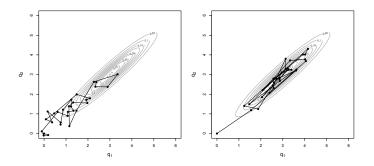
Split Hamiltonian Monte Carlo

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Objective

• Compare the random walk Metropolis algorithm (left) with Hamiltonian Monte Carlo (right).



- HMC explores the parameter space more efficiently but requires costly gradient evaluations.
- Our objective is to reduce the computational cost of HMC.



The Metropolis algorithm

- 1. Given our current state $q^{(n)}$, we propose a new state q^* .
- 2. Calculated the acceptance probability

$$a(q^{(n)},q^*) = \min(1,rac{f(q^*)}{f(q^{(n)})})$$

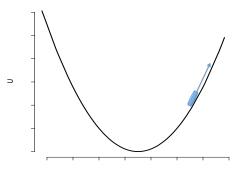
3. With probability $a(q^{(n)},q^*)$, accept the proposed state q^* as the new state, $q^{(n+1)}=q^*$, or remain at the current state $q^{(n+1)}=q^{(n)}$.

Hamiltonian Monte Carlo

- Hamiltonian Monte Carlo (HMC) reduces the random walk behavior of Metropolis.
- It proposes states that are distant from the current state, but nevertheless have a high probability of acceptance.
- These distant proposals are found by numerically simulating Hamiltonian dynamics for some specified amount of fictitious time.
- Simulation involves costly evaluation of the gradient of the log density.

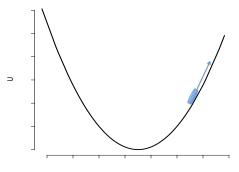
Physical interpretation of HMC

- Consider a frictionless hockey puck that slides on a surface of varying height.
- The state space of this dynamical system consists of its position q and its momentum (i.e., mv) denoted as p.



Physical interpretation of HMC

- Based on q, we define the potential energy, U(q), such that U(q) is proportional to the height of the surface at position q.
- Based on p, we define the kinetic energy, K(p); the kinetic energy is $m|v|^2/2$, so $K(p)=|p|^2/(2m)$.



Hamiltonian's equations

• The dynamic system can be represented by the *Hamiltonian* function:

$$H(q,p) = U(q) + K(p)$$

 Hamilton's equations determine how q and p change over time:

$$\frac{dq_j}{dt} = \frac{\partial H}{\partial p_j}$$
$$\frac{dp_j}{dt} = -\frac{\partial H}{\partial a_i}$$

• They define a mapping, T_s , from the state at some time t to the state at time t + s.

Sampling

In Bayesian statistics, q consists of the model parameters,

$$U(q) = -\log(P(q)L(q|D))$$

- We also introduce fictitious momentum variables p.
- Typically, we set $p \sim N(0, M)$ so

$$K(p) = \sum_{i} p_i^2/2m_i$$

where M is called the *mass matrix* and is usually set to I.

• The joint density of q and p is

$$P(q,p) = \frac{1}{Z} \exp(-H(q,p))$$
$$= \frac{1}{Z} \exp(-U(q)) \exp(-K(p))$$



Properties of HMC

- Reversibility; the target distribution remains invariant
- Conservation of the Hamiltonian; the acceptance probability is one.
- **Volume preservation**; the determinant of the Jacobian matrix for the mapping is one.
- See Neal (2010) for detailed discussion.

Euler's method

- In practice, we need to approximate these equations by discretizing time, using some small step size ε .
- We can use Euler's method,

$$p_j(t+\varepsilon) = p_j(t) + \varepsilon \frac{dp_i}{dt}(t) = p_j(t) - \varepsilon \frac{\partial U}{\partial q_j}(q(t))$$

$$q_j(t+\varepsilon) = q_j(t) + \varepsilon \frac{dq_j}{dt}(t) = q_j(t) + \varepsilon \frac{\partial K}{\partial \rho_j}(\rho(t))$$

· However, the approximation error is high.

