Vehicle Dynamics and Simulation

Using Eigenvalues and Eigenvectors
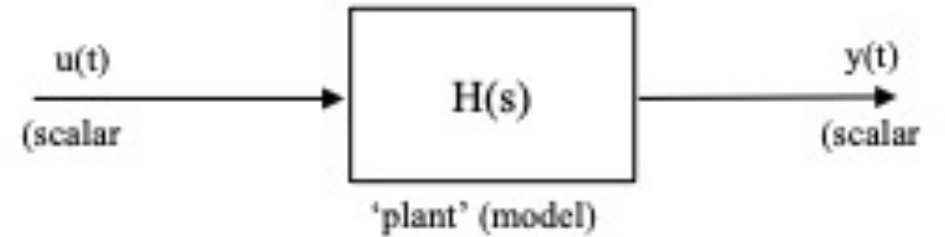
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Lecture overview

- Transfer functions

- Modal motion in free vibration

  - Eigenvalues

  - Eigenvectors

Transfer Functions

- Transfer functions relate input to output

- The roots of the characteristic equation determine the dynamics of the system (free vibration response)

- Eigen structures provide more information about mode shapes

$$H(s) = \frac{b_m s^m + b_{m-1} s^{m-1} + \cdots + b_1 s + b_0}{a_n s^n + a_{n-1} s^{n-1} + \cdots + a_1 s + a_0}$$

$$a_n s^n + a_{n-1} s^{n-1} + \cdots + a_1 s + a_0 = 0$$

Simple Example

- Using the simple suspension example from Section 2b
  - Assume zero initial conditions
  - Take Laplace transform
  - Write transfer function
  - Enter parameter values
- The roots of the characteristic equation in this example are;

$$-1.875 \pm 6.372i$$

- The nature of the roots e.g. complex, repeated, distinct and real determine the general solution approach.
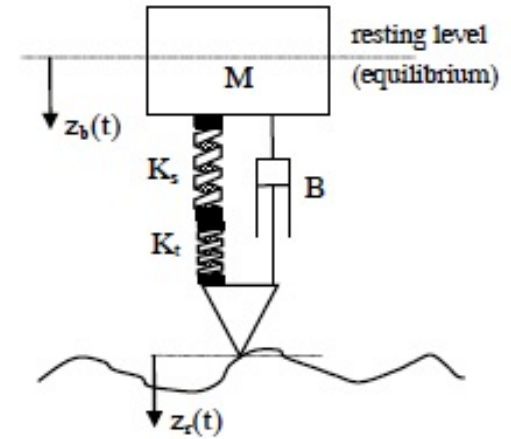
$$H(s) = \frac{Y(s)}{U(s)} = \frac{Bs + K}{Ms^2 + Bs + K}$$

$$H(s) = \frac{3.75s + 44.1}{s^2 + 3.75s + 44.1}$$

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Laplace Transform and the Transfer Function

- State space representation;

$$\dot{X} = AX + BU \qquad\qquad [1]$$

$$Y = CX + DU \qquad\qquad [2]$$

- Assuming zero initial conditions and taking the Laplace transform of [1];

$$sX = AX + BU$$

$$(sI - A)X = BU$$

$$X = (sI - A)^{-1}BU$$

- Substituting into [2];

$$Y = C(sI - A)^{-1}BU + DU$$

$$\boxed{H(s) = C(sI - A)^{-1}B + D} \qquad\qquad [3]$$

Laplace Transform and the Transfer Function

- Equation [3] provides a general solution in terms of the transfer function, $H(s)$ and is an alternate form to the State Space Representation.

- Comparing the denominator of [3] $(sI - A)$ with the definition for the eigenvalues;

$$Av = \lambda v \qquad\qquad [4]$$

- It can be shown that the transfer function poles are the roots of the Characteristic Equation and the eigenvalues of A which are found using $\det(A - \lambda I) = 0$

- Matrix A governs the fundamental modes of vibration i.e. how the system will freely vibrate as it settles after some initial disturbance – not the inputs.

