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**by**

**Andrew William Ritchie, B.S. Chem**

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The University of Texas at Austin, 2015

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<Abstract: May not exceed 350 words. It should be a continuous description, not disconnected notes or an outline.>

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## **Chapter 1 Introduction**

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## Chapter 2 Simulation Methods

### 2.1 LABELING AND MUTATING PROTEINS *IN SILICO*

Simulation parameters for cyanocysteine and GDPNP were obtained from previous studies<sup>1,2</sup>. Simulations were performed using the Amber03 force field in GOMACS.<sup>3</sup> Ral starting structures came from the 2RGF<sup>4</sup> crystal structure, Ras starting structures came from 1LFD<sup>5</sup>, and Rap starting structures came from 1GUA<sup>6</sup>.

Generation of Ral $\beta$  starting structures has previously been reported.<sup>2</sup> In short, all cysteine sidechain atoms except for C $\beta$  were deleted from 2RGF and the sidechain was renamed to alanine. The C-terminal residues RTFT were taken from the pdb structure 1RAX (deposited in the PDB but unpublished) which ends in KKRTFT, and pasted onto the 2RGF structure, which ends in KKRT, by aligning the backbone atoms of the common KKRT residues and adding the FT coordinates onto the Ral structure file. The first residues in the 2RGF structure are ALA and LEU; these were changed to SER and HIS by renaming the residues in the pdb file, deleting hydrogen atoms, and renaming, using CD1 of LEU as ND1 of HIS. The N-terminal glycine (GLY 4) was modeled using Avogadro, completing the sequence. All missing heavy atoms were added using the tleap utility of Amber Tools.<sup>7</sup> The end resultant sequence is, starting from the N-terminal, GSH+Ral(2RGF,Cys $\rightarrow$ Ala)+FT.

To make Ras, the GAMGS sequence from chain B of 4K81<sup>8</sup> was used due to previous work which showed it to be the lowest energy conformation of this sequence available in the Protein Data Bank when bound to Ras.<sup>9</sup> The N-terminal methionine of 1LFD was aligned to the methionine immediately following the GAMGS sequence in

4K81, after which the GAMGS coordinates were added to the Ras structure file. The end resultant sequence is GAMGS+Ras(1LFD).

Rap starting structures and mutations have also already been reported.<sup>10</sup> To start, a GSH tag left on the N-terminal methionine after cleavage of the hexa-histidine affinity label during protein purification was modeled onto the N-terminus. This was done by searching the Protein Data Bank for proteins starting with the sequence GSHM (Met is the first Rap residue present in 1GUA). 87 NMR structures and 1 crystal structure were obtained from pdb 1AQ5 (20 NMR structures)<sup>11</sup>, 1W9R (19 NMR structures)<sup>12</sup>, 2WCY (48 NMR structures)<sup>13</sup>, and 2VKJ (1 crystal structure)<sup>14</sup>. Rap was aligned to the methionine backbone of each GSHM structure using VMD<sup>15</sup>, creating 88 structures containing the GSHM N-terminus. After adding hydrogen atoms using the GROMACS utility pdb2gmx<sup>16</sup>, an energy minimization was performed and the lowest energy structure from this collection was chosen as the Rap model to be used for further calculations. The end resultant sequence is GSHM+Rap(1GUA). It was observed that in some structures, the N-terminal residues of Ral $\beta$  protruded between bonded atoms of Rap. To eliminate this nonphysical steric overlap, heavy atom restraints were placed on all but the Ral $\beta$  N-terminal GSHM residues and a 500 step gentle minimization was performed in Gromacs<sup>16</sup>.

All side chain mutations were generated in the same manner using Amber Tools.<sup>7</sup> All side chain atoms except for shared heavy atoms were deleted from the mutation residue. The wild type residue was renamed to the desired residue and the resulting structure was passed to the tleap utility in Amber Tools<sup>7</sup> to model back in the missing atoms. In this way, starting from 1LFD Ras D30/E31K and 1GUA Rap E30D/K31E the additional Ras constructs D30/E31, D30E/E31, and D30E/E31K, and the additional Rap constructs E30/K31, E30D/K31, and E30/K31E were each constructed.

To generate cyanylated Ral structures N27C<sub>SCN</sub> , G28C<sub>SCN</sub> , N29C<sub>SCN</sub> , Y31C<sub>SCN</sub> , K32C<sub>SCN</sub> , and N54C<sub>SCN</sub> , residues were renamed to MET in the pdb file (because methionine contains the same number of heavy atoms as cyanocysteine), retaining and renaming any atoms common to the native residue and methionine. Missing methionine atoms were added using tleap.<sup>7</sup> To complete the mutation to cyanocysteine, the MET atoms CG, SD, and CE were renamed SG, CD, and NE and a short energy minimization was performed.

To dock Rap to Ral, the alpha carbons of the GTPase units of 1LFD and 1GUA were first aligned using the Smith-Waterman algorithm<sup>17</sup> with a gap penalty of -3 (chosen because it gave results most consistent with the STAMP<sup>18</sup> structural alignment in VMD<sup>15</sup>), and the Ral coordinates were merged with the Rap coordinates of 1GUA and saved as a reference structure, Rap(1GUA)+Ral(1LFD). The mutated Rap structure was then aligned to the Rap unit of the Rap(1GUA)+Ral(1LFD) reference structure and the coordinates of the resulting Rap(mutant)+Ral(1LFD) were saved. To dock mutated Ras structures to Ral, the mutated Ras was aligned to the 1LFD Ras and the Ral coordinates were merged with the mutated Ras structure and the coordinates of the resulting Rap(mutant)+Ral(1LFD) were saved. To introduce the probe to the docked system, the cyanylated Ral was aligned to the Ral of each GTPase(mutant)+Ral(1LFD) reference structure, and the GTPase(mutant) coordinates were merged with the cyanylated Ral coordinates to create each GTPase+Probe construct: Ral+N27C<sub>SCN</sub> , Ral+G28C<sub>SCN</sub> , Ral+N29C<sub>SCN</sub> , Ral+Y31C<sub>SCN</sub> , Ral+K32C<sub>SCN</sub> , Ral+N54C<sub>SCN</sub> , Rap E30/K31+N27C<sub>SCN</sub> , Rap E30/K31+G28C<sub>SCN</sub> , Rap E30/K31+N29C<sub>SCN</sub> , Rap E30/K31+Y31C<sub>SCN</sub> , Rap E30/K31+K32C<sub>SCN</sub> , Rap E30/K31+N54C<sub>SCN</sub> , Rap E30/K31E+N27C<sub>SCN</sub> , Rap E30/K31E+G28C<sub>SCN</sub> , Rap E30/K31E+N29C<sub>SCN</sub> , Rap E30/K31E+Y31C<sub>SCN</sub> , Rap E30/K31E+K32C<sub>SCN</sub> , Rap E30/K31E+N54C<sub>SCN</sub> , Rap E30D/K31+N27C<sub>SCN</sub> , Rap

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 D30E/E31K+K32C<sub>SCN</sub> , and Ras D30E/E31K+N54C<sub>SCN</sub> .

## 2.2 ENHANCED MOLECULAR DYNAMICS IN AMBER03: N-DIMENSIONAL UMBRELLA SAMPLING AND WEIGHTED HISTOGRAM ANALYSIS METHOD

An umbrella sampling strategy was used to obtain a Boltzmann-weighted statistical ensemble of thiocyanate probe orientations for all MD sampling. Through examining simulated protein structures, it became increasingly apparent that a second degree of freedom,  $\chi_1$ , was relevant to our probe conformational distributions. Therefore, two different umbrella sampling strategies were tested: one-dimensional sampling about the thiocyanate  $\chi_2$  dihedral angle and two-dimensional umbrella sampling about the thiocyanate  $\chi_1$  and  $\chi_2$  dihedral angles, shown in Figure 2-1. All molecular dynamics were completed using the GROMACS<sup>16</sup> software package at 300 K with the AMBER03<sup>19</sup> force field and periodic boundary conditions.

Six probe locations on RalGDS were examined: N27C<sub>SCN</sub>, G28C<sub>SCN</sub>, N29C<sub>SCN</sub>, Y31C<sub>SCN</sub>, K32C<sub>SCN</sub>, and N54C<sub>SCN</sub>, in the monomeric state and docked to each GTPase system examined. We have therefore examined all probe locations and mutated constructs for which experimental data are available. Six structures for each system modeled were generated by fixing the thiocyanate  $\chi_2$  dihedral angle from 0° to 300° in 60° increments. Each structure was sampled with a dihedral potential that was flat within  $\pm 30^\circ$  of the fixed-dihedral position and quadratic with a force constant of 1000 kJ mol<sup>-1</sup> rad<sup>-2</sup> otherwise. These restraining potentials were carried through for the duration of the system set-up and simulation. Each structure was energy minimized with cut-off electrostatics, solvated with tip3p water<sup>11</sup> in a dodecahedron box, charge balanced by randomly replacing the appropriate number of water molecules with sodium or chloride ions using the genion GROMACS utility, and solvent relaxed by sampling for 20 ps with position restraints on all non-solvent heavy atoms with a force constant of 1000 kJ mol<sup>-1</sup>



nm<sup>-2</sup> using PME<sup>20,21</sup> electrostatics with a real-space cut-off of 0.9 nm, spacing of 0.12 nm, and interpolation order 4. Each rotamer of each system was then sampled using the GROMACS stochastic dynamics integrator, constraints on hydrogen-bonds using the LINCS algorithm<sup>22</sup>, and PME electrostatics for 3 ns, recording snapshots every 5 ps, for a total of 18 ns of simulation and 3606 frames for each system. To test for structural rearrangements at positions 30 and 31 of Rap that would significantly influence the measured electrostatic field, the  $\chi_1$  dihedral angle on Rap residues 30 and 31 were each individually biased to create a Boltzmann-weighted ensemble of  $\chi_1$  rotomers. Rotamer libraries of Rap E30/K31, Rap E30D/K31, Rap E30/K31E, and the double mutant Rap E30D/K31E were created using the strategy described above. Once again, structures containing dihedral angles from 0° to 300° in 60° degree increments were generated for each mutant; these were energy minimized, solvated in tip3p water, relaxed, and then sampled for 3 ns each, generating 18 ns of simulation for each biased torsional angle. To summarize, for 24 mutant-containing rotamer structures, MD sampling included 18 ns of simulation biasing the nitrile probe, 18 ns of simulation biasing of Rap position 30, and 18 ns of simulation biasing of Rap position 31, for a total of 1296 ns of MD simulation of this system

Each frame was assigned to one of 72 5° bins from -180° to 175° based on the  $\chi_2$  dihedral angle. The weighted histogram analysis method (WHAM)<sup>23,24</sup> was then used to calculate a torsional potential of mean force (PMF) for each of  $N$  bins  $i$ , which is related to the torsional probability distribution for each bin  $i$  ( $P_i$ ) described by equation (2-1):

$$P_i = \frac{e^{-\beta \cdot PMF_i}}{\sum_{j=0}^N e^{-\beta \cdot PMF_j}}, \beta = \frac{1}{k_b T} \quad (2-1)$$

which is the typical Boltzmann distribution function for a state  $i$  divided by the partition function, where  $T$  is the temperature in Kelvin,  $k_b$  is the Boltzmann constant, and  $PMF_j$  is the PMF for some state  $j$ . In this way, a Boltzmann-weighted distribution of structures was assembled in order to determine the average orientation of both the thiocyanate residues and the side chains at Rap positions 30 and 31 for comparing to experimental data.

The two-dimensional umbrella sampling was done in much the same way as the one-dimensional sampling, with few minor changes. The  $\chi_1$  angle was fixed from  $0^\circ$  to  $300^\circ$  in  $30^\circ$  increments, resulting in 12 structures. Each of these structures then had the  $\chi_2$  angle fixed from  $0^\circ$  to  $300^\circ$  in  $30^\circ$  increments, resulting in 144 total structures. To avoid steric clashes in the starting structures, for each  $\chi_2$  rotation, the distance between the center of mass coordinate of each rotated atom and every non-rotated atom was calculated. If a distance was found to be less than  $1.5 \text{ \AA}$ , the  $\chi_2$  angle was rotated  $\pm 1.5^\circ$  from the dihedral center and the distances were recalculated. This was done until all non-bonded atoms were at least  $1.5 \text{ \AA}$  from each rotated atom. Next, two harmonic dihedral restraining potentials were generated for each structure, one for the  $\chi_1$  dihedral angle and one for the  $\chi_2$  dihedral angle. Following the same set-up strategy used in the one-dimensional sampling, the system was then energy minimized with cut-off electrostatics and dihedral force constants of  $1000 \text{ kJ mol}^{-1} \text{ rad}^{-2}$ , solvated in tip3p water in a dodecahedron box, and charge balanced as described above. The system underwent solvent relaxation using PME electrostatics for 20 ps with a force constant of  $1000 \text{ kJ mol}^{-1} \text{ nm}^{-2}$  on heavy backbone atoms, dihedral force constants of  $150 \text{ kJ mol}^{-1} \text{ rad}^{-2}$ , and unrestrained side-chain atoms. Each of the 144  $\chi_1$ - $\chi_2$  rotamers were then sampled for 400 ps using the GROMACS stochastic dynamics integrator with PME electrostatics (again, with a real-space cut-off of 0.9 nm, spacing of 0.12 nm, and interpolation order 4) and

dihedral restraining potentials of  $70 \text{ kJ mol}^{-1} \text{ rad}^{-2}$ , for a total of 57.6 ns of simulation and 11664 frames for each system. It is worth reiterating that each step used a progressively smaller dihedral restraining potential. This was done to ensure that the dihedral angles of the final structure before sampling were as close as possible to the umbrella-sampling window while still allowing nearby residues to relax to orientations that accommodate the inclusion of our probe. Starting with a large restraining potential fixes the probe to a specific location orientation and forcibly moves nearby residues to accommodate the probe to minimize interaction energies. Subsequent weakening of the restraining potential allows the probe to respond to its surroundings in a manner more typical of MD, allowing both the probe as well as the residues near the probe to relax to energy minimized orientations. Without this subsequent weakening, many simulations resulted in dihedral forces becoming larger than the integrator can or is expected to handle, which were usually caused by steric clashes between the probe and side-chain atoms.

