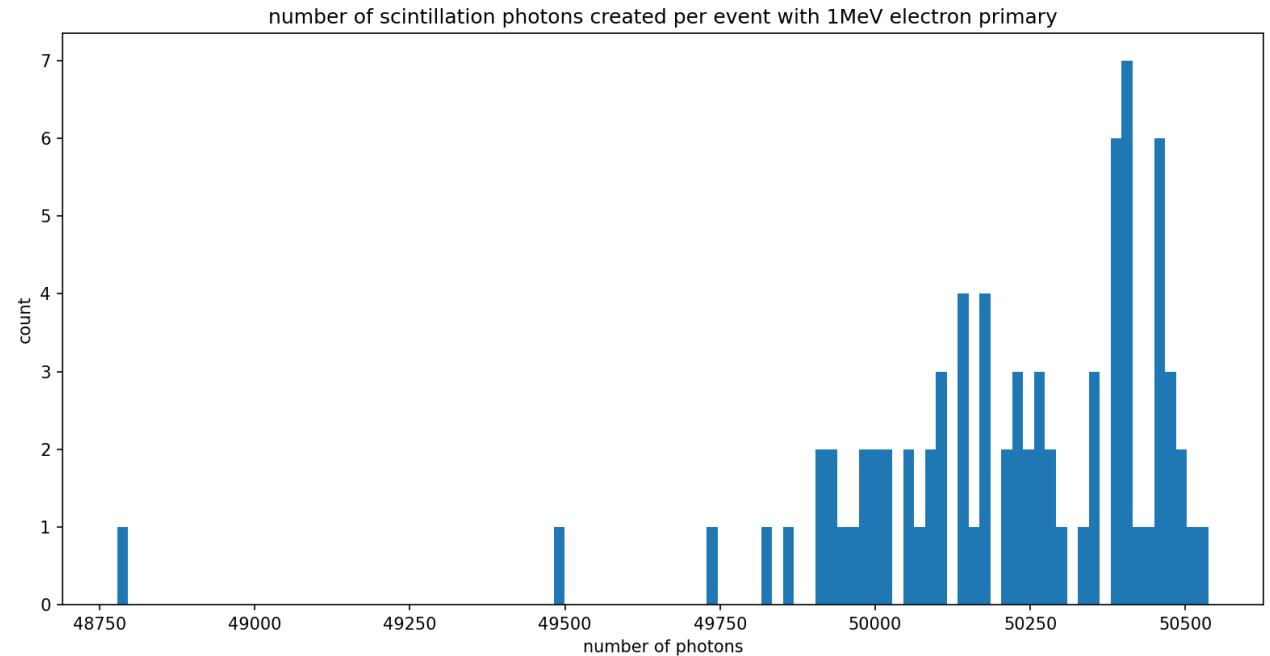
This comes out to be around 1.012 when checking each photon individually by its track and 1.008 using the process vector->GetNumPhotons method.



