

# An Introduction to Radio Interferometry

2-2 Diffraction pattern of a single slit



You can find relevant material  
on my personal webpage

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Radio Interferometer is an optical device

1. Light is EM-wave.
2. **EM-wave at long-distance limit: plane wave  $E = E_0 \cos(kx - \omega t + \phi_0)$**   
 $\phi_0$ : phase offset;  $\omega = 2\pi\nu$ : angular frequency;  $\nu$ : frequency
3. The EM-wave propagation can be described by Huygens' principle  
(i) every point on a wave front can be regarded as a new source  
(ii) the wave front is spherically symmetric with respect to the source
4. The solutions of EM-wave equation satisfy the principle of superposition.
5. In vacuum, the speed of EM wave propagation is  $c = \frac{1}{\sqrt{\epsilon_0 \mu_0}}$ ; in the medium that the index of refraction is  $n$ , the speed of EM-wave propagation is  $c/n$ , and the wavelength becomes  $\lambda_0/n$ , where  $\lambda_0$  is the wavelength in vacuum.

1. EM-wave at long-distance limit: plane wave  $E = E_0 \cos(kx - \omega t + \phi_0)$

## Lecture Unit 1-1

1. EM-wave at long-distance limit: plane wave  $E = E_0 \cos(kx - \omega t + \phi_0)$

# Flux density

Energy is proportional to the square of electric or magnetic field strength.

The **energy flux density**, which is the amount of energy passing a unit area in a unit of frequency and time,

Is also proportional to the **square of electric and/or magnetic field strength**.

Electric field:  $\vec{E} = -\nabla\Phi$

Total electric potential energy :  $W_E = \frac{1}{2} \int \rho(\vec{x})\Phi(\vec{x})d^3x = \frac{\epsilon_0}{2} \int (\nabla \cdot \vec{E})\Phi(\vec{x})d^3x$

$$= \frac{\epsilon_0}{2} \int \vec{E} \cdot (-\nabla\Phi(\vec{x}))d^3x = \frac{\epsilon_0}{2} \int |\vec{E}|^2 d^3x$$

From the symmetry of Maxwell's equations we expect magnetic field energy to have a similar form  $W_B = \frac{1}{2\mu_0} \int |\vec{B}|^2 d^3x$

