Fall 2016 Notes

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Chapter 1

Probability

We will devote this chapter to the material that is covered in MA 51900 (discrete probability) as it was covered in DasGupta's class. We will, for the most part, reference Feller's An introduction to probability theory and its applications, Volume 1 [3] (especially for the discrete noncalculus portion of the class) and DasGupta's own book Fundamentals of Probability: A First Course [1].

1.1 Counting

Some counting

Chapter 2

Introduction to Partial Differential Equations

Here we summarize some important points about PDEs. The material is mostly taken from Evans's Partial Differential Equations [2] with occasional detours to Strauss's Partial Differential Equations: An Introduction [8]. We will be following Dr. Petrosyan's Course Log which can be found here https://www.math.purdue.edu/~arshak/F16/MA523/courselog/, i.e., summarizing the appropriate chapters from [2].

2.1 First-Order PDEs

The transport equation

In this section, we consider the simplest first-order PDE, the *transport equation** with constant coefficients, i.e., the PDE

$$u_t + b \cdot Du = 0 \quad \text{in } \mathbb{R} \times (0, \infty),$$
 (2.1)

where b is a fixed vector in \mathbb{R}^n , and $u: \mathbb{R}^n \times [0, \infty) \to \mathbb{R}$ is the solution to the PDE. Our task is to find solutions u which satisfy the equation (2.1).

To address this task, let us suppose for a moment that we have a (smooth) solution u and try to compute it using the PDE (2.1). First, note that (2.1) asserts that the directional derivative $D_{(b,1)}u = 0$. Fix a point $(x,t) \in \mathbb{R}^n \times (0,\infty)$ and define

$$z(s) := u(x + sb, t + s)$$

for $s \in \mathbb{R}$. Then

$$\dot{z}(s) = Du(x + sb, t + s) \cdot b + u_t(x + sb, t + s) = 0.$$

Thus, z is a constant function of s and, consequently for each (x,t), u is constant on the line through (x,t) with direction $(b,1) \in \mathbb{R}^{n+1}$. Hence if we know the value of u at any point on each such line, we know its value everywhere in $\mathbb{R}^n \times (0,\infty)$.

^{*}For more details, refer to $[2, \S 2.1]$.

Initial-value problem

Let's now look at the transport equation with initial conditions

$$\begin{cases} u_t + b \cdot Du = 0 & \text{in } \mathbb{R}^n \times (0, \infty), \\ u = g & \text{on } \mathbb{R}^n \times \{t = 0\}. \end{cases}$$
 (2.2)

Here $b \in \mathbb{R}^n$ and $g : \mathbb{R}^n \to \mathbb{R}$ are known, and u is the unknown. Given (x,t), the line through (x,t) with direction (b,1) is represented parametrically by (x+sb,t+s) for $s \in \mathbb{R}$. This line hits the plane $\Gamma := \mathbb{R}^n \times \{t=0\}$ when s=-t, at the point (x-tb,0). Since u is constant on the line and u(x-tb,0)=g(x-tb), we deduce

$$u(x,t) = q(x-tb) (2.3)$$

for $x \in \mathbb{R}^n$, $t \ge 0$. So if (2.2) has a sufficiently regular solution u (at least C^1), it must certainly be given by (2.3).

2.2 Characteristics

We now turn our attention to a very important method for solving first-order PDEs, the method of *characteristics*. This section is paraphrased if copied not verbatim from [2, §3.2].

Derivation of characteristic ODEs

Consider the first-order (possibly non-linear) PDE

$$F(Du, u, x) = 0 \quad \text{in } U, \tag{2.4}$$

subject to the boundary condition

$$u = g \quad \text{on } \Gamma,$$
 (2.5)

where $\Gamma \subset \partial U$ and $g \colon \Gamma \to \mathbb{R}$ are known. We shall assume, for simplicity, that F and g are smooth. We now develop the *method of characteristics* to solve (2.4), (2.5) by converting the PDE into a system of ODEs. We proceed as follows: Suppose u solves (2.4), (2.5) and fix a point $x \in U$. We would like to calculate u(x) by finding some curve lying within U, connecting x with a point $x^0 \in \Gamma$ and along which we can compute u. Since (2.5) says u = g on Γ , we know the value of u at one end x^0 . We hope to be able to calculate u all along the curve, and so in particular at x.

