

# An one dimensional model for naturally curved composite ribbons

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### **CONTENT**

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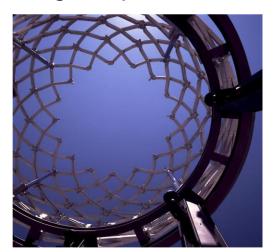
- ☐ Elastic ribbons
- ☐ Applications
- ☐ Motivation
- ☐ Ribbon model
- ☐ Benchmark problems
- ☐ Composite lattices
- Conclusions



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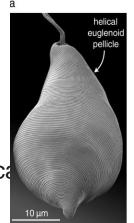
### **ELASTIC RIBBONS**

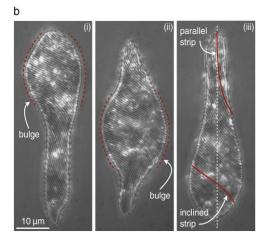
- Ribbons are thin elastic objects whose thickness t, width a and length  $\ell$  are all very different,  $t \ll a \ll \ell$
- Applications in designing lightweight structures that carespond to actuation
- Their flexibility and their thin geometry help turn the relatively small strain produced by actuation into a large-amplitude motion

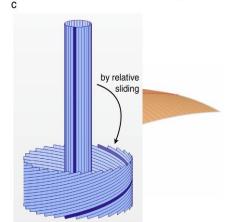




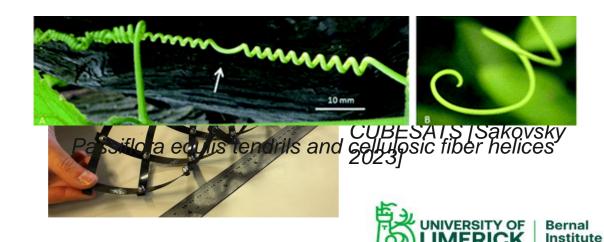
Adaptive architecture [Hobbermann 2012]







The pelliga of auglonids or vides excellent as insufficient insufficie



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### MULTISTABLE LATTICES





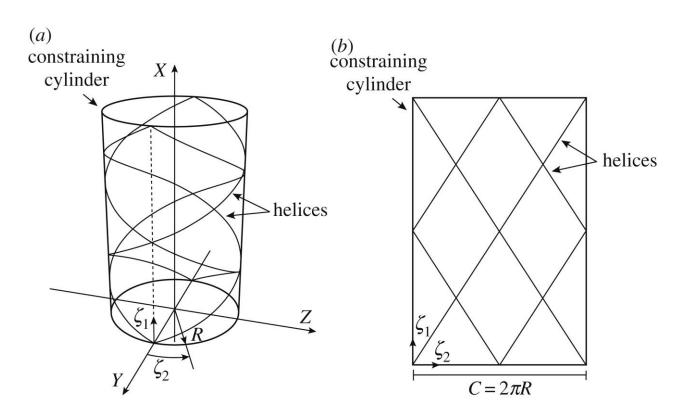


Stable



Pirrera et al. 2013

Unstable



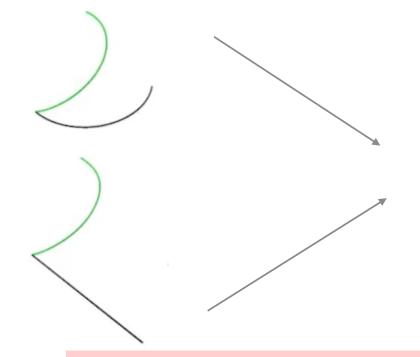
### Behaviour of lattice structures from non-cylindrical geometry



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### GENERAL MODEL

### INDIVIDUAL RIBBON ELEMENTS









Stable shape-1



Stable shape-2



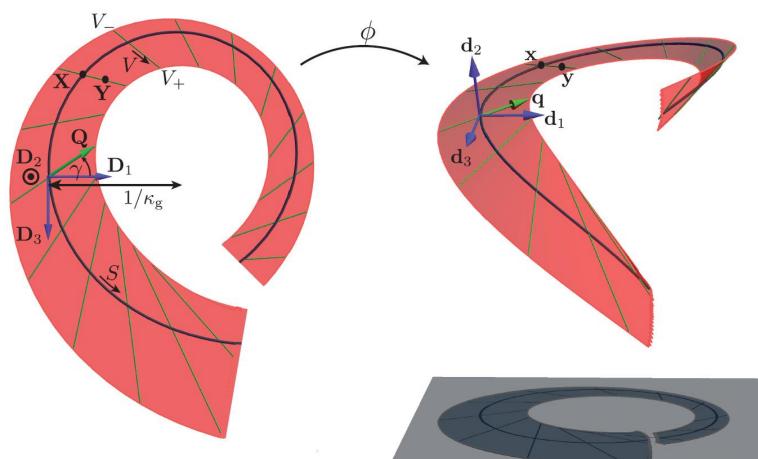
- **Developing a general model** for lattice with arbitrary shapes allowing large deformation extendable to anisotropic materials
- Linking lattices changes the system's energy in a manner that resembles a **series expansion**, allowing us to approximate any continuous energy to desired accuracy
- Alternative geometry -> Higher design space and tailored response