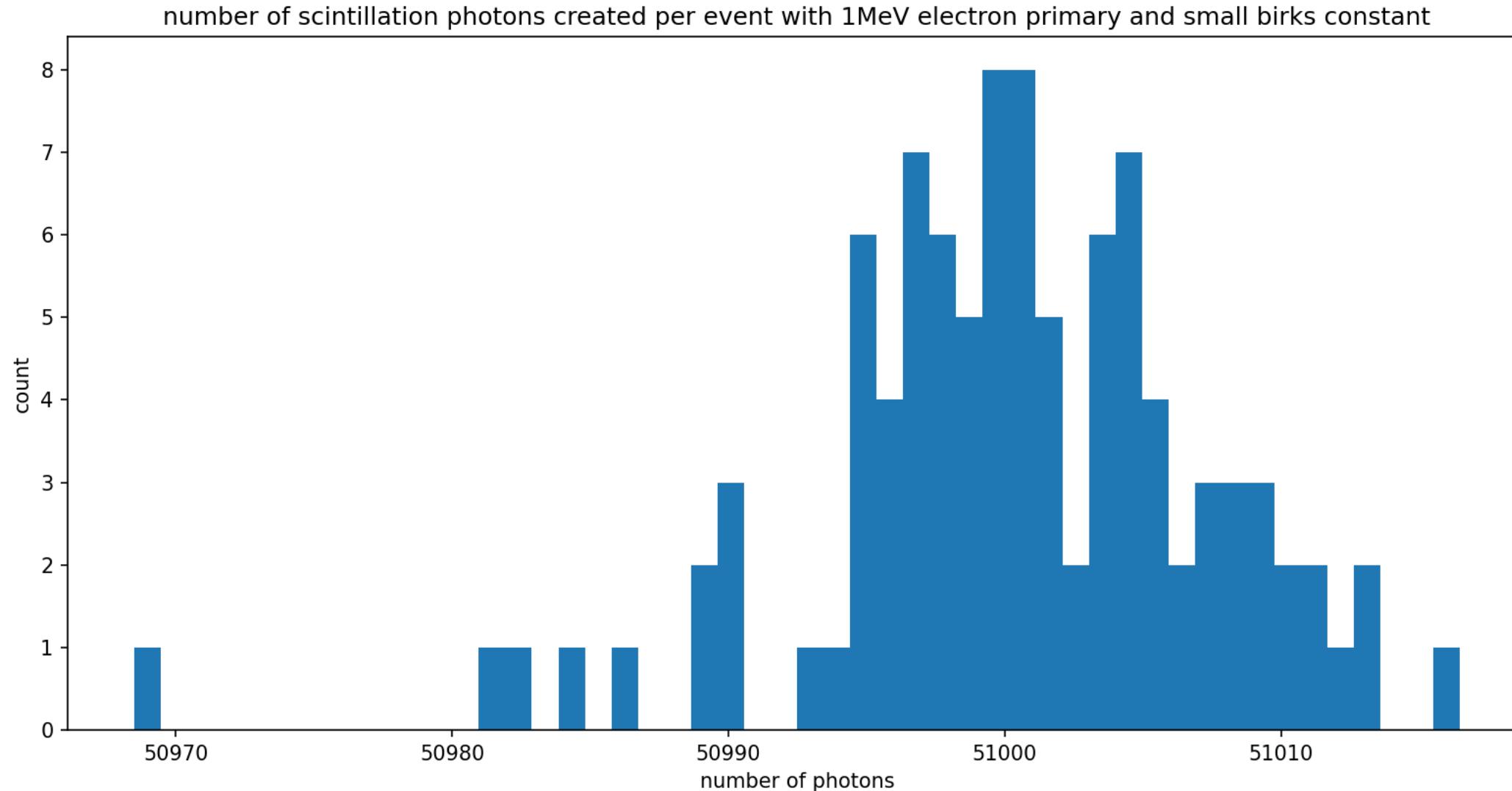
Using processVector->getNumPhotons



Checking each track individually

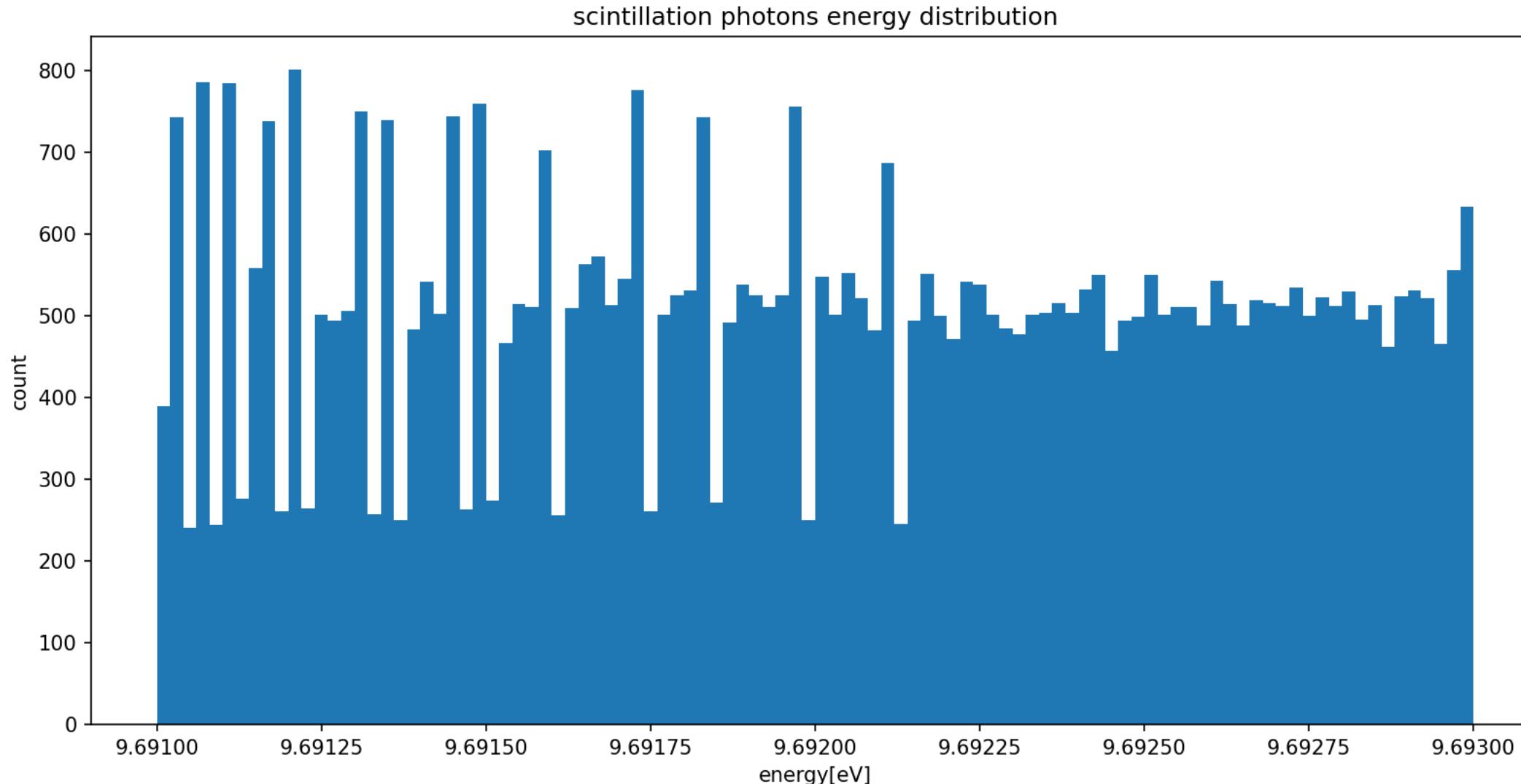


Finally, when setting birks constant to a really small number, but not 0, so that the physics package does not override it, you see a distribution centered around 51000 photons per event as expected:

