

# Voltage pulse unveils a far away perturbation

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(Dated: September 11, 2020)

## Abstract

Here we present the unveiling of dynamical interferences in a Fermi liquid by a voltage pulse propagating through a quantum wire. Time-dependent simulations were performed with the open-source Python package Kwant and its extension Tkwant.

## INTRODUCTION

### MODEL

We consider a  $N$  sites one dimensional wire connected to two semi-infinite leads. We model the system by the following tight-binding Hamiltonian:

$$\mathbf{H} = -\gamma \sum_i c_i^\dagger c_{i+1} + \sum_i w(t) \theta(-i) c_i^\dagger c_i + \sum_i V_{QPC}(i) c_i^\dagger c_i \quad (1)$$

The scattering region will be indexed by  $i \in \llbracket 1, N \rrbracket$ , whereas left (right) lead sites are indexed with  $i \leq 0$  ( $i \geq N + 1$ ). The voltage pulse is modeled by the second term, with  $\theta(x)$  being the Heaviside function and  $w(t)$ :

$$w(t) = V_P e^{-2((t-t_0)/\tau)^2} \quad (2)$$

$V_P$  being the amplitude,  $\tau$  the width and  $t_0$  the starting time of the perturbation. We introduced a Quantum Point Contact (QPC) which consists in a filtering potential barrier:

$$V_{QPC}(i) = V_0 e^{-((x-x_0)/\xi)^2} \quad (3)$$

$V_0$  being the amplitude,  $\xi$  the width and  $x_0$  the position of the QPC. The time-dependent perturbation can be absorbed by the gauge transforming leading to a redefinition of the hopping parameter the left lead where the pulse is applied and the scattering region:

$$\gamma \rightarrow \gamma e^{-i\phi(t)} \quad (4)$$

with:

$$\phi(t) = \frac{e}{\hbar} \int_{-\infty}^t w(u) du \quad (5)$$

### SIMULATION

Tkwant allows to evolve the states in time. The density is obtained by integrating over the energy of all onebody states:

$$n(i, t) = \int \frac{dE}{2\pi} f(E) |\psi(i, E)|^2 \quad (6)$$

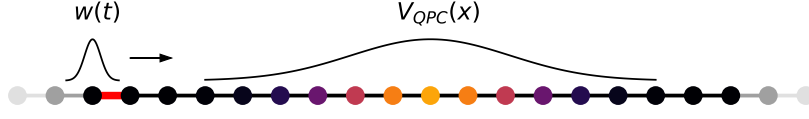


FIG. 1: A schematic of the system. The wire is connected to two semi-infinite leads. The site color reflects the QPC's voltage amplitude. The time-dependent hopping parameter is displayed in red (here  $L = 18$ ,  $V_0 = 1$ ,  $x_0 = L/2$ ,  $\xi = 3$ ).

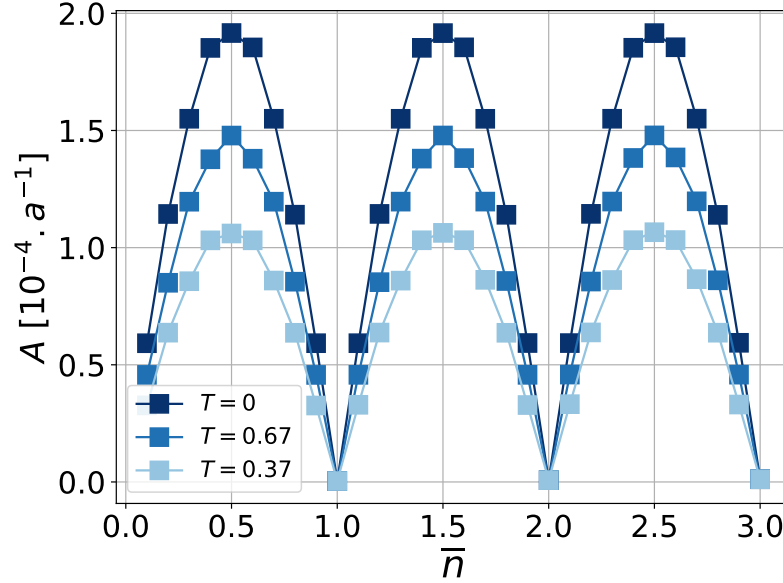


FIG. 2: The oscillation amplitude  $A$  vs the injected charge *overlined* for different values for transmission  $T$ . Here the parameters were:  $W = 1$ ,  $L = 3000$ ,  $\xi = 50$ ,  $\tau = 500$ ,  $t_0 = 1200$ .