### RIBBON THEORY





Deformed Configuration



Considering ribbon is inextensible

S: Arc length measured along the centerline of undeformed conf

 $D_3 = X'(S)$  is the unit tangent along the centerline

 $m{D_1}$  ,  $m{D_2}$  ,  $m{D_3}$  form an orthonormal frame

 $d_3 = \mathbf{x}'(s)$  is the tangent of the deformed configuration

$$q(\eta, S) = \eta(S)d_3(S) + d_1(S)$$

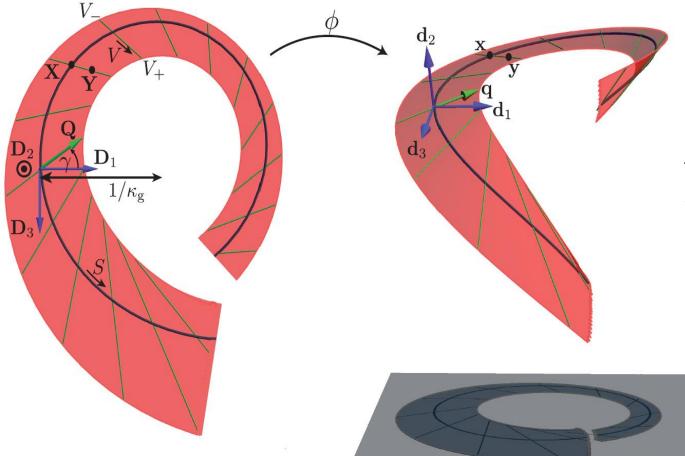


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### RIBBON THEORY

### Undeformed Configuration

Deformed Configuration



Transformation from reference to deformed configuration is expressed by the mapping  $\phi$ :

$$\phi: \mathbf{Y} = \mathbf{X}(S) + V \mathbf{Q}(\eta, S) \mapsto \mathbf{y} = \mathbf{x}(S) + V \mathbf{q}(\eta, S)$$

V is a coordinate along the generatrix

Therefore, **S** and **V** are the longitudinal and transverse coordinate to parameterize ribbon surface.

Strains can be calculated from the Darboux vector:

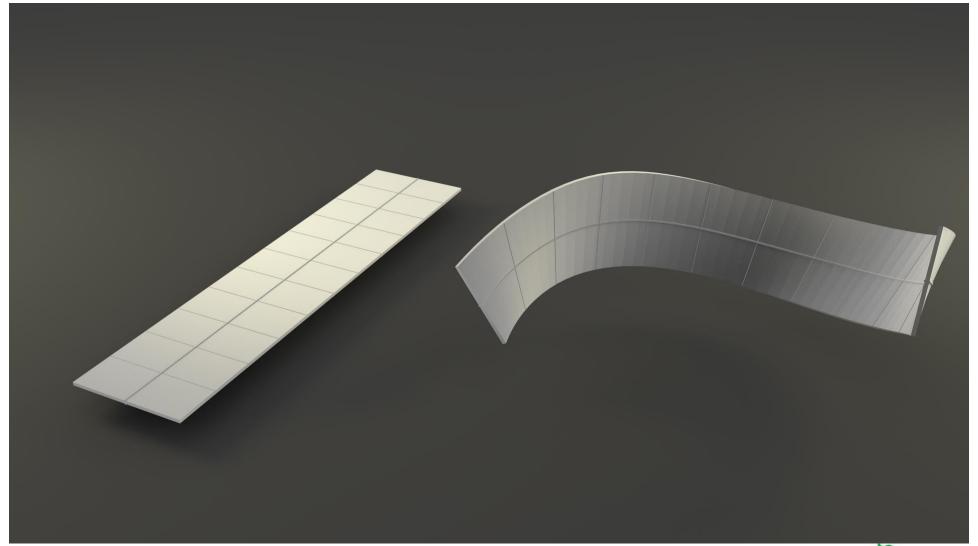
$$\mathbf{d}_i'(S) = \boldsymbol{\omega}(S) \times \mathbf{d}_i(S)$$

 $\omega_1(S), \omega_2(S), \omega_3(S)$ : are the normal, geodesic and the twisting curvatures respectively



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### DEVELOPABLE RECTANGULAR RIBBONS