Finding the characteristic curve

But how do we choose a path in U so all of this will work? Suppose the curve is described parametrically by the function $\mathbf{x}(x) = (x^1(s), \dots, x^n(s))$, the parameter s lying in some subinterval $I \subset \mathbb{R}$. Assuming u is a C^2 solution of (2.4), we define also

$$z(s) := u(\mathbf{x}(s)). \tag{2.6}$$

In addition, set

$$\mathbf{p}(s) := Du(\mathbf{x}(s)); \tag{2.7}$$

that is, $\mathbf{p}(s) = (p^{1}(s), \dots, p^{n}(s))$, where

$$p^{i}(s) = u_{xi}(\mathbf{x}(s)), \qquad 1 \le i \le n..$$
 (2.8)

So $z(\cdot)$ gives us the values of u along the curve and $\mathbf{p}(\cdot)$ records the values of gradient Du. We must choose a function $\mathbf{x}(\cdot)$ that will allow us to compute $z(\cdot)$ and $\mathbf{p}(\cdot)$.

Differentiating (2.8), we have

$$\dot{p}^{i}(s) = \sum_{j=1}^{n} u_{x_{i}x_{j}}(\mathbf{x}(s))\dot{x}^{j}(s). \tag{2.9}$$

But this expression is not too promising since it involves second order derivatives of u which we do not know (in fact, our solution need not be so regular as to have second order derivatives). On the other hand, if we differentiate (2.4) with respect to x_i , we have

$$\sum_{i=1}^{n} F_{p_j}(Du, u, x) u_{x_j x_i} + F_z(Du, u, x) u_{x_i} + F_{x_i}(Du, u, x) = 0.$$
(2.10)

We can use this identity to get rid of the second order derivatives in (2.12), provided we first set

$$\dot{x}^{j}(s) = F_{p_{j}}(\mathbf{p}(s), z(s), \mathbf{x}(s)), \qquad 1 \le j \le n.$$
 (2.11)

Assuming (2.11) holds, we evaluate (2.10) at $x = \mathbf{x}(s)$, thereby obtaining from (2.6) and (2.7) the identity

$$\sum_{j=1}^{n} F_{p_{j}}(\mathbf{p}(s), z(s), \mathbf{x}(s)) u_{x_{i}x_{j}}(\mathbf{x}(s)) + F_{z}(\mathbf{p}(s), z(s), \mathbf{x}(s)) p^{i}(s) + F_{x_{i}}(\mathbf{p}(s), z(s), \mathbf{x}(s)) = 0$$

Finally, we differentiate (2.6) to give us

$$\dot{z}(s) = \sum_{j=1}^{n} u_{x_j}(\mathbf{x}(s))\dot{x}^j(s)$$

$$= \sum_{j=1}^{n} p^j(s) F_{p_j}(\mathbf{p}(s), z(s), \mathbf{x}(s)),$$
(2.12)

the second equality holding by (2.8) and (2.9).

We summarize our results by rewriting equations (2.11), (2.11), and (2.12) as

$$\begin{cases}
(a) \ \dot{\mathbf{p}}(s) = -D_x F(\mathbf{p}(s), z(s), \mathbf{x}(s)) - D_z F(\mathbf{p}(s), z(s), \mathbf{x}(s)) \mathbf{p}(s), \\
(b) \ \dot{z}(s) = D_p F(\mathbf{p}(s), z(s), \mathbf{x}(s)) \cdot \mathbf{p}(s), \\
(c) \ \dot{\mathbf{x}}(s) = D_p F(\mathbf{p}(s), z(s), \mathbf{x}(s)).
\end{cases} (2.13)$$

Furthermore,

$$F(\mathbf{p}(s), z(s), \mathbf{x}(s)) = 0. \tag{2.14}$$

These identities hold for $s \in I$.

The system (2.13) of 2n+1 first order ODEs comprises the *characteristic equotions/ODEs* of the nonlinear first-order PDE (2.4). The functions $\mathbf{p}(\cdot)$, $\mathbf{x}(\cdot)$ are called *characteristics* and $\mathbf{x}(\cdot)$ is called the *projected characteristic* (it is the projection of the full characteristics $(\mathbf{p}, z, \mathbf{x}) \subset \mathbb{R}^{2n+1}$ onto the physical region $U \subset \mathbb{R}^n$).

Theorem 2.1 (Structure of characteristic ODEs). Let $u \in C^2(U)$ solve the nonlinear, first-order partial differential equation (2.4) in U. Assume $\mathbf{x}(\cdot)$ solves the ODE (2.13)(c), where $\mathbf{p}(\cdot) = Du(\mathbf{x}(\cdot))$, $z(\cdot) = u(\mathbf{x}(\cdot))$. Then $\mathbf{p}(\cdot)$ solves the ODE (2.13)(a) and $z(\cdot)$ solves the ODE (2.13)(b), for those s such that $\mathbf{x}(s) \in U$.

We still need to discover appropriate initial conditions for (2.13) to be useful. We do that in the following section.

Examples

But before we move on, we look at some examples to show you how to use (2.13) to find solutions to (2.4).

