

2016~2023

日本综合研究院大学

前：物理；后：数学

KEK 素粒子专攻

一般入试过去问试题尝试解答

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by *heyingqiu*

主页：heyingqiu.github.io

邮箱：he.yingqiu@hotmail.com

物理

第 1 問

[問 1]

$$[\hat{x}, \hat{p}] = \hat{x}\hat{p} - \hat{p}\hat{x} = i\hbar$$

$$\hat{a}^{\dagger} \hat{a} = \frac{mw}{2\hbar} (\hat{x}^2 + \frac{\hat{p}^2}{m^2 w^2} - \frac{i}{mw} \hat{x}\hat{p} + \frac{i}{mw} \hat{p}\hat{x}) = \frac{mw}{2\hbar} (\hat{x}^2 + \frac{\hat{p}^2}{m^2 w^2} - \frac{\hbar}{mw})$$

$$\Rightarrow \hat{x}^2 + \frac{\hat{p}^2}{m^2 w^2} = \frac{2\hbar}{mw} \hat{a}^{\dagger} \hat{a} + \frac{\hbar}{mw}$$

$$\hat{H} = \frac{1}{2} mw^2 (\hat{x}^2 + \frac{\hat{p}^2}{m^2 w^2}) = \frac{1}{2} mw^2 \left(\frac{2\hbar}{mw} \hat{a}^{\dagger} \hat{a} + \frac{\hbar}{mw} \right) = \hbar w (\hat{a}^{\dagger} \hat{a} + \frac{1}{2})$$

[問 2]

$$\hat{H}\phi_n = E_n \phi_n \rightarrow \hbar w (\hat{a}^{\dagger} \hat{a} + \frac{1}{2}) \phi_n = E_n \phi_n$$

$$\begin{aligned} \text{l.h.s.} &= \hbar w \hat{a}^{\dagger} \hat{a} \phi_n + \frac{1}{2} \phi_n = \hbar w (\sqrt{n} \hat{a}^{\dagger} \phi_{n-1} + \frac{1}{2} \phi_n) = \hbar w (\sqrt{n} \sqrt{n} \phi_n + \frac{1}{2} \phi_n) \\ &= \hbar w (n + \frac{1}{2}) \phi_n \end{aligned}$$

$$\text{Therefore, } E_n = \hbar w (n + \frac{1}{2})$$

[問 3]

$$\begin{aligned} \hat{a}\phi(x) &= \sqrt{\frac{mw}{2\hbar}} \left[x + \frac{i}{mw} (-i\hbar) \frac{\partial}{\partial x} \right] A \exp(-\frac{mw}{2\hbar} x^2 + ikx) \\ &= \sqrt{\frac{mw}{2\hbar}} \left[x + \frac{\hbar}{mw} \left(-\frac{mwx}{\hbar} + ik \right) \right] \phi(x) \\ &= \sqrt{\frac{mw}{2\hbar}} \cdot ik \phi(x) \end{aligned}$$

$$\Rightarrow \alpha = ik \sqrt{\frac{mw}{2\hbar}}$$

[問 4]

Let $C^0 = 1$, and we have $\hat{a}\phi_0(x) = 0$

$$\phi(x) = B \left[\sum_{n=0}^{\infty} \frac{1}{(n+1)!} C^{n+1} \phi_{n+1}(x) + C^0 \phi_0(x) \right]$$

$$\hat{a}\phi(x) = B \left[\sum_{n=0}^{\infty} \frac{1}{(n+1)!} C^{n+1} \sqrt{n+1} \phi_n(x) \right]$$

Since $\hat{a}\phi(x) = B \left[\sum_{n=0}^{\infty} \frac{1}{n!} C^n x \phi_n(x) \right]$, therefore $\frac{1}{(n+1)!} C^{n+1} = \frac{1}{n!} C^n x$,

$$C^{n+1} = \sqrt{n+1} \propto C^n \Rightarrow C^n = \sqrt{n!} \propto n^{n-1} \quad (n > 1)$$

$$\phi(x) = B \left[\sum_{n=1}^{\infty} n^{-1} \phi_n(x) + \phi_0(x) \right]$$

.....

第 2 問

[問 1]

$$\hat{H}\psi(x, y, z) = E\psi(x, y, z)$$

$$\begin{aligned} \text{l.h.s.} &= -\frac{\hbar^2}{2m} \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \right) \left(\frac{2}{L} \right)^{\frac{3}{2}} \sin\left(\frac{\pi n_x x}{L}\right) \sin\left(\frac{\pi n_y y}{L}\right) \sin\left(\frac{\pi n_z z}{L}\right) \\ &= \frac{\hbar^2}{2m} \left(\frac{2}{L} \right)^{\frac{3}{2}} \left[\left(\frac{\pi n_x}{L} \right)^2 + \left(\frac{\pi n_y}{L} \right)^2 + \left(\frac{\pi n_z}{L} \right)^2 \right] \sin\left(\frac{\pi n_x x}{L}\right) \sin\left(\frac{\pi n_y y}{L}\right) \sin\left(\frac{\pi n_z z}{L}\right) \\ &= \frac{\hbar^2 \pi^2}{2m L^2} (n_x^2 + n_y^2 + n_z^2) \psi(x, y, z) \end{aligned}$$

$$\text{Therefore, } E(n_x, n_y, n_z) = \frac{\hbar^2 \pi^2}{2m L^2} (n_x^2 + n_y^2 + n_z^2)$$

[問 2]

$$E = \sum_{j=1}^N E^{(j)} = \frac{\hbar^2 \pi^2}{2m L^2} \sum_{j=1}^N (n_x^{(j)} + n_y^{(j)} + n_z^{(j)})^2 = \frac{\hbar^2 \pi^2}{2m L^2} \sum_{j=1}^{3N} n^{(j)} \text{, assume } n_x = n_y = n_z = N$$

$$\begin{aligned} Z_{V,N}(T) &= \frac{1}{N!} \sum_{n=1}^{\infty} \exp \left(-\frac{\frac{\hbar^2 \pi^2}{2m L^2} \sum_{j=1}^{3N} n^{(j)} \text{ }}{k_B T} \right) \\ &= \frac{1}{N!} \prod_{j=1}^{3N} \sum_{n=1}^{\infty} \exp \left(-\frac{\frac{\hbar^2 \pi^2}{2m L^2} n^{(j)} \text{ }}{k_B T} \right) \\ &= \frac{1}{N!} \left[\prod_{n=1}^{\infty} \exp \left(-\frac{\frac{\hbar^2 \pi^2}{2m L^2} n^2 \text{ }}{k_B T} \right) \right]^{3N} \end{aligned}$$

" $\frac{1}{N!}$ " is caused by particles are identical with each other

$$\text{Besides, } \varepsilon_0 = \frac{\hbar^2 \pi^2}{2m L^2}$$

[問 3]

$$\begin{aligned} F(T, V, N) &= -k_B T \ln Z_{V,N}(T) \\ &= -k_B T \cdot 3N \left[\ln \left(\frac{1}{N!} \right) + \ln \sum_{n=1}^{\infty} e^{-\frac{\frac{\hbar^2 \pi^2}{2m L^2} n^2}{k_B T}} \right] \\ &\approx 3k_B T N \left[N(\ln N - 1) - \ln \sum_{n=1}^{\infty} e^{-\frac{\frac{\hbar^2 \pi^2}{2m L^2} n^2}{k_B T}} \right] \end{aligned}$$

$$F(T, \lambda V, \lambda N) = ?$$

第3問

[P6] 1]

$$(1) \begin{cases} \dot{x} = r \cos\theta - r\dot{\theta} \sin\theta \\ \dot{y} = r \sin\theta + r\dot{\theta} \cos\theta \end{cases} \quad T = \frac{1}{2}m(x^2 + y^2) = \frac{1}{2}m(r^2 + r^2\dot{\theta}^2)$$

$$(2) \mathcal{L} = \frac{1}{2}m(r^2 + r^2\dot{\theta}^2) - U(r)$$

$$\frac{\partial \mathcal{L}}{\partial r} = m\ddot{r}, \frac{\partial \mathcal{L}}{\partial \dot{r}} = mr\dot{\theta}^2 - \frac{dU}{dr} \Rightarrow m\ddot{r} = mr\dot{\theta}^2 - \frac{dU}{dr} = \frac{\ell^2}{mr^3} - \frac{dU}{dr}$$

$$\frac{\partial \mathcal{L}}{\partial \theta} = mr^2\dot{\theta}, \frac{\partial \mathcal{L}}{\partial \dot{\theta}} = 0 \Rightarrow mr\ddot{\theta} = 0 \Rightarrow \frac{d\ell}{dt} = 0 \Rightarrow \ell \text{ is conserved}$$

$$\text{And } \frac{dE}{dt} = mr\ddot{r} + \frac{\ell^2}{mr^3}\dot{r} \stackrel{(2)}{=} \frac{\ell^2}{mr^3} + \frac{dU}{dr} \cdot \dot{r} \\ = \dot{r} \left[m\ddot{r} - \left(\frac{\ell^2}{mr^3} - \frac{dU}{dr} \right) \right] \\ = 0 \Rightarrow E \text{ is conserved}$$

[P6] 2]

$$\ell = mr^2 \frac{d\theta}{dt} \rightarrow d\theta = \frac{\ell}{m} u^2 dt \quad du = -u^2 dr \rightarrow \frac{dr}{dt} = -\frac{du}{u^2 dt}$$

$$E = \frac{1}{2}m \left(-\frac{du}{u^2 dt} \right)^2 + \frac{\ell^2 u^2}{2m} - k\ell u \quad \frac{du}{dt} = \frac{du}{d\theta} \frac{d\theta}{dt} = \frac{\ell}{m} u^2 \frac{du}{d\theta}$$

$$= \frac{1}{2}m \left(\frac{\ell}{m} u^2 \frac{du}{d\theta} \right)^2 + \frac{\ell^2 u^2}{2m} - k\ell u$$

$$\Rightarrow \frac{2mE}{\ell^2} + \frac{2mk}{\ell^2} u - u^2 = \frac{2m}{\ell^2} \cdot \frac{m}{2} \cdot \frac{\ell^2}{m^2} \left(\frac{du}{d\theta} \right)^2$$

$$\Rightarrow \left(\frac{du}{d\theta} \right)^2 = -u^2 + \frac{2mk}{\ell^2} u + \frac{2mE}{\ell^2} \quad \text{Since } \dot{r} > 0, \dot{\theta} < 0$$

$$\Rightarrow d\theta = -\frac{du}{\sqrt{-u^2 + \frac{2mk}{\ell^2} u + \frac{2mE}{\ell^2}}}$$

(2)

$$\int \frac{du}{\sqrt{-u^2 + \frac{2mk}{\ell^2} u + \frac{2mE}{\ell^2}}} = -\arccos \left(\frac{2u - \frac{2mk}{\ell^2}}{\sqrt{\left(\frac{2mk}{\ell^2}\right)^2 + \frac{8mE}{\ell^2}}} \right) + C$$

$$\text{We can have } \cos(\theta - \theta_0) = \frac{u - \frac{mk}{\ell^2}}{\sqrt{\left(\frac{mk}{\ell^2}\right)^2 + \frac{2mE}{\ell^2}}}$$

$$\Rightarrow \frac{mk}{\ell^2} + \frac{mk}{\ell^2} \sqrt{1 + \frac{2mE}{mk^2}} \cos(\theta - \theta_0) = \ell \dot{\theta}$$

$$\Rightarrow r = \frac{\ell^2}{mk} \frac{1}{1 + \sqrt{1 + \frac{2mE}{mk^2}} \cos(\theta - \theta_0)}$$

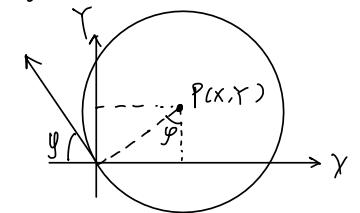
$$(3) E = \sqrt{1 + \frac{2\ell^2}{mk^2} - E}$$

(i) $E = 0, E = \frac{-mk^2}{2\ell^2}$, orbit: circle

(ii) $0 < E < 1, -\frac{mk^2}{2\ell^2} < E < 0$, orbit: ellipse

(iii) $E = 1, E = 0$, orbit: parabola

(iv) $E > 1, E > 0$, orbit: hyperbola



第4問

[P7] 1]

$$(1) \oint \vec{B} \cdot d\vec{l} = 2\pi \sqrt{x^2 + y^2} |\vec{B}^\circ| = \mu_0 I$$

$$\Rightarrow |\vec{B}^\circ| = \frac{\mu_0 I}{2\pi \sqrt{x^2 + y^2}}$$

$$(2) B_x^\circ = -|\vec{B}^\circ| \cos\psi = -\frac{y}{\sqrt{x^2 + y^2}} |\vec{B}^\circ| = -\frac{\mu_0 y I}{2\pi \sqrt{x^2 + y^2}}$$

$$B_y^\circ = |\vec{B}^\circ| \sin\psi = \frac{x}{\sqrt{x^2 + y^2}} |\vec{B}^\circ| = \frac{\mu_0 x I}{2\pi \sqrt{x^2 + y^2}}$$

[P7] 2]

$$(1) -\frac{d\vec{B}}{dt} = k$$

$$\int \vec{r} \times \vec{E} \cdot d\vec{l} = k \cdot \pi r^2$$

$$\oint \vec{E} \cdot d\vec{l} = |\vec{E}| \cdot 2\pi r$$

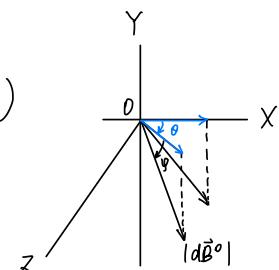
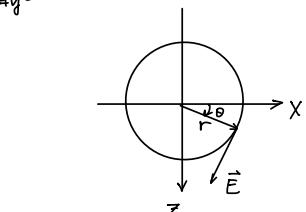
$$\Rightarrow |\vec{E}| = \frac{kr}{2}, \vec{E} = \left(-\frac{kr}{2} \sin\theta, -\frac{kr}{2} \cos\theta \right)$$

$$(2) \vec{I} = \alpha \vec{j} = \alpha \sigma \vec{E} = \left(-\frac{kra\sigma}{2} \sin\theta, -\frac{kra\sigma}{2} \cos\theta \right)$$

$$(3) |\vec{I}| = \frac{\alpha \sigma kr}{2}$$

$$\begin{aligned} dB_x^\circ &= |\vec{dB}^\circ| \cos\psi \cdot \sin\theta \\ &= \frac{1}{4\pi} \frac{|\vec{I}| r d\theta}{r^2 + y^2} \frac{y}{\sqrt{r^2 + y^2}} \sin\theta \\ &= \frac{|\vec{I}| r y}{4\pi (r^2 + y^2)^{\frac{3}{2}}} \sin\theta d\theta \end{aligned}$$

$$dB_y^\circ = -|\vec{dB}^\circ| \sin\psi = -\frac{|\vec{I}| r^2}{4\pi (r^2 + y^2)^{\frac{3}{2}}} d\theta$$



$$dB_z^o = (dB^o) \cos\theta \cos\phi = \frac{\mu_0 |\vec{I}| r y}{4\pi(r^2+y^2)^{\frac{3}{2}}} \cos\theta d\theta$$

$$(4) B_x^o = \frac{\mu_0 |\vec{I}| r y}{4\pi(r^2+y^2)^{\frac{3}{2}}} \int_0^{2\pi} \sin\theta d\theta = 0, \quad B_z^o = 0$$

$$B_y^o = -\frac{\mu_0 |\vec{I}| r^2}{4\pi(r^2+y^2)^{\frac{3}{2}}} \int_0^{2\pi} d\theta = -\frac{\mu_0 |\vec{I}| r^2}{2(r^2+y^2)^{\frac{3}{2}}} = -\frac{\alpha \sigma \mu_0 r^3}{4(r^2+y^2)^{\frac{3}{2}}}$$

$$\vec{B} = (0, -\frac{\alpha \sigma \mu_0 r^3}{4(r^2+y^2)^{\frac{3}{2}}}, 0)$$

數學

[P9|1]

$$A^2 = \begin{pmatrix} 2 & 3 & 2 \\ 2 & 4 & -2 \\ 2 & 5 & -2 \end{pmatrix} \begin{pmatrix} 2 & -3 & 2 \\ 2 & 4 & -2 \\ 2 & 5 & -2 \end{pmatrix} = \begin{pmatrix} 2 & 4 & -2 \\ 0 & 0 & 0 \\ 2 & 4 & -2 \end{pmatrix}$$

$$A^3 = \begin{pmatrix} 2 & 3 & 2 \\ 2 & 4 & -2 \\ 2 & 5 & -2 \end{pmatrix} \begin{pmatrix} 2 & 4 & 2 \\ 0 & 0 & 0 \\ 2 & 4 & 2 \end{pmatrix} = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

[P9|2]

$$\begin{aligned} \exp(xA) &= \sum_{n=0}^{\infty} \frac{1}{n!} (Ax)^n = I + Ax + \frac{A^2}{2}x^2 \\ &= \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix} + \begin{pmatrix} 2 & 3 & 2 \\ 2 & 4 & -2 \\ 2 & 5 & -2 \end{pmatrix}x + \begin{pmatrix} 1 & 2 & -1 \\ 0 & 0 & 0 \\ 1 & 2 & -1 \end{pmatrix}x^2 \\ &= \begin{pmatrix} 1-2x+x^2 & -3x+2x^2 & 2x-x^2 \\ 2x & 1+4x & -2x \\ 2x+x^2 & 5x+2x^2 & 1-2x-x^2 \end{pmatrix} \end{aligned}$$

[P9|3]

Let $F(x) = \begin{pmatrix} f(x) \\ g(x) \\ h(x) \end{pmatrix}$, we know $\frac{dF(x)}{dt} = AF$, then $F(x) = e^{Ax} C$

Since $\begin{pmatrix} f(0) \\ g(0) \\ h(0) \end{pmatrix} = \begin{pmatrix} a \\ b \\ c \end{pmatrix}$, therefore $C = \begin{pmatrix} a \\ b \\ c \end{pmatrix}$, and

$$\begin{pmatrix} f(x) \\ g(x) \\ h(x) \end{pmatrix} = \begin{pmatrix} 1-2x+x^2 & -3x+2x^2 & 2x-x^2 \\ 2x & 1+4x & -2x \\ 2x+x^2 & 5x+2x^2 & 1-2x-x^2 \end{pmatrix} \begin{pmatrix} a \\ b \\ c \end{pmatrix}$$

$$\Rightarrow \begin{cases} f(x) = a + (-2a-3b+2c)x + (a+2b-c)x^2 \\ g(x) = b + (2a+4b-2c)x \\ h(x) = c + (2a+5b-2c)x + (a+2b-c)x^2 \end{cases}$$

第2題

[P9|1]

$$\mathcal{L}[e^{at}](s) = \int_0^\infty dt e^{at} e^{-st} dt = \frac{1}{s-a} e^{(a-s)t} \Big|_0^\infty = \frac{1}{s-a} \quad (s>a)$$

$$\mathcal{L}[\sin t](s) = \int_0^\infty dt \sin t e^{-st} dt = -\frac{1}{s} e^{-st} \sin t \Big|_0^\infty + \frac{1}{s} \int_0^\infty \cos t e^{-st} dt$$

$$= \frac{1}{s} \left[-\frac{1}{s} e^{-st} \cos t \Big|_0^\infty + \frac{1}{s} \int_0^\infty -\sin t e^{-st} dt \right]$$

$$\Rightarrow \int_0^\infty dt \sin t e^{-st} dt (1+\frac{1}{s^2}) = -\frac{1}{s^2} (0-1)$$

$$\Rightarrow \int_0^\infty dt \sin t e^{-st} dt = \frac{1}{1+s^2}$$

$$\mathcal{L}[\cos t](s) = \int_0^\infty dt \cos t e^{-st} dt = \sin t e^{-st} \Big|_0^\infty - \int_0^\infty (-s) \sin t e^{-st} dt$$

$$= s \left[-\cos t e^{-st} \Big|_0^\infty + (-s) \int_0^\infty \cos t e^{-st} dt \right]$$

$$\Rightarrow \int_0^\infty dt \cos t e^{-st} dt (1+s^2) = (-s)(0-1)$$

$$\Rightarrow \int_0^\infty dt \cos t e^{-st} dt = \frac{s}{1+s^2}$$

[P9|2]

$$\mathcal{L}[f'(t)] = \int_0^\infty dt f'(t) e^{-st} = f(t) e^{-st} \Big|_0^\infty - (-s) \int_0^\infty dt f(t) e^{-st} = -f(0) + s \mathcal{L}[f(t)]$$

$$\begin{aligned} \mathcal{L}[f''(t)] &= -f'(0) + s \mathcal{L}[f(t)] = -f'(0) + s[-f(0) + s \mathcal{L}[f(t)]] \\ &= -f'(0) - s f(0) + s^2 \mathcal{L}[f(t)] \end{aligned}$$

