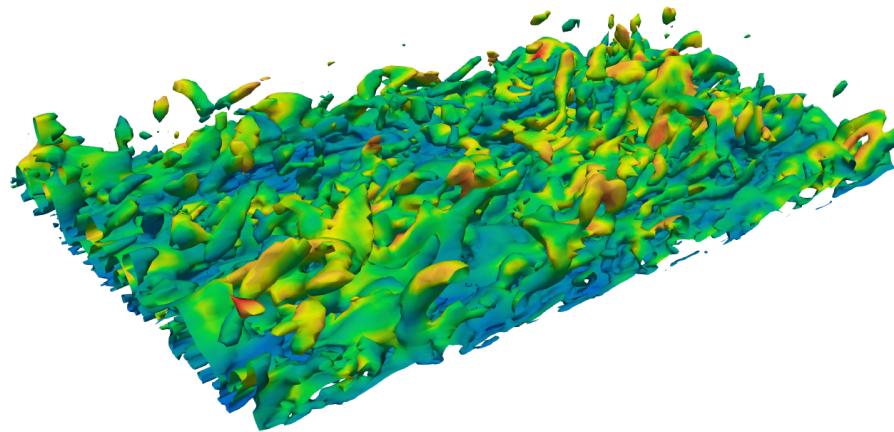


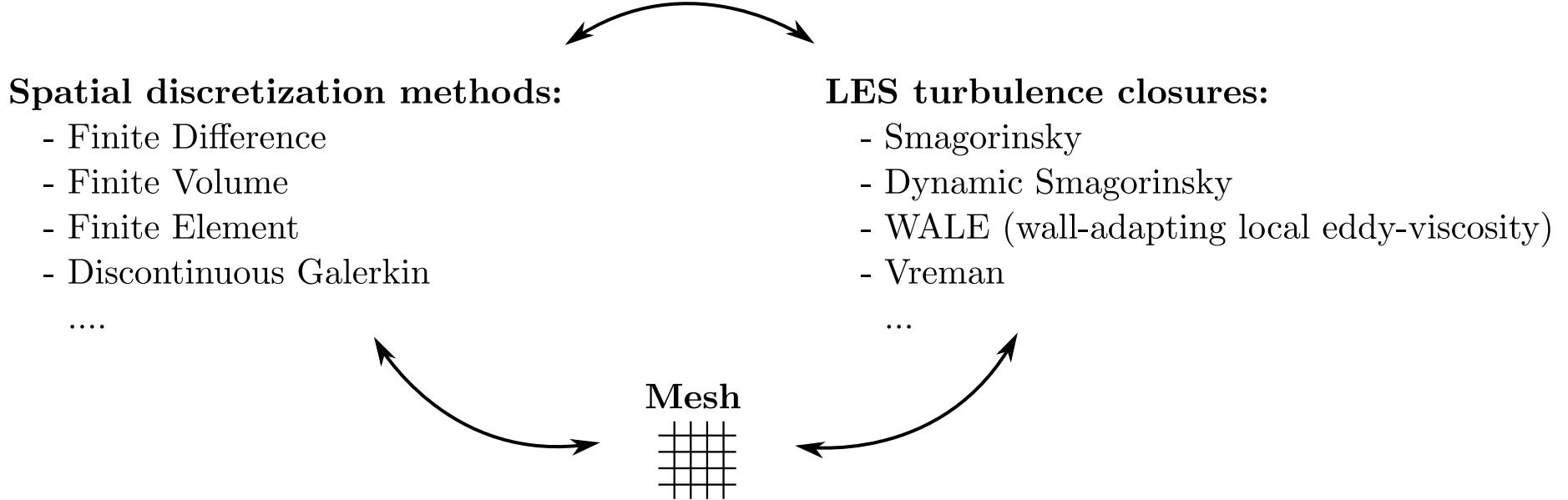
Progress in high-order large-eddy simulation of aeronautical flows using the integral-length scale approximation turbulence model

B. Font*, S. Pezzano, F. Naddei, M. Montagnac, O. Lehmkuhl



HiFiLeD Symposium 2022

Discretization methods, turbulence models, and mesh types, affect convergence



- There is no obvious choice to match discretization method and turbulence model.
- Turbulence models usually are case- (and mesh-) dependent too.

$$\tau_{ij}^a = -2\nu_{\text{SFS}} \bar{S}_{ij}$$

$$\phi = \langle \phi \rangle + \phi'$$

The integral length scale approximation (ILSA) model

The filter width Δ depends on an estimated integral scale, and not on the local cell size.
The integral scale is calculated with temporal-averaged quantities.

$$\nu_{\text{SFS}} = (C_k L_{\text{est}})^2 |\bar{S}| \quad \Delta = C_\Delta L_{\text{est}} \quad L_{\text{est}} = \frac{\langle \mathcal{K}_{\text{res}} \rangle^{3/2}}{\langle \varepsilon_{\text{tot}} \rangle}$$

C_k is determined via s_τ , ILSA's single parameter which encapsulates the amount of turbulent kinetic energy aimed to be modelled.

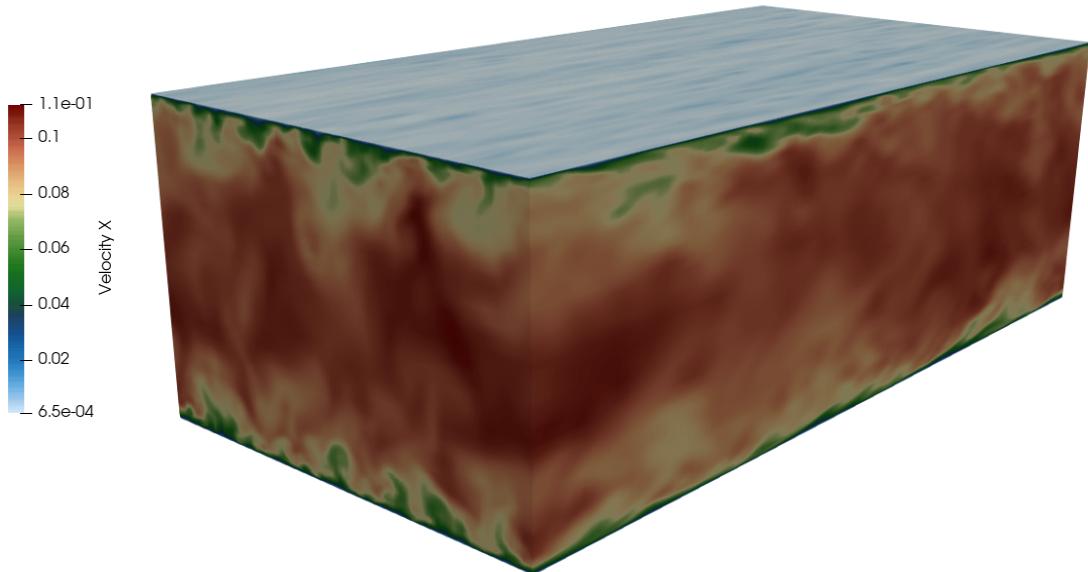
$$s_\tau = \left[\frac{\langle \tau_{ij}^a \tau_{ij}^a \rangle}{\langle (\tau_{mn}^a + R_{mn}^a)(\tau_{mn}^a + R_{mn}^a) \rangle} \right]^{1/2} \quad \mathcal{X}_1[1 - (1/s_\tau)^2]C_k^4 - \mathcal{X}_2C_k^2 + \mathcal{X}_3 = 0$$