A Boltzmann-weighted statistical ensemble of structures obtained from two-dimensional sampling was then assembled using WHAM. Each frame was assigned to one of 5184  $5^\circ$  by  $5^\circ$  bins. The bins were assigned based on equation (2-2), where  $b_1$  is the one-dimensional  $\chi_1$  bin number,  $b_2$  is the one-dimensional  $\chi_2$  bin number,  $B_i$  is the total number of bins in degree of freedom  $i$  (72 for all two-dimensional sampling done), and  $i_{index}$  refers to the subscript on  $b$ . Conditional probability was assumed.

$$\text{Bin}(b_1, b_2) = \sum_{i=b_1, b_2} B_i^{2-i_{index}} \cdot i \quad (2-2)$$

To validate our 2D WHAM code, we examined a Ryckaert-Bellemans dihedral potential, shown in equation (2-3), where  $\phi$  is some angle and  $n$  and  $C_n$  are some example parameters obtained from the GROMACS manual, shown in Table 2-1. This is a very simple potential function with no contributions from any other source, unlike the

potential energy calculation in a protein, which will be influenced by various force field parameters such as bond force constants. However in WHAM, the sources of the potentials are not distinguished, and we are able to use the simple Ryckaert-Bellemans model to validate the code. We constructed the PMF from the sum of the potentials for a given pair of coordinates and the unbiased probability distribution,  $p^\circ(\chi_1, \chi_2)$ , using the ratio of the Boltzmann distribution function to the partition function, shown in equation (2-4), over an array of 2D dihedral angles ranging from  $-180^\circ$  to  $180^\circ$ .

$$V_{rb} = \sum_{n=0}^5 C_n (\cos(\phi - \pi))^n \quad (2-3)$$

Next, we constructed the PMF and probability distribution landscapes, shown in Figure 2-1a. We then applied biasing potential windows to each dimension in a manner that mimics the methods used in GROMACS. The biased probability,  $p(\chi_1, \chi_2)$ , is given by equation (2-5). We then performed a Monte Carlo simulation centered on each of 144 biased windows, each with dimensions of  $30^\circ \times 30^\circ$ , with a probability  $p(\chi_1, \chi_2)$  of sampling a given pair of dihedral angles, which was then analyzed using our 2D WHAM code to return the unbiased PMF and  $p^\circ(\chi_1, \chi_2)$ .

$$p^\circ(\chi_1, \chi_2) = \frac{e^{-\beta(V_{rb}(\chi_1) + V_{rb}(\chi_2))}}{\iint e^{-\beta(V_{rb}(\chi_1) + V_{rb}(\chi_2))} d\chi_1 d\chi_2} \quad (2-4)$$

$$p(\chi_1, \chi_2) = \frac{e^{-\beta(V_{rb}(\chi_1) + V_1 + V_{rb}(\chi_2) + V_2)}}{\iint e^{-\beta(V_{rb}(\chi_1) + V_1 + V_{rb}(\chi_2) + V_2)} d\chi_1 d\chi_2} \quad (2-5)$$

Figure 2-1 shows the analytical PMF and probability distributions (a), WHAM PMF and probability distributions on 144 windows each containing 40 frames (b), WHAM PMF and probability distributions on 144 windows each containing 80 frames (c), WHAM PMF and probability distributions on 144 windows each containing 160 frames (d), and WHAM PMF and probability distributions on 144 windows each containing 1000 frames (e). The major features of the probability distribution become

clear after only 40 frames; after 80 frames the probability distributions look very similar to the analytical distribution, and after 160 frames very little improvement is seen. We can also see that the high probability regions, representing the staggered orientations, have PMF landscapes that look like the analytical PMF, although the gauche regions appear to be ill characterized still. After 1000 frames the probability distribution is nearly identical to the analytical distribution and the moderate ranges of the PMF (light blue, > 17.92 kJ mol<sup>-1</sup>) now quantitatively approach the predictions of the analytical expression. Exact analytical PMF matching of low probability regions is very slow, requiring  $\geq 10000$  frames. From these results we conclude that the PMF converges to the analytical expression slowly for regions of low probability and quickly for regions of high probability. Table 2-2 summarizes the umbrella sampling setups for each project, listed by chapter.

Table 2-1: Sample Parameters for Ryckaert Bellemans dihedral potential function used for validating 2D WHAM code

	<b>Constraint (kJ mol<sup>-1</sup>)</b>		<b>Constraint (kJ mol<sup>-1</sup>)</b>
<b>C<sub>0</sub></b>	9.28	<b>C<sub>3</sub></b>	-3.06
<b>C<sub>1</sub></b>	12.16	<b>C<sub>4</sub></b>	-26.24
<b>C<sub>2</sub></b>	-13.12	<b>C<sub>5</sub></b>	-31.5

Table 2-2: Umbrella Sampling Parameters, grouped by chapter

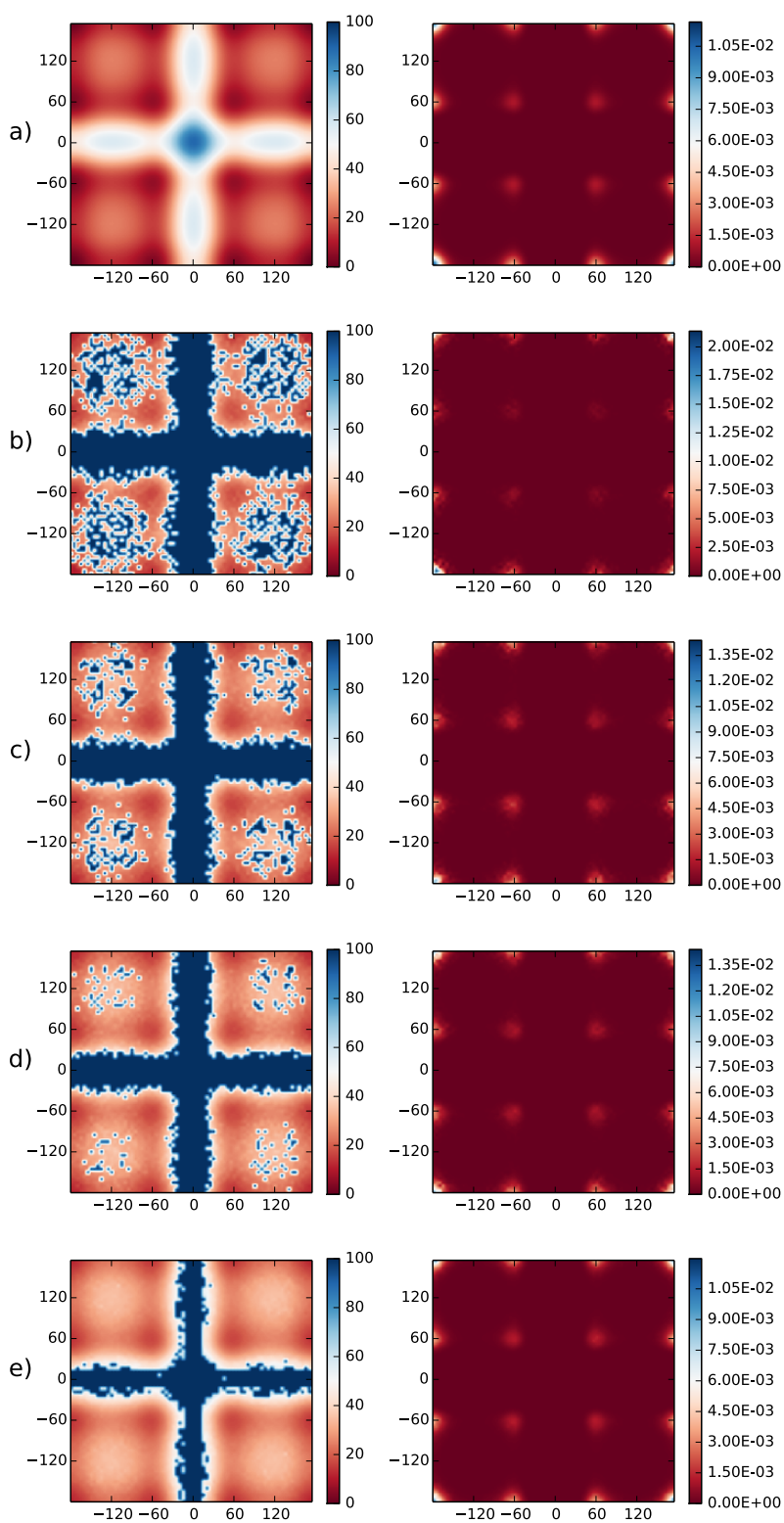
Chapters	Biasing Coordinate	Umbrella Window Size	Biasing Potential Shape	Windows per Experiment	Simulation Time per Experiment	Total Simulation Time
4, 5	CNC $\chi_2$	60°	Quadratic- Flat-	6	3.0 ns	18.0 ns
4	Position 30 $\chi_1$	60°	Quadratic- Flat-	6	3.0 ns	18.0 ns
4	Position 31 $\chi_1$	60°	Quadratic- Flat-	6	3.0 ns	18.0 ns
5, 6	CNC $\chi_1$ , $\chi_2$	30°	Quadratic	144	0.4 ns	57.6 ns
7	CNC $\chi_1$ , $\chi_2$	30 °	Quadratic	144	2.0 ns	288.0 ns
8	N/A	N/A	N/A	N/A	4-7 ns	4-7 ns





PMF

Probability Distribution



### Figure 2-1: 2D WHAM Validation

Comparison between the (left) potential of mean force (PMF) and the (right) probability distributions of a) an analytic Ryckaert-Bellemans dihedral potential and Monte Carlo 2D umbrella sampling for b) 40 frames/biasing window; c) 80 frames/biasing window; d) 160 frames/biasing window; e) 1000 frames/biasing window. Units on the PMF are  $\text{kJ mol}^{-1}$ .

## 2.3 ELECTROSTATIC CLUSTERING IN VIBRATIONAL CHROMOPHORE DIHEDRAL SPACE

The largest bottleneck for these sorts of calculations we do are the electrostatics. A single node on Stampede can generate >10 ns of simulation per day. That number can be increased (logarithmically) by using additional nodes. However, the continuum solvent electrostatics calculations take anywhere from 45-60 seconds (APBS) to  $\approx$  20 minutes (AMOEBA) per frame. If we keep every 4 ps and collect 250 frames per nanosecond, then the electrostatics require 5-8 ns/day for APBS calculations and approximately 0.3 ns/day for AMOEBA calculations. This can be decreased further by running the serial calculations in parallel. Regardless, it would be convenient to find some method of pruning the total number of frames for continuum electrostatics calculations while ensuring that the average field does not differ significantly from the average using every frame.

For convenience, the vacuum electrostatic field at the nitrile due to solute only was chosen as an indicator of total electrostatic field. In the absence of solvent, this is trivial to calculate for both point charge force fields (Amber03) as well as multipole force fields (AMOEBA). This was chosen because it was 1) intuitive and 2) there is consistently a good correlation between the solute Coulomb field and the PB solvent reaction field, as seen in Figure 2-2,git suggesting that frames which well represent the Coulomb field also well represent the reaction field.

We then took advantage of the weighted averaging over binned data. The Boltzmann weighted average is calculated as in equation (2-6),

$$\langle x \rangle = \sum_{i=1}^{n_{\text{bins}}} \rho_i \sum_{j=1}^{c_i} \frac{x_{ij}}{c_i} \quad (2-6)$$

where the probability of being in each bin  $i$  is  $\rho_i$ , the number of times bin  $i$  is visited is  $c_i$ , and each value in bin  $i$  is  $x_{ij}$  for  $j = 1$  to  $j = c_i$ . There exists some subarray of values in bin  $i$  that has  $k_i$  values, where  $k_i \geq c_i$  entries and  $\left| \sum_{j=1}^{k_i} \frac{x_{ij}}{k_i} - \sum_{j=1}^{c_i} \frac{x_{ij}}{c_i} \right| \leq \chi$ , where  $\chi$  is some threshold. As  $\chi$  approaches 0,  $k_i$  approaches  $c_i$ , and the subarray is the full array and the averages are identical. Using the a set of test data where umbrella windows were centered every 120 degrees (at 60°, 180°, and 300°, the expected alkane maximum probability torsions) with a flat biasing potential  $\pm 60^\circ$  of the window center and a force constant of **\*\*XXX**, the clustered average field for each Ral probe in the monomeric state, docked to each of the four Rap1a mutants and each of the four Ras mutants is plotted against their average field for various values of  $\chi$ , indicated in the upper-left corner of each subplot, from the full data sets in Figure 2-3. From this, it's clear that the clustered averages are linearly correlated to the full averages. The correlation coefficients and best-fit slopes have also been plotted as a function of the cutoff,  $\chi$ , in Figure 2-4. Even for a relatively large  $\chi = 1$ , the clustered correlation coefficient and slope is *very* close to 1.0. In general, as  $\chi$  approaches zero, the correlation coefficient and slope also approach zero. Moreover, at a  $\chi = 0.01$ , only approximately 20% of all frames are used, which is a significant decrease in computation requirement. This method also has the advantage of guaranteeing that the property the cluster is based on always has a clustered average nearly identical to the full average, which is a useful sanity check.

