MA 523: Homework, Midterms and Practice Problems Solutions

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1 Homework Solutions

These are my (corrected) solutions to Petrosyan's Math 523 homework for the fall semester of 2016.

1.1 Homework 1

PROBLEM 1.1.1 (Taylor's formula). Let $f: \mathbf{R}^n \to \mathbf{R}$ be smooth, $n \geq 2$. Prove that

$$f(x) = \sum_{|\alpha| \le k} \frac{1}{\alpha!} D^{\alpha} f(0) x^{\alpha} + O(|x|^{k+1})$$

as $x \to \mathbf{0}$ for each $k = 1, 2, \ldots$, assuming that you know this formula for n = 1.

Hint: Fix $x \in \mathbb{R}^n$ and consider the function of one variable g(t) := f(tx). Prove that

$$\frac{d^k}{dt^k}g(t) = \sum_{|\alpha|=k} \frac{k!}{\alpha!} D^{\alpha} f(tx) x^{\alpha},$$

by induction on m.

SOLUTION. Taking the hint, let us consider the function in one variable g(t) := f(tx) for $x \in \mathbb{R}^n$ fixed. We show by induction on k that

$$\frac{d^k}{dt^k}g(t) = \sum_{|\alpha|=k} \frac{k!}{\alpha!} D^{\alpha} f(tx) x^{\alpha}.$$
(1.1.1)

Once we have shown (1.1.1) holds, evaluating g at t=1 gives us the desired equality; i.e,

$$f(x) = g(1)$$

which, by Taylor's formula in one variable, is

$$= \sum_{j=0}^{k} \frac{g^{(j)}(0)}{j!} 1^{j} + O(|x|^{k+1})$$

applying (1.1.1) here gives us

$$= \sum_{j=0}^{k} \frac{1}{j!} \left[\sum_{|\alpha|=j} \frac{j!}{\alpha!} D^{\alpha} f(tx) x^{\alpha} \right] + O(|x|^{k+1})$$

$$= \sum_{j=0}^{k} \left[\sum_{|\alpha|=j} \frac{1}{\alpha!} D^{\alpha} f(0) x^{\alpha} \right] + O(|x|^{k+1})$$

$$= \sum_{|\alpha| \le k} \frac{1}{\alpha!} D^{\alpha} f(0) x^{\alpha} + O(|x|^{k+1})$$

as desired.

Let us now show that (1.1.1) holds. To prove this we consider the algebra on the differential operator d/dt. By the chain rule, we have

$$\frac{d}{dt}(\,\cdot\,) = \sum_{j=1}^{n} x_j \frac{\partial}{\partial x_j}(\,\cdot\,).$$

Since f is smooth by Schwartz's theorem the differential operators $\partial/\partial x_j$ and $\partial/\partial x_l$ commute for all $1 \leq j, l \leq n$. Therefore, by the multinomial theorem,

$$\frac{d^k}{dt^k}(\,\cdot\,) = \left(\sum_{j=1}^n x_j \frac{\partial}{\partial x_j}(\,\cdot\,)\right)^k = \sum_{|\alpha|=k} \frac{m!}{\alpha!} x^\alpha D^\alpha(\,\cdot\,).$$

PROBLEM 1.1.2. Write down the characteristic equation for the PDE

$$u_t + b \cdot Du = f \tag{*}$$

on $\mathbf{R}^n \times (0, \infty)$, where $b \in \mathbf{R}^n$. Using the characteristic equation, solve (*) subject to the initial condition u = g on $\mathbf{R}^n \times \{t = 0\}$. Make sure the answer agrees with formula (5) in §2.1.2 of [E].

SOLUTION. Write

$$F(p, z, x, t) := (b, 1) \cdot p - f = 0.$$

Then the characteristic equations to the problem (*) with the initial value $u(\cdot,0)=g(\cdot)$ are given by

$$\begin{cases} \dot{p} = -D_{x,t}F - D_{z}Fp = 0, \\ \dot{z} = D_{p}F \cdot p = (b,1) \cdot p = f, \\ (\dot{x},\dot{t}) = D_{p}F = (b,1). \end{cases}$$

Now let us solve the characteristic equations above subject to the initial values $(x(0), t(0)) = (x^0, 0) \in \mathbf{R}^n \times (0, \infty)$; these are

$$\begin{cases} x(s) = x^{0} + bs, & t(s) = s, \\ z(s) = z(0) + \int_{0}^{s} f(x(\tau), t(\tau)) d\tau \\ = g(x^{0}) + \int_{0}^{s} f(x^{0} + b\tau, \tau) d\tau. \end{cases}$$

Solving back, we have t = s, $x^0 = x - bs = x - bt$, and therefore

$$u(x,t) = z(s) = g(x - bt, t) + \int_0^t f(x - b\tau, \tau) d\tau$$

solves the transport equation (*) with initial value $u(\cdot,0)=g(\cdot)$. This verifies formula [5] from [E, §2.1.2].

Problem 1.1.3. Solve using the characteristics:

- (a) $x_1^2 u_{x_1} + x_2^2 u_{x_2} = u^2$, u = 1 on the line $x_2 = 2x_1$.
- (b) $uu_{x_1} + u_{x_2} = 1$, $u(x_1, x_1) = x_1/2$.
- (c) $x_1u_{x_1} + 2x_2u_{x_2} + u_{x_3} = 3u$, $u(x_1, x_2, 0) = g(x_1, x_2)$.

SOLUTION. For part (a): Write $F := (x_1^2, x_2^2) \cdot p - z^2 = 0$. We have the characteristic equations

$$\begin{cases} \dot{p} = -D_x F - D_z F p = 2((x_1 - z)p_1, (x_2 - z)p_2), \\ \dot{z} = D_p \cdot p = z^2, \\ \dot{x} = (x_1^2, x_2^2). \end{cases}$$

We can solve the characteristic equations with respect to the initial conditions $x(0) = (x^0, 2x^0)$, z(0) = 1 on the line $x_1 = 2x_2$; these are

$$\begin{cases} x_1(s) = \frac{x^0}{1+x^0s}, & x_2(s) = \frac{2x^0}{1+2x^0s}, \\ z(s) = \frac{1}{1-s}. & \end{cases}$$

Now we solve these in terms of the coordinates (x_1, x_2) . Assuming $x^0 \neq 0$, we have

$$s = \frac{1}{x^0} - \frac{1}{x_1}$$
 and $s = \frac{1}{2x^0} - \frac{1}{x_2}$.

Therefore,

$$s = 2\left(\frac{1}{2x^0} - \frac{1}{x_2}\right) - \left(\frac{1}{x^0} - \frac{1}{x_1}\right)$$
$$= \frac{1}{x_1} - \frac{2}{x_2}.$$

Thus,

$$u(x_1, x_2) = \frac{1}{1 - \left(\frac{1}{x_1} - \frac{2}{x_2}\right)} = \frac{x_1 x_2}{x_1 x_2 - x_2 - 2x_1}$$

solves the PDE F for (x_1, x_2) on the line $x_1 = 2x_2$ away from the origin.

For part (b): Write $F = (z, 1) \cdot p - 1 = 0$. Then we have the characteristic equations

$$\begin{cases} \dot{p} = -D_x F - D_z p = -(p_1, 0) \\ \dot{z} = D_p \cdot p = 1 \\ \dot{x} = D_p F = (z, 1) \end{cases}$$

Next we solve the characteristic equations subject to the initial conditions $x(0) = (x^0, x^0)$, $z(0) = x^0/2$ on the line $x_1 = x_2$; these are

$$\begin{cases} z(s) = \frac{1}{2}x^0 + s, \\ x_1(s) = x^0 + \frac{1}{2}(x^0s + s^2), & x_2(s) = x^0 + s. \end{cases}$$

Then, solving in terms of the coordinates (x_1, x_2) , we have

$$x^0 = 2(x_2 - z)$$
 and $s = 2z - x_2$.

Therefore,

$$x_1 = 2(x_2 - z) + (x_2 - z)(2z - x_2) + \frac{1}{2}(2z - x_2)^2$$

= $-\frac{1}{2}x_2(x_2 - 4) + (x_2 - 2)z$.

Hence,

$$u(x_1, x_2) = \frac{2x_1 + x_2^2 - 4x_2}{2(x_2 - 2)}$$

solves the PDE F subject to the condition $u(x_1, x_1) = x_1/2$ provided $x_2 \neq 2$.

For part (c): Write $F := (x_1, 2x_2, 1) \cdot p - 3z = 0$. Then the characteristic equations are

$$\begin{cases} \dot{p} = -D_x F - D_z p = (2p_1, p_2, 3p_3) \\ \dot{z} = D_p F \cdot p = 3z \\ \dot{x} = D_p F = (x_1, 2x_2, 1) \end{cases}$$

Next we sole the characteristic equations subject to the initial conditions $x(0) = (x_1^0, x_2^0, 0), z(s) = g(x_1^0, x_2^0);$ these are

$$\begin{cases} x_1(s) = x_1^0 e^s, & x_2(s) = x_2^0 e^{2s}, & x_3(s) = s, \\ z(s) = g(x_1^0, x_2^0) e^{3s}. \end{cases}$$

Then, solving for u in terms of the coordinates (x_1, x_2, x_3) , we have

$$s = x_3$$
, $x_1^0 = x_1 e^{-s}$, and $x_2^0 = x_2 e^{-2s}$.

Thus,

$$u(x_1, x_2, x_3) = g(x_1 e^{-x_3}, x_2 e^{-2x_3})e^{3x_3}$$

solves the PDE F subject to the condition $u(x_1, x_2, 0) = g(x_1, x_2)$.

PROBLEM 1.1.4. For the equation

$$u = x_1 u_{x_1} + x_2 u_{x_2} + \frac{1}{2} (u_{x_1}^2 + u_{x_2}^2)$$

find a solution with $u(x_1, 0) = (1 - x_1^2)/2$.