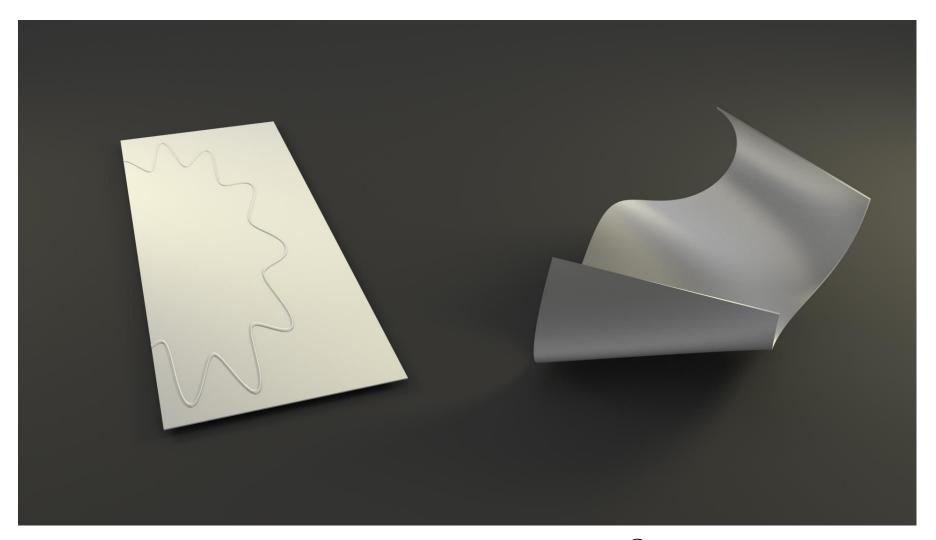
 $V_{\pm}(\eta, S) = \pm \frac{w}{2}$  (rectangular ribbon)
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 $\eta(S) = \tan \gamma = \tan \left( \mathbf{d}_1 \measuredangle \mathbf{q} \right)$ 



## RECTANGULAR RIBBONS WITH ARBITARY CENTERLINE





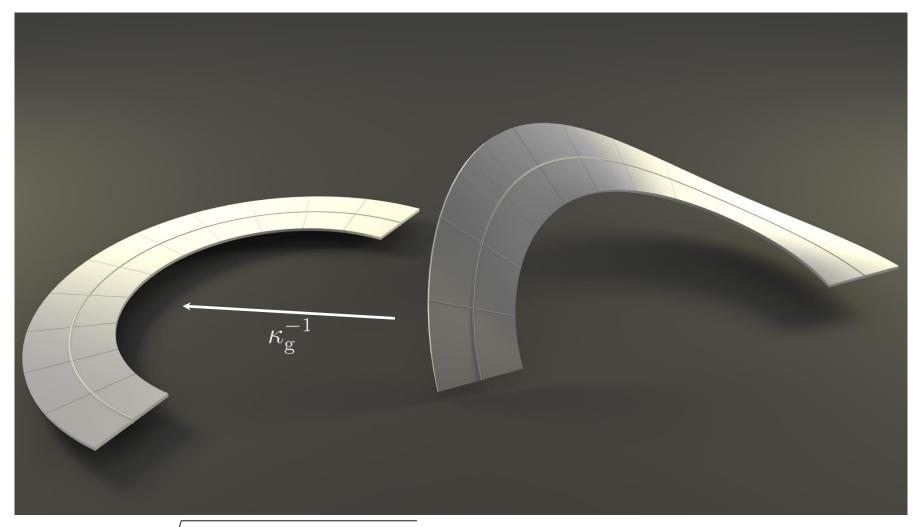
$$V_{\pm}(\eta,S) = ?$$



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### DEVELOPABLE CURVED RIBBONS





$$V_{\pm}(\eta, S) = \frac{1}{\kappa_{\rm g}} \frac{1 - \sqrt{1 \mp (1 + \eta^2) w \kappa_{\rm g} (1 \mp \frac{w \kappa_{\rm g}}{4})}}{1 + \eta^2}$$

(annular ribbon)

$$\eta(S) = \tan \gamma = \tan \left( \mathbf{d}_1 \angle \mathbf{q} \right)$$

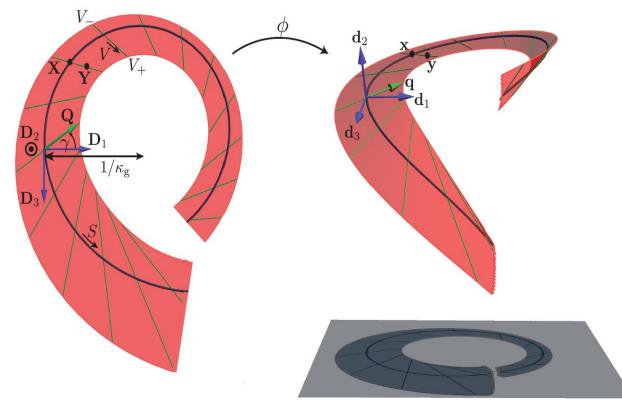


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### RIBBON THEORY

### Undeformed Configuration

### Deformed Configuration



### Rectangular Ribbon

$$V_{\pm}(\eta, S) = \pm \frac{w}{2}$$

### Annular Ribbon

$$V_{\pm}(\eta, S) = \frac{1}{\kappa_{\rm g}} \frac{1 - \sqrt{1 \mp (1 + \eta^2) w \kappa_{\rm g} (1 \mp \frac{w \kappa_{\rm g}}{4})}}{1 + \eta^2}$$