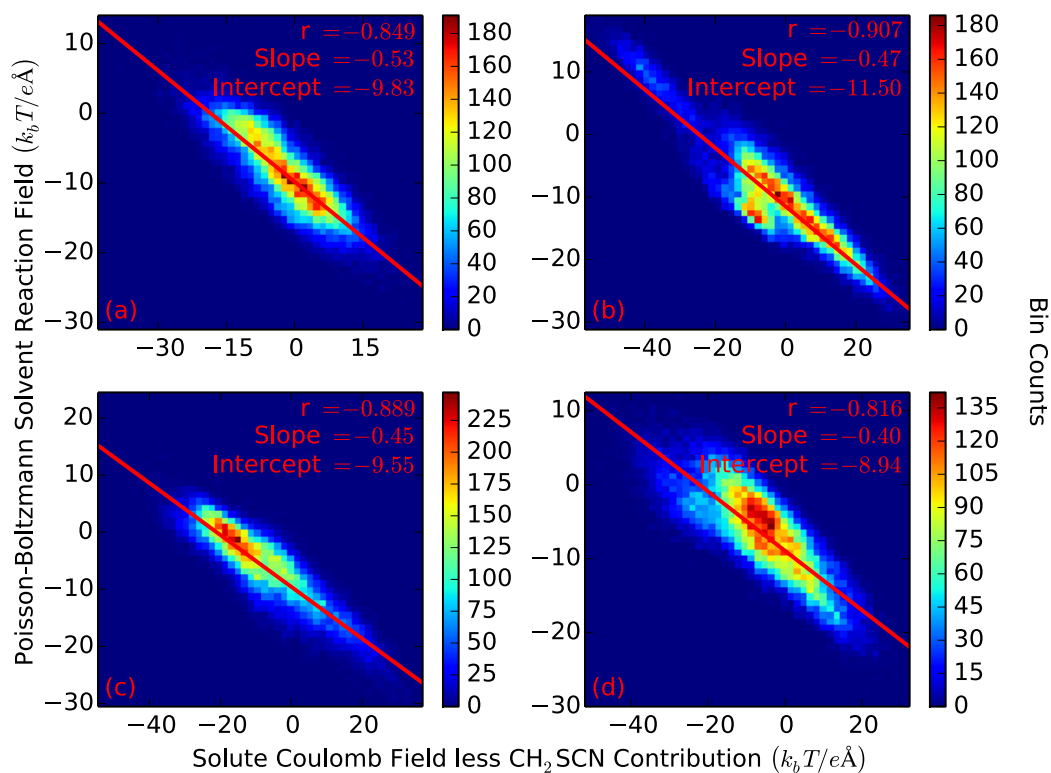


Figure 2-2: PB Solvent Reaction Field vs. Solute Analytic Coulomb Field

Comparison between the analytic Coulomb field at the nitrile bond midpoint due to solute (less the contributions due to the probe itself) (x-axis) and the solvent reaction field at the nitrile bond midpoint (y-axis) using Amber03 point charges for a) Ral G28C<sub>SCN</sub> monomer; b) Ral N54C<sub>SCN</sub> monomer; c) Ral G28C<sub>SCN</sub> docked to wild type Rap; d) Ral N54C<sub>SCN</sub> docked to wild type Rap. Correlation coefficients ( $r$ ), slopes ( $m$ ), and y-intercepts ( $\text{int}$ ), are indicated in the upper-right corner of each figure.

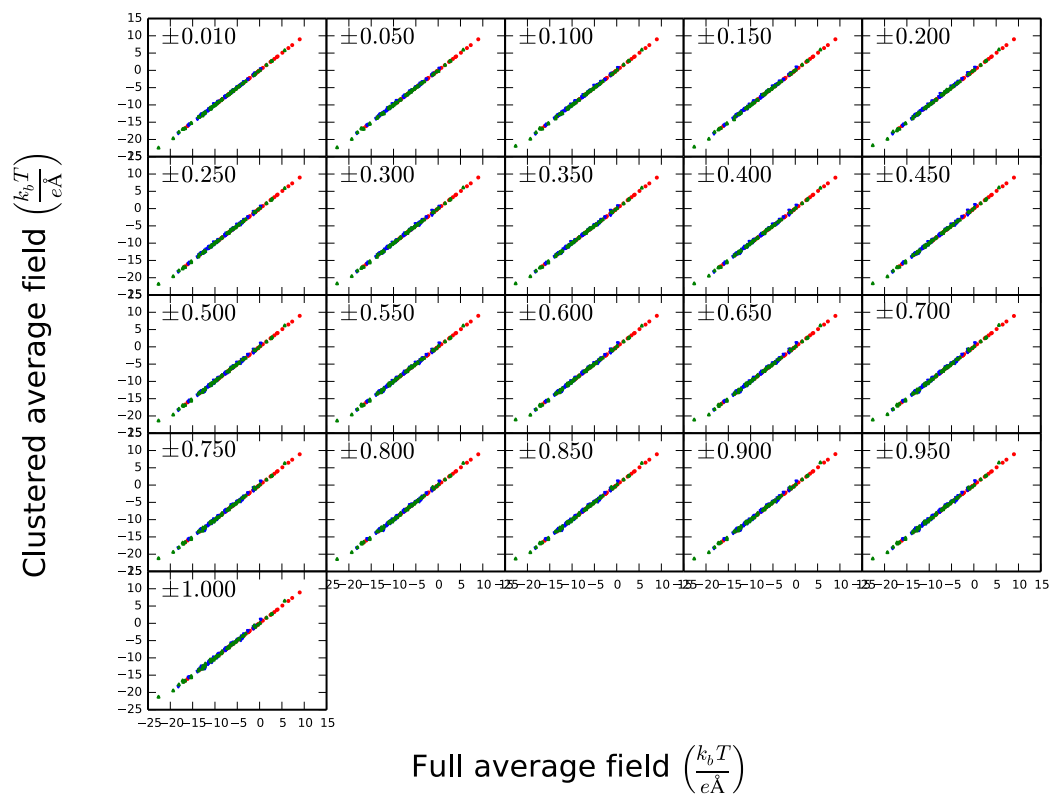


Figure 2-3: Field Values using Clustering Vs. Field Values using All Frames

Average Coulomb electrostatic field (red), solvent reaction field (blue), and the electrostatic field calculated using the AMOEBA force field (green) from clustered frames versus the respective full averages for various cutoff values,  $\chi$ , indicated in the top left corner of each box.

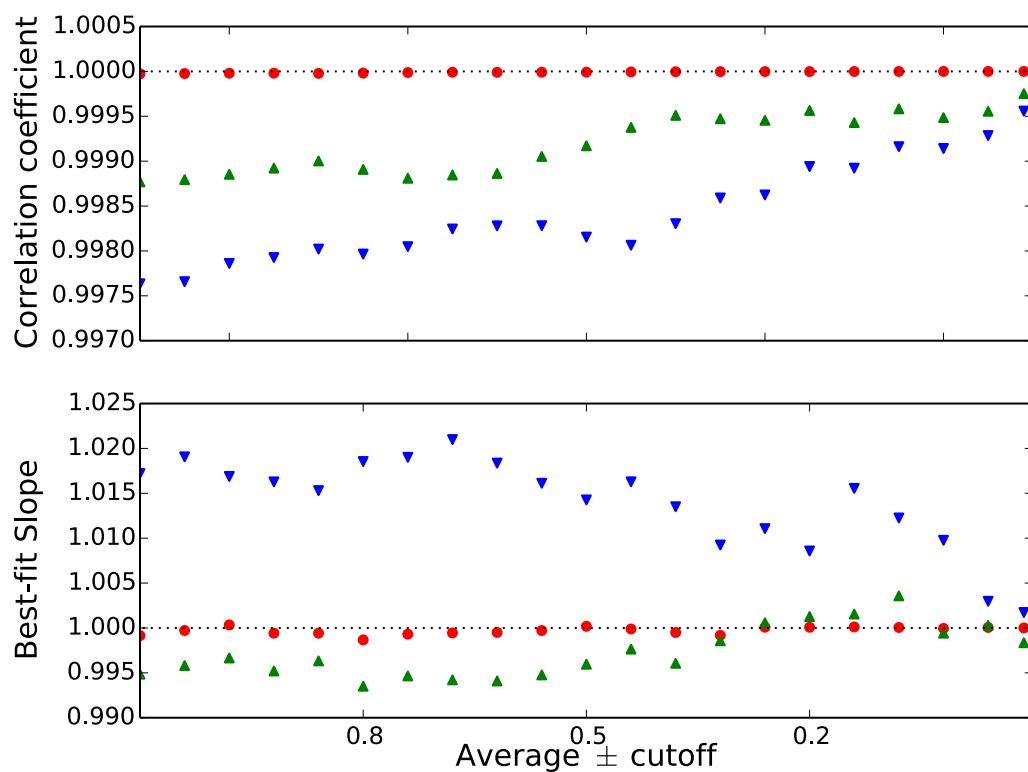


Figure 2-4: Correlations and Slopes at Various Cutoff Values

Correlation coefficients of the Coulomb electrostatic field (red), solvent reaction field (blue), and the electrostatic field calculated using the AMOEBA force field (green) as a function of the cutoff,  $\chi$ . (Bottom) Best-fit slopes of the Coulomb electrostatic field (red), solvent reaction field (blue), and the electrostatic field calculated using the AMOEBA force field (green) as a function of the cutoff,  $\chi$ .

## 2.4 SIDE CHAIN ORIENTATION: AZIMUTHAL AND POLAR ANGLES

To analyze our molecular dynamics simulations of the torsional distribution of the thiocyanate residue on each Ral $\beta$  mutant and on the side chains at Rap positions 30 and 31, we defined two angles for each side chain with respect to the Rap-Ral $\beta$  surface, which we term azimuthal and polar angles. The azimuthal angle ( $\theta$ ) was defined by measuring the angle between the side chain of interest and the average plane between the Rap and Ral $\beta$  protein surfaces. This plane was generated in the same manner as Ensign *et al.*<sup>9</sup> The positions of the C $\alpha$  atoms of the Rap surface residues Q25, D33, P34, T35, I37, E37, D38, Y40, R41, and K42 from the Rap-Ral $\beta$  starting structure were fit with a least squares regression to obtain the Rap surface plane. Similarly, the Ral $\beta$  surface plane was obtained by a least squares fit of the C $\alpha$  atoms of the Ral $\beta$  surface residues I18, R20, N27, G28, N29, M30, Y31, K32, S33, K52, and N54 in the same structure. These two planes were averaged to define the surface plane of the Rap-Ral $\beta$  interface and a normalized vector,  $\vec{N}$ , orthogonal to this plane. The vector,  $\vec{v}$ , defining the orientation of the side chain was the C $\delta$ -N $\epsilon$  bond vector for cyanocysteine, the C $\alpha$ -C $\gamma$  vector for aspartate, the C $\alpha$ -C $\delta$  vector for glutamate, and the C $\alpha$ -N $\zeta$  vector for lysine. For each MD snapshot, the heavy atoms of each frame were aligned to the starting structure, and the azimuthal angle was calculated from Equation (2-7):

$$\theta = 90^\circ - \cos^{-1}(\vec{N} \cdot \vec{v}) \quad (2-7)$$

To generate the polar angle ( $\phi$ ), a reference axis called the polar axis was first constructed from the intersection of the surface plane defined above with a second plane defined by the least squares fit of all alpha carbons in starting structure (note: this plane is not necessarily orthogonal to the surface plane). The vector orthogonal to both the surface plane as well as this vertical plane was calculated from the cross product. This



vector,  $\vec{X}$ , is the polar axis. The angle between the polar axis and the vector in plane along the x-axis was calculated from the inverse cosine of the dot product, where the sign of the angle was determined by the sign of the y coordinate (since inverse cosine only returns the absolute value of the angle from 0° to 180°). The angle between the nitrile Cδ-Ne vector and the vector along the x-axis was calculated in the same manner,

including sign determination. The polar angle ( $\phi$ ) was then determined by equation (2-8),

$$\phi = \cos^{-1}(\vec{x} \cdot \vec{B}_p) - \cos^{-1}(\vec{X} \cdot \vec{x}) - 90^\circ \quad (2-8)$$

where  $\vec{B}_p$  is the projection of the normalized bond vector onto the surface plane, and  $\vec{x}$  is a reference vector on the average surface plane along the  $x = 0$  axis. In this way, a complete 360° perspective of polar angles were determined. Figure 2-5 is a visual representation of these two angles superimposed on the structure of the Rap-Ralβ interface. Figure 2-5A defines a surface plane at the interface of Rap (above the plane) and Ral (below the plane); azimuthal angles are reported relative to this plane. When the cross hairs on Figure 2-5B are translated to the Cα atom of each simulated residue on Rap or Ralβ (represented as spheres), they become the origin of the polar angles shown in Figure 4-1 and Figure 4-2.

To obtain Boltzmann weighted angle averages, the cyclic boundaries had to be considered to avoid obtaining averages that would be affected by an arbitrary decision for the position of our torsional angle at 0°. We calculated a cyclic average angle for  $n$  equally weighted values of  $\theta$  from equation (2-9). Each frame was assigned to a bin  $j$ , with each bin  $j$  having a probability of  $p_j$ , resulting in 72 bins, each having  $m_j$  equally probable members.

$$\langle \theta \rangle = \text{atan2} \left( \frac{1}{n} \sum_{j=1}^n \sin \theta_j, \frac{1}{n} \sum_{j=1}^n \cos \theta_j \right) \quad (2-9)$$

When all probabilities  $p_j$  are equal, then dividing the sum of values of  $\sin \theta$  expressed in Equation (2-9) by the quantity in each bin ( $m_j$ ) is equivalent to the special case of summing the products of each value times the probability of that value, shown in equation (2-10):

$$\langle \theta \rangle = \text{atan2} \left[ \sum_{j=1}^{72} \left( p_j \cdot \frac{1}{m_j} \sum_{i=1}^{m_j} \sin \theta_i \right), \sum_{j=1}^{72} \left( p_j \cdot \frac{1}{m_j} \sum_{i=1}^{m_j} \cos \theta_i \right) \right] \quad (2-10)$$

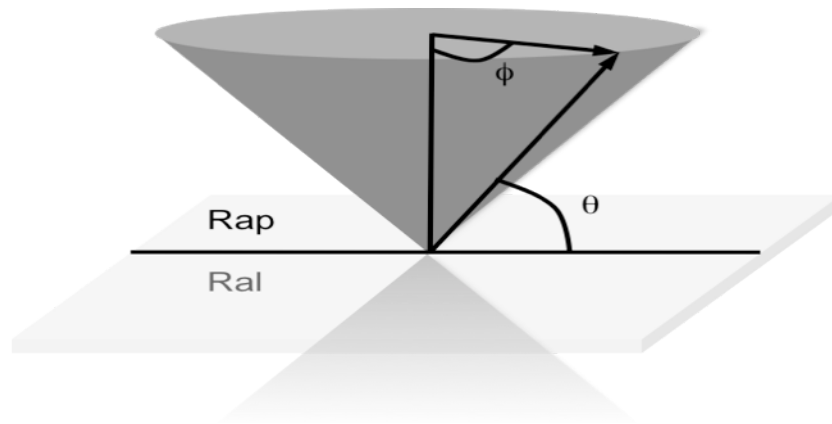
Mardia, *et al.*<sup>25</sup> have previously demonstrated that the variance for a cyclic average can be described by equation (2-11):

$$\langle \theta^2 \rangle - \langle \theta \rangle^2 = 1 - \frac{1}{n} \sqrt{\left( \sum_{j=1}^n \cos \theta_j \right)^2 + \left( \sum_{j=1}^n \sin \theta_j \right)^2} \quad (2-11)$$

which was then rearranged to accommodate weighting, yielding equation (2-12):

$$\langle \theta^2 \rangle - \langle \theta \rangle^2 = 1 - \sqrt{\left[ \sum_{j=1}^{72} \left( p_j \cdot \frac{1}{m_j} \sum_{i=1}^{m_j} \sin \theta_j \right) \right]^2 + \left[ \sum_{j=1}^{72} \left( p_j \cdot \frac{1}{m_j} \sum_{i=1}^{m_j} \cos \theta_j \right) \right]^2} \quad (2-12)$$