SOLUTION. The equation is nonlinear and therefore, we do not expect the method of characteristics to yield a unique solution to the PDE

$$F := x_1 p_1 + x_2 p_2 + \frac{1}{2} (p_1^2 + p_2^2) - z.$$

Let us find the characteristic equations for F; these are

$$\begin{cases} \dot{p} = -D_x F - D_z F p = -(p_1, p_2) - (-1)(p_1, p_2) = 0, \\ \dot{z} = D_p F \cdot p = (x_1 + p_1, x_2 + p_2) \cdot (p_1, p_2) = (x_1 + p_1) p_1 + (x_2 + p_2) p_2, \\ \dot{x} = D_p F = (x_1 + p_1, x_2 + p_2), \end{cases}$$

Next we solve the characteristic equations subject to the initial values $x(0) = (x^0, 0), z(0) = \frac{1}{2}(1 - (x^0)^2)$ and, after revisiting the equation F, $p_1(0) = -x^0$ and

$$p_2(0)^2 = 2(-(x^0)^2 + \frac{1}{2}(x^0)^2 + \frac{1}{2}(1 - (x^0)^2)) = 1$$

so $p_2(0) = \pm 1$. Therefore, the solution to the characteristic equations subject to these initial values is

$$\begin{cases} p_1(s) = -x^0, & p_2(s) = \pm 1, \\ x_1(s) = x^0, & x_2(s) = \pm 1(e^s - 1), \\ z(s) = \frac{1}{2}(1 - (x^0)^2) + (e^s - 1). \end{cases}$$

Thus, solving for s and x^0 in terms of the coordinates (x_1, x_2) , we have

$$u(x_1, x_2) = \frac{1}{2}(1 - x_1^2) \pm x_2.$$

1.2 Homework 2

PROBLEM 1.2.1. Verify assertion (36) in [E, §3.2.3], that when Γ is not flat near x^0 the noncharacteristic condition is

$$D_p F(p^0, z^0, x^0) \cdot \nu(x^0) \neq 0.$$

(Here $\nu(x^0)$ denotes the normal to the hypersurface Γ at x^0).

SOLUTION. Throughout this, let (p^0, z^0, x^0) denote an admissible triple to the PDE F at some point x^0 in its domain. First, note that the condition

$$D_p F(p^0, z^0, x^0) \cdot \nu(x^0) \neq 0$$

reduces to the standard noncharacteristic boundary condition if Γ is flat near x^0 since in that case the normal to the hypersurface at x^0 will be $(0, \ldots, 0, 1)$; i.e.,

$$0 \neq D_p F(p^0, z^0, x^0) \cdot (0, \dots, 0, 1)$$

= $F_{p_n}(p^0, z^0, x^0)$.

We shall therefore proceed to flatten the hypersurface Γ near x^0 and apply the standard noncharacteristic boundary condition.

Assuming Γ is reasonably tame, by the implicit function theorem, it can be written as the graph $\{x_n = \varphi(x_1, \dots, x_{n-1})\}$ on some neighborhood U of x^0 for φ smooth. Now consider the smooth mapping $\Phi \colon U \to V$ given by

$$\begin{cases} y_j := \Phi^j(x) := x_j, & 1 \le j \le n - 1, \\ y_n := \Phi^n(x) := x_n - \varphi(x_1, \dots, x_{n-1}), \end{cases}$$

where we use y to denote new coordinates on the image of Φ . Note that $\nu(x^0)$ is parallel to the gradient $D_x \Phi^n = (-\varphi_{x_1}, \ldots, -\varphi_{x_{n-1}}, 1)$ so the inner product of the latter with $F_{p_n}(p^0, z^0, x^0)$ is nonzero if and only if the inner product of $\nu(x^0)$ with $F_{p_n}(p^0, z^0, x^0)$ is nonzero.

Set $\Delta := \Phi(\Gamma)$ and define $v(y) := u(\Phi^{-1}(y))$. Then $u(x) = v(\Phi(x))$. Moreover, by the chain rule we have

$$D_{x_i}u = \sum_{j=1}^n D_{y_j}vD_{x_j}\Phi^j, \quad 1 \le j \le n;$$

i.e., $D_x u = D_y v D_x \Phi$. Thus, v satisfies the PDE

$$G(D_yv,v,y):=F(D_yvD_x\varPhi,v,\varPhi^{-1}(y))=0$$

in Δ and, since Δ has been flattened near $y^0 := \Phi(x^0)$, applying the noncharacteristic condition, we have

$$D_{p_n}G = (D_{p_1}F)(D_{x_1}\Phi^n) + \dots + (D_{p_n}F)(D_{x_n}\Phi)$$

= $D_pF \cdot D_x\Phi^n$.

Therefore, if (p^0, z^0, x^0) is a compatible triple for F and $(q^0, z^0, y^0) = (p^0 D_x \Phi(x^0), z^0, \Phi(x^0))$ is the corresponding for G, then

$$D_{p_n}G(q^0, z^0, y^0) = D_p F(p^0, z^0, x^0) \cdot D_x \Phi^n(x^0);$$

i.e.,

$$D_p F(p^0, z^0, x^0) \cdot \nu(x^0) \neq 0$$

where $\nu(x^0)$ is the normal vector to Γ at x^0 .

PROBLEM 1.2.2. Show that the solution of the quasilinear PDE

$$u_t + a(u)u_x = 0$$

with initial conditions u(x,0) = g(x) is given implicitly by

$$u = g(x - a(u)t).$$

Show that the solution develops a shock (becomes singular) for some t > 0, unless a(g(x)) is a nondecreasing function of x.

SOLUTION.

PROBLEM 1.2.3. Show that the function u(x,t) defined for $t \geq 0$ by

$$u(x,t) = \begin{cases} -\frac{2}{3} \left(t + \sqrt{3x + t^2} \right) & \text{for } 4x + t^2 > 0\\ 0 & \text{for } 4x + t^2 < 0 \end{cases}$$

is an (unbounded) entropy solution of the conservation law $u_t + (u^2/2)_x = 0$ (inviscid Burgers' equation).

1.3 Homework 3

PROBLEM 1.3.1. Consider the initial value problem

$$\begin{cases} u_t = \sin u_x, \\ u(x,0) = \frac{\pi}{4}x. \end{cases}$$

Verify that the assumptions of the Cauchy–Kovalevskaya theorem are satisfied and obtain the taylor series of the solution about the origin.

SOLUTION.

PROBLEM 1.3.2. Consider the Cauchy problem for u(x,y)

$$\begin{cases} u_y = a(x, y, u)u_x + b(x, y, u), \\ u(x, 0) = 0, \end{cases}$$

let a and b be analytic functions of their arguments. Assume that $D^{\alpha}a(0,0,0) \geq 0$ and $D^{\alpha}b(0,0,0) \geq 0$ for all α . (Remember by definition, if $\alpha = 0$ then $D^{\alpha}f = f$.)

- (a) Show that $D^{\beta}u(0,0) \geq 0$ for all $|\beta| \leq 2$.
- (b) Prove that $D^{\beta}u(0,0) \geq 0$ for all $\beta = (\beta_1, \beta_2)$.

Hint: Argue as in the proof of the Cauchy–Kovalevskaya theorem; i.e., use induction in β_2 .

SOLUTION.

PROBLEM 1.3.3. (Kovalevskaya's example) show that the line $\{t=0\}$ is characteristic for the heat equation $u_t = u_{xx}$. Show there does not exist an analytic solution u of the heat equation in $\mathbf{R} \times \mathbf{R}$, with $u = 1/(1+x^2)$ on $\{t=0\}$.

Hint: Assume there is an analytic solution, compute its coefficients, and show instead that the resulting power series diverges in any neighborhood of (0,0).

Solution.

1.4 Homework 4

PROBLEM 1.4.1 (Legendre transform). Let $u(x_1, x_2)$ be a solution of the quasilinear equation

$$a^{11}(Du)u_{x_1x_1} + 2a^{12}(Du)u_{x_1x_2} + a^{22}(Du)u_{x_2x_2} = 0$$

in some region of \mathbb{R}^2 , where we can invert the relations

$$p^1 = u_{x_1}(x_1, x_2), \quad p^2 = u_{x_2}(x_1, x_2)$$

to solve for

$$x^1 = x^1(p_1, p_2), \quad x^2 = x^2(p_1, p_2).$$

Define then

$$v(p) := \mathbf{x}(p) \cdot p - u(\mathbf{x}(p)),$$

where $\mathbf{x} = (x^1, x^2)$, $p = (p_1, p_2)$. Show that v satisfies the *linear* equation

$$a^{22}(p)v_{p_1p_2} - 2a^{12}(p)v_{p_1p_2} + a^{11}(p)v_{p_1p_2} = 0.$$

Hint: See [Evans, 4.4.3b], prove the identities (29).

SOLUTION.

PROBLEM 1.4.2. Find the solution u(x,t) of the one-dimensional wave equation

$$u_{tt} - u_{xx} = 0$$

in the quadrant x > 0, t > 0 for which

$$\begin{cases} u(x,0) = f(x), & u_t(x,0) = g(x), & \text{for } x > 0, \\ u_t(0,t) = \alpha u_x(0,t), & \text{for } t > 0, \end{cases}$$

where $\alpha \neq -1$ is a given constant. Show that generally no solution exists when $\alpha = -1$.

Hint: Use a representation u(x,t) = F(x-t) + G(x+t) for the solution.

SOLUTION.

PROBLEM 1.4.3. (a) Let u be a solution of the wave equation $u_{tt} - c^2 u_{xx} = 0$ for $0 < x < \pi$, t > 0 such that $u(0,t) = u(\pi,t) = 0$. Show that the energy

$$E(t) = \frac{1}{2} \int_0^{\pi} \left(u_t^2 + c^2 u_x^2 \right) dx, \quad t > 0$$

is independent of t; i.e., dE/dt = 0 for t > 0. Assume that u is C^2 up to the boundary.

(b) Express the energy E of the Fourier series solution

$$u(x,t) = \sum_{n=1}^{\infty} (a_n \cos(nct) + b_n \sin(nct)) \sin(nx)$$

in terms of coefficients a_n , b_n .

1.5 Homework 5

PROBLEM 1.5.1. Prove that Laplace's equation $\Delta u = 0$ is rotation invariant; that is, if O is an orthogonal $n \times n$ matrix and we define $v(x) := u(Ox), x \in \mathbf{R}^n$, then $\Delta v = 0$.

Solution.

PROBLEM 1.5.2. Let n=2 and U be the halfplane $\{x_2>0\}$. Prove that

$$\sup_{U} u = \sup_{\partial U} u$$

for $u \in C^2(U) \cap C(\bar{U})$ which are harmonic in U under the additional assumption that u is bounded from above in \bar{U} . (The additional assumption is needed to exclude examples like $u = x_2$.)