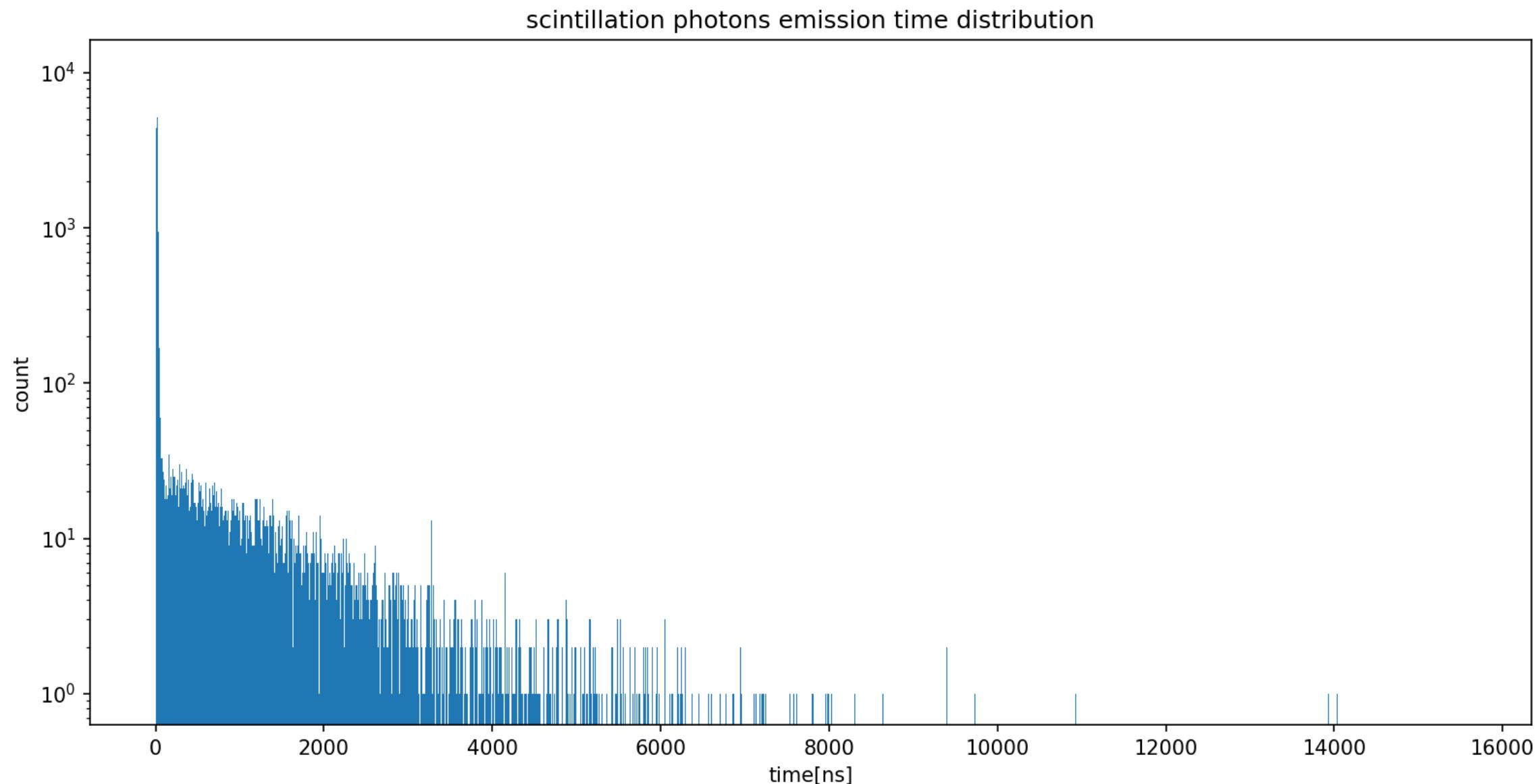


Moving on, we want to perform extra sanity checks. We start by firing a 1 MeV electron directly in the center of our LAr cylinder.

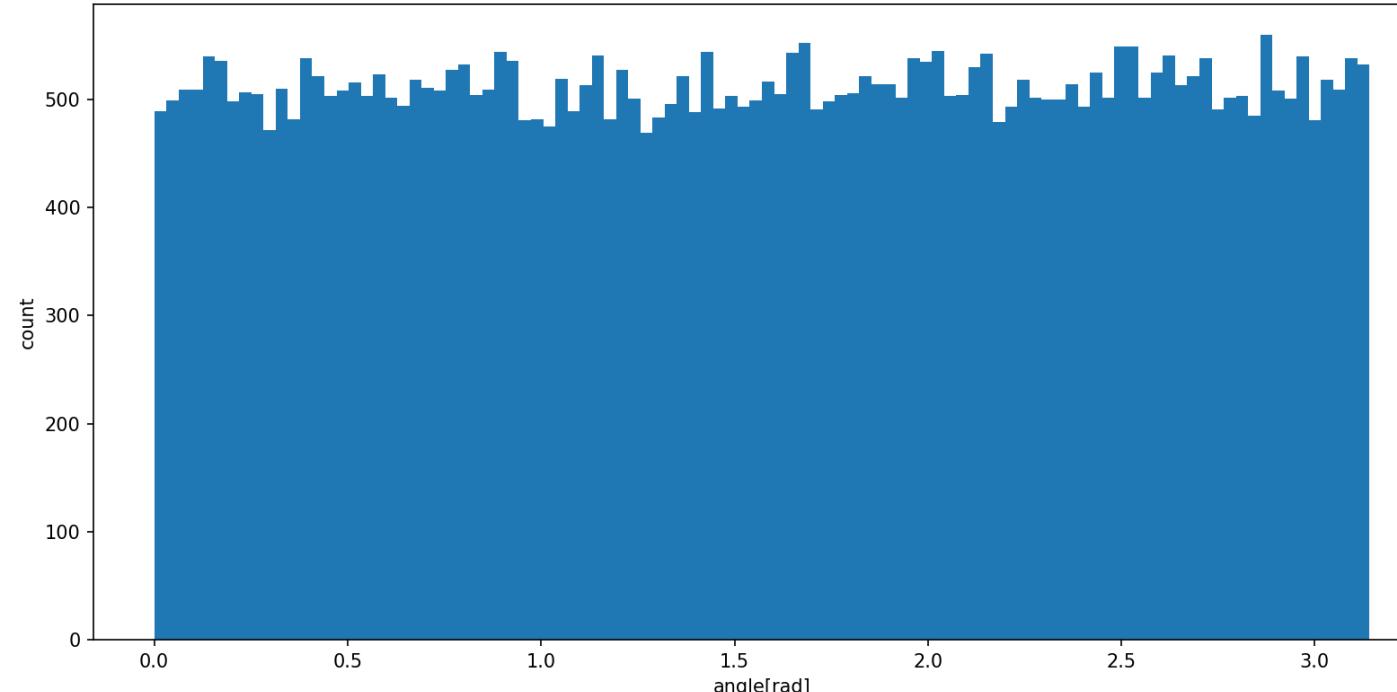
Firstly, the emitted scintillation photons' wavelength – this distribution matches the parameters we set – centered around 128nm wavelength for the emission spectra.