Note

Since $f(0)=0, f'(0)=1$, then $\mathcal{L}[f'(t)] = s \mathcal{L}[f(t)]$, $\mathcal{L}[f''(t)] = -1 + s^2 \mathcal{L}[f(t)]$

$$\mathcal{L}[f''(t) - 3f'(t) + 2f(t)] = -1 + s^2 \mathcal{L}[f(t)] - 3s \mathcal{L}[f(t)] + 2 \mathcal{L}[f(t)]$$

$$= (s^2 - 3s + 2) \mathcal{L}[f(t)] + 1$$

$$\text{And } \mathcal{L}\left[\frac{1}{s^2+1}\right] = \frac{1}{s^2+1}, \text{ then } \mathcal{L}[f(t)] = \frac{-s^2}{(s^2+1)(s-1)(s-2)} = \frac{\frac{3}{10}s + \frac{1}{10}}{s^2+1} + \frac{\frac{1}{2}}{s-1} + \frac{-\frac{4}{5}}{s-2}$$

$$A_1 = \frac{1}{2}, A_2 = -\frac{4}{5}, B = \frac{3}{10}, C = \frac{1}{10}, S_1 = 1, S_2 = 2$$

[P9|3]

$$f(t) = A_1 e^{st} + A_2 e^{S_2 t} + B \cos t + C \sin t$$

$$= \frac{1}{2} e^{st} - \frac{4}{5} e^{S_2 t} + \frac{3}{10} \cos t + \frac{1}{10} \sin t$$

第3題

$$[\text{問}1] \quad \bar{\partial}f = \bar{\partial}g = g(z+\bar{z}\bar{z}), \bar{g}(z, \bar{z}) - g(z, \bar{z}) \quad d(re^{i\theta}) = e^{i\theta}dr + ire^{i\theta}d\theta$$

$$S = \lim_{r \rightarrow 0} \frac{\bar{\partial}f}{\bar{\partial}z} = \lim_{r \rightarrow 0} \frac{\bar{\partial}u + i\bar{\partial}v}{e^{i\theta} \bar{\partial}r} = e^{-i\theta} \frac{df}{dr}$$

[問2]

$$\begin{aligned} \lim_{\substack{x \rightarrow 0 \\ y=0}} \frac{\bar{\partial}f}{\bar{\partial}z} &= \lim_{x \rightarrow 0} \frac{\bar{\partial}u + i\bar{\partial}v}{\bar{\partial}x} = \frac{\partial u}{\partial x} + i \frac{\partial v}{\partial x} \\ \lim_{\substack{y \rightarrow 0 \\ x=0}} \frac{\bar{\partial}f}{\bar{\partial}z} &= \lim_{y \rightarrow 0} \frac{\bar{\partial}u + i\bar{\partial}v}{i\bar{\partial}y} = -i \frac{\partial u}{\partial y} + \frac{\partial v}{\partial y} \end{aligned} \Rightarrow \begin{cases} \frac{\partial u}{\partial x} = \frac{\partial v}{\partial y} \\ \frac{\partial u}{\partial y} = -\frac{\partial v}{\partial x} \end{cases}$$

[問3]

Equate $\frac{df}{dz}$ in x and y , $\frac{dx+idz}{dx} = \frac{du+idv}{-idy}$,
 $\frac{df}{dz}$ does not exist unless the equation are satisfied.

[問4]

$$\frac{\partial u}{\partial x} = 2x+3 \Rightarrow v(x, y) = (2x+3)y + q_0, \quad q_0 = q_0(x)$$

$$\frac{\partial u}{\partial y} = -2y+1 \Rightarrow v(x, y) = (-2y+1)x + p, \quad p = p(y)$$

$$\text{Then, } v(x, y) = 2xy + 3y + x$$

$$\text{And } h(z) = x^2 - y^2 + 3xy + y + i(2xy + x + 3y)$$

[問5]

$$\frac{\partial u}{\partial x} = \cos x \Rightarrow v(x, y) = (\cos x)y + l_0$$

$$\frac{\partial u}{\partial y} = 0 \Rightarrow v(x, y) = x + p \quad > \text{obviously, impossible.}$$

第4問

[問1]

$$dx = \frac{\partial x}{\partial r} dr + \frac{\partial x}{\partial \theta} d\theta = \cos \theta dr - r \sin \theta d\theta$$

$$dy = \frac{\partial y}{\partial r} dr + \frac{\partial y}{\partial \theta} d\theta = \sin \theta dr + r \cos \theta d\theta$$

$$Ax(x, y) dx + Ay(x, y) dy = (r^2 \cos^2 \theta - r^2 \sin^2 \theta - 3r \cos \theta)(\cos \theta dr - r \sin \theta d\theta) + (2r^2 \cos \theta \sin \theta - 3r \sin \theta)(\sin \theta dr + r \cos \theta d\theta)$$

$$\Rightarrow A_r(r, \theta) = r^2 \cos \theta - 3r, \quad A_\theta(r, \theta) = r^3 \sin \theta$$

[問2]

$$\begin{aligned} I_C &= \int_{C_1} d\vec{s} \cdot \vec{A} + \int_{C_2} d\vec{s} \cdot \vec{A} \\ &= \int_0^\pi R^3 \sin \theta d\theta + \int_{-R}^R (r^2 - 3r) dr = 2R^3 + \frac{2}{3}R^3 = \frac{8}{3}R^3 \end{aligned}$$

[問3]

$$\frac{\partial A_\theta}{\partial r} = 3r^2 \sin \theta, \quad \frac{\partial A_r}{\partial \theta} = -r^2 \sin \theta$$

$$I_D = \iint_D dr d\theta (3r^2 \sin \theta + r^2 \sin \theta) = 4 \int_0^\pi \sin \theta d\theta \int_0^R r^2 dr = \frac{8}{3}R^3$$

物理

第1問

[問1]

$$(1) \frac{\partial L}{\partial q} = mq, \quad \frac{\partial L}{\partial \dot{q}} = 0 \Rightarrow mq = 0$$

$$(2) \hat{p} = \frac{\partial L}{\partial \dot{q}} = mq, \quad H = p\dot{q} - L = mq^2 - \frac{m}{2}\dot{q}^2 = \frac{p^2}{2m}$$

$$(3) \hat{p} = -i\hbar \frac{\partial}{\partial q}, \quad \hat{H} \Psi(q) = -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial q^2} \Psi(q) \Rightarrow -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial q^2} \Psi(q) = E \Psi(q)$$

$$(4) \hat{H} \Psi(q) = -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial q^2} C_n \exp(i\frac{q}{R}) = -\frac{\hbar^2}{2m} \left(\frac{n}{R}\right)^2 \Psi(q)$$

$$\Rightarrow E = \hbar^2 \frac{n^2}{2mR^2}$$

$$\int_0^{2\pi} dq |\Psi_n|^2 = \int_0^{2\pi} dq |C_n|^2 = 1 \Rightarrow C_n = \sqrt{\frac{1}{2\pi R}}$$

[問2]

$$(1) \frac{\partial L}{\partial \dot{q}} = mq + \frac{i\hbar \theta}{2\pi R} \cdot \frac{\partial L}{\partial q} = 0 \Rightarrow mq = 0, \text{ not depend on } \theta.$$

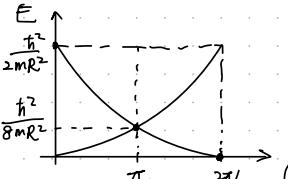
$$(2) \hat{p} = mq + \frac{i\hbar \theta}{2\pi R}, \quad H = mq^2 + \frac{i\hbar \theta}{2\pi R} \dot{q} - \frac{m}{2}\dot{q}^2 - \frac{i\hbar}{2\pi R} \theta \dot{q} = \frac{m}{2}\dot{q}^2 = \frac{1}{2m} (\hat{p} - \frac{i\hbar \theta}{2\pi R})^2$$

$$\hat{H} \Psi = \frac{1}{2m} \left(\hat{p}^2 - \frac{i\hbar \theta}{2\pi R} \hat{p} + \left(\frac{i\hbar \theta}{2\pi R} \right)^2 \right) \Psi = \left[-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial q^2} + \frac{i\hbar^2 \theta}{2m\pi R} \frac{\partial}{\partial q} + \frac{(i\hbar \theta)^2}{2m(2\pi R)^2} \right] \Psi = E \Psi$$

$$(3) ① -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial q^2} \Psi = E' \Psi \Rightarrow E' = \frac{\hbar^2 k^2}{2mR^2}$$

$$② \frac{i\hbar^2 \theta}{2m\pi R} \frac{\partial}{\partial q} \Psi = \frac{i\hbar^2 \theta}{2m\pi R} \left(\frac{i\hbar}{R} \right) \Psi = -\frac{n\hbar^2 \theta}{2m\pi R^2} \Psi, \Rightarrow E'' = -\frac{n\hbar^2 \theta}{2m\pi R^2}$$

$$\Rightarrow E_n = \frac{\hbar^2}{2mR^2} \left(n^2 - \frac{\theta}{\pi} n + \frac{\theta^2}{(\pi)^2} \right)$$



$$(4) E_0 = \frac{\hbar^2 \theta^2}{2m\pi^2(2\pi)^2}, \quad E_1 = \frac{\hbar^2}{2mR^2} \left(1 - \frac{\theta}{\pi} + \frac{\theta^2}{(\pi)^2} \right)$$

(5-1) no ideal

$$(5-2) \hat{H}_{\theta+2\pi} = \frac{1}{2m} \left[\hat{p} - \frac{i\hbar}{2\pi R} (\theta+2\pi) \right]^2 = \frac{1}{2m} \left[(\hat{p} - \frac{i\hbar}{R}) - \frac{\hbar^2}{2mR^2} \theta \right]^2$$

$$e^{-i\theta/R} \hat{H}_\theta e^{-i\theta/R} = e^{-i\theta/R} \frac{1}{2m} \left(\hat{p} - \frac{i\hbar}{2\pi R} \theta \right)^2 e^{-i\theta/R}$$

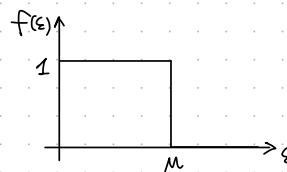
$$\begin{aligned} &= \frac{1}{2m} e^{-i\theta/R} \left(\hat{p}^2 - \hat{p} \frac{i}{\pi R} \theta + \left(\frac{i\hbar}{2\pi R} \theta \right)^2 \right) e^{i\theta/R} \\ &= \frac{1}{2m} \left[\left(\hat{p} - \frac{i\hbar}{R} \right)^2 - \frac{\hbar^2}{\pi R^2} \theta \left(\hat{p} - \frac{i\hbar}{R} \right) + \left(\frac{i\hbar}{2\pi R} \theta \right)^2 \right] \\ &= \frac{1}{2m} \left[\left(\hat{p} - \frac{i\hbar}{R} \right)^2 - \frac{\hbar^2}{2mR^2} \theta \right]^2 \quad \square \end{aligned}$$

$$\begin{aligned} (5-3) \hat{H}_{\theta+2\pi} \Psi &= e^{i\theta/R} \hat{H}_\theta e^{-i\theta/R} \Psi \\ &= \hat{H}_\theta e^{-i\theta/R} \Psi e^{i\theta/R} \quad ? \\ &= \hat{H}_\theta \Psi \end{aligned}$$

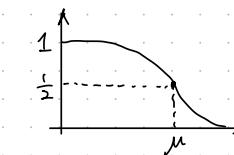
第2問

[問1]

① when T=0



② when T > 0



[問2]

$$D(\epsilon) \Delta \epsilon = \int_{\Delta \epsilon} \frac{d\epsilon}{h^3} = \frac{1}{h^3} \iiint dx dy dz \iiint dp_x dp_y dp_z = \frac{4\pi V P^2}{h^3} \Delta \epsilon$$

$$\Sigma = \frac{P^2}{2m}, \quad \Delta \epsilon = \frac{P}{m} \Delta p \Rightarrow P = \sqrt{2m\Sigma}, \quad \Delta p = \frac{m}{\sqrt{2m\Sigma}} \Delta \epsilon$$

$$\therefore D(\epsilon) \Delta \epsilon = \frac{4\pi V}{(2\pi h)^3} 2m \epsilon \frac{m}{\sqrt{2m\Sigma}} \Delta \epsilon \Rightarrow D(\epsilon) = \frac{V}{4\pi^2} \left(\frac{2m}{h^2} \right)^{\frac{3}{2}} \epsilon^{\frac{1}{2}}$$

$$[問3] \text{ when } T=0, \quad N = 2 \int_0^{\epsilon_F} d\epsilon \frac{V}{4\pi^2} \left(\frac{2m}{h^2} \right)^{\frac{3}{2}} \epsilon^{\frac{1}{2}}$$

$$\text{let } \frac{2}{4\pi^2} \left(\frac{2m}{h^2} \right)^{\frac{3}{2}} = A, \text{ then } N = AV \int_0^{\epsilon_F} \epsilon^{\frac{1}{2}} d\epsilon = AV \frac{2}{3} \epsilon_F^{\frac{3}{2}} \\ \Rightarrow \epsilon_F = \left(\frac{2n}{2A} \right)^{\frac{2}{3}} = \frac{h^2}{2m} (3\pi^2)^{\frac{2}{3}}$$

$$[問4] T_F = \frac{\epsilon_F}{k_B} = \frac{h^2}{2mk_B} (3\pi^2)^{\frac{2}{3}} \approx \frac{(1.0546 \times 10^{-34})^2}{2 \times 9.11 \times 10^{-31} \times 1.38 \times 10^{-23}} \times 3.1 \times 10^{20} \approx 1.37 \times 10^5 \text{ K}$$

$$[問5] D(\epsilon) = \frac{A}{2} V \epsilon^{\frac{1}{2}}, \quad \frac{dD(\epsilon)}{d\epsilon} = \frac{A}{4} V \epsilon^{-\frac{1}{2}}, \quad \int_0^M d\epsilon D(\epsilon) = \frac{A}{3} V M^{\frac{3}{2}}$$

$$N = 2 \int_0^\infty d\epsilon D(\epsilon) f(\epsilon) = 2 \left(\frac{A}{3} V M^{\frac{3}{2}} + \frac{\pi^2}{6} (k_B T)^2 \frac{A}{4} V M^{-\frac{1}{2}} \right) = \frac{2A}{3} V M^{\frac{3}{2}} \left(1 + \frac{\pi^2}{8} \left(\frac{k_B T}{h} \right)^2 \right)$$

$$\Rightarrow \mu = \left(\frac{3n}{2A} \right)^{\frac{2}{3}} \left(1 + \frac{\pi^2}{8} \left(\frac{kT}{\mu} \right)^2 \right)^{-\frac{2}{3}}$$

$$= \mu_0 \left(1 + \frac{\pi^2}{8} \left(\frac{kT}{\mu_0} \right)^2 \right)^{-\frac{2}{3}}$$

$$\approx \mu_0 \left(1 + \frac{\pi^2}{8} \left(\frac{kT}{\mu_0} \right)^2 \right)^{-\frac{2}{3}}$$

↓ Since μ is changing around μ_0 .

using Taylor expansion around $T=0$, we get $\mu \approx \mu_0 \left[1 - \frac{\pi^2}{6} \left(\frac{kT}{\mu_0} \right)^2 \right]$

[PQ6]

$$C = \frac{\pi^2}{4\varepsilon_F} k_B N \cdot 2T = \frac{\pi^2}{2} N k_B \left(\frac{k_B T}{\varepsilon_F} \right)$$

[PQ7]

Quantum statistics and classical statistics differences.

第3 PQ

$$[\text{PQ1}] \quad \begin{aligned} \mathcal{L} &= \frac{1}{2} m \dot{x}^2 - \frac{1}{2} k x^2 \quad \frac{\partial \mathcal{L}}{\partial x} = m \ddot{x} \quad \frac{\partial \mathcal{L}}{\partial \dot{x}} = -kx \quad \Rightarrow \quad m \ddot{x} = -kx \\ (1) \quad p &= \frac{\partial \mathcal{L}}{\partial \dot{x}} = m \dot{x} \quad H(x, p) = m \dot{x}^2 - \frac{1}{2} m x^2 + \frac{1}{2} k x^2 = \frac{p^2}{2m} + \frac{1}{2} k x^2 \\ (2) \quad \dot{x} &= \frac{p}{m}, \quad \dot{p} = -kx \end{aligned}$$

[PQ2]

(1) when $k=0$, $\dot{p}=0 \Rightarrow p(t)=p(t=0)=p(0)$

$$\dot{x} = \frac{p(0)}{m} \Rightarrow x(t) = \frac{p(0)}{m} t + x(0)$$

$$M = \begin{pmatrix} 1 & \frac{t}{m} \\ 0 & 1 \end{pmatrix}$$

(2) when $k=m\omega^2 > 0$

$$\ddot{x} = \frac{p}{m} = -\omega^2 x \Rightarrow x = A_1 \cos(\omega t) + A_2 \sin(\omega t)$$

$$\Rightarrow p = m \dot{x} = -m\omega A_1 \sin(\omega t) + m\omega A_2 \cos(\omega t)$$

take into initial condition, we have

$$\begin{cases} A_1 = x(0) \\ m\omega A_2 = p(0), \text{ therefore } \end{cases} \quad \begin{cases} x = x(0) \cos(\omega t) + \frac{p(0)}{m\omega} \sin(\omega t) \\ p = -m\omega x(0) \sin(\omega t) + p(0) \cos(\omega t) \end{cases}$$

$$\therefore M = \begin{pmatrix} \cos(\omega t) & \frac{\sin(\omega t)}{m\omega} \\ -m\omega \sin(\omega t) & \cos(\omega t) \end{pmatrix}$$

(3) when $k = -m\omega^2 < 0$

$$\ddot{x} = \omega^2 x \Rightarrow x = B_1 \cosh(\omega t) + B_2 \sinh(\omega t)$$

$$\Rightarrow p = -m\omega B_1 \sinh(\omega t) + m\omega B_2 \cosh(\omega t)$$

take into initial condition, we have

$$\begin{cases} B_1 = x(0) \\ m\omega B_2 = p(0) \end{cases}, \text{ therefore } \begin{cases} x = x(0) \cosh(\omega t) + \frac{p(0)}{m\omega} \sinh(\omega t) \\ p = -m\omega x(0) \sinh(\omega t) + p(0) \cosh(\omega t) \end{cases}$$

$$\therefore M = \begin{pmatrix} \cosh(\omega t) & \frac{\sinh(\omega t)}{m\omega} \\ -m\omega \sinh(\omega t) & \cosh(\omega t) \end{pmatrix}$$

[PQ3]

$$(1) \quad W(t) = \dot{x}_2 p_1 + \dot{x}_2 \dot{p}_1 - \dot{p}_2 x_1 - p_2 \dot{x}_1$$

$$= \frac{p_2}{m} p_1 + x_2 (-k x_1) - (-k x_2) x_1 - p_2 \frac{p_1}{m}$$

$$= 0$$

$$(2) \quad \begin{cases} x_n(t) = a x_n(0) + b p_n(0) \\ p_n(t) = c x_n(0) + d p_n(0) \end{cases}, \quad n=1, 2.$$

$$\begin{aligned} W(t) &= (a x_2(0) + b p_2(0)) (c x_1(0) + d p_1(0)) - (c x_2(0) + d p_2(0)) (a x_1(0) + b p_1(0)) \\ &= ac x_1(0) x_2(0) + ad x_2(0) p_1(0) + bc p_2(0) x_1(0) + bd p_1(0) p_2(0) \\ &\quad - ac x_2(0) x_2(0) - bc p_1(0) x_2(0) - da x_1(0) p_2(0) - db p_1(0) p_2(0) \\ &= (ad - bc) x_2(0) p_1(0) + (bc - da) x_1(0) p_2(0) \\ &= \det T \cdot W(t=0) \end{aligned}$$

$$\Rightarrow \det T = 1$$

第4 PQ

$$(1) \quad \begin{array}{c} \text{Gaussian Surface} \\ \text{cylinder} \end{array}, \quad \iint_S \vec{E} \cdot d\vec{S} = 2\pi r l E \quad \Rightarrow E = \frac{\alpha \lambda}{\epsilon_0 r}$$

$$\frac{Q}{\epsilon_0} = \frac{\lambda l \cdot 2\pi r}{\epsilon_0}$$

$$(2) \quad E_A(r) = \frac{\alpha \lambda}{\epsilon_0 r}, \quad E_B(r) = \frac{-\alpha \lambda}{\epsilon_0 (d-r)}$$

$$V_{AB} = \int_{d-a}^{d+a} (E_A + E_B) dr = \frac{\alpha \lambda}{\epsilon_0} l \ln \frac{d+a}{d-a}$$

$$(3) \quad C = \frac{Q'}{V_{AB}} = \frac{2\lambda \cdot 2\pi A}{\alpha \lambda l \ln \frac{d+a}{d-a}} = \frac{4\pi \epsilon_0}{l \ln \frac{d+a}{d-a}}$$

$$(4) \quad \text{For A: } \oint_C \vec{H} \cdot d\vec{s} = 2\pi r H_A \quad \Rightarrow H_A = \frac{I}{2\pi r}, \text{ same as } H_B = \frac{-I}{2\pi(d-r)}$$