Hint: Take for $\epsilon > 0$ the harmonic function

$$u(x_1, x_2) - \epsilon \ln \sqrt{x_1^2 + (x_2 + 1)^2}.$$

Apply the maximum principle to a region $\{x_1^2 + (x_2 + 1)^2 < a^2, x_2 > 0\}$ with large a. Let $\epsilon \to 0$.

SOLUTION.

PROBLEM 1.5.3. Let $U \subseteq \mathbf{R}^n$ be an open set. We say $v \in C^2(U)$ is subharmonic if

$$-\Delta v < 0$$
 in U .

(a) Let $\varphi \colon \mathbf{R}^m \to \mathbf{R}$ be smooth and convex. Assume u^1, \dots, u^m are harmonic in U and

$$v := \varphi(u_1, \dots, u_m).$$

Prove v is subharmonic.

Hint: Convexity for a smooth function $\varphi(z)$ is equivalent to $\sum_{j,k=1}^{m} \varphi_{z_j,z_k}(z)\xi_j\xi_k \geq 0$ for any $\xi \in \mathbf{R}^m$.

(b) Prove $v := |Du|^2$ is subharmonic, whenever u is harmonic. (Assume that harmonic functions are C^{∞} .)

1.6 Homework 6

PROBLEM 1.6.1. For n = 2 find Green's function for the quadrant $U := \{x_1, x_2 > 0\}$ by repeated reflection.

SOLUTION. Taking the hit, set $x' := (x_1, -x_2), x'' := (-x_1, x_2), x''' := (-x_1, -x_2),$ and define

$$\varphi^{x}(y) := \Phi(y - x') + \Phi(y - x'') - \Phi(y - x'''). \tag{1.6.1}$$

We claim that φ^x , as defined above, solves

$$\begin{cases} \Delta \varphi^x = 0 & \text{in } U, \\ \varphi^x(y) = \Phi(y - x) & \text{on } \partial U. \end{cases}$$

It is clear that $\Delta \varphi^x = 0$ since it is built up from the fundamental solutions on \mathbf{R}^n (this follows from the linearity of the Laplace operator). To see that $\varphi^x(y) = \Phi(x-y)$ on ∂U , we do a case by case analysis.

Note that on $\{x_1 = 0\} \subseteq \partial U$, we have

$$\varphi^x(y_1,0) = \Phi(-x_1, y_2 + x_2) + \Phi(-x_1, y_2 - x_2) - \Phi(x_1, y_2 + x_2),$$

where, since the fundamental solution is radial, we have $\Phi(-x_1, y_2 + x_2) = \Phi(x_1, y_2 + x_2)$, and hence the above equals

$$= \Phi(-x_1, y_2 - x_2)$$
$$= \Phi(y - x)$$

and on $\{x_2 = 0\} \subseteq \partial U$, we have

$$\varphi^x(0, y_2) = \Phi(y_1 - x_1, x_2) + \Phi(y_1 + x_1, -x_2) - \Phi(y_1 + x_1, x_2)$$

where, again because Φ is radial, $\Phi(y_1 + x_1, -x_2) = \Phi(y_1 + x_1, x_2)$, thus the above equals

$$= \Phi(y_1 - x_1, x_2)$$
$$= \Phi(y - x).$$

Thus, $\varphi^x(y) = \Phi(y - x)$ on ∂U .

Therefore, Green's function on U is

$$G(x,y) = \Phi(y-x) - \varphi^{x}(y) = \Phi(y-x) - \Phi(y-x') - \Phi(y-x'') + \Phi(y-x''').$$

PROBLEM 1.6.2. (Precise form of Harnack's inequality) Use Poisson's formula for the ball to prove

$$\frac{r^{n-2}(r-|x|)}{(r+|x|)^{n-1}}u(0) \le u(x) \le \frac{r^{n-2}(r+|x|)}{(r-|x|)^{n-1}}u(0)$$

whenever u is positive and harmonic in $B(0,r) = \{ x \in \mathbf{R}^n : |x| < r \}.$

SOLUTION. Recall Poisson's formula for the ball

$$u(x) = \frac{r^2 - |x|^2}{n\alpha_n r} \int_{\partial B(0,r)} \frac{g(y)}{|x - y|^n} dS(y), \tag{1.6.2}$$

where $x \in B(0,r)$ and u solves the boundary-value problem

$$\begin{cases} \Delta u = 0 & \text{in } B(0, r), \\ u = g & \text{on } \partial B(0, r). \end{cases}$$

For fixed $x \in B(0, r)$, write

$$u(x) = r^{n-2}(r+|x|)(r-|x|) \left[\frac{1}{n\alpha_n r^{n-1}} \int_{\partial B(0,r)} \frac{g(y)}{|x-y|^n} \, dS(y) \right].$$

Now, since $r + |x| \ge |x - y| \ge r - |x|$ for all $y \in \partial B(0, r)$, we have

$$\frac{r^{n-2}(r-|x|)}{(r+|x|)^{n-1}} \oint_{\partial B(0,r)} g(y) \, dS(y) \le u(x) \le \frac{r^{n-2}(r+|x|)}{(r-|x|)^{n-1}} \oint_{\partial B(0,r)} g(y) \, dS(y). \tag{1.6.3}$$

Since u = g on the boundary $\partial B(0, r)$, by applying the mean-value property on (1.6.3) we have

$$\frac{r^{n-2}(r-|x|)}{(r+|x|)^{n-1}}u(0) \le u(x) \le \frac{r^{n-2}(r+|x|)}{(r-|x|)^{n-1}}u(0),$$

as desired.

PROBLEM 1.6.3. Let $P_k(x)$ and $P_m(x)$ be homogeneous harmonic polynomials in \mathbb{R}^n of degrees k and m respectively; i.e.,

$$\begin{cases} P_k(\lambda x) = \lambda^k P_k(x), & P_m(\lambda x) = \lambda^m P_m(x) & \text{for every } x \in \mathbf{R}^n, \, \lambda > 0, \\ \Delta P_k = 0, & \Delta P_m = 0 & \text{in } \mathbf{R}^n. \end{cases}$$

(a) Show that

$$\left\{ \begin{array}{l} \displaystyle \frac{\partial P_k}{\partial \nu} = k P_k(x), \quad \frac{\partial P_m}{\partial \nu} = m P_m(x) \quad \text{on } \partial B(0,1), \end{array} \right.$$

where $B(0,1) = \{ x \in \mathbf{R}^n : |x| < 1 \}$ and ν is the outward normal on $\partial B(0,1)$.

(b) Use (a) and Green's formula to prove that

$$\int_{\partial B(0,1)} P_k(x) P_m(x) d\sigma = 0, \quad \text{if } k \neq m.$$

SOLUTION. For part (a), let

$$P_k(x) = \sum_{|\alpha|=k} a_{\alpha} x^{\alpha}.$$

Then, since $\nu = (x_1, \dots, x_n)$, the derivative along ν is given by

$$\frac{\partial P_k(x)}{\partial \nu} = \sum_{j=1}^n (P_k)_{x_j} x_j$$

$$= \sum_{j=1}^n \left(\sum_{|\alpha|=k} a_{\alpha} x^{\alpha} \right)_{x_j} x_j$$

$$= \sum_{j=1}^n \left(\sum_{l=1}^m a_{\alpha} x_1^{\alpha_1^l} \cdots x^{\alpha_j^l} \cdots x^{\alpha_n^l} \right)_{x_j} x_j$$

where $\sum_{j=1}^{n} \alpha_j^l = k$ and $1 \leq j \leq \binom{n+k-1}{n} =: m$ (by the stars and bars theorem)

$$= \sum_{j=1}^{n} \sum_{l=1}^{m} \left(\alpha_j^l a_{\alpha} x_1^{\alpha_1^l} \cdots x_{\alpha_j^{l-1}}^{\alpha_j^l} \cdots x_{\alpha_n^{l}}^{\alpha_n^l} \right) x_j$$

$$= \sum_{j=1}^{n} \sum_{l=1}^{m} \alpha_j^l a_{\alpha} x_1^{\alpha_1^l} \cdots x_{\alpha_j^{l}}^{\alpha_j^l} \cdots x_{\alpha_n^{l}}^{\alpha_n^l}$$

$$= \sum_{j=1}^{n} \sum_{l=1}^{m} \alpha_j^l a_{\alpha} x_1^{\alpha_n^l}$$

switching the order of summation, we have

$$= \sum_{l=1}^{m} \sum_{j=1}^{n} \alpha_{j}^{l} a_{\alpha} x^{\alpha}$$

$$= \sum_{l=1}^{m} k a_{\alpha} x^{\alpha}$$

$$= k \sum_{l=1}^{m} a_{\alpha} x^{\alpha}$$

$$= k P_{k}(x).$$

This argument, of course, applies to every $k \in \mathbb{N}$. For part (b), by Green's theorem, we have

$$\int_{B(0,r)} P_k(x) \Delta P_m(x) - (\Delta P_k(x)) P_m(x) dx = \int_{\partial B(0,r)} P_k(x) \frac{\partial}{\partial \nu} P_m(x) - \frac{\partial}{\partial \nu} P_k(x) P_m(x) dS(x)$$

$$= \int_{\partial B(0,r)} (m-k) P_k(x) P_m(x) dS(x),$$

where the left-hand side is equal to zero since both ΔP_k and ΔP_m are zero. Since $m \neq k$, it must be the case that

$$\int_{\partial B(0,r)} P_k(x) P_m(x) dS(x) = 0.$$

1.7 Homework 7

PROBLEM 1.7.1. Solve the Dirichlet problem for the Laplace equation in \mathbb{R}^2

$$\begin{cases} \Delta u = 0 & \text{in } 1 < |x| < 2 \\ u = x_1 & \text{on } |x| = 1, \\ u = 1 + x_1 x_2 & \text{on } |x| = 2. \end{cases}$$

Hint: Use Laurent series.