The leapfrog method

• It is more common to use the *leapfrog* method instead:

$$p_{j}(t+\varepsilon/2) = p_{j}(t) - (\varepsilon/2)\frac{\partial U}{\partial q_{j}}(q(t))$$

$$q_{j}(t+\varepsilon) = q_{j}(t) + \varepsilon\frac{\partial K}{\partial p_{j}}(p(t+\varepsilon/2))$$

$$p_{j}(t+\varepsilon) = p_{j}(t+\varepsilon/2) - (\varepsilon/2)\frac{\partial U}{\partial q_{j}}(q(t+\varepsilon))$$

• At the end of the L leapfog steps, we have a propose a new state (q^*, p^*) , which is accepted with probability

$$\min[1, \exp(-H(q^*, p^*) + H(q, p))]$$

The leapfrog method

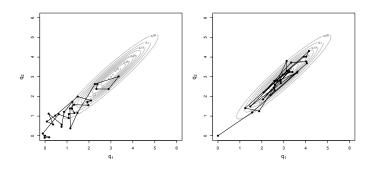
Sample initial values for p from N(0, I)

$$\begin{aligned} & \textbf{for } \ell = 1 \text{ to } L \quad \textbf{do} \\ & p \leftarrow p - (\varepsilon/2) \frac{\partial U}{\partial q} \\ & q \leftarrow q + \varepsilon p \\ & p \leftarrow p - (\varepsilon/2) \frac{\partial U}{\partial q} \end{aligned}$$

end for

Ilustration

• Consider sampling from a bivariate normal distribution.



• **Left plot**: The first 30 iterations of RWM with 20 updates per iterations. **Right plot**: The first 30 iterations of HMC with 20 leapfrog steps.

Improving HMC

- We still need to find optimum values of ε and L (i.e., optimum trajectory length, $\varepsilon \times L$).
- We also need to choose an appropriate mass matrix, M.
- The simulation requires a costly evaluation of the gradient of the energy function.

NUTS

- Hoffman and Gelman (2011) have tackled the first problem: optimum trajectory length.
- They have proposed a new approach called No-U-Turn Sampler (NUTS).
- NUTS uses a recursive algorithm to build a set of likely candidate points that spans a wide swath of the target distribution, stopping automatically when the trajectory starts to double back and retrace its steps.
- Their method is similar to the doubling procedure proposed by Neal (2003) for slice sampling.

RMHMC

- Girolami and Calderhead (2011) have tackled the second problem: optimum mass matrix.
- They have introduced a new method, called Riemannian Manifold HMC (RMHMC).
- It is more natural to put the Hamiltonian dynamic on Riemannian manifold of distributions rather than Euclidean space.
- They follow Amari (2000) and use the Fisher information matrix, G(q), as a metric on the manifold.
- That is, they use position specific mass matrix, M = G(q)

RMHMC

- This way, they reduce autocorrelation by exploiting the geometric properties of the parameter space.
- However, using position specific mass matrix, however, leads to non-separability of Hamiltonian on Riemannian Manifold.
- As a result, RMHMC becomes computationally intensive because simulating from Hamiltonian dynamics involves solving implicit equations, which require additional iterative numerical techniques (e.g., fixed-point iteration).

SGLD

- Further improvement in HMC efficiency can be achieved by using optimization routines within MCMC.
- Recently Welling and Teh (2011) proposed a new approach, called Stochastic Gradient Langevin Dynamics (SGLD), to reduce the computational cost of HMC.
- They combine Langevin dynamics (which can be considered as HMC with only one leapfrog step) with stochastic approximation theory.
- Their approach allows efficient use of mini-batches of data to take advantage of data redundancy.