Therefore $H = \frac{I}{2\pi} \left(\frac{1}{r} - \frac{1}{d-r} \right)$

$$(5) B = \mu_0 H$$

$$\Phi = \oint \vec{B} \cdot d\vec{S} = l \int_{d-a}^{d+a} \mu_0 \frac{I}{2\pi} \left(\frac{1}{r} - \frac{1}{d-r} \right) dr = \frac{\mu_0 I}{2\pi} \ln \frac{d+a}{d-a}$$

$$(6) z \perp I = \frac{\Phi}{l} \Rightarrow L = \frac{\Phi}{z \cdot l I} = \frac{\mu_0}{2\pi} \ln \frac{d+a}{d-a}$$

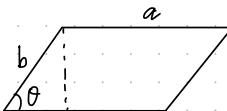
↑ ← unit length

數學

第1問

[PQ1]

$$S = b \sin \theta \cdot a = |\vec{a} \times \vec{b}|$$



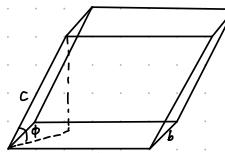
[PQ2]

$$V = S \cdot c \cdot \cos \phi = |\vec{a} \times \vec{b} \cdot \vec{c}|$$

[PQ3]

$$M = \vec{a} \times \vec{b} \cdot \vec{c} = \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ ax & ay & az \\ bx & by & bz \end{vmatrix} \cdot (cx\hat{x} + cy\hat{y} + cz\hat{z})$$

$$= (aybz - azby)Cx + (azbx - axbz)Cy + (axby - aybx)Cz$$



第2問

[PQ1]

$$\text{Take (2) into (1), we have } \nabla^2 G(\vec{x}) = \int \frac{\partial^2 \vec{k}}{\partial \vec{x}^2} \frac{1}{|\vec{k}|^2} \nabla^2 e^{i\vec{k} \cdot \vec{x}} = \int \frac{d^3 \vec{k}}{G^2} \frac{(i\vec{k})^2}{|\vec{k}|^2} e^{i\vec{k} \cdot \vec{x}}$$

$$= - \int \frac{d^3 \vec{k}}{G^2} e^{i\vec{k} \cdot \vec{x}}$$

$$\text{For } \delta^{(1)}(x) = \frac{dk}{2\pi} e^{-ikx}, \text{ therefore } \nabla^2 G(x) = -\delta^{(3)}(\vec{x})$$

[PQ2]

Green's theorem: $\oint_{C_1 + C_2 + C_3 + C_4} f(z) dz = 0$

$$\int_0^\infty dz \frac{\sin z}{z} = \int_0^\infty dz \frac{e^{iz} - e^{-iz}}{2iz} = \int_{-\infty}^\infty dz \frac{e^{iz}}{2iz}, \text{ let } f(z) = \frac{e^{iz}}{z}, z = Re^{i\theta}$$

$$\lim_{R \rightarrow \infty} \int_{C_{111}} dz \frac{e^{iz}}{z} = \lim_{R \rightarrow \infty} \int_0^\pi \frac{e^{iR(\cos \theta + i\sin \theta)}}{Re^{i\theta}} iRe^{i\theta} d\theta \rightarrow 0$$

$$\lim_{\epsilon \rightarrow 0} \int_{C_\epsilon} dz \frac{e^{iz}}{z} = \lim_{\epsilon \rightarrow 0} \int_\pi^0 \frac{e^{i(\epsilon e^{i\theta})}}{\epsilon e^{i\theta}} i\epsilon e^{i\theta} d\theta = i \int_\pi^0 d\theta = -i\pi$$

$$\lim_{\substack{R \rightarrow \infty \\ \epsilon \rightarrow 0}} \int_{C_1 + C_2 + C_3 + C_4} dz \frac{e^{iz}}{z} = \lim_{\substack{R \rightarrow \infty \\ \epsilon \rightarrow 0}} \int_{-\epsilon}^\epsilon \frac{e^{iz}}{z} dz + \int_\epsilon^R \frac{e^{iz}}{z} dz = \int_{-\infty}^{+\infty} \frac{e^{iz}}{z} dz$$

$$\therefore -i\pi + \int_{-\infty}^{+\infty} \frac{e^{iz}}{z} dz = 0 \Rightarrow \int_0^\infty dz \frac{\sin z}{z} = \frac{i}{2} \int_{-\infty}^{+\infty} dz \frac{e^{iz}}{z} = \frac{\pi}{2}$$

[PQ3, PQ4]. No idea

第3問

[PQ1]

$$\frac{d^2 f(x)}{dx^2} = \frac{2x \frac{df}{dx} - 2f}{1-x^2}$$

$$(a) \frac{dW(x)}{dx} = \frac{df_1}{dx} \frac{df_2}{dx} + f_1 \frac{d^2 f_2}{dx^2} - \frac{df_2}{dx} \frac{df_1}{dx} - f_2 \frac{d^2 f_1}{dx^2}$$

$$= f_1 \frac{2}{1-x^2} \left(x \frac{df_2}{dx} - f_2 \right) - f_2 \frac{2}{1-x^2} \left(x \frac{df_1}{dx} - f_1 \right)$$

$$= \frac{2x}{1-x^2} (f_1 \frac{df_2}{dx} - f_2 \frac{df_1}{dx})$$

$$= \frac{2x}{1-x^2} W(x)$$

$$(b) \int \frac{1}{W(x)} dW(x) = \int \frac{2x}{1-x^2} dx \Rightarrow \ln |W(x)| = \ln |1-x^2|^{-1} + C$$

$$\Rightarrow W(x) = C \frac{1}{1-x^2}$$

$$[PQ2] 2x \frac{df}{dx} - 2f = 0 \Rightarrow \int \frac{1}{f} df = \int \frac{1}{x} dx \Rightarrow f_1 = C_1 x$$

take condition of $f_1(1) = 1$ into, we got $f_1 = x$

[PQ3]

$$W(x) = x \frac{df_2}{dx} - f_2 = \frac{C}{1-x^2} \Rightarrow \frac{df_2}{dx} - \frac{1}{x} f_2 - \frac{C}{x(1-x^2)} = 0$$

$$[PQ4] p(x) = -\frac{1}{x}, q(x) = \frac{-C}{x(1-x^2)}$$

$$\alpha(x) = \int p(x) dx = \frac{1}{x}$$

$$f_2(x) = \frac{1}{\alpha(x)} \left[\int_x^\infty \alpha(x) q(x) dx + C_1 \right] = x \left[-C_1 \int_x^\infty \left(\frac{1}{x^2} - \frac{1}{1-x^2} \right) dx + C_1 \right]$$

$$= \pi \left\{ -C_1 \left[\frac{1}{x} - \frac{1}{2} \ln \frac{1+x}{1-x} \right] + C_1 \right\}$$

$$= \frac{C_1 x}{2} \ln \frac{1+x}{1-x} + C_1 x + C$$

$$\therefore f(x) = C_2 f_1(x) + C_3 f_2(x) = C_2 x + \frac{CC_3 x}{2} \ln \frac{1+x}{1-x} + C_3 C_1 x + C C_3$$

$$= A_1 \pi \ln \frac{1+x}{1-x} + A_2 x + A_3$$

第4問

[P4] 1]

$$U^*U = I \Rightarrow \det(U^*U) = \det I \Rightarrow (\det U)^2 = 1 \Rightarrow \det U = 1$$

[P4] 2]

$$UX = \lambda X \Rightarrow (U - \lambda I)X = 0 \quad \det(U - \lambda I) = 0 \Rightarrow \det U - \lambda \det I = 0 \Rightarrow \lambda = 1$$

[P4] 3]

$$\begin{cases} \det U = ad - bc = e^{i\varphi} \\ a+d = 1 \end{cases} \Rightarrow \begin{cases} c = \frac{a-a^2-e^{i\varphi}}{b} \\ d = 1-a \end{cases}$$

[P4] 4]

$$\begin{vmatrix} \frac{1}{\sqrt{2}} - \lambda & \frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} - \lambda \end{vmatrix} = \left(\frac{1}{\sqrt{2}} - \lambda\right)\left(\frac{1}{\sqrt{2}} - \lambda\right) + \frac{1}{2} = 0 \Rightarrow \lambda^2 + \left(-\frac{1}{\sqrt{2}} - \frac{1}{\sqrt{2}}\right)\lambda + \frac{1}{2} = 0$$

$$\Rightarrow \lambda = \frac{(1+i) \pm \sqrt{-6i}}{2\sqrt{2}}, \quad (1-i)^2 = -2i \Rightarrow \sqrt{-2i} = 1-i \Rightarrow \sqrt{-6i} = (1-i)\sqrt{3}$$

$$= \frac{(1+i) \pm (1-i)\sqrt{3}}{2\sqrt{2}}$$

$$\therefore \lambda_1 = \frac{1+\sqrt{3}}{2\sqrt{2}} + i \frac{1-\sqrt{3}}{2\sqrt{2}}, \quad \lambda_2 = \frac{1-\sqrt{3}}{2\sqrt{2}} + i \frac{1+\sqrt{3}}{2\sqrt{2}}$$

SOKENDAI KEK 入試 2021 解答

物理題

第1問

[P1]

$$\hbar = \sqrt{[q][\epsilon][L]^3[m]}$$

[P2]

$$\begin{aligned} i\hbar \hat{\pi} \Psi(x,t) &= \hat{H} \Psi(x,t) \Rightarrow \frac{d}{dt} \Psi(x,t) = \frac{\hat{H}}{i\hbar} \Psi(x,t) \\ -i\hbar \frac{d}{dt} \Psi^*(x,t) &= \hat{H} \Psi^*(x,t) \Rightarrow \frac{d}{dt} \Psi^*(x,t) = \frac{\hat{H}}{i\hbar} \Psi^*(x,t) \end{aligned}$$

$$\begin{aligned} \frac{d\langle x \rangle}{dt} &= \frac{d}{dt} \langle \hat{x} | \Psi | \Psi \rangle = \langle \frac{d\hat{x}}{dt} | \Psi | \frac{d\Psi}{dt} \rangle + \langle \Psi | \frac{d\hat{x}}{dt} | \Psi \rangle + \langle \Psi | \frac{d\hat{x}}{dt} | \frac{d\Psi}{dt} \rangle \\ &= \frac{1}{i\hbar} \langle \Psi | \hat{H} \hat{x} | \Psi \rangle + \frac{1}{i\hbar} \langle \Psi | \hat{x} \hat{H} | \Psi \rangle \\ &= \frac{1}{i\hbar} \langle [\hat{x}, \hat{H}] \rangle \\ &= \frac{1}{i\hbar} \left\langle \left[\hat{x}, \frac{\hat{p}^2}{2m} + V(x) \right] \right\rangle \quad \Rightarrow \quad \langle a \rangle = \frac{d}{dt} \left(\frac{d\langle x \rangle}{dt} \right) = \frac{1}{m} \frac{d\langle \hat{p} \rangle}{dt} \\ &= \frac{1}{i\hbar 2m} \langle [\hat{x}, \hat{p}^2] \rangle \\ &= \frac{1}{i\hbar 2m} \cdot 2i\hbar \langle \hat{p} \rangle \\ &= \frac{\langle \hat{p} \rangle}{m} \end{aligned}$$

[P3]

$$\frac{d}{dt} (e^{iS(x)/\hbar} e^{-iEt/\hbar}) = \frac{i}{\hbar} \left(-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} - g\epsilon x \right) (e^{iS(x)/\hbar} e^{-iEt/\hbar})$$

$$l.h.s = e^{iS(x)/\hbar} \left(-\frac{iE}{\hbar} \right) e^{-iEt/\hbar} = \frac{-iE}{\hbar} \Psi_E(x,t)$$

$$r.h.s = \frac{i}{\hbar} \left[-\frac{\hbar^2}{2m} \left(\frac{i}{\hbar} \frac{\partial^2 S}{\partial x^2} + \left(\frac{i}{\hbar} \frac{\partial S}{\partial x} \right)^2 \right) - g\epsilon x \right] \Psi_E(x,t)$$

$$\Rightarrow E = -\frac{\hbar^2}{2m} \left[\frac{i}{\hbar} \frac{\partial^2 S}{\partial x^2} + \left(\frac{i}{\hbar} \frac{\partial S}{\partial x} \right)^2 \right] - g\epsilon x$$

$$\Rightarrow \frac{d^2 S(x)}{dx^2} + \frac{i}{\hbar} \left(\frac{dS(x)}{dx} \right)^2 - \frac{2mi}{\hbar} (E + g\epsilon x) = 0$$

[P4]

$$\frac{dS(x)}{dx} = \frac{dS_0(x)}{dx} + \hbar \frac{dS_1(x)}{dx}, \quad \frac{d^2 S(x)}{dx^2} = \frac{d^2 S_0(x)}{dx^2} + \hbar \frac{d^2 S_1(x)}{dx^2}$$

$$\frac{d^2 S_0}{dx^2} + \hbar \frac{d^2 S_1}{dx^2} + \frac{i}{\hbar} \left[\left(\frac{dS_0}{dx} \right)^2 + 2\hbar \frac{dS_0}{dx} \frac{dS_1}{dx} + \hbar \left(\frac{dS_1}{dx} \right)^2 \right] - \frac{2mi}{\hbar} (E + g\epsilon x) = 0$$

Multiply \hbar in both sides, and ignore the \hbar^2 terms, we get

$$\hbar \frac{d^2S_0}{dx^2} + i \left(\frac{dS_0}{dx} \right)^2 + 2i\hbar \frac{dS_0}{dx} \frac{dS_1}{dx} - 2m(E + \beta \epsilon x) = 0$$

$$\Rightarrow \hbar \left(\frac{d^2S_0}{dx^2} + 2i \frac{dS_0}{dx} \frac{dS_1}{dx} \right) + i \left[\left(\frac{dS_0}{dx} \right)^2 - 2m(E + \beta \epsilon x) \right] = 0$$

$$\Rightarrow \begin{cases} E = \frac{1}{2m} \left(\frac{dS_0(x)}{dx} \right)^2 - \beta \epsilon x \\ 0 = \frac{dS_0(x)}{dx} \frac{dS_1(x)}{dx} - \frac{i}{\hbar} \frac{d^2S_0(x)}{dx^2} \end{cases}$$

[P9] 5]

$$\frac{dS_0}{dx} = \pm \sqrt{2m(E + \beta \epsilon x)} \Rightarrow S_0(x) = \frac{\sqrt{2m}}{3\beta \epsilon} \left[(E + \beta \epsilon L)^{\frac{3}{2}} - E^{\frac{3}{2}} \right]$$

$$\frac{d^2S_0}{dx^2} = \pm \sqrt{2m} \frac{9\epsilon}{2} (E + \beta \epsilon x)^{-\frac{1}{2}}$$

$$0 = \sqrt{2m(E + \beta \epsilon x)} \frac{dS_0(x)}{dx} - \frac{i}{\hbar} \sqrt{2m} \frac{9\epsilon}{2} (E + \beta \epsilon x)^{-\frac{1}{2}}$$

$$\Rightarrow \frac{dS_1(x)}{dx} = \frac{i\beta \epsilon}{4} (E + \beta \epsilon x)^{-1} \Rightarrow S_1(x) = \frac{i}{4} \ln(1 + \frac{\beta \epsilon L}{E} x) \Big|_0^L = \frac{i}{\pi} \ln(1 + \frac{\beta \epsilon L}{E})$$

$$\therefore \psi(x) = e^{i(S_0(x) + S_1(x))/\hbar} = e^{i \frac{\sqrt{2m}}{3\beta \epsilon \hbar} \left[(E + \beta \epsilon L)^{\frac{3}{2}} - E^{\frac{3}{2}} \right] (1 + \frac{\beta \epsilon L}{E})^{-\frac{1}{2}}}$$

[P9] 6] no idea

第 27 題

[P9] 1]

$$(1) Z = \sum_{\sigma_i=\pm 1} e^{-\beta E_i / (k_B T)} = \sum_{\sigma_i=\pm 1} e^{\frac{N}{2} M \sigma_i H / (k_B T)}$$

$$= \prod_{i=1}^N \left(e^{\mu H / (k_B T)} + e^{-\mu H / (k_B T)} \right) = \left[2 \cosh \left(\frac{\mu H}{k_B T} \right) \right]^N$$

$$F = -k_B T \ln Z = -k_B T N \left[\ln 2 + \ln \left(\cosh \left(\frac{\mu H}{k_B T} \right) \right) \right]$$

$$\begin{aligned} S &= -\frac{\partial F}{\partial T} = k_B N \left[\ln 2 + \ln \left(\cosh \left(\frac{\mu H}{k_B T} \right) \right) \right] + k_B T N \frac{1}{\cosh \frac{\mu H}{k_B T}} \sinh \left(\frac{\mu H}{k_B T} \right) \left(-\frac{\mu H}{k_B} \right) T^{-2} \\ &= k_B N \left\{ \ln 2 + \ln \left(\cosh \left(\frac{\mu H}{k_B T} \right) \right) \right\} - \tanh \left(\frac{\mu H}{k_B T} \right) \frac{\mu H}{k_B T} \end{aligned}$$

$$\bar{E} = k_B T^2 \frac{\partial}{\partial T} \left(\frac{F}{k_B T} \right) = -k_B T^2 \frac{\partial}{\partial T} \left(N \cdot \ln \cosh \left(\frac{\mu H}{k_B T} \right) \right)$$

$$= k_B T^2 \cdot N \cdot \tanh \left(\frac{\mu H}{k_B T} \right) \left(\frac{\mu H}{k_B T^2} \right) = N H N \tanh \frac{\mu H}{k_B T}$$

$$\begin{aligned} (2) M &= \mu \sum_{\sigma_i=\pm 1} \sigma_i = -\frac{1}{V} \left(\frac{\partial F}{\partial H} \right)_T = k_B T n \cdot \tanh \left(\frac{\mu H}{k_B T} \right) \cdot \frac{M}{k_B T} \\ &= n \mu \tanh (\beta \mu H) \end{aligned}$$

(3) ?

$$M_0 = n \mu \tanh (\beta \mu a M_0) \approx n \mu \left[\beta \mu a M_0 - (\beta \mu a M_0)^3 / 3 \right]$$

$$= M_0 \left[n \mu^2 a \beta - n \mu^4 a^3 \beta^3 M_0^2 / 3 \right]$$

$$\Rightarrow M_0 \sim \sqrt{\frac{n \mu^2 a \beta - 1}{n \mu^4 a^3 \beta^3}} = \sqrt{\frac{(T_c - T)^{\frac{1}{2}}}{\alpha T_c^2 / \hbar k_B T^3}}$$

$$d\varepsilon = \frac{P}{m} d\phi$$

$$\varepsilon = \frac{P^2}{2m}$$

$$P = \sqrt{2m\varepsilon}$$

[P9] 2]

$$(1) D(\varepsilon) d\varepsilon = \int \frac{dw}{h^3} = \frac{V}{h^3} \int dp_x dp_y dp_z = \frac{4\pi V}{h^3} P^2 dp$$

$$= \frac{c \pi V}{(2\pi\hbar)^3} (2m\varepsilon) \cdot \frac{m}{\sqrt{2m\varepsilon}} d\varepsilon$$

$$\Rightarrow D(\varepsilon) = \frac{m^{\frac{3}{2}} \varepsilon V}{\sqrt{2\pi^2 h^3}} \varepsilon^{\frac{1}{2}}$$

$$\sqrt{(\varepsilon + \mu \sigma H) 2m} = \phi$$

$$(2) D_+(\varepsilon) d\varepsilon = \frac{4\pi V}{(2\pi\hbar)^3} P^2 dp$$

$$d\phi = (\varepsilon + \mu \sigma H)^{-\frac{1}{2}} (2m)^{\frac{1}{2}} \frac{1}{2} d\varepsilon$$

$$= \frac{V}{2\pi^2 h^3} (\varepsilon + \mu H) \cdot 2m \cdot 2^{\frac{1}{2}} \cdot m^{\frac{1}{2}} (\varepsilon + \mu H)^{-\frac{1}{2}} d\varepsilon$$

$$= \frac{m^{\frac{3}{2}} V}{\sqrt{2\pi^2 h^3}} (\varepsilon + \mu H)^{\frac{1}{2}} d\varepsilon$$

$$D_+(\varepsilon) = \frac{m^{\frac{3}{2}} V}{\sqrt{2\pi^2 h^3}} (\varepsilon + \mu H)^{\frac{1}{2}}, D_-(\varepsilon) = \frac{m^{\frac{3}{2}} V}{\sqrt{2\pi^2 h^3}} (\varepsilon - \mu H)^{\frac{1}{2}}$$

$$(3) N_+ = \int_0^{\varepsilon_F} D_+(\varepsilon) d\varepsilon = \frac{m^{\frac{3}{2}} V}{\sqrt{2\pi^2 h^3}} \int_0^{\varepsilon_F} (\varepsilon + \mu H)^{\frac{1}{2}} d\varepsilon$$

$$= \frac{2 m^{\frac{3}{2}} V}{3 \sqrt{2\pi^2 h^3}} \left[(\varepsilon_F + \mu H)^{\frac{3}{2}} - (\mu H)^{\frac{3}{2}} \right]$$

$$N_- = \frac{m^{\frac{3}{2}} N}{\sqrt{2\pi^2 h^3}} \int_{\mu H}^{\varepsilon_F} (\varepsilon - \mu H)^{\frac{1}{2}} d\varepsilon = \frac{2 m^{\frac{3}{2}} V}{3 \sqrt{2\pi^2 h^3}} (\varepsilon_F - \mu H)^{\frac{3}{2}}$$

$$N = N_+ \sigma_+ + N_- \sigma_- = N_+ - N_-$$

$$= \frac{2m^{\frac{3}{2}}V}{3\sqrt{2}\pi^2 h^3} \left[(\varepsilon_F + \mu H)^{\frac{3}{2}} - (\varepsilon_F - \mu H)^{\frac{3}{2}} - (\mu H)^{\frac{3}{2}} \right]$$