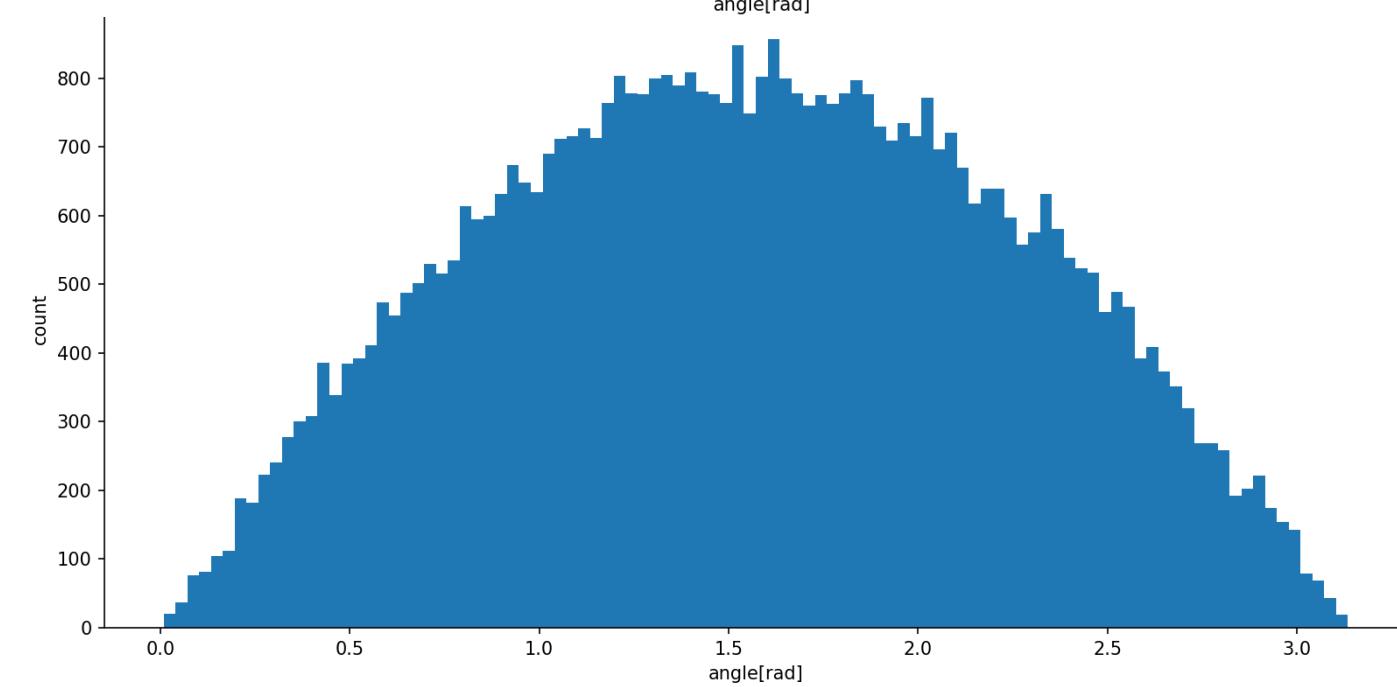
Next, looking at the emission timing of the scintillation photons, with fast component set to 7ns and slow component set to 1500ns:



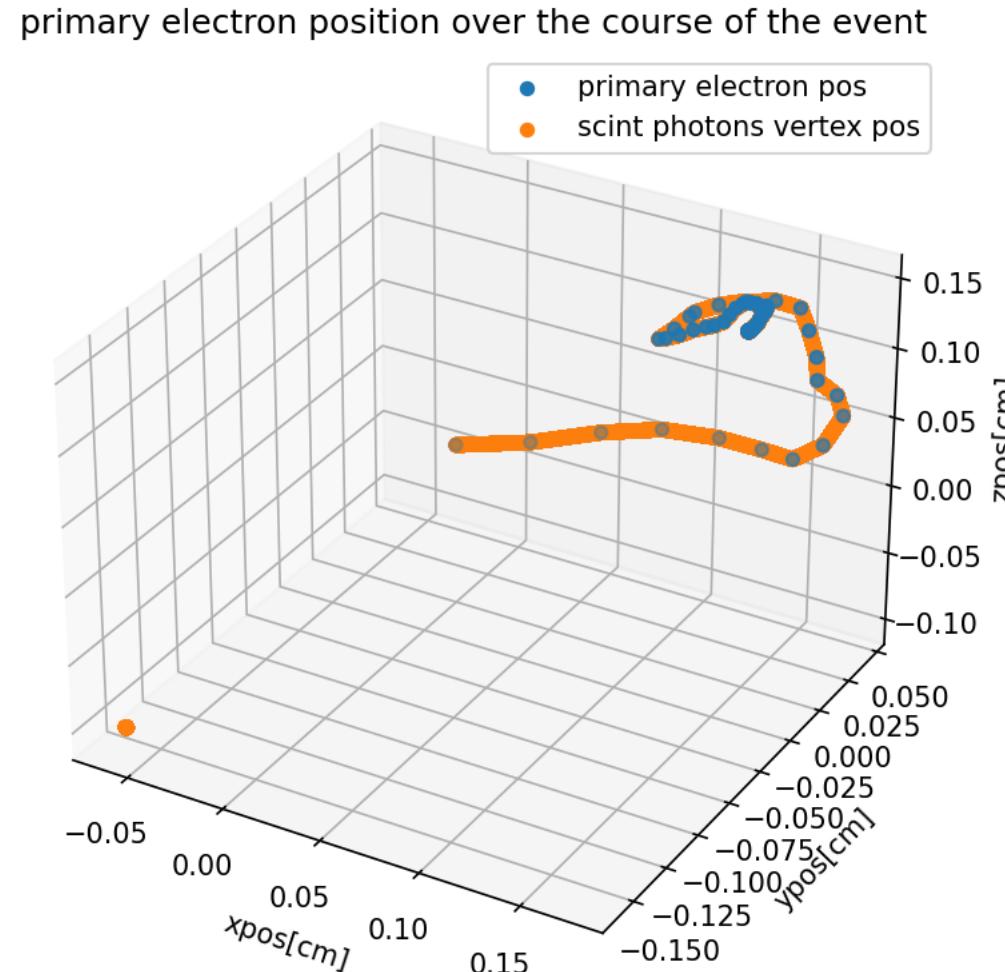
scintillation photons phi distribution



Next, the scintillation photons' angular distribution, separated to spherical angles  $\theta$  and  $\phi$ .