General Case: with geodesic curvature  $\kappa_g$ 

$$V_{\rm c}(\eta, \eta', S) = \frac{1}{(1+\eta^2)\kappa_{\rm g}(S) - \eta'}$$

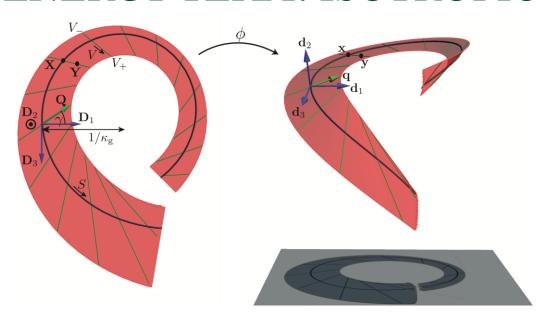
$$\left| \frac{V}{V_{\rm c}} \right| < 1$$
  $\frac{1 - V_{+}/V_{\rm c}}{1 - V_{-}/V_{\rm c}} > 0$ 



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### **ENERGY TERM: ISOTROPIC CASE**





$$\mathbf{K}(\eta, \eta', \omega_1, S, V) = \frac{\mathbf{K}_0(\eta, \omega_1)}{1 - \frac{V}{V_c(\eta, \eta', S)}}$$

$$\mathbf{K}_0(\eta, \omega_1) = -\omega_1 \begin{pmatrix} 1 & -\eta \\ -\eta & \eta^2 \end{pmatrix}_{(\mathbf{d_3}, \mathbf{d_1})}$$

$$da = |\partial_S \mathbf{y} \times \partial_V \mathbf{y}| \, dS \, dV = \left(1 - \frac{V}{V_c(\eta, S)}\right) dS \, dV$$

$$E(\eta, \eta', \omega_1) = \frac{D}{2} \int_0^L \left[ \left( -V_c \ln \frac{1 - V_+ / V_c}{1 - V_- / V_c} \right) \text{tr}(\mathbf{K}_0^2) - 2(V_+ - V_-) \mathbf{Q}_r : \mathbf{K}_0 \right] dS$$

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### **COMPOSITE RIBBONS**

### Considering inextensibility

$$U = \frac{1}{2} \boldsymbol{\kappa}^T \boldsymbol{D} \boldsymbol{\kappa}$$

$$D=D_{11}egin{bmatrix} 1 & 
u & \psi \\ 
\nu & eta & \phi \\ 
\psi & \phi & 
ho \end{bmatrix}$$
 Anisotropy with fully populated bending stiffness

$$U = \int \int \frac{1}{2} D_{11} \left( \kappa_{11}^2 + \beta \kappa_{22}^2 + 4\kappa_{12} \left( \kappa_{12} \rho + \kappa_{22} \phi \right) + 2\kappa_{11} \left( \kappa_{22} \nu + 2\kappa_{12} \psi \right) \right) da$$

### Curvature tensor with developability constraint

$$\mathbf{K}(\eta, \eta', \omega_1, S, V) = \frac{\mathbf{K}_0(\eta, \omega_1)}{1 - \frac{V}{V_c(\eta, \eta', S)}}$$

$$\mathbf{K}(\eta, \eta', \omega_1, S, V) = \frac{\mathbf{K}_0(\eta, \omega_1)}{1 - \frac{V}{V_0(\eta, \eta', S)}} \qquad \mathbf{K}_0(\eta, \omega_1) = -\omega_1 \left( \mathbf{d}_3 \otimes \mathbf{d}_3 - \eta \left( \mathbf{d}_3 \otimes \mathbf{d}_1 + \mathbf{d}_1 \otimes \mathbf{d}_3 \right) + \eta^2 \mathbf{d}_1 \otimes \mathbf{d}_1 \right)$$

$$U = \int \int \frac{1}{2} D_{11} \frac{\omega_1^2}{\left(1 - \frac{V}{V_c}\right)^2} \left(1 + \beta \eta^4 + 2\eta^2 \left(\nu + 2\rho\right) - 4\eta^3 \phi - 4\eta \psi\right) da$$

$$V_{\rm c}(\eta, \eta', S) = \frac{1}{(1 + \eta^2) \, \kappa_{\rm g}(S) - \eta'}$$



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### TOTAL ENERGY: COMPOSITE RIBBONS