- Because mesh independence, ILSA yields grid converged results for a range of s_τ .
- Suitable for h -adaptivity, and p -adaptivity in high-order methods.

1. A. Rouhi, and U. Piomelli, 2016. Phys. Rev. Fluids 1, 044401.

2. O. Lehmkuhl, U Piomelli, and G Houzeaux, 2019. Int. J. Heat and Fluid Flow.

Turbulent channel flow at $Re_\tau = 950$



Geometry: $6h \times 2h \times 3h$

$Ma = 0.08$

BC: Top & bottom walls: no-slip.
Others: periodic.

IC: Perturbed analytical solution.

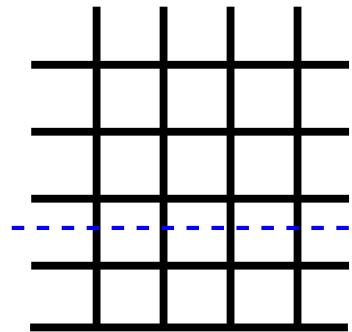
Statistics averaged for $100 T$.

Mesh WMLES (DOFs): $54 \times 62 \times 54$

$$y^+ = 100$$

Numerical methods: a) DGSEM p3, SLAU2 convective scheme, RK4 time scheme, equilibrium wall model
b) FV, Kok convective scheme, RK4 time scheme, equilibrium wall model

Equilibrium wall model implementation

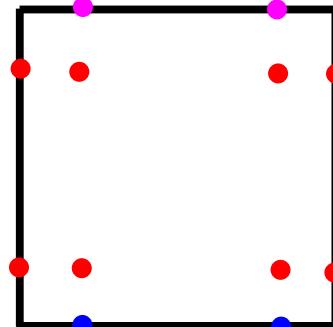


Exchange location: $u(y), y$

Wall model (WM):

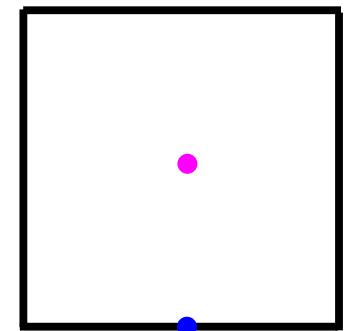
$$\frac{u(y)}{u_\tau} = \frac{1}{k} \ln \left(\frac{u_\tau y}{\nu} \right) + B$$

DGSEM p3



$$y_1^+ = 100$$

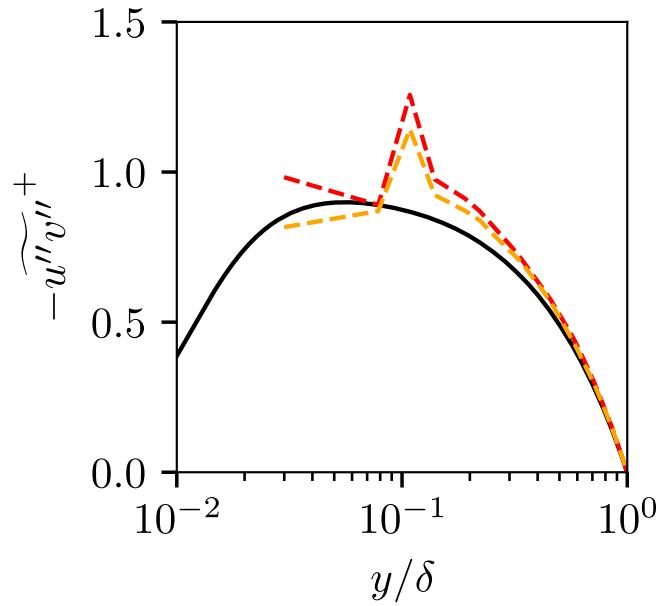
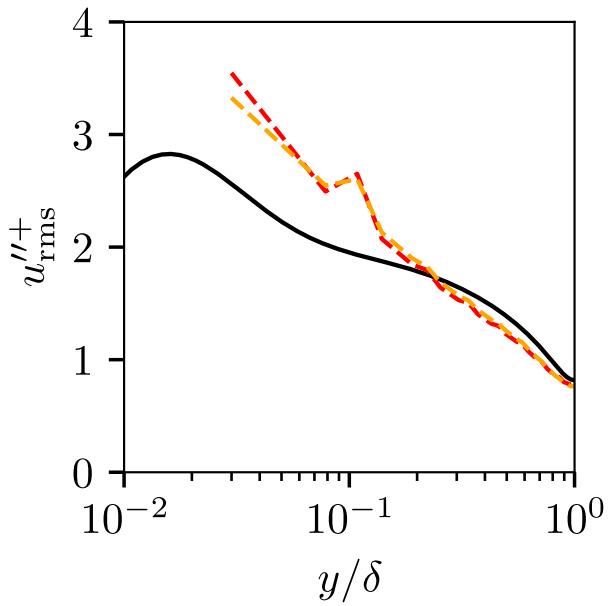
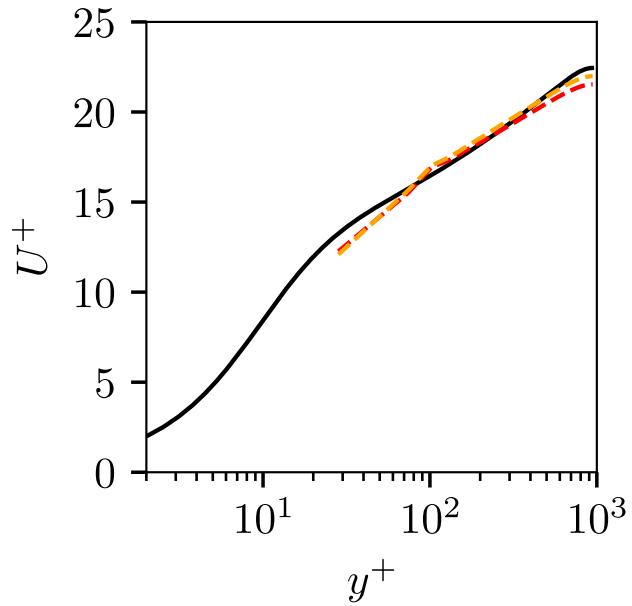
FV