And finally, looking at the positions in which the scintillation photons were created and seeing they roughly match the track of the primary electron:

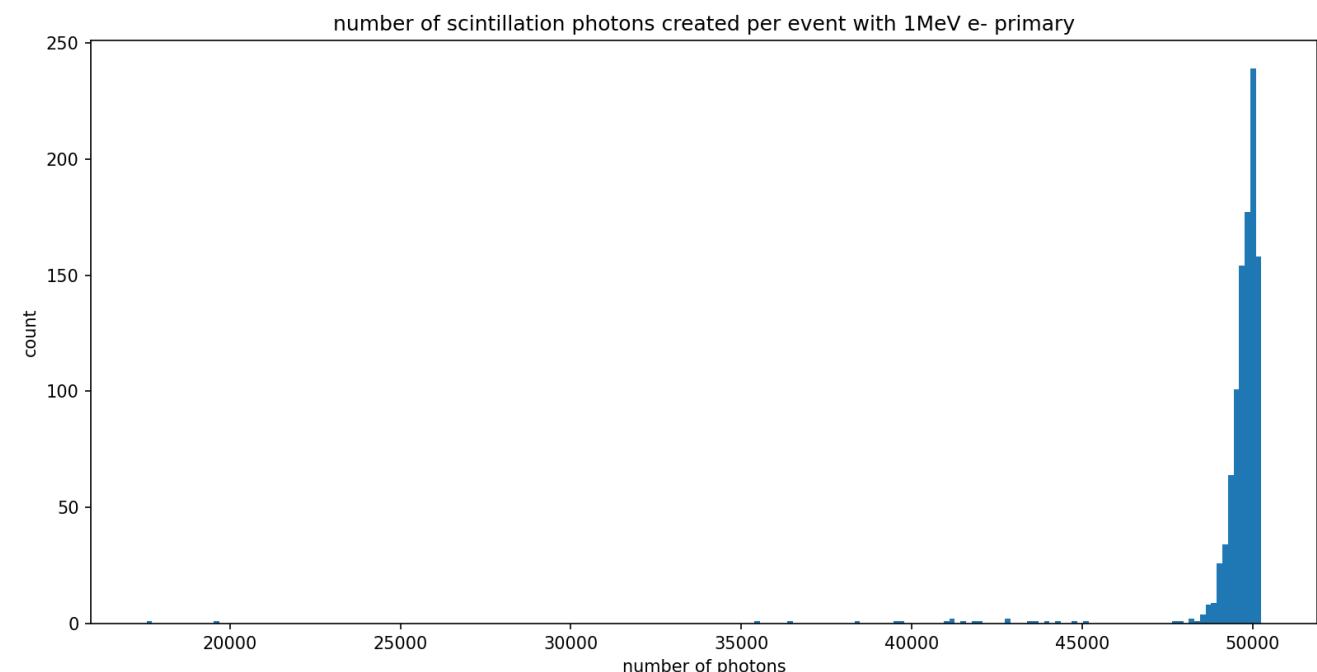
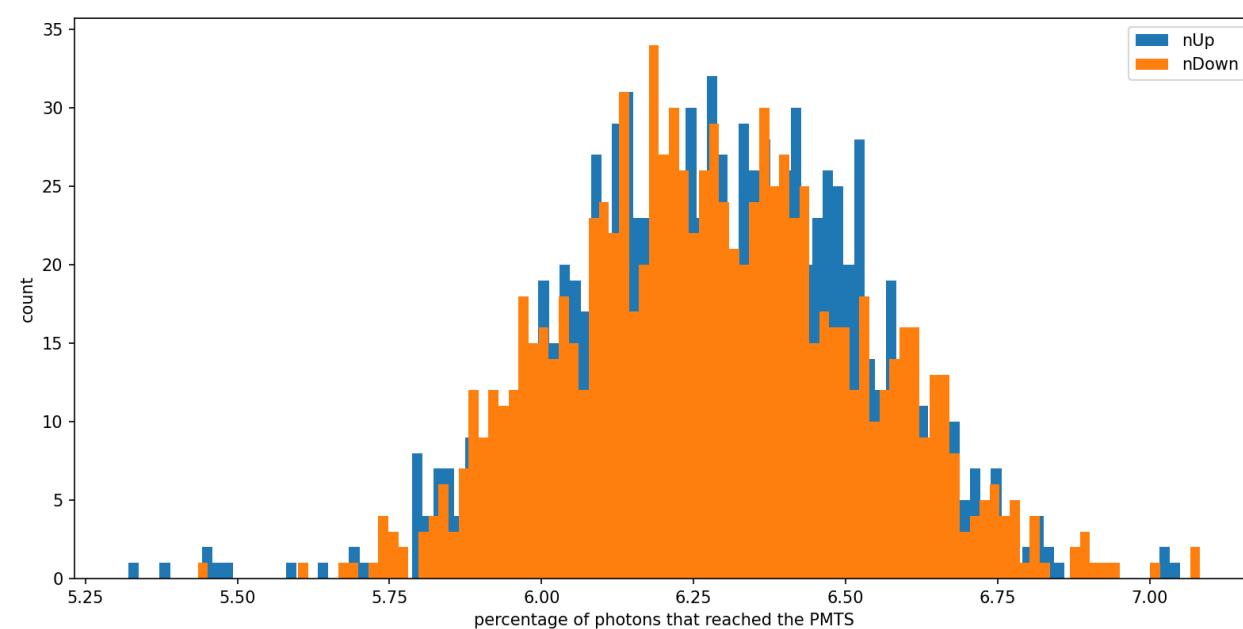