**A**



**B**

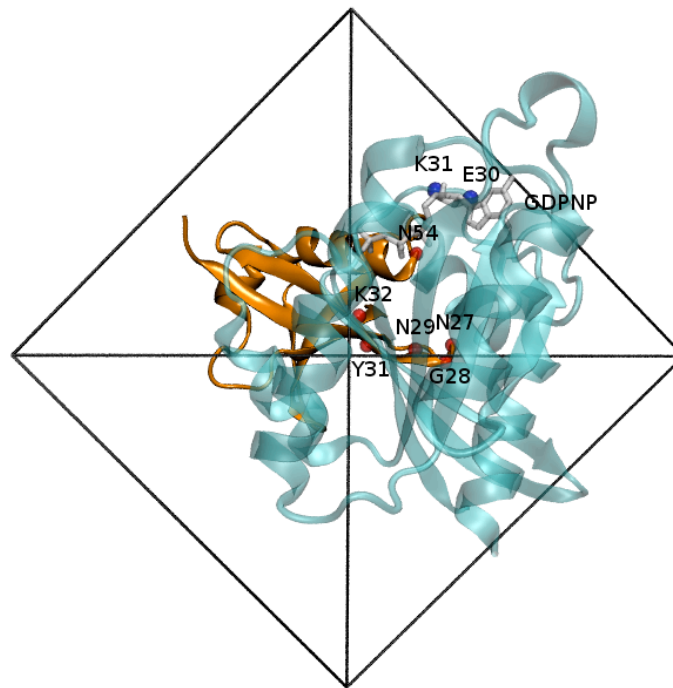


Figure 2-5: Representation of azimuthal and polar angles

Representation of the azimuthal ( $\theta$ ) and polar ( $\phi$ ) angles. (A) The surface plane was defined by the average of the plane fit with all C $\alpha$  atoms. Azimuthal angles are shown

relative to this plane, where  $R_{ap}$  is above the plane and  $R_{al}$  is below the plane. (B)  
Representation of the polar angle. Translating the black cross hairs to the  $C\alpha$  of each residue (represented by spheres) determines the origin of the polar angles presented in Figure 4-1 and Figure 4-2.  $R_{ap}$  is shown in blue and  $R_{al}$  in orange.

## **2.5 PROBE PARAMETERIZATION FOR AMOEBA**

<Body text to begin here.>

1.	atom	259	55	CB	"CNC CB"	6	12.011	4
2.	atom	260	56	SG	"CNC SG"	16	32.066	2
3.	atom	261	57	CD	"CNC CD"	6	12.011	2
4.	atom	262	58	NE	"CNC NE"	7	14.007	1
5.	atom	263	59	HB	"CNC HB"	1	1.008	1
6.								
7.	multipole	259	8	260	-0.15228			
8.					0.07407	0.00000	0.31740	
9.					-0.15117			
10.					0.00000	-0.21945		
11.					-0.19282	0.00000	0.37062	
12.	multipole	260	259	261	0.33074			
13.					0.44389	0.00000	0.28736	
14.					1.22369			
15.					0.00000	-2.16613		
16.					-0.31636	0.00000	0.94244	
17.	multipole	261	260	262	0.24556			
18.					0.06457	0.00000	-0.41800	
19.					0.15740			
20.					0.00000	0.22633		
21.					0.26601	0.00000	-0.38373	
22.	multipole	12	8	259	0.12898			
23.					0.02551	0.00000	0.07014	
24.					0.19051			
25.					0.00000	0.17012		
26.					-0.01582	0.00000	-0.36063	
27.	multipole	263	259	8	0.09179			
28.					-0.07114	0.00000	-0.02080	
29.					-0.16061			
30.					0.00000	-0.02258		
31.					-0.03267	0.00000	0.18319	
32.	multipole	262	261	260	-0.58843			
33.					-0.00512	0.00000	-0.24238	
34.					0.32332			
35.					0.00000	0.16836		
36.					-0.08202	0.00000	-0.49168	
37.								
38.	polarize	259		1.3340	0.3900	263		
39.	polarize	260		3.3000	0.3900	261		
40.	polarize	261		1.3340	0.3900	260	262	
41.	polarize	262		1.0730	0.3900	261		
42.	polarize	263		0.4960	0.3900	259		
43.								
44.	vdw	58		3.7100	0.1050			
45.	vdw	55		3.7800	0.1060			
46.	vdw	56		4.0050	0.3550			
47.	vdw	57		3.7800	0.1060			
48.	vdw	59		2.8700	0.0330	0.900		
49.	bond	55	7	323.0000	1.5317			
50.	bond	55	56	235.8000	1.8353			
51.	bond	55	59	341.0000	1.0817			
52.	bond	56	57	235.8000	1.7060			
53.	bond	57	58	450.0000	1.1374			
54.	angle	7	55	56	53.2000	109.6948		

Code and Parameters 2-1: Cyanocysteine AMOEBA Parameters

1.	atom	271	67	H1	"MeSCN H1"	1	1.008	1
2.	atom	272	68	C1	"MeSCN C1"	6	12.011	4
3.	atom	273	69	S	"MeSCN S"	16	32.066	2
4.	atom	274	70	C	"MeSCN C"	6	12.011	2
5.	atom	275	71	N	"MeSCN N"	7	14.007	1
6.								
7.	multipole	272	273	271	-0.22754			
8.					0.00000	0.00000	-0.02449	
9.					0.72145			
10.					0.00000	0.72145		
11.					0.00000	0.00000	-1.44290	
12.	multipole	273	272	274	0.33074			
13.					0.44389	0.00000	0.28736	
14.					1.22369			
15.					0.00000	-2.16613		
16.					-0.31636	0.00000	0.94244	
17.	multipole	274	273	275	0.24556			
18.					0.06457	0.00000	-0.41800	
19.					0.15740			
20.					0.00000	0.22633		
21.					0.26601	0.00000	-0.38373	
22.	multipole	271	272	273	0.07989			
23.					0.02678	0.00000	-0.31473	
24.					0.66454			
25.					0.00000	-0.00998		
26.					0.17231	0.00000	-0.65456	
27.	multipole	275	274	273	-0.58843			
28.					-0.00512	0.00000	-0.24238	
29.					0.32332			
30.					0.00000	0.16836		
31.					-0.08202	0.00000	-0.49168	
32.								
33.	polarize	271			0.4960	0.3900	272	
34.	polarize	272			1.3340	0.3900	271 273	
35.	polarize	273			3.3000	0.3900	272 274	
36.	polarize	274			1.3340	0.3900	273 275	
37.	polarize	275			1.0730	0.3900	274	
38.								
39.	vdw	68			3.7800	0.1060		
40.	vdw	69			4.0050	0.3550		
41.	vdw	70			3.7800	0.1060		
42.	vdw	67			2.8700	0.0330	0.900	
43.	vdw	71			3.7100	0.1050		
44.	bond	68	69		235.8000	1.8209		
45.	bond	68	67		341.0000	1.0794		
46.	bond	69	70		235.8000	1.7068		
47.	bond	70	71		450.0000	1.1373		
48.	angle	69	68	67	60.2400	110.5747		
49.	angle	67	68	67	39.5700	110.6263		
50.	angle	68	69	70	60.4300	99.4461		
51.	angle	69	70	71	60.0000	178.9042		
52.	strbnd	69	68	67	11.5000	11.5000		
53.	strbnd	68	69	70	-5.7500	-5.7500		
54.	torsion	67	68	69 70	0.0000 0.0 1	0.0000 180.0 2	0.6600 0.0 3	



```
55. torsion      68   69   70   71      0.0000 0.0 1  0.0000 180.0 2  0.5000 0.0 3
```

Code and Parameters 2-2: Methyl Thiocyanate AMOEBA Parameters

1.	atom	264	60	C1	"EtSCN C1"	6	12.011	4
2.	atom	265	61	H1	"EtSCN H1"	1	1.008	1
3.	atom	266	62	C2	"EtSCN C2"	6	12.011	4
4.	atom	267	63	H2	"EtSCN H2"	1	1.008	1
5.	atom	268	64	S	"EtSCN S"	16	32.066	2
6.	atom	269	65	C	"EtSCN C"	6	12.011	2
7.	atom	270	66	N	"EtSCN N"	7	14.007	1
8.								
9.	multipole	264	266	265	-0.18333			
10.					0.00000	0.00000	0.32541	
11.					-0.25019			
12.					0.00000	-0.25019		
13.					0.00000	0.00000	0.50038	
14.	multipole	266	264	268	-0.23503			
15.					0.06819	0.00000	0.15097	
16.					-0.26732			
17.					0.00000	-0.08487		
18.					-0.08743	0.00000	0.35219	
19.	multipole	267	266	264	0.09244			
20.					-0.00077	0.00000	-0.03413	
21.					0.02305			
22.					0.00000	0.01348		
23.					-0.00070	0.00000	-0.03653	
24.	multipole	265	264	266	0.08187			
25.					0.01969	0.00000	-0.08459	
26.					0.06966			
27.					0.00000	0.02253		
28.					-0.00032	0.00000	-0.09219	
29.	multipole	268	266	269	0.33074			
30.					0.44389	0.00000	0.28736	
31.					1.22369			
32.					0.00000	-2.16613		
33.					-0.31636	0.00000	0.94244	
34.	multipole	269	268	270	0.24556			
35.					0.06457	0.00000	-0.41800	
36.					0.15740			
37.					0.00000	0.22633		
38.					0.26601	0.00000	-0.38373	
39.	multipole	270	269	268	-0.58843			
40.					-0.00512	0.00000	-0.24238	
41.					0.32332			
42.					0.00000	0.16836		
43.					-0.08202	0.00000	-0.49168	
44.								
45.	polarize	264		1.3340	0.3900	265	266	
46.	polarize	265		0.4960	0.3900	264		
47.	polarize	266		1.3340	0.3900	264	267	
48.	polarize	267		0.4960	0.3900	266		
49.	polarize	268		3.3000	0.3900	269		
50.	polarize	269		1.3340	0.3900	268	270	
51.	polarize	270		1.0730	0.3900	269		

[illegible]

### Code and Parameters 2-3: Ethyl Thiocyanate AMOEBA Parameters

1.	atom	276	72	C1	"HxSCN C1"	6	12.011	4
2.	atom	277	73	C2	"HxSCN C2"	6	12.011	4
3.	atom	278	74	H1	"HxSCN H1"	1	1.008	1
4.	atom	279	75	C3	"HxSCN C3"	6	12.011	4
5.	atom	280	76	H2	"HxSCN H2"	1	1.008	1
6.	atom	281	77	H3	"HxSCN H3"	1	1.008	1
7.	atom	282	78	C4	"HxSCN C4"	6	12.011	4
8.	atom	283	79	C5	"HxSCN C5"	6	12.011	4
9.	atom	284	80	H4	"HxSCN H4"	1	1.008	1
10.	atom	285	81	H5	"HxSCN H5"	1	1.008	1
11.	atom	286	82	C6	"HxSCN C6"	6	12.011	4
12.	atom	287	83	H6	"HxSCN H6"	1	1.008	1
13.	atom	288	84	S	"HxSCN S"	16	32.066	2
14.	atom	289	85	C	"HxSCN C"	6	12.011	2
15.	atom	290	86	N	"HxSCN N"	7	14.007	1
16.								
17.	multipole	279	282	277	-0.12665			
18.					0.16365	0.00000	0.14302	
19.					0.06092			
20.					0.00000	-0.43628		
21.					-0.27661	0.00000	0.37536	
22.	multipole	276	277	278	-0.15938			
23.					0.00000	0.00000	0.26734	
24.					-0.20136			
25.					0.00000	-0.20136		
26.					0.00000	0.00000	0.40272	
27.	multipole	283	282	286	-0.11656			
28.					0.24230	0.00000	0.09285	
29.					0.15205			
30.					0.00000	-0.43470		
31.					-0.43635	0.00000	0.28265	
32.	multipole	282	279	283	-0.11327			
33.					0.22022	0.00000	0.08479	
34.					0.15602			
35.					0.00000	-0.44910		
36.					-0.36114	0.00000	0.29308	
37.	multipole	277	279	276	-0.12195			
38.					0.19367	0.00000	0.15795	
39.					-0.01138			
40.					0.00000	-0.31282		
41.					-0.22362	0.00000	0.32420	
42.	multipole	286	283	288	-0.18161			
43.					0.04843	0.00000	0.18545	
44.					-0.01034			
45.					0.00000	-0.10804		
46.					-0.18618	0.00000	0.11838	
47.	multipole	280	277	279	0.05998			
48.					0.00441	0.00000	-0.07636	

49.				0.04595		
50.				0.00000	0.03338	
51.				-0.00973	0.00000	-0.07933
52.	multipole	284	282	279	0.06214	
53.				0.02649	0.00000	-0.05387
54.				0.01977		
55.				0.00000	0.03305	
56.				-0.00324	0.00000	-0.05282
57.	multipole	281	279	282	0.05805	
58.				0.00916	0.00000	-0.05550
59.				0.03877		
60.				0.00000	0.02937	
61.				-0.06123	0.00000	-0.06814
62.	multipole	285	283	282	0.07555	
63.				0.01936	0.00000	-0.04776
64.				0.05942		
65.				0.00000	0.03376	
66.				-0.01589	0.00000	-0.09318
67.	multipole	278	276	277	0.05847	
68.				0.01125	0.00000	-0.10159
69.				0.06569		
70.				0.00000	0.03435	
71.				-0.03952	0.00000	-0.10004
72.	multipole	287	286	283	0.07235	
73.				0.02117	0.00000	-0.04358
74.				0.04575		
75.				0.00000	0.03869	
76.				-0.01006	0.00000	-0.08444
77.	multipole	288	286	289	0.33074	
78.				0.44389	0.00000	0.28736
79.				1.22369		
80.				0.00000	-2.16613	
81.				-0.31636	0.00000	0.94244
82.	multipole	289	288	290	0.24556	
83.				0.06457	0.00000	-0.41800
84.				0.15740		
85.				0.00000	0.22633	
86.				0.26601	0.00000	-0.38373
87.	multipole	290	289	288	-0.58843	
88.				-0.00512	0.00000	-0.24238
89.				0.32332		
90.				0.00000	0.16836	
91.				-0.08202	0.00000	-0.49168
92.						
93.	polarize	276		1.3340	0.3900	277 278
94.	polarize	277		1.3340	0.3900	276 280
95.	polarize	278		0.4960	0.3900	276
96.	polarize	279		1.3340	0.3900	281
97.	polarize	280		0.4960	0.3900	277
98.	polarize	281		0.4960	0.3900	279
99.	polarize	282		1.3340	0.3900	284
100.	polarize	283		1.3340	0.3900	285
101.	polarize	284		0.4960	0.3900	282
102.	polarize	285		0.4960	0.3900	283
103.	polarize	286		1.3340	0.3900	287
104.	polarize	287		0.4960	0.3900	286
105.	polarize	288		3.3000	0.3900	289

