

✦ 電位 (Electric Potential)

➤ 電位與位能有密切關係，即每單位電荷的位能。

➤ 電位為純量，運算不需考慮向量問題。

$$W_{EXT} = +\Delta U = U_f - U_i = -W_c \quad (\because v = \text{const.} \Rightarrow W_{net} = \Delta K = 0 = W_{EXT} + W_c)$$

● 電位(electric potential)之定義：

$$\Delta V = \frac{\Delta U}{q} \Rightarrow \Delta V = \frac{W_{EXT}}{q} = -\frac{W_c}{q} = \frac{\vec{F}_c \cdot \Delta \vec{s}}{q} \quad (\text{SI單位} \Rightarrow 1 \text{ V} = 1 \text{ J/C})$$

$$\text{考慮微量變化} \Rightarrow dV = \frac{dU}{q} = -\frac{W_c}{q} = -\frac{\vec{F}_c \cdot d\vec{s}}{q}$$

$$W_c \text{ 以靜電力作功表示} \Rightarrow dV = -\frac{q\vec{E} \cdot d\vec{s}}{q} = -\vec{E} \cdot d\vec{s}$$

$$\Rightarrow \Delta V = V_B - V_A = -\int_A^B \vec{E} \cdot d\vec{s}$$

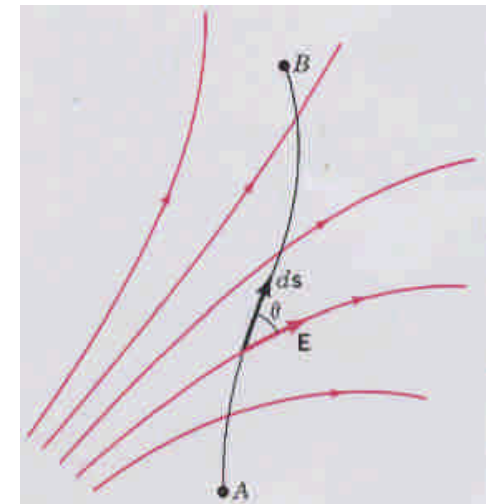


Fig.25.2

重力位能 $\Rightarrow U_g = mgy$; 電位能 $\Rightarrow U_E = qEy$

重力位 $\Rightarrow U_g = gy$; 電位 $\Rightarrow V_E = Ey$

● 均勻電場(uniform field)

$$\Delta V = -\int \vec{E} \cdot d\vec{s} = -\vec{E} \cdot \int d\vec{s} = -\vec{E} \cdot \Delta\vec{s}$$

$\Rightarrow \left\{ \begin{array}{l} \text{電荷沿等位線移動不會造成電位差。} \\ \text{電位沿電場方向呈線性遞減。} \end{array} \right.$

如 Fig. 25.3 $\Rightarrow \Delta V = \pm Ed$

電場 E 的等效單位 $\Rightarrow 1 \text{ V/m} = 1 \text{ N/C}$

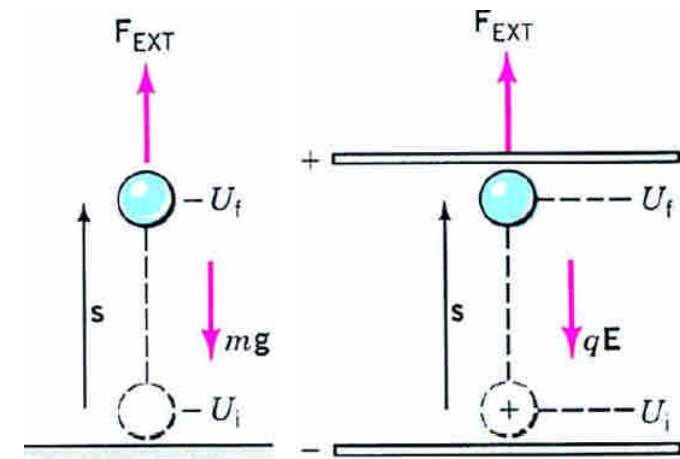


Fig.25.2

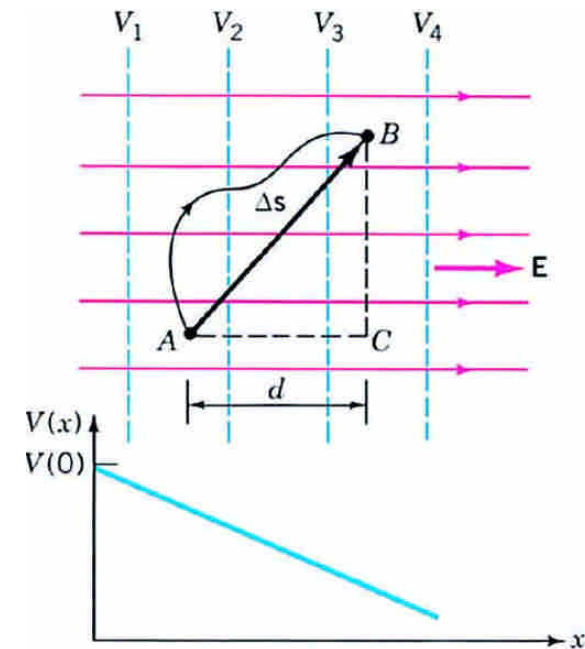


Fig.25.3

- 等位線或等位面(equipotential)

- ⇒ { 1. 電場線必垂直於等位線。
2. 沿等位線移動的電荷不作功。

- 帶電粒子運動的能量

$$\Delta K + \Delta U = 0 \Rightarrow \Delta K = -\Delta U = -q\Delta V$$

$$\Rightarrow \Delta K = e\Delta V \quad (1 \text{ eV} = 1.602 \times 10^{-19} \text{ J})$$

- 點電荷的電位(potential of point charge)

$$\vec{E} = E_r \hat{r} = \frac{kQ}{r^2} \hat{r}$$

$$V_B - V_A = -\int_A^B \vec{E} \cdot d\vec{s} = -\int_A^B E_r dr = -\int_A^B \frac{kQ}{r^2} dr$$

$$= -\left[-\frac{kQ}{r} \right]_A^B = kQ \left(\frac{1}{r_B} - \frac{1}{r_A} \right)$$

考慮 $V_B = 0$ at $r \rightarrow \infty$, 則： $V = \frac{kQ}{r}$

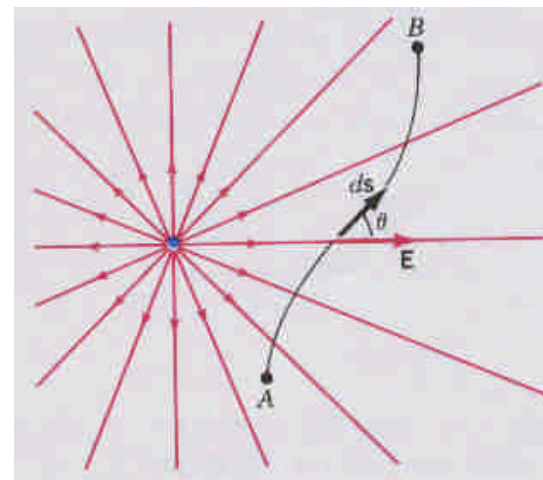


Fig.25.7

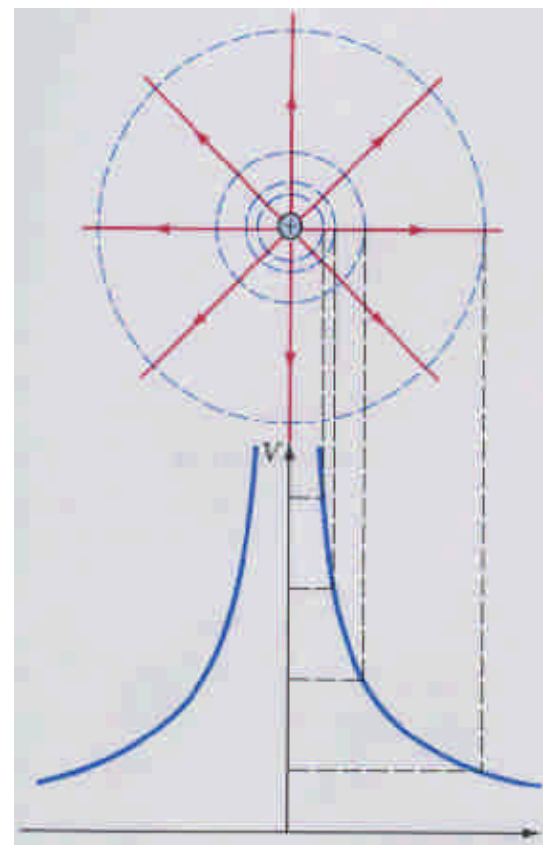


Fig.25.8

- 多個點電荷的電位

— 利用疊加原理(電位為純量，可直接進行加減) $\Rightarrow V = \sum \frac{kQ_i}{r_i}$

➤ Case (A) $\Rightarrow V=0$ (如Fig.25.9)，但 $E \neq 0$ (如Fig.25.10)

