Unit V - Wave Mechanics

Q.1 Write three forms of Heisenberg's Uncertainty principle and give any two examples of its application.

It is **impossible to measure precisely** and **simultaneously** both the **pair of** certain **canonical conjugate variables** that describe the behavior of an atomic system.

The product can never be smaller than the number of the order of $\hbar/2$ ($\hbar = \hbar/2\pi$), **h (Planck's constant)=** 6.6 X 10⁻³⁴ J-sec

- 1. $\Delta P. \Delta X \ge \hbar/2$ or change in **momentum x Displacement** $\ge \hbar/2$
- 2. $\Delta E. \Delta t \ge \hbar/2$ or change in **Energy x** change in **time** $\ge \hbar/2$
- 3. $\Delta J. \Delta \phi \ge \hbar/2$ or change in **Angular momentum x** Change in **Angle** $\ge \hbar/2$

Applications: 1. Non existence of electron in nucleus (or electron cannot be found in nucleus)

- **2.** Finite width of spectral lines.
- **3.** The radius of Bohr's first Orbit.
- 4. The binding energy of an electron in an atom
- 5. Zero point energy of a Harmonic Oscillator.

Q.2 Derive the Schrödinger's Time independent (TI) and Time dependent (TD) wave equations.

Schrödinger wave equations [(Time independent-**TI** and Time dependent-**TD**) enable us to establish the De-Broglie wave function (ψ) of a wave associated with a moving particle.

(Time independent-TI)

The general differential equation of motion is given by $\frac{\partial^2 \psi}{\partial x^2} + \frac{\partial^2 \psi}{\partial y^2} + \frac{\partial^2 \psi}{\partial z^2} = \frac{1}{v^2} \frac{\partial^2 \psi}{\partial t^2}$ where ψ is function of x, y, z and t. v is velocity of the particle in medium.

or
$$\left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2}\right)\psi = \frac{1}{v^2}\frac{\partial^2\psi}{\partial t^2}$$

since
$$\nabla$$
(Del Operator) = $\left(i\frac{\partial}{\partial x} + j\frac{\partial}{\partial y} + k\frac{\partial}{\partial z}\right)$

$$therefore \ \pmb{\nabla}^2(\text{Laplacian Operator}) = \nabla \cdot \nabla = \left(i\frac{\partial}{\partial x} + j\frac{\partial}{\partial y} + k\frac{\partial}{\partial z}\right) \cdot \left(i\frac{\partial}{\partial x} + j\frac{\partial}{\partial y} + k\frac{\partial}{\partial z}\right) \Longrightarrow \left(\frac{\pmb{\partial}^2}{\pmb{\partial} x^2} + \frac{\pmb{\partial}^2}{\pmb{\partial} y^2} + \frac{\pmb{\partial}^2}{\pmb{\partial} z^2}\right)$$

or
$$\nabla^2 \psi = \frac{1}{v^2} \frac{\partial^2 \psi}{\partial t^2} \dots \dots [1]$$

Let the solution of above differential equation be $\pmb{\psi} = \pmb{\psi}_0 \pmb{e}^{-i\omega t}$

Double differentiating it w.r.t. t

$$\frac{\partial \psi}{\partial t} = -i\omega \, \psi_0 e^{-i\omega t}$$

$$\frac{\partial \psi}{\partial t^2} = i^2 \omega^2 \psi$$

$$\frac{\partial \psi}{\partial t^2} = -\omega^2 \psi \qquad \dots \dots [2]$$

Substituting [2] into [1] gives us

$$\nabla^2 \psi = -\frac{\omega^2}{v^2} \, \psi$$

$$\nabla^2 \psi + \frac{\omega^2}{v^2} \, \psi = 0$$

$$\nabla^2 \psi + \frac{4\pi^2 \mathbf{v}^2}{v^2} \psi = 0$$

$$\nabla^2 \psi + \frac{4\pi^2 \mathbf{v}^2}{\mathbf{v}^2 \lambda^2} \psi = 0 \quad \dots \dots [3]$$

 $\omega = 2\pi v$ where \mathbf{v} is frequency of wave

 $v = v\lambda$ where v is velocity if wave

According to **De-Broglie hypothesis**, a wave is associated of wavelength λ with any moving particle. The relation is $\lambda = h/p$ (where h = Planck's constant (6.6X10⁻³⁴ joule-sec) and p = momentum of moving particle with velocity v (p = mv). If we introduce $\lambda = h/p$ into [3] then wave function equation ψ become the wave function ψ of a **De-Broglie wave** associated with moving particle.