$$y_1^+ = 100$$

DGSEM converged profiles for 100T

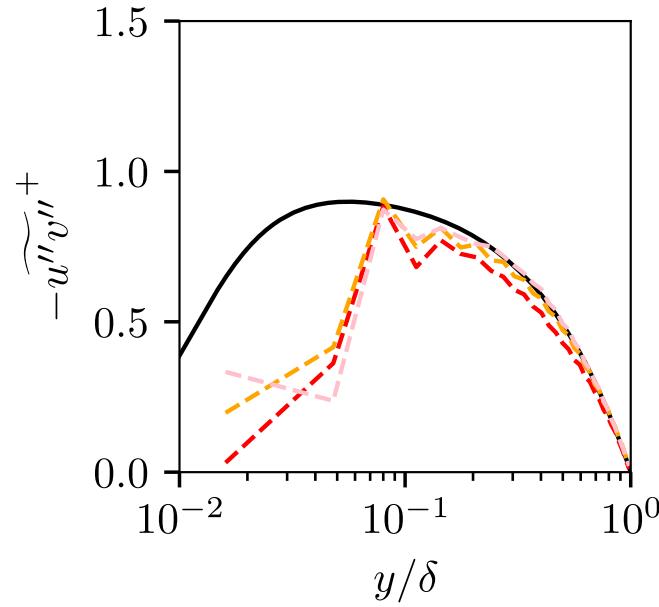
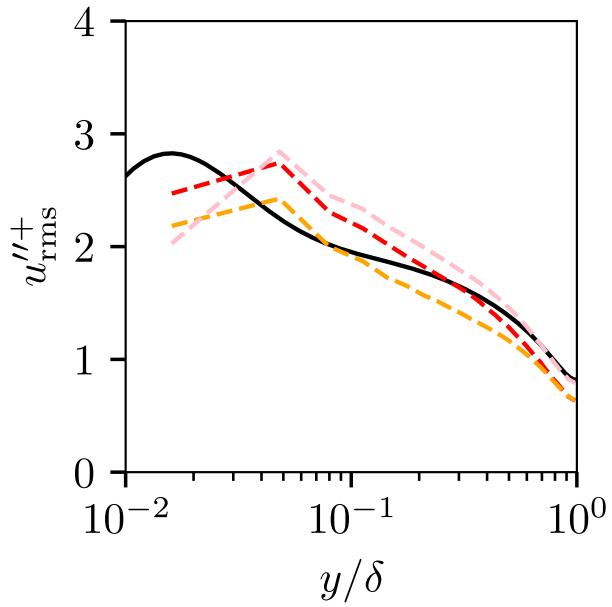
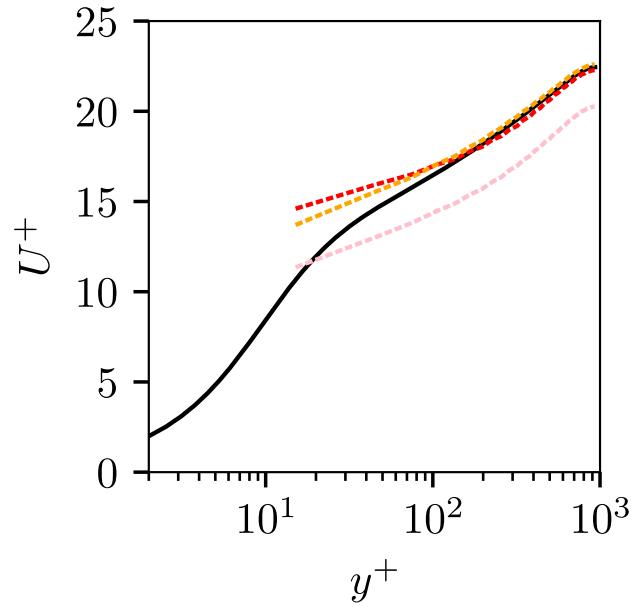
— DNS (3)
- - - DGSEM p3 SLAU2 WM iLES
- - - DGSEM p3 SLAU2 WM ILSA 0.022



- Good agreement above $y^+ = 100$
- ILSA slightly improves the mean profile and the fluctuations.