The linear case

Suppose (2.4) is linear, i.e., has the form

$$F(D, u, x) = \mathbf{x} \cdot Du(x) + c(x)u(x) = 0, \qquad x \in U.$$
(2.15)

Then, rewriting (2.15) in terms of p, z, and x, we have $F(p, z, x) = \mathbf{b}(x) \cdot p + c(x)z$, so

$$D_p F = \mathbf{b}(x)$$

so (2.13)(c) becomes

$$\dot{\mathbf{x}}(s) = \mathbf{x}(\mathbf{x}(s)),\tag{2.16}$$

an ODE involving only the function $\mathbf{x}(\cdot)$. Furthermore (2.13)(b) becomes

$$\dot{z}(s) = \mathbf{b}(\mathbf{x}(s)) \cdot \mathbf{p}(s). \tag{2.17}$$

Then equation (2.14) simplifies (2.17), yielding

$$\dot{z}(s) = -c(\mathbf{x}(s))z(s).$$

This ODE is linear in $z(\cdot)$, noce we know the function $\mathbf{x}(\cdot)$ by solving (2.16). In summary, we have

$$\begin{cases} (\mathbf{a}) \ \dot{\mathbf{x}}(s) = \mathbf{b}(\mathbf{x}(s)), \\ (\mathbf{b}) \ \dot{z}(s) = -c(\mathbf{x}(s))z(s). \end{cases}$$
 (2.18)

Example 2.2. Let's now look at a simple example to see how to use (2.24) to solve a PDE. Consider the PDE

$$\begin{cases} x_1 u_{x_2} - x_2 u_{x_1} = u & \text{in } U, \\ u = g & \text{on } \Gamma, \end{cases}$$
 (*)

where $U = \{x_1 > 0, x_2 > 0\}$ and $\Gamma = \{x_1 > 0, x_2 = 0\} \subset \partial U$. The PDE (*) is of the form (2.15) with $\mathbf{b} = (-x_2, x_1)$ and c = -1. Thus the equations (2.24) read

$$\begin{cases} \dot{x}^1 = -x^2, & \dot{x}^2 = x^1, \\ \dot{z} = z. \end{cases}$$
 (**)

Solving this system of ODEs we have

$$\begin{cases} x^{1}(s) = x^{0} \cos s, & x^{2}(s) = x^{0} \sin x, \\ z(s) = z^{0} e^{s} \\ = g(x^{0}) e^{s}, \end{cases}$$

where $x^0 \ge 0$, $0 \le s \le \pi/2$. Now, fix $(x_1, x_2) \in U$. Select s > 0, $x^0 > 0$ so that $(x_1, x_2) = (x^1(s), x^2(s)) = (x^0 \cos s, x^0 \sin s)$ and solve for x^0 , in this case, $x^0 = \sqrt{x_1^2 + x_2^2}$, $s = \arctan(x_2/x_1)$, and therefore

$$u(x) = u(x^{1}(s), x^{2}(s))$$

$$= z(s)$$

$$= g(x^{0})e^{s}$$

$$= g(\sqrt{x_{1}^{2} + x_{2}^{2}})e^{\arctan(x_{2}/x_{1})}.$$

The quasilinear case

Let's look at the quasilinear case now, i.e., (2.4) with the form

$$F(Du, u, x) = \mathbf{b}(x, u(x)) \cdot Du(x) + c(x, u(x)) = 0.$$
(2.19)

In this circumstance $F(p, z, x) = \mathbf{b}(x, z) \cdot p + c(x, z)$, whence

$$D_p F = \mathbf{b}(x, z).$$

Hence equation (2.13)(c) reads

$$\dot{\mathbf{x}}(s) = \mathbf{b}(\mathbf{x}(s), z(s)),\tag{2.20}$$

an ODE involving only the function \mathbf{x} . Furthermore, (2.13)(c) becomes

$$\dot{z}(s) = \mathbf{b}(\mathbf{x}(s)) \cdot \mathbf{p}(s), \tag{2.21}$$

which, after applying (2.14), turns into

$$\dot{z}(s) = -c(\mathbf{x}(s))z(s). \tag{2.22}$$

In summary, we have

$$\begin{cases} (\mathbf{a}) \ \dot{\mathbf{x}}(s) = \mathbf{b}(\mathbf{x}(s)), \\ (\mathbf{b}) \ \dot{z}(s) = -c(\mathbf{x}(s))z(s). \end{cases}$$
 (2.23)

We will see later that the equation for $\mathbf{p}(\cdot)$ is in fact not needed (at least in the linear and quasilinear cases).