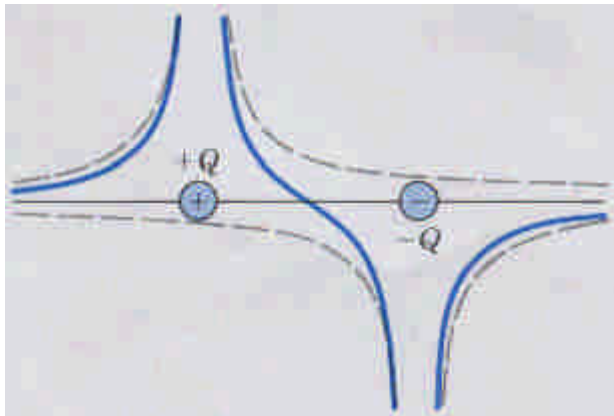


Fig.25.9

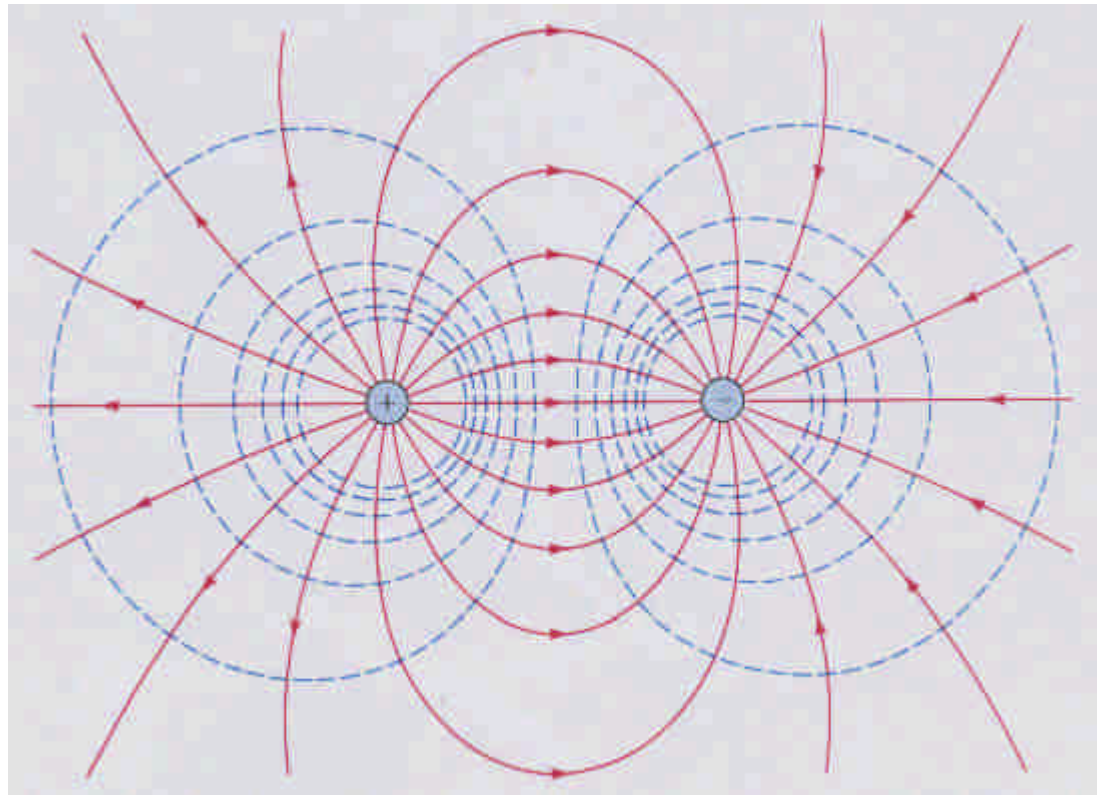


Fig.25.10

➤ Case (B) $\Rightarrow E = 0$ (如 Fig.25.11), 但 $V \neq 0$ (如 Fig.25.12)

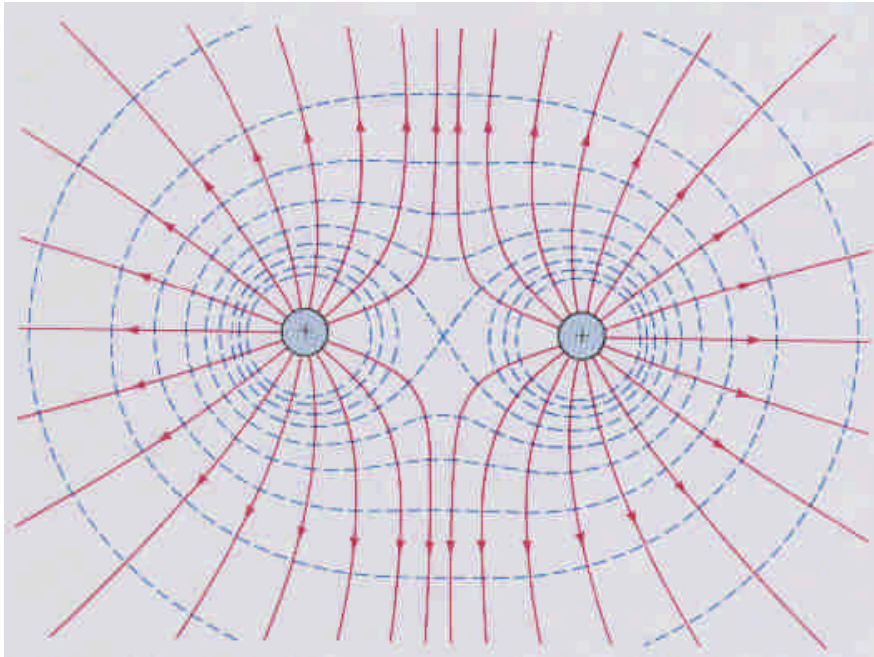


Fig.25.11

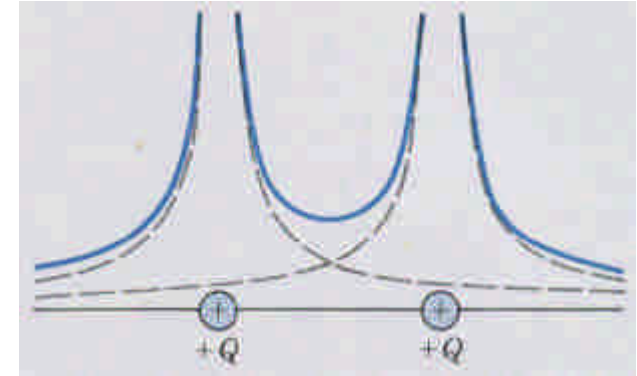


Fig.25.12

● 兩點電荷的電位能

$$\Rightarrow U = qV = q\left(\frac{kQ}{r}\right) = \frac{kqQ}{r} \begin{cases} \text{若極性相同, 則 } U > 0, \text{ 表外力作正功(即保守力作負功)} \\ \text{若極性相異, 則 } U < 0, \text{ 表外力作負功(即保守力作正功)} \end{cases}$$

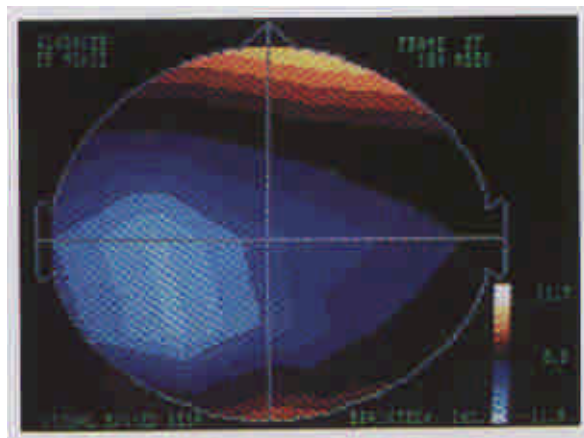
(相當於將兩電荷自無窮遠處等速移至相距 r 處外力所作的功)

$$\Delta U = U(r) - U(\infty) = q\Delta V = q\left(-\int_{\infty}^r \vec{E} \cdot d\vec{r}\right) = -q \int_{\infty}^r \frac{kQ}{r^2} dr = \frac{kqQ}{r} \Rightarrow U(r) = \frac{kqQ}{r} \quad (\because U(\infty)=0)$$

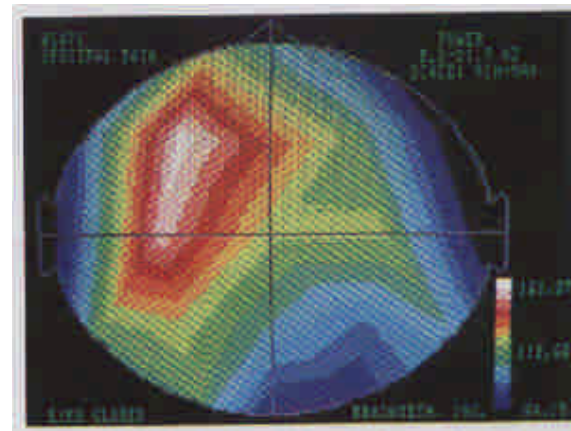
- 多個點電荷的電位能

$$U_{ij} = \frac{kq_i q_j}{r_{ij}}, \quad \text{but } i \neq j \quad (\text{※計算總電位能須注意 } i \neq j)$$