SOLUTION. First, let us make the change of variables $(x_1, x_2) \mapsto re^{i\theta}$ to the Dirichlet problem in question:

$$\begin{cases} \Delta u = 0 & \text{in } 1 < r < 2, \\ u = \frac{1}{2} (e^{i\theta} + e^{-i\theta}) & \text{on } r = 1, \\ u = 1 + \frac{1}{i} (e^{i2\theta} - e^{-i2\theta}) & \text{on } r = 2. \end{cases}$$
(1.7.1)

Now, suppose u is a solution, of the form

$$u(re^{i\theta}) = b \ln r + \sum_{n=-\infty}^{\infty} (a_n r^n + \overline{a_{-n}} r^{-n}) e^{in\theta},$$

to the problem (1.7.1). It is easy to see that u is in fact harmonic:

$$\Delta u = u_{rr} + \frac{1}{r}u_r + \frac{1}{r^2}u_{\theta\theta}$$

$$= -br^{-2} + br^{-2} + \sum_{n=-\infty}^{\infty} \left[\left(n(n-1) + n - n^2 \right) a_n r^n + \left(n(n-1) + n - n^2 \right) \overline{a_{-n}} r^{-n} \right] e^{in\theta}$$

$$= 0.$$

Next we use the boundary data to compute the coefficients a_n , $n \in \mathbb{Z}$. Using the data (1.7.1), on r = 1 we have

$$\frac{1}{2}(e^{i\theta} + e^{-i\theta}) = \sum_{n=-\infty}^{\infty} (a_n + \overline{a_{-n}})e^{in\theta},$$

and on r=2

$$1 + \frac{1}{i} (e^{i2\theta} - e^{-i2\theta}) = b \ln 2 + \sum_{n = -\infty}^{\infty} (2^n a_n + 2^{-n} a_{-n}) e^{in\theta}.$$

These equations immediately tell us that $b = 1/\ln 2$. Moreover, the following relations on the coefficients hold

$$\begin{cases} \frac{1}{2} = a_1 + \overline{a_{-1}} & \frac{1}{2} = a_{-1} + \overline{a_1}, \\ \frac{1}{i} = 2^2 a_2 + 2^{-2} \overline{a_{-2}}, & -\frac{1}{i} = 2^2 a_{-2} + 2^{-2} \overline{a_2}, \\ 0 = a_n + \overline{a_{-n}} & \text{for } n \neq \pm 1, \\ 0 = 2^n a_n + 2^{-n} \overline{a_{-n}} & \text{for } n \neq \pm 2. \end{cases}$$

A little calculation shows that

$$\begin{cases} a_1 = -\frac{1}{6}, & a_{-1} = \frac{2}{3}, \\ a_2 = -\frac{4i}{15}, & a_{-2} = -\frac{4i}{15}, \\ a_n = 0 & \text{for } n \neq \pm 1, \pm 2. \end{cases}$$

Thus,

$$\begin{split} u(r\mathrm{e}^{\mathrm{i}\theta}) &= \frac{1}{\ln 2} \ln r + \left(-\frac{4\mathrm{i}}{15} r^{-2} + \frac{4\mathrm{i}}{15} r^2 \right) \mathrm{e}^{-\mathrm{i}2\theta} + \left(\frac{2}{3} r^{-1} - \frac{1}{6} r \right) \mathrm{e}^{-\mathrm{i}\theta} \\ &\quad + \left(-\frac{1}{6} r + \frac{2}{3} r^{-1} \right) \mathrm{e}^{\mathrm{i}\theta} + \left(-\frac{4\mathrm{i}}{15} r^2 + \frac{4\mathrm{i}}{15} r^{-2} \right) \mathrm{e}^{\mathrm{i}2\theta} \\ &= \frac{1}{\ln 2} \ln r - \frac{8}{15} r^{-4} \left(\frac{r^2 \mathrm{e}^{\mathrm{i}2\theta} - r^2 \mathrm{e}^{-\mathrm{i}2\theta}}{2\mathrm{i}} \right) + \frac{8}{15} \left(\frac{r^2 \mathrm{e}^{\mathrm{i}2\theta} - r^2 \mathrm{e}^{-\mathrm{i}2\theta}}{2\mathrm{i}} \right) \\ &\quad + \frac{4}{3} r^{-2} \left(\frac{r \mathrm{e}^{\mathrm{i}\theta} + r \mathrm{e}^{-\mathrm{i}\theta}}{2} \right) - \frac{1}{3} \left(\frac{r \mathrm{e}^{\mathrm{i}\theta} + r \mathrm{e}^{-\mathrm{i}\theta}}{2} \right). \end{split}$$

Substituting back, we have

$$u(x_1, x_2) = \frac{1}{\ln 2} \ln(x_1^2 + x_2^2) - \frac{16x_1x_2}{15(x_1^2 + x_2^2)^2} + \frac{16x_1x_2}{15} + \frac{4x_1}{3(x_1^2 + x_2^2)} - \frac{x_1}{3}.$$
 (1.7.2)

It is easy to see that (1.7.2) satisfies the boundary data at |x| = 1 and |x| = 2.

PROBLEM 1.7.2. Let Ω be a bounded domain with a C^1 boundary, $g \in C^2(\partial \Omega)$ and $f \in C(\bar{\Omega})$. Consider the so called *Neumann problem*

$$\begin{cases}
-\Delta u = f & \text{in } \Omega, \\
\frac{\partial u}{\partial \nu} = g & \text{on } \partial \Omega,
\end{cases}$$
(*)

where ν is the outer normal on $\partial\Omega$. Show that the solution of (*) in $C^2(\Omega) \cap C^1(\bar{\Omega})$ is unique up to a constant; i.e., if u_1 and u_2 are both solutions of (*), then $u_2 = u_1 + \text{const.}$ in Ω .

Hint: Look at the proof of the uniqueness for the Dirichlet problem by energy methods [E, 2.2.5a].

SOLUTION. Suppose u_1 and u_2 are solutions to the Neumann problem (*). Define $v := u_1 - u_2$. Then v is harmonic in Ω and $\partial v/\partial \nu = 0$ on $\partial \Omega$. Consider the energy functional

$$E[v] = \frac{1}{2} \int_{\Omega} |Dv|^2 dx.$$

By Green's formula version (ii),

$$E[v] = \frac{1}{2} \int_{\Omega} |Dv|^2 dx$$
$$= -\frac{1}{2} \int_{\Omega} v \Delta v \, dx + \int_{\partial U} \frac{\partial v}{\partial \nu} v \, dS(x)$$
$$= 0.$$

This implies that $|Dv|^2 = Dv \cdot Dv = 0$ which, since the standard inner product on \mathbb{R}^n is positive-definite, implies that $Dw \equiv 0$. It follows that $u_1 = u_2 + \text{const}$, i.e., the solution u to (*) is unique up to a constant.

PROBLEM 1.7.3. Write down an explicit formula for a solution of

$$\begin{cases} u_t - \Delta_x u + cu = f & \text{in } \mathbf{R}^n \times (0, \infty), \\ u = g & \text{on } \mathbf{R}^n \times \{t = 0\}, \end{cases}$$

where $c \in \mathbf{R}$.

Hint: Rewrite the problem in terms of $v(x,t) := e^{ct}u(x,t)$.

SOLUTION. Taking the hint, let us rewrite the problem in terms of $v(x,t) = e^{ct}u(x,t)$:

$$\begin{cases} v_t - \Delta_x v = e^{ct} f & \text{in } \mathbf{R}^n \times (0, \infty), \\ v = g & \text{on } \mathbf{R}^n \times \{t = 0\}. \end{cases}$$
 (1.7.3)

By Duhamel's principle, the problem (1.7.3) is solved by

$$v(x,t) = \int_{\mathbf{R}^n} \Phi(x-y,t)g(y) \, dy + \int_0^t \int_{\mathbf{R}^n} \Phi(x-y,t-s)e^{cs} f(y,s) \, dy ds,$$

where Φ is the fundamental solution to the heat equation. Thus, the formula

$$u(x,t) = e^{-ct} v(x,t) = e^{-ct} \int_{\mathbf{R}^n} \Phi(x-y,t) g(y) \, dy + e^{-ct} \int_0^t \int_{\mathbf{R}^n} \Phi(x-y,t-s) e^{cs} f(y,s) \, dy ds$$

solves the original problem.

2 Homework 8

PROBLEM 2.0.1. Show that the function

$$u(x,t) := \sum_{k=-\infty}^{\infty} (-1)^k \Phi(x-2k,t)$$

where

$$\Phi(x,t) = \frac{e^{-\frac{x^2}{4t}}}{\sqrt{4\pi t}}$$

is positive for |x| < 1, t > 0.

Hint: Show that u satisfies $u_t = u_{xx}$ for t > 0,

$$\begin{cases} u = 0 & \text{on } \{ |x| = 1 \} \times \{ t \ge 0 \}, \\ u = \delta_0 & \text{on } \{ |x| \le 1 \} \times \{ t = 0 \}. \end{cases}$$

Then, carefully apply the maximum/minimum principle in a domain $\{|x| \le 1\} \times \{\epsilon \le t \le T\}$ for small $\epsilon > 0$ and large T > 0 pass to the limit as $\epsilon \to 0+$ and $T \to \infty$.

SOLUTION. Taking the hint, let us verify that $u_t = u_{xx}$, for t > 0. By direct computation, we have

$$\Phi_{x}(x,t) = \frac{\partial}{\partial x} \left(\frac{e^{-\frac{x^{2}}{4t}}}{\sqrt{4\pi t}} \right) \qquad \Phi_{xx}(x,t) = \frac{\partial}{\partial x} \left(-\frac{xe^{-\frac{x^{2}}{4t}}}{2\sqrt{4\pi}t^{\frac{3}{2}}} \right) \\
= -\frac{xe^{-\frac{x^{2}}{4t}}}{2\sqrt{4\pi}t^{\frac{3}{2}}}, \qquad \qquad = \frac{x^{2}e^{-\frac{x^{2}}{4t}}}{4\sqrt{4\pi}t^{\frac{5}{2}}} - \frac{e^{-\frac{x^{2}}{4t}}}{2\sqrt{4\pi}t^{\frac{3}{2}}} \\
= \frac{(x^{2} - 2t)e^{-\frac{x^{2}}{4t}}}{4\sqrt{4\pi}t^{\frac{5}{2}}},$$

and

$$\begin{split} \varPhi_t(x,t) &= \frac{\partial}{\partial t} \left(\frac{\mathrm{e}^{-\frac{x^2}{4t}}}{\sqrt{4\pi t}} \right) \\ &= \frac{x^2 \mathrm{e}^{-\frac{x^2}{4t}}}{4\sqrt{4\pi t^{\frac{5}{2}}}} - \frac{\mathrm{e}^{-\frac{x^2}{4t}}}{2\sqrt{4\pi t^{\frac{3}{2}}}} \\ &= \frac{(x^2 - 2t)\mathrm{e}^{-\frac{x^2}{4t}}}{4\sqrt{4\pi t^{\frac{5}{2}}}}. \end{split}$$

Since $\Phi_t = \Phi_{xx}$ it follows that $u_t = u_{xx}$ (assuming uniform convergence of u).