Example 2.3. Let's look at an example of a quasilinear PDE. Consider the PDE

$$\begin{cases} u_{x_2} + u_{x_1} = u & \text{in } U, \\ u = g & \text{on } \Gamma. \end{cases}$$
 (*)

Here $U = \{x_2 > 0\}$ and $\Gamma = \{x_2 = 0\} = \partial U$ with $\mathbf{b} = (1,1)$ and $c = -z^2$. Thus, the equations (2.23) yield

$$\begin{cases} \dot{x}^1 = 1, & \dot{x}^2 = 1, \\ \dot{z} = z^2. \end{cases}$$

Consequently

$$\begin{cases} x^{1}(s) = x^{0} + s, & x^{2}(s) = s, \\ z(s) = \frac{z^{0}}{1 - sz^{0}} = \frac{g(x^{0})}{1 - sg(x^{0})}, & \end{cases}$$

where $x^0 \in \mathbb{R}$, $s \ge 0$, provided the denominator is not zero.

Now, fix a point $(x_1, x_2) \in U$ and select s > 0 and $x^0 \in \mathbb{R}$ so $(x_1, x_2) = (x^1(s), x^2(s)) = (x^0 + s, s)$, i.e., $x^0 = x_1 - x_2$, $s = x_2$. Then

$$u(x) = u(x^{1}(s), x^{2}(s))$$

$$= z(s)$$

$$= \frac{g(x^{0})}{1 - sg(x^{0})}$$

$$= \frac{g(x_{1} - x_{2})}{1 - x_{2}g(x_{1} - x_{2})}.$$

This solution of course only makes sense if $1 - x_2 g(x_1 - x_2) \neq 0$.

The fully nonlinear case

In the general case, we must integrate the full characteristics (2.13) if possible. In this case, we cannot generally reduce (2.13) and we must look at the PDE on a case-by-case basis.

Example 2.4. Let's look at an example. Consider the fully nonlinear PDE

$$\begin{cases} u_{x_1} u_{x_2} = u & \text{in } U, \\ u = x_2^2 & \text{on } \Gamma, \end{cases}$$
 (*)

where $U = \{x_2 > 0\}, \Gamma = \partial U = \{x_1 = 0\}.$ Here, $F(p, z, x) = p_1 p_2 - z$, and hence (2.13) yield

$$\begin{cases} \dot{p}^1 = p^1, & \dot{p}^2 = p^2, \\ \dot{z} = 2p^1p^2, & \\ \dot{x}^1 = p^2, & \dot{x}^2 = p^1. \end{cases}$$

After integrating these equations, we have

$$\begin{cases} x^{1}(s) = p_{2}^{0}(e^{s} - 1), & x^{2}(s) = x^{0} + p_{1}^{0}(e^{s} - 1), \\ z(s) = z^{0} + p_{1}^{0}p_{2}^{0}(e^{2s} - 1), \\ p^{1}(s) = p_{1}^{0}e^{s}, & p^{2}(s) = p_{2}^{0}e^{s}, \end{cases}$$

where $x^0 \in \mathbb{R}$, $s \in \mathbb{R}$, and $z^0 = (x^0)^2$.

We must determine $p^0=(p_1^0,p_2^0)$. Since $u=x_2^2$ on Γ , $p_2^0=u_{x_2}(0,x^0)=2x^0$. Furthermore the PDE $u_{x_1x_2}=u$ itself implies $p_1^0p_2^0=z^0=(x^0)^2$, and so $p_1^0=x^0/2$. Consequently the formulas above become

$$\begin{cases} x^{1}(s) = 2x^{0}(e^{s} - 1), & x^{2}(s) = \frac{x^{0}}{2}(e^{s} + 1), \\ z(s) = (x^{0})^{2}e^{2s}, & p^{1}(s) = \frac{x^{0}}{2}e^{s}, & p^{2}(s) = 2x^{0}e^{s}. \end{cases}$$

Fix a point $(x_1, x_2) \in U$. Select s and x^0 so that $(x_1, x_2) = (x^1(s), x^2(s)) = (2x^0(e^s - 1), \frac{x^0}{2}(e^s + 1))$. This equality implies

$$x^0 = \frac{4x_2 - x_1}{4};$$
 $e^s = \frac{x_1 + 4x_2}{4x_2 - x_1};$

so

$$u(x) = u(x^{1}(s), x^{2}(s))$$

$$= z(s)$$

$$= (x^{0})^{2} e^{2s}$$

$$= \frac{(x_{1} + 4x_{2})^{2}}{16}.$$

Compatibility conditions on boundary data

Let $x^0 \in \Gamma$. We intend to use the characteristic ODEs (2.13) to construct a solution u to (2.4), (2.5), at least near x^0 , and for this, we must determine appropriate initial conditions

$$\mathbf{p}(0) = p^0, \qquad z(0) = z^0, \qquad \mathbf{x}(0) = x^0.$$
 (2.24)