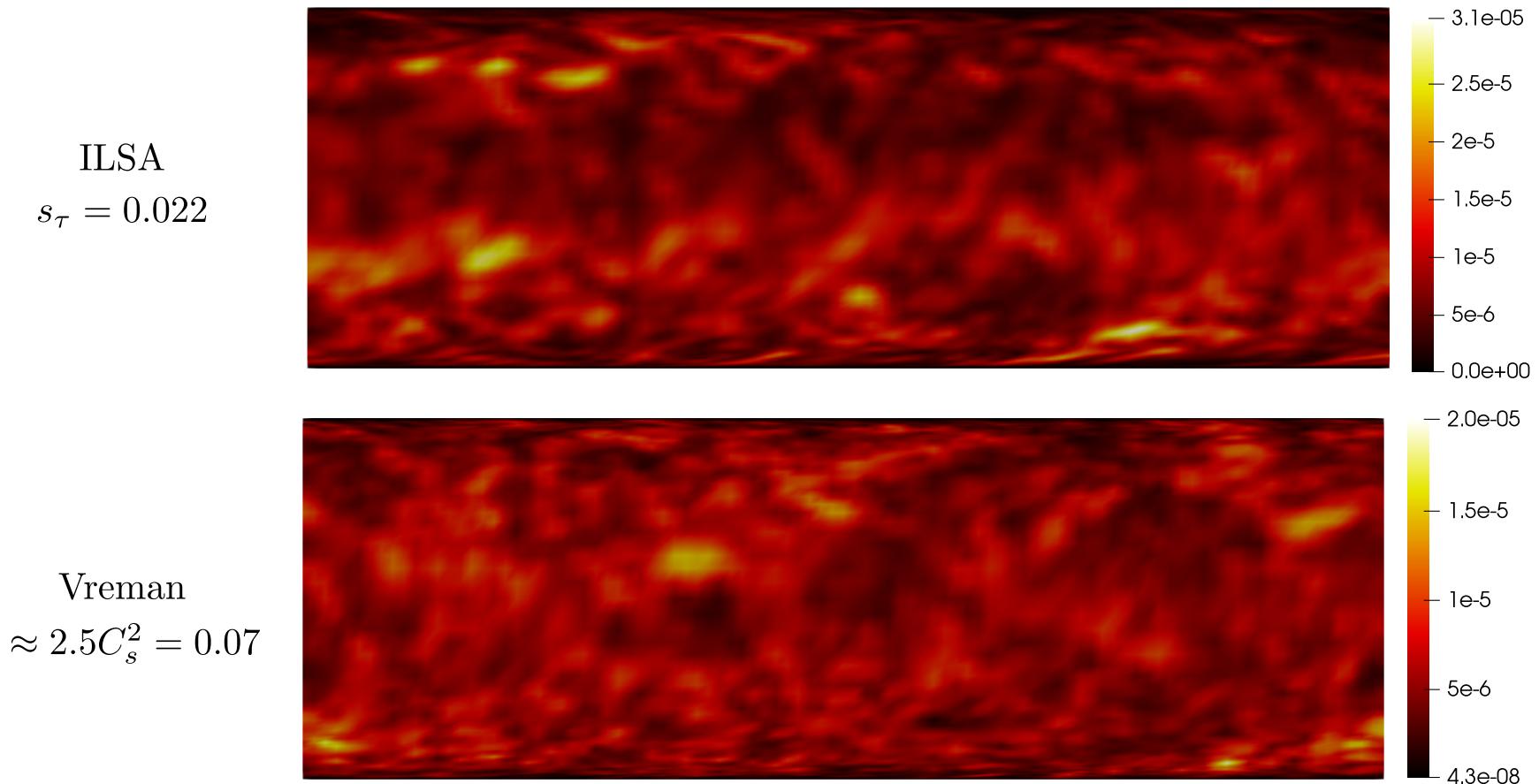
FV converged profiles for 100T

— DNS (3)
- - - FV Kok WM iLES
- - - FV Kok WM iLES $y^+ = 15$
- - - FV Kok WM ILSA 0.022

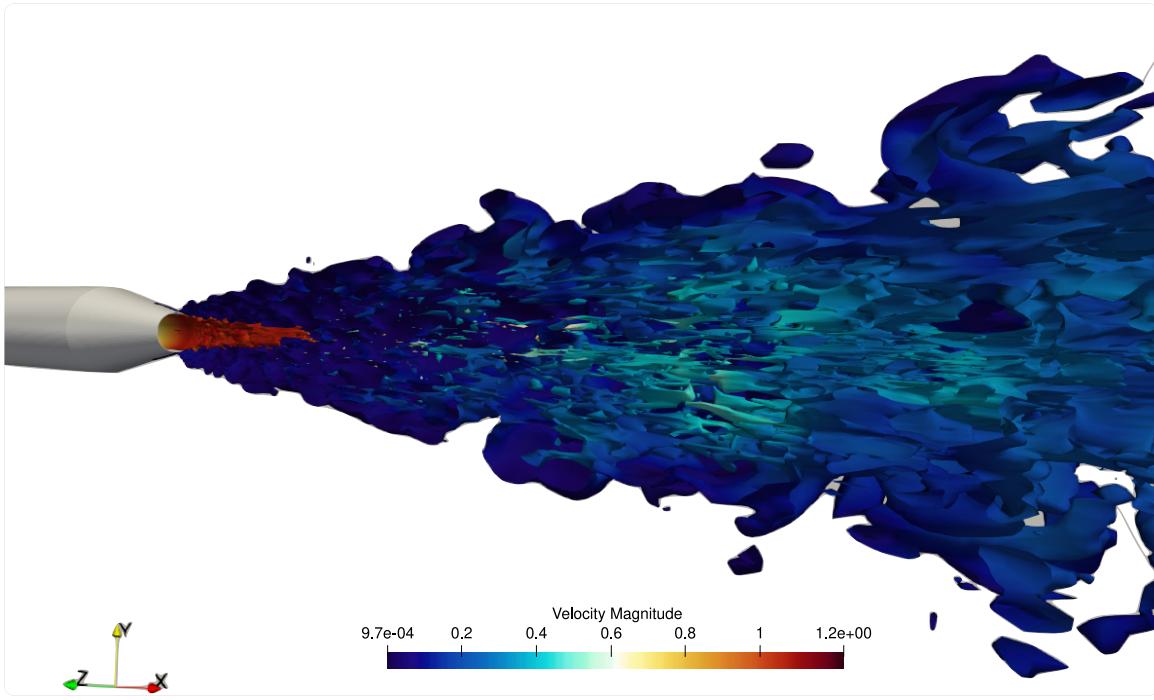


- Log-layer mismatch when using $y^+ = 15$ exchange location
- Good agreement above $y^+ = 100$
- ILSA slightly improves the mean profile and the fluctuations.