➤ 腦部經閃光刺激後，可激發其電位的分佈



(tumor)



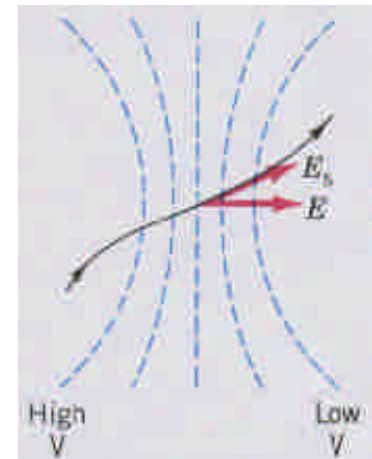
(Epilepsy)

Fig.25.14

- 電位推導電場

$$dV = -\vec{E} \cdot d\vec{s} = -Eds \cos \theta = -E \cos \theta ds = -E_s ds$$

$$\Rightarrow E_s = -\frac{dV}{ds}$$



$$\begin{aligned}\text{考慮三維直角座標系} \Rightarrow dV &= -\vec{E} \cdot d\vec{s} = -(E_x \hat{i} + E_y \hat{j} + E_z \hat{k}) \cdot (dx \hat{i} + dy \hat{j} + dz \hat{k}) \\ &= -(E_x dx + E_y dy + E_z dz)\end{aligned}$$

$$\text{其中 } E_x = -\left(\frac{dV}{dx}\right)_{y,z=\text{const.}} = -\left(\frac{\partial V}{\partial x}\right) \xrightarrow{\text{同理可知}} \vec{E} = -\left(\frac{\partial V}{\partial x}\right)\hat{i} - \left(\frac{\partial V}{\partial y}\right)\hat{j} - \left(\frac{\partial V}{\partial z}\right)\hat{k}$$

●連續電荷分佈(continuous charge distribution)

電
位
計
算
方
式

$$\Rightarrow \left\{ \begin{array}{l} (1) dV = \frac{k dq}{r} \Rightarrow V = k \int \frac{dq}{r} \\ \quad \text{(零電位假設在無窮遠)} \\ (2) V_B - V_A = -\int_A^B \vec{E} \cdot d\vec{s} \\ \quad \text{(零電位可任意假設，但電場須藉由Gauss's Law求出)} \end{array} \right.$$

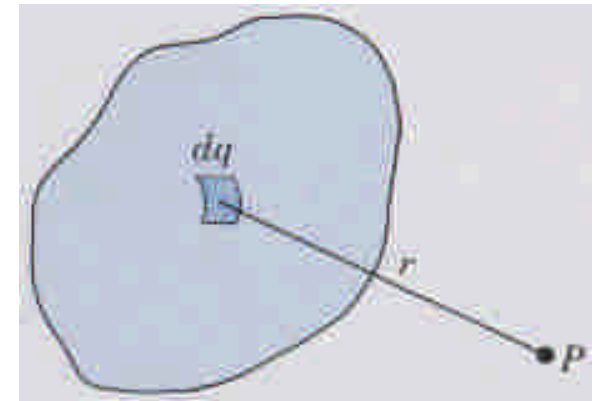


Fig.25.15

Example 25.5 :

單位環形面積所帶電荷 $\Rightarrow dq = \sigma(2\pi x dx)$

$$dV = \frac{k dq}{r} = \frac{k\sigma(2\pi x dx)}{(x^2 + y^2)^{1/2}}$$

$$V = 2\pi k\sigma \int_0^a \frac{x dx}{(x^2 + y^2)^{1/2}} = 2\pi k\sigma \int_0^a \frac{(1/2) dx^2}{(x^2 + y^2)^{1/2}}$$

$$= 2\pi k\sigma \left[(x^2 + y^2)^{1/2} \right]_0^a$$

$$= 2\pi k\sigma \left[(a^2 + y^2)^{1/2} - y \right]$$

$$(\because \text{令 } H = x^2 + y^2 \text{ \& } dH = dx^2 \Rightarrow \int (\frac{1}{2}) H^{-1/2} dH = H^{1/2} + c)$$

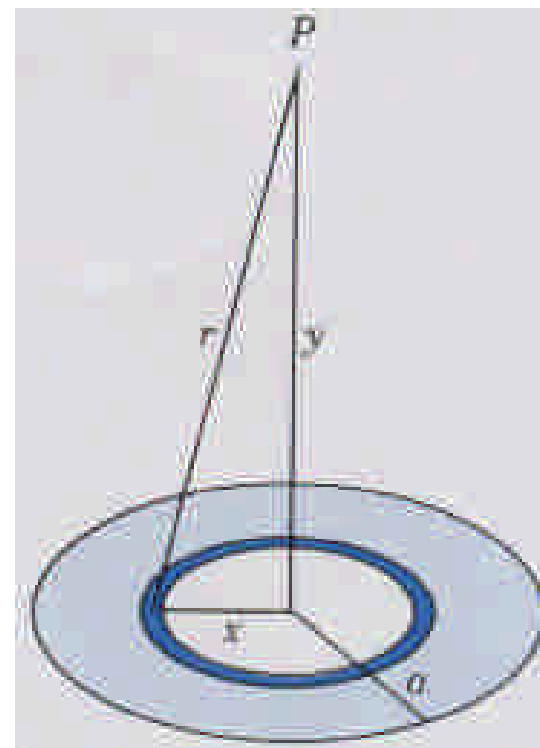


Fig.25.16

討論：

1.有限圓盤在無窮遠處可近似點電荷，即 $y \gg a \rightarrow V \approx \frac{kQ}{y}$

$$V = 2\pi k\sigma \left[(a^2 + y^2)^{1/2} - y \right] = 2\pi k\sigma \left[y + \frac{a^2}{2y} - y \right] = \pi k\sigma \cdot \frac{a^2}{y} = \frac{k\sigma\pi a^2}{y} = \frac{kQ}{y}$$

其中 $(a^2 + y^2)^{1/2} = y \left(1 + \frac{a^2}{y^2} \right)^{1/2} \approx y \left(1 + \frac{a^2}{2y^2} + \dots \right)$ ($\because \sigma\pi a^2 = Q$)
(二項式展開)

$$(\because y \gg a \Rightarrow a/y \ll 1)$$

2.推導電場(習題 Ex.47)

$$\begin{aligned} E_y &= -\frac{\partial V}{\partial y} = -\frac{\partial}{\partial y} \left\{ 2\pi k\sigma \left[(a^2 + y^2)^{1/2} - y \right] \right\} = -2\pi k\sigma \frac{\partial}{\partial y} \left[(a^2 + y^2)^{1/2} - y \right] \\ &= -2\pi k\sigma \left[y(a^2 + y^2)^{-1/2} - 1 \right] = 2\pi k\sigma \left[1 - y(a^2 + y^2)^{-1/2} \right] \end{aligned}$$

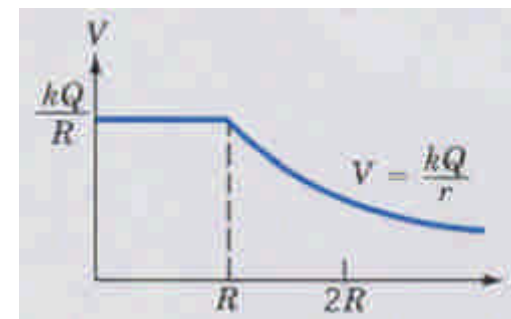
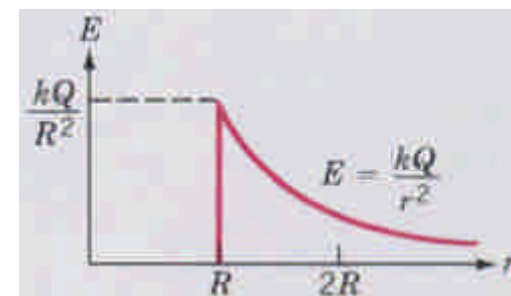
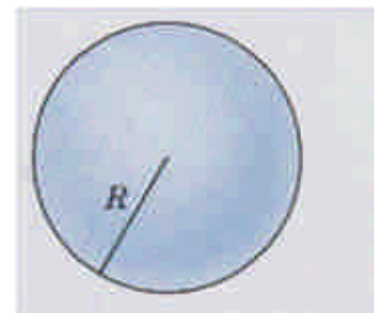
(同Example 23.8結果)

Example 25.6 :

$$\begin{aligned}\text{When } r > R &\Rightarrow E = \frac{kQ}{r^2} ; \quad V(r) - V(\infty) = -\int_{\infty}^r \vec{E} \cdot d\vec{r} \\ &= -\int_{\infty}^r \left(\frac{kQ}{r^2} \hat{r}\right) \cdot d\vec{r} = -\int_{\infty}^r \frac{kQ}{r^2} dr = -kQ \left[-\frac{1}{r} \right]_{\infty}^r = \frac{kQ}{r} \\ \therefore V(\infty) &= 0 \Rightarrow V(r) = \frac{kQ}{r}\end{aligned}$$

$$\text{When } r < R \Rightarrow E = 0 ; \quad V(r) - V(R) = 0$$

$$\therefore V(R) = \frac{kQ}{R} \Rightarrow V(r) = \frac{kQ}{R}$$



Example 25.7 :

$$\text{導體球的電位能} \Rightarrow dW = Vdq = \left(\frac{kq}{R}\right) dq \Rightarrow W = \int_0^Q \frac{kq}{R} dq = \frac{kQ^2}{2R} = \frac{1}{2} QV$$