We will assume throughout that Γ is flat (i.e., isometric to $\{x_n = 0\}$) at least near x^0 .[†] Clearly if the curve $\mathbf{x}(\cdot)$ passes through x^0 , we should insist that

$$z^0 = g(x^0). (2.25)$$

What should we require concerning $\mathbf{p}(0) = p^0$? Since (2.5) implies

$$u(x_1,\ldots,x_{n-1},0)=g(x_1,\ldots,x_{n-1})$$

near x^0 , we may differentiate this to find

$$u_{x_i}(x^0) = g_{x_i}(x^0).$$

This, along with the PDE (2.4), forces p^0 to satisfy

$$\begin{cases} p_i^0 = g_{x_i}(x^0) & 1 \le i \le n - 1, \\ F(p^0, z^0, x^0) = 0. \end{cases}$$
 (2.26)

 $^{^\}dagger \mathrm{This}$ can always be achieved, assuming some regularity of $\Gamma.$

These identities provide n equations for the n quantities $p^0 = (p_1^0, \dots, p_n^0)$. We call (2.25) and (2.26) compatibility conditions. A triple $p^0, z^0, x^0 \in \mathbb{R}^{2n+1}$ satisfying (2.25), (2.26), is called *admissible*. Note z^0 is uniquely determined by the boundary condition and our choice of x^0 , but p^0 satisfying (2.26) may not exist or be unique.

Having ascertained what are appropriate boundary conditions for the characteristic ODEs with $\mathbf{x}(\cdot)$ intersecting Γ at x^0 , we proceed to construct a solution to (2.4), (2.5), near x^0 . We now ask, can we somehow appropriately perturb (p^0, z^0, x^0) , keeping the compatibility conditions?

In other words, given a point $y = (y_1, \dots, y_{n-1}, 0) \in \Gamma$, with y close enough to x^0 , we intend to solve the ODEs (2.13) with initial conditions

$$\mathbf{p}(0) = \mathbf{q}(y), \qquad z(0) = g(y), \qquad \mathbf{x}(0) = y.$$
 (2.27)

Our task is now to find a function $\mathbf{q}(\cdot) = (q^1(\cdot), \dots, q^n(\cdot))$, so that

$$\mathbf{q}(x^0) = p^0 \tag{2.28}$$

and $\mathbf{q}(y), g(y), y$ is admissible; i.e., so

$$\begin{cases} q^{i}(y) = g_{x_{i}}(y) & 1 \le i \le n - 1, \\ F(\mathbf{q}(y), g(y), y) = 0, \end{cases}$$
 (2.29)

hold for all $y \in \Gamma$ close to x^0 .

Lemma 2.5 (Noncharacteristic boundary conditions). There exists a unique solution $\mathbf{q}(\cdot)$ of (2.28), (2.29), for all $y \in \Gamma$ sufficiently close to x^0 , provided

$$F_{p_n}(p^0, z^0, x^0) \neq 0.$$

See $[2, \S 3.2 \text{ c}]$ for proof.

More generally,

$$D_p F(p^0, z^0, x^0) \cdot \boldsymbol{\nu}(x^0) \neq 0,$$

 $\nu(x^0)$ denoting the outward unit normal to ∂U at x^0 .

Now, given a point $y = (y_1, \dots, y_{n-1}, 0)$ sufficiently close to x^0 we solve the characteristic ODEs (2.28) subject to (2.29).

Write

$$\begin{cases} \mathbf{p}(s) = \mathbf{p}(y, s) = \mathbf{p}(y_1, \dots, y_{n-1}, s), \\ z(s) = z(y, s) = z(y_1, \dots, y_{n-1}, s), \\ \mathbf{x}(s) = \mathbf{x}(y, s) = \mathbf{x}(y_1, \dots, y_{n-1}, s), \end{cases}$$

displaying the dependence of the solution of (2.28), (2.29), with respect to s and y.

Lemma 2.6 (Local invertibility). Assume we have the noncharacteristic condition $F_{p_n}(p^0, z^0, x^0) \neq$ 0. Then there exists an open interval $I \subset \mathbb{R}$ containing 0, a neighborhood W of x^0 in $\Gamma \subset \mathbb{R}^{n-1}$ and a neighborhood V of x^0 in \mathbb{R}^n , such that for each $x \in V$ there exists a unique $s \in I$, $y \in W$ such that

$$x = \mathbf{x}(y, s).$$

The mappings $x \mapsto s, y$ are C^2 .

In view of this lemma, for each $x \in V$, we can uniquely solve (at least locally) the equation

$$\begin{cases} x = \mathbf{x}(y, s), \\ y = \mathbf{y}(x), \quad s = s(x). \end{cases}$$
 (2.30)

Define

$$\begin{cases} u(x) := z(\mathbf{y}(x), s(x)), \\ \mathbf{p}(x) = \mathbf{p}(\mathbf{y}(x), s(x)), \end{cases}$$
 (2.31)

for $x \in V$, s and y as in (2.30).