Now we can safely move on to try and fully recreate the findings of Creus et al.

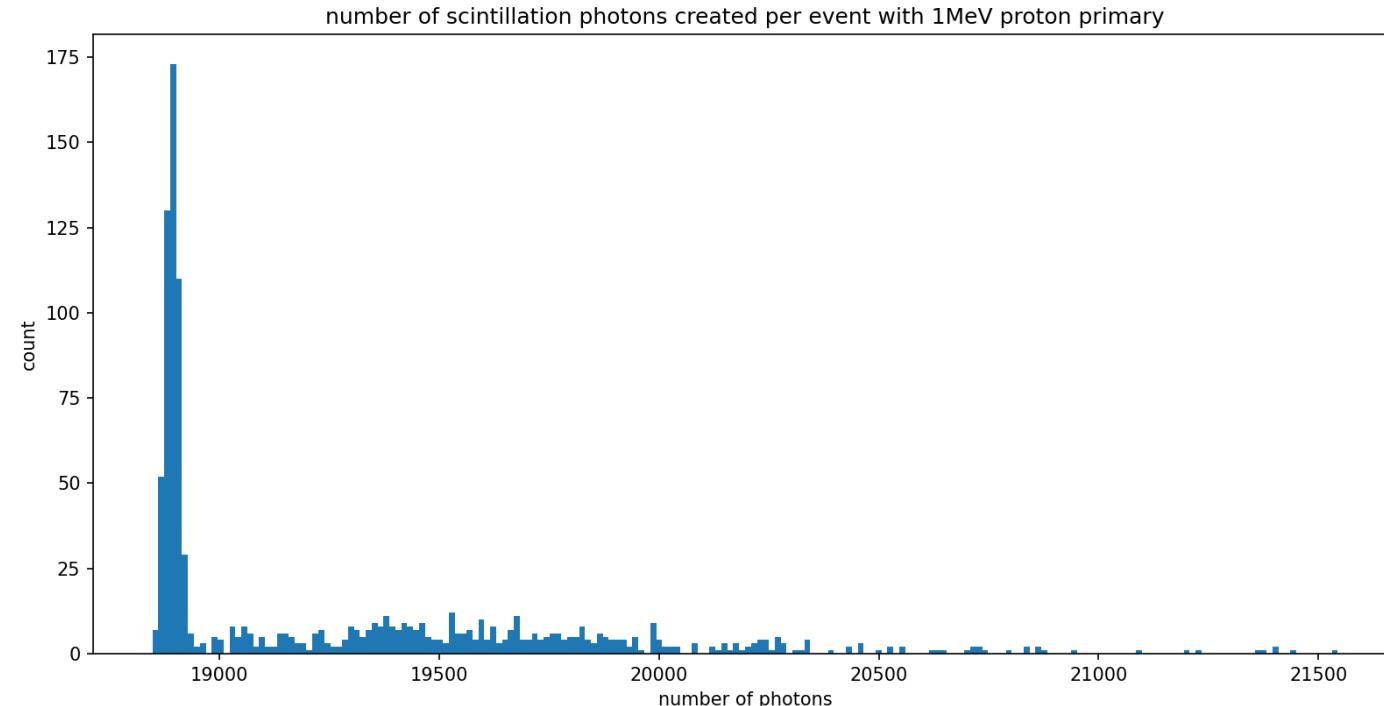
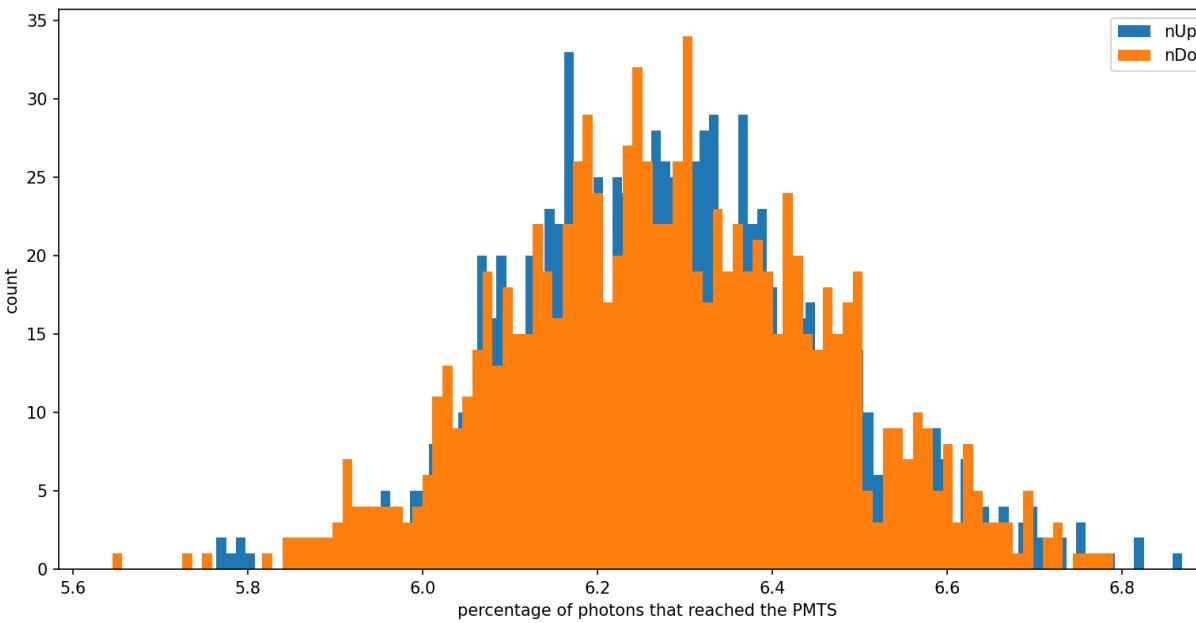
We first try to emulate PMTS and coincidence between them in the following way – we ask how many photons reached each base of the LAr cylinder (in their research there are 2 PMTS surrounding each base), and in order to take into account the QE of the TPB and PMT we roll a random number in a uniform distribution in the range (0,1) and if it is  $< 0.25$  we say the photon was absorbed inside the pmt.

In order to ask if there was coincidence between the PMTs, we look at the time difference of photons reaching the top PMT and the bottom PMT, and if in an event we get a time difference that is smaller than  $\frac{c}{n*h}$  (where c is the speed of light, n is the refractive index of LAr , and h is the height of the cylinder), we say that there is coincidence in the event.

The following plots are for a primary electron:



The following plots are for a primary proton:



The rest of the plots are pretty similar to the electron plots. The main difference is in the positions in which the scintillation photons are created – a proton primary deposits most of its energy over a very short distance relative to the electron, and so all scintillation photons are created very close to the center of the LAr cylinder (that is also the reason the scintillation yield is smaller for a proton relative to an electron)