Since $(\varepsilon_F + \mu H)^{\frac{3}{2}} \approx \varepsilon_F^{\frac{3}{2}} + \frac{3}{2}\mu (\varepsilon_F + \mu H)^{\frac{1}{2}}$ $\Big|_{H=0}$ $H + O(H^2)$

 $(\varepsilon_F - \mu H)^{\frac{3}{2}} \approx \varepsilon_F^{\frac{3}{2}} - \frac{3}{2}\mu (\varepsilon_F - \mu H)^{\frac{1}{2}} \Big|_{H=0} H + O(H^2)$.

$$\Rightarrow M \approx \frac{2m^{\frac{3}{2}}V}{3\sqrt{2}\pi^2 h^3} M \varepsilon_F H + o(H) \sim H \quad \square$$

第2問

[Pb] 1]

$$(1) \dot{x} = a(2\dot{\theta} + 2\dot{\theta} \cos 2\theta) \quad \dot{z} = 2a\dot{\theta} \sin 2\theta$$

$$K = \frac{1}{2}m(\dot{x}^2 + \dot{z}^2) = \frac{m}{2}[4a^2\dot{\theta}^2(1+\cos 2\theta)^2 + 4a^2\dot{\theta}^2(\sin 2\theta)^2]$$
 $= 2a^2\dot{\theta}^2 m [4\cos^4\theta + 4\cos^2\theta(1-\cos^2\theta)]$
 $= 8ma^2\dot{\theta}^2 \cos^2\theta$

$$(2) U = mgz = 2mg a \sin^2\theta$$

$$(3) \frac{dx}{d\theta} = a(2 + 2\cos 2\theta) \quad \frac{dz}{d\theta} = 2a\sin 2\theta$$

$$s = \int_0^\theta \sqrt{(1+\cos 2\theta)^2 + \sin^2 2\theta} d\theta = 2a \int_0^\theta \sqrt{2 + 2\cos 2\theta} d\theta \quad (\theta < 0) \\ = -2a \sin \theta$$

$$(4) \dot{s} = \frac{ds}{d\theta} = -4a \cos \theta$$

$$\mathcal{L} = K - U = 8ma^2\dot{\theta}^2 \cos^2\theta - 2mg a \sin^2\theta$$
 $= \frac{1}{2}m\dot{s}^2 - \frac{mg}{8a} s^2$

$$(5) \frac{\partial \mathcal{L}}{\partial \dot{s}} = m\ddot{s} \quad \frac{\partial \mathcal{L}}{\partial s} = -\frac{mg}{4a}s \quad \Rightarrow \quad \ddot{s} = -\frac{gs}{4a}$$

$$(6) S(t) = A_1 \cos \sqrt{\frac{g}{4a}}t + A_2 \sin \sqrt{\frac{g}{4a}}t \quad T = \frac{2\pi}{\sqrt{\frac{g}{4a}}} = 4\pi \sqrt{\frac{a}{g}}$$

$$(7) \text{when } \theta = \frac{\pi}{2}, x = 2a\theta = a\pi = 1.54/2$$

$$t = \frac{T}{2} = 2\pi \sqrt{\frac{a}{g}} = 4\pi/\sqrt{10}$$

Pb] 2

(1)

$$x = \sin \phi a(2\theta + \sin 2\theta)$$

$$y = \cos \phi a(2\theta + \sin 2\theta)$$

$$\dot{x} = a\dot{\phi} \cos \phi (2\theta + \sin 2\theta) + 2a\sin \phi \dot{\theta} (1 + \cos 2\theta)$$

$$\dot{y} = -a\dot{\phi} \sin \phi (2\theta + \sin 2\theta) + 2a\cos \phi \dot{\theta} (1 + \cos 2\theta)$$

$$\dot{z} = 2a\dot{\theta} \sin 2\theta$$

$$K = \frac{1}{2}m(\dot{x}^2 + \dot{y}^2 + \dot{z}^2)$$

$$= \frac{1}{2}m \left\{ a^2\dot{\phi}^2(2\theta + \sin 2\theta)^2 + 4a^2\dot{\theta}^2 \sin^2 \phi (1 + \cos 2\theta)^2 + 4a^2\dot{\theta}^2 \cos^2 \phi (1 + \cos 2\theta)^2 \right. \\ \left. + 4a^2\dot{\theta}^2 \sin^2 \theta \right\}$$

$$= \frac{1}{2}m a^2 \left\{ \dot{\phi}^2 (4\theta^2 + 4\theta \sin 2\theta + \sin^2 2\theta) + 8\dot{\theta}^2 + 8\dot{\theta}^2 \cos 2\theta \right\}$$

$$\text{we have } s = 4a \sin \theta \quad \dot{s} = 4a\dot{\theta} \cos \theta \quad \frac{s}{a} = 4 \sin \theta \sim \alpha$$

$$K = \frac{1}{2}ma^2 \cdot 16\dot{\theta}^2 \cos^2 \theta + \frac{1}{2}ma^2\dot{\phi}^2(4\theta^2 + 8\theta^2 + 4\theta^2)$$

$$= \frac{1}{2}m\dot{s}^2 + \frac{1}{2}m\dot{\phi}^2 s^2$$

$$\mathcal{L} = K - U = \frac{1}{2}m\dot{s}^2 - \frac{mg}{8a}s^2 + \frac{1}{2}m\dot{\phi}^2 s^2$$

$$(2) \frac{\partial \mathcal{L}}{\partial \dot{\phi}} = 0 \quad \frac{\partial \mathcal{L}}{\partial \phi} = l \quad \Rightarrow \quad \frac{dl}{dt} = 0, \quad : l \text{ is conserved}$$

$$(3) \quad l = \frac{\partial L}{\partial \dot{\phi}} = m s^2 \dot{\phi} \Rightarrow \dot{\phi} = \frac{l}{ms^2}$$

$$E = \frac{1}{2} m \dot{s}^2 + \frac{mg}{8\alpha} s^2 + \frac{l^2}{2ms^2}$$

$$\frac{\partial L}{\partial s} = -\frac{mg}{2\alpha} s + m\dot{\phi}^2 s, \quad \frac{\partial L}{\partial \dot{s}} = m\dot{s} \Rightarrow \ddot{s} = -\frac{g s}{2\alpha} + \frac{l^2}{m s^3}$$

$$\frac{dE}{dt} = m\dot{s}\ddot{s} + \frac{mgs}{4\alpha} \dot{s} + \frac{l^2}{2m}(-2)s^{-3} \dot{s}$$

$$= \dot{s} \left\{ -\frac{gms}{4\alpha} + \frac{l^2}{ms^3} + \frac{mgs}{4\alpha} - \frac{l^2}{ms^3} \right\} = 0$$

$\Rightarrow E$ is conserved.

(4)(5) no idea

第4題

[P6]1]

$$(1) \quad Q = -\epsilon_0 S_x$$

$$(2) \quad E \cdot \vec{S} = -2\epsilon_0 S_x / \epsilon_0 \Rightarrow E = -\frac{\epsilon_0 x}{\epsilon_0}$$

$$(2) \quad M_e \frac{d\vec{V}}{dt} = \epsilon \vec{E}$$

$$(3) \quad \frac{dx^2}{dt^2} + \frac{\epsilon^2 \mu_0}{\epsilon_0 m_e} x = 0, \text{ compare to } \frac{dx^2}{dt^2} + \omega_p^2 x = 0$$

$$\Rightarrow \omega_p = \sqrt{\frac{\mu_0}{\epsilon_0 m_e}}$$

[P6]2]

$$(1) \quad \nabla \times (\nabla \times \vec{E}) = \nabla(\nabla \cdot \vec{E}) - \nabla^2 \vec{E} = \nabla \left(\frac{\rho}{\epsilon_0} \right) - \nabla^2 \vec{E}$$

$$\nabla \times \left(-\frac{\partial \vec{B}}{\partial t} \right) = -\frac{\partial}{\partial t} \nabla \times \vec{B} = -\frac{\partial}{\partial t} [\mu_0 (\vec{j} + \epsilon_0 \frac{\partial}{\partial t} \vec{E})]$$

$$\Rightarrow \nabla^2 \vec{E} - \epsilon_0 \mu_0 \frac{\partial^2 \vec{E}}{\partial t^2} - \nabla \left(\frac{\rho}{\epsilon_0} \right) - \frac{\partial}{\partial t} [\mu_0 \vec{j}] = 0$$

$$\text{Since } \rho = 0, \quad \frac{d}{dt} \vec{j} = -\epsilon_0 \mu_0 \frac{\partial \vec{E}}{\partial t} = -\epsilon_0 \omega_p^2 \vec{E}$$

$$\therefore \nabla^2 \vec{E} - \epsilon_0 \mu_0 (\omega_p^2 \vec{E} + \frac{\partial^2 \vec{E}}{\partial t^2}) = 0$$

$$(2) \quad \vec{E} = \vec{E}_0 \cos(kz - \omega t + \phi) \quad \vec{r} = (x, y, z) \\ = E_0 \cos(kz - \omega t + \phi) \quad \vec{k} = (0, 0, k)$$

Take it into the equation, then

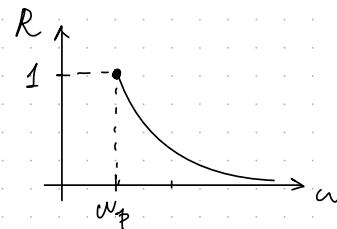
$$-E_0 k^2 \cos(kz - \omega t + \phi) - \epsilon_0 \mu_0 (\omega_p^2 E_0 \cos(kz - \omega t + \phi)) \\ + \epsilon_0 \mu_0 \omega^2 E_0 \cos(kz - \omega t + \phi) = 0$$

$$\Rightarrow -k^2 - \epsilon_0 \mu_0 (\omega_p^2 - \omega^2) = 0 \Rightarrow \omega = \sqrt{\omega_p^2 + \frac{k^2}{\epsilon_0 \mu_0}}$$

$$(3) \quad V_p = \frac{\omega}{k} = \left[\epsilon_0 \mu_0 \left(1 - \left(\frac{\omega_p}{\omega} \right)^2 \right) \right]^{-\frac{1}{2}} \quad k = \sqrt{\epsilon_0 \mu_0 (\omega^2 - \omega_p^2)}$$

$$m = \frac{c}{V_p} = \sqrt{1 - \left(\frac{\omega_p}{\omega} \right)^2}$$

$$(4) \quad R = \left| \frac{1 - \sqrt{1 - \left(\frac{\omega_p}{\omega} \right)^2}}{1 + \sqrt{1 - \left(\frac{\omega_p}{\omega} \right)^2}} \right|$$



數學

第一題

$$[\text{P6}1] \quad \int_{-\infty}^{+\infty} dx e^{-ax^2} \int_{-\infty}^{+\infty} dy e^{-ay^2} = \int_0^\infty d\theta \int_0^\infty dr r e^{-ar^2} = \pi r \int_0^\infty \frac{1}{2} dr r^2 e^{-ar^2} = \frac{\pi}{a}$$

$$\Rightarrow \int_{-\infty}^{+\infty} dx e^{-ax^2} = \sqrt{\frac{\pi}{a}}$$

$$[\text{P6}2] \quad \text{let } x-b=t, I_1 = \int_{-\infty}^{+\infty} dt e^{-at^2} (t+b)^2 = \int_{-\infty}^{+\infty} dt e^{-at^2} t^2 + 2bt + b^2 = \int_{-\infty}^{+\infty} e^{-at^2} t^2 dt + b^2 \int_{-\infty}^{+\infty} e^{-at^2} dt$$

$$\frac{d}{da} \int_{-\infty}^{+\infty} dx e^{-ax^2} = \int_{-\infty}^{+\infty} -x^2 e^{-ax^2} dx = -\frac{\pi}{2} a^{-\frac{3}{2}}$$

$$\therefore I_1 = \frac{\sqrt{\pi}}{2} a^{-\frac{3}{2}} + b^2 \pi a^{-\frac{1}{2}}$$

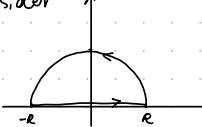
$$[P9] 3] I_2 = \int dz e^{-\alpha(x^2+y^2+z^2)} [(x+b_1)^2 + (y+b_2)^2 + (z+b_3)^2]$$

$$\int dx e^{-\alpha x^2} (x+b_1)^2 \int dy e^{-\alpha y^2} \int dz e^{-\alpha z^2} = I_1^2 \frac{\pi}{\alpha} \Rightarrow I_2 = 8\pi^2 \left(\frac{1}{\alpha} + b^2 \right)^2$$

P9) 4 ~

$$\text{第2 P9)} \quad \text{let } f(z) = \frac{z}{z^2 - (p+iq)^2}$$

$$[P9] 1] \text{ The pole of function are at } z=p+iq \text{ and } z=-p-iq, \text{ consider}$$

$\oint dz \frac{z}{z^2 - (p+iq)^2} e^{iaz}$	$z=Re^{i\theta}$	
	$dz = e^{i\theta} d\theta$	
$I_R = \int_0^\pi iRe^{i\theta} f(z) e^{iaR\cos\theta - iaR\sin\theta} d\theta$	$dz = iRe^{i\theta} d\theta$	

when $R \rightarrow \infty$ $|f(z)| = |f(Re^{i\theta})| < \epsilon$ for all θ within the integration range.

$\epsilon \rightarrow 0$. Then,

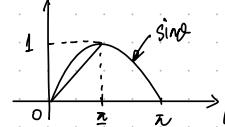
$$|I_R| \leq \epsilon R \int_0^\pi e^{-aR\sin\theta} d\theta = 2\epsilon R \int_0^{\frac{\pi}{2}} e^{-aR\sin\theta} d\theta, \quad (\alpha > 0)$$

since $\frac{2}{\pi}\theta \leq \sin\theta$

$$\leq 2\epsilon R \int_0^{\frac{\pi}{2}} e^{-2aR\theta/\pi} d\theta$$

$$= 2\epsilon/a (1 - e^{-2aR}) < \frac{\pi}{\alpha} \epsilon$$

$$\Rightarrow \lim_{R \rightarrow \infty} I_R = 0$$



$$\text{Therefore } \int_{-\infty}^{+\infty} f(x) e^{iax} = 2\pi i \operatorname{Res}[f(x)e^{iax}, p+iq] = 2\pi i \lim_{x \rightarrow p+iq} \frac{x e^{-iax}}{x^2 - (p+iq)^2} (x - (p+iq))$$

$$= \pi i e^{-ip+aq}$$

$$[P9] 2] \quad f(z) = e^{iz^2} = \cos(z^2) + i\sin(z^2),$$

$$\text{consider } \oint_{C_1+C_2+C_3} f(z) dz = 0$$

$$C_1 = \int_{C_1} f(z) dz = \int_0^R \cos(r^2) dr + i \int_0^R \sin(r^2) dr \quad z = re^{i\theta} = r, dz = dr$$

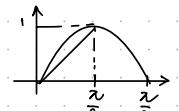
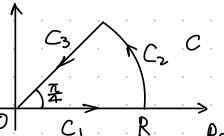
$$= \int_0^R \cos(r^2) dr + i \int_0^R \sin(r^2) dr$$

$$\lim_{R \rightarrow \infty} C_1 = \frac{1}{2} (I_2 + iI_3)$$

$$C_2 = \int_{C_2} f(z) dz = \int_0^{\frac{\pi}{4}} e^{iR^2} e^{i\theta} iRe^{i\theta} d\theta = iR \int_0^{\frac{\pi}{4}} e^{iR^2} e^{i\theta} e^{iR^2 \cos(2\theta) - R^2 \sin(2\theta)} d\theta, dz = iRe^{i\theta} d\theta$$

$$= iR \int_0^{\frac{\pi}{4}} e^{-R^2 \sin(2\theta)} d\theta \leq iR \int_0^{\frac{\pi}{4}} e^{-R^2 \frac{\pi}{4}} d\theta = iR \cdot \frac{\pi}{4R^2} (1 - e^{-R^2})$$

$$\lim_{R \rightarrow \infty} |C_2| \rightarrow 0$$



$$C_3 = \int_{C_3} f(z) dz = \int_R^0 e^{i(r e^{i\frac{\pi}{4}})^2} e^{i\frac{\pi}{4}} dr$$

$$= \int_R^0 \left(\frac{\sqrt{2}}{2} + i\frac{\sqrt{2}}{2} \right) e^{-r^2} dr \quad \lim_{R \rightarrow \infty} C_3 = \frac{\sqrt{2}}{2} (1+i)(-i)$$

$$z = re^{i\frac{\pi}{4}}, dz = e^{i\frac{\pi}{4}} dr$$

$$\text{Therefore: } \frac{1}{2}(I_2 + iI_3) + 0 - \frac{\sqrt{2}}{4} (1+i) = 0$$

$$\Rightarrow I_2 = \frac{\pi}{2}, I_3 = \frac{\pi}{2}$$

第3問

$$[P9] 1] |A - \lambda I| = \begin{vmatrix} -\lambda & \frac{1}{2} & 0 \\ \frac{1}{2} & -\lambda & \frac{1}{2} \\ 0 & \frac{1}{2} & -\lambda \end{vmatrix} = -\lambda(\lambda^2 - \frac{1}{4}) - \frac{1}{2} \left(\frac{1}{\lambda} \right) = 0 \Rightarrow \begin{cases} \lambda_1 = 0 \\ \lambda_2 = 1 \\ \lambda_3 = -1 \end{cases}$$

$$\textcircled{1} \lambda_1 = 0, (A - \lambda I) \vec{x} = \begin{pmatrix} 0 & \frac{1}{2} & 0 \\ \frac{1}{2} & 0 & \frac{1}{2} \\ 0 & \frac{1}{2} & 0 \end{pmatrix} \begin{pmatrix} x_1 \\ x_2 \\ x_3 \end{pmatrix} = 0 \Rightarrow \vec{x}_1 = \begin{pmatrix} \frac{1}{2} \\ 0 \\ -\frac{1}{2} \end{pmatrix}$$

$$\textcircled{2} \lambda_2 = 1, (A - \lambda I) \vec{x} = \begin{pmatrix} -1 & \frac{1}{2} & 0 \\ \frac{1}{2} & -1 & \frac{1}{2} \\ 0 & \frac{1}{2} & -1 \end{pmatrix} \begin{pmatrix} x_1 \\ x_2 \\ x_3 \end{pmatrix} = 0 \Rightarrow \vec{x}_2 = \begin{pmatrix} \frac{1}{2} \\ \frac{1}{2} \\ \frac{1}{2} \end{pmatrix}$$

$$\textcircled{3} \lambda_3 = -1, (A - \lambda I) \vec{x} = \begin{pmatrix} 1 & \frac{1}{2} & 0 \\ \frac{1}{2} & 1 & \frac{1}{2} \\ 0 & \frac{1}{2} & 1 \end{pmatrix} \begin{pmatrix} x_1 \\ x_2 \\ x_3 \end{pmatrix} = 0 \Rightarrow \vec{x}_3 = \begin{pmatrix} \frac{1}{2} \\ -\frac{1}{2} \\ \frac{1}{2} \end{pmatrix}$$

$$[P9] 2] \quad A = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & -1 \end{pmatrix} \quad O^{-1}AO = A, \quad O = \begin{pmatrix} \frac{1}{2} & \frac{i}{2} & \frac{1}{2} \\ \frac{i}{2} & \frac{1}{2} & -\frac{i}{2} \\ \frac{1}{2} & -\frac{i}{2} & \frac{1}{2} \end{pmatrix}$$

$$O^{-1} = \begin{pmatrix} 0 & 0 & 0 \end{pmatrix}$$

$$[P9] 3] O^{-1}AOO^{-1}A \cdots AO = O^{-1}A^n O = A^n$$

$$\Rightarrow A^n = O \Lambda^n O^{-1} \quad \dots$$

第4問

$$[P9] 1] e^{ax} \frac{d}{dx} (e^{-ax} f(x)) = e^{ax} (-ae^{-ax} f(x) + \frac{df}{dx} e^{-ax}) = -af(x) + \frac{df}{dx} = (\hat{D} - a)f(x)$$

$$[P9] 2] \quad z^2 + 2pz + q = 0 \quad \text{solution: } \alpha \cdot \beta = \frac{-2p \pm \sqrt{4p^2 - 4q}}{2}, \quad \alpha = -p + \sqrt{p^2 - q}, \quad \beta = -p - \sqrt{p^2 - q}$$

$$y_1 = C_1 e^{\alpha x} + C_2 e^{\beta x} = C_1 e^{(-p + \sqrt{p^2 - q})x} + C_2 e^{(-p - \sqrt{p^2 - q})x}$$