Putting all of this together, we have the following result:

Theorem 2.7 (Local existence theorem). The function u defined above is C^2 and solves the PDE

$$F(Du(x), u(x), x) = 0$$
 $x \in V$,

with the boundary condition

$$u(x) = g(x)$$
 $x \in \Gamma \cap V$.

See [2, §3.2.4] for more details.

2.3 Power Series

We now turn our attention to an important class of solutions to PDEs, analytic solutions. We begin this section with a brief discussion on noncharacteristic surfaces.

Noncharacteristic surfaces

Consider the k^{th} -order quasilinear PDE

$$\sum_{|\alpha|=k} a_{\alpha}(D^{k-1}u, \dots, u, x)D^{\alpha}u + a_0(D^{k-1}u, \dots, u, x) = 0,$$
(2.32)

in some region in $U \subset \mathbb{R}^n$. Let us assume for simplicity that Γ is a smooth (n-1)-dimensional hypersurface in U. For any $x^0 \in \Gamma$ let $\boldsymbol{\nu}(x^0) = \boldsymbol{\nu} = (\nu_1, \dots, \nu_n)$ denote the unit normal to Γ at x^0 . Furthermore, we define the j^{th} normal derivative of u at $x^0 \in \Gamma$ by

$$\frac{\partial^{j} u}{\partial \nu^{j}} := \sum_{|\alpha|=j} \binom{j}{\alpha} D^{\alpha} u \boldsymbol{\nu}^{\alpha} = \sum_{\alpha_{1} + \dots + \alpha_{n} = j} \binom{j}{\alpha} \frac{\partial^{j} u}{\partial^{\alpha_{1}} x_{1} \cdots \partial^{\alpha_{n}} x_{n}} \nu_{1}^{\alpha_{1}} \cdots \nu_{n}^{\alpha_{n}}.$$

Let $g_0, \ldots, g_{k-1} \colon \Gamma \to \mathbb{R}$ be k given functions. The *Cauchy problem* is to find a function u solving the PDE (2.32), subject to the boundary conditions

$$\begin{cases}
 u = g_0, \\
 \frac{\partial u}{\partial \nu} = g_1, \\
 \vdots \\
 \frac{\partial^{k-1} u}{\partial \nu^{k-1}} = g_{k-1}
\end{cases}$$
(2.33)

We say that (2.33) prescribe the Cauchy data on Γ .

Evans demonstrates how to extend the notion of noncharacteristic curve/surface to the PDE (2.32), (2.33), first in the case that Γ is flat and then makes some remarks about the case that Γ is nonflat. The most important result from this section is the notion of noncharacteristic surface for (2.32), (2.33) which is summarized in the following definition:

Definition 2.8. We say the surface Γ is noncharacteristic for the PDE (2.32) provided

$$\sum_{|\alpha|=k} a_{\alpha} \boldsymbol{\nu}^{\alpha} \neq 0. \tag{2.34}$$

for all values of the arguments of the coefficients $a_{\alpha}(|\alpha|=k)$.

Theorem 2.9 (Cauchy data and noncharacteristic surfaces). Assume that Γ is noncharacteristic for the PDE (2.32). Then if u is a smooth solution of (2.32) satisfying the Cauchy data (2.33), we can uniquely compute all of the partial derivatives of u along Γ in terms of Γ , the functions g_0, \ldots, g_{k-1} , and the coefficients $a_{\alpha}(|\alpha| = k)$, a_0 .

See $[2, \S4.6.b]$ for a proof.

Real analytic functions

We now briefly review the concept of analytic functions.

Definition 2.10. A function $f: \mathbb{R}^n \to \mathbb{R}$ is said to be *(real) analytic near* x_0 if there exists r > 0 and constants $\{f_{\alpha}\}$ such that

$$f(x) = \sum_{\alpha} f_{\alpha}(x - x_0)^{\alpha} \qquad |x - x_0| < r,$$

the sum taken over all multiindices α .

Remarks 2.11.

- (i) Remember that when we write x^{α} we mean $x^{\alpha} = x_1^{\alpha_1} \cdots x_n^{\alpha_n}$.
- (ii) If f is analytic near x_0 , then f is C^{∞} near x_0 . Furthermore the coefficients f_{α} are computed as $f_{\alpha} = D^{\alpha} f(x_0)/\alpha!$, where $\alpha! = \alpha_1! \cdots \alpha_n!$. Thus f equals its Taylor formula about x_0

$$f(x) = \sum_{\alpha} \frac{1}{\alpha!} D^{\alpha} f(x_0) (x - x_0)^{\alpha} \qquad |x - x_0| < r.$$

To simplify, we hereafter take x_0 to be 0.