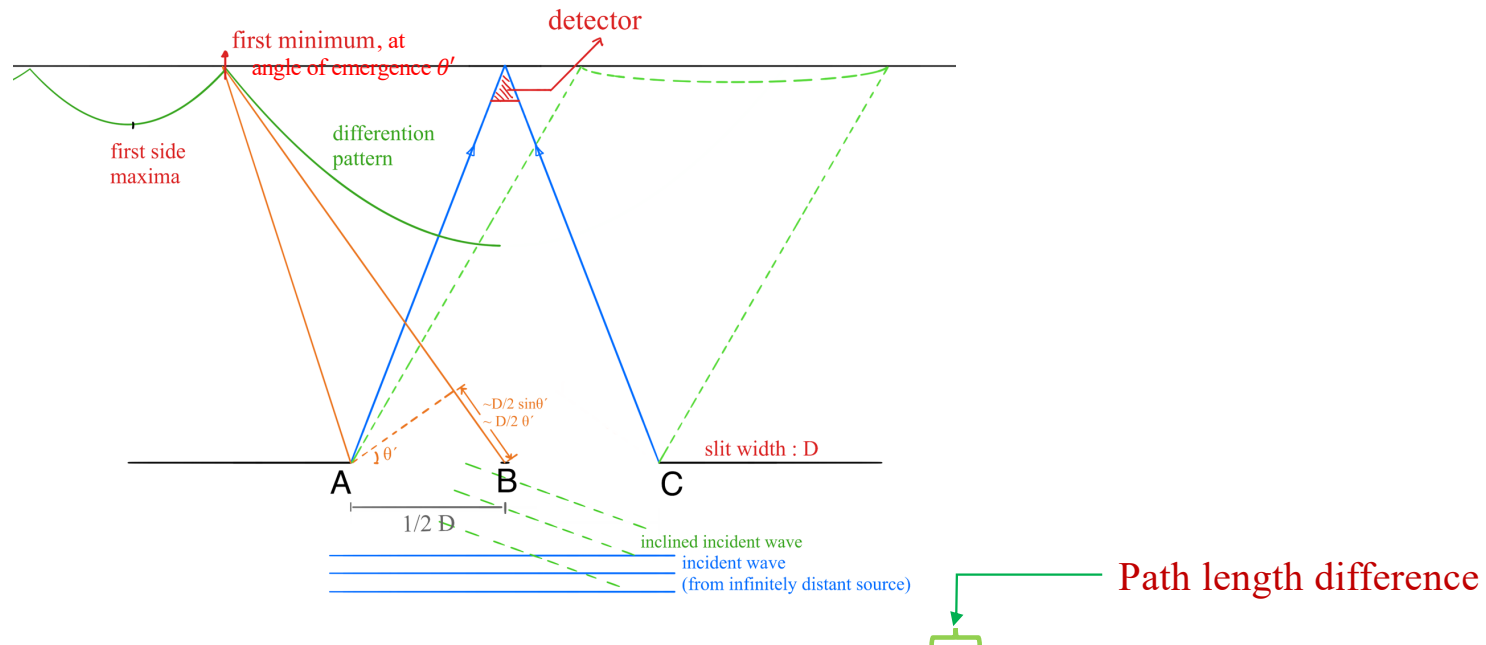
In EM wave

$$\frac{E_m}{B_m} = \frac{1}{\mu_0 \epsilon_0 c}$$

1. EM-wave at long-distance limit: plane wave  $E = E_0 \cos(kx - \omega t + \phi_0)$
2. Energy flux density  $\propto E_0^2$

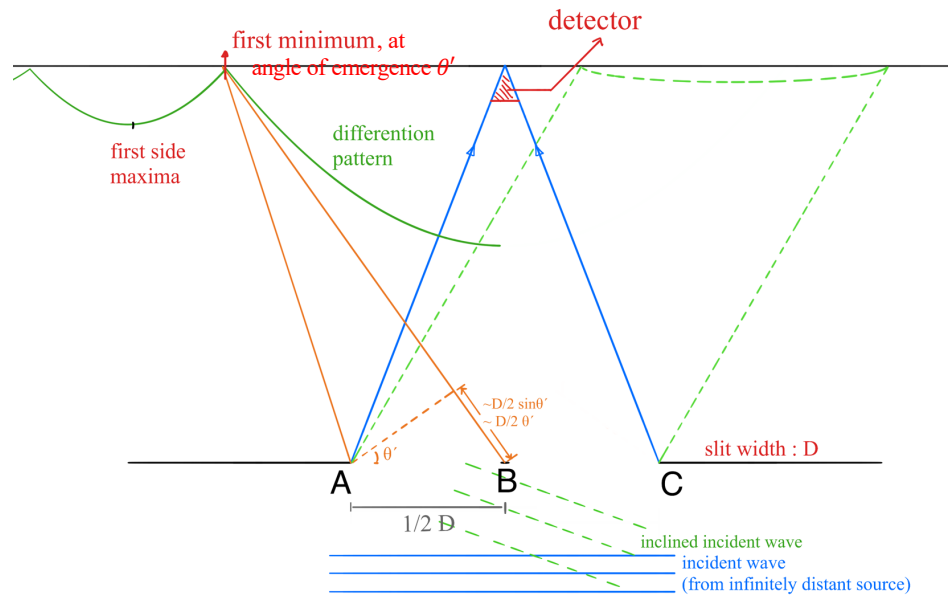
## Lecture Unit 1-2

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$$(E_A + E_B)^2 = E_0^2 [\cos(kx - \omega t) + \cos(k(x + \Delta x) - \omega t)]^2$$

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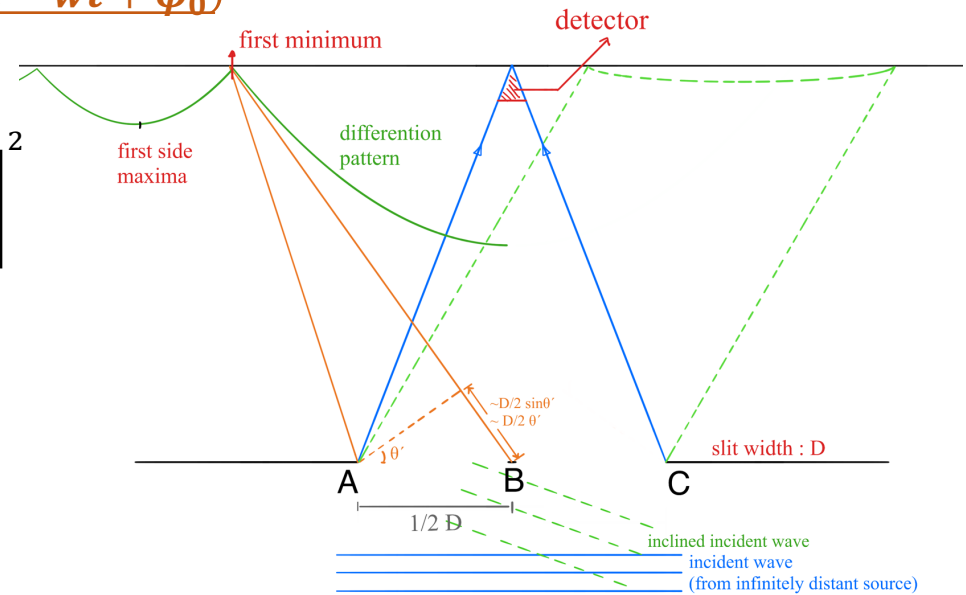
$$k\Delta x = \frac{2\pi}{\lambda} \frac{\lambda}{2} = \pi \Rightarrow \frac{D}{2} \sin \theta' \sim \frac{D}{2} \theta' = \frac{\lambda}{2} \Rightarrow \theta' \sim \lambda/D$$