The ILSA SGS viscosity follows physical structures in an anisotropic WRLES mesh



Jet exhaust aerodynamics and noise (JEAN)* test case



$$Re = 10^6$$

$$Ma = 0.9$$

IC: RANS-SA using FV + p staging

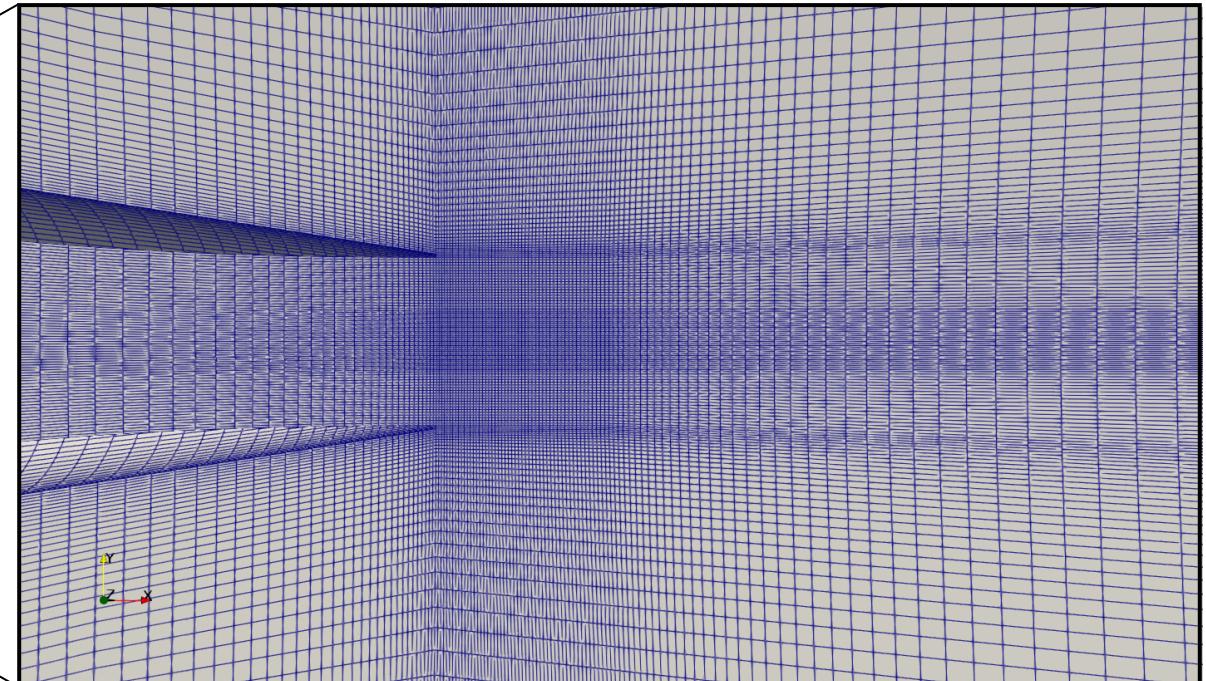
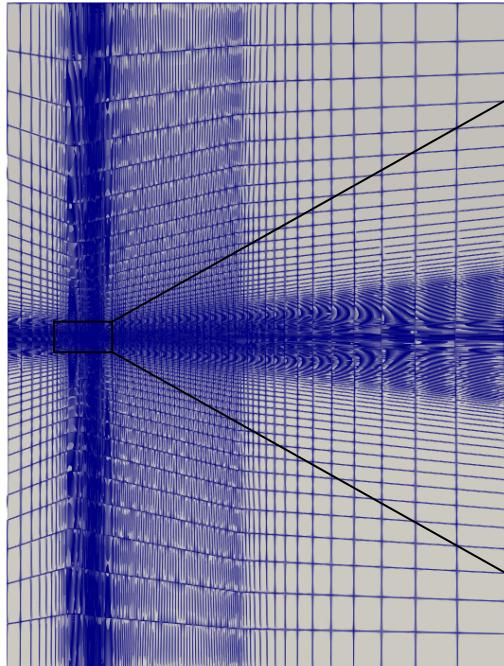
Linear WM mesh: 138M DOFs (2.1M Hex64)
HO WM mesh: 130M DOFs (2.0M Hex64)

Numerical methods:

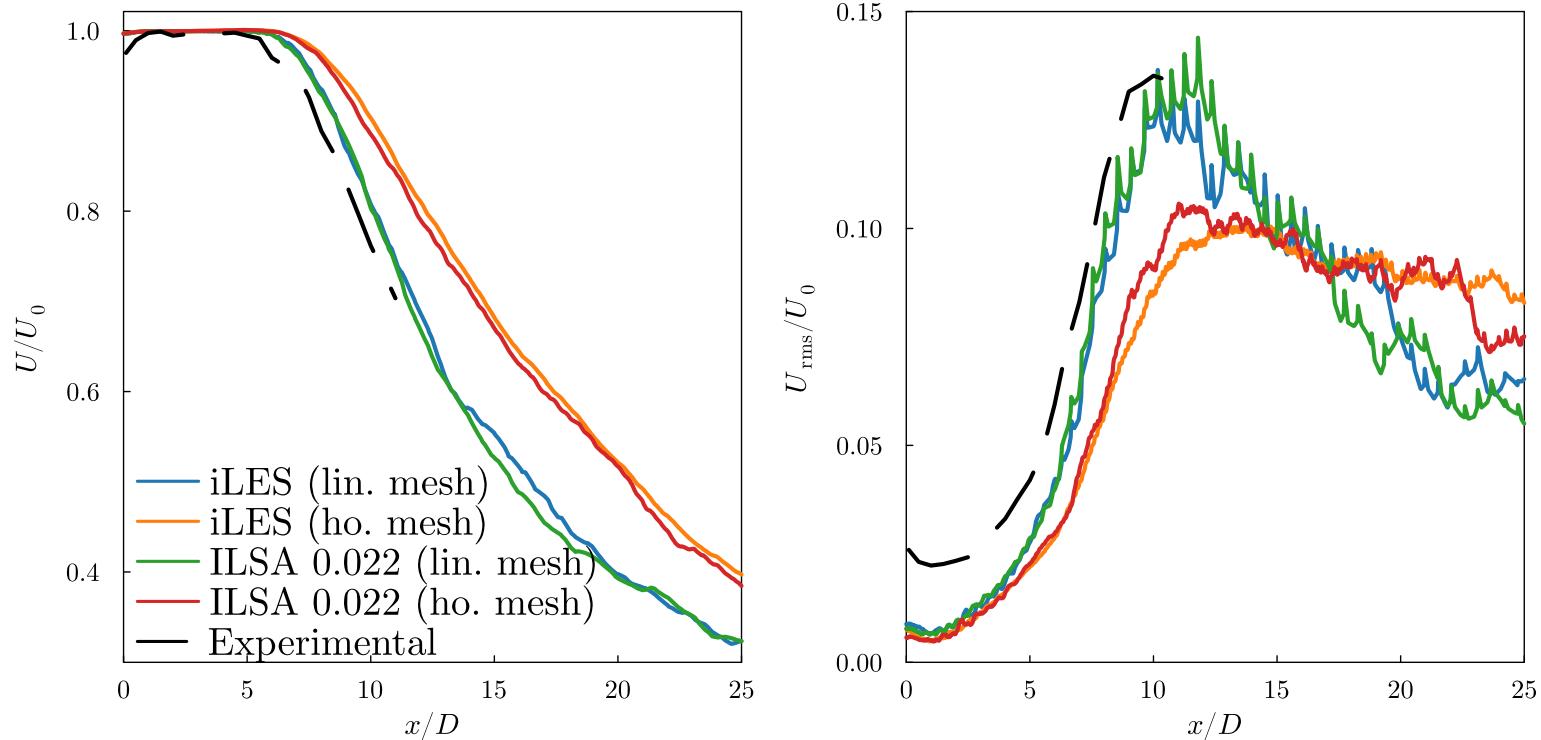
- DGSEM p=3
- SLAU2 convective scheme
- RK3 time scheme
- Equilibrium wall model