Next we show that u = 0 on $\{|x| = 1\} \times \{t \ge 0\}$ and $u = \delta_0$ on $\{|x| = 1\} \times \{t = 0\}$. To show u = 0 fix a $t \ge 0$ and, after relabeling if necessary, assume that x = 1 which gives us

$$u(1,t) = \sum_{k=-\infty}^{\infty} (-1)^k \frac{e^{-\frac{(1-2k)^2}{4t}}}{\sqrt{4\pi t}}$$
$$= \frac{1}{\sqrt{4\pi t}} \left(\dots - e^{-\frac{9}{4t}} + e^{-\frac{1}{4t}} - e^{-\frac{1}{4t}} + e^{-\frac{9}{4t}} + \dots \right)$$
$$= 0.$$

Similarly for u(-1,t) = 0. For $u(|x| \le 1,0)$, we have a

$$u(|x| \le 1, 0) = \sum_{k=-\infty}^{\infty} (-1)^k \lim_{t \to 0+} \left[e^{-\frac{(x-2k)^2}{4t}} / \sqrt{4\pi t} \right]$$
$$= \sum_{k=-\infty}^{\infty} (-1)^k \delta_0(x-2k)$$
$$= \delta_0(x)$$

since $|x| \leq 1$ and values δ_0 is zero for values x - 2k outside of the interval [-1, 1].

At last we show that u is positive for |x| < 1, t > 0. Seeking a contradiction, suppose u is negative on some point (x_0, t_0) in $\{|x| < 1\} \times \{\epsilon \le t \le T\}$. Then by the minimum principle, u achieves its minimum somewhere on the bottom boundary $\{|x| = 1\} \times \{t = \epsilon\}$. Therefore, there exists a sequence $(x_n, t_n +) \to (x, 0)$, where $|x_n|, |x| < 1$, such that u(x, 0) < 0. However, we have shown above that $u(x, 0) = \delta_0(x)$ for |x| < 1; i.e., either u(x, 0) = 0 or $u(x, 0) = +\infty$. This is a contradiction. Therefore, it must be the case that $u \ge 0$ for |x| < 1, t > 0.

Problem 2.0.2 (Tikhonov's example). Let

$$g(t) := \begin{cases} e^{-t^{-2}} & t > 0, \\ 0 & t \le 0. \end{cases}$$

Then $g \in C^{\infty}(\mathbf{R})$ and we define

$$u(x,t) := \sum_{k=0}^{\infty} \frac{g^{(k)}(t)}{(2k)!} x^{2k}.$$

Assuming that the series is convergent, show that u(x,t) solves the heat equation in $\mathbf{R} \times (0,\infty)$ with the initial condition u(x,0) = 0, $x \in \mathbf{R}$. Why doesn't this contradict the uniqueness theorem for the initial value problem?

Solution. Let u be as above. Then

$$u_t(x,t) = \frac{\partial}{\partial t} \left(\sum_{k=0}^{\infty} \frac{g^{(k)}(t)}{(2k)!} x^{2k} \right)$$
$$= \sum_{k=0}^{\infty} \frac{g^{(k+1)}(t)}{(2k)!} x^{2k}$$
$$= \sum_{k=2}^{\infty} \frac{g^{(k)}(t)}{(2k-2)!} x^{2k-2},$$

and

$$u_{x}(x,t) = \frac{\partial}{\partial x} \left(\sum_{k=0}^{\infty} \frac{g^{(k)}(t)}{(2k)!} x^{2k} \right) \qquad u_{xx}(x,t) = \frac{\partial}{\partial x} \left(\sum_{k=0}^{\infty} \frac{g^{(k)}(t)}{(2k-1)!} x^{2k-1} \right)$$

$$= \sum_{k=0}^{\infty} \frac{g^{(k)}(t)}{(2k)!} 2kx^{2k-1} \qquad \qquad = \sum_{k=1}^{\infty} \frac{g^{(k)}(t)}{(2k-1)!} (2k-1)x^{2k-2} + \frac{\partial}{\partial x} g^{(0)}(t)$$

$$= \sum_{k=1}^{\infty} \frac{g^{(k)}(t)}{(2k-1)!} x^{2k-1}, \qquad \qquad = \sum_{k=2}^{\infty} \frac{g^{(k)}(t)}{(2k-2)!} x^{2k-2}.$$

Thus, $u_t - \Delta u = 0$; i.e., u solves the heat equation. As this example shows, unless some assumptions on u such as subexponential (cf. [E §2.3], Theorem 7) growth is assumed.

PROBLEM 2.0.3. Evaluate the integral

$$\int_{-\infty}^{\infty} \cos(ax) e^{-x^2} dx, \qquad (a > 0).$$

Hint: Use the separation of variables to find the solution of the corresponding initial-value problem for the heat equation.

SOLUTION. By separation of variables,

$$u(x,t) = \cos(ax)e^{-a^2t}$$

is a solution to the Cauchy problem

$$\begin{cases} u_t - \Delta u = 0 & \text{in } \mathbf{R} \times (0, \infty), \\ u = \cos(ax) & \text{on } \mathbf{R} \times \{t = 0\}. \end{cases}$$

However, the convolution

$$u(x,t) = \frac{1}{\sqrt{4\pi t}} \int_{-\infty}^{\infty} \cos(ay) e^{-\frac{|x-y|^2}{4t}} dy$$

is also a solution to the Cauchy problem. Now note that

$$\int_{-\infty}^{\infty} \cos(ay) e^{-y^2} dy = \sqrt{\pi} \cdot u(0, \frac{1}{4})$$
$$= \sqrt{\pi} e^{-\frac{a^2}{4}}.$$

PROBLEM 2.0.4. (a) Show that for n = 3 the general solution to the wave equation $u_{tt} - \Delta u = 0$ with spherical symmetry about the origin has the form

$$u = \frac{1}{r}F(r+t) + \frac{1}{r}G(r-t), \quad r = |x|,$$

with suitable F and G.

(b) Show that the solution with initial data of the form

$$u(r,0) = 0, \quad u_t(r,0) = h(r)$$

(h is an even function of r) is given by

$$u = \frac{1}{2r} \int_{r-t}^{r+t} \rho h(\rho) \, d\rho.$$

SOLUTION.

PROBLEM 2.0.5. Show that the solution $w(x_1,t)$ of the initial-value problem for the Klein-Gordon equation

$$\begin{cases} w_{tt} = w_{x_1 x_1} - \lambda^2 w, \\ w(x_1, 0) = 0, & w_t(x_1, 0) = h(x_1) \end{cases}$$
 (2.0.1)

is given by

$$w(x_1,t) = \frac{1}{2} \int_{x_1-t}^{x_1+t} J_0(\lambda s) h(y_1) \, dy_1.$$

Here $s^2 = t^2 - (x_1 - y_1)^2$, while J_0 denotes the Bessel function defined by

$$J_0(z) := \frac{2}{\pi} \int_0^{\frac{\pi}{2}} \cos(z \sin \theta) d\theta.$$

Hint: Descend to (2.0.1) from the two-dimensional wave equation satisfied by

$$u(x_1, x_2, t) = \cos(\lambda x_2) w(x_1, t).$$

Solution.

Problem 2.0.6. Let u solve

$$\begin{cases} u_{tt} - \Delta u = 0 & \text{in } \mathbf{R}^3 \times (0, \infty), \\ u = g, & u_t = h & \text{on } \mathbf{R}^3 \times \{t = 0\} \end{cases}$$

where g and h are smooth and have compact support. Show there exists a constant C such that

$$|u(x,t)| \le Ct^{-1} \quad (x \in \mathbf{R}^3, t > 0).$$

Solution.

3 Exams

3.1 Midterm Practice Problems

Problem 3.1.1. Solve $u_{x_1}^2 + x_2 u_{x_2} = u$ with initial conditions $u(x_1, 1) = x_1^2/4 + 1$.

SOLUTION. By inspection, we may suspect that $v(x_1, x_2) = x_1^2/4 + x_2$ is a solution to the PDE. It certainly satisfies the boundary condition. A routine calculation shows that v is in fact a solution to the PDE. Lucky guess!

More formally, let us solve this problem using the method of characteristics. First, write

$$F(p, z, x) = (p^{1}(s))^{2} + x^{2}(s)p^{2}(s) - z(s) = 0.$$

Then, the characteristic ODEs are

$$\begin{cases} \left(\dot{p}^{1}(s),\dot{p}^{2}(s)\right) = -(0,p^{2}(s)) + (p^{1}(s),p^{2}(s)) \\ = (p^{1}(s),0), \\ \dot{z}(s) = (2p^{1}(s),x^{2}(s)) \cdot (p^{1}(s),p^{2}(s)) \\ = 2p^{1}(s)^{2} + x^{2}(s)p^{2}(s), \\ \left(\dot{x}^{1}(s),\dot{x}^{2}(s)\right) = (2p^{1}(s),x^{2}(s)). \end{cases}$$

Now, for $(x^1(0), x^2(0)) = (x^0, 1)$, integrating the characteristics, we get

$$\begin{cases} (p^{1}(s), p^{2}(s)) = (p_{0}^{1}e^{s}, p_{0}^{2}), \\ (x^{1}(s), x^{2}(s)) = (2p_{0}^{1}e^{s} + x_{0}^{1}, x_{0}^{2}e^{s}), \\ z(s) = \frac{(x^{0})^{2}}{4}e^{2s} + p_{0}^{2}e^{s} + z^{0} \end{cases}$$

Using the initial condition and the PDE, we find that

$$p_0^1 = \frac{x^0}{2}, \quad p_0^2 = \frac{\left(x^0\right)^2}{4} + 1 - \frac{\left(x^0\right)^2}{4} = 1,$$

$$x_0^1 = 0, \qquad x_0^2 = 1$$

$$z^0 = 0,$$

and consequently

$$\begin{cases} (x^{1}(s), x^{2}(s)) = (x^{0}e^{s}, e^{s}), \\ z(s) = \frac{(x^{0})^{2}}{4}e^{2s} + e^{s} \end{cases}$$

so, rewriting z in terms of (x^1, x^2) , we have

$$z(s) = \frac{(x^0)^2}{4} e^{2s} + e^s$$
$$= \frac{(x^1(s))^2}{4} + x^2(s),$$

so the solution in terms of (x_1, x_2) , is

$$u(x_1, x_2) = \frac{x_1^2}{4} + x_2,$$

just as we suspected.