$$\nabla^2 \psi + \frac{4\pi^2 p^2}{h^2} \psi = 0$$

[momentum (p) = mv]

$$\nabla^2 \psi + \frac{4\pi^2 m^2 v^2}{h^2} \ \psi \ = 0$$

[since
$$\hbar = \frac{h}{2\pi}$$
 therfore $\hbar^2 = \frac{h^2}{4\pi^2}$]

$$\nabla^2 \psi + \frac{m^2 v^2}{\left(\frac{h^2}{4\pi^2}\right)} \psi = 0 \Longrightarrow \nabla^2 \psi + \frac{m^2 v^2}{\hbar^2} \psi = 0 \quad \dots [4]$$

Since total energy of particle is $E = \frac{1}{2}mv^2 + V$

$$or \frac{1}{2}mv^2 = 2(E - V)$$

$$mv^2 = 2(E - V)$$

$$m^2v^2 = 2m(E - V) \dots [5]$$

Substituting equation [5] into [4] we get

$$\nabla^2 \psi + \frac{2m}{\hbar^2} (E - V) \psi = 0$$

This equation is known as Schrödinger's time-independent (TI) equation because the time derivative $(\frac{\partial \psi}{\partial t^2})$ in equation [1] was eliminated by substituting its value from equation [2].

(Time-dependent-TD):

To find the time dependent-equation (TD) from time-independent (TI) [i.e. $\nabla^2 \psi + \frac{2m}{\hbar^2} (E - V) \psi = 0$, we substitute the value of E ψ again into time independent-equation. The value E ψ is obtained from first time derivative of solution of ψ (i.e. $\psi = \psi_0 e^{-i\omega t}$).

Differentiating it w.r.t. t

$$\begin{split} \frac{\partial \psi}{\partial t} &= -i\omega \ \psi_0 e^{-i\omega t} \\ \frac{\partial \psi}{\partial t} &= -i \omega \ \psi_0 e^{-i\omega t} \\ \frac{\partial \psi}{\partial t} &= -\frac{i^2}{i} \frac{2\pi E}{h} \psi \\ \frac{\partial \psi}{\partial t} &= -\frac{i}{i} \frac{E}{\left(\frac{h}{2\pi}\right)} \psi \\ \frac{\partial \psi}{\partial t} &= \frac{E}{i} \frac{\psi}{\left(\frac{h}{2\pi}\right)} \\ \frac{\partial \psi}{\partial t} &= \frac{E}{i} \frac{\psi}{h} \quad \left[\text{since } h = \frac{h}{2\pi} \right] \\ E\psi &= i \hbar \frac{\partial \psi}{\partial t} \qquad \dots \dots [1] \end{split}$$

Here **E** equivalent to $i\hbar \frac{\partial}{\partial t}$. Therefore, $i\hbar \frac{\partial}{\partial t}$ is known as energy operator.

The same value of Eψ can also be found out from Schrödinger's time-independent (TI) equation as

since
$$\nabla^2 \psi + \frac{2m}{\hbar^2} (E - V) \psi = 0$$

$$\frac{2m}{\hbar^2}(E-V)\,\psi = -\nabla^2\psi$$

$$E\psi - V\psi = -\nabla^2 \psi \frac{\hbar^2}{2m}$$

$$E\psi = -\nabla^2\psi \frac{\hbar^2}{2m} + V\psi$$

$$E\psi = \left(-\nabla^2 \frac{\hbar^2}{2m} + V\right)\psi \quad \dots \dots [2]$$

Here the quantity $\left(-\nabla^2\frac{\hbar^2}{2m}+V\right)$ is known as **Hamiltonian operator (H).** It can be total energy operator for conservative force field and for non-conservative force field it is not the total energy operator.

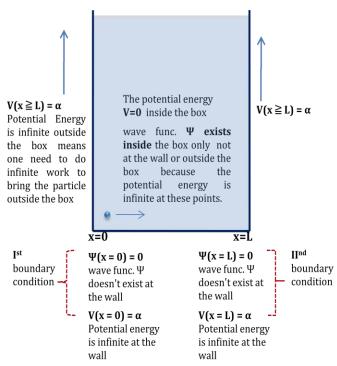
The equation [1] and [2] can be written as $E\psi=i\hbar\frac{\partial\psi}{\partial t}=\left(-\nabla^2\frac{\hbar^2}{2m}+V\right)\psi$ or

$$i\hbar\frac{\partial\psi}{\partial t} = \left(-\nabla^2\frac{\hbar^2}{2m} + V\right)\psi$$

This equation is known as Schrödinger's time-dependent (TD) equation.

Q.3 Establish the wave function ψ of a particle confined in 1-D (one dimensional) box of infinite potential barrier height.

Q. Discuss the Application of Schrödinger wave equation or Establish the wave function ψ of a Particle confined in 1-D (one dimensional) box of infinite potential barrier height.