$$\begin{aligned} (E_A + E_B)^2 &= E_0^2 [\cos(kx - \omega t) + \cos(k(x + \Delta x) - \omega t)]^2 \\ &= E_0^2 [\cos(kx - \omega t) + \cos(kx + \pi - \omega t)]^2 \end{aligned}$$

Two cosine functions that cancel each other completely.

1. EM-wave at long-distance limit: plane wave  $E = E_0 \cos(kx - \omega t + \phi_0)$
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$$\left[ \sum_{n=0}^{N-1} \cos(k(x + \Delta x_n) - \omega t) \right]^2 = \left[ \sum_{n=0}^{N-1} \cos\left(kx + k \frac{nD}{N} \sin \theta' - \omega t\right) \right]^2$$

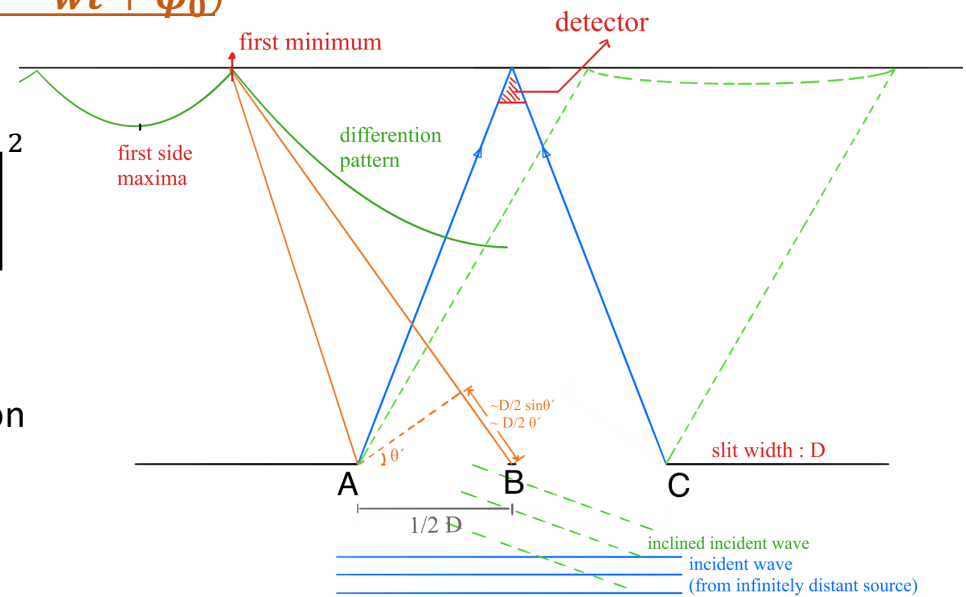




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In the limit of  $N \rightarrow \infty$ , the summation becomes integration



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$$\sin \theta - \sin \varphi = 2 \sin \left( \frac{\theta - \varphi}{2} \right) \cos \left( \frac{\theta + \varphi}{2} \right)$$

↓

$$= - \frac{[\sin(kx + kD \sin \theta' - \omega t) - \sin(kx - \omega t)]}{k \sin \theta'}$$

$$= - \frac{\sin \left( \frac{1}{2} kD \sin \theta' \right) \cos \left( kx - \omega t + \frac{1}{2} kD \sin \theta' \right)}{\frac{1}{2} k \sin \theta'}$$

