

## Julia-Scripts

A small collection of Julia utility scripts and their associated wrappers that I use. The shell scripts all use the environmental variable JPROGRAM to find the julia scripts, so make sure these are in a location given by the directory to which you set JPROGRAM.

### Convolving data with a Gaussian or Cauchy (Lorentz) distribution

For two functions  $f, g: R \rightarrow R$ , we can define their convolution as

$$(f * g)(x) \equiv \int_{-\infty}^{\infty} f(x') g(x - x') dx'$$

The assumption here is that the data to convolve represents the output of some function that is zero for non-positive arguments. The lower integration limit can then be truncated to 0.

$$(f * g)(x) \equiv \int_0^{\infty} f(x') g(x - x') dx'$$

I made this for functions  $f$  being scattering cross section functions  $\sigma(E)$ , where  $E$  is some positive energy. There are two possibilities coded for  $g$

#### The Gaussian Distribution

The gaussian distribution is given by

$$e^{-\frac{(x' - x)^2}{2\gamma^2}}$$

Cross sections  $f(x)$  can be convolved with this distribution:

$$\begin{aligned} \tilde{f}(x) &= \\ &= \frac{\int_0^{\infty} dx' \sigma(x') e^{-\frac{(x - x')^2}{2\gamma^2}}}{\int_0^{\infty} dx' e^{-\frac{(x - x')^2}{2\gamma^2}}} \\ &= \frac{1}{\gamma \sqrt{2\pi}} \int_0^{\infty} dx' \sigma(x') e^{-\frac{(x - x')^2}{2\gamma^2}} \end{aligned}$$

In implementation, the normalization factor  $\int_0^{\infty} dx' e^{-\frac{(x - x')^2}{2\gamma^2}}$  is computed numerically

#### The Cauchy (Lorentz) Distribution

The Cauchy — or Lorentz — distribution is given by

$$\frac{\gamma}{\pi} \frac{1}{\left[ (x - x')^2 + \gamma^2 \right]^{3/2}}$$

where  $\gamma$  is the distribution width.

#### Implementation

In practice, the upper integration limit is taken to be the last available value of  $x$ . This should not be an issue if the Gaussian width,  $\gamma$ , is small enough compared to this. The lower limit, 0, may not be in the domain of  $f$  or may not return a desirable value. Convolving such data must be done carefully, depending on the function  $f(x)$ .

The shell script `convolve` is a wrapper script for the Julia script `convolve.jl`. Script arguments

### Using `convolve`

`convolve`

Wrapper for my julia script located at : `\$JPROGRAM`

usage: `convolve [operation] [operand]`

operations:	operand:	function:
<code>-h</code>	none	show this message.
<code>-i, --input</code>	file	specify input file
<code>-o, --output</code>	file	specify output file
<code>--dx, --width</code>	float	specify gaussian width
<code>--nx</code>	integer	number of convolution x-grid points
<code>--logx</code>	none	if supplied, a logarithmic grid will be used
<code>-t, --ct, --convtype</code>	(gauss cauchy)	determines convolution function

The following is a valid implementation of `convolve` :

```
convolve -i data_to_convolve.dat -o convolved_data.dat --dx=1.3e-3 --nx=1000
```

### Convolve scattering electron collision cross sections with a Maxwell-Boltzmann distribution

In the situation where we have scattering cross sections (electron-atom or electron-molecule) as a function of collision energy and want to obtain thermally averaged rate coefficients as a function of the kinetic temperature, the script `thermal` can be used. Just as above, this reduces to convolving a function  $f(x)$  with a function  $g(x)$ . Here, the function  $f$  is assumed to be the function  $\sigma(E)$  — a scattering cross section as a function of collision energy ( $E$ ). The function  $\sigma(E)$  is zero for negative  $E$ , but may be zero or undefined at  $E = 0$ . In practice, the value  $E = 0$  is approached but not passed exactly. Some cross sections are 0 at threshold ( $E = 0$ ), while others behave as  $1/E$ . The choice of assumed behavior can be given to the script, whose usage is given below.

#### The Maxwell-Boltzmann Distribution

The Maxwell-Boltzmann distribution is a probability distribution, typically used for describing particle speeds in an ideal gas. In three dimensions, It is given by

$$\left( \frac{m}{2\pi kT} \right)^{\frac{3}{2}} e^{-\frac{mv^2}{kT}}$$

where  $v$  is the particle speed,  $m$  is the particle mass,  $k$  is Boltzmann's constant, and  $T$  is the kinetic temperature. Given cross sections  $\sigma(E)$ , we can use the Maxwell-Boltzmann distribution to obtain thermal rate coefficients (rate coefficients in the case where we have a gas of particles at a temperature  $T$ ). The cross sections are first convolved with the Maxwell-Boltzmann distribution to obtain rate coefficients at a given temperature, given by

$$\alpha(T) = \frac{\int_0^{\infty} \sigma(E) \sqrt{2E/m} \sqrt{E} e^{-E/kT} dE}{\int_0^{\infty} \sqrt{2E/m} \sqrt{E} e^{-E/kT} dE}$$

```

\int\limits_{0}^{\infty} \sqrt{E} e^{-E/kT} dE
}

```

which should then be averaged over initial states. However, this is not done by the script (this requires information like state energies and rotational quantum numbers, so this can be done separately [1]) The normalization factor is also computed numerically.