(考慮外力將dq自無窮遠移至導體球上所作的功)

Discussion : $U = QV$ 表單一電荷 ; $U = 1/2 QV$ 表系統電荷

✦ 導體(conductors)

- 在靜電平衡狀態下，導體內部與表面的電位皆相同。

➤導體內部空腔之電場為0，即所謂的屏蔽效應。

$$V_B - V_A = -\int_A^B \vec{E} \cdot d\vec{s} \quad ; \quad \text{If } V_B = V_A \Rightarrow \vec{E} = 0$$

➤在均勻電場中的電中性導體，接近球面等位面為圓形，電力線呈輻射狀。

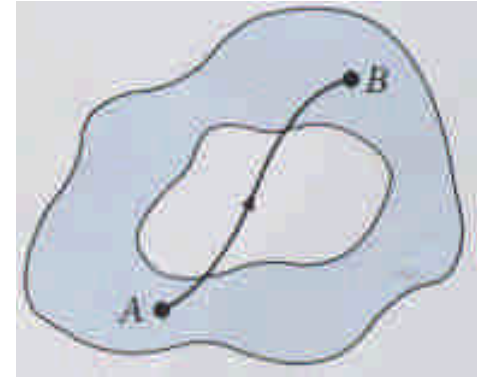


Fig.25.18

- 導體曲率半徑愈小，則電荷密度愈大。

➤導線相連的導體球

$$\begin{aligned} V_1 = V_2 &\Rightarrow \frac{Q_1}{R_1} = \frac{Q_2}{R_2} \Rightarrow \frac{4\pi R_1^2 \sigma_1}{R_1} = \frac{4\pi R_2^2 \sigma_2}{R_2} \\ &\Rightarrow \frac{\sigma_1}{\sigma_2} = \frac{R_2}{R_1} \Rightarrow \sigma_1 R_1 = \sigma_2 R_2 \end{aligned}$$

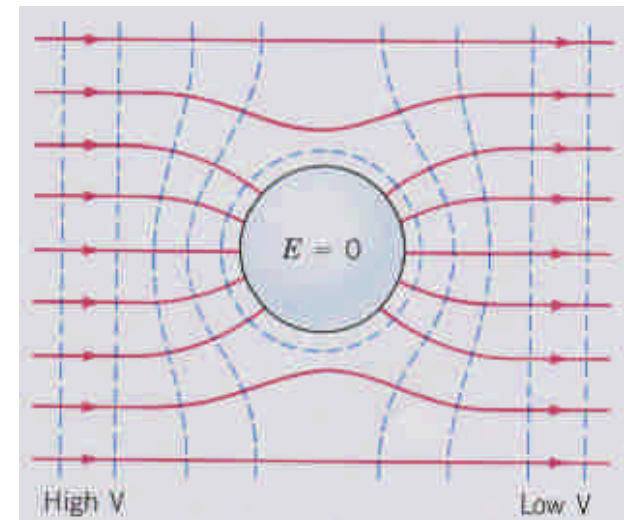


Fig.25.19



Fig.25.20

➤ 尖端放電(discharge)

因靠近導體表面 $E = \sigma/\epsilon_0$ ，由此推知，愈尖銳處會有較大的電場 E ，若電場足夠大 ($\sim 3 \times 10^6$ V/m)，則在空氣中造成放電現象。

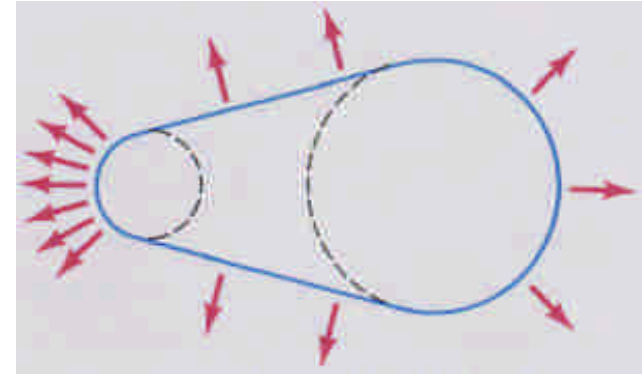


Fig.25.21

崩潰(breakdown)電場的發生

— 因宇宙射線及大地輻射導致空氣分子部份解離，而電場作用會導致已解離的電子加速撞擊其他分子形成更多離子，致使空氣失去絕緣特性變成導體，造成電暈放電(corona discharge)。

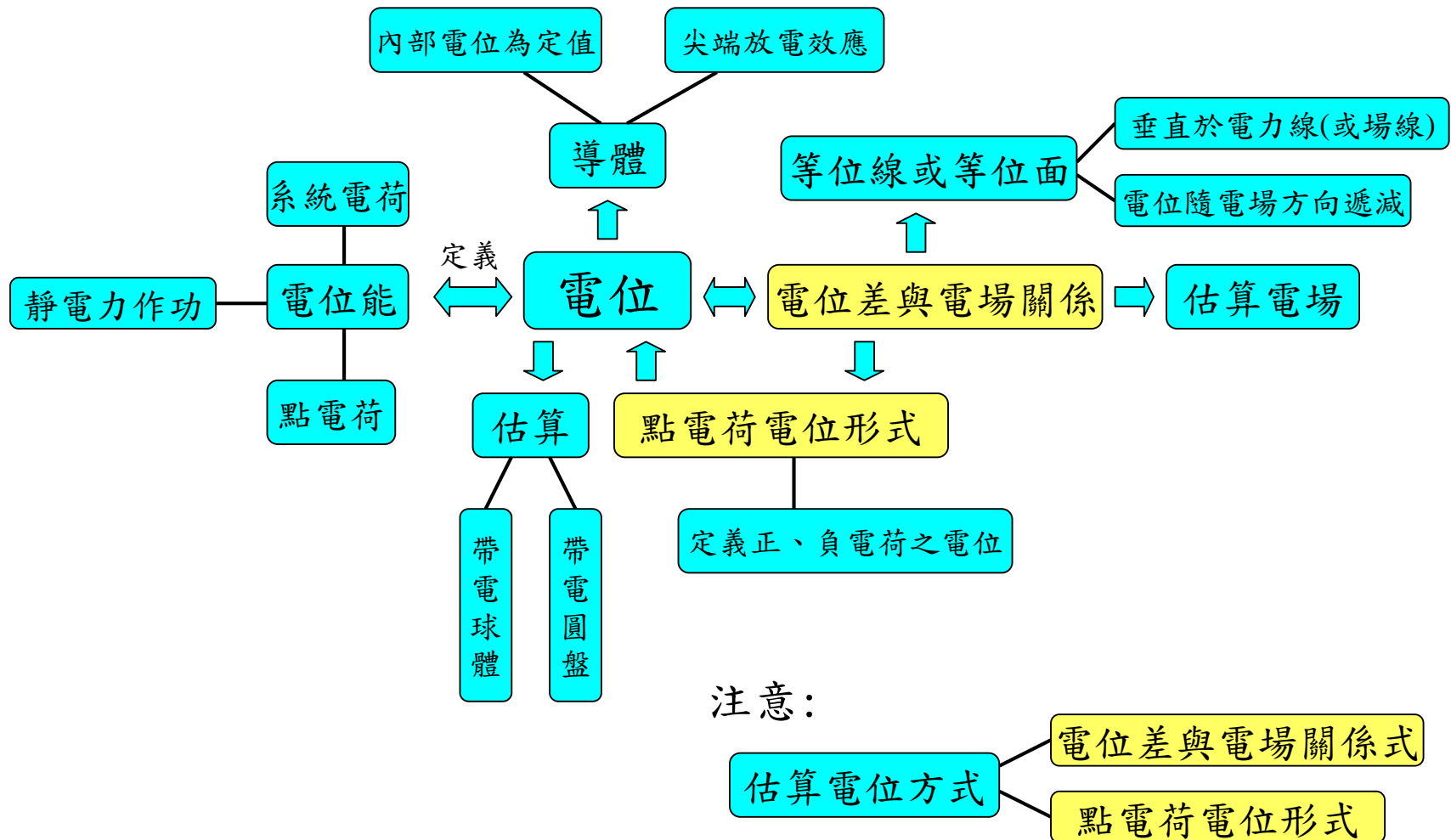
應用

避雷針(lightning rod)、飛機機翼的拖曳短線、場離子顯微鏡(field-ion microscopy)、高壓設備的平緩表面(防止放電)。

塵爆(dust explosion)—半徑0.05 mm的灰塵粒子在150V即能放電。

$$V = kQ/R, E = kQ/R^2 \Rightarrow V = ER \Rightarrow V \propto R$$

本章重要觀念發展脈絡彙整



習題

- 教科書習題 (p.503~p.508)