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In the limit of  $N \rightarrow \infty$ , the summation becomes integration

$$\begin{aligned} & \sin \theta - \sin \varphi = 2 \sin \left( \frac{\theta - \varphi}{2} \right) \cos \left( \frac{\theta + \varphi}{2} \right) \\ & \quad \downarrow \\ & - \frac{[\sin(kx + kD \sin \theta' - \omega t) - \sin(kx - \omega t)]}{k \sin \theta'} \\ & = - \underbrace{\frac{\sin\left(\frac{1}{2}kD \sin \theta'\right)}{\frac{1}{2}k \sin \theta'}}_{\text{Amplitude modulation}} \cos \left( kx - \omega t + \underbrace{\frac{1}{2}kD \sin \theta'}_{\text{Phase modulation}} \right) \end{aligned}$$

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$$\sqrt{\tilde{P}(\theta)} = \frac{\sin\left(\frac{1}{2}kD \sin \theta'\right)}{\frac{1}{2}k \sin \theta'} \sim \frac{\sin\left(\frac{1}{2}kD \theta'\right)}{\frac{1}{2}k \theta'}$$

Diffraction pattern:  $\tilde{P}(\theta)$

$$\sin \theta - \sin \varphi = 2 \sin\left(\frac{\theta - \varphi}{2}\right) \cos\left(\frac{\theta + \varphi}{2}\right)$$

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$$= - \frac{[\sin(kx + kD \sin \theta' - \omega t) - \sin(kx - \omega t)]}{k \sin \theta'}$$

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Diffraction pattern:  $\tilde{P}(\theta)$

First zero:  $\frac{1}{2}kD\theta' = \pi \Rightarrow \frac{1}{2}\frac{2\pi}{\lambda}D\theta' = \pi \Rightarrow \theta' = \frac{\lambda}{D}$

$$\sin \theta - \sin \varphi = 2 \sin\left(\frac{\theta - \varphi}{2}\right) \cos\left(\frac{\theta + \varphi}{2}\right)$$



$$= - \frac{[\sin(kx + kD \sin \theta' - \omega t) - \sin(kx - \omega t)]}{k \sin \theta'}$$