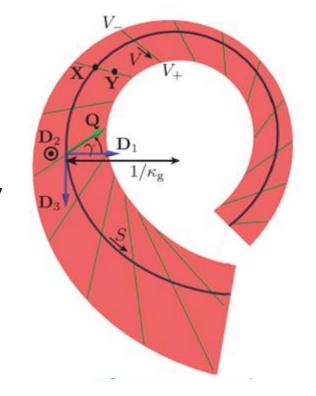
$$U = \int \int \frac{1}{2} D_{11} \frac{\omega_1^2}{\left(1 - \frac{V}{V_c}\right)^2} \left(1 + \beta \eta^4 + 2\eta^2 \left(\nu + 2\rho\right) - 4\eta^3 \phi - 4\eta \psi\right) da$$

$$da = |\partial_S \mathbf{y} \times \partial_V \mathbf{y}| dS dV = \left(1 - \frac{V}{V_c(\eta, S)}\right) dS dV,$$

$$V_{\rm c}(\eta, \eta', S) = \frac{1}{(1 + \eta^2) \, \kappa_{\rm g}(S) - \eta'}.$$

$$U = \int_{0}^{L} \frac{1}{2} D_{11} \left( \kappa_{11}^{2} + \beta \kappa_{22}^{2} + 4 \kappa_{12} \left( \kappa_{12} \rho + \kappa_{22} \phi \right) + 2 \kappa_{11} \left( \kappa_{22} \nu + 2 \kappa_{12} \psi \right) \right) \log \frac{V^{-} - V_{c}}{V^{+} - V_{c}} \mathrm{d}S$$

Using virtual work principle, the weak form can be derived for solving numerically



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# **EQUATIONS OF EQUILIBIRUM**



Principle of the virtual work:

$$\mathcal{W}_{\eta} + \mathcal{W}_{\mathbf{i}}^{\text{rod}} + \mathcal{W}_{\mathbf{e}} + \mathcal{W}_{\text{cEB}} = 0$$

$$\mathcal{W}_{\text{cEB}} = -\int_{0}^{L} \mathbf{N} \cdot (\hat{\mathbf{x}}' - \hat{\mathbf{d}}_{3}) \, \mathrm{d}S$$

$$\mathcal{W}_{e} = \int_{0}^{L} (\mathbf{p} \cdot \hat{\mathbf{x}} + \mathbf{c} \cdot \hat{\boldsymbol{\psi}}) \, \mathrm{d}S$$

$$\mathcal{C}_{g}(\omega_{2}, S) = \kappa_{g}(S) - \omega_{2} = 0$$

$$\mathcal{C}_{d}(\omega_{1}, \omega_{3}, \eta) = \eta \, \omega_{1} - \omega_{3} = 0$$

$$\mathcal{C}_{d}(\omega_{1}, \omega_{3}, \eta) = \eta \, \omega_{1} - \omega_{3} = 0$$

$$\mathcal{W}_{\eta} = -\int_{0}^{L} \left[ \frac{\partial \mathcal{E}}{\partial v'} \, \hat{\eta}' + \left( \frac{\partial \mathcal{E}}{\partial v} - \lambda_{d} \, \frac{\partial \mathcal{C}_{d}}{\partial v} \right) \, \hat{\eta} \right] \, \mathrm{d}S$$
Balance equations:
$$\mathcal{W}_{\eta} = -\int_{0}^{L} \left[ \frac{\partial \mathcal{E}}{\partial v'} \, \hat{\eta}' + \left( \frac{\partial \mathcal{E}}{\partial v} - \lambda_{d} \, \frac{\partial \mathcal{C}_{d}}{\partial v} \right) \, \hat{\eta} \right] \, \mathrm{d}S$$

$$\mathbf{N}'(S) + \mathbf{p}(S) = \mathbf{0},$$

$$\mathbf{M}'(S) + \mathbf{x}'(S) \times \mathbf{N}(S) + \mathbf{c}(S) = \mathbf{0}$$

Constitutive law:

$$\mathbf{M} = \sum_{i=1}^{3} \left( \frac{\partial \mathcal{E}}{\partial \omega_{i}} - \lambda_{d} \frac{\partial \mathcal{C}_{d}}{\partial \omega_{i}} - \lambda_{g} \frac{\partial \mathcal{C}_{g}}{\partial \omega_{i}} \right) \mathbf{d}_{i}$$



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# LIMITING CASES

### Sadowsky's Limit: Narrow Rectangular Ribbons

$$E_S = \frac{Dw}{2} \int_0^L \omega_1^2 (1 + \eta^2)^2 \, dS$$

Sadowsky, M, Sitzungsber. Preuss. Akad. Wiss. 22, 412–415 (1930).