For cross sections behaving as  $1/E$ , the low-energy (near-0) part of the integral is very important, especially at lower temperatures. In the case where the cross section data is not available at a low enough energy, the cross sections can be extrapolated assuming a  $1/E$  behavior based on the datum corresponding to the lowest available energy. This is sensitive whether that point is a resonance and obviously does not take into account resonances that might be present at lower energies. Ideally, this is done with data that was obtained at low enough scattering energies such that there are no resonances. See the following for more detail on the usage.

### Using thermal

This script requires my script units (<https://github.com/banana-bred/units>).

thermal

```

Wrapper for my julia script (\$JPROGRAM), which produces state-selected
kinetic rate coefficients from cross sections behaving as 1/E (E being electron energy)
at the E=0 threshold.

```

```
usage: thermal [operation] [operand]
```

operations:	operand:	function:
-h	none	show this message.
-i,--input	file	specify input file
-o,--output	file	specify output file
--logx	none	if supplied, a logarithmic grid will be used
--Ti	float	lowest kinetic temperature (K)
--Tf	float	highest kinetic temperature (K)
--nT	integer	number of kinetic temperatures
--input-xs-units	string	the unit type of the input data (e.g., "cm^3 / s") Output rates will be in these units
--input-energy-units	string	the unit type of the input energy (e.g., "eV")
--extrap	none	extrapolate cross sections to E = 0 threshold assume it is zero.
--electron-energy-min	float	lowest electron energy (should be closer to 0)
--num-extrap-energies	integer	number of extrapolation energies (extrapolation points)

The following would be a valid implementation of thermal :

```

thermal -i data_to_convolve.dat -o convolved_data.dat \
    --logx \
    --Ti=1e-6 \
    --Tf=1e3 \
    --extrap \
    --electron-energy-min=1e-8 \
    --num-extrap-energies=1000 \

```

```
--nT=500 \
--input-xs-units=cm \
--input-energy-units=eV
```

Input file: a file of electron energies ( $E_{el}$ ) and cross sections ( $\sigma$ ) in any units of energy and length<sup>2</sup>, formatted as

```
Eel  σ
.   .
.   .
.   .
```

Output file: a file of temperatures ( $T$ ) in K and rate coefficients ( $\alpha$ ) in cm<sup>3</sup>/s

```
T  α
.  .
.  .
.  .
```

### Fitting interatomic / intermolecular potentials

Given data that (hopefully) resembles an interatomic or intermolecular potential, the script `morseFit` can be used to fit the data to a Morse potential. One case in which this might be useful is when such a potential is calculated and needs to be extrapolated to larger distances.

#### The Morse Potential

The Morse potential is often used as an improved approximation to the harmonic oscillator model for molecular vibration. It is given by

$$V(r) = D_e \left( e^{-2a(r - r_e)} - 2e^{-a(r - r_e)} \right) + E_{\text{diss}}$$

where  $r$  is a separation distance,  $D_e$  is the well depth with respect to the dissociation limit,  $r_e$  is the equilibrium distance, and  $E_{\text{diss}}$  is the dissociation limit of the potential.

#### Implementation

The scripts take as input a file containing space separated  $x(r)$ ,  $y(V(r))$  data. A fit can be generate of all of the data or only a part of the data for  $r > r_{\min}$ . This is useful for fitting only the tail of the data. The optimized parameters (using the method of least squares) can then be used to generate data over a user-defined range of  $r$  values. Should the `--tailonly` option be supplied, the script will append the fitted data to the supplied data. A decent guess for  $E_{\text{diss}}$  should be supplied, otherwise the least squares method might not work.

#### Using `morseFit`.

`morseFit`

Wrapper for my julia script located at (`\$JPROGRAM`)

usage: `morseFit [operation] [operand]`

operations:                      operand:                      function:

-h	none	show this message.
-i,--input	file	specify input file
-o,--output	file	specify output file
--tailonly	none	if specified, only fit the tail of t
--rcut	number	cutoff distance for fitting only t
--rmin	number	smallest distance for the fit. Ignor
--rmax	number	largest distance for the fit
--rstep	number	linear step size for the fit
-a	number	initial guess for the 'a' parameter
-r0	number	initial guess for the 'r0' parameter
-D	number	initial guess for the 'D' parameter
--dissociation-limit	number	initial guess for the absolute value
--positive-dissoc-limit	none	if specified, the dissociation limit
--print-parameters	none	if specified, print the fit paramete

The following is a valid implementation of morseFit :

```
morseFit -i input.dat -o output.dat --tailonly --rcut=2.6 --rstep=0.01 --rma
```

For a full fit: morseFit -i input.dat -o output.dat --rmin=0.8  
--rstep=0.01 --rmax=10 --dissociation-limit 75 --positive-dissoc-  
limit

For a tail fit: morseFit -i input.dat -o output.dat --rcut=2.6  
--rstep=0.01 --rmax=10 --dissociation-limit 75 --positive-dissoc-  
limit

## References

[1] Forer, J. et. al. (2023) *Kinetic rate coefficients for electron-driven collisions with CH: dissociative recombination and rovibronic excitation*. Monthly Notices of the Royal Astronomical Society 527, 5238-5234