Exercise: 5,11,15,17,21,23,25,31,35,37,43,45,47,51,57,61,65,69,73

Problem: 5,7,10,11,12,13

※Hint: Problem 12 Ans. (b) $E_r = 2kp \cos\theta/r^3$; $E_\theta = kp \sin\theta/r^3$

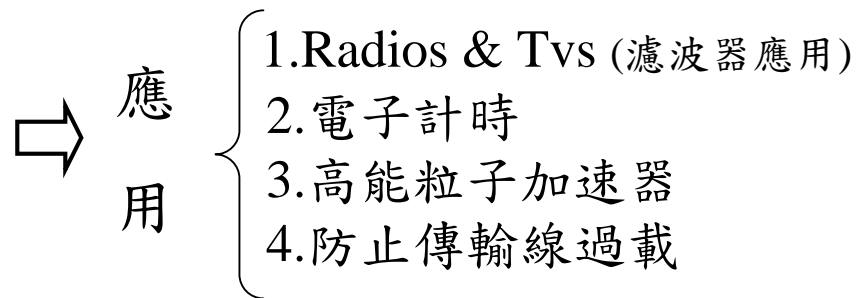
- 基本觀念問題：

- 1.請寫出電位差與電場關係式
- 2.請說明導體尖端放電原理。

- 延伸思考問題：

- 1.請探討Example 25-5帶電圓盤表面的電位。

✦ 電容器(Capacitor)



(Von Kleist)



Fig.26.1

- 電容 (Capacitance) — 相當於一比例常數

$Q = C\Delta V$ — 儲存電荷(Q)正比於兩板間的電位差 ΔV

$= CV$ (將 ΔV 視為一參數 V)

$$\Rightarrow C = \frac{Q}{V} \quad \Rightarrow \text{SI unit: } 1 \text{ farad} = 1 \text{ coulomb/volt}$$

一般電容器的電容值約 $1\text{pF}(10^{-12}\text{F})$ 或 $1\mu\text{F}(10^{-6}\text{F})$ 。

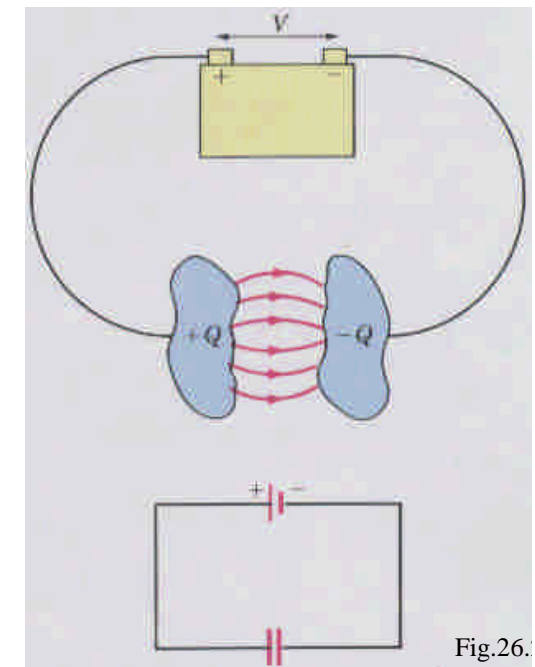


Fig.26.

- 平行電板電容(Parallel-Plate Capacitor)

$$E = \frac{\sigma}{\epsilon_0} = \frac{Q}{\epsilon_0 A} \quad (\because \sigma = \frac{Q}{A})$$

$$C = \frac{Q}{V} = \frac{Q}{Ed} = \frac{Q}{(\frac{Q}{\epsilon_0 A})d} = \frac{\epsilon_0 A}{d}$$

電容C與電容的幾何性質有關：

$$\left\{ \begin{array}{l} \text{If } V = \text{const.} \Rightarrow C \propto Q \Rightarrow \begin{cases} Q = \sigma A \Rightarrow Q \propto A \\ Q = \epsilon_0 EA = \epsilon_0 VA/d \Rightarrow Q \propto 1/d \end{cases} \\ \text{If } Q = \text{const.} \Rightarrow C \propto 1/V \Rightarrow V = Ed \quad (\because E = \frac{Q}{\epsilon_0 A} = \text{const.}) \Rightarrow V \propto d \Rightarrow C \propto 1/d \end{array} \right.$$

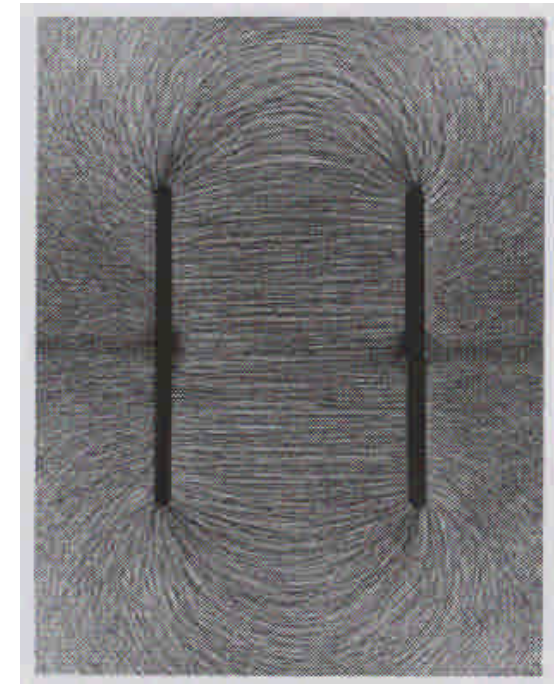


Fig.26.3



Fig.26.4

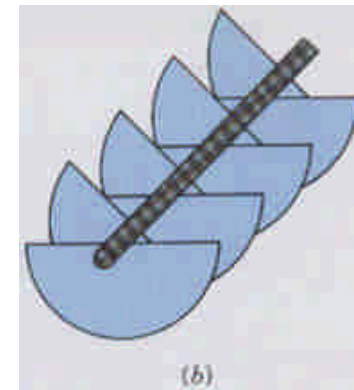
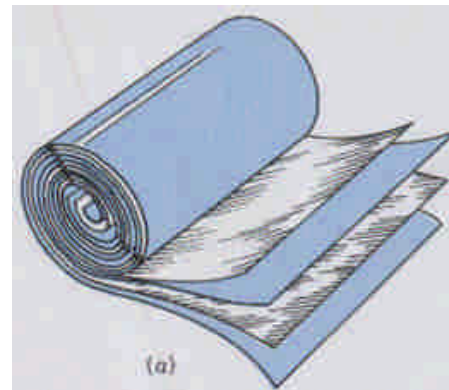


Fig.26.5

Example 26.3：孤立導體球的電容？

假設導體球表面帶電+Q，而地面為電容的另一面

$$C = \frac{Q}{V} = \frac{Q}{|V(R) - V(\infty)|} = \frac{Q}{kQ/R} = \frac{R}{k} = 4\pi\epsilon_0 R \quad (\because k = \frac{1}{4\pi\epsilon_0})$$

Example 26.4：球形電容器(spherical capacitor)

$$V_2 - V_1 = -\int_{R_1}^{R_2} E_r dr = -\int_{R_1}^{R_2} \frac{kQ}{r^2} dr = -\left[-\frac{kQ}{r}\right]_{R_1}^{R_2} = kQ\left(\frac{1}{R_2} - \frac{1}{R_1}\right)$$

$$C = \frac{Q}{V} = \frac{Q}{|V_2 - V_1|} = \frac{Q}{kQ\left(\frac{1}{R_1} - \frac{1}{R_2}\right)} = \frac{R_1 R_2}{k(R_2 - R_1)}$$

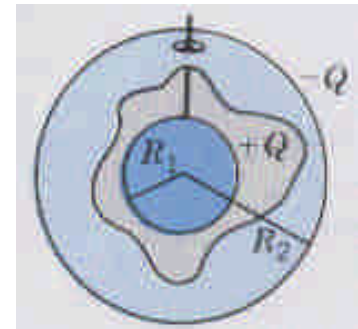


Fig.26.6

Discussion:

$$(a) \text{ If } R_2 \gg R_1 \Rightarrow C = \frac{R_1 R_2}{k(R_2 - R_1)} \approx \frac{R_1 R_2}{kR_2} = \frac{R_1}{k} = 4\pi\epsilon_0 R_1 \quad (\text{近似孤立導體球電容})$$

$$(b) \text{ If } R_2 \approx R_1 = R, \quad R_2 - R_1 = d \Rightarrow C = \frac{R_1 R_2}{k(R_2 - R_1)} \approx \frac{R^2}{kd} = \frac{4\pi\epsilon_0 R^2}{d} = \frac{\epsilon_0 A}{d} \quad (\because A = 4\pi R^2)$$

(近似平行電板電容)

Example 26.5 : 圓柱形電容器(cylindrical capacitor)

$$C = \frac{Q}{|V_b - V_a|} = \frac{\lambda L}{2k\lambda \ln(b/a)} = \frac{L}{2k \ln(b/a)} = \frac{2\pi\epsilon_0 L}{\ln(b/a)}$$

$$V_b - V_a = -\int_a^b E_r dr = -\int_a^b \left(\frac{2k\lambda}{r}\right) dr = -2k\lambda \int_a^b \frac{dr}{r} = -2k\lambda \ln \frac{b}{a}$$