worked example

An Example (written)
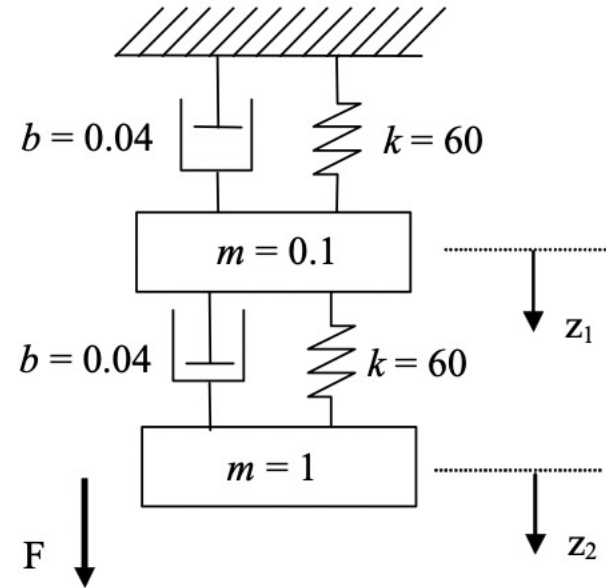
The equations of the above system are (expressed in terms of two second order differential equations);

$$\ddot{z}_2 = F = k(z_2 - z_1) - b(\dot{z}_2 - \dot{z}_1)$$

$$m\ddot{z}_1 = k(z_2 - z_1) + b(\dot{z}_2 - \dot{z}_1) - kz_1 - b\dot{z}_1$$

and the input, $u = F$. So that the set of equations describing the state derivatives becomes;

$$\dot{x}_1 = x_3$$

$$\dot{x}_2 = x_4$$

$$\dot{x}_3 = -\frac{2k}{m}x_1 + \frac{k}{m}x_2 - \frac{2b}{m}x_3 + \frac{b}{m}x_4$$

$$\dot{x}_4 = kx_1 - kx_2 + bx_3 - bx_4 + u$$

So that the state space representation becomes;

$$\begin{bmatrix} \dot{x}_1 \\ \dot{x}_2 \\ \dot{x}_3 \\ \dot{x}_4 \end{bmatrix} = \begin{bmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \\ -2k/m & k/m & -2b/m & b/m \\ k & -k & b & -b \end{bmatrix} \begin{bmatrix} x_1 \\ x_2 \\ x_3 \\ x_4 \end{bmatrix} + \begin{bmatrix} 0 \\ 0 \\ 0 \\ 1 \end{bmatrix} u \qquad [3]$$

$$\dot{X} = AX + BU$$

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Modal motion in free vibration – Eigenvalues

- The vector of deflections, $\boldsymbol{z(t)} = [x_1 \ x_2]^T$ for our LTI example may be written;

$$\boldsymbol{z(t)} = Re\{\boldsymbol{u}_1 e^{\lambda_1 t} + \boldsymbol{u}_2 e^{\lambda_2 t}\}$$

where each term $\boldsymbol{v}_i e^{\lambda_i t}$ represents a single vibrational mode, $\boldsymbol{v}_i$ are complex constants [2x1 vector in this example], λ_i are complex scalars and $\max(i) = n$ with n being the number of states.

- Evaluating a single term in the above, split λ_1 into real and imaginary parts;