106.	polarize	289		1.3340	0.3900	288	290
107.	polarize	290		1.0730	0.3900	289	
108.							
109.	vdw	75		3.8200	0.1010		
110.	vdw	72		3.8200	0.1010		
111.	vdw	79		3.8200	0.1010		
112.	vdw	78		3.8200	0.1010		
113.	vdw	73		3.8200	0.1010		
114.	vdw	82		3.7800	0.1060		
115.	vdw	84		4.0050	0.3550		
116.	vdw	85		3.7800	0.1060		
117.	vdw	76		2.9800	0.0240	0.940	
118.	vdw	80		2.9800	0.0240	0.940	
119.	vdw	77		2.9800	0.0240	0.940	
120.	vdw	81		2.9800	0.0240	0.940	
121.	vdw	74		2.9600	0.0240	0.920	
122.	vdw	83		2.8700	0.0330	0.900	
123.	vdw	86		3.7100	0.1050		
124.	bond	75	78	453.0000	1.5299		
125.	bond	75	73	453.0000	1.5298		
126.	bond	75	77	341.0000	1.0885		
127.	bond	72	73	323.0000	1.5282		
128.	bond	72	74	341.0000	1.0863		
129.	bond	79	78	453.0000	1.5312		
130.	bond	79	82	345.3000	1.5276		
131.	bond	79	81	341.0000	1.0866		
132.	bond	78	80	341.0000	1.0879		
133.	bond	73	76	341.0000	1.0878		
134.	bond	82	84	235.8000	1.8351		
135.	bond	82	83	341.0000	1.0810		
136.	bond	84	85	235.8000	1.7063		
137.	bond	85	86	450.0000	1.1375		
138.	angle	78	75	73	48.2000	113.1072	
139.	angle	78	75	77	38.0000	109.3223	
140.	angle	73	75	77	38.0000	109.2787	
141.	angle	77	75	77	39.5700	106.3036	
142.	angle	73	72	74	42.4400	111.1869	
143.	angle	74	72	74	39.5700	107.7267	
144.	angle	78	79	82	48.2000	111.5264	
145.	angle	78	79	81	38.0000	109.6162	
146.	angle	82	79	81	38.0000	109.5120	
147.	angle	81	79	81	39.5700	106.9493	
148.	angle	75	78	79	48.2000	112.7564	
149.	angle	75	78	80	38.0000	109.3923	
150.	angle	79	78	80	38.0000	109.3217	
151.	angle	80	78	80	39.5700	106.4572	
152.	angle	75	73	72	48.2000	112.9179	
153.	angle	75	73	76	38.0000	109.3170	
154.	angle	72	73	76	38.0000	109.3845	
155.	angle	76	73	76	45.5700	106.3035	
156.	angle	79	82	84	53.2000	109.0812	
157.	angle	79	82	83	42.4400	111.1744	
158.	angle	84	82	83	60.2400	108.2095	
159.	angle	83	82	83	39.5700	108.8985	
160.	angle	82	84	85	60.4300	99.5693	
161.	angle	84	85	86	60.0000	179.0730	
162.	strbnd	78	75	73	18.7000	18.7000	

163.	strbnd	78	75	77	11.5000	18.7000							
164.	strbnd	73	75	77	11.5000	18.7000							
165.	strbnd	73	72	74	11.5000	11.5000							
166.	strbnd	78	79	82	18.7000	18.7000							
167.	strbnd	78	79	81	11.5000	18.7000							
168.	strbnd	82	79	81	11.5000	18.7000							
169.	strbnd	75	78	79	18.7000	18.7000							
170.	strbnd	75	78	80	11.5000	18.7000							
171.	strbnd	79	78	80	11.5000	18.7000							
172.	strbnd	75	73	72	18.7000	18.7000							
173.	strbnd	75	73	76	11.5000	18.7000							
174.	strbnd	72	73	76	11.5000	18.7000							
175.	strbnd	79	82	84	18.7000	18.7000							
176.	strbnd	79	82	83	11.5000	11.5000							
177.	strbnd	84	82	83	11.5000	11.5000							
178.	strbnd	82	84	85	-5.7500	-5.7500							
179.	torsion	73	75	78	79	0.576	0.0	1	-				
0.017	180.0	2	2.031	0.0	3								
180.	torsion	73	75	78	80	0.000	0.0	1	0.000	180.0	2	0.000	
0.0	3												
181.	torsion	77	75	78	79	0.000	0.0	1	0.000	180.0	2	0.000	
0.0	3												
182.	torsion	77	75	78	80	0.000	0.0	1	0.000	180.0	2	0.000	
0.0	3												
183.	torsion	78	75	73	72	0.484	0.0	1	0.014	180.0	2	2.221	
0.0	3												
184.	torsion	78	75	73	76	0.000	0.0	1	0.000	180.0	2	0.000	
0.0	3												
185.	torsion	77	75	73	72	0.000	0.0	1	0.000	180.0	2	0.000	
0.0	3												
186.	torsion	77	75	73	76	0.000	0.0	1	0.000	180.0	2	0.000	
0.0	3												
187.	torsion	74	72	73	75	0.0000	0.0	1	0.0000	180.0	2	0.3410	
0.0	3												
188.	torsion	74	72	73	76	0.0000	0.0	1	0.0000	180.0	2	0.2990	
0.0	3												
189.	torsion	82	79	78	75	0.364	0.0	1	-				
0.024	180.0	2	1.958	0.0	3								
190.	torsion	82	79	78	80	0.000	0.0	1	0.000	180.0	2	0.000	
0.0	3												
191.	torsion	81	79	78	75	0.000	0.0	1	0.000	180.0	2	0.000	
0.0	3												
192.	torsion	81	79	78	80	0.000	0.0	1	0.000	180.0	2	0.000	
0.0	3												
193.	torsion	78	79	82	84	-							
0.688	0.0	1	0.489	180.0	2	1.957	0.0	3					
194.	torsion	78	79	82	83	0.000	0.0	1	0.000	180.0	2	0.000	
0.0	3												
195.	torsion	81	79	82	84	0.000	0.0	1	0.000	180.0	2	0.000	
0.0	3												
196.	torsion	81	79	82	83	0.000	0.0	1	0.000	180.0	2	0.000	
0.0	3												
197.	torsion	79	82	84	85	-							
2.643	0.0	1	0.709	180.0	2	1.072	0.0	3					
198.	torsion	83	82	84	85	0.000	0.0	1	0.000	180.0	2	0.000	
0.0	3												

```
199.      torsion      82   84   85   86      0.0000 0.0 1  0.0000 180.0 2  0.5000
0.0 3
```

Code and Parameters 2-4: Hexyl Thiocyanate AMOEBA Parameters



## **2.6 SMALL MOLECULE SIMULATIONS IN AMOEBA**

<Body text to begin here.>

## **Chapter 3 Electrostatic Field Methods**

### **3.1 AMBER03 WITH EXPLICIT TIP3P WATER**

#### **3.1.1 Reaction Field Electrostatics**

#### **3.1.2 Hybrid Solvent Reaction Field Electrostatics and Solute Coulomb Field**

<Body text to begin here.>

## **3.2 AMBER03 WITH POISSON-BOLTZMANN CONTINUUM SOLVENT**

### **3.2.1 Reaction Field Method**

<Body text to begin here.>

### **3.2.2 Grid spacing and size**

<Body text to begin here.>

### **3.2.3 Box Location**

<Body text to begin here.>

### **3.3 AMBER03 WITH POISSON-BOLTZMANN CONTINUUM SOLVENT AND SELECT EXPLICIT TIP3P WATER MOLECULES**

#### **3.3.1 5 Å Water Sphere Around the Vibrational Chromophore**

<Body text to begin here.>

#### **3.3.2 Single Water Molecule Nearest the Vibrational Chromophore**

<Body text to begin here.>

#### **3.3.3 Water Molecular Hydrogen Bonding to the Vibrational Chromophore**

<Body text to begin here.>

## **3.4 AMOEBA**

### **3.4.1 Poisson-Boltzmann Continuum Solvent**

<Body text to begin here.>

### **3.4.2 Explicit AMOEBA Water**

<Body text to begin here.>

### **3.4.3 Charge Penetration Field Corrections**

<Body text to begin here.>

## **Chapter 4 The Role of Electrostatics in Differential Binding of RalGDS to Rap Mutations E30D and K31E Investigated by Vibrational Spectroscopy of Thiocyanate Probes**

### **4.1 INTRODUCTION**

The human proteins p21<sup>Ras</sup> (hereafter Ras) and Rap1A (hereafter Rap) are members of the Ras family of guanosine triphosphate (GTP)-hydrolyzing proteins that switch between an ON state when bound to GTP and an OFF state when bound to guanosine diphosphate (GDP) in the regulation of signal transduction pathways<sup>26</sup>. When in the GTP-bound ON state, both proteins dock to the Ras binding domain (RBD) of downstream effector proteins to propagate a signal transduction cascade. As members of the Ras superfamily of GTPases, Ras and Rap are similar in both sequence and structure; they share 50% amino acid identity<sup>6</sup>, 80% amino acid homology, and have nearly identical structure and effector binding surfaces (rmsd of 0.7 Å for homologous residues).

Despite these similarities, Ras and Rap have very different functions within the cell. Ras is involved in the propagation of chemical signals that regulate cell division, cell survival and apoptosis, and has been a focus of research in molecular oncology for many years<sup>27-29</sup>. The biological function of Rap is less well understood, but it has recently been shown to play a role in cellular adhesion<sup>30</sup> and has been implicated in cancer metastasis<sup>31,32</sup>. The origin of the functional specificity of these two similar proteins lies in the ability of Ras and Rap to bind to different downstream effectors while in the GTP-bound ON state, thus initiating different signaling cascades. One of the best studied downstream effector proteins of Ras is the human protein c-Raf-1 (hereafter Raf), while Rap is thought to interact most strongly with the Ras binding domain (RBD) of the downstream effector Ral guanosine dissociation stimulator (RalGDS, hereafter Ral).

While the structures of the RBD of the downstream effector proteins are very similar<sup>5,6</sup>, measurements of thermodynamic<sup>2,33-35</sup> parameters of the protein-protein binding events ( $DG_{dock}$  and  $K_d$ ) indicate high specificity between the two GTPases and their appropriate downstream effector. Given the structural similarities of these protein-protein interfaces, a possible mechanism for the binding specificity that each GTPase shows for its appropriate downstream effector could be based on differences in electrostatic complementarity at the protein-protein interface. To investigate this hypothesis, the binding of these two GTPases to a variant of Ral have previously been investigated by our laboratory through both experimental and computational methods<sup>1,2</sup> in order to explore the contribution of electrostatic fields to interface formation in a system in which structural factors are controlled as much as possible.

In 1995, Herrmann and coworkers<sup>35,36</sup> demonstrated the importance of the amino acids at positions 30 and 31 on Ras and Rap in discriminating downstream effector partners. These amino acids, Ras Asp 30 and Glu 31 and Rap Glu 30 and Lys 31, are positioned at the protein-protein binding interface, and studies involving reversion mutations of Rap to Ras have shown that these residues are involved with the interactions with downstream binding partners. Those studies found that the charge reversion mutation Rap K31E and the double mutation Rap E30D/K31E cause the dissociation constant with the downstream effector Raf to resemble Ras D30/E31 more closely than Rap E30/K31<sup>36</sup>. Furthermore, the double mutant Rap E30D/K31E co-crystallized with the downstream effector Raf, which usually binds more strongly to Ras. Similarly, the mutation Ras E31K was used to co-crystallize Ras with Ral, normally Rap's immediate downstream effector<sup>5</sup>, and the double reversion mutant Ras D30E/E31K has significantly reduced binding affinity with Raf<sup>37</sup>. The role of these mutations in creating strong binding

affinity for an alternative downstream effector suggested an electrostatic mechanism caused by altering the charge of the residue at position 31 on either GTPase.

The arrangement of amino acid charges from secondary and tertiary protein structure can generate large and heterogeneous electrostatic fields that affect all protein function, including protein-protein interactions.<sup>38-43</sup> Vibrational Stark effect (VSE) spectroscopy is a recently developed experimental technique capable of measuring electrostatic fields in proteins<sup>9,44-50</sup>. In VSE spectroscopy, the intrinsic response of a probe vibrational oscillator to its local electrostatic environment is measured spectroscopically and is used to quantitate the magnitude and direction of the local electric field to which the probe is exposed. After calibration, the probe is inserted into a known position of a protein where it becomes a highly local, sensitive, and directional reporter of fluctuations of the protein's electrostatic field caused by structural or chemical perturbations to the protein. For example, the formation of a protein-protein interface through a docking interaction may induce changes in the absorption energy of a probe located on each protein surface as that probe is exchanged from an electrostatic field generated by the protein-water interface to one generated by the protein-protein interface.