[問3] $y_2 = C_1 e^{i(-\gamma + \sqrt{\gamma^2 - q})x} + C_2 e^{i(-\gamma - \sqrt{\gamma^2 - q})x}$

[問4] $y_3 = C_1 e^{-\gamma x}$

[問5] $\bullet \bullet \bullet$

SOKENDAI KEK 入試 2019 解答

物理量

第1問

[問1]

$$(1) \hat{H} = -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} + \frac{m\omega^2}{2} x^2$$

$$(2) \hat{H}|\psi_0\rangle = \left(-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} + \frac{m\omega^2}{2} x^2 \right) \alpha e^{-bx^2} = \left\{ -\frac{\hbar^2}{2m} [-2b + (-2bx)^2] + \frac{m\omega^2}{2} x^2 \right\} |\psi_0\rangle = E_0 |\psi_0\rangle$$

$$-\frac{\hbar^2}{2m} 4b^2 + \frac{m\omega^2}{2} = 0 \Rightarrow b = \frac{m\omega}{2\hbar}, \quad E_0 = \frac{\hbar^2 b}{m} = \frac{\hbar\omega}{2}$$

$$\hat{H}|\psi_1\rangle = \left(-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} + \frac{m\omega^2}{2} x^2 \right) cx e^{-bx^2} = \left\{ -\frac{\hbar^2}{2m} [-bb + (-2bx)^2] + \frac{m\omega^2}{2} x^2 \right\} |\psi_1\rangle = E_1 |\psi_1\rangle,$$

$$E_1 = \frac{3\hbar^2 b}{m} = \frac{3}{2}\hbar\omega$$

$$(3) \int \psi_0^* \psi_0 dx = \alpha^2 \int e^{-2bx^2} dx = \alpha^2 \sqrt{\frac{\pi}{2b}} = 1 \Rightarrow \alpha = \left(\frac{2b}{\pi}\right)^{\frac{1}{4}}$$

$$\int \psi_1^* \psi_1 dx = c^2 \int x^2 e^{-2bx^2} dx = \frac{1}{4} \sqrt{\frac{\pi}{2}} b^{-\frac{3}{2}} c^2 = 1 \Rightarrow c = 2 \left(\frac{2b^3}{\pi}\right)^{\frac{1}{4}}$$

[問2]

$$(1) \hat{H}' = -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} + \frac{m\omega'^2}{2} x^2 \quad |\psi'_0\rangle = \left(\frac{m\omega'}{\pi\hbar}\right)^{\frac{1}{4}} e^{-\frac{m\omega'}{2\hbar}x^2}$$

$$(2) \int_0^\infty |\psi'_0|^2 dx = 1 \Rightarrow \alpha = \left(\frac{m\omega}{4\hbar\pi}\right)^{\frac{1}{4}} \quad \tilde{\psi}_0 = \left(\frac{m\omega}{4\hbar\pi}\right)^{\frac{1}{4}} e^{-\frac{m\omega}{2\hbar}x^2}, \quad E_0 = \frac{\hbar\omega}{2}$$

[問3]

$$\begin{cases} i\hbar \dot{\phi}(x,t) = \hat{H}(\phi(x)) \phi(t) \\ i\hbar \frac{d}{dt} \phi(x,t) = \hat{H}\phi(x,t) \end{cases} \Rightarrow i\hbar \dot{\phi}(x) \frac{d\phi(t)}{dt} = \hat{H}\phi(x)\phi(t)$$

$$i\hbar \frac{1}{\phi(t)} \frac{d\phi(t)}{dt} = \frac{1}{\phi(x)} \hat{H}\phi(x) = E$$

$$\Rightarrow \begin{cases} \frac{d\phi(t)}{dt} = \frac{i}{\hbar} E \phi(t) \\ \hat{H}\phi(x) = E\phi(x) \end{cases} \quad \therefore \phi(t) = e^{-iEt/\hbar}$$

$$\therefore \psi(x,t) = \sum_{n=0}^{\infty} a_n \psi_n(x) e^{-iEnt/\hbar}$$

$$(2) e^{-iEnt/\hbar} = e^{-i(Ent/\hbar + 2\pi)} \Rightarrow T = 2\pi\hbar/E_h$$

(3) ?

第2問

[PB1]

$$H = \sum_{i=1}^N \frac{p_i^2}{2m} + V(r) + f \cdot r$$

[PB2]

$$\begin{aligned} Z &= (2\pi k_B T)^N \int d\mathbf{r} d^N p e^{-H/k_B T} = (2\pi k_B T)^N \int d\mathbf{r} d^N p e^{-\sum \frac{p_i^2}{2mk_B T}} \int d^3 r e^{-\frac{V(r)+f \cdot r}{k_B T}} \\ &= (2\pi k_B T)^N (2\pi k_B T N)^{\frac{N}{2}} \left[\int_0^{+\infty} e^{-\frac{(b+f)r-ab}{k_B T}} dp \right]^N \\ &= \left(\frac{2\pi k_B T}{2\pi k_B T^2} \right)^{\frac{N}{2}} \left(\frac{k_B T}{b+f} \right)^N e^{\frac{abN}{k_B T}} e^{-\frac{a(b+f)N}{k_B T}} \\ &= \left(\frac{m}{2\pi k_B T} \right)^{\frac{N}{2}} \left(\frac{k_B T}{b+f} \right)^{\frac{N}{2}} e^{-\frac{Nfa}{k_B T}} \end{aligned}$$

[PB3]

$$F = -k_B T \ln Z = -k_B T N \left\{ \frac{1}{2} \ln \frac{m}{2\pi k_B T^2} + \ln \frac{(k_B T)^{\frac{3}{2}}}{b+f} - \frac{fa}{k_B T} \right\}$$

$$U = -k_B T^2 \frac{\partial}{\partial T} \left(\frac{F}{k_B T} \right) = k_B T^2 N \frac{\partial}{\partial T} \left(\frac{3}{2} \ln k_B T - \frac{fa}{k_B T} \right) = -\frac{3}{2} N k_B T + fa N$$

$$\begin{aligned} S &= -\frac{\partial F}{\partial T} = k_B N \left\{ \frac{1}{2} \ln \frac{m}{2\pi k_B T^2} + \ln \frac{(k_B T)^{\frac{3}{2}}}{b+f} - \frac{fa}{k_B T} \right\} + T \left[\frac{3}{2} \frac{1}{T} + \frac{fa}{k_B T^2} \right] \\ &= k_B N \left\{ \frac{1}{2} \ln \frac{m}{2\pi k_B T^2} + \ln \frac{(k_B T)^{\frac{3}{2}}}{b+f} + \frac{3}{2} \right\} \end{aligned}$$

[PB4]

[PB5]

第3問

[PB1]

$$\dot{x} = r \cos\theta - \dot{\theta} r \sin\theta, \quad \dot{y} = r \sin\theta + \dot{\theta} r \cos\theta$$

$$\mathcal{L} = \frac{1}{2} m \{ (r \cos\theta - \dot{\theta} r \sin\theta)^2 + (r \sin\theta + \dot{\theta} r \cos\theta)^2 + z^2 \} - mgz$$

$$= \frac{1}{2} m (r^2 + \dot{\theta}^2 r^2 + z^2) - mgz$$

[PB2]

$$(1) \quad \dot{r} = -\dot{\theta} r \sin\theta, \quad \dot{z} = \dot{\theta} r \cos\theta$$

$$\mathcal{L} = \frac{1}{2} m (\dot{r}^2 + \dot{\theta}^2 r^2 + z^2) - mgz$$

$$= \frac{1}{2} m [\dot{r}^2 + \dot{\theta}^2 (R + a \cos\theta)^2] - mg a \sin\theta$$

$$(2) \quad L_z = r \dot{\theta} y - \dot{\theta} r x = m (r \cos\theta \cdot (r \sin\theta + \dot{\theta} r \cos\theta) - r \sin\theta (r \cos\theta - \dot{\theta} r \sin\theta))$$

$$= m r^2 \dot{\theta} = m (R + a \cos\theta)^2 \dot{\theta}$$

$$\frac{\partial \mathcal{L}}{\partial \dot{\theta}} = 0, \quad \frac{\partial \mathcal{L}}{\partial \theta} = m \dot{\theta} (R + a \cos\theta)^2 \Rightarrow \frac{d}{dt} \left(\frac{\partial \mathcal{L}}{\partial \dot{\theta}} \right) = \frac{d}{dt} L_z = 0$$

Therefore, L_z is conserved.

$$(3) \quad \frac{\partial \mathcal{L}}{\partial \dot{\theta}} = -m a \dot{\theta}^2 \sin\theta (R + a \cos\theta) - m g a \cos\theta, \quad \frac{\partial \mathcal{L}}{\partial \theta} = m a^2 \dot{\theta}$$

$$\Rightarrow m a^2 \ddot{\theta} = -m a \dot{\theta}^2 \sin\theta (R + a \cos\theta) - m g a \cos\theta$$

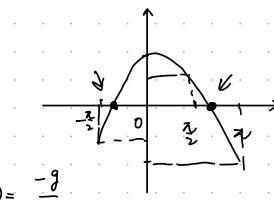
$$\text{Since } m \frac{d}{dt} (R + a \cos\theta)^2 \dot{\theta} = 0 \Rightarrow m \dot{\theta} (R + a \cos\theta)^2 = C_1$$

$$\therefore f(\theta) = \frac{1}{a} \cdot \frac{C_1 \sin\theta}{m(R+a\cos\theta)} + \frac{g}{a} \cos\theta$$

[PB3]

$$(1) \quad \frac{df}{d\theta} = 0 \Rightarrow f(\theta) = 0$$

$$f(0) = \frac{g}{a}, \quad f(-\frac{\pi}{2}) = \frac{-C_1}{mRa}, \quad f(\frac{\pi}{2}) = \frac{C_1}{mRa}, \quad f(\pi) = -\frac{g}{a}$$



$$(2) \quad f(\phi_0) = 0$$

$$\frac{d^2(f(\phi_0 + \delta))}{d\delta^2} = \ddot{f} = -f(\phi_0 + \delta) = -f(\phi_0) - f'(\phi_0) \delta$$

$$\Rightarrow \ddot{\delta} = -f'(\phi_0) \delta$$

(3)

$$\ddot{\delta} = A_1 \cos[\sqrt{f'(\phi_0)} t] + A_2 \sin[\sqrt{f'(\phi_0)} t]$$

$$\ddot{\delta} = -A_1 \sqrt{f'(\phi_0)} \sin[\sqrt{f'(\phi_0)} t] + A_2 \sqrt{f'(\phi_0)} \cos[\sqrt{f'(\phi_0)} t]$$

$$\ddot{\delta}(0) = A_2 \sqrt{f'(\phi_0)} = 0 \Rightarrow A_2 = 0$$

$$\ddot{\delta}(t) = A_1 \cos[\sqrt{f'(\phi_0)} t]$$

$|\ddot{\delta}(0)| = |A_1| \ll 1$, Therefore δ is oscillating around zero.

第4問

[PB1]

$$\text{rot rot } \vec{E} = \text{grad div } \vec{E} - \Delta \vec{E} = -\Delta \vec{E}$$

$$\text{rot} \left(-\frac{2}{\sigma \epsilon_0} \vec{B} \right) = -\epsilon_0 \mu_0 \frac{\partial}{\partial t} \vec{E} \Rightarrow \Delta \vec{E} - \epsilon_0 \mu_0 \frac{\partial^2}{\partial t^2} \vec{E} = 0,$$

take $E_x(x, y, z, t) = E_x(y, z) e^{i(\omega t - kx)}$ into,

$$\left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2}\right) E_x(y, z) e^{i(wt-kx)} - \epsilon_0 \mu_0 \frac{\partial^2}{\partial t^2} E_x(y, z) e^{i(wt-kx)} = 0$$

$$\left(\frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2}\right) E_x + (-k^2) E_x + \epsilon_0 \mu_0 \omega^2 E_x = 0 \Rightarrow \frac{\partial^2 E_x}{\partial y^2} + \frac{\partial^2 E_x}{\partial z^2} + k^2 E_x = 0$$

[P92]

$$E_x|_{y=0, a=0} = 0 \quad (1)$$

$$E_x|_{z=0, b=0} = 0 \quad (2)$$

[P93]

$$-\left[\left(\frac{m\pi}{a}\right)^2 + \left(\frac{n\pi}{b}\right)^2\right] E_x + k^2 E_x = 0 \Rightarrow \frac{\omega^2}{c^2} - k^2 = \left(\frac{m\pi}{b}\right)^2 + \left(\frac{n\pi}{a}\right)^2$$

$$\text{when } m=n=1, \omega = \sqrt{k^2 + \pi^2 \left(\frac{1}{a^2} + \frac{1}{b^2}\right)}$$

[P94]

$$\nabla \times \vec{E} = \begin{vmatrix} \hat{e}_x & \hat{e}_y & \hat{e}_z \\ \partial_x & \partial_y & \partial_z \\ E_x & E_y & E_z \end{vmatrix} = (\partial_y E_z - \partial_z E_y) \hat{e}_x + (\partial_z E_x - \partial_x E_z) \hat{e}_y + (\partial_x E_y - \partial_y E_x) \hat{e}_z$$

$$-\frac{\partial \vec{B}}{\partial t} = -i\omega \vec{B}, B_x = 0 \Rightarrow \partial_y E_z - \partial_z E_y = 0$$

$$\left\{ \begin{array}{l} -i\omega B_y = \partial_z E_x - \partial_x E_z = \partial_z E_x + ik E_z \\ -i\omega B_z = \partial_x E_y - \partial_y E_x = -ik E_y - \partial_y E_x \end{array} \right. \quad (1)$$

$$\left\{ \begin{array}{l} -i\omega B_2 = \partial_x E_y - \partial_y E_x = -ik E_y - \partial_y E_x \\ \nabla \times \vec{B} = (\partial_y B_z - \partial_z B_y) \hat{e}_x + (\partial_z B_x - \partial_x B_z) \hat{e}_y + (\partial_x B_y - \partial_y B_x) \hat{e}_z \end{array} \right. \quad (2)$$

$$\mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t} = i\omega \mu_0 \epsilon_0 \vec{E}$$

$$\left\{ \begin{array}{l} i\omega \mu_0 \epsilon_0 E_y = -\partial_x B_2 = ik B_z \\ i\omega \mu_0 \epsilon_0 E_z = \partial_x B_y = -ik B_y \end{array} \right. \quad (3) \quad (4)$$

From (1) (2) (3) (4), we can have

$$-i\omega B_y = \partial_z E_x + ik \frac{-k B_y}{\omega \mu_0 \epsilon_0} \Rightarrow B_y = \frac{\partial_z E_x}{-i\omega + \frac{ik^2}{\omega \mu_0 \epsilon_0}}$$

$$-i\omega B_z = -ik \frac{k B_z}{\omega \mu_0 \epsilon_0} - \partial_y E_x \Rightarrow B_z = \frac{-\partial_y E_x}{-i\omega + \frac{ik^2}{\omega \mu_0 \epsilon_0}}$$

$$E_y = \frac{k}{\omega \mu_0 \epsilon_0} \cdot \frac{-\partial_y E_x}{-i\omega + \frac{ik^2}{\omega \mu_0 \epsilon_0}} = \frac{-k \partial_y E_x}{-i\omega^2 \mu_0 \epsilon_0 + ik^2}$$

$$E_z = \frac{-k}{\omega \mu_0 \epsilon_0} \cdot \frac{\partial_z E_x}{-i\omega + \frac{ik^2}{\omega \mu_0 \epsilon_0}} = \frac{-k \partial_z E_x}{-i\omega^2 \mu_0 \epsilon_0 + ik^2}$$

数学

第1問

[P91]

$$|A - \lambda I| = \begin{vmatrix} -\lambda & m \\ m & M - \lambda \end{vmatrix} = -\lambda(M - \lambda) - m^2 = 0 \Rightarrow \lambda = \frac{M \pm \sqrt{M^2 + 4m^2}}{2}$$

[P92]

$$(A - \lambda I) \vec{x} = 0$$

$$\begin{pmatrix} -\lambda & m \\ m & M - \lambda \end{pmatrix} \begin{pmatrix} x_1 \\ x_2 \end{pmatrix} = 0 \Rightarrow x_1 + x_2 \begin{pmatrix} m \\ \lambda_+ \end{pmatrix}, x_1 - x_2 \begin{pmatrix} m \\ \lambda_- \end{pmatrix} \Rightarrow P = \begin{pmatrix} m & m \\ \lambda_+ & \lambda_- \end{pmatrix}$$

[P93]

$$m = M\epsilon, \lambda_- = \frac{M - \sqrt{M^2 + 4m^2 \epsilon^2}}{2} = \frac{M}{2}(1 - \sqrt{1 + 4\epsilon^2})$$

$$\frac{d\lambda_-}{d\epsilon} = -\frac{M}{2} \frac{1}{2} \cdot 8\epsilon (1 + 4\epsilon^2)^{-\frac{1}{2}} = -2M\epsilon(1 + 4\epsilon^2)^{-\frac{1}{2}}$$

$$\frac{d^2\lambda_-}{d\epsilon^2} = -2M(1 + 4\epsilon^2)^{-\frac{1}{2}} + 8M\epsilon^2(1 + 4\epsilon^2)^{-\frac{3}{2}}$$

$$\lambda_- = \frac{M}{2} - M\epsilon$$

第2問

[P91]

$$\int_{-\infty}^{\infty} dx \int_{\alpha}^{\infty} dy e^{-\frac{(x-y)^2}{4}} = \int_0^{\infty} dr \int_0^{\infty} d\theta r e^{-r^2} = 2\pi \frac{1}{2} \int_0^{\infty} dr (r^2) e^{-r^2} = \pi$$

$$[P92] \int_{-\infty}^{+\infty} dx e^{-x^2} = \sqrt{\pi} \Rightarrow \int_{-\infty}^{+\infty} dx e^{-\frac{1}{4}x^2} = \frac{2}{\sqrt{\pi}} \int_{-\infty}^{\infty} d\left(\frac{1}{2}x\right) e^{-\left(\frac{1}{2}x\right)^2} = \sqrt{\pi}$$

[P93]

$$\frac{\partial^n}{\partial x^n} \int_{-\infty}^{+\infty} e^{-\alpha x^2} dx = (-1)^n \int_{-\infty}^{+\infty} x^{2n} e^{-\alpha x^2} dx \Rightarrow \int_{-\infty}^{+\infty} x^{2n} e^{-\alpha x^2} dx = \sqrt{\pi} \frac{(2n-1)!!}{2^n} \alpha^{-\frac{1}{2}(n+1)}$$

$$\frac{\partial^n}{\partial x^n} \left[\frac{\pi}{\alpha} \right] = \sqrt{\pi} (-1)^n \frac{(2n-1)!!}{2^n} \alpha^{-\frac{1}{2}(n+1)}$$

$$\therefore \int_{-\infty}^{+\infty} dx x^{2n} e^{-\frac{1}{2}x^2} = \sqrt{\pi} \frac{(2n-1)!!}{2^n} \left(\frac{1}{2}\right)^{-\frac{1}{2}(n+1)}$$