### Wunderlich's Limit: Finite width Rectangular Ribbons

$$E_{W} = \frac{Dw}{2} \int_{0}^{L} \omega_{1}^{2} (1 + \eta^{2})^{2} \frac{1}{\eta' w} \ln \left( \frac{1 + \eta' w/2}{1 - \eta' w/2} \right) dS$$

Wunderlich, W.: Über ein abwickelbares Möbiusband. Monatshefte für Mathematik 66(3), 266–289 (1962). Starostin and van der Heijden, PRL 101 (2008).

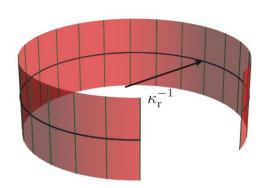
### Audoly & Seffen's Limit: Cylindrical Ribbon

$$E = \frac{Dw}{2} \int_0^L \left[ \omega_1^2 (1 + \eta^2)^2 - 2\kappa_r \omega_1 (1 + \nu \eta^2) \right] dS$$

Basile Audoly and Keith A. Seffen, Journal of Elasticity (2015)



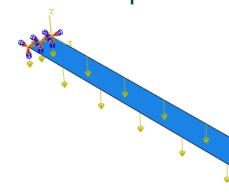
E. L. Starostin et al., Nature Materials (2007)





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# BENDING $[O_{4}]$



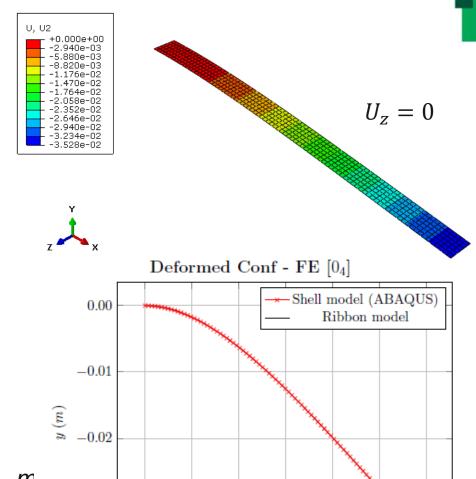
$$D = D_{11} \begin{bmatrix} 1 & \nu & \psi \\ \nu & \beta & \phi \\ \psi & \phi & \rho \end{bmatrix}$$

G<sub>12</sub>: 5.2 GPa

$$D = D_{11} \begin{pmatrix} 1. & 0.0179 \\ 0.0179 & 0.0598 \\ 0. & 0. \end{pmatrix}$$

$$E_{11}$$
: 157 GPa  $D = D_{11} \begin{pmatrix} 1. & 0.0179 & 0. \\ 0.0179 & 0.0598 & 0. \\ 0. & 0. & 0.0329 \end{pmatrix}$  GPa – m

$$D_{11} = 0.05387, \beta = 0.0598, \nu = 0.0179,$$
  
 $\psi = 0, \ \phi = 0, \ \rho = 0.0331$ 



-0.03

0.00

0.05

0.10

0.15

x(m)



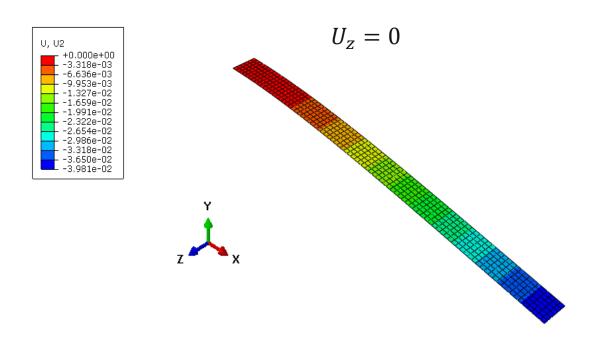
0.30

0.25

0.20

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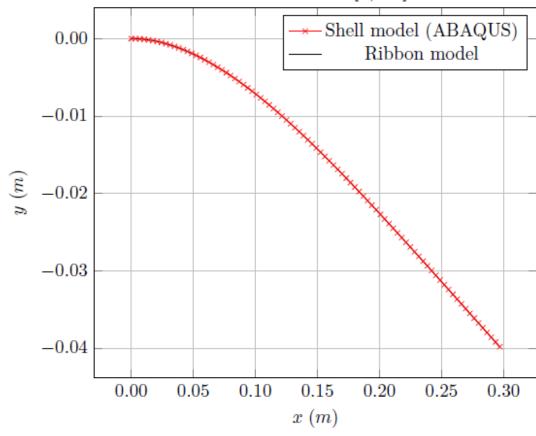
# BENDING [0/90/90/0]



$$D = D_{11} \begin{pmatrix} 1 & 0.0203 & 0.0 \\ 0.0203 & 0.2010 & 0.0 \\ 0.0 & 0.0 & 0.0373 \end{pmatrix} GPa - m^3$$

$$D_{11} = 0.0475, \beta = 0.2010, \nu = 0.0203,$$
  
 $\psi = 0, \ \phi = 0, \ \rho = 0.0373$ 

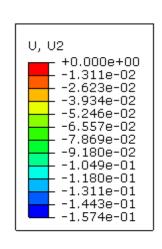
### Deformed Conf - FE $[0/90]_S$

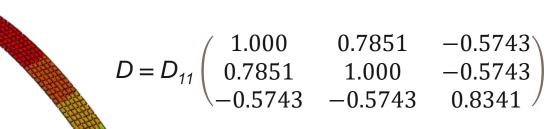




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# BENDING [-45/45/45/-45]