PROBLEM 3.1.2. Find the maximal $t_0 > 0$ for which the (classical) solution of the Cauchy problem

$$\begin{cases} uu_x + u_t = 0, \\ u(x,0) = e^{-\frac{x^2}{2}}, \end{cases}$$

exists in $\mathbf{R} \times [0, t)$; i.e., the first time $t = t_0$ when the shock develops.

SOLUTION. First, let us find a solution to the PDE using the method of characteristics. Write

$$F(p, z, x) = z(s)p^{1}(s) + p^{2}(s).$$

Then, the characteristic ODEs are

$$\begin{cases} \left(\dot{p}^{1}(s), \dot{p}^{2}(s)\right) = -(0,0) - p^{1}(p^{1}(s), p^{2}(s)) \\ = \left(-p^{1}(s)^{2}, -p^{1}(s)p^{2}(s)\right), \\ \dot{z}(s) = \left(z(s), 1\right) \cdot \left(p^{1}(s), p^{2}(s)\right) \\ = z(s)p^{1}(s) + p^{2}(s) \\ = 0, \\ \left(\dot{x}(s), \dot{t}(s)\right) = \left(z(s), 1\right). \end{cases}$$

Thus, integrating the characteristic ODEs from $(x^0, 0)$, we have

$$\begin{cases} \dot{z}(s) = z^{0}, \\ (x(s), t(s)) = (z^{0}s + x^{0}, s); \end{cases}$$

since the PDE is quasilinear, we disregard (p^1, p^2) .

Applying the boundary conditions, we see that

$$z^0 = u(x^0, 0) = e^{-(x^0)^2/2}$$
.

Here's where it gets tricky. After a little struggling, we see that there is really no way to solve for z in terms of (x(s), t(s)). However, we can solve for the projected characteristics:

$$(x(t,y),t) = (e^{-y^2/2}t + y,t);$$

and this is really all that matters for us to find the time t_0 when the shock develops, i.e., the time when the projected characteristic fails to be injective.

A little calculation shows that this happens precisely when $t = e^{-1/2}$.

PROBLEM 3.1.3. If ρ_0 denotes the maximum density of cars on a highway (i.e., under bumpet-to-bumper conditions), then a reasonable model for traffic density ρ is given by

$$\begin{cases} \rho_t + (F(\rho))_x = 0, \\ F(\rho) = c\rho \left(1 - \frac{\rho}{\rho_0}\right), \end{cases}$$

where c > 0 is a constant (free speed of highway). Suppose the initial density is

$$\rho(x,0) = \begin{cases} \frac{1}{2}\rho_0 & \text{if } x < 0, \\ \rho_0 & \text{if } x > 0. \end{cases}$$

Find the shock curve and describe the weak solution. (Interpret your result for the traffic flow.)

SOLUTION. First, note that

$$(F(\rho))_x = F'(\rho)\rho_x$$

$$= \left[-c\frac{\rho}{\rho_0} + c\left(1 - \frac{\rho}{\rho_0}\right) \right]\rho_x$$

$$= \left(c - \frac{2c\rho}{\rho_0}\right)\rho_x.$$

Let us find a solution to the PDE using the method of characteristics. Write

$$G(p, z, x) = p^{2}(s) + F'(z(s))p^{1}(s).$$

Then, the characteristic ODEs are

$$\begin{cases} \left(\dot{p}^{1}(s), \dot{p}^{2}(s)\right) = \left(-F''(z(s))p^{1}(s), -F''(z(s))p^{2}(s)\right), \\ \dot{z}(s) = F'(z(s))p^{1}(s) + p^{2}(s) \\ = 0, \\ \left(\dot{x}^{1}(s), \dot{x}^{2}(s)\right) = \left(F'(z(s)), 1\right). \end{cases}$$

Now, integrating the characteristics, we have

$$\begin{cases} z(s) = z^0, \\ \left(x^1(s), x^2(s)\right) = \left(F'(z^0)s + x^0, s\right). \end{cases}$$

We have two cases to consider, $x^0 < 0$ or $x^0 > 0$. For $x^0 < 0$, $z^0 = \frac{\rho_0}{2}$ and the projected characteristics look like

$$(F'(\frac{\rho_0}{2})t + x^0, t) = \left(\left[c - \frac{2c(\frac{\rho_0}{2})}{\rho_0} \right] t + x^0, t \right)$$

= $(0 \cdot t + x^0, t)$
= (x^0, t)

(where we have replaced s with the more appropriate t). Whereas for $x^0 > 0$, we have

$$(F'(\rho_0)t + x^0, t) = \left(\left[c - \frac{2c\rho_0}{\rho_0} \right] t + x^0, t \right)$$

= $(-ct + x^0, t)$.

These characteristics intersect precisely when

$$t = \frac{x_1^0 - x_2^0}{c},$$

where $x_1^0 > 0$, $x_2^0 < 0$.

PROBLEM 3.1.4. Find the characteristics of the second order equation

$$u_{xx} - (2\cos x)u_{xy} - (3 + \sin^2 x)u_{yy} - yu_y = 0,$$

and transform it to the canonical form.

SOLUTION. First, writing the PDE in the form

$$Au_{xx} + 2Bu_{xy} + Cu_{yy} + 2Du_x + 2Eu_y + Fu = 0,$$

we see that $A=1, B=-\cos x, C=-3\sin^2 x$, and $E=-\frac{y}{2}$. We solve for the characteristic curve by find a solution to the ODEs

$$\frac{dy}{dx} = \frac{B \pm \sqrt{B^2 - AC}}{A}$$
$$= -\cos x \pm \sqrt{\cos^2 x + 3 + \sin^2 x}$$
$$= -\cos x \pm 2.$$

The solutions give us the following ODEs

$$\begin{cases} y = -\sin x + 2x + \xi(x, y), \\ y = -\sin x - 2x + \eta(x, y). \end{cases}$$

Integrating these equations, we have

$$\begin{cases} \xi(x,y) = y + \sin x - 2x, \\ \eta(x,y) = y + \sin x + 2x. \end{cases}$$

These are the characteristic strips for the PDE.

To put this PDE in canonical form, we first compute the following partial derivatives

$$u_{x} = u_{\xi}\xi_{x} + u_{\eta}\eta_{x},$$

$$u_{y} = u_{\xi}\xi_{y} + u_{\eta}\eta_{y},$$

$$u_{xx} = u_{\xi}\xi_{xx} + u_{\eta}\eta_{xx} + (u_{\xi\xi}\xi_{x} + u_{\xi\eta}\eta_{x})\xi_{x} + (u_{\xi\eta}\xi_{x} + u_{\eta\eta}\eta_{x})\eta_{x}$$

$$= u_{\xi\xi}(\xi_{x})^{2} + u_{\eta\eta}(\eta_{x})^{2} + 2u_{\xi\eta}\xi_{x}\eta_{x} + u_{\xi}\xi_{xx} + u_{\eta\eta}\eta_{xx},$$

exploiting symmetry, we can find u_{yy} by replacing x with y above

$$u_{yy} = u_{\xi\xi}(\xi_y)^2 + u_{\eta\eta}(\eta_y)^2 + 2u_{\xi\eta}\xi_y\eta_y + u_{\xi}\xi_{yy} + u_{\eta}\eta_{yy},$$

the last thing we need to figure out is the mixed partial

$$u_{xy} = u_{\xi}\xi_{xy} + u_{\eta}\eta_{xy} + (u_{\xi\xi}\xi_y + u_{\xi\eta}\eta_y)\xi_x + (u_{\xi\eta}\xi_y + u_{\eta\eta}\eta_y)\eta_x$$

= $u_{\xi\xi}\xi_x\xi_y + u_{\eta\eta}\eta_x\eta_y + u_{\xi\eta}(\xi_x\eta_y + \xi_y\eta_x) + u_{\xi}\xi_{xy} + u_{\eta}\eta_{xy}.$

Now find the partials $\xi_x, \eta_x, \xi_y, \eta_y, \xi_{xy}, \dots$, etc.

$$\xi_x = \cos x - 2,$$
 $\eta_x = \cos x + 2,$ $\xi_{xx} = -\sin x,$ $\eta_{xx} = -\sin x,$ $\xi_{xy} = 0,$ $\eta_{xy} = 0,$ $\eta_{y} = 1,$ $\eta_{yy} = 0.$

Thus,

$$\begin{cases} u_x = (\cos x - 2)u_{\xi} + (\cos x + 2)u_{\eta}, \\ u_y = u_{\xi} + u_{\eta}, \\ u_{xx} = (\cos x - 2)^2 u_{\xi\xi} + (\cos x + 2)^2 u_{\eta\eta} \\ + 2(\cos x + 2)(\cos x - 2)u_{\xi\eta} - (\sin x)u_{\xi} - (\sin x)u_{\eta} \\ = (\cos^2 x - 4\cos x + 4)u_{\xi\xi} + (\cos^2 x + 4\cos x + 4)u_{\eta\eta} \\ + 2(\cos^2 x - 4)u_{\xi\eta} - (\sin x)u_{\xi} - (\sin x)u_{\eta} \\ u_{yy} = u_{\xi\xi} + u_{\eta\eta} + 2u_{\xi\eta}, \\ u_{xy} = (\cos x - 2)u_{\xi\xi} + (\cos x + 2)u_{\eta\eta} + 2(\cos x)u_{\xi\eta}, \end{cases}$$

so the canonical form is

$$0 = u_{xx} - (2\cos x)u_{xy} - (3\sin^2 x)u_{yy} - yu_y$$

$$= \xi^2 u_{\xi\xi} + \eta^2 u_{\eta\eta}$$

$$+ 2\xi \eta u_{\xi\eta} - (\sin x)u_{\xi} - (\sin x)u_{\eta}$$

$$- (2\cos x) ((\cos x - 2)u_{\xi\xi} + (\cos x + 2)u_{\eta\eta} + 2(\cos x)u_{\xi\eta})$$

$$- (3\sin^2 x)(u_{\xi\xi} + u_{\eta\eta} + 2u_{\xi\eta})$$

$$- y(u_{\xi} + u_{\eta})$$

Who cares.

PROBLEM 3.1.5. Let $Lu := u_{xx} - 4u_{yy} + \sin(y + 2x)u_x = 0$.

