a) 
$$\frac{du}{dt} = -u + H(u - \frac{1}{2}) - \omega$$

$$\frac{dw}{dt} = \mathcal{E}(2u - \omega)$$
where 
$$H(u - \frac{1}{2}) = \begin{cases} 0, & u \leq \frac{1}{2} \\ 1, & u \geq \frac{1}{2} \end{cases}$$
Set 
$$u' = 0 \quad \text{ff} \quad \omega' = 0$$

$$0 = -u + H(u - \frac{1}{2}) - \omega$$

$$0 = -u + H(u-1/2) - \omega$$

$$0 = \mathcal{E}(\mathcal{E}u - \omega)$$

$$\omega = \mathcal{H}(u - 1/2) - \mathcal{U}$$

$$\omega = 2u$$

Solving for fix prints:  $2u = \mathcal{H}(u - 1/2) - u \qquad w = 2u$ 

Jacobian associated with its linear stability

$$\int (\bar{u}, \bar{\omega})^{2} \int \left( \begin{array}{ccc} F_{u} & F_{\omega} \\ G_{u} & G_{\omega} \end{array} \right) = \left( \begin{array}{ccc} \frac{\partial \mathcal{U}(u - \frac{1}{2})}{\partial u} - | \cdot \cdot \cdot - | \cdot | \\ 2\mathcal{E} & -\mathcal{E} \end{array} \right)$$

When 
$$u < \frac{1}{2}$$
:
$$2u = 0 - u \qquad \overline{w} = 0$$

$$\overline{u} = 0$$

$$2(0,0) = \begin{pmatrix} -1 & -1 \\ 2\mathcal{E} & -\mathcal{E} \end{pmatrix} = \begin{pmatrix} \lambda+1 & -1 \\ 2\mathcal{E} & \lambda+\mathcal{E} \end{pmatrix}$$

$$= (\lambda+1)(\lambda+\mathcal{E}) + 2\mathcal{E}$$

$$= \lambda^2 + \mathcal{E}\lambda + \lambda + \mathcal{E} + 2\mathcal{E}$$

$$= \lambda^2 + \lambda(1+\mathcal{E}) + 3\mathcal{E}$$

$$\lambda = -(1+\mathcal{E}) \pm \sqrt{(1+\mathcal{E})^2 - 12\mathcal{E}}$$

When 
$$u = 1/2$$
, we have a contradictory  $u$ .  
 $2u = 1 - u = 9$   $u = \frac{9}{3}$  when we said  $u = \frac{1}{2}$ , not provible

b) Assume 
$$u(0) = \frac{3}{4}$$
 and  $w(0) = 0$ .

$$\frac{du}{dt} = -u + H(u - \frac{1}{2}) \quad \frac{dw}{dt} = E(2u - w)$$
Since  $u(0) > \frac{1}{2}$ ,
$$\frac{du}{dt} = \frac{1}{4} \quad \text{and} \quad \frac{1}{4} \quad \text{acts much fister}$$
So,  $u = H(u - \frac{1}{2}) - w$ 

 $u = H(u - \frac{1}{2}) - \omega$   $w' = \mathcal{E}(2(1 - \omega) - \omega)$   $u = \frac{1}{2} + \frac{1}{2} + \frac{1}{2} + \frac{1}{2} = \frac{1}{2} + \frac{1}{2} = \frac{1}{2} + \frac{1}{2} = \frac{1}{2} = \frac{1}{2} + \frac{1}{2} = \frac{1}{2} =$ 

We see that U-7 1 b/c  $\frac{1}{4}$  acts much faster than E and our heaviside function will set  $U=4-\omega$ , where w(0)=0. Since  $\omega$  relies on E<<1,  $\omega$ ,  $\omega$ -71.

Once u reaches I, 'w' begins to climb slowly ble of E.

Our slow equation becomes:  

$$w' = \mathcal{E}\left[2\left(1-\omega\right)-\omega\right)$$

$$w' + 3\mathcal{E}\omega = \mathcal{E}2 \quad , \quad M = e^{3\mathcal{E}\mathcal{E}}$$

$$\int e^{3\mathcal{E}\mathcal{E}}w \frac{1}{dt} = \int \mathcal{C}2e^{3\mathcal{E}\mathcal{E}}dt$$

$$e^{3\mathcal{E}\mathcal{E}}w \frac{1}{dt} = \int \mathcal{C}2e^{3\mathcal{E}\mathcal{E}}dt$$

$$e^{3\mathcal{E}\mathcal{E}}w \frac{1}{dt} = \mathcal{E}2e^{3\mathcal{E}\mathcal{E}}dt$$

$$e^{3\mathcal{E}\mathcal{E}}w = \frac{\mathcal{E}2}{\mathcal{E}3}e^{3\mathcal{E}\mathcal{E}}dt$$

$$w = \frac{2}{3} + Ce^{3\mathcal{E}\mathcal{E}}, \quad w(0) = 0$$

$$0 = \frac{2}{3} + C = 7 \quad C = -\frac{2}{3}$$
So the slow equation for  $w$ :
$$w = \frac{2}{3}\left(1-e^{3\mathcal{E}\mathcal{E}}\right)$$
Since for  $w(T) = \frac{1}{2}$ 

$$\frac{1}{2} = \frac{2}{3}\left(1-e^{3\mathcal{E}\mathcal{E}}\right)$$

$$\frac{3}{4} = 1-e^{-3\mathcal{E}\mathcal{E}} = 1 \quad e^{3\mathcal{E}\mathcal{E}} = \frac{1}{4}$$

$$=7 - 3ET = \ln(\frac{1}{4})$$

$$T = \ln(\frac{1}{4}) \cdot \frac{1}{-3E}$$

$$T = \frac{\ln(4)}{3E}$$

When E is decreased, Treaches its ending of the spike stowly. This makes sense ble at W(T)=1/2 we shoot back to our w' and the newson slowly settles at its fixed pt.

when E is decreased we reach ending of spike slowly.