## MA 523: Homework 9

Carlos Salinas

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## Problem 9.1

(a) Show that for n = 3 the general solution to the wave equation  $u_{tt} - \Delta u = 0$  with spherical symmetry about the origin has the form

$$u = \frac{1}{r}F(r+t) + \frac{1}{r}G(r-t), \quad r = |x|,$$

with suitable F and G.

(b) Show that the solution with initial data of the form

$$u(r,0) = 0, \quad u_t(r,0) = h(r)$$

(h is an even function of r) is given by

$$u = \frac{1}{2r} \int_{r-t}^{r+t} \rho h(\rho) \, d\rho.$$

SOLUTION. For part (a): We show that

$$u = \frac{1}{r}F(r+t) + \frac{1}{r}G(r-t)$$

is in fact a solution to the wave equation in  $\mathbb{R}^3 \times (0, \infty)$ . By direct calculation, we have

$$u_{tt} = \frac{1}{r}F''(r+t) + \frac{1}{r}G''(r-t),$$

$$u_{r} = \frac{1}{r}F'(r+t) + \frac{1}{r}G'(r-t) - \frac{1}{r^{2}}F(r+t) - \frac{1}{r^{2}}G(r-t),$$

$$u_{rr} = \frac{1}{r}F''(r+t) + \frac{1}{r}G''(r-t) - \frac{2}{r^{2}}F'(r+t) - \frac{2}{r^{2}}G'(r-t) + \frac{2}{r^{3}}F(r+t) + \frac{2}{r^{3}}G(r-t),$$

$$\Delta_{S^{2}}u = 0.$$

and lastly,

$$\Delta u = \frac{\partial}{\partial r^2} u + \frac{2}{r} \frac{\partial}{\partial r} u + \frac{1}{r^2} \Delta_{S^2} u$$

$$= \frac{1}{r} F''(r+t) + \frac{1}{r} G''(r-t)$$

$$- \frac{2}{r^2} F'(r+t) - \frac{2}{r^2} G'(r-t) + \frac{2}{r^3} F(r+t) + \frac{2}{r^3} G(r-t)$$

$$+ \frac{2}{r^2} F'(r+t) + \frac{2}{r^2} G'(r-t) - \frac{2}{r^3} F(r+t) - \frac{2}{r^3} G(r-t)$$

$$+ 0$$

$$= \frac{1}{r} F''(r+t) + \frac{1}{r} G''(r-t).$$

Therefore, looking at the equation for  $u_{tt}$ , we see that

$$u_{tt} - \Delta u = 0;$$

i.e.,  $u(r,t):=\frac{1}{r}F(r+t)+\frac{1}{r}G(r-t)$  is a solution to the wave equation with F and G at least twice differentiable and such that they satisfy the initial conditions of the (nonhomogeneous) wave equation.

We still need to show that if u is a solution to the wave equation with spherical symmetry it has the form prescribed above. We trust this can be done for now and return to this problem as time permits.

For part (b): We proceed as in [E, §2.4.1(b)]. Set  $v := \frac{1}{r}u$ . Then the v satisfies the initial-value problem

## Problem 9.2

Show that the solution  $w(x_1,t)$  of the initial-value problem for the Klein-Gordon equation

$$\begin{cases} w_{tt} = w_{x_1 x_1} - \lambda^2 w, \\ w(x_1, 0) = 0, \quad w_t(x_1, 0) = h(x_1) \end{cases}$$
(9.1)

is given by

$$w(x_1,t) = \frac{1}{2} \int_{x_1-t}^{x_1+t} J_0(\lambda s) h(y_1) \, dy_1.$$

Here  $s^2 = t^2 - (x_1 - y_1)^2$ , while  $J_0$  denotes the Bessel function defined by

$$J_0(z) := \frac{2}{\pi} \int_0^{\frac{\pi}{2}} \cos(z \sin \theta) d\theta.$$

(Hint: Descend to (9.1) from the two-dimensional wave equation satisfied by

$$u(x_1, x_2, t) = \cos(\lambda x_2) w(x_1, t).$$

SOLUTION. Taking the hint, the equation  $u(x_1, x_2, t) := \cos(\lambda x_2)w(x_1, t)$  satisfies the two-dimensional wave equation as we will shortly see. First, let us compute find the necessary partial derivatives of u,

$$u_{tt} = \cos(\lambda x_2) w_{tt}(x_1, t),$$

$$u_{x_1} = \cos(\lambda x_2) w_{x_1}(x_1, t),$$

$$u_{x_1 x_1} = \cos(\lambda x_2) w_{x_1 x_1}(x_1, t),$$

$$u_{x_2} = -\lambda \sin(\lambda x_2) w(x_1, t),$$

$$u_{x_2 x_2} = -\lambda^2 \cos(\lambda x_2) w(x_1, t).$$

Then, by (9.1) together with the equations above we have

$$\begin{cases} u_{tt} - \Delta u = \cos(\lambda x_2)(w_{tt} - w_{x_1 x_1} + \lambda^2 w) = 0 & \text{in } \mathbb{R}^3 \times (0, \infty), \\ u = 0, \quad u = \cos(\lambda x_2)h(x_1) & \text{on } \mathbb{R}^3 \times \{t = 0\}. \end{cases}$$

By Kirchhoff's formula,

$$u(x_1, x_2, t) = \int_{\partial B(x_1, x_2, t)} t \cos(\lambda y_2) h(y_1) dS(y_1, y_2)$$

solves the above initial-value problem.

## Problem 9.3

Let u solve

$$\begin{cases} u_{tt} - \Delta u = 0 & \text{in } \mathbb{R}^3 \times (0, \infty), \\ u = g, & u_t = h & \text{on } \mathbb{R}^3 \times \{t = 0\} \end{cases}$$

where g and h are smooth and have compact support. Show there exists a constant C such that

$$|u(x,t)| \le Ct^{-1} \quad (x \in \mathbb{R}^3, t > 0).$$

Solution.