- (a) Consider the level curve $\Gamma = \{(x,y) : \varphi(x,y) = C\}$ where $|D\varphi| \neq 0$ on Γ . Define what it means for Γ to be characteristic with respect to L at a point $(x_0, y_0) \in \Gamma$.
- (b) Find the points at which the curve $x^2 + y^2 = 5$ is characteristic.

(c) Is it true that every smooth simple closed curve Γ in \mathbf{R}^2 has at least one point at which it is characteristic with respect to L?

Solution.

PROBLEM 3.1.6. Consider the second order equation

$$Lu := u_{xx} - 2xu_{xy} + x^2u_{yy} - 2u_y = 0.$$

- (a) Find the characteristic curves of Lu = 0. What is the type of this equation?
- (b) Find the points on the line $\Gamma := \{ (x, y) \in \mathbf{R}^2 : x + y = 1 \}$ at which Γ is characteristic with respect to Lu = 0.

SOLUTION.

PROBLEM 3.1.7. Solve the initial boundary value problem for the equation $u_{tt} = u_{xx}$ in $\{x > 0, t > 0\}$ satisfying

$$\begin{cases} u(x,0) = \sin^2 x, & u_t(x,0) = \sin x, \\ u(0,t) = 0. \end{cases}$$

SOLUTION.

PROBLEM 3.1.8. Consider the initial/boundary value problem

$$\begin{cases} u_{tt} - u_{xx} = 0 & \text{for } 0 < x < \pi, \ t > 0, \\ u(x,0) = x, & u_t(x,0) = 0 & \text{for } 0 < x < \pi, \\ u_x(0,t) = 0, & u_x(\pi,t) = 0 & \text{for } t > 0. \end{cases}$$

- (a) Find a weak solution of the problem.
- (b) Is the solution unique? Continuous? C^1 ?

SOLUTION.

PROBLEM 3.1.9. Let B_1^+ denote the open half-ball $\{x \in \mathbf{R}^n : |x| < 1, x_n > 0\}$. Assume $u \in C(\bar{B}_1^+)$ is harmonic in B_1^+ with u = 0 on $\partial B_1^+ \cap \{x_n = 0\}$. Set

$$v(x) := \begin{cases} u(x) & \text{if } x_n \ge 0, \\ -u(x_1, \dots, x_{n-1}, -x_n) & \text{if } x_n < 0, \end{cases}$$

for $x \in B_1$. Prove v is harmonic in B_1 .

Hint: It will be enough to prove that $\int_B \nabla v \nabla \eta \, dx = 0$ for any test function $\eta \in C_0^{\infty}(B_1)$. Split $\int_{B_1} = \int_{B_1^+} + \int_{B_1^-}$ and apply the integration by parts formula to each of $\int_{B_1^{\pm}}$.

SOLUTION.

PROBLEM 3.1.10. Let u and v be harmonic functions in the unit ball $B_1 \subseteq \mathbf{R}^n$. What can you conclude about u and v if

- (a) $D^{\alpha}u(0) = D^{\alpha}v(0)$ for every multiindex α ?
- (b) $u(x) \leq v(x)$ for every $x \in B_1$ and u(0) = v(0)?

Justify your answer in each case.

SOLUTION.

PROBLEM 3.1.11. Let Φ be the fundamental solution of the Laplace equation in \mathbf{R}^n and $f \in C_0^{\infty}(\mathbf{R}^n)$. Then the convolution

$$u(x) := (\Phi * f)(x) = \int_{\mathbf{R}^n} \Phi(x - y) f(y) \, dy$$

is a solution to the Poisson equation $-\Delta u = f$ in \mathbf{R}^n . Show that if f is radial, i.e., f(y) = f(|y|), and supported in $B_R := \{ |x| < R \}$, then

$$u(x) = c\Phi(x)$$

for any $x \in \mathbf{R}^n \setminus B_R$, where

$$c = \int_{\mathbf{R}^n} f(y) \, dy.$$

Hint: Use polar (spherical) coordinates and apply the mean value property for harmonic functions.

3.2 Final Practice Problems

PROBLEM 3.2.1. Let Ω be a bounded domain in \mathbf{R}^n with a smooth boundary. Show that the problem

$$\begin{cases}
-\Delta u = f & \text{in } \Omega, \\
u + \alpha \frac{\partial u}{\partial \nu} = g & \text{on } \partial \Omega
\end{cases}$$

has at most one solution in $C^2(\Omega) \cap C(\bar{\Omega})$ if $\alpha > 0$. Here ν is the outward normal on $\partial \Omega$ and f, g are assumed to be smooth.

SOLUTION. We will use energy methods to show that there is only one solution to the problem

$$\begin{cases}
-\Delta u = f & \text{in } \Omega, \\
u + \alpha \frac{\partial u}{\partial \nu} = g & \text{on } \partial \Omega
\end{cases}$$
(3.2.1)

Suppose u_1 and u_2 are two distinct solutions to the problem (3.2.1). Define $v := u_1 - u_2$. Then v solves the problem

$$\begin{cases} -\Delta v = 0 & \text{in } \Omega', \\ v = -\alpha \frac{\partial v}{\partial \nu} & \text{on } \partial \Omega'. \end{cases}$$

Consider the energy

$$E[v] = \frac{1}{2} \int_{\Omega'} |Dv|^2 dx.$$

PROBLEM 3.2.2. Let g be a continuous function with compact support in \mathbb{R}^n . Write the formula for the bounded solution of

$$\begin{cases} u_t - \Delta u = 0 & \text{for } x \in \mathbf{R}^n, \ t > 0, \\ u(x,0) = g(x) & \text{for } x \in \mathbf{R}^n. \end{cases}$$

Prove that $\lim_{t\to\infty} u(x,t) = 0$, where the convergence is uniform in $x \in \mathbf{R}^n$.

PROBLEM 3.2.3. Find an explicit solution to the problem

$$\begin{cases} u_t - u_{xx} = 0 & \text{for } x \in \mathbf{R}^n, \ t > 0, \\ u(x, 0) = e^{3x} & \text{for } x \in \mathbf{R}. \end{cases}$$

PROBLEM 3.2.4. Find a formula for the solution of

$$\begin{cases} u_{tt} - u_{xx} + u = 0 & \text{in } \mathbf{R} \times (0, \infty), \\ u(x, 0) = f(x), & u_t(x, 0) = g(x), \end{cases}$$

where $f, g \in C_0^{\infty}(\mathbf{R})$.

(*Hint:* Method I: Use Hadamard's method of descent. Namely, find h(y) such that v(x, y, t) := h(y)u(x, t) solves

$$v_{tt} - (v_{xx} + v_{yy}) = 0.$$

Method II: Use the Fourier transform.)

Solution.

PROBLEM 3.2.5. Let $u \in C^2(\mathbf{R}^n \times [0, \infty))$ satisfy

$$\begin{cases} u_{tt} - \Delta u = 0 & \text{in } \mathbf{R}^n \times (0, \infty), \\ u(x, 0) = g(x), & u_t(x, 0) = h(x). \end{cases}$$

Show that if both g and h are radial, then so is $u(\cdot,t)$ for any t>0. (Recall that the function f is called radial if f(x)=f(|x|).)

SOLUTION.

PROBLEM 3.2.6. Find the value of the solution u of the initial value problem

$$\begin{cases} u_{tt} - \Delta u = 0 & \text{for } x \in \mathbf{R}^3, \ t > 0, \\ u(x,0) = 0, u_t(x,0) = \varphi(x), \end{cases}$$

where

$$\varphi(x) := \begin{cases} 1 & \text{for } |x| < a, \\ 0 & \text{for } |x| \ge a \end{cases}$$

at a point (x, t) such that |x| + t < a.

Solution.

PROBLEM 3.2.7. LEt u be a nonzero harmonic function in $B(0,R) := \{ x \in \mathbb{R}^n : |x| < R \}$. Define

$$E(r) := \int_{\partial B(0,r)} u^2(y) \, d\sigma_y.$$

Show that $\ln E(r)$ is a convex function of $\ln r$; i.e.,

$$E(\sqrt{ab})^2 \le E(a)E(b)$$
, for $a, b > 0$,

for any $0 < a \le c < R$.

(*Hint*: Use the representation of u as a uniformly convergent series

$$u(x) = \sum_{k=0}^{\infty} p_k(x), \qquad |x| < R,$$

where $p_k(x)$ is a homogeneous harmonic polynomial of order k.)

SOLUTION.

PROBLEM 3.2.8. Use Kirchhoff's formula and Duhamel's principle to obtain an integral representation of the solution to the following Cauchy problem,

$$\begin{cases} u_{tt} - \Delta u = e^{-t} f(x) & x \in \mathbf{R}^3, \ t > 0, \\ u(x,0) = u_t(x,0) = 0, & x \in \mathbf{R}^3. \end{cases}$$

Verify that the integral representation reduces to the obvious solution $u = e^{-t} + t - 1$ when f(x) = 1.

Solution.

PROBLEM 3.2.9. Let $f(x) = e^{-|x|^2}$, $x \in \mathbb{R}^n$. Find f * f. (*Hint:* Use either the heat equation or the Fourier transform.)

SOLUTION.

PROBLEM 3.2.10. Recall that a solution to the heat equation

$$\begin{cases} u_t - \Delta u = 0 & \text{in } \mathbf{R}^n \times (0, \infty), \\ u = g & \text{on } \mathbf{R}^n \times \{t = 0\} \end{cases}$$

is given by

$$u(x,t) = \int_{\mathbf{R}^n} \Phi(x - y, t) g(y) dt,$$

where, for t > 0,

$$\Phi(z,t) = (4\pi t)^{-\frac{n}{2}} e^{-\frac{|x|^2}{4t}}.$$

Assume that g is continuous and compactly supported. Show that there exists a C > 0 such that

$$|Du(x,t)| \le Ct^{-\frac{1}{2}} ||g||_{L^{\infty}}.$$

4 Qualifying Exams

4.1 Qualifying Exam, August '04

PROBLEM 4.1.1. Consider the initial value problem

$$\begin{cases} a(x,y)u_x + b(x,y)u_y = -u, \\ u = f & \text{on } S^1 = \{x^2 + y^2 = 1\}, \end{cases}$$

where a and b satisfy

$$a(x,y) + b(x,y)y > 0$$

for any $x, y \in \mathbf{R}^n \setminus \{(0,0)\}.$

- (a) Show that the initial value problem has a unique solution in a neighborhood of S^1 . Assume that a, b, and f are smooth.
- (b) Show that the solution of the initial value problem actually exists in $\mathbb{R}^2 \setminus \{(0,0)\}$.