The nitrile stretching vibration is one of several useful VSE probes that have been identified<sup>47</sup>. Several characteristics of the nitrile oscillator makes it particularly attractive; its absorption energy of  $\sim 2100\text{-}2250\text{ cm}^{-1}$  is in a region of the spectrum that is outside the vibrational background of a complex biomolecule, it is reasonably sensitive to fields of the magnitude thought to be present in proteins, and it can be incorporated into proteins with relative ease<sup>51</sup>. There are a growing number of reports describing VSE spectroscopy of nitrile probes to study the function of electrostatic fields in enzyme active sites<sup>9,46,48-50</sup>. Recently in our laboratory, VSE spectroscopy has been used to study the change in absorption energy of the nitrile probe incorporated at eleven locations on



the surface of Ral that become buried in the protein-protein interface upon binding to either Ras D30/E31 or Rap E30/K31<sup>2</sup>. This study determined that at some probe locations the change in absorption upon binding to Ras D30/E31 and Rap E30/K31 was the same, while at other locations the change was different, suggesting that these latter amino acids may participate in an electrostatic mechanism that enables Ral to distinguish Ras from Rap. Furthermore, measurement of the dissociation constant of the docking of the GTPases with nitrile-containing Ral mutants showed no deleterious effect of the spectroscopic probe on the formation of the docked complex. These measurements were later confirmed through extensive molecular dynamics sampling of the protein-protein complex<sup>1</sup>. We hypothesized that several of these Ral-based nitrile probes would be ideal for investigating our hypothesis of an electrostatic mechanism for binding discrimination mediated by the charge on position 31 of these two GTPases.

In this report, we describe the systematic investigation of the effect of the E30D and K31E mutations to Rap on the local electrostatic fields formed by docking to the downstream effector and measured by the nitrile probe at six locations on the surface of the RBD of Ral using VSE, molecular dynamics (MD) simulations, and dissociation constant ( $K_d$ ) measurements. We selected 6 amino acids on Ral that are positioned in the protein-protein interface near positions 30 and 31 of Ras and Rap when the docked complex is formed: N27, G28, N29, Y31, K32, and N54. The locations of these amino acids within the Rap-Ral interface, as well as Rap positions 30 and 31, are shown in Figure 1. Positions N27 and Y31 were selected because at these probe locations, we previously measured a significant difference between docking of Ras D30/E31 versus Rap E30/K31<sup>2</sup>, and wanted to investigate whether these observations were due to the different amino acid identity at positions 30 and 31. Position N29 was chosen because, along with position N27, molecular dynamics (MD) structural sampling of the

thiocyanate side chain within the docked complex showed that the nitrile group has the largest angle with respect to the Ras-Ral interfacial plane of the collection of Ral-based probes we have investigated, approximately  $45^\circ$  from the plane of the Ras-Ral interface<sup>1</sup>. Because the effect of the change in electrostatic field is on the projection of the field vector onto the nitrile probe bond axis, these positions seemed the most likely to feel the effect of the field vector at positions 30 and 31 of Rap, which appear from the crystal structures 1LFD<sup>5</sup> and 1GUA<sup>6</sup> to be approximately perpendicular to the interface. Gly28 was chosen as intermediate between these two residues. Finally, positions K32 and N54 were selected based on their physical proximity to the side chains of positions 30 and 31 on the GTPase once the docked complex formed.

Each of the six Ral amino acids were mutated to cysteine, then chemically modified to introduce the nitrile VSE probe in the form of the thiocyanate. These mutants were then bound to Rap E30/K31, Rap E30D, Rap K31E, and the double mutant Rap E30D/K31E. The difference in vibrational absorption energy,  $Dn_{obs}$ , between the Ral monomer and the docked complex was measured by Fourier transform infrared spectroscopy (FTIR). The dissociation constant of the docked complex,  $K_d$ , was determined through a fluorescence assay. Finally, extensive MD simulations on docked complexes of all Rap and Ral variants were conducted to determine Boltzmann-weighted orientational data for the Ral-based nitrile probes and for the side chains at Rap positions 30 and 31. These simulations demonstrate that the mutation K31E is almost exclusively responsible for changes in side chain orientations at Rap positions 30 and 31 that cause the observed change in  $K_d$ . VSE spectroscopy demonstrates that only two of the probe locations examined, at N27C and N29C displayed a change in the absorption energy upon binding the Ras-like Rap double mutants that strongly resembled the change for Ras D30/E31. However, several of these probes did respond in an additive manner to the

individual single mutations. In some cases, these effects are replicated in subtle structural changes at positions 30 and 31 observed in molecular dynamics sampling. These studies support both a structural and electrostatic mechanism to explain observed differences in GTPase-effector binding.

## **4.2 RESULTS**

### **4.2.1 Dissociation Constant Measurements**

Measured dissociation constants of both WT and SCN-labeled Ral $\beta$  mutants docking with all GTPases studied here are reported in Table 4-1. Binding of WT Ral to Rap E30/K31 was approximately 10-fold faster than binding to Ras D30/E31, as has been reported before<sup>2, 35, 36</sup>. Any effect of the thiocyanate probe on the formation of the interface can be investigated by comparing  $K_d$  values obtained for wild type (WT) Ral versus the SCN-labeled Ral $\beta$  mutants docking to Ras D30/E31 and Rap E30/K31 in Table 4-1. Although there are some variations in the absolute value of these numbers, they are for the most part small. The two largest deviations, caused by Ral $\beta$  G28C<sub>SCN</sub> (reduced  $K_d$  by an order of magnitude) and Ral $\beta$  Y31C<sub>SCN</sub> (increased  $K_d$  by an order of magnitude), still showed an order of magnitude increase in  $K_d$  when binding to Ras D30/E31 as opposed to Rap E30/K31, as is expected from our previous work. The dissociation constants presented in Table 4-1 demonstrate that the presence of the SCN probe on Ral $\beta$  mutants did not substantially affect binding to either Ras D30/E31 or Rap E30/K31, as has been observed before<sup>2</sup>. Along with results from molecular dynamics sampling of this system, described below, this is strong circumstantial evidence that the nitrile VSE probe does not significantly alter the interface formed between Rap and the SCN-labeled Ral $\beta$  mutants compared to the WT interaction. Our experimental

mutagenesis and chemical labeling strategy therefore does not destroy the protein-protein interaction that we are attempting to measure.

The results in Table 4-1 confirm previous reports that the reversion mutation at Rap position 31 alters the binding interaction between Rap and Ral to resemble that of Ras and Ral. All SCN-labeled Ral $\beta$  mutants interacted with Rap E30/K31E with a  $K_d$  10-fold higher than with Rap E30/K31, and this effect was preserved in the double mutant Rap E30D/K31E. The single mutation Rap E30D/K31 had no effect on binding, and all  $K_d$  values measured with that construct were essentially identical to Rap E30/K31. It is therefore clear that the amino acid located at position 31 of the GTPase is critical in the mechanism that enables Ral to distinguish structurally similar but functionally distinct GTPases for appropriate binding. Exploring the structural and electrostatic components of that mechanism is the subject of the MD sampling and VSE spectroscopy discussed here.

#### **4.2.2 Molecular Dynamics Simulations**

Structural details of the interface formed between Rap and Ral have been investigated with two high-quality crystal structures, 1GUA and 1LFD<sup>5,52</sup>. These are crucial starting points for any comprehensive biochemical understanding of the Rap-Ral interactions, but the static structures of these proteins might not represent the full range of side chain conformations of either our SCN probe or positions 30 and 31 of the GTPase, and so are not sufficient for our purposes. Furthermore, our experiments are on an unnatural variant of the system, the SCN-labeled Ral $\beta$  mutant. Although our  $K_d$  measurements indicate that the presence of the thiocyanate probe does not significantly disrupt interface formation, knowledge of the structure and orientation of the probe within the interface is necessary to interpret vibrational absorption data. To address these

issues, we performed extensive molecular dynamics sampling of each SCN-labeled Ral $\beta$  mutant docked with Rap E30/K31 and the Rap mutants E30D/K31, E30/K31E, and E30D/K31E to accumulate a Boltzmann-weighted ensemble of the orientations of the thiocyanate group and the side chains at Rap positions 30 and 31. The WHAM-derived  $\chi_2$  torsional probability distribution for all thiocyanate groups and  $\chi_1$  torsional probability distribution side chains at Rap positions 30 and 31 are shown in Figures S1-S3. Similar to what we have seen previously<sup>53</sup>, the torsional distribution for all studied side chains was characteristic of an unhindered alkyl group, with three probability maxima separated by  $\sim 120^\circ$  and essentially no significant difference depending on the chemical identity of either the SCN-labeled Ral $\beta$  mutant in the docked complex (Figure S1) or the Rap variant (Figures S2 and S3). Our MD sampling strategy therefore provides us with a comprehensive ensemble of structures of each docked Rap-Ral $\beta$  variant studied with VSE spectroscopy to aid in interpreting our spectroscopic results.

To analyze our molecular dynamics simulations of the torsional distribution of the thiocyanate residue on each Ral $\beta$  mutant and on the side chains at Rap positions 30 and 31, we defined two angles for each side chain with respect to the Rap-Ral $\beta$  surface, which we term azimuthal and polar angles. These two angles are described in the Supplemental Materials, and are shown schematically in Figure 2. When the cross hairs on Figure 2B are translated to the C $\alpha$  atom of each simulated residue on Rap or Ral $\beta$  (represented as spheres), they become the origin of the polar angles shown in Figures 3 and S4. The average azimuthal and polar angles of the nitrile probe relative to the Rap-Ral $\beta$  binding interface are shown in Supplemental Figure S4. These simulations were compared with previous simulations of the Ras-Ral $\beta$  binding interface<sup>53</sup>, and in all cases, the orientation of the nitrile at the docked interface for Ras D30/E31 and Rap E30/K31 are identical, within the error distribution of the Boltzmann-weighted ensemble. This

indicates that the structure of the probe at the interface of these two GTPases is very similar, as expected by the structural similarities of the crystal structures, the amino acid homology of the GTPases, and the fact that the measured  $K_d$  values of the SCN-labeled Ral $\beta$  mutants dock with Ras D30/E31 and Rap with the same dissociation constant as WT Ral. Furthermore, as seen in Figure S4, mutations to positions 30 and 31 of Rap did not substantially alter the orientation of the nitrile probe at the interface, with all 6 nitrile positions remaining approximately constant with the same dihedral distribution for Rap E30/K31, or the E30D, K31E, E30D/K31E mutations. The consistency of the probe orientation in these Boltzmann-weighted ensembles is further indication of the structural robustness of these interacting proteins, making them suitable for systematic measurements of electrostatic fields at the protein-protein interface without compromising structural integrity of the interface. This observation is important to justify using VSE spectroscopy to explore the biochemistry of Rap-Ral binding with a spectroscopic technique that doesn't compromise the very interaction we are interesting in studying.

The azimuthal and polar angles of the side chains at Rap positions 30 and 31 for each mutant studied here are shown in Figure 4-2. When looking at position 30 (Figure 4-2A), the azimuthal angles showed essentially no difference depending on either the Rap mutant or the SCN-labeled Ral $\beta$  mutant to which it was docked. In most cases, the aspartate (for E30D and E30D/K31E) or glutamate vector (for WT and Rap K31E) pointed approximately parallel and slightly below the Rap-Ral surface plane. The single exception to this was when Rap E30D was docked with Ral $\beta$  K32C<sub>SCN</sub>, where it was observed that the residue oriented itself 5° above the Rap-Ral $\beta$  K32C<sub>SCN</sub> surface plane. The polar angles of the side chain at position 30 (Figure 4-2B) showed significantly more fluctuations, but these fluctuations were essentially random and all demonstrated very

large WHAM-derived standard deviations. This suggests that both aspartate and glutamate side chains at this position undergo large polar motions that are not influenced by the chemical identity of position 31 or the location of the thiocyanate probe.

The side chains at Rap position 31 behaved very differently. The azimuthal angles of the side chain at Rap position 31 (Figure 4-2C) were pointed significantly further below the Rap-Ral $\beta$  mutant plane than at position 30 (55-60° below the surface plane for position 31 versus < 25° below the surface plane for position 30), but with the exception of docked to Ral $\beta$  Y31C<sub>SCN</sub>, all Rap mutants behaved identically. The distribution of polar angles of mutants of Rap position 31, however, displayed a significant dependence on the identity of the residue at position 31 (Figure 4-2D). Although again the distribution of observed angles for each Rap mutant was large, an obvious trend appeared in the data. With the sole exception Rap mutants docked to Ral $\beta$  N54C<sub>SCN</sub>, Rap E30/K31 and the Rap E30D behaved identically, while the Rap mutants K31E and E30D/K31E also behaved identically. This strongly supports the evidence that mutations at Rap K31 are most important for determining binding to the downstream effector Ral, and that mutations to Rap E30 do not significantly affect the behavior of either Rap E30/K31 or Rap K31E. Both of these observations are consistent with our observations of differences in dissociation constants between WT and mutated Rap binding to the downstream effector Ral.

We investigated the causes of the difference in orientation between the lysine and glutamate side chains at Rap position 31 through inspection of representative MD snapshots. This clearly revealed that when a lysine is at position 31, it pointed towards a hydrogen bonding acceptor pocket formed by Ral $\beta$  D51, N54, and E58. When this side chain was mutated to Rap K31E in either the single or double mutant, the negatively charged Glu reoriented to avoid electrostatic repulsion with the Ral $\beta$  surface caused by

the hydrogen bond accepting pocket created by these three residues. This appears to be the central cause for both the dissociation constant and electrostatic differences between Rap-Ral $\beta$  binding and Ras-Ral $\beta$  binding described here. As discussed above, this behavior was not observed when the Rap constructs were docked with Ral $\beta$  N54C<sub>SCN</sub>; investigation of MD structures showed that this probe location disrupted the hydrogen bonding pocket sufficiently to cause this portion of the Ral $\beta$  surface retracted slightly from the Rap-Ral $\beta$  interface, leaving K31 without the hydrogen bonding pocket. Although the importance of this hydrogen bonding pocket can be hypothesized from the 1LFD crystal structure (of Ras variant E31K docked with RalGDS), our MD sampling has provided the first confirmation of the observed biochemical behavior of Ras and Rap with direct structural evidence.