$$\lambda_1 = a + bi$$

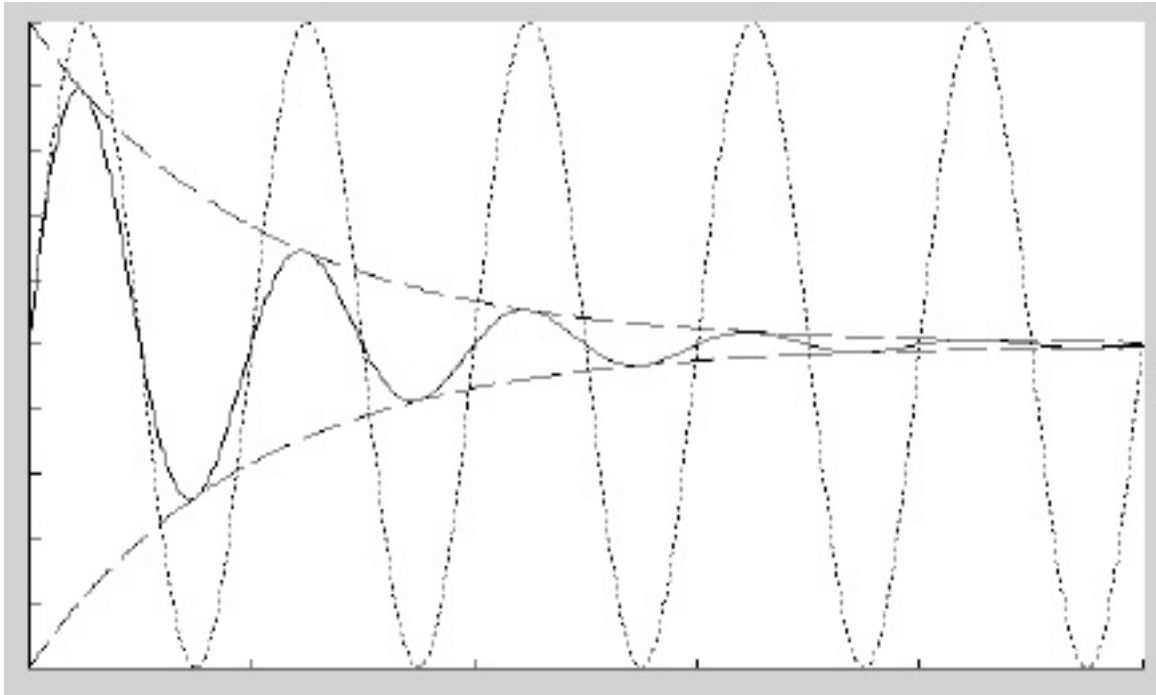- Using the above and Euler's formula we can better evaluate what is happening;

$$\boldsymbol{u}_1 e^{\lambda_i t} = \boldsymbol{u}_1 e^{(a+bi)t} = \boldsymbol{u}_1 e^{at} e^{ibt}$$

$$\boxed{\mathbf{v}_1 e^{\lambda_i t} = \boldsymbol{u}_1 e^{at}(\cos(bt) + i\sin(bt))}$$ 　　　[5]

Modal motion in free vibration - Eigenvalues

- From [5], a should be -ve bounding the response to a decaying exponential, b gives the frequency of the sinusoidal component, u_i (complex) determines the magnitude and the relative phase of each mode.

Modal decomposition of response
Solid line = total response
Short dash = sinusoidal component
Long dash = exponential decay

Modal motion in free vibration – Eigenvalues

- Eigenvalues come in (complex conjugate) pairs and can be written;

$$\lambda_{1,2} = \lambda_{1\&2} = \sigma \pm j\omega_d$$

where σ is the modal damping factor and ω_d is the damped natural frequency.

Check for yourself

- Use the eig function in MATLAB to determine the eigenvalues of matrix A from the previous example.

- How are the complex conjugate pairs placed within the resulting (vector) answer?

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Modal motion in free vibration – Eigenvalues

- From the eigenvalues it is possible to tell
  - Damped natural frequency, ω_d
  - Natural frequency [Hz], $\omega_n/2\pi$
  - Damping factor, σ
  - Damping ratio, $\zeta = \cos(\theta)$
  - Settling time (within 2%), $T_s = \dfrac{4}{|\sigma|}$
  - Percent overshoot, $100e^{\frac{-\pi\zeta}{\sqrt{1-\zeta^2}}}$
- Note: $\lambda = 0$ corresponds to the steady-state response of the system (not dynamics)
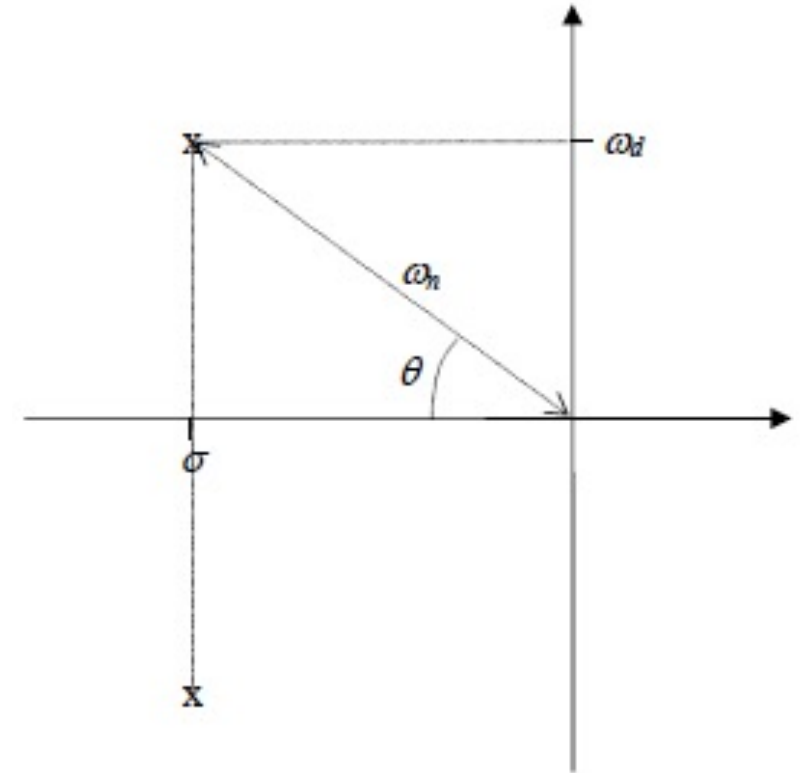
Check for yourself