第3問

$$[P91] z = e^{i\theta}, z^6 + 1 = e^{6i\theta} + 1 = 0, \theta = \frac{\pi + 2k\pi}{6}, z = e^{i\frac{\pi + 2k\pi}{6}}, k=0,1,2$$

$$[P92] \theta = \frac{2\pi}{n}, |z|=R, z = R e^{i\theta}, \text{ consider } f(z) = \frac{1}{z^n+1}$$

$$\oint_{C_1+C_2+C_3} f(z) dz = 2\pi i \operatorname{Res}[f(z), z = e^{i\pi/n}]$$

Along C_2 , as $R \rightarrow \infty$, $\int_{C_2} f(z) dz \rightarrow 0$

Along C_1 , $\int_1^R \frac{1}{x^n + 1} dx$

Along C_3 , $\int_R^0 \frac{1}{(xe^{i\frac{\pi}{n}})^n + 1} e^{i\frac{2\pi}{n}} dx = -e^{i\frac{2\pi}{n}} \int_0^R \frac{1}{x^n + 1} dx$

when $R \rightarrow \infty$, therefore $(1 - e^{i\frac{2\pi}{n}}) \int_0^\infty \frac{1}{x^n + 1} dx = 2\pi i \operatorname{Res}[f(z), z = e^{i\frac{\pi}{n}}]$

$$\text{let } I(n) = \int_1^\infty \frac{1}{x^n + 1} dx \Rightarrow I(n) = 2\pi i \frac{1}{1 - e^{2\pi i/n}} \lim_{z \rightarrow e^{i\frac{\pi}{n}}} \frac{z - e^{i\frac{\pi}{n}}}{z^n + 1} = \frac{2\pi i}{-n} \frac{1}{e^{-in\pi} - e^{in\pi}} = \frac{\pi}{n \sin(\frac{\pi}{n})}$$

$$\text{when } n=6, \int_{-\infty}^{+\infty} \frac{1}{x^6 + 1} dx = \frac{2\pi}{6 \sin(\frac{\pi}{6})} = \frac{2\pi}{3}$$

第4問

$$\begin{aligned} (\text{P01}) \quad y'(x) &= e^{fx} g(x) e^{fx} + \left(-\frac{df(x)}{dx}\right) e^{-fx} / dx g(x) e^{fx} \\ &= g(x) - f(x)y \end{aligned}$$

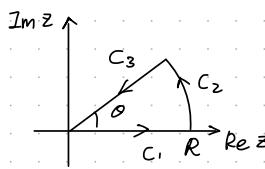
[P02]

$$f(x) = -x \Rightarrow f(x) = -\frac{1}{2}x^2 + C_1$$

$$g(x) = x$$

$$\begin{aligned} \therefore y &= e^{\frac{1}{2}x^2 - C_1} \int dx \cdot x e^{-\frac{1}{2}x^2 + C_1} = e^{\frac{1}{2}x^2} \int dx \cdot x e^{-\frac{1}{2}x^2} = e^{\frac{1}{2}x^2} - \int d(-\frac{1}{2}x^2) e^{-\frac{1}{2}x^2} \\ &= e^{\frac{1}{2}x^2} \left(-e^{-\frac{1}{2}x^2} + C_1\right) \end{aligned}$$

$$y(0) = (1 + C_1) = 0 \Rightarrow C_1 = 1 \quad \therefore y = e^{\frac{1}{2}x^2} - 1$$



SOKENDAI KEK 2018 解答

物理

第1問

[問1]

$$(1) \quad q \vec{v} \times \vec{B} = q \begin{pmatrix} \vec{e}_x & \vec{e}_y & \vec{e}_z \\ v_x & v_y & v_z \\ 0 & 0 & B \end{pmatrix} = q v_y B \vec{e}_x - q v_x B \vec{e}_y$$

$$\begin{cases} m \frac{dv_x}{dt} = q v_y B \\ m \frac{dv_y}{dt} = -q v_x B \end{cases} \Rightarrow \frac{d^2 v_x}{dt^2} = m \frac{dv_y}{dt} = -m^2 v_x \Rightarrow \frac{dx}{dt} = v_x = A \cos(\omega t + \alpha) \\ \frac{dy}{dt} = v_y = \frac{1}{\omega} \frac{dv_x}{dt} = -A \sin(\omega t + \alpha) \end{math>$$

$$\therefore \begin{cases} x = R \sin(\omega t + \alpha) + x_c \\ y = R \cos(\omega t + \alpha) + y_c \end{cases}$$

$$\begin{aligned} (2) \quad &\begin{cases} v_x = \omega R \cos(\omega t + \alpha) \\ v_y = -\omega R \sin(\omega t + \alpha) \end{cases} \\ &\begin{cases} x_c = x - R \sin(\omega t + \alpha) = \pi + \frac{1}{\omega} v_y \\ y_c = y - \frac{1}{\omega} v_x \end{cases} \end{aligned}$$

$$\begin{aligned} \frac{dx_c}{dt} &= \frac{dx}{dt} + \frac{1}{\omega} \frac{dv_y}{dt} \\ &= v_x + \frac{1}{\omega} \cdot (-\omega v_x) \\ &= 0 \\ \frac{dy_c}{dt} &= 0 \end{aligned}$$

$$\star [\hat{A}, \hat{f}(\hat{B})] = [\hat{A}, \hat{B}] \frac{\partial f}{\partial B}$$

[問2]

$$\begin{aligned} (1) \quad [\hat{p}_x, \hat{p}_y] &= [p_x - qA_x, p_y - qA_y] = [p_x - qA_x, p_y] - [p_x - qA_x, qA_y] \\ &= -q[A_x, p_y] - q[p_x, A_y] \\ &= q[p_y, q] \frac{\partial A_x}{\partial y} - q[p_x, q] \frac{\partial A_y}{\partial x} \\ &= -i\hbar q \frac{\partial A_x}{\partial y} + i\hbar q \frac{\partial A_y}{\partial x} \\ &= i\hbar q_B B \end{aligned}$$

$$(2) \quad \hat{H}_1 = \frac{1}{2m} \hat{p}_{\perp}^2 = \frac{1}{2m} (\hat{p}_x^2 + \hat{p}_y^2) = \frac{qB}{2m} (\hat{x}^2 + \hat{p}^2)$$

$$(3) \quad \hat{a}^{\dagger} \hat{a} = \frac{1}{2\hbar} (\hat{x}^2 + i\hat{x}\hat{p} - i\hat{p}\hat{x} + \hat{p}^2) = \frac{1}{2\hbar} (\hat{x}^2 + \hat{p}^2) - \frac{1}{2}$$

$$\hat{H}_1 = \hbar\omega (\hat{a}^{\dagger} \hat{a} + \frac{1}{2}), \quad \text{Since } \hat{a}^{\dagger} \hat{a} \hat{\psi}_n = \hbar\omega_n, \quad \hat{H}_1 \hat{\psi}_n = E_n \hat{\psi}_n, \\ \text{Therefore, } E_n = \hbar\omega (n + \frac{1}{2})$$

$$(4) O = \langle \hat{x} | \hat{a} | 0 \rangle = \langle \hat{x} | \frac{i}{\hbar \hbar} (\hat{x} + i \hat{p}) | 0 \rangle = \frac{i}{\hbar \hbar} [\langle \hat{x} | \hat{x} | 0 \rangle + i \langle \hat{x} | \hat{p} | 0 \rangle] \\ = \frac{i}{\hbar \hbar} [x \psi_0(x) + i (-i \hbar \frac{d}{dx}) \psi_0(x)]$$

$$\Rightarrow x \psi_0(x) = -\hbar \frac{d}{dx} \psi_0(x) \Rightarrow \frac{1}{\hbar} x d\psi_0 = -\frac{1}{\hbar} \chi dx \Rightarrow \psi_0(x) = N \exp \left\{ -\frac{\chi^2}{2\hbar} \right\}$$

$$(5) [\hat{x}_c, \hat{H}_\perp] = [\hat{x} + \frac{1}{qB} \hat{T}y, \frac{1}{2m} (\hat{p}_x^2 + \hat{p}_y^2)] = \frac{1}{2m} \{ [\hat{x}, \hat{p}_x^2 + \hat{p}_y^2] + \frac{1}{qB} [\hat{T}y, \hat{p}_x^2 + \hat{p}_y^2] \}$$

$$[\hat{x}, \hat{p}_x^2] = 2[\hat{x}, \hat{p}_x] \hat{p}_x = 2[\hat{x}, \hat{p}_x - qA_x] \hat{p}_x = 2i\hbar \hat{T}x$$

$$[\hat{x}, \hat{p}_y^2] = 2[\hat{x}, \hat{p}_y - qA_y] \hat{p}_y = 0$$

$$[\hat{T}y, \hat{p}_x^2] = 2[\hat{T}y, \hat{p}_x] \hat{p}_x = -2i\hbar qB \hat{T}x$$

$$\text{Then, } [\hat{x}_c, \hat{H}_\perp] = \frac{1}{2m} \{ 2i\hbar \hat{T}x + \frac{1}{qB} (-2i\hbar qB) \hat{T}x \} = 0$$

$$\text{Same as } [\hat{y}_c, \hat{H}_\perp] = 0$$

$$(6) \hat{R}^2 \hat{n} = (\hat{x} - \hat{x}_c)^2 \hat{n} + (\hat{y} - \hat{y}_c)^2 \hat{n} = (\frac{1}{qB})^2 (\hat{p}_y^2 + \hat{p}_x^2) \hat{n} = 2m (\frac{1}{qB})^2 \hat{H}_\perp \hat{n}$$

$$\Rightarrow \lambda_n = 2m (\frac{1}{qB})^2 + \frac{qB}{m} (n + \frac{1}{2}) = \frac{2\hbar}{qB} (n + \frac{1}{2})$$

第2題

$$[問1] \quad \text{Introduce } \Xi = \prod_i (1 - e^{-E_i/kT})^{-w_i}$$

$$E = PC, D(E) dE = \int_{\infty}^E \frac{d\omega}{h^3 c^3} = \frac{4\pi V}{h^3 c^3} \Xi^2 dE, \text{ let } x = \frac{E}{kT}$$

$$\ln \Xi_1 = -2 \int_{-\infty}^{+\infty} \frac{d\omega}{h^3} \ln (1 - e^{-E/kT}) = -\frac{8\pi V}{h^3 c^3} (kT)^3 \int_0^\infty x^2 \ln(1 - e^{-x}) dx$$

$$= -\frac{8\pi V}{h^3 c^3} (kT)^3 \left\{ \frac{1}{3} x^3 \ln(1 - e^{-x}) \Big|_0^\infty - \int_0^\infty \frac{1}{3} x^3 \frac{e^{-x}}{1 - e^{-x}} dx \right\}$$

$$= \frac{8\pi V}{h^3 c^3} (kT)^3 \cdot \frac{1}{3} \cdot \frac{\pi^4}{45} = \frac{8\pi^5 V}{45} \left(\frac{kT}{hc}\right)^3$$

$$F_1 = -kT \ln \Xi_1 = -\frac{1}{3} \alpha V T^4. \text{ Let } \alpha = \frac{8\pi^5 k^4}{45 (hc)^3}$$

$$\mathcal{U}_1 = -\frac{1}{V} T^2 \frac{\partial}{\partial T} \left(\frac{F_1}{T} \right) = \alpha T^4, P_1 = -\left(\frac{\partial F}{\partial V} \right)_T = \frac{1}{3} \alpha T^4$$

$$(2) \ln \Xi_2 = \frac{8\pi V}{h^3 c^3} (kT)^3 \cdot \frac{2}{8} \frac{\pi^4}{45} = \frac{2\pi^5 V}{45} \left(\frac{kT}{hc}\right)^3$$

$$F_2 = \frac{7}{8} F_1 \quad \therefore \mathcal{U}_2 = \frac{7}{8} \mathcal{U}_1 = \frac{7}{8} \alpha V T^4, P_2 = \frac{7}{8} P_1 = \frac{7}{24} \alpha V T^4$$

$$[問2] \quad S_r = -\left(\frac{\partial F}{\partial T} \right)_V \cdot \frac{1}{V} = CT^3 \quad S_c = \frac{1}{V} S_r = \frac{7}{8} CT^3$$

第3題

[問1]

$$(1) T = \frac{1}{2} m (\dot{x}^2 + \dot{y}^2 + \dot{z}^2)$$

$$(2) L = \frac{1}{2} m (\dot{x}^2 + \dot{y}^2 + \dot{z}^2) - U(x, y, z) \quad P_x = \frac{\partial L}{\partial \dot{x}} = m \dot{x}$$

$$H(x, y, z, P_x, P_y, P_z) = \frac{P_x^2 + P_y^2 + P_z^2}{2m} + U(x, y, z)$$

(3)

$$\begin{cases} x = r \sin \theta \cos \phi \\ y = r \sin \theta \sin \phi \\ z = r \cos \theta \end{cases} \quad \begin{cases} P_r = P_x \sin \theta \cos \phi + P_y \sin \theta \sin \phi - P_z \cos \theta \\ P_\theta = P_x r \cos \theta \cos \phi + P_y r \cos \theta \sin \phi - P_z r \sin \theta \\ P_\phi = -P_x r \sin \theta \sin \phi + P_y r \sin \theta \cos \phi \end{cases}$$

From Cartesian coordinate,
To spherical coordinate

$$(4) \dot{r} = \frac{\partial \mathcal{L}}{\partial P_r} = \frac{P_r}{m} \Rightarrow P_r = m \dot{r}$$

$$\dot{\theta} = \frac{P_\theta}{m r^2} \Rightarrow P_\theta = m r^2 \dot{\theta}$$

$$\dot{\phi} = \frac{P_\phi}{m r^2 \sin \theta} \Rightarrow P_\phi = m r^2 \sin \theta \dot{\phi}$$

$$\frac{dP_\phi}{dt} = -\frac{\partial \mathcal{L}}{\partial \dot{\phi}} = 0 \Rightarrow P_\phi \text{ is conserved}. \quad \frac{dP_\theta}{dt} = -\frac{1}{2m} P_\phi^2 \frac{1}{r^2} (-2) \cos \theta (\sin \theta)^{-3}$$

$$\frac{dL^2}{dt} = 2P_\theta \frac{dP_\theta}{dt} + \frac{d}{dt} \left(\frac{P_\theta^2}{r^2} \right)$$

$$= 2P_\theta \cdot \frac{dP_\theta}{dt} \cdot \frac{1}{\sin \theta} + \dot{\theta}^2 \frac{(-2)(\sin \theta)^{-3}}{\sin \theta} \cos \theta \dot{\theta} = 0$$

[問2]

$$(1) \frac{dP_r}{dt} = m \ddot{r}, P_\theta = m r^2 \dot{\theta} = 0, P_\phi = h, \sin \theta = 1$$

$$-\frac{\partial \mathcal{L}}{\partial r} = \frac{1}{2m} (P_\theta^2 \cos^2 \theta + \frac{P_\phi^2}{\sin^2 \theta} \cos^2 \theta) - \frac{dU}{dr} = \frac{h^2}{mr^3} - \frac{M}{r^2}$$

$$\Rightarrow m \ddot{r} = \frac{h^2}{mr^3} - \frac{M}{r^2}$$

$$(2) |r| = a, m \ddot{r} = \frac{h^2}{ma^3} - \frac{M}{a^2} \Rightarrow r = \frac{1}{2} \left(\frac{h^2}{ma^3} - \frac{M}{ma^2} \right) t^2 + A_1 t + A_2$$

$$\text{when } x = 2\pi a = \frac{1}{2} \left(\frac{h^2}{ma^3} - \frac{M}{ma^2} \right) T^2$$

$$\Rightarrow T = \sqrt{\frac{1}{4\pi} \left(\frac{h^2}{ma^4} - \frac{M}{ma^3} \right)}$$

?

$$(3) -\frac{d\dot{x}}{dt} = \frac{\hbar^2}{mr^3} - kr \Rightarrow m\ddot{r} = \frac{\hbar^2}{mr^3} - kr$$

$$\begin{cases} x = r \cos \phi \\ y = r \sin \phi \end{cases} \quad //$$

第4問

[PQ1]

$$(1) 2\pi r \cdot l \cdot E(r) = \lambda l \cdot 2\pi a / \epsilon_0 \Rightarrow E(r) = \frac{\alpha \lambda}{\epsilon_0 r}$$

$$\int_a^b E(r) dr = \frac{\alpha \lambda}{\epsilon_0} \int_a^b \frac{1}{r} dr = \frac{\alpha \lambda}{\epsilon_0} \ln \frac{b}{a} = V \Rightarrow \lambda = \frac{\epsilon_0 V}{\alpha \ln \frac{b}{a}}$$

$$\therefore E(r) = \frac{V}{r \ln \frac{b}{a}}$$

$$(2) C_0 = \frac{Q_0}{V} = \frac{2\pi a \lambda}{V} \quad \therefore U = \frac{1}{2} \frac{Q_0}{V} V^2 = \frac{\pi \epsilon_0 V^2}{\ln \frac{b}{a}} \quad (\text{quant length energy})$$

$$(3) \Delta U = F \Delta z, \text{ when } \Delta z = 1.$$

$$F = \frac{\pi \epsilon_0 V^2}{\ln \frac{b}{a}} - \frac{\pi \epsilon_0 V^2}{\ln \frac{b}{a}} = \frac{\pi V^2}{\ln \frac{b}{a}} (\epsilon - \epsilon_0)$$

[PQ2]

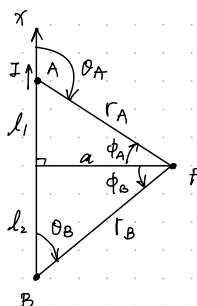
$$(1) \vec{\phi}_c \vec{H} d\vec{s} = H \cdot 2\pi r$$

$$\iint_S \vec{J} \cdot d\vec{S} = I \Rightarrow H = \frac{I}{2\pi r}$$

$$(2) H = \int dH = \int \frac{I}{4\pi} \frac{dx \sin \theta}{r^2} = \frac{I}{2\pi} \int_{\phi_B}^{\phi_A} \frac{\cos \phi \cdot a \, d\phi}{\cos^2 \phi \cdot (\frac{a}{\cos \phi})^2}$$

$$= \frac{I}{4\pi a} \int_{\phi_B}^{\phi_A} \cos \phi \, d\phi$$

$$= \frac{I}{4\pi a} (\sin \phi_A - \sin \phi_B) = \frac{I}{4\pi a} \left(\frac{a}{\sqrt{a^2 + l_1^2}} - \frac{a}{\sqrt{a^2 + l_2^2}} \right) = \frac{I}{2\pi a} \left(\frac{1}{\sqrt{a^2 + l_1^2}} - \frac{1}{\sqrt{a^2 + l_2^2}} \right)$$



$$\sin \theta = \cos \phi = \frac{a}{r}$$

$$\therefore \frac{dx}{dt} = \pm \sqrt{\frac{2}{x}} \Rightarrow \int \sqrt{x} dx = \pm \sqrt{2} \int dt \Rightarrow x = \pm \left(\frac{3\sqrt{2}}{2} t + C_1 \right)^{\frac{2}{3}}$$

$$x = \pm \left(\frac{3\sqrt{2}}{2} t + 1 \right)^{\frac{2}{3}} \stackrel{d\dot{x}|_0 = \sqrt{2}}{=} \sqrt{2} \left(\frac{3\sqrt{2}}{2} t + 1 \right)^{\frac{2}{3}}$$

[PQ3]

$$\frac{dx}{dt} = \frac{3\sqrt{2}}{2} \cdot \frac{2}{3} \left(\frac{3\sqrt{2}}{2} t + 1 \right)^{-\frac{1}{3}} = \sqrt{2} \left(\frac{3\sqrt{2}}{2} t + 1 \right)^{-\frac{1}{3}}$$

$$\frac{d^2x}{dt^2} = \sqrt{2} \cdot \frac{3\sqrt{2}}{2} \left(-\frac{1}{3} \right) \left(\frac{3\sqrt{2}}{2} t + 1 \right)^{-\frac{4}{3}} = -\left(\frac{3\sqrt{2}}{2} t + 1 \right)^{-\frac{4}{3}}$$

$$\frac{d^3x}{dt^3} = (-1) \frac{3\sqrt{2}}{2} \left(-\frac{4}{3} \right) \left(\frac{3\sqrt{2}}{2} t + 1 \right)^{-\frac{7}{3}} = 2\sqrt{2} \left(\frac{3\sqrt{2}}{2} t + 1 \right)^{-\frac{7}{3}}$$

$$\therefore x = 1 + \sqrt{2}t + \frac{1}{2}t^2 + \frac{2\sqrt{2}}{6}t^3 + O(t^3)$$

$$= 1 + \sqrt{2}t - \frac{1}{2}t^2 + \frac{\sqrt{2}}{3}t^3 + O(t^3)$$

[PQ4]

$$(a) \frac{\partial}{\partial x} A_y - \frac{\partial}{\partial y} A_x = -\frac{1}{y} \quad (b) \frac{\partial}{\partial x} A_y - \frac{\partial}{\partial y} A_x = 0$$