SOLUTION.

PROBLEM 4.1.2. Let $u \in C^2(\mathbf{R} \times [0, \infty))$ be a solution of the initial value problem for the onedimensional wave equation

$$\begin{cases} u_{tt} - u_{xx} = 0 & \text{on } \mathbf{R} \times (0, \infty), \\ u = f, & u_t = g & \text{in } \mathbf{R} \times 0, \end{cases}$$

where f and g have compact support. Define the kinetic energy by

$$K(t) = \frac{1}{2} \int_{-\infty}^{\infty} u_t^2(x, t) dx,$$

and the potential energy by

$$P(t) = \frac{1}{2} \int_{-\infty}^{\infty} u_x^2(x,t) \, dx.$$

Show that

- (a) K(t) + P(t) is constant in t,
- (b) K(t) = P(t) for all large enough times t.

SOLUTION.

PROBLEM 4.1.3. Use Kirchhoff's formula and Duhamel's principle to obtain an integral representation of the solution of the following Cauchy problem

$$\begin{cases} u_{tt} - \Delta u = e^{-t} g(x) & \text{for } x \in \mathbf{R}^3, \ t > 0, \\ u(x,0) = u_t(x,0) = 0 & \text{for } x \in \mathbf{R}^3. \end{cases}$$

Verify that the integral representation reduces to the obvious solution $u = e^{-t} + t - 1$ when g(x) = 1.

SOLUTION.

PROBLEM 4.1.4. Let Ω be a bounded open set in \mathbf{R}^n and $g \in C_0^{\infty}(\Omega)$. Consider the solutions of the initial boundary value problem

$$\begin{cases} \Delta u - u_t = 0 & \text{for } x \in \Omega, \ t > 0, \\ u(x, 0) = g(x) & \text{for } x \in \Omega, \\ u(x, t) = 0 & \text{for } xi \in \partial\Omega, \ t \geq 0, \end{cases}$$

and the Cauchy problem

$$\begin{cases} \Delta v - v_t = 0 & \text{for } x \in \mathbf{R}^n, \ t > 0, \\ v(x, 0) = |g(x)| & \text{for } x \in \mathbf{R}^n, \end{cases}$$

where we put g = 0 outside Ω .

(a) Show that

$$-v(x,t) \le u(x,t) \le v(x,t)$$

for any $x \in \Omega$, t > 0.

(b) Use (a) to conclude that

$$\lim_{t \to \infty} u(x, t) = 0,$$

for any $x \in \Omega$.

SOLUTION.

PROBLEM 4.1.5. Let $P_k(x)$ and $P_m(x)$ be homogeneous harmonic polynomials in \mathbb{R}^n of degrees k and m respectively; i.e.,

$$P_k(\lambda x) = \lambda^k P_k(x), \qquad P_m(\lambda x) = \lambda^m P_m(x),$$

for any $x \in \mathbf{R}^n$, $\lambda > 0$,

$$\Delta P_k = 0, \qquad \Delta P_m = 0$$

in \mathbf{R}^n .

(a) Show that

$$\frac{\partial P_k(x)}{\partial \nu} = kP_k(x), \qquad \frac{\partial P_m(x)}{\partial \nu} = mP_m(x)$$

on ∂B_1 , where $B_1 = \{ |x| < 1 \}$ and ν is the outward normal on ∂B_1 .

(b) Use (a) and Green's second identity to prove that

$$\int_{\partial B_1} P_k(x) P_m(x) \, dS = 0,$$

if $k \neq m$.

4.2 Qualifying Exam, August '05

Problem 4.2.1.

(a) Find a solution of the Cauchy problem

$$\begin{cases} yu_x + xu_y = xy, \\ u = 1 & \text{on } S^1 = \{ x^2 + y^2 = 1 \}. \end{cases}$$

(b) Is the solution unique in a neighborhood of the point (1,0)? Justify your answer.

SOLUTION. The solution to teh first part is

$$u(x,y) = \frac{x^2 + y^2}{4} + \frac{3}{4}.$$

PROBLEM 4.2.2. Consider the second order PDE in $\{x > 0, y > 0\} \subseteq \mathbb{R}^2$

$$x^2 u_{xx} - y^2 u_{yy} = 0.$$

- (a) Classify the equation and reduce it to the canonical form.
- (b) Show that the general solution of the equation is given by the formula

$$u(x,y) = F(x,y) + \sqrt{xy}G(x/y).$$

SOLUTION.

PROBLEM 4.2.3. Let Φ be the fundamental solution of the Laplace equation in \mathbf{R}^3 and $f \in C_0^{\infty}(\mathbf{R}^n)$. Then the convolution

$$u(x) := (\Phi * f)(x) = \int_{\mathbf{R}^n} \Phi(x - y) f(y) \, dy$$

is a solution of the Poisson equation $-\Delta u = f$ in \mathbf{R}^n . Show that if f is radial (i.e., f(y) = f(|y|)) and supported in $B_R = \{ |x| < R \}$, then

$$u(x) = c\Phi(x),$$

for any $x \in \mathbf{R}^n \setminus B_R$, where

$$c = \int_{\mathbf{R}^n} f(y) \, dy.$$

[Hint: Use spherical (polar) coordinates and the mean value property.]

PROBLEM 4.2.4. Consider the so-called 2-dimensional wave equation with dissipation

$$\begin{cases} u_{tt} - \Delta u + \alpha u_t = 0 & \text{in } \mathbf{R}^2 \times (0, \infty), \\ u(x, 0) = g(x), & u_t(x, 0) = h(x) & \text{for } x \in \mathbf{R}^2, \end{cases}$$

where $g, h \in C_0^{\infty}(\mathbf{R}^2)$ and $\alpha \geq 0$ is a constant.

(a) Show that for an appropriate choice of constant λ and μ the function

$$v(x_1, x_2, x_3, t) := e^{\lambda t + \mu x_3} u(x_1, x_2, t)$$

solves the 3-dimensional wave equation $v_{tt} - \Delta v = 0$.

(b) Use (a) to prove the following domain of dependence result: for any point $(x_0, t_0) \in \mathbf{R}^2 \times (0, \infty)$ the value $u(x_0, t_0)$ is uniquely determined by values of g and h in $\overline{B_{t_0}(x_0)} := \{ |x - x_0| \le t_0 \}$. (You may use the corresponding result for the wave equation without proof.)

SOLUTION.

PROBLEM 4.2.5. Let u(x,t) be a bounded solution of the heat equation $u_t = u_{xx}$ in $\mathbf{R} \times (0,\infty)$ with the initial condition

$$u(x,0) = u_0(x)$$

for $x \in \mathbf{R}$, where $u_0 \in C^{\infty}$ is 2π -periodic, i.e., $u_0(x+2\pi) = u_0(x)$. Show that

$$\lim_{t \to \infty} u(x, t) = a_0,$$

uniformly in $x \in \mathbf{R}$, where

$$a_0 := \frac{1}{2\pi} \int_0^{2\pi} u_0(x) \, dx.$$

4.3 Qualifying Exam, January '14

PROBLEM 4.3.1. Consider the first order equation in \mathbb{R}^2

$$x_2 u_{x_1} + x_1 u_{x_2} = 0.$$

- (a) Find the characteristic curves of the equation.
- (b) Consider the Cauchy problem for this equation prescribed on the line $x_1 = 1$:

$$u(1, x_2) = f(x_2).$$

Find a necessary condition on f so that the proble is solvable in a neighborhood of the point (1,0).

SOLUTION.

PROBLEM 4.3.2. Let u be a continuous bounded solution of the initial value problem for the Laplace equation

$$\begin{cases} \Delta u = 0 & \text{in } \{x_n > 0\}, \\ u(x', 0) = g(x') & \text{for } x' \in \mathbf{R}^{n-1}, \end{cases}$$

where g is a continuous function with compact support in \mathbb{R}^{n-1} . Here $n \geq 2$. Prove that

$$u(x) \longrightarrow 0,$$
 as $|x| \longrightarrow \infty$,

for $x \in \{x_n > 0\}$.

SOLUTION.

PROBLEM 4.3.3. Let u be a bounded solution of the heat equation

$$\Delta u - u_t = 0$$
 in $\mathbf{R} \times (0, \infty)$,

with the initial conditions u(x,0) = g(x), where g is a bounded continuous function on **R** satisfying the Hölder condition

$$|g(x) - g(y)| \le M|x - y|^{\alpha}, \quad x, y \in \mathbf{R}$$

with a constant $\alpha \in (0,1]$. Show that

$$|u(x,t) - u(y,t)| \le M|x - y|^{\alpha}, \quad x, y \in \mathbf{R}, t > 0,$$

 $|u(x,t) - u(x,s)| \le C_{\alpha}M|t - s|^{\alpha/2}, \quad x \in \mathbf{R}, t, s > 0.$

[Hint: For the last inequality, in the representation formula of u(x,t) as a convolution with the heat kernel $\Phi(y,t)$, make a change of variables $z=y/\sqrt{t}$ and use that $\left|\sqrt{t}-\sqrt{s}\right|\leq\sqrt{|t-s|}$.]

PROBLEM 4.3.4. Let u be a positive harmonic function in the unit ball B_1 in \mathbb{R}^n . Show that

$$|D(\ln u)| \le M \qquad \text{in } B_{1/2}$$

for a constant M depending only on the dimension n.

[Hint: Use the interior derivative estimate $|Du(x)| \leq (C_n/r) \sup_{B_r(x)} |u|$ for $B_r(x) \subseteq B_1$ as well as the Harnack inequality for harmonic functions.]

SOLUTION.

PROBLEM 4.3.5. Let u be a C^2 solution of the initial value problem

$$\begin{cases} u_{tt} - \Delta u = |x|^k & \text{in } \mathbf{R}^n \times (0, \infty), \\ u = 0, & u_t = 0 & \text{on } \mathbf{R}^n \times \{0\}. \end{cases}$$

for some $k \geq 0$. Prove that there exists a function $\varphi(r)$ such that

$$u(x,t) = t^{k+2}\varphi(|x|/t).$$

[Hint: As one of the steps show that u is (k+2)-homogeneous in (x,t) variables, i.e., $u(\lambda x, \lambda t) = \lambda^{k+2} u(x,t)$ for any $\lambda > 0$.]