- Using the previous example find the eigenvalues of the system and hence determine;
  - Damped natural frequency, ω_d
  - Natural frequency [Hz], $\omega_n / 2\pi$
  - Damping factor, σ
  - Damping ratio, $\zeta = \cos(\theta)$
  - Settling time (within 2%), $T_s = \dfrac{4}{|\sigma|}$
  - Percent overshoot, $100e^{\frac{-\pi\zeta}{\sqrt{1-\zeta^2}}}$

Modal motion in free vibration – Eigenvectors

- Eigenvectors can show the magnitudes at which the states vibrate in relation to one another.

- Writing eigenvalues and eigenvectors together in matrix form;

$$AV = VD$$

where;

$$V = \begin{bmatrix} u_1 & u_2 & \cdots & u_n \\ \lambda_1 u_1 & \lambda_2 u_2 & \cdots & \lambda_n u_n \end{bmatrix} \text{ and } D = \begin{bmatrix} \lambda_1 & & & \\ & \lambda_2 & & \\ & & \ddots & \\ & & & \lambda_n \end{bmatrix}$$

Modal motion in free vibration – Eigenvectors

- Using MATLAB to find the eigenvalues of the example system, A matrix;

$$V = \begin{bmatrix} -0.0284 - 0.0017i & -0.0284 + 0.0017i & 0.0665 + 0.0495i & 0.0665 - 0.0495i \\ 0.0015 + 0.0001i & 0.0015 - 0.0001i & 0.1298 + 0.0966i & 0.1298 - 0.0966i \\ 0.0700 + 0.9958i & 0.0700 - 0.9958i & -0.2685 + 0.3593i & -0.2685 - 0.3593i \\ -0.0036 - 0.0510i & -0.0036 + 0.0510i & -0.5236 + 0.7011i & -0.5236 - 0.7011i \end{bmatrix}$$

- Note
    - The second and fourth columns are the complex conjugates of the first and third columns respectively
    - Rows three and four are the first and second rows multiplied by their respective eigenvalues
    - The system can then be characterized by considerably less 'unique information'

Modal motion in free vibration – Eigenvectors

- Dividing through by the largest magnitude eigenvector ($-0.0284 - 0.017i$) to normalize the eigenvectors.

- Plot the eigenvector components (first mode)