[PQ5]

$$\oint_C \vec{A}(x, y') d\vec{r}' = \iint \left(\frac{\partial A_y}{\partial x} - \frac{\partial A_x}{\partial y} \right) dx dy$$

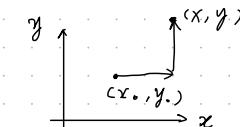
(a) $\oint_C \vec{A}(x, y') d\vec{r}' \neq 0$, depend on trace

(b) $\oint_C \vec{A}(x, y') d\vec{r}' = 0$, independent on trace

[PQ6]

$$S(x, y) = \int_{(x_0, y_0)}^{(x, y_0)} \frac{1}{x} dx + \int_{(x, y_0)}^{(x, y)} -\frac{1}{y} dy$$

$$= \ln \frac{x}{x_0} + \left(-\ln \frac{y}{y_0} \right) = \ln \frac{y_0}{x_0} \frac{x}{y}$$



[PQ7]

[PQ1]

$$\int_{-\infty}^{\infty} dx e^{ikx} \frac{\partial}{\partial x} f(x, t) = f(x, t) e^{ikx} \Big|_{-\infty}^{+\infty} - \int_{-\infty}^{+\infty} ik e^{ikx} f(x, t) = -ik \hat{f}(k)$$

$$\therefore \left(\frac{\partial^2}{\partial x^2} - \frac{\partial}{\partial t} \right) G(x, t) \rightarrow (-k^2 - i\omega) \hat{G}(k, \omega) \quad ; \quad \hat{G}(k, \omega) = \frac{1}{k^2 + i\omega}$$

$$-i\delta(x) \delta(t) \rightarrow -1$$

$$[PQ2] G(k, t) = \frac{1}{2\pi} \int_{-\infty}^{+\infty} dw e^{itw} \frac{1}{k^2 + iw} = \frac{1}{2\pi} \cdot 2\pi i \cdot \text{Res} \left[\frac{e^{itw}}{k^2 + iw}, w = ik^2 \right]$$

$$= i \lim_{w \rightarrow ik^2} (w - ik^2) \frac{e^{itw}}{k^2 + iw} = e^{-k^2 t}$$

数学

第1問

[PQ1]

$$\frac{dE}{dt} = \frac{dx(t)}{dt} \cdot \frac{d^2x(t)}{dt^2} + \frac{1}{x(t)^2} \frac{d^2x(t)}{dt^2} = \frac{dx}{dt} \left(\frac{d^2x(t)}{dt^2} + \frac{1}{x(t)^2} \right) = 0$$

$$[PQ2] \frac{1}{2} \left(\frac{dx}{dt} \right)^2 - \frac{1}{x} = C, \text{ when } t=0, \frac{1}{2} \cdot 2 - 1 = 0 = C$$

$$[PQ3] G(x,t) = \frac{1}{2\pi} \int_{-\infty}^{+\infty} dk e^{-ixk} e^{-k^2 t} = \frac{1}{2\pi} \int_{-\infty}^{+\infty} dk e^{-t(k + \frac{ix}{2t})^2 + t(\frac{ix}{2t})^2} \\ = \frac{1}{2\pi} e^{-\frac{x^2}{4t}} \sqrt{\frac{\pi}{t}}$$

第4問

[PQ1]

$$\alpha_i \alpha_j + \alpha_j \alpha_i = \begin{pmatrix} 0 & \sigma_i \\ \sigma_i & 0 \end{pmatrix} \begin{pmatrix} 0 & \sigma_j \\ \sigma_j & 0 \end{pmatrix} + \begin{pmatrix} 0 & \sigma_i \\ \sigma_i & 0 \end{pmatrix} \begin{pmatrix} 0 & \sigma_j \\ \sigma_j & 0 \end{pmatrix} = \begin{pmatrix} 0 & \sigma_i \\ 0 & \sigma_i \sigma_j \end{pmatrix} + \begin{pmatrix} \sigma_i \sigma_j & 0 \\ 0 & 0 \end{pmatrix} = \begin{pmatrix} \sigma_i \sigma_j + \sigma_j \sigma_i & 0 \\ 0 & \sigma_i \sigma_j + \sigma_j \sigma_i \end{pmatrix} \\ = \begin{pmatrix} 2\delta_{ij} I & 0 \\ 0 & 2\delta_{ij} I \end{pmatrix} = 2\delta_{ij} \mathbb{I}_{4 \times 4} \\ \alpha_i \beta + \beta \alpha_i = \begin{pmatrix} 0 & \sigma_i \\ \sigma_i & 0 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} + \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} 0 & \sigma_i \\ \sigma_i & 0 \end{pmatrix} = \begin{pmatrix} 0 & -\sigma_i I \\ \sigma_i I & 0 \end{pmatrix} + \begin{pmatrix} 0 & \sigma_i I \\ -\sigma_i I & 0 \end{pmatrix} = 0_{4 \times 4} \\ \beta^2 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} = \mathbb{I}_{4 \times 4}$$

[PQ2]

$$A^2 = \sum_i A_i A_i = \sum_i (\alpha_i n_i \sin\theta + \beta \cos\theta)(\alpha_i n_i \sin\theta + \beta \cos\theta) \\ = \sum_i \left[(\alpha_i^2 \sin^2\theta + \sin\theta \cos\theta (\alpha_i \beta + \beta \alpha_i)) + \beta^2 \cos^2\theta \right] n_i^2 \\ = \sum_i \mathbb{I} n_i^2 \\ = \mathbb{I}$$

[PQ3]

$$|A - \lambda \mathbb{I}_{4 \times 4}| = \left| \sin\theta \begin{pmatrix} 0 & \sigma \\ \sigma & 0 \end{pmatrix} n + \cos\theta \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} - \begin{pmatrix} \lambda I & 0 \\ 0 & \lambda I \end{pmatrix} \right| \\ = \begin{vmatrix} (\cos\theta - \lambda)I & \sin\theta \sigma n \\ \sin\theta \sigma n & (-\cos\theta - \lambda)I \end{vmatrix} = ?$$

SOKENDAI KEK 入試 2017 解答

物理Ⅲ

第1問

第2問

[PQ1]

$$[PQ2] F = \sum_{\nu} P_{\nu} E_{\nu} + T \sum_{\nu} P_{\nu} \ln P_{\nu} + \lambda \left(\sum_{\nu} P_{\nu} - 1 \right) = \sum_{\nu} (E_{\nu} + T \ln P_{\nu} + \lambda) P_{\nu} - 1$$

$$\Delta F = \sum_{\nu} (E_{\nu} + T \ln P_{\nu} + \lambda) \Delta P_{\nu} = 0$$

$$\because E_{\nu} + T \ln P_{\nu} + \lambda = 0 \Rightarrow P_{\nu} = e^{\frac{-E_{\nu}}{T}} e^{-E_{\nu}/T}$$

$$\text{Since } \sum_{\mu} P_{\mu} = e^{\frac{-E_1}{T}} \sum_{\nu} e^{\frac{-E_{\nu}}{T}} = 1 \quad \therefore e^{\frac{-E_1}{T}} = \frac{1}{\sum_{\nu} e^{-E_{\nu}/T}}$$

$$\therefore P_{\nu} = \frac{e^{-E_{\nu}/T}}{\sum_{\mu} e^{-E_{\mu}/T}}$$

[PQ3]

$$Z = \sum_{\{\sigma_i\}} e^{-\frac{H}{kT}} = \sum_{\{\sigma_i\}} e^{\sum_{i,j} \sigma_i \sigma_j / kT}, \quad Z^d = e^{-dJN/kT}$$

$$F_0^d = -kT \ln Z^d = -kT (\partial JN/kT) = -dJN$$

[PQ4] [PQ5]

第3問

$$[PQ1] (1) \quad \frac{\partial U}{\partial x} = U_0 \{ -k \sin(kx + \theta_0) \} = 0, \text{ when } x=0 \Rightarrow \theta_0 = 0, \pi$$

$$(2) \quad L = \frac{m \dot{x}^2}{2} - U_0 \{ \cos(kx + \theta_0) - \cos\theta_0 \}$$

$$\frac{\partial L}{\partial x} = k U_0 \sin(kx + \theta_0), \quad \frac{\partial L}{\partial \dot{x}} = m \dot{x} \Rightarrow m \ddot{x} = k U_0 \sin(kx + \theta_0)$$

$$(3) \quad \theta_0 = 0$$

$$m \ddot{x} \approx k U_0 (kx) = k^2 U_0 x \quad ?$$

[PQ2]

第4問

$$[PQ1] \quad I - C' \Delta x \frac{d(V + \Delta V)}{dt} - (I + \Delta I) = 0, \quad V - (-L' \Delta x \frac{dI}{dt}) - (V + \Delta V) = 0$$

[P8] 2

$$\frac{dI}{dx} = \lim_{\Delta x \rightarrow 0} \frac{(I + \Delta I) - I}{\Delta x} = -C' \frac{d(V + \Delta V)}{dt} = -C' \frac{dV}{dt}$$

$$\frac{dV}{dx} = \lim_{\Delta x \rightarrow 0} \frac{(V + \Delta V) - V}{\Delta x} = -L' \frac{dI}{dt}$$

[P8] 3

$$\frac{d^2 I}{dx^2} = -C' \frac{d}{dt} \frac{dV}{dx} = L'C' \frac{d^2 I}{dt^2} \Rightarrow C = \frac{1}{L'C'}$$

$$[P8] \int_s \vec{E} \cdot d\vec{s} = \frac{Q}{\epsilon_0} \Rightarrow 2\pi r |\vec{E}| = \frac{Q}{\epsilon_0} \Rightarrow |\vec{E}| = \frac{Q}{2\pi r \epsilon_0}$$

$$[P8] V_{ab} = \int_a^b |E| dr = \frac{\mu_0}{2\pi \epsilon_0} \int_a^b \frac{1}{r} dr = \frac{Q}{2\pi \epsilon_0} \ln \frac{b}{a}$$

$$C' = \frac{Q}{V_{ab}} = \frac{2\pi \epsilon_0}{\ln \frac{b}{a}}$$

$$[P8] \frac{1}{\mu_0} \oint_C \vec{B} \cdot d\vec{s} = I \Rightarrow \frac{1}{\mu_0} 2\pi r |\vec{B}| = I \Rightarrow |\vec{B}| = \frac{\mu_0 I}{2\pi r}$$

$$[P8] U_B = \frac{1}{2\mu_0} \int_V |\vec{B}|^2 dV = \frac{1}{2\mu_0} \left(\frac{\mu_0 I}{2\pi r} \right)^2 \iiint \frac{1}{r^2} r \sin \theta dr d\theta dz \\ = \frac{\mu_0 I^2}{4\pi} \ln \frac{b}{a} \Rightarrow L' = \frac{2U_B}{I^2} = \frac{\mu_0}{2\pi} \ln \frac{b}{a}$$

$$[P8] \frac{1}{L'C'} = \frac{1}{\sqrt{2\pi \mu_0}}$$

数学

第1問

$$[P8] 1 T_1^2 = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 0 \end{pmatrix}, T_1^3 = \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \quad T_1^n = \begin{cases} \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 0 \end{pmatrix}, & n=2,4,6,\dots \\ \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}, & n=1,3,5,\dots \end{cases}, \quad n=2,4,6,\dots$$

[P8] 2

$$\begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \begin{pmatrix} a_{11} & a_{12} & a_{13} \\ a_{21} & a_{22} & a_{23} \\ a_{31} & a_{32} & a_{33} \end{pmatrix} - \begin{pmatrix} a_{11} & a_{12} & a_{13} \\ a_{21} & a_{22} & a_{23} \\ a_{31} & a_{32} & a_{33} \end{pmatrix} \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

$$= \begin{pmatrix} a_{21} & a_{22} & a_{23} \\ a_{11} & a_{12} & a_{13} \\ 0 & 0 & 0 \end{pmatrix} - \begin{pmatrix} a_{12} & a_{11} & 0 \\ a_{22} & a_{21} & 0 \\ a_{32} & a_{31} & 0 \end{pmatrix} = \begin{pmatrix} a_{21}-a_{12} & a_{22}-a_{11} & a_{23} \\ a_{11}-a_{21} & a_{12}-a_{22} & a_{13} \\ -a_{32} & -a_{31} & 0 \end{pmatrix} = 0$$

{T8} is changeable.

[P8] 2

$$e^{ix} = I + A(ix) + \frac{A^2}{2}(ix)^2$$

第2回

[P8] 1

$$\lambda^2 - 4 = 0 \quad \lambda = \pm 2 \Rightarrow f_1(x) = e^{2x}, f_2 = e^{-2x}$$

[P8] 2

$$\frac{dC_1}{dx} e^{2x} + \frac{dC_2}{dx} e^{-2x} = 0$$

$$f(x) = C_1(x) e^{2x} + C_2(x) e^{-2x}$$

$$\frac{df(x)}{dx} = e^{2x} \frac{dC_1}{dx} + 2C_1 e^{2x} + e^{-2x} \frac{dC_2}{dx} - 2C_2 e^{-2x} = 2C_1 e^{2x} - 2C_2 e^{-2x}$$

$$\frac{d^2 f(x)}{dx^2} = 2 \left[e^{2x} \frac{dC_1}{dx} + 2C_1 e^{2x} - e^{-2x} \frac{dC_2}{dx} + 2C_2 e^{-2x} \right]$$

$$\Rightarrow 2 \left[e^{2x} \left(\frac{dC_1}{dx} + 2C_1 \right) + e^{-2x} \left(-\frac{dC_2}{dx} + 2C_2 \right) \right] - 4C_1 e^{2x} - 4C_2 e^{-2x}$$

$$= 2 \left(e^{2x} \frac{dC_1}{dx} - e^{-2x} \frac{dC_2}{dx} \right) = \frac{8}{e^{2x} + e^{-2x}} \Rightarrow (e^{2x} \frac{dC_1}{dx} - e^{-2x} \frac{dC_2}{dx})(e^{2x} + e^{-2x}) = 4$$

$$\textcircled{1} \quad 2e^{2x} \frac{dC_1}{dx} (e^{2x} + e^{-2x}) = 4 \Rightarrow \frac{dC_1}{dx} (e^{4x} + 1) = 2$$

$$\textcircled{2} \quad -2e^{-2x} \frac{dC_2}{dx} (e^{2x} + e^{-2x}) = 4 \Rightarrow \frac{dC_2}{dx} (1 + e^{-4x}) = -2$$

[P8] 3

$$C_1 = \int \frac{1}{e^{4x} + 1} dx$$

第3回

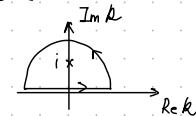
[P8] 1

$$F(k) = \frac{1}{2\pi} \int_{-\infty}^{\infty} dx e^{-ikx} e^{-ix} = \frac{1}{2\pi} \left[\int_0^{+\infty} dx e^{-ikx} e^{-ix} + \int_{-\infty}^0 dx e^{-ikx} e^{-ix} \right] \\ = \frac{1}{2\pi} \left[\frac{1}{1-i k} e^{(1-i k)x} \Big|_0^\infty + \frac{1}{1-i k} e^{(1-i k)x} \Big|_{-\infty}^0 \right] \\ = \frac{1}{2\pi} \left[\frac{-1}{1-i k} + \frac{1}{1-i k} \right] \\ = \frac{1}{\pi(1+k^2)}$$

[P8] 2

$$f(x) = \int_{-\infty}^{\infty} \frac{i}{\pi(i+k)} e^{ikx} dk = \frac{1}{2\pi} \int_{-\infty}^{+\infty} \left(\frac{1}{1+ik} + \frac{1}{1-ik} \right) e^{ikx} dk$$

\textcircled{1} $x > 0$



$$\int_{-\infty}^{\infty} \frac{e^{ikx}}{i+ik} dk = 2\pi i \operatorname{Res} \left[\frac{e^{ikx}}{1+ik}, k=i \right] \\ = 2\pi i \lim_{k \rightarrow i} \frac{e^{ikx}}{i+ik} (k-i) \\ = 2\pi e^{-x}$$

$$\textcircled{2} \quad x < 0 \quad : \quad \int_{-\infty}^{+\infty} \frac{e^{ikx}}{1-ik} dk = 2\pi i \operatorname{Res}\left[\frac{e^{ikx}}{1-ik}, k=-i\right] = 2\pi e^x$$

$$\therefore f(x) = \frac{1}{2\pi} \cdot 2\pi e^{-|x|} = e^{-|x|}$$

第4問

$$[\text{P4-1}] \quad (x-1)^2 + \sin^2 t = 1 \Rightarrow x = \pm \cos t + 1$$

$$x = \cos t + 1$$

$$z(t) = 0$$

〔P4-2〕

$$\begin{aligned} \oint_C d\vec{r} \cdot \vec{A} &= \oint_C y dx = \int_2^0 \sqrt{1-(x-1)^2} dx - \int_0^2 \sqrt{1-(x-1)^2} dx = 2 \int_2^0 \sqrt{1-(x-1)^2} dx \\ &= 2 \int_1^2 \sqrt{1-x^2} dx \quad \text{let } x = \sin \theta \quad dx = \cos \theta d\theta \\ &= 2 \int_{\frac{\pi}{2}}^{\frac{\pi}{2}} (\cos \theta)^2 d\theta = 2 \int_{\frac{\pi}{2}}^{\frac{\pi}{2}} \frac{\cos 2\theta + 1}{2} d\theta = -2\pi \end{aligned}$$

$$[\text{P4-3}] \quad \nabla \times \vec{A} = \begin{pmatrix} \hat{e}_x & \hat{e}_y & \hat{e}_z \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ y & z & x \end{pmatrix} = -\hat{e}_x - \hat{e}_y - \hat{e}_z$$

$$\frac{\partial y}{\partial \theta} \frac{\partial z}{\partial \varphi} - \frac{\partial z}{\partial \theta} \frac{\partial y}{\partial \varphi} = \sin \theta \sin \varphi \cos \psi = \sin \theta \cos \psi$$

$$\frac{\partial z}{\partial \theta} \frac{\partial x}{\partial \varphi} - \frac{\partial x}{\partial \theta} \frac{\partial z}{\partial \varphi} = -\sin \theta \sin \varphi (-\sin \psi) = \sin^2 \theta \sin \psi$$

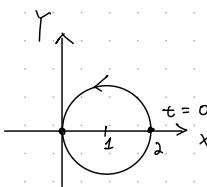
$$\frac{\partial x}{\partial \theta} \frac{\partial y}{\partial \varphi} - \frac{\partial y}{\partial \theta} \frac{\partial x}{\partial \varphi} = \cos \theta \cos \varphi \sin \theta \cos \psi - \cos \theta \sin \varphi \sin \theta (-\sin \psi) = \cos \theta \sin \psi$$

〔P4-4〕

$$\int_S dS \cdot \operatorname{rot} \vec{A} = - \int_B [1 \cdot \cos \varphi (\cos \psi + \sin \psi) + \cos \theta \sin \theta] d\theta d\varphi$$

$$\int_{\frac{\pi}{2}}^{\frac{\pi}{2}} \int_0^{2\pi} (\cos \varphi + \sin \psi) d\varphi d\theta = 0 \quad \Rightarrow \quad \int_S dS \cdot \operatorname{rot} \vec{A} = -2\pi$$

$$\int_0^{\frac{\pi}{2}} \cos \theta \sin \theta d\theta \int_0^{2\pi} d\varphi = 2\pi$$