Split HMC

- We also focus on reducing the computational cost of HMC.
- The computation is mainly dominated by gradient evaluations.
- We show how the technique of "splitting" the Hamiltonian (Leimkuhler and Reich, 2004) can be used to reduce the computational cost of HMC,

$$H(q,p) = H_1(q,p) + H_2(q,p) + \cdots + H_K(q,p)$$

• The leapfrog method in fact can be regarded as a symmetric splitting of the Hamiltonian H(q,p)=U(q)+K(p) as

$$H(q,p) = U(q)/2 + K(p) + U(q)/2$$

Split HMC with a partial analytic solution

- Suppose $U(q) = U_0(q) + U_1(q)$.
- Then, we can split H as

$$H(q,p) = U_1(q)/2 + [U_0(q) + K(p)] + U_1(q)/2$$

- Suppose the middle part can be handled analytically.
- Then its simulation introduces no error.
- We should be able to use a larger step size and fewer steps.

Split HMC with a partial analytic solution

- We approximate U(q) by $U_0(q)$.
- We set $U_1(q) = U(q) U_0(q)$, the error in this approximation.
- Specifically, we set $U_0(q)$ to the energy function for $N(\hat{q}, \mathcal{J}^{-1}(\hat{q}))$.
- Here, \hat{q} is the MAP estimate, and $\mathcal{J}(\hat{q})$ is the Hessian matrix of U at \hat{q} .

Split HMC with a partial analytic solution

- We set $H_2(q, p) = U_0(q) + K(p)$.
- The corresponding Hamilton's equations will be a system of first-order linear differential equations that can be handled analytically.
- We set

$$A = \begin{bmatrix} 0 & I \\ -\mathcal{J}(\hat{q}) & 0 \end{bmatrix}$$

• Then, we diagonalize $A = \Gamma D \Gamma^{-1}$.

Algorithm 1

$$\begin{split} R &\leftarrow \Gamma e^{D\varepsilon} \Gamma^{-1} \\ \text{Sample initial values for } p \text{ from } N(0,I) \\ \textbf{for } \ell &= 1 \text{ to } L \text{ do} \\ p &\leftarrow p - (\varepsilon/2) \frac{\partial U_1}{\partial q} \\ q^* &\leftarrow q - \hat{q} \\ X_0 &\leftarrow (q^*,p) \\ (q^*,p) &\leftarrow RX_0 \\ q &\leftarrow q^* + \hat{q} \\ p &\leftarrow p - (\varepsilon/2) \frac{\partial U_1}{\partial q} \end{split}$$

Split HMC by splitting the data

- Next, suppose $H_2(q, p)$ cannot be handled analytically.
- However, suppose that the computational cost for $U_0(q)$ is still substantially lower than for U(q).
- In these situations, we can use the following split:

$$H(q,p) = U_1(q)/2 + \sum_{m=1}^{M} [U_0(p)/2M + K(p)/M + U_0(p)/2M] + U_1(q)/2$$

Split HMC by splitting the data

- For example, suppose our statistical analysis involves a large data set with many observations.
- We can construct $U_0(q)$ based on a small part of the observed data, R_0 .
- We use the remaining observations, R_1 , to construct $U_1(q)$.

$$U(\theta) = U_0(\theta) + U_1(\theta)$$

$$U_0(\theta) = -\log[P(\theta)] - \sum_{i \in R_0} \log[P(y_i|\theta)]$$

$$U_1(\theta) = -\sum_{i' \in R_1} \log[P(y_{i'}|\theta)]$$

Algorithm 2

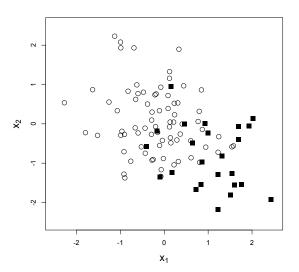
Sample initial values for p from N(0, I)

$$\begin{aligned} & \textbf{for } \ell = 1 \text{ to } L \quad \textbf{do} \\ & p \leftarrow p - (\varepsilon/2) \frac{\partial U_1}{\partial q} \\ & \textbf{for } m = 1 \text{ to } M \quad \textbf{do} \\ & p \leftarrow p - (\varepsilon/2M) \frac{\partial U_0}{\partial q} \\ & q \leftarrow q + (\varepsilon/M)p \\ & p \leftarrow p - (\varepsilon/2M) \frac{\partial U_0}{\partial q} \end{aligned}$$