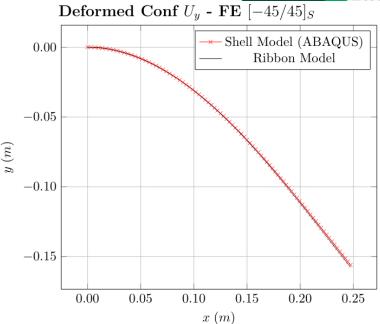


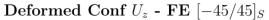


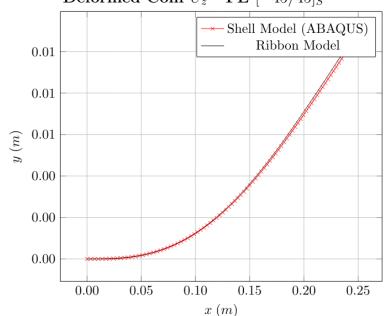


$$D_{11} = 0.0165, \beta = 1, \nu = 0.7851,$$
  
 $\psi = -0.5743, \quad \phi = -0.5743,$   
 $\rho = 2.1762$ 

Captures well the Bending-Twisting coupling of angle plies

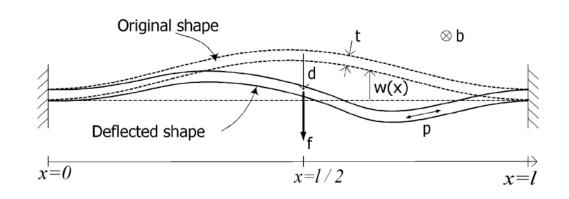


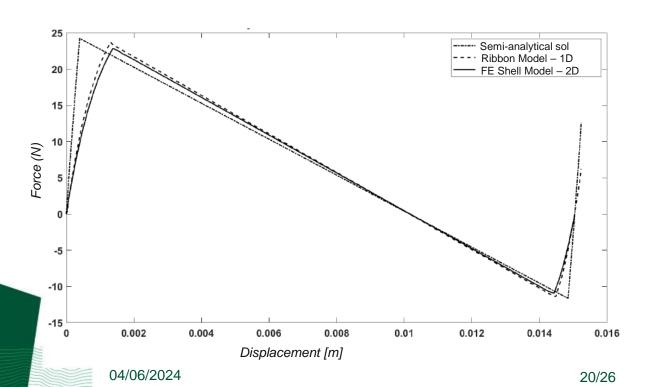


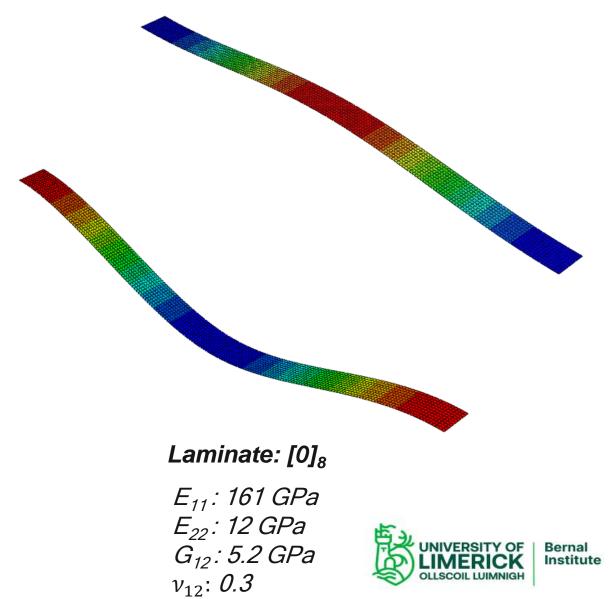


### SNAP-THROUGH OF COMPOSITE RIBBON

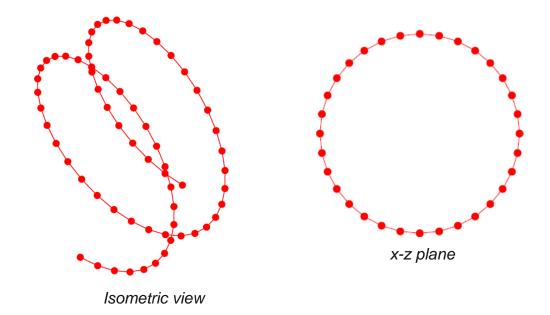


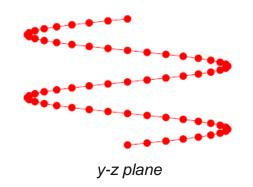


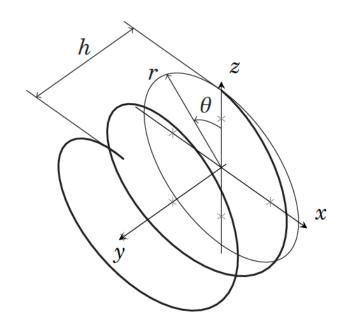




### CYLINDRICAL LATTICES





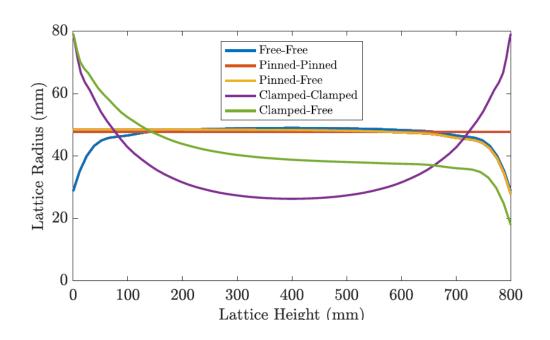


- Length = 1000 mm, r =79.2 mm, h= 100mm
- Discretization: 32 nodes.
- The helices are hinged at intersection points
- At intersection, relative angle between helices can change

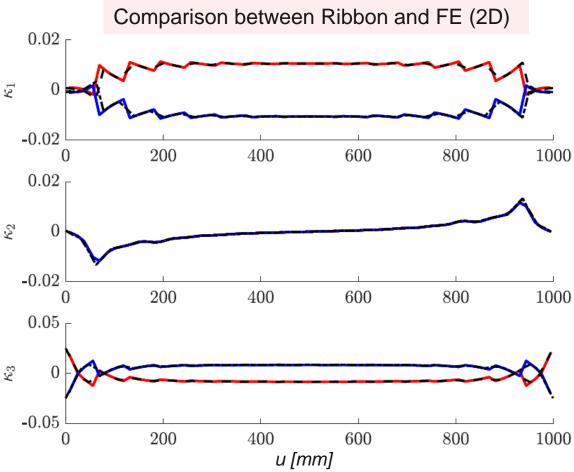


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### EFFECT OF BOUNDARY CONDITIONS



 $\kappa_1, \kappa_2$  and  $\kappa_3$  are the two normal and twisting curvatures respectively



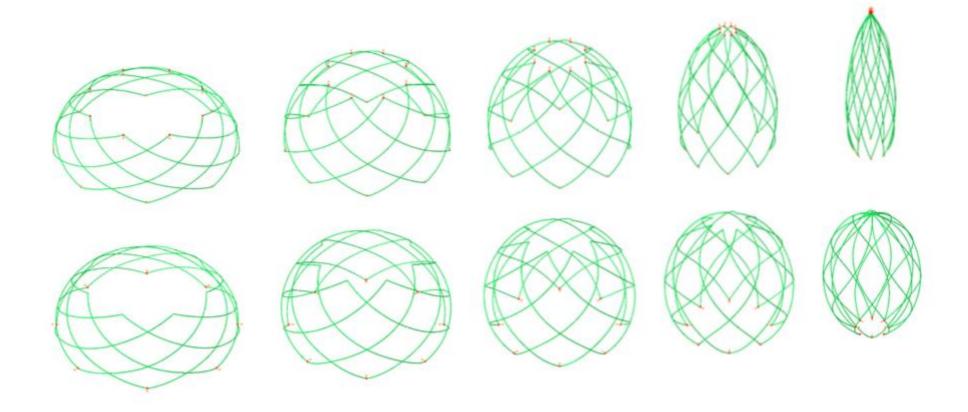
Curvature profile for the lattice at an axial extension of 700mm subjected to clamped-clamped BCs
Solid red and blue lines represent positive and negative handed helices

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# Future Work: Non-cylindrical lattices





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### **CONCLUSIONS**

- A new model for anisotropic ribbons with reduced dimensionality has been developed.
- The model captures geometric nonlinearities associated with large deformations and rotations in anisotropic materials well.
- The formulation relies on two variables:  $\eta$  (related to vector defining the generatrix) and  $\omega_1$  (normal curvature).  $\eta$  is resolved after integration across the width.
- Various test problems were solved to demonstrate the model's ability to capture the anisotropic behaviour of composites
- Cylindrical lattices exhibiting non-uniform curvatures due to boundary effects were modelled, demonstrating the effectiveness of the new formulation.
- Future works will involve investigating general non-linear lattices having applications in deployable structures for satellites and spacecrafts.



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