* <https://cordis.europa.eu/project/id/G4RD-CT-2000-00313>

A highly anisotropic mesh can be challenging for turbulence models

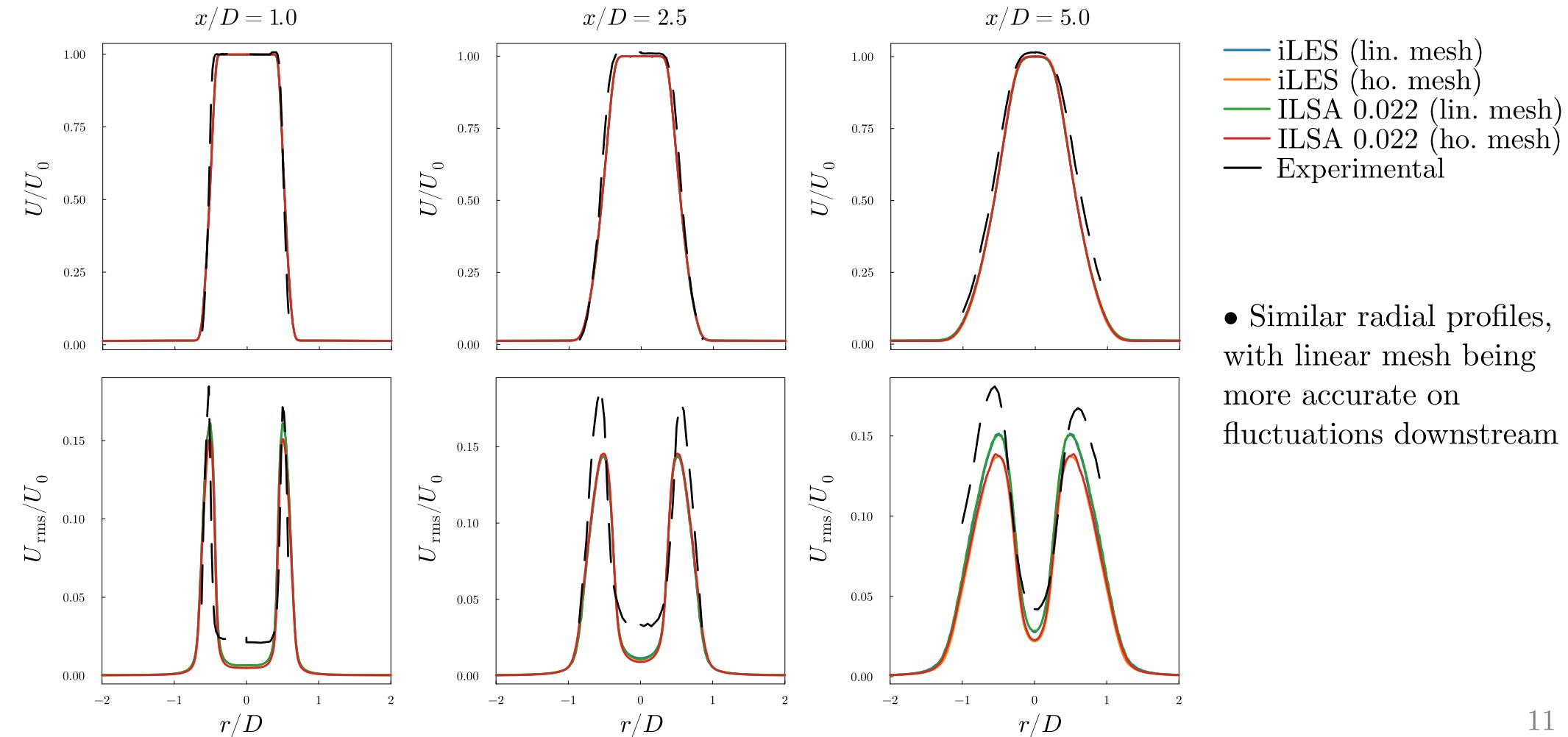


Jet aerodynamics comparison for iLES/ILSA and linear/high-order mesh

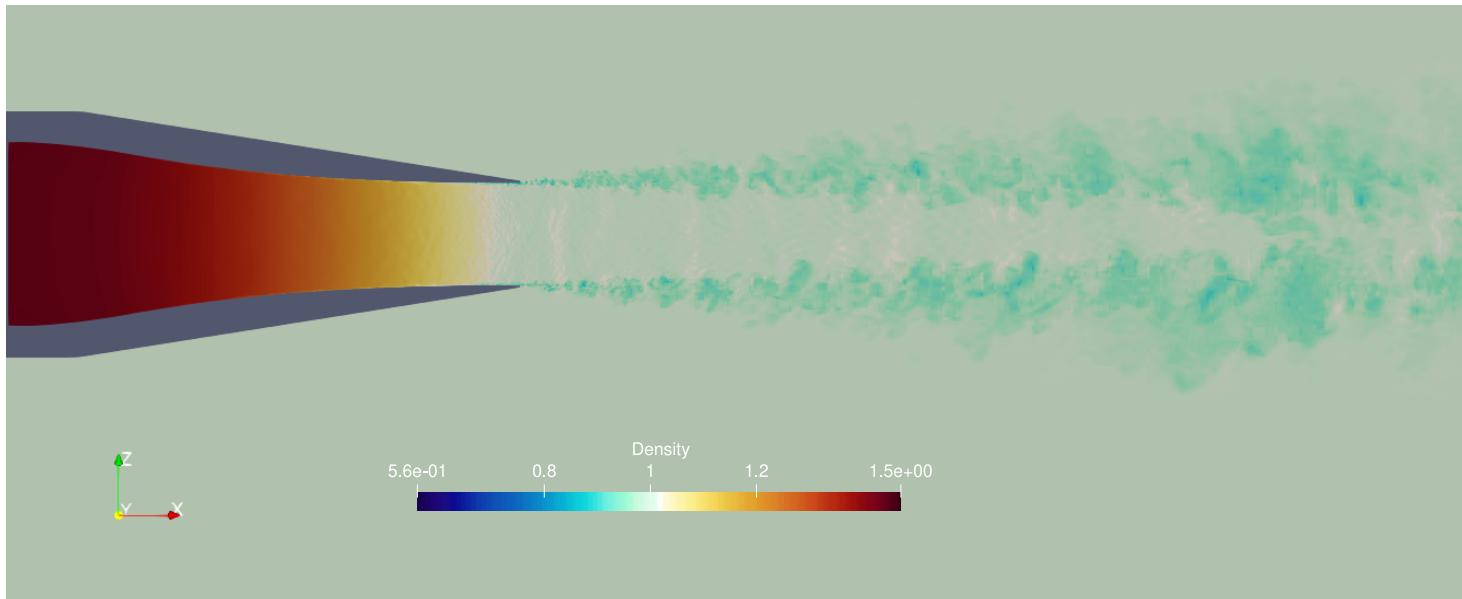


- Linear mesh provides a better general trend, but is more noisy.
- HO mesh mitigates the noise, but more dissipation appears.
- ILSA slightly improves the profiles for both mesh cases.