Example 2.12. Let us look at an example of an analytic function. For r > 0, set

$$f(x) := \frac{r}{r - (x_1 + \dots + x_n)}$$
 for $|x| < r/\sqrt{n}$.

Then

$$f(x) = \frac{1}{1 - \frac{x_1 + \dots + x_n}{r}}$$

$$= \sum_{k=0}^{\infty} \left(\frac{x_1 + \dots + x_n}{r}\right)^k$$

$$= \sum_{k=0}^{\infty} \frac{1}{r^k} \sum_{|\alpha| = k} {|\alpha| \choose \alpha} x^{\alpha}$$

$$= \sum_{\alpha} \frac{|\alpha|!}{r^{|\alpha|} \alpha!} x^{\alpha}.$$

We employed the multinomial theorem for the third equality above and recalled that $\binom{|\alpha|}{\alpha} = |\alpha|!/\alpha!$. This power series is absolutely convergent for $|x| < r/\sqrt{n}$. Indeed,

$$\sum_{\alpha} \frac{|\alpha|!}{r^{|\alpha|} \alpha!} |x^{\alpha}| = \sum_{k=0}^{\infty} \left(\frac{|x_1| + \dots + |x_n|}{r} \right)^k < \infty,$$

since $|x_1| + \cdots + |x_n| \le |x|\sqrt{n} < r$.

The analytic equation we have just considered is quite important as it allows us to *majorize*, and therefore confirm the convergence of, other power series.

Definition 2.13. Let

$$f = \sum_{\alpha} f_{\alpha} x^{\alpha}, \qquad g = \sum_{\alpha} g_{\alpha} x^{\alpha}$$

be two power series. We say that g majorizes f, written $g \gg f$, if

$$g_{\alpha} \geq |f_{\alpha}|$$
 for all α .

Here is an important result:

Lemma 2.14 (Majorants).

- (i) If $g \gg f$ and g converges for |x| < r, then f also converges for |x| < r.
- (ii) If $f = \sum_{\alpha} f_{\alpha} x^{\alpha}$ converges for |x| < r and $0 < s\sqrt{n} < r$, then f has a majorant for $|x| < s/\sqrt{n}$.

Cauchy-Kovalevskaya theorem

We now turn our attention to a very important theorem regarding analytic solutions to (2.32) with analytic Cauchy data (2.33) specified on an analytic, noncharacteristic hypersurface Γ .

Reduction to a first-order system

We intend to construct a solution u as a power series, but we must first transform (2.32), (2.33) into a more convenient form.

First of all, upon flattening out the boundary by an analytic mapping, we can reduce to the situation that $\Gamma \subset \{x_n = 0\}$. Additionally, by subtracting off appropriate analytic functions, we may assume that the Cauchy data are identically zero. Consequently, we may assume, without loss of generality, that our problem reads:

enerality, that our problem reads:
$$\begin{cases} \sum_{|\alpha|=k} a_{\alpha}(D^{k-1}u,\ldots,u,x)D^{\alpha}u + a_{0}(D^{k-1}u,\ldots,u,x) = 0 & \text{for } |x| < r, \\ u = 0, \\ \frac{\partial u}{\partial x_{n}} = 0, & \text{for } |x'| < r, x_{n} = 0, \\ \vdots \\ \frac{\partial^{k-1}u}{\partial x_{n}^{k-1}} = 0, \end{cases}$$
 (2.35)

r > 0 to be found. Here $a_{\alpha}(|\alpha| = k)$ and a_0 are analytic are analytic, and as usual we write $x' = (x_1, \dots, x_{n-1})$.

Finally, we this to a first-order system. To do so, we set

$$\mathbf{u} := \left(u, \frac{\partial}{\partial x_1} u, \dots, \frac{\partial^{k-1}}{\partial x_n^{k-1}} u\right),$$

the components of which are all the partial derivatives of u of order less than k. Let m hereafter denote the number of components of \mathbf{u} , so $\mathbf{u} \colon \mathbb{R}^n \to \mathbb{R}^m$, $\mathbf{u} = (u^1, \dots, u^m)$. Observe from the boundary condition in (2.35) that $\mathbf{u} = 0$ for |x'| < r, $x_n = 0$.

For $1 \le k \le m-1$, we can compute $u_{x_n}^k$ in terms o $\{\mathbf{u}_{x_j}\}_{j=1}^{n-1}$. Furthermore in view of the noncharacteristic condition $a_{(0,\dots,0,k)} \ne 0$ near 0, we can utilize the PDE in (2.35) to solve for $u_{x_n}^m$ in terms of \mathbf{u} and $\{\mathbf{u}_{x_j}\}_{j=1}^{n-1}$.