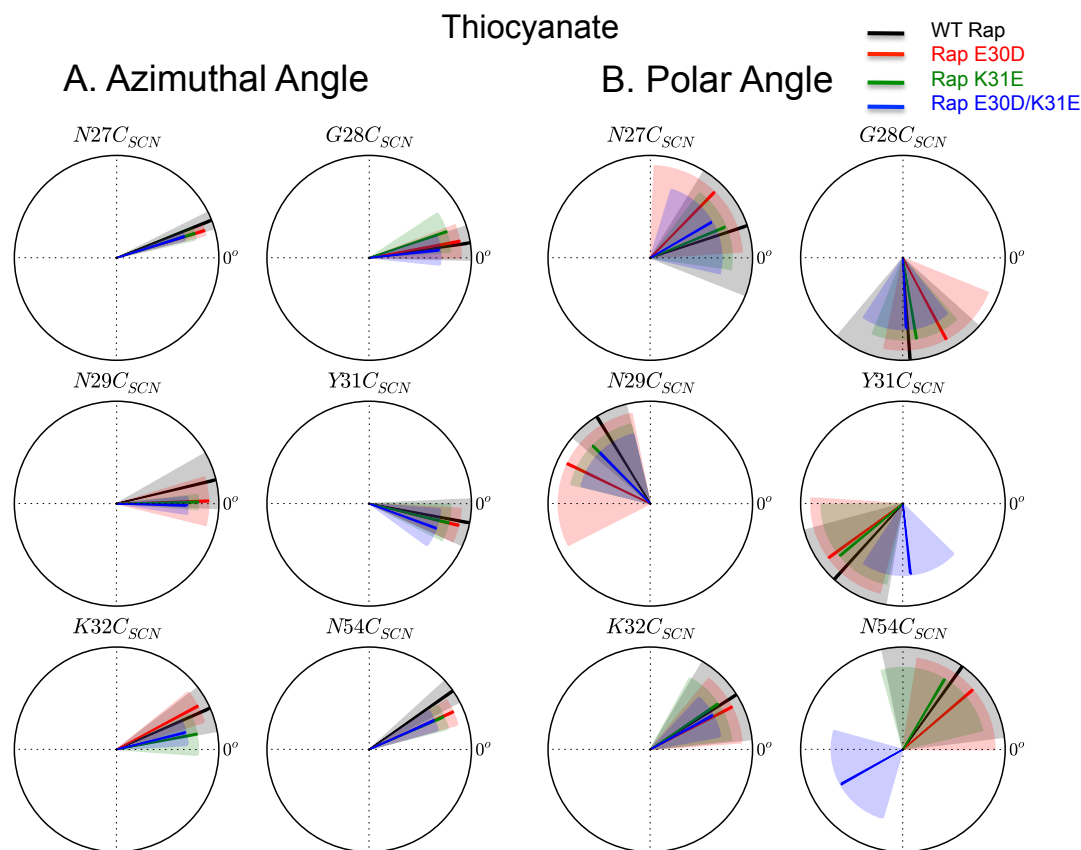


Figure 4-1: Azimuthal and Polar Angles of Simulated CNC Probes

(A) Azimuthal and (B) Polar angles of the thiocyanate on SCN-labeled Ral $\beta$  mutants calculated from each mutant docked with Rap E30/K31 (black), Rap E30D/K31 (red), Rap E30/K31E (green), and Rap E30D/K31E (blue). Azimuthal angles are shown relative to the Rap-Ral $\beta$  surface plane. Polar angles are shown relative to the coordinate system described in Figure 2-5B. The shaded area represents one standard deviation on the calculated angle from the Boltzmann-weighted ensemble of structures derived from Equation \*\*\*S6.

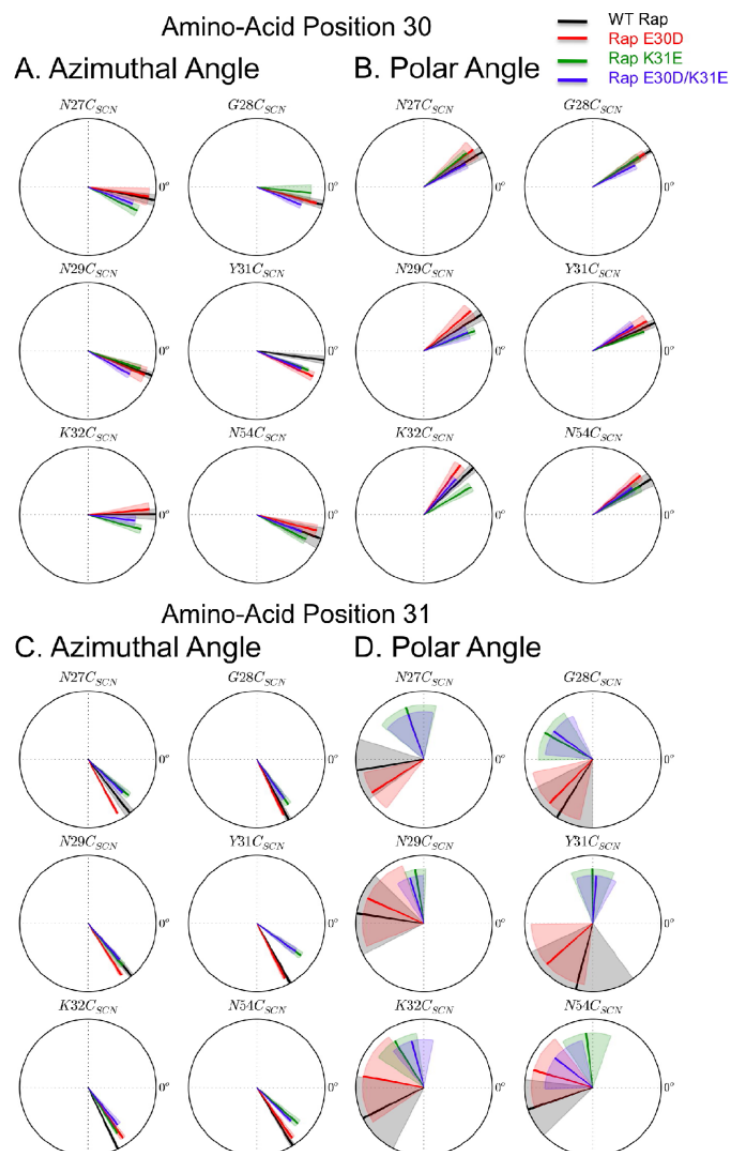


Figure 4-2: Azimuthal and Polar Angles of Simulated Rap Position 30 and 31 Sidechains

Azimuthal and polar angles of the side chain at Rap positions 30 and 31 in Rap E30/K31 (black), Rap E30D/K31 (red), Rap E30/K31E (green), and Rap E30D/K31E (blue). (A) Azimuthal angle at Rap position 30; (B) polar angle at Rap position 30; (C) azimuthal angle at Rap position 31; (D) polar angle at Rap position 31. Azimuthal angles are shown relative the Rap-Ral surface plane. Polar angles are shown relative to the

coordinate system shown in Figure 2-5. Dashed lines represent one standard deviation on the calculated angle from the Boltzmann-weighted ensemble of structures. \*\*\*Find better quality image

### 4.2.3 VSE Spectroscopy of the Docked Protein-Protein Complex

A principle goal of the work described here is to deconvolute the effects of structural and electrostatic influences on protein-protein interface formation. Our systematic molecular dynamics simulations allowed us to highlight differences between WT and mutant Rap that appeared to be due entirely to subtle changes in structure. With that information, a systematic exploration of electrostatic factors was conducted by measuring changes in electrostatic interactions between the GTPase and the downstream effector Ral by VSE spectroscopy of vibrational probes placed at the GTPase-effector interface. In this section we systematically review the spectroscopic results from each VSE probe, then draw general conclusions in the Discussion section.

*Ral $\beta$  N27C<sub>SCN</sub> and N29C<sub>SCN</sub>*: Selection of each of the 6 SCN-labeled Ral $\beta$  probes was based on consideration of each probes' orientation compared to the Rap-Ral $\beta$  interface, proximity to the Rap positions 30 and 31, and large differences in vibrational absorption energy upon binding to Ras D30/E31 and Rap E30/K31 measured in a previous study<sup>2</sup>. N27C<sub>SCN</sub> and N29C<sub>SCN</sub> were selected for probe placement because our Boltzmann-weighted molecular dynamics simulations of the orientations of the 6 thiocyanate probes when docked with each of the Rap constructs determined that thiocyanate group on these two Ral $\beta$  mutants had some of the largest angles with respect to the plane of the Ras-Ral $\beta$  interface of any of our probes, approximately 20-30° above the surface plane (Figure S4A)<sup>53</sup>. Because VSE spectroscopy is only sensitive to changes in the electrostatic field vector projected onto the nitrile bond axis, and because, as shown in Figure 4-2C, mutations to Rap K31 were themselves close to perpendicular to the Rap-Ral $\beta$  surface plane, probes perpendicular to the Rap-Ral $\beta$  plane would be most sensitive

to mutations to Rap K31. A Ral $\beta$ -based SCN probe that was itself as close to perpendicular to the Rap-Ral $\beta$  surface plane would therefore be in an appropriate position to observe changes caused by the charge reversion mutation at position 31.

An example of the VSE data collected here is shown in Figure 4. The Ral $\beta$  N29C<sub>SCN</sub> mutant was incubated with each Rap mutant, concentrated, and the absorption energy of the nitrile probe was recorded and compared with the measured absorption energy when bound to Rap E30/K31 from a previous study<sup>2</sup>. When docked with Rap E30D, the thiocyanate absorption energy was 2160.8 cm<sup>-1</sup>, identical to the observed absorption energy when Ral $\beta$  N29C<sub>SCN</sub> was bound to Rap E30/K31. When Ral $\beta$  N29C<sub>SCN</sub> was incubated with Rap K31E, the absorption energy was 0.6 cm<sup>-1</sup> higher in energy than when docked with Rap E30/K31, 2161.4 cm<sup>-1</sup>. However, when docked with the double mutant, Rap E30D/K31E, the absorption energy of the thiocyanate shifted 0.8 cm<sup>-1</sup> higher in energy (2161.6 cm<sup>-1</sup>). Both Rap mutants containing K31E were therefore more similar to the observed absorption energy when Ral $\beta$  N29C<sub>SCN</sub> is docked with Ras D30/E31 (2161.1 cm<sup>-1</sup>) than with Rap E30/K31. This means that the probe is experiencing an electrostatic environment in the double mutant E30D/K31E that is more like that of Ras D30/E31 than to Rap E30/K31, although the effect is small.

Changes in the absorption energy ( $\Delta\nu_{obs}$ ) of the thiocyanate on each SCN-labeled Ral $\beta$  mutant due to binding for each SCN-labeled Ral $\beta$  mutant studied here are summarized in Table 4-2: Measured vibrational frequencies of SCN-labeled Ral $\beta$  mutants docked with Rap E30/K31; the observed changes in vibrational frequency upon docking each probe to Ras D30/E31 and each Rap mutant and Figure 5. In these figures, all absorption energies are referenced to the absorption energy of the nitrile probe when docked with Rap E30/K31 (i.e.  $\Delta\nu_{obs} = 0$  represents no change from the thiocyanate absorption energy when docked with Rap E30/K31 reported previously)<sup>2</sup>. These shifts

are related to the change in electrostatic field through the known value of the Stark tuning rate of the cyanocysteine probe, of  $0.7 \text{ cm}^{-1}/(\text{MV}/\text{cm})$ ; the values of determined from equation 1 are reported in Table 4-2: Measured vibrational frequencies of SCN-labeled Ral $\beta$  mutants docked with Rap E30/K31; the observed changes in vibrational frequency upon docking each probe to Ras D30/E31 and each Rap mutant.

When the probe is located at Ral $\beta$  N27C<sub>SCN</sub> and N29C<sub>SCN</sub>, the data in Figure 5 clearly show that both the single and double reversion mutants Rap K31E and E30D/K31E have VSE absorption energies that are similar to Ras D30/E31, not Rap E30/K31. In both cases, the single mutation Rap E30D appears to have a negligible effect on the electrostatic environment of the probe. At these particular probe locations, therefore, the double reversion mutation does indeed revert the electrostatic field of Rap back to that found in Ras, supporting our hypothesis. Although these energy shifts are small, the trend towards higher absorption energies when bound to Ras D30/E31 and the double mutant Rap E30D/K31E than when bound to Rap E30/K31 are clear. The relatively small effect of mutations on the magnitude of the absorption energy of Ral $\beta$  N29C<sub>SCN</sub> in particular may be due to the fact that calculations have indicated that the nitrile at this location is actually most sensitive to changes in electrostatic field caused by amino acid mutations at the Ral surface, thus screening the effect of mutations made to the Ras surface<sup>53</sup>. This effect is currently being investigated in our laboratory and will be described in a future report.

*K32C<sub>SCN</sub> and N54C<sub>SCN</sub>*: The closest Ral $\beta$  amino acid to positions 30 and 31 is N54, (a distance from backbone atoms of approximately 11-16 Å in our Boltzmann-weighted ensemble of structures). Because MD simulations indicated that this also was oriented above the Rap-Ral $\beta$  surface plane, this probe location was also selected for this study. As shown in Figure 5, the two single mutations E30D and K31E have small but opposite

effects on the absorption energy of N54C<sub>SCN</sub> when compared to Rap E30/K31; Rap E30D is 0.4 cm<sup>-1</sup> higher in absorption energy, while Rap K31E is 0.1 cm<sup>-1</sup> lower in absorption energy. The combined effect of the double mutant Rap E30D/K31E, however, gave a VSE shift 0.2 cm<sup>-1</sup> higher in energy than Rap E30/K31, resulting in an absorption energy that was approximately the sum of the behavior of the two single mutations. As discussed above, because N54C<sub>SCN</sub> visibly disrupted the hydrogen-bonding interaction with Rap K31 in the docked complex, it is likely that the structural disruption of this interface means that this position is not appropriate for deconvoluting structural and electrostatic effects. This observation demonstrates the importance of investigation both structural and electrostatic contributions to the biochemical question of interest.

The nitrile vibrational probe was also placed at Ralβ K32C<sub>SCN</sub> because of its proximity to positions 30 and 31 on the GTPase in the docked complex. As can be seen in Table 4-2: Measured vibrational frequencies of SCN-labeled Ralβ mutants docked with Rap E30/K31; the observed changes in vibrational frequency upon docking each probe to Ras D30/E31 and each Rap mutant and Figure 5, although the measured error in  $\Delta\nu_{obs}$  was larger than other positions, there was a dramatic effect of the mutation Rap K31E on the absorption energy of the thiocyanate compared to Rap E30/K31 (-1.0 cm<sup>-1</sup>), while Rap E30D caused only a small perturbation (-0.2 cm<sup>-1</sup>). The behavior of the double mutant Rap E30D/K31E was the sum of these two shifts, -1.2 cm<sup>-1</sup>. This was very different from the response on binding to Ras D30/E31, which showed a shift in absorption energy of +0.7 cm<sup>-1</sup> compared to Rap E30/K31. Thus, again, the differences in the electrostatic fields between Ras D30/E31 and Rap E30/K31 experienced by the probe at Ralβ K32C<sub>SCN</sub> did not appear to be caused by the reversion mutations at positions 30 and 31 of Rap.