Jet aerodynamics comparison for iLES/ILSA and linear/high-order mesh

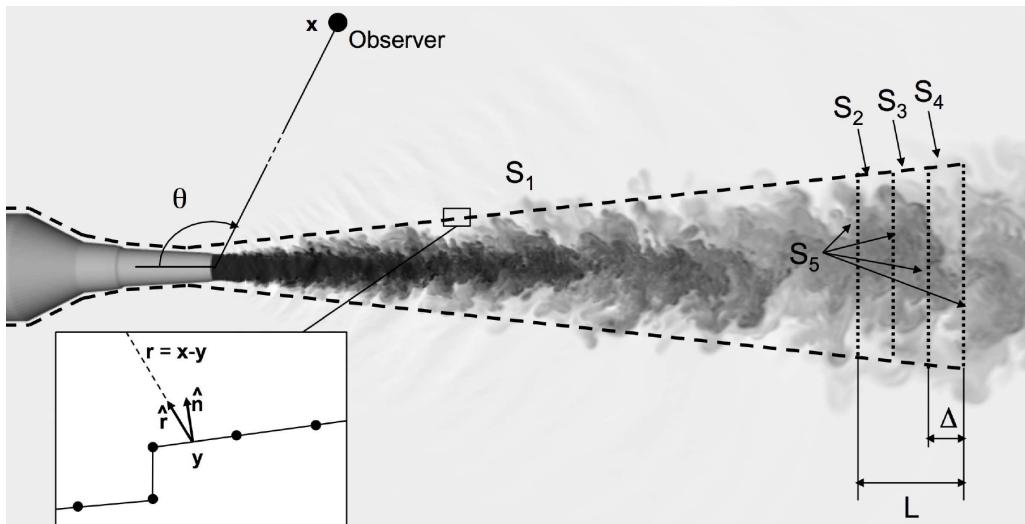


Tripping or a turbulent inlet condition might be necessary...



- If no wall model is used, unphysical artifacts can develop inside the nozzle.
- Adding an equilibrium wall model yields laminar flow inside the nozzle hence impacting the jet exit
- Boundary-layer tripping or turbulent inlet might improve the nozzle outlet flow state.