Employing these relations, we can consequently transform (2.35) into a boundary-value problem for a first-order system \mathbf{u} , the coefficients of which are analytic functions. This system is of the form

$$\begin{cases}
\mathbf{u}_{x_n} = \sum_{j=1}^{n-1} \mathbf{B}_j(\mathbf{u}, x') \mathbf{u}_{x_j} + \mathbf{c}(\mathbf{u}, x') & \text{for } |x| < r, \\
\mathbf{u} = 0 & \text{for } |x'| < r, x_n = 0,
\end{cases}$$
(2.36)

where we are given the analytic function $\mathbf{B}_j \colon \mathbb{R}^m \times \mathbb{R}^{n-1} \to \operatorname{Mat}(m,\mathbb{R}), \ 1 \leq j \leq n-1$, and $\mathbf{c} \colon \mathbb{R}^m \times \mathbb{R}^{n-1} \to \mathbb{R}^m$. We will write $\mathbf{B}_j = (b_j^{k\ell})$ and $\mathbf{c} = (c^1, \dots, c^m)$. Carefully note that we have assumed $\{\mathbf{B}_j\}_{j=1}^{n-1}$ and \mathbf{c} do not depend on $x_n > We$ can always reduce to this situation by introducing if necessary a new component u^{m+1} of the unknown \mathbf{u} , with $u^{m+1} \equiv x_n$.

In particular, the components of the system of partial differential equations in (2.36) read

$$u_{x_n}^k = \sum_{j=1}^{n-1} \sum_{\ell=1}^m b_j^{k\ell}(\mathbf{u}, x') u_{x_j}^\ell + c^k(\mathbf{u}, x'), \qquad 1 \le k \le m.$$
 (2.37)

Having reduced to the special form (2.36), we can now expand \mathbf{u} into a power series near 0.

Theorem 2.15 (Cauchy–Kovalevskaya theorem). Assume $\{B_j\}_{j=1}^{n-1}$ and \mathbf{c} are real analytic functions. Then there exists r > 0 and a real analytic function

$$\mathbf{u} = \sum_{\alpha} \mathbf{u}_{\alpha} x^{\alpha} \tag{2.38}$$

solving the boundary-value problem (2.36).

2.4 Second-order PDEs

We now turn our attention to second-order partial differential equations. For this section, we will mainly refer to [5, §3] starting with a minor detour to [2, §7.2.5].

Equations in two variables

In this section, we consider second-order hyperbolic PDEs involving only two variables. The rough idea is that since a function in two variables has only three second order partial derivatives, algebraic and analytic simplifications in the structure of the PDE may be possible, which are unavailable for more than two variables.

We begin by considering a general linear second-order PDE in two variables

$$\sum_{i,j=1}^{2} a^{ij} u_{x_i x_j} + \sum_{i=1}^{2} b^i u_{x_i} + cu = 0,$$
(2.39)

where the coefficients a^{ij} , b^i , c, $1 \le i, j \le 2$, with $a^{ij} = a^{ji}$, and the unknown u are functions of the two variables x_1 and x_2 in some region $U \subset \mathbb{R}^2$. Note that for the moment, and in contrast to the theory developed above, we do *not* identify ether x_1 or x_2 wit the variable t denoting time.

We now pose the following question:

Is it possible to simplify the structure of the PDE (2.39) by introducing new independent variables?

In other words, can we expect to turn the PDE into a nicer form by rewriting in terms of new variables $y = \Phi(x)$?

More precisely, set

$$y_1 = \Phi^1(x_1, x_2)$$

$$y_2 = \Phi^2(x_1, x_2)$$
(2.40)

2.5 Second-order linear PDEs

That last section was a bit useless, now we'll get down to developing some useful theory.

Characteristics for linear and quasilinear second-order equations

We start with the general quasilinear second-order equation for a function with u(x,y)

$$au_{xx} + 2bu_{xy} + cu_{yy} - d = 0, (2.41)$$

where a, b, c, and d depend on x, y, u, u_x, u_y . Here the Cauchy problem consists of finding a solution u of (2.41) with given (compatible) values of u, u_x, u_y on a curve Γ in the xy-plane. Thus, for Γ given parametrically by

$$x = f(s), \quad y = g(s),$$
 (2.42)

we prescribe on Γ

$$u = h(s), u_x = \varphi(s), u_y = \psi(s). \tag{2.43}$$

The values of any function v(x, y) and of its first derivative $v_x(v, y, v_y(x, y))$ along the curve Γ are connected by the compatibility condition

$$\dot{v} = v_x f'(s) + v_y g'(s)$$

which follows by differentiating v(f(s), g(s)) with respect to s.

Supóngase que $f: X \to \mathbb{R}$ es una aplicación biyectiva

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