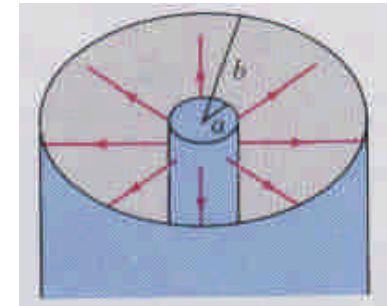


Fig.26.7

●電容串聯(series) $\begin{cases} Q = Q_1 = Q_2 \\ V = V_1 + V_2 \end{cases}$

$$\Rightarrow V = V_1 + V_2 \Rightarrow \frac{Q}{C_{eq}} = \frac{Q_1}{C_1} + \frac{Q_2}{C_2} \Rightarrow \frac{1}{C_{eq}} = \frac{1}{C_1} + \frac{1}{C_2}$$

(series) $\frac{1}{C_{eq}} = \frac{1}{C_1} + \frac{1}{C_2} + \dots + \frac{1}{C_N}$

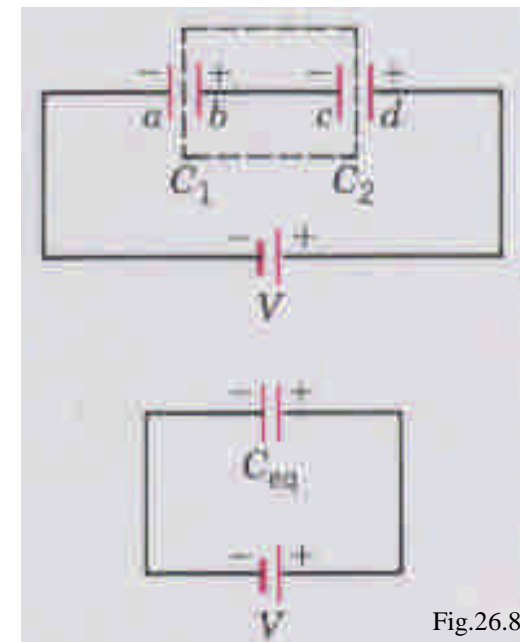


Fig.26.8

- 電容並聯(parallel) $\begin{cases} V = V_1 = V_2 \\ Q = Q_1 + Q_2 \end{cases}$

$$\Rightarrow Q = Q_1 + Q_2 \Rightarrow C_{eq} V = C_1 V_1 + C_2 V_2 \Rightarrow C_{eq} = C_1 + C_2$$

$$(parallel) \quad C_{eq} = C_1 + C_2 + \dots + C_N$$

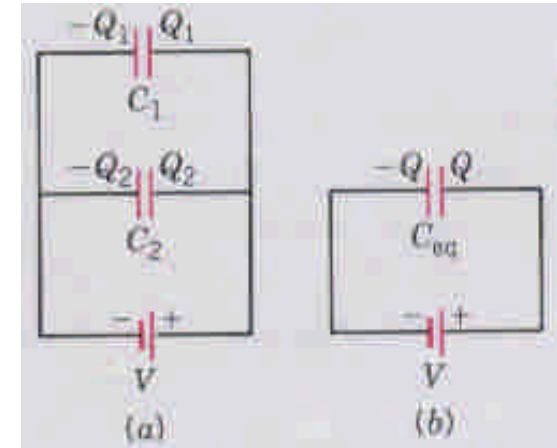


Fig.26.9

- 電容儲存的能量(Energy stored in a capacitor)

$$dW = Vdq = (q/C)dq \text{ (假設} dq \text{ 經導線自負極板移至正極板)}$$

$$\Rightarrow W = \int_0^Q \frac{q}{C} dq = \frac{Q^2}{2C} = \frac{1}{2} QV = \frac{1}{2} CV^2 = U_E$$

$$\text{考慮平行電板電容} \Rightarrow C = \frac{\epsilon_0 A}{d}, V = Ed$$

$$U_E = \frac{1}{2} CV^2 = \frac{1}{2} \frac{\epsilon_0 A}{d} \cdot (Ed)^2 = \frac{1}{2} \epsilon_0 E^2 (Ad)$$

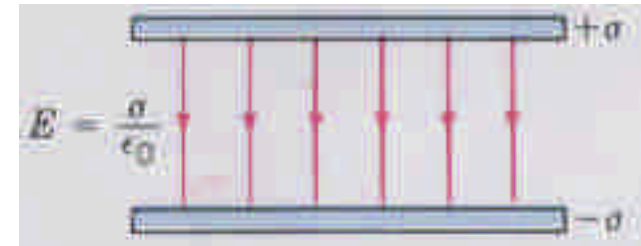


Fig.26.12

$$u_E (\text{energy density}) = U_E / Ad (\text{體積}) = \frac{1}{2} \epsilon_0 E^2 \quad (\text{能量以電場形式儲存，亦即電容係儲存電能。})$$

Example 26.5 : 如Fig.26.11，求初始狀態(a)及末狀態(b)的兩電容器上的電荷、電位差及其儲存能量大小。

初
狀
態
(a)

$$\begin{cases} Q_1 = C_1 V_1 = 5\mu \times 12 = 60\mu C \\ Q_2 = C_2 V_2 = 3\mu \times 12 = 36\mu C \end{cases}$$

$$V_1 = V_2 = 12V$$

$$\begin{cases} U_1 = \frac{1}{2} Q_1 V_1 = \frac{1}{2} (60\mu C)(12V) = 360\mu J \\ U_2 = \frac{1}{2} Q_2 V_2 = \frac{1}{2} (36\mu C)(12V) = 216\mu J \end{cases}$$

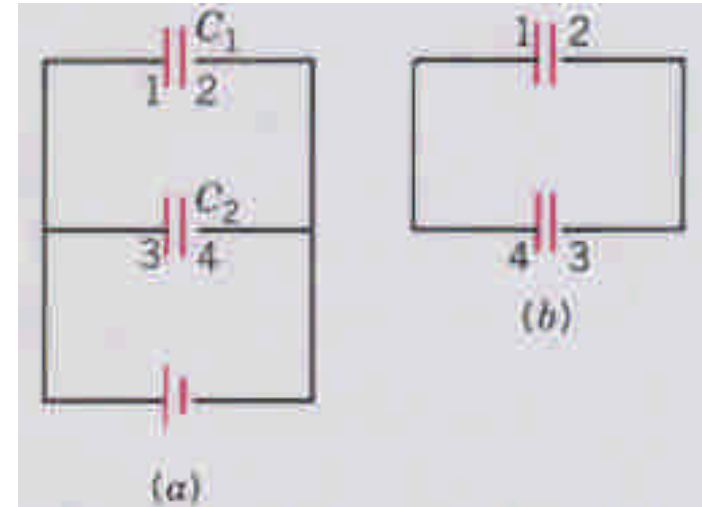


Fig.26.11

重新組合並聯後，假設兩電容的電荷分別為 Q'_1, Q'_2

末
狀
態
(b)

$$\begin{cases} Q'_1 + Q'_2 = 60\mu C - 36\mu C = 24\mu C & (1) \\ V'_1 = V'_2 \Rightarrow \frac{Q'_1}{C_1} = \frac{Q'_2}{C_2} \Rightarrow \frac{Q'_1}{5\mu F} = \frac{Q'_2}{3\mu F} \Rightarrow 3Q'_1 - 5Q'_2 = 0 & (2) \end{cases} \Rightarrow Q'_1 = 15\mu C, Q'_2 = 9\mu C$$

$$V'_1 = V'_2 = \frac{Q'_2}{C_2} = \frac{9\mu C}{3\mu F} = 3V \quad ; \quad \begin{cases} U_1 = \frac{1}{2} Q'_1 V'_1 = \frac{1}{2} (15\mu C)(3V) = 22.5\mu J \\ U_2 = \frac{1}{2} Q'_2 V'_2 = \frac{1}{2} (9\mu C)(3V) = 13.5\mu J \end{cases}$$

Example 26.9 : 孤立金屬球的電位能

$$dU_E = u_E(dV_{\text{volume}}) = \frac{1}{2}\epsilon_0 E^2(4\pi r^2 dr)$$

$$= \frac{1}{2}\epsilon_0 \left(\frac{kQ}{r^2}\right)^2 (4\pi r^2 dr) = \frac{kQ^2}{2r^2} dr$$

$$\left[\because E = \frac{kQ}{r^2} \quad , \quad \text{as } r > R \right]$$

$$U_E = \frac{kQ^2}{2} \int_R^\infty r^{-2} dr = \frac{kQ^2}{2} \left[-\frac{1}{r} \right]_R^\infty = \frac{kQ^2}{2R}$$