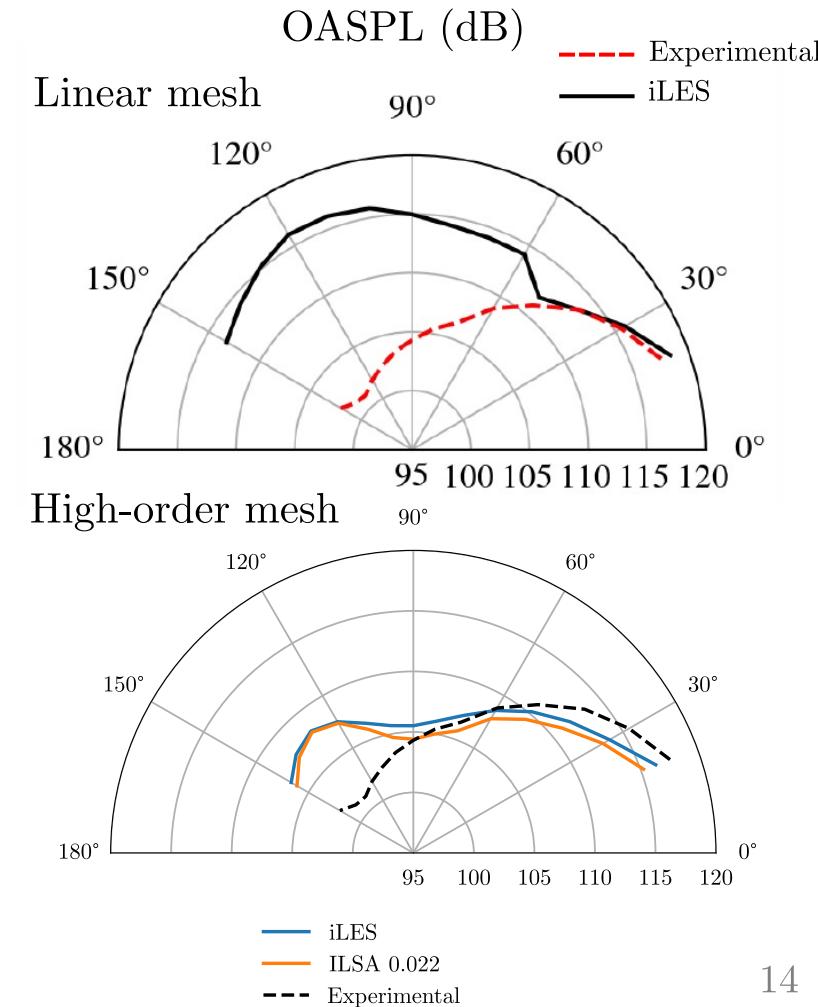
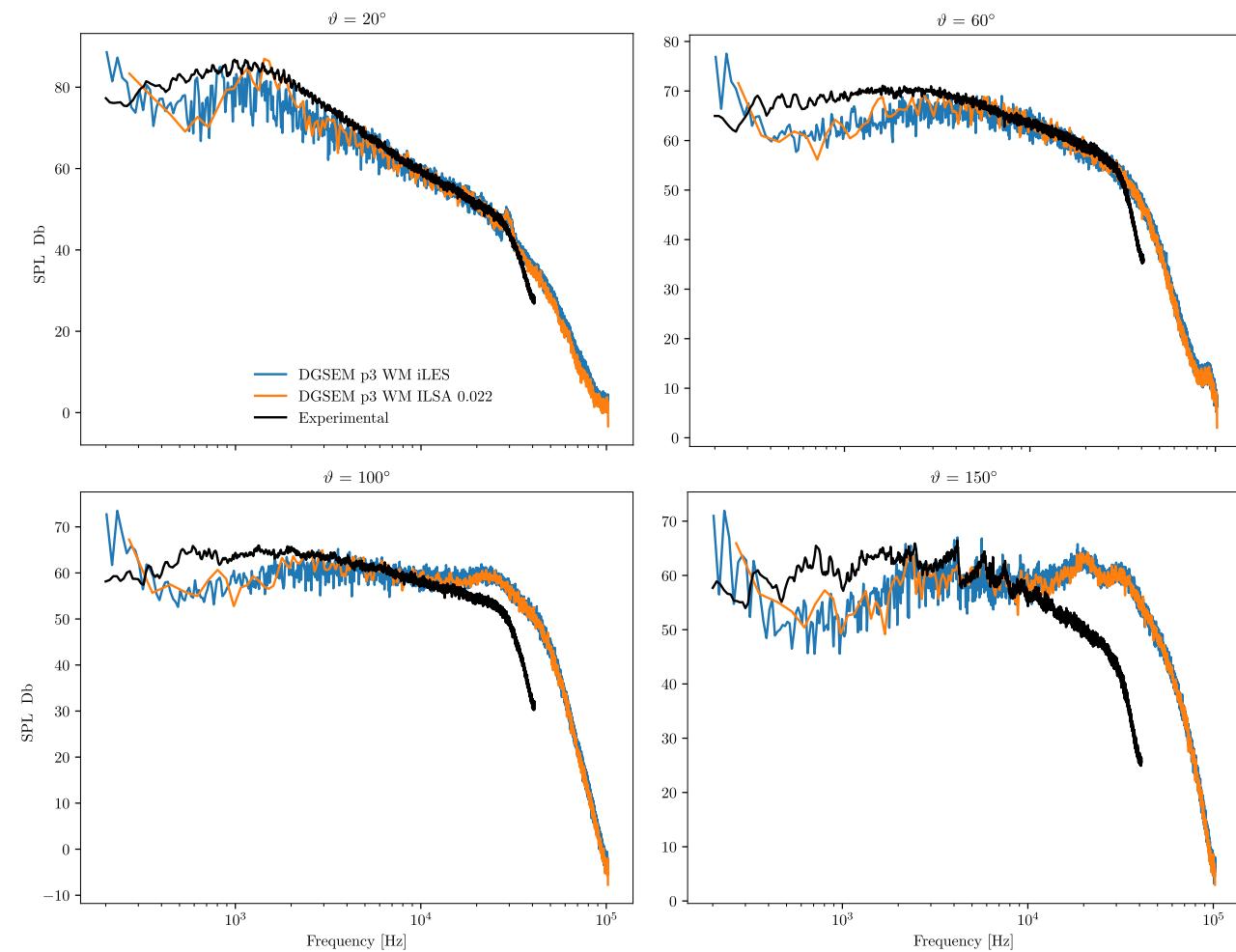
Acoustic analysis comparison for iLES/ILSA on the high-order mesh



Mendez et al., 2013. Int J. Aeroacoustics 12, 1-20

- Far-field sound calculated through the Ffowcs Williams-Hawkings (FWH) equations.
- 1) Flow state (U) temporal evolution saved over S at frequency f during T time.
 - 2) Compute the pressure fluctuating source terms temporal evolution, $F1(\mathbf{x}, \mathbf{y}, t)$ and $F2(\mathbf{x}, \mathbf{y}, t)$, from U , where:
 - \mathbf{x} is the observer location
 - \mathbf{y} is the probe location
 - 3) Integrate $F1$ and $F2$ over S to obtain the sound pressure level (SPL):
$$\text{SPL}(\mathbf{x}, St) [\text{dB}]$$
 - 4) Integrate SPL over St to obtain the overall SPL:
$$\text{OASPL}(\mathbf{x}) [\text{dB}]$$

Acoustic analysis comparison for iLES/ILSA on the high-order mesh



Summary

Channel test case

- iLES and ILSA have been compared in a WMLES turbulent channel.
- ILSA improves the turbulence statistics wrt. iLES in both DGSEM and FV discretizations.
- The ILSA eddy viscosity follows the flow structures even in highly anisotropic WRLES grids.

JEAN nozzle test case

- A highly anisotropic linear mesh and a quadratic mesh have been compared.
- Large mesh-related oscillations appear on the aerodynamic profiles of the linear mesh.
- While the quadratic mesh helps mitigating the oscillations, results are not improved.
- ILSA slightly improves the iLES results.

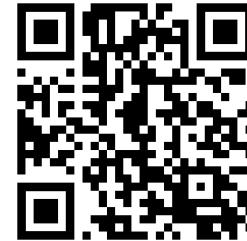
Future work

- Add tripping/turbulent inlet in the JEAN test case.
- CRM high-lift geometry with DGSEM/FV using WM and ILSA.

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Thanks!



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