- The relative magnitude and phase is seen on the two plots
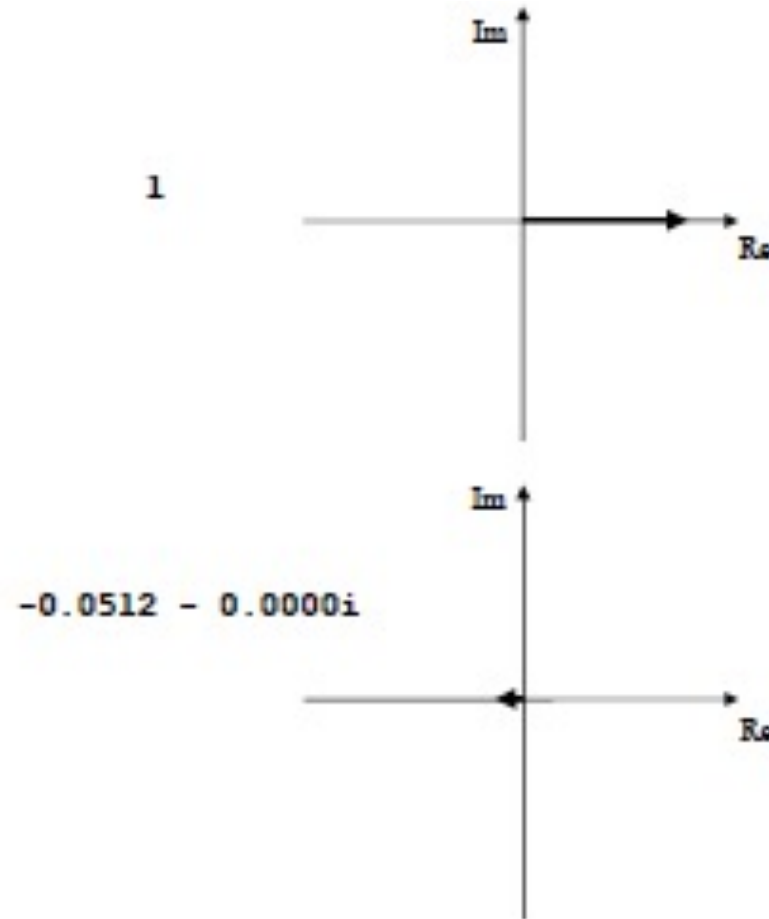
Modal motion in free vibration – Eigenvectors

- Similarly for the second (non-conjugate) mode of interest

- The relative magnitude and phase is seen on the two plots

- Note the differences between first and second modes of vibration

0.5125+0.0000i
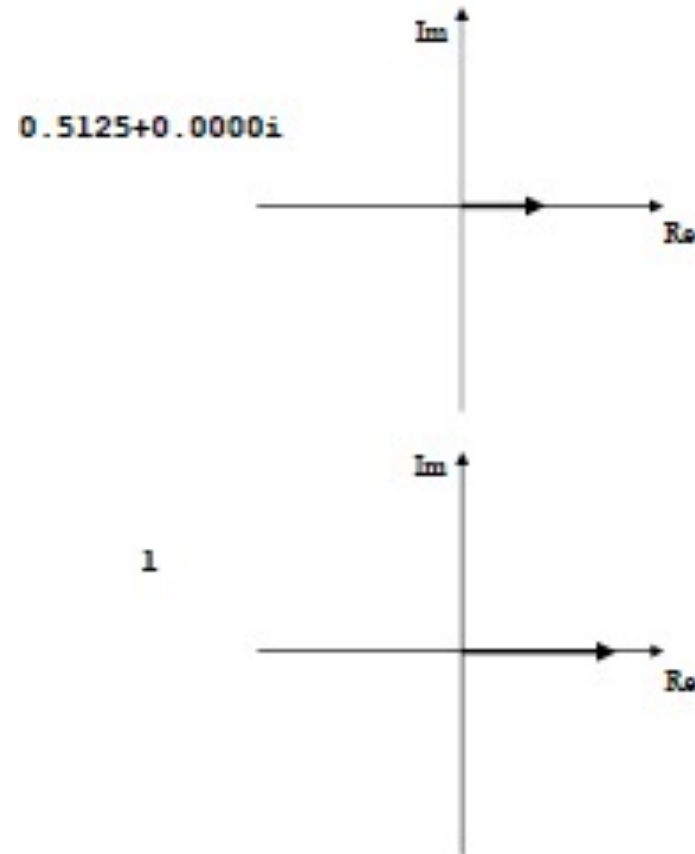
Conclusions

- Transfer function vs state space representation
- Eigenvalues tell us;
  - Damped natural frequency
  - Natural frequency
  - Damping factor
  - Damping ratio
  - Settling time
  - Percent overshoot
- Eigenvectors help us to understand vibration of the modes relative to one another