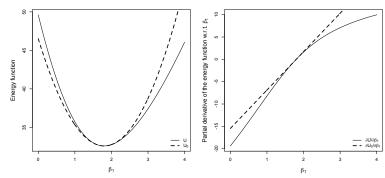
end for

$$p \leftarrow p - (\varepsilon/2) \frac{\partial U_1}{\partial q}$$

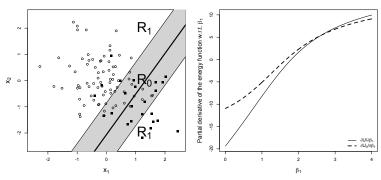
end for



- For Algorithm 1, we use the MAP estimates, $\hat{\theta}$, for model parameters θ .
- $U_0(\theta)$ is the potential energy function of the normal distribution $N(\hat{\theta}, \mathcal{J}^{-1}(\hat{\theta}))$



- For Algorithm 2, we use the MAP estimates, $\hat{\theta}$ and define U_0 based on the data points with high entropy.
- These are the points that are close to the classification boundary defined based on $\hat{\theta}$.



- Note that U_0 is not used to approximate U.
- Rather, $\partial U_0/\partial \beta_j$ is used to approximate $\partial U/\partial \beta_j$.
- Recall that

$$\frac{\partial U}{\partial \beta_j} = \frac{\beta_j}{\sigma_{\beta}^2} - \sum_{i=1}^n x_{ij} \left[y_i - \frac{\exp(\alpha + x_i^T \beta)}{1 + \exp(\alpha + x_i^T \beta)} \right]$$

- The term $\exp(\alpha + x_i^T \beta)/(1 + \exp(\alpha + x_i^T \beta))$ is in fact $P(y_i = 1 | x_i, \alpha, \beta)$.
- For high entropy data points, this estimated probability is close to 0.5.

Experiments

- We set the number of leapfrog steps to L=20 for the standard HMC, and find ε such that the acceptance probability (AP) is close to 0.65.
- We set L and ε for the Split HMC methods such that the trajectory length, εL , remains the same, but with a larger stepsize and hence a smaller number of steps.
- To measure the efficiency of each sampling method, we use the autocorrelation time (ACT).
- ACT can be roughly interpreted as the number of MCMC transitions required to produce samples that can be considered as independent.

Experiments- Simulated data

	НМС	Split HMC	
		Normal Appr.	Data Splitting
L	20	10	3
g	20	10	12.6
S	0.187	0.087	0.096
AP	0.69	0.74	0.74
τ	4.6	3.2	3.0
au imes g	92	32	38
$\begin{array}{c c} \tau \times g \\ \tau \times s \end{array}$	0.864	0.284	0.287

Experiments- StatLog

	НМС	Split HMC	
		Normal Appr.	Data Splitting
L	20	14	3
g	20	14	13.8
S	0.033	0.026	0.023
AP	0.69	0.74	0.85
τ	5.6	6.0	4.0
au imes g	112	84	55
au imes s	0.190	0.144	0.095

Experiments- Chess

	НМС	Split HMC	
		Normal Appr.	Data Splitting
L	20	9	2
g	20	13	11.8
S	0.022	0.011	0.013
AP	0.62	0.73	0.62
τ	10.7	12.8	12.1
au imes g	214	115	143
au imes s	0.234	0.144	0.161

Future directions

- Finding tractable approximations to the posterior distribution other than normal.
- Other methods for splitting the Hamiltonian dynamics by splitting the data,
- Combining our method with Riemannian Manifold HMC (Girolami and Calderhead, 2011) and No-U-Turn sampler (Hoffman and Gelman, 2011).