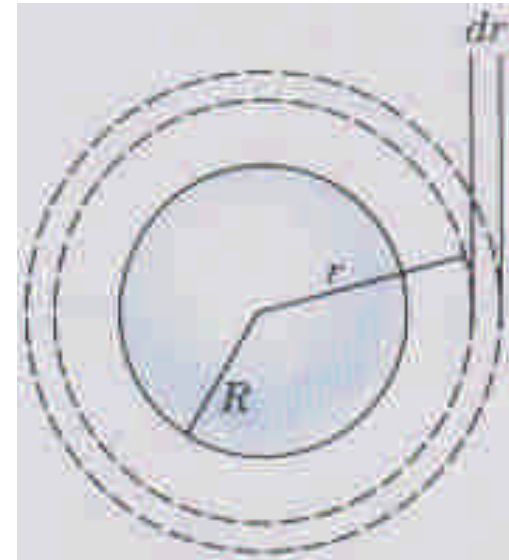


Fig.26.13

✦ 介電質 (Dielectrics)

— 將電容兩平板間插入**非導電物質**會導致電容增加，而此非導電物質即介電質，但也有例外，如：水。

(i) 未接電池 \Rightarrow 電位差 V 會改變，電荷 Q 不變。

$$\Rightarrow V_D = \frac{V_0}{\kappa} \quad (\kappa \text{ 為介電質常數, } \kappa \geq 1)$$

$$\Rightarrow E_D = \frac{E_0}{\kappa} \quad (\because V = Ed)$$

$$\Rightarrow C_D = \frac{Q_0}{V_D} = \frac{Q_0}{V_0/\kappa} = \kappa C_0$$

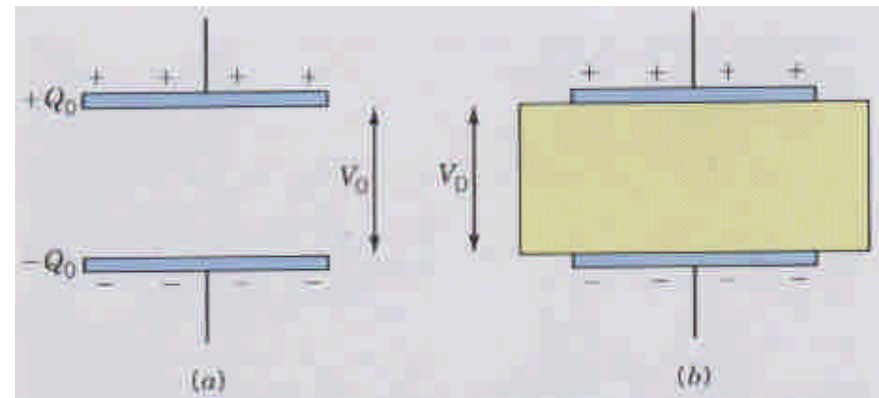


Fig.26.14

(ii) 連接電池 \Rightarrow 電位差 V 不變，電荷 Q 會改變。

$$\Rightarrow Q_D = \kappa Q_0$$

$$\Rightarrow C_D = \frac{Q_D}{V_0} = \frac{\kappa Q_0}{V_0} = \kappa C_0$$

$$(\because E = \sigma / \epsilon_0 \Rightarrow E = \kappa E_0 ,$$

$$\therefore E_D = E / \kappa = \kappa E_0 / \kappa = E_0)$$

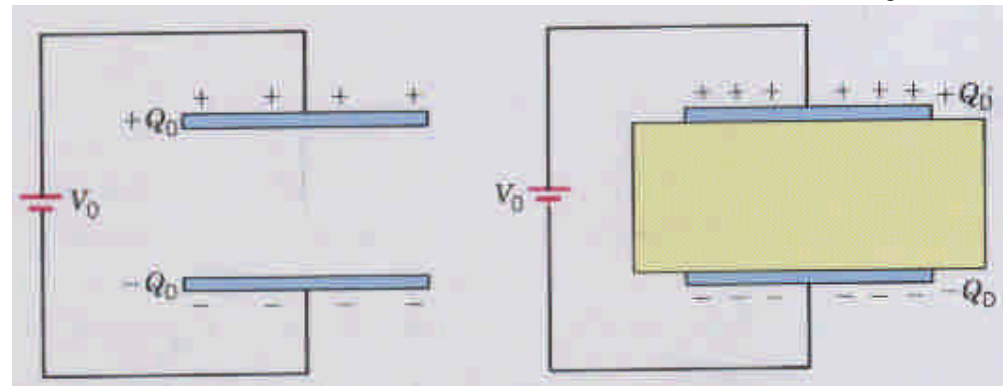


Fig.26.15

●介電質的優點：

- 1.可增加電容。
- 2.可減小電容的體積。
- 3.可增加臨界電位差(或介電強度)，使電容不致放電崩潰。

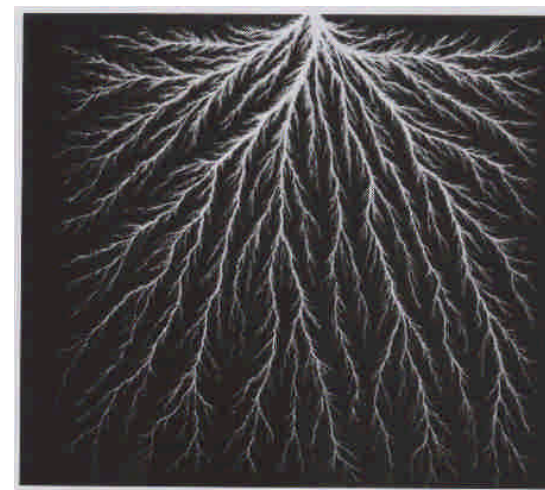


Fig.26.16

●介電質強度(dielectric strength)

—使介電質喪失絕緣特性而崩潰放電的最大電場強度。

●介電質的原子觀點(Atomic view of dielectrics)

- 介電質常數(Dielectric Constant) κ 相當於該電介質內部電荷對外部電場的反應程度。
- 一般分子密度較低，其介電質常數 κ 亦較小，如氣體。
- 介電質若為非極性分子，則由於外加電場的作用會形成感應電偶極矩(induced dipole moment)。

TABLE 26.1 DIELECTRIC
CONSTANTS AND
STRENGTHS

Material	Dielectric Constant	Dielectric Strength (10^6 V/m)
Air	1.00059	3
Paper	3.7	16
Glass	4–6	9
Paraffin	2.3	11
Rubber	2–3.5	30
Mica	6	150
Water	80	—

➤ 感應電偶極矩及永久電偶極矩皆會沿外加電場方向排列，最後在介電質兩端造成電荷區隔現象，即所謂的電極化(polarization)現象，如Fig.26.17。其中區隔的感應電荷極性與平板極性相反，在介電質內部形成與外部電場 E_0 指向相反的感應電場 E_i ，導致介電質的淨電場減小，即：

$$E_D = E_0 - E_i = \frac{E_0}{\kappa}$$

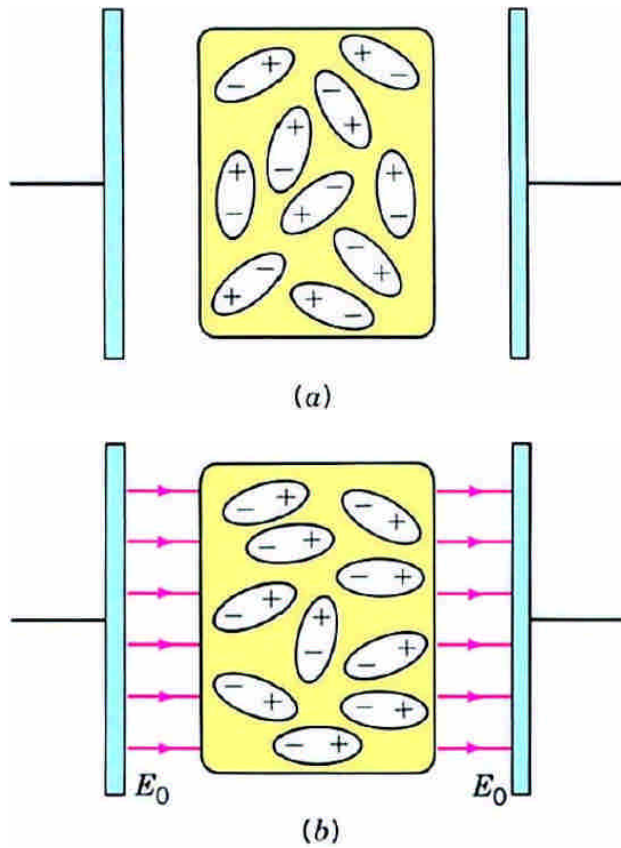


Fig.26.17

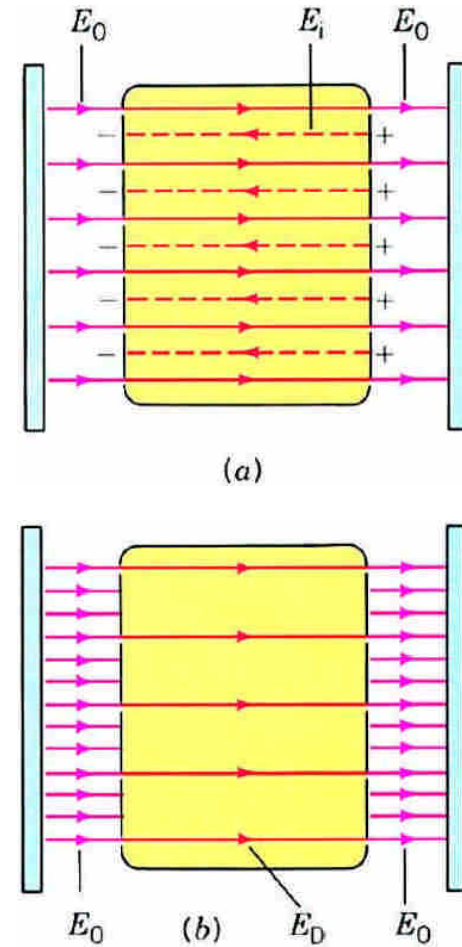


Fig.26.18

Example 26.10 已知介電質的厚度 t 、介電質常數 κ 、平行電板的面積 A 及分隔距離 d ，求 Fig.26.19 的電容？

$$E_0 = \frac{\sigma}{\epsilon_0} \quad \text{and} \quad E_D = \frac{E_0}{\kappa}$$