Since the TI Schrodinger equation:

$$\nabla^2 \psi + \frac{2m}{\hbar^2} (E - V) \psi = 0$$

The wave function (ψ) exist only inside the box where V=0. Therefore the above expression reduces to

$$\nabla^2 \psi + \frac{2m}{\hbar^2} E \psi = 0 \quad \dots [1]$$

$$Let \frac{2mE}{\hbar^2} = k^2 \qquad \dots \dots [2]$$

$$\nabla^2 \psi + k^2 \psi = 0$$

Let the solution of above differential equation be

$$\psi_n = \mathbf{A} \sin kx + \mathbf{B} \cos kx$$
[3] Here values of A and B is to be determined.

To find the values of constants, we apply I^{st} boundary condition (i.e. $\psi(x=0) = 0$ means ψ doesn't exist at x=0 or at the wall)

$$0 = \mathbf{A}\sin k. \, 0 + \mathbf{B}\cos k. \, 0$$
$$\mathbf{or} \, \mathbf{B} = 0$$

Substituting it in eq. [3] we get

$$\psi_n = \mathbf{A} \sin kx \dots [4]$$

Now, to find the value of constant A, we apply II^{nd} boundary condition (i.e. $\psi(x=L) = 0$ means ψ[4] doesn't exist at x=L or at the wall). Subs. x=L and $\psi=0$ in eq. [4],

 $0 = A \sin kL$, we find that A \neq 0 then Sin kL =0

 $\sin kL = 0$

 $kL = \pm n\pi$

$$k = \frac{n\pi}{I}$$
 or $k^2 = \frac{n^2\pi^2}{I^2}$ [5]

Since eq. [2] = [5] therefore, $k^2 = \frac{n^2 \pi^2}{L^2} = \frac{2mE}{\hbar^2}$

$$E = \frac{n^2 \pi^2 \hbar^2}{2mL^2} \rightarrow \frac{n^2 \pi^2 h^2}{4.\pi^2.2mL^2} \rightarrow \frac{n^2 \pi^2 h^2}{4.\pi^2.2mL^2} \rightarrow \frac{n^2 h^2}{8mL^2}$$

$$E=\frac{n^2h^2}{8mL^2}$$

Since *n* is an integer or a variable value in above expression and other parameter are constants therefore, we infer that energy of a particle in 1D box is quantized or has fixed values.

$$E_n = \frac{n^2 h^2}{8mL^2} \to E_1 = \frac{1^2 h^2}{8mL^2} = \frac{h^2}{8mL^2} \to E_1$$

$$E_2 = \frac{2^2 h^2}{8mL^2} \to 4E_1$$

$$E_3 = \frac{3^2 h^2}{8mL^2} \to 9E_1$$

As we found above that on applying the IInd boundary condition we won't be able to determine the value of constant A, therefore we use a mathematical tool called **normalization condition**.

$$\int_{-\infty}^{+\infty} \psi_n \cdot \psi_n^* dx = 1$$
, here ψ_n^* is complex conjugate of function ψ_n . Therefore, $\int_{-\infty}^{+\infty} |\psi_n|^2 dx = 1$

It is to mention that $|\psi_n|^2$ represents the probability density of finding the particle in length L.

Since $\psi_n = A \sin kx$ eq. [4] therefore,

$$\int_{-\infty}^{+\infty} |\psi_n|^2 dx = \int_{0}^{L} |\psi_n|^2 dx = A^2 \int_{0}^{L} \sin^2 \left(\frac{n\pi x}{L}\right) dx = 1$$

$$\frac{A^2}{2} \left[\int_0^L dx - \int_0^L \cos\left(\frac{2n\pi x}{L}\right) dx \right] = 1$$

$$\frac{A^2}{2} \left[x - \left(\frac{L}{2n\pi} \right) \sin \frac{2n\pi x}{L} \right]_0^L = 1$$

$$\frac{A^2}{2}L = 1$$
 or $A = \sqrt{\frac{2}{L}} = 1$

$$\psi_n = \sqrt{\frac{2}{L}} \sin kx$$

This is the required wave function of particle confined in 1D box of infinite potential barrier height.

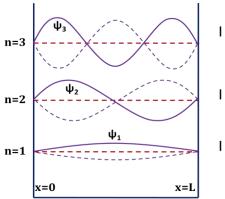


Fig1 represents the waveforms at n=1,n=2 and n=3

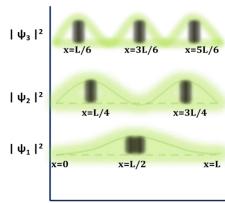
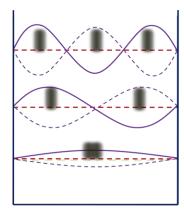


Fig2 Black shaded portion represents the probability density of finding the particle in length L



combined figure showing waveform and position of finding the particle (black shaded portion)

Try numerical on Energy $E_n=rac{n^2h^2}{8mL^2}$ and wave function the particle confined in 1D box $\,\psi_n=$

 $\sqrt{\frac{L}{2}}\sin kx.$