These two probe locations responded to the double mutant Rap E30D/K31E essentially as the addition of each single E30D and K31E mutation. When the probe was at N54C<sub>SCN</sub>, the mutation E30D caused an increase in absorption energy relative to Rap E30/K31, while K31E caused a decrease in absorption energy. The double mutant E30D/K31E shifted approximately halfway between these two extremes. In the case of the probe at location Ralβ K32C<sub>SCN</sub>, both single mutations caused a reduction in the absorption energy of the nitrile probe, while the double mutation caused a shift that was the sum of the two. In these two cases, it appears that we are measuring additive electrostatic effects as would be predicted by a simple model of adding electrostatic fields together. However, in neither case did the double reversion mutation to Rap produce an electrostatic environment measured by the probe that was similar to Ras.

*G28C<sub>SCN</sub> and Y31C<sub>SCN</sub>*: We chose Ralβ G28C<sub>SCN</sub> as a probe location because it was between positions N27 and N29, even though it did not meet other desirable criteria. Previous studies had shown very little sensitivity to differences in binding to Ras D30/E31 versus Rap E30/K31,<sup>2</sup> and a position in the docked complex that was more consistently parallel to the surface than either N27C<sub>SCN</sub> or N29C<sub>SCN</sub><sup>53</sup>. Even still, this probe did respond strongly to mutants Rap E30D and K31E (+1.0 cm<sup>-1</sup> and +0.5 cm<sup>-1</sup> versus Rap E30/K31, respectively), compared to a shift of -0.8 cm<sup>-1</sup> when bound to Ras D30/E31. The double mutant, Rap E30D/K31E, however, only demonstrated a shift of +0.2 cm<sup>-1</sup> compared to Rap E30/K31. The probe at this location thus reacted in a manner in which the two single mutations appear to cancel each other out.

Position Ralβ Y31C<sub>SCN</sub> was selected for study because the change in electrostatic field upon docking to Rap E30/K31 differed from Ras D30/E31 by a large amount, 0.8 cm<sup>-1</sup>, and resulted in changes in field in the opposite direction (sign of  ). As shown in Table 4-2: Measured vibrational frequencies of SCN-labeled Ralβ mutants docked



with Rap E30/K31; the observed changes in vibrational frequency upon docking each probe to Ras D30/E31 and each Rap mutant, both single mutants Rap E30D and Rap K31E shifted the absorption energy of the nitrile probe on Ral $\beta$  Y31C<sub>SCN</sub> to lower energy, by -2.0 cm<sup>-1</sup> and -1.0 cm<sup>-1</sup> respectively. However, the combined effect of the double mutant was to shift the absorption energy lower by only -0.2 cm<sup>-1</sup> compared to Rap E30/K31, substantially less than either single mutant. As shown in our molecular dynamics simulations, there is a significant reorientation of the side chain at position Rap 31 above (azimuthal angle) and away from (polar angle) Ral $\beta$  Y31C<sub>SCN</sub> in the mutants Rap K31 and E30D/K31E which are likely causing this probe to become significantly less sensitive to change in electrostatic field in the binding region in which we are interested.

At both of these locations on the protein surface, Ral $\beta$  G28C<sub>SCN</sub> and Ral $\beta$  Y31C<sub>SCN</sub>, although each single Rap mutant caused a large change in absorption energy (higher in energy at Ral $\beta$  G28C<sub>SCN</sub>, and lower in energy at Ral $\beta$  Y31C<sub>SCN</sub>), the double mutant effectively canceled out those changes in both cases. While MD sampling of position Ral $\beta$  Y31C<sub>SCN</sub> makes it clear that this is likely due to a large structural reorientation of the Rap sidechain position 31, the influences of the probe at position Ral $\beta$  G28C<sub>SCN</sub> are much less clear. After extensive inspection of our MD simulations, we have found no significant structural differences near G28C<sub>SCN</sub> to explain this observation. It could be that the distance between G28 and the region of the binding surface we are investigating by mutations at Rap E30 and K31 make Ral $\beta$  G28C<sub>SCN</sub> an ineffective probe for this study. This result is noteworthy, however, because the distance over which linear VSE effects can be accurately measured has still not been experimentally established.

### 4.3 DISCUSSION

This study was motivated by functional observations that positions 30 and 31 of Ras and Rap, which are among the few chemical differences between the Ras and Rap interfaces, could help discriminate appropriate downstream effectors for each GTPase. Because of the structural similarities of these two protein surfaces, it is possible that downstream effector binding selectivity could be caused by changes in the electrostatic fields at the GTPase-effector interface caused by mutations at position 30 and 31, and in particular by the reversal of charge caused by the K31E mutation. Measuring this effect from the perspective of the downstream effector Ral $\beta$ , which can bind to both GTPases, proved useful to explore this aspect of GTPase function. Previous experimental work and extensive MD sampling provided us with criteria to guide the selection of useful places to position the nitrile probe on the Ral $\beta$  surface. Three criteria that were used were 1) angle of the nitrile with respect to the GTPase-effector interface, with probes perpendicular to the plane of the interface preferred (N27C<sub>SCN</sub> and N29C<sub>SCN</sub>); 2) proximity to positions 30 and 31 when the docked complex is formed (K32C<sub>SCN</sub> and N54C<sub>SCN</sub>), and 3) previously measured discrimination in absorption energy when docking to Ras as opposed to Rap, indicating those probes end up in significantly different electrostatic environments after the docked complex is formed (N27C<sub>SCN</sub>, N29C<sub>SCN</sub>, and Y31C<sub>SCN</sub>). The location G28C<sub>SCN</sub> was chosen because of its position between N27C<sub>SCN</sub> and N29C<sub>SCN</sub>, not because of any favorable selection criteria, and could be considered as a control location on the Ral $\beta$  surface. No probe displayed ideal behavior in all three selection criteria, and some probes were favorable from one aspect but unfavorable from another. For example, N27C<sub>SCN</sub> and N29C<sub>SCN</sub> were the most perpendicular of all of the Ral $\beta$  mutant, but were far removed from Rap positions 30 and 31 in the docked complex. Thus even with detailed structural and electrostatic studies we have carried out on these GTPases, working within

the limitations of the biological system does not allow us to design a perfect experiment. All probes will experience a convolution of favorable and unfavorable interactions, and all probes will be unique. Electrostatic fields must therefore be explored from the perspectives of multiple probes, and information taken from the aggregation of the data.

We observed three general trends from the six probe locations studied here. The first was seen with N27C<sub>SCN</sub> and N29C<sub>SCN</sub>, which displayed the Rap-to-Ras reversion behavior caused by the double mutation Rap E30D/K31E. The absorption energy of this double mutant essentially matched that of Ras D30/E31, and in both cases was caused almost exclusively by the mutation K31E, not E30D. This is direct confirmation of hypotheses proposed from previous crystallographic and docking studies that the difference between Ras and Rap at these two positions leads to the functional discrimination these two GTPases have for different downstream effectors. The second observed trend was that of the double mutant displaying a difference in absorption energy that was simply the sum of the effect caused by the two single mutations. This was seen at positions K32C<sub>SCN</sub> and N54C<sub>SCN</sub>. This is strong experimental evidence that these mutations cause very little disruption in the docked complex, and experience additive changes in electrostatic field that can be measured by VSE spectroscopy of appropriately placed probes. However, it is important to study these effects from as many different probe locations as possible, because the third trend was less easily interpreted. For two probe locations, G28C<sub>SCN</sub> and Y31C<sub>SCN</sub>, while each single Rap mutation caused a large shift in vibrational absorption energy, the double mutant Rap E30D/K31E behaved essentially identically to Rap E30/K31. We have found no structural cause for this behavior. These residues may simply be unimportant for the formation of the Rap-Ral interface, and changes in electrostatic field measured by probes at these locations may not be correlated with changes in Rap-Ral interface formation.

In summary, this study demonstrates that while VSE spectroscopy is a useful tool for examining molecular-level mechanism of electrostatic events in complex biological systems, the convolution of distance, orientation, and change in determining the change in local electrostatic field actually experienced by the nitrile probe needs to be carefully interpreted. Observation of reversion behavior in the double mutant Rap E30D/K31E with the probes Ral $\beta$  N27C<sub>SCN</sub> and N29C<sub>SCN</sub> strongly supports the role these two positions play in supporting an electrostatic mechanism of functional discrimination in GTPase-effector binding. Although these probe locations displayed the reversion mutation behavior that we were predicting when designing these experiments, the convolution of structure and sensitivity to mutations on the surface of Rap demonstrate that unambiguous measurement of electrostatic effects at the Rap-Ral interface will be difficult to achieve. Molecular dynamics sampling appears to be particularly useful in selection of appropriately placed VSE probes. Further experiments on WT and mutant Rap will be guided by extensive MD and continuum electrostatic calculations.

Table 4-1: Dissociation constant of the formation of docked complexes of SCN-labeled Ra $\beta$  mutants with GTPases

GTPase	Ra $\beta$ Mutation K <sub>d</sub> ( $\mu$ M)						
	WT	N27C <sub>SCN</sub>	G28C <sub>SCN</sub>	N29C <sub>SCN</sub>	Y31C <sub>SCN</sub>	K32C <sub>SCN</sub>	N54C <sub>SCN</sub>
Rap E30/K31	0.26 $\pm$ 0.09	0.30 $\pm$ 0.05	0.04 $\pm$ 0.004	0.15 $\pm$ 0.08	1.0 $\pm$ 0.1	0.69 $\pm$ 0.18	0.91 $\pm$ 0.20
Ras D30/E31	1.4 $\pm$ 0.2	7.3 $\pm$ 2.2	4.8 $\pm$ 1.1	3.1 $\pm$ 0.6	12.9 $\pm$ 4.8	6.0 $\pm$ 2.3	4.7 $\pm$ 0.4
Rap E30D	0.24 $\pm$ 0.06	0.12 $\pm$ 0.01	0.03 $\pm$ 0.015	0.10 $\pm$ 0.02	0.89 $\pm$ 0.35	0.34 $\pm$ 0.09	0.79 $\pm$ 0.10
Rap K31E	1.8 $\pm$ 0.7	2.9 $\pm$ 0.7	8.1 $\pm$ 1.5	6.0 $\pm$ 1.1	6.2 $\pm$ 2.0	1.9 $\pm$ 0.8	2.2 $\pm$ 1.5
Rap E30D/K31E	1.0 $\pm$ 0.1	5.4 $\pm$ 1.0	2.5 $\pm$ 0.3	5.0 $\pm$ 1.7	2.0 $\pm$ 0.2	1.8 $\pm$ 0.5	1.1 $\pm$ 0.3

Table 4-2: Measured vibrational frequencies of SCN-labeled Ral $\beta$  mutants docked with Rap E30/K31; the observed changes in vibrational frequency upon docking each probe to Ras D30/E31 and each Rap mutant

		Ral $\beta$ Mutant					
GTPase		N27C	G28C	N29C	Y31C	K32C	N54C
Rap E30/K31	$\tilde{\nu}$ (cm <sup>-1</sup> )	2162.6 (0.4)	2161.8 (0.1)	2160.8 (0.2)	2161.5 (0.2)	2160.9 (0.2)	2161.4 (0.2)
		Difference Compared to Rap E30/K31					
Ras D30/E31	$\Delta\tilde{\nu}$ (cm <sup>-1</sup> )	-0.5 (1.0)	-0.8 (0.1)	0.3 (0.4)	-0.8 (0.4)	0.7 (0.8)	-0.5 (0.4)
	$\Delta F$ (MV/cm)	0.7 (1.4)	1.1 (0.1)	-0.4 (0.6)	1.1 (0.6)	-1.0 (1.1)	0.7 (0.6)
Rap E30D	$\Delta\tilde{\nu}$ (cm <sup>-1</sup> )	-0.1 (0.4)	1.0 (0.5)	0.0 (0.2)	-2.0 (0.4)	-0.2 (0.6)	0.4 (0.3)
	$\Delta F$ (MV/cm)	0.1 (0.6)	-1.4 (0.7)	0.0 (0.3)	2.9 (0.6)	0.3 (0.8)	-0.6 (0.4)
Rap K31E	$\Delta\tilde{\nu}$ (cm <sup>-1</sup> )	-0.6 (0.6)	0.5 (0.2)	0.6 (0.2)	-1.0 (0.4)	-1.0 (0.2)	-0.1 (0.3)
	$\Delta F$ (MV/cm)	0.9 (0.9)	-0.7 (0.3)	-0.9 (0.3)	1.4 (0.6)	1.4 (0.3)	0.1 (0.4)
Rap E30D/K31E	(cm <sup>-1</sup> )	-0.3 (0.5)	0.2 (0.3)	0.8 (0.2)	-0.2 (0.4)	-1.2 (0.4)	0.2 (0.3)
	$\Delta F$ (MV/cm)	0.4 (0.7)	-0.3 (0.4)	-1.1 (0.3)	0.3 (0.6)	1.7 (0.6)	-0.3 (0.4)

## **Chapter 5 Optimizing Electrostatic Field Calculations with the Adaptive Poisson-Boltzmann Solver to Predict Electric Fields at Protein-Protein Interfaces I: Sampling and Focusing**

### **5.1 INTRODUCTION**

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### **5.2 RESULTS**

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### **5.3 DISCUSSION**

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# **Chapter 6 Optimizing Electrostatic Field Calculations with the Adaptive Poisson-Boltzmann Solver to Predict Electric Fields at Protein-Protein Interfaces II: Explicit Near-Probe and Hydrogen Bonding Water Molecules**

## **6.1 INTRODUCTION**

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## **6.2 RESULTS**

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## **6.3 DISCUSSION**

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## **Chapter 7 Electrostatic Fields at Protein-Protein Interfaces: Increased Sampling Time and Various Electrostatic Methods: A Case for Simulating in Polarizable Force Fields**

### **7.1 INTRODUCTION**

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### **7.2 RESULTS**

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### **7.3 DISCUSSION**

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## **Chapter 8 Electrostatic Fields in Small Thiocyanate Molecules with Ensembles Generated using the AMOEBA Force Field**

### **8.1 INTRODUCTION**

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### **8.2 RESULTS**

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### **8.3 DISCUSSION**

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## Appendix

## **Glossary**

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