$$\Rightarrow V = E_0(d - t) + E_D t$$

$$= \frac{\sigma}{\epsilon_0} \left[(d - t) + \frac{t}{\kappa} \right]$$

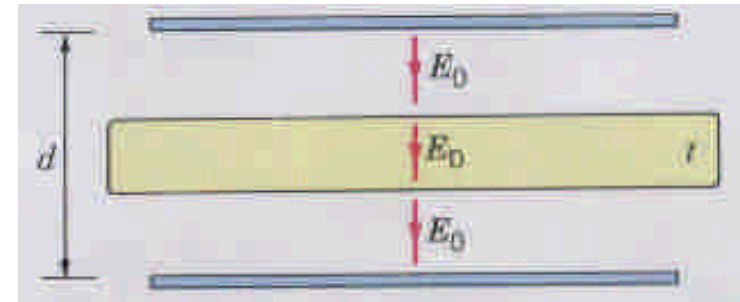


Fig.26.19

$$C = \frac{Q}{V} = \frac{\sigma A}{\frac{\sigma}{\epsilon_0} \left[(d - t) + \frac{t}{\kappa} \right]} = \frac{\epsilon_0 A}{d + t \left(\frac{1}{\kappa} - 1 \right)}$$

(note: If $\kappa = 1$, then $C = \frac{\epsilon_0 A}{d}$)

- 考慮介電質的高斯定律 $\Rightarrow \oint \vec{E} \cdot d\vec{A} = \frac{Q_f}{\epsilon}$ (※以下推導證明僅供參考，不列入考試範圍)

$$E_D = E_0 - E_i = \frac{E_0}{\kappa} \Rightarrow \frac{\sigma_f}{\epsilon_0} - \frac{\sigma_b}{\epsilon_0} = \frac{\sigma_f}{\kappa \epsilon_0} \Rightarrow \sigma_f A - \sigma_b A = \frac{\sigma_f A}{\kappa} \Rightarrow Q_f - Q_b = \frac{Q_f}{\kappa}$$

考慮如右圖虛線的封閉曲面積分：

$$\begin{aligned} \oint \vec{E} \cdot d\vec{A} &= \int \vec{E}_D \cdot d\vec{A} = (E_0 - E_i)A = \left(\frac{\sigma_f}{\epsilon_0} - \frac{\sigma_b}{\epsilon_0} \right) A \\ &= \frac{1}{\epsilon_0} (\sigma_f A - \sigma_b A) = \frac{1}{\epsilon_0} (Q_f - Q_b) = \frac{Q_f}{\kappa \epsilon_0} = \frac{Q_f}{\epsilon} \end{aligned}$$

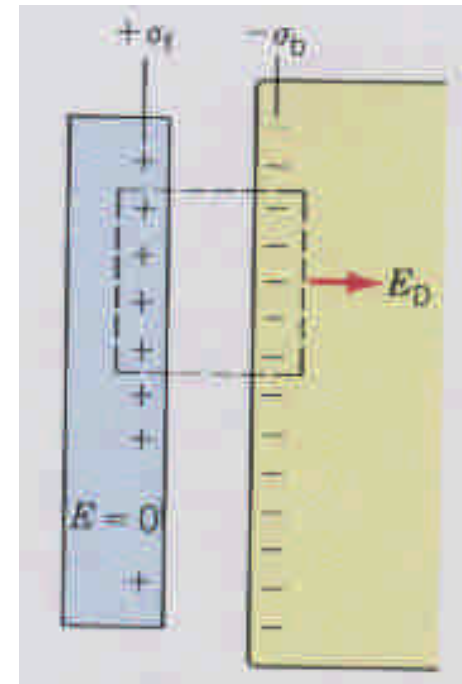
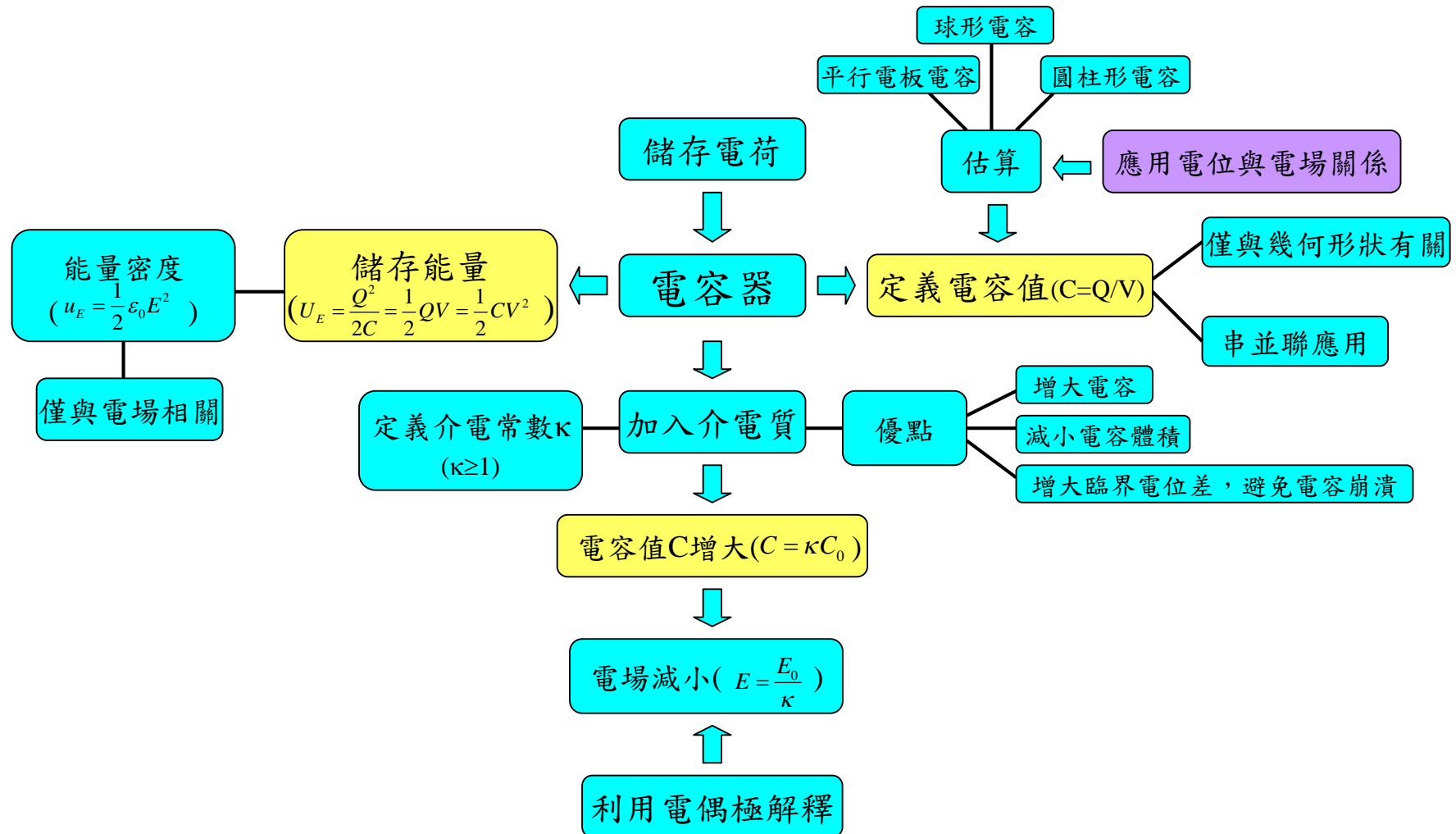


Fig.26.20

本章重要觀念發展脈絡彙整



習題

●教科書習題 (p.525~p.529)

Exercise:

7,11,13,15,19,21,25,29,31,33,34,35,41,42,43,47,53,57,59,63

Problem: 1,8,9,11,13

※Hint:

Ex34 Ans. (a)None; (b)Halved; (c)Halved

Ex42 Ans. $\frac{2\kappa_1\kappa_2C_0}{\kappa_1 + \kappa_2}$

Problem 8 Ans. $CV^2/2\kappa$ (題目修訂：....constant $\kappa\lambda$ 修訂為 κ)

●基本觀念問題：

- 1.請問電容器加入介電質的優點為何？
- 2.請問電容器崩潰放電的原因為何？