SOKENDAI KEK 入試 2016 解答

物理

第1問

[問題1]

$$\begin{aligned} \text{(1)} \quad H|\psi_n\rangle &= (H_0 + \lambda H')|n\rangle + \lambda|\psi_n^{(1)}\rangle + \lambda^2|\psi_n^{(2)}\rangle + \dots \\ &= H_0|n\rangle + H_0\lambda|\psi_n^{(1)}\rangle + H'\lambda|n\rangle + \dots \\ &= H_0|n\rangle + \lambda(H_0|\psi_n^{(1)}\rangle + H'|n\rangle) + \dots \end{aligned}$$

$$\begin{aligned} E_n|\psi_n\rangle &= (\sum_n + \lambda E_n^{(1)} + \dots)(|n\rangle + \lambda|\psi_n^{(1)}\rangle + \dots) \\ &= \sum_n |n\rangle + \lambda(E_n^{(1)}|n\rangle + \sum_n |\psi_n^{(1)}\rangle + \dots) \end{aligned}$$

multiply each with $\langle k|$

$$\begin{aligned} \langle k|H_0|n\rangle + \lambda(\langle k|H'|n\rangle + \langle k|H_0|\psi_n^{(1)}\rangle) + \dots &= \langle k|E_n^{(1)}|n\rangle + \lambda(\langle k|E_n^{(1)}|n\rangle + \\ &\quad \langle k|E_n|\psi_n^{(1)}\rangle) \end{aligned}$$

$$\text{for } k=n, \quad \langle n|H_0|1\rangle = \langle n|E_0, \quad \therefore E_n^{(1)} = \langle n|H'|n\rangle \quad \text{--- ①}$$

$$\begin{aligned} \text{(2)} \quad H_0|\psi_n^{(1)}\rangle &= H_0|n\rangle + \sum_{m \neq n} H_0 C_m^{(1)}|m\rangle \quad \langle k|E_n|\psi_n^{(1)}\rangle = E_n \delta_{kn} + \sum_{m \neq n} C_m^{(1)} E_n \delta_{mk} \\ \langle k|H_0|\psi_n^{(1)}\rangle &= C_k^{(1)} E_n \delta_{kn} + \sum_{m \neq n} C_m^{(1)} E_n \delta_{mk} \end{aligned}$$

$$\text{when } m=k, m \neq n, k \neq n, \quad \text{①} \Rightarrow \langle k|H'|n\rangle + C_k^{(1)} E_k = C_k^{(1)} E_n + B_n^{(1)} \delta_{kn}$$

$$\therefore C_m^{(1)} = C_k^{(1)} = \frac{\langle k|H'|n\rangle}{E_n - E_m}$$

$$\langle n|\psi_n^{(1)}|\psi_n\rangle = 1$$

$$\Rightarrow (\langle n| + \lambda(\langle n|\psi_n^{(1)}\rangle + \lambda^2 \langle n|\psi_n^{(2)}\rangle + \dots))(|n\rangle + \lambda|\psi_n^{(1)}\rangle + \lambda^2|\psi_n^{(2)}\rangle + \dots) = 1$$

$$\Rightarrow \langle n|n\rangle + \lambda(\langle n|\psi_n^{(1)}\rangle + \langle n|\psi_n^{(2)}\rangle) + \dots = 1$$

$$\therefore \langle n|\psi_n^{(1)}\rangle + \langle n|\psi_n^{(2)}\rangle = 0$$

$$\Rightarrow \langle n|(C_n^{(1)}|n\rangle + \sum_{m \neq n} C_m^{(1)}|m\rangle) + (\langle n|C_n^{(1)} + \sum_{m \neq n} \langle m|C_n^{(1)}\rangle)|n\rangle = 0 \Rightarrow C_n^{(1)} = 0 \quad \text{②}$$

$$\text{(4)} \quad O|\psi_n\rangle = O|n\rangle + \lambda O|\psi_n^{(1)}\rangle = O|n\rangle + \lambda C_n^{(1)} O|n\rangle + \lambda \sum_{m \neq n} C_m^{(1)} O|m\rangle$$

$$\begin{aligned} \langle \psi_n^{(1)}|O|\psi_n\rangle &= (\langle n| + \lambda \langle n|C_n^{(1)} + \sum_{m \neq n} \langle m|C_n^{(1)}\rangle)(O|n\rangle + \lambda C_n^{(1)} O|n\rangle + \lambda \sum_{m \neq n} C_m^{(1)} O|m\rangle) \\ &= \langle n|O|n\rangle + \lambda \langle n|C_n^{(1)} O|n\rangle + \lambda \sum_{m \neq n} \langle m|C_n^{(1)} O|m\rangle + \lambda \langle n|C_n^{(1)} O|n\rangle + \lambda \sum_{m \neq n} \langle m|C_n^{(1)} O|m\rangle \end{aligned}$$

[P0] 2

$$1) \Delta \psi(r, \theta, \phi) = \frac{1}{r^2} \frac{\partial^2}{\partial r^2} \left(r^2 \frac{\partial}{\partial r} A \exp(-\frac{r}{a_0}) \right) = \frac{1}{a_0 r} \left(\frac{r}{a_0} - 2 \right) \Delta \psi(r, \theta, \phi)$$

$$\left[\frac{-h^2}{8m} \left(\frac{1}{a_0 r} \left(\frac{r}{a_0} - 2 \right) \right) - \frac{C}{r} \right] \Delta \psi(r, \theta, \phi) = E \Delta \psi(r, \theta, \phi) \Rightarrow \begin{cases} a_0 = \frac{h^2}{mC} \\ E = -\frac{mC^2}{2h^2} \end{cases}$$

$$(2) \int \Delta \psi \Delta \psi dxdydz = \int_0^R dr \int_0^\pi d\theta \int_0^{2\pi} r^2 \sin^2 \theta e^{-\frac{2r}{a_0}} = 4\pi R^2 \int_0^R r^2 e^{-\frac{2r}{a_0}}$$

$$(2) \text{ (X)} \quad = 8\pi A^2 \left(\frac{a_0}{2}\right)^3 = 1 \Rightarrow A = \frac{\sqrt{a_0^3}}{8\pi}$$

[P0] 2

[P0] 1

$$(1) Z = \frac{1}{h} \iint e^{-\beta \left(\frac{p_x^2}{2m} + \frac{p_y^2}{2} \right)} dp dx = \frac{1}{h} \cdot \frac{\sqrt{2\pi m}}{\beta} \sqrt{\frac{2\pi}{\alpha\beta}} = \frac{2\pi}{\alpha\beta} \sqrt{\frac{m}{\alpha}}$$

$$\int_{-\infty}^{+\infty} dx e^{-\beta^2 x^2} = \frac{1}{2} \pi^{\frac{1}{2}} \alpha^{\frac{3}{2}}$$

$$\langle \frac{p_x^2}{2m} \rangle = \frac{1}{h} \iint \frac{p_x^2}{2m} e^{-\beta \left(\frac{p_x^2}{2m} + \frac{p_y^2}{2} \right)} dp dx = \frac{1}{h} \cdot \frac{1}{2m} \cdot \frac{1}{2} \pi^{\frac{1}{2}} \cdot \left(\frac{2\pi}{\beta} \right)^{\frac{3}{2}} \sqrt{\frac{m}{\alpha\beta}} \Rightarrow \langle \frac{p_x^2}{2m} \rangle = \langle \frac{p_y^2}{2m} \rangle$$

$$\langle \frac{p_z^2}{2} \rangle = \frac{1}{h} \iint \frac{p_z^2}{2} e^{-\beta \left(\frac{p_x^2}{2m} + \frac{p_y^2}{2} \right)} dp dx = \frac{1}{h} \cdot \frac{1}{2} \pi^{\frac{1}{2}} \cdot \left(\frac{2\pi}{\beta} \right)^{\frac{3}{2}} \sqrt{\frac{m}{\alpha\beta}}$$

$$(2) Z = \frac{1}{h} \iint e^{-\beta \sum_{i=1}^3 \left(p_i^2 / (2m) + (x_i^2 + y_i^2 + z_i^2) / (2m) + \alpha(x_i^2 + y_i^2 + z_i^2) / (2m) \right)} dp dx = \left(\frac{2\pi}{\beta} \sqrt{\frac{m}{\alpha}} \right)^{3N}$$

$$\langle E \rangle = \frac{1}{h} \iint H e^{-\beta H} dp dx = -\frac{\partial \ln Z}{\partial \beta} = -3N \frac{1}{2} \cdot \frac{\partial}{\partial \beta} \left(-\frac{1}{2} \right) \Rightarrow C = \frac{\partial \langle E \rangle}{\partial T} = -3NkT$$

[P0] 2

$$(1) Z = \sum_{n_x} \sum_{n_y} \sum_{n_z} e^{-\beta(n_x \hbar \omega_x + n_y \hbar \omega_y + n_z \hbar \omega_z + \frac{3}{2})/\hbar \omega} = e^{-\beta \frac{3}{2} \hbar \omega} \left(\frac{1}{1 - e^{-\beta \hbar \omega}} \right)^3$$

$$\begin{aligned} \langle E \rangle &= -N \frac{\partial \ln Z}{\partial \beta} = -N \left\{ \frac{3}{2} \left(-\frac{3}{2} \hbar \omega \beta - 3 \ln(1 - e^{-\beta \hbar \omega}) \right) \right\} \\ &= \frac{3}{2} \hbar \omega N + 3N \frac{1}{1 - e^{-\beta \hbar \omega}} (\hbar \omega e^{-\beta \hbar \omega}) \\ &= \frac{3}{2} \hbar \omega N + 3N \hbar \omega \frac{1}{e^{\beta \hbar \omega} - 1} \end{aligned}$$

$$(2) \lim_{T \rightarrow \infty} \frac{\partial \langle E \rangle}{\partial T} = \lim_{T \rightarrow \infty} 3N \hbar \omega \left(\frac{-1}{e^{\beta \hbar \omega / (kT)} - 1} \right)^2 \cdot e^{\frac{\hbar \omega}{kT}} \cdot \left(\frac{-\hbar \omega}{kT} \right)^2, \text{ let } \frac{\hbar \omega}{kT} = x \rightarrow 0$$

$$= \lim_{x \rightarrow 0} 3Nk \frac{x^2 e^x}{(e^x - 1)^2} = 3Nk \lim_{x \rightarrow 0} \frac{x^2 e^x + 2x e^x}{2(e^x - 1)e^x} = 3Nk$$

[P0] 3

$$(1) Z = \frac{1}{h^3} \iint e^{-\beta H} dp dx = \frac{L^3}{h^3} \iint e^{-\beta(p_x^2 + p_y^2 + p_z^2) / (2m)} dp = \left(\frac{L}{h} \right)^3 \left(2\pi \hbar \omega \frac{1}{\beta} \right)^{\frac{3}{2}}$$

$$\langle E \rangle = -\frac{\partial \ln Z}{\partial \beta} = -\frac{1}{2} \left(\frac{L}{h} \right)^3 \frac{3}{2} \left(2\pi \hbar \omega \frac{1}{\beta} \right)^{\frac{1}{2}} \cdot 2\pi \hbar \omega \cdot (-1) \beta^{-2} = \frac{3}{2} kT$$

$$C = \frac{\partial \langle E \rangle}{\partial T} = \frac{3}{2} k$$

$$(2) D(\epsilon) d\epsilon = \frac{L^3}{h^3} 4\pi p^2 dp = \frac{4\pi L^3}{Q^2 h^3} (2\pi \hbar \omega)^{\frac{1}{2}} m d\epsilon$$

$$N = 2 \int_0^M D(\epsilon) d\epsilon = \frac{L^3}{\pi h^3} 2^{\frac{1}{2}} m^{\frac{3}{2}} \frac{2}{3} M^{\frac{3}{2}} \Rightarrow M = \frac{h^2}{2m} (3\pi^2 N_e)^{\frac{2}{3}}$$

(3) (X)

第3問

[P0] 1

$$(1) L = \frac{1}{2} m \dot{x}^2 - \frac{1}{2} k x^2$$

$$(2) \frac{dL}{dx} = -kx, \frac{dL}{dx} = m \dot{x}, \Rightarrow \ddot{x} + \omega^2 x = 0, \omega_0 = \sqrt{k/m}$$

$$\text{Solution: } x = A \cos \omega_0 t + A_2 \sin \omega_0 t = A \cos(\omega_0 t + \alpha)$$

$$(3) \frac{dE}{dt} = m \dot{x} \ddot{x} + k x \dot{x} = 0$$

[P0] 2

$$(1) \text{ Assume } x = a_1 e^{i\omega t} + a_2 e^{i\omega t}, \text{ take into } \ddot{x} + 2\lambda \dot{x} + \omega^2 x = 0,$$

$$(-a_1 i \omega^2 e^{i\omega t} - a_2 i \omega^2 e^{i\omega t}) + 2\lambda(i a_1 e^{i\omega t} + i a_2 e^{i\omega t}) + \omega^2(a_1 e^{i\omega t} + a_2 e^{i\omega t}) = 0$$

$$\Rightarrow a_1(-\omega_0^2 + 2i\lambda\omega_0 + \omega_0^2) e^{i\omega t} + a_2(-\omega_0^2 + 2i\lambda\omega_0 + \omega_0^2) e^{i\omega t} = 0$$

$$\text{consider } r^2 - 2i\lambda r - \omega_0^2 = 0 \Rightarrow r = \frac{2i\lambda \pm \sqrt{-4\lambda^2 + 4\omega_0^2}}{2} = i\lambda \pm \omega_0$$

$$\therefore x = a_1 e^{-(i\lambda + \omega_0)t} + a_2 e^{(i\lambda - \omega_0)t}$$

$$= e^{-\lambda t} (a_1 e^{i\omega_0 t} + a_2 e^{-i\omega_0 t})$$

$$= A e^{-\lambda t} \cos(\omega_0 t + \alpha)$$

$$(2) \bar{U}(t) = \frac{1}{t} \int_t^{t+\alpha} \frac{m \omega_0^2 A^2 e^{-2\lambda t'} \cos^2(\omega_0 t' + \alpha)}{2} dt' = \frac{1}{2} m \omega_0^2 A^2 e^{-2\lambda t}$$

$$\int_t^{t+\alpha} e^{-2\lambda t'} \cos^2(\omega_0 t' + \alpha) dt' \approx e^{-2\lambda t} \int_t^{t+\alpha} \frac{1 + \cos(2\omega_0 t' + 2\alpha)}{2} dt' = \frac{1}{2} e^{-2\lambda t} \frac{t+\alpha-t}{2} = \frac{1}{2} e^{-2\lambda t} \frac{\alpha}{2}$$

$$(3) \dot{x} = -A e^{-\lambda t} \{ \lambda \cos(\omega_0 t + \alpha) + \omega_0 \sin(\omega_0 t + \alpha) \}$$

$$\bar{T} = \frac{1}{t} \int_t^{t+\alpha} \frac{m \omega_0^2 e^{-2\lambda t'} \{ \lambda \cos(\omega_0 t' + \alpha) + \omega_0 \sin(\omega_0 t' + \alpha) \}^2}{2} dt' = \frac{1}{2} m \omega_0^2 A^2 e^{-2\lambda t}$$

$$(4) \quad \frac{d}{dt} [U + \bar{T}] = -\lambda m w_0^2 A^2 e^{-2\lambda t}$$

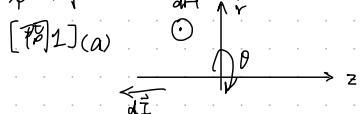
$$\bar{D}(t) = \frac{1}{t} \int_t^{t+\tau_0} \frac{\alpha x^2}{z} = \lambda m w_0^2 A^2 e^{-2\lambda t}$$

$$(5) \quad f_T = -\alpha \dot{x} = -\frac{\partial D}{\partial x} \Rightarrow D = \frac{1}{2} \alpha x^2 = \frac{1}{2} \alpha \frac{P^2}{m^2}$$

$$\frac{dE}{dt} = \frac{P}{m} \dot{p} + \frac{P}{m} kx = \left\{ \left(\frac{\partial L}{\partial x} + f_T \right) + kx \right\} \frac{P}{m} = f_T \frac{P}{m} = -\frac{\partial D}{\partial x} \frac{P}{m} = -\alpha \dot{x} \frac{P}{m}$$

$$= -\alpha \left(\frac{P}{m} \right)^2 = -2D$$

第4問



(b)

[問2]

$$(a) \oint_C \vec{E} \cdot d\vec{s} = \int_0^R (E_r - E_r) dr + \int_{z_1}^{z_2} (E_{z1} - E_{z2}) dz = 0$$

$$(b) E_r = \frac{\lambda}{2\pi r \epsilon_0} \quad H_0(r) = \frac{\partial V}{\partial r} \quad , \quad \frac{E_r}{H_0} = \frac{1}{\epsilon_0} V$$

$$(c) \lim_{r \rightarrow \infty} \frac{E_r}{H_0} = \sqrt{\frac{\mu_0}{\epsilon_0}} = C M_0 = (2\pi r) M N / (s \cdot A^2)$$

数学

[問1]

$$\vec{A}\vec{x} = \lambda \vec{x}$$

$$\frac{1}{3} \begin{pmatrix} 2\alpha + \beta & \alpha - \beta \\ 2\alpha - 2\beta & \alpha + 2\beta \end{pmatrix} \begin{pmatrix} 1 \\ 1 \end{pmatrix} = \alpha \begin{pmatrix} 1 \\ 1 \end{pmatrix} \Rightarrow \text{eigenvalue } \alpha$$

$$\frac{1}{3} \begin{pmatrix} 2\alpha + \beta & \alpha - \beta \\ 2\alpha - 2\beta & \alpha + 2\beta \end{pmatrix} \begin{pmatrix} -1 \\ 2 \end{pmatrix} = \beta \begin{pmatrix} -1 \\ 2 \end{pmatrix} \Rightarrow \text{eigenvalue } \beta$$

[問3]

$$P = \begin{pmatrix} 1 & -1 \\ 1 & 2 \end{pmatrix} \quad P^{-1} = \frac{1}{3} \begin{pmatrix} 2 & 1 \\ -1 & 1 \end{pmatrix} \quad A^n = P \begin{pmatrix} \alpha^n & 0 \\ 0 & \beta^n \end{pmatrix} P^{-1}$$

$$A^n \begin{pmatrix} -3 \\ 18 \end{pmatrix} = \frac{1}{3} \begin{pmatrix} 1 & -1 \\ 1 & 2 \end{pmatrix} \begin{pmatrix} \alpha^n & 0 \\ 0 & \beta^n \end{pmatrix} \begin{pmatrix} 2 & -1 \\ -1 & 1 \end{pmatrix} \begin{pmatrix} -3 \\ 18 \end{pmatrix} = \begin{pmatrix} 4\alpha^n - 7\beta^n \\ 4\alpha^n + 14\beta^n \end{pmatrix} = \begin{pmatrix} a_n - b_n \\ a_n + 2b_n \end{pmatrix}$$

$$\Rightarrow \begin{cases} a_n = 4\alpha^n \\ b_n = 7\beta^n \end{cases}$$

第2問

[問1]

$$I_{n+1} = \int_0^\infty dx e^{-x} x^{n+1} = -e^{-x} x^{n+1} \Big|_0^\infty - \int_0^\infty -e^{-x} (n+1) x^n = (n+1) I_n$$

[問2]

$$\text{let } x^{\frac{1}{2}} = t \quad dx = 2t dt$$

$$\int_0^\infty e^{-t^2} 2t^2 dt = 2 \frac{1}{2} \frac{1}{2} \sqrt{\pi} = \frac{\sqrt{\pi}}{2}$$

$$\int_{-\infty}^\infty x^2 e^{-\alpha x^2} = \frac{1}{2} \sqrt{\pi} \alpha^{-\frac{3}{2}}$$

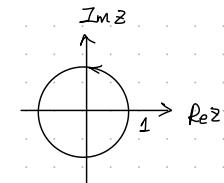
第3問

$$[問1] \frac{df(z)}{dz} = -2\pi \sin(2\pi z) \quad \frac{d^2 f(z)}{dz^2} = -4\pi^2 \cos(2\pi z) \quad \frac{d^3 f(z)}{dz^3} = 8\pi^3 \sin(2\pi z)$$

$$f(z) = -1 + \frac{1}{2} \alpha \pi^2 (z - \frac{1}{2})^2 = -1 + 2\pi^2 (z - \frac{1}{2})^2$$

[問2]

$$\oint_{|z|=1} dz \frac{\cos 2\pi z}{(z - \frac{1}{2})^4} = 2\pi i \frac{1}{3!} \left[\frac{d^3}{dz^3} \cos 2\pi z \right]_{z=\frac{1}{2}} = 0$$



第4問

[問1]

$$z(y(x), g(x)) = y(x) \cdot g(x), \quad g(x) = \frac{1}{x}, \quad y = xz$$

$$\begin{aligned} \frac{dy}{dx} &= \frac{\partial z}{\partial y} \frac{dy}{dx} + \frac{\partial z}{\partial g} \frac{dg}{dx} = \frac{1}{x} \frac{dy}{dx} + y(x) \left(-\frac{1}{x^2} \right) = \frac{1}{x} \cdot \frac{y^2 - x^2}{2xz} + y \left(-\frac{1}{x^2} \right) \\ &= \frac{1}{x^2} \left[\frac{x^2 - x^2}{2xz} - xz \right] = \frac{z^2 - 1}{2xz} - \frac{z}{x} = \frac{1}{x} \left[\frac{z^2 - 1}{z^2} - z \right] \end{aligned}$$

[問2]

$$\frac{dz}{z^2 - 1} = \frac{1}{x} dx \Rightarrow - \int \frac{2z dz}{z^2 + 1} = \int \frac{1}{x} dx \Rightarrow -\ln(z^2 + 1) = \ln x + C$$

$$(z^2 + 1)^{-1} = Cx \Rightarrow y = \pm \sqrt{Cx - x^2}$$

keyring

2024.5.20.0