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Amplitude modulation

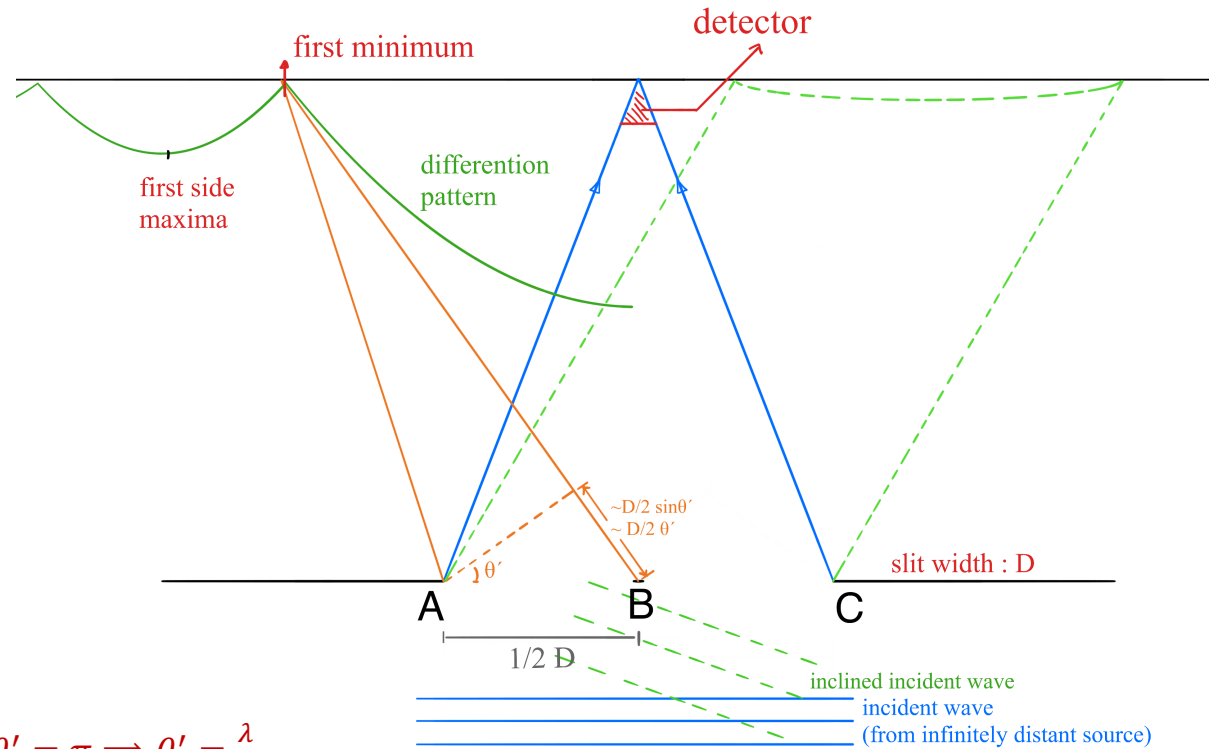
Phase modulation

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$$\sqrt{\tilde{P}(\theta)} = \frac{\sin\left(\frac{1}{2}kD \sin \theta'\right)}{\frac{1}{2}k \sin \theta'} \sim \frac{\sin\left(\frac{1}{2}kD \theta'\right)}{\frac{1}{2}k \theta'}$$

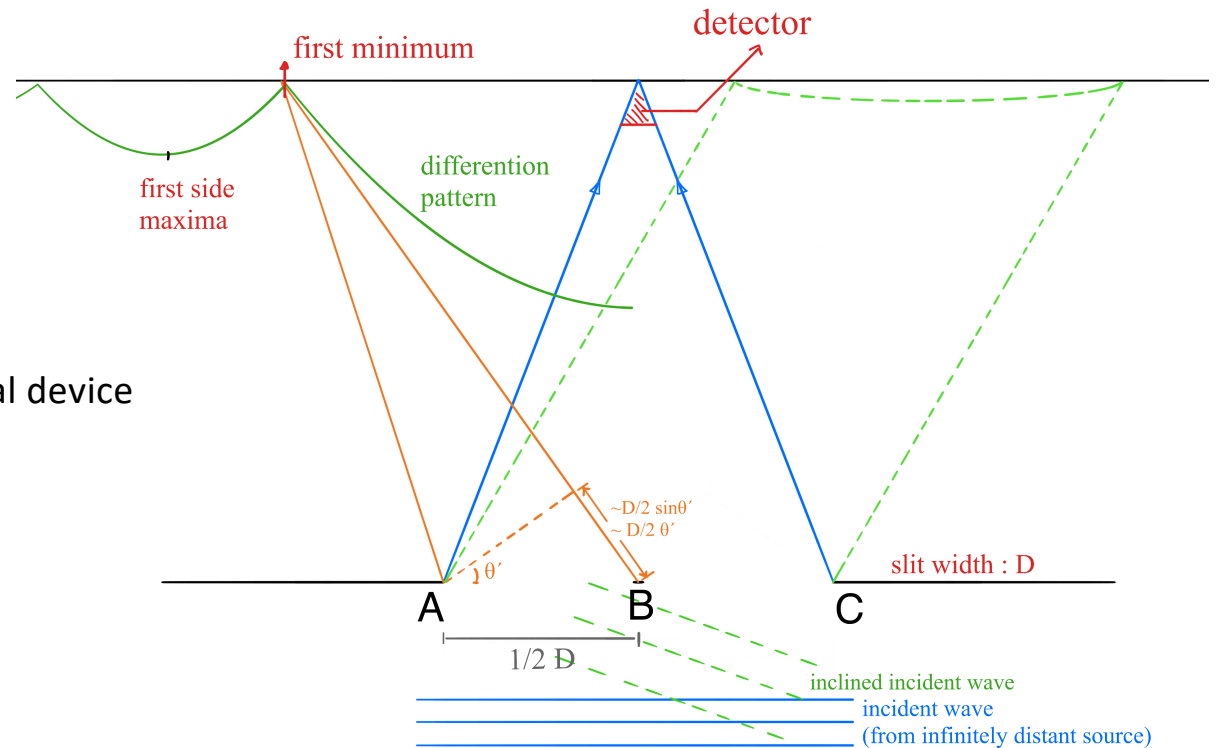
Diffraction pattern:  $\tilde{P}(\theta)$

First zero:  $\frac{1}{2}kD\theta' = \pi \Rightarrow \frac{1}{2}\frac{2\pi}{\lambda}D\theta' = \pi \Rightarrow \theta' = \frac{\lambda}{D}$



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The slit and the detector make an observational device that the response function is the same as the diffraction pattern of the single slit.



1. The diffraction pattern of a single slit:  $\tilde{P}(\theta)$ ,  $\sqrt{p(\theta)} \sim \frac{\sin\left(\frac{1}{2}kD\theta\right)}{\frac{1}{2}k\theta}$

2. The first zero of the diffraction pattern is at the angle of emergence  $\theta' = \frac{\lambda}{D}$