For  $T = 0$ ,  $V_0 = 3$ ,  $\xi = 50$ , for  $T = 0.67$  ( $T = 0.37$ ),  $V_0 = 0.4$  ( $V_0 = 0.24$ ) and  $\xi = 1$ .

We define the injected charge  $\bar{n}$  as:

$$\bar{n} = \int I(t) dt \quad (7)$$

By using the Landauer-Büttiker formula  $I = eV/h$ , we obtain the following definition of  $\bar{n}$ :

$$\bar{n} = \frac{eV_P\tau}{h} \sqrt{\frac{\pi}{2}} \quad (8)$$

## INTERFERENCES

For one-dimensional chain, the scattering states have the simple form:

$$\psi(x, t) = \frac{1}{\sqrt{v(k)}} e^{ikx - iEt} \quad (9)$$

In our case, the QPC reflects the incoming states, and we obtain the following wavefunction:

$$\Psi(x, t) = \frac{1}{\sqrt{v(k)}} e^{ikx - iEt} + r \frac{1}{\sqrt{v(k)}} e^{-ikx - iEt} \quad (10)$$

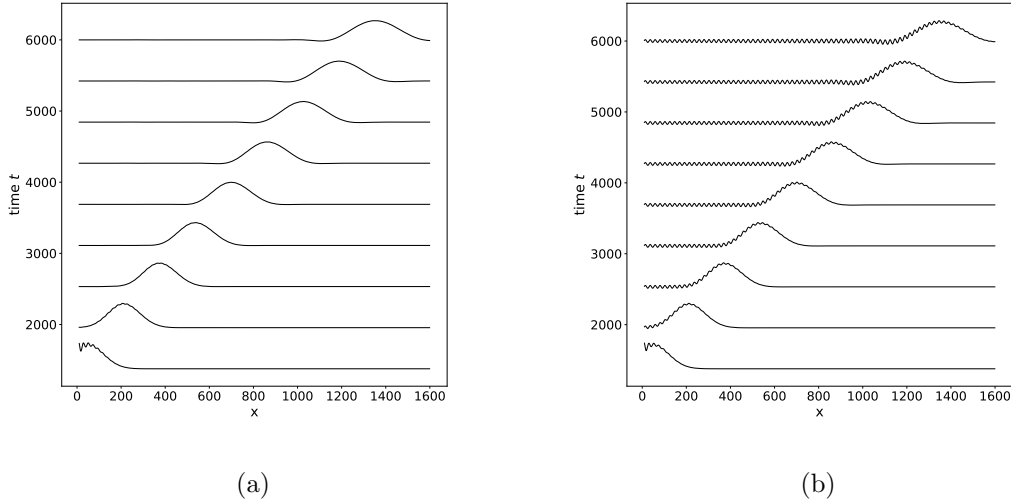


FIG. 3: The charge density as a function of time  $t$  and position  $x$  (a) without (b) QPC.

The parameters were  $\bar{n} = 0.01$ ,  $\tau = 500$ ,  $t_0 = 1200$ , (a)  $V_0 = 0$  (b)  $V_0 = 3$ ,  $\xi = 50$ ,  
 $E_F = 2.02$ .

This leads to destructive interferences. However, as the voltage pulse propagates, the incoming state phase shift dynamically:

$$\psi(x, t) = \frac{1}{\sqrt{v(k)}} e^{ikx - iEt + i\phi(t)} \quad (11)$$

which unveils the presence of interferences by changing the phase differences between the incoming and reflected states by  $\phi(t \gg \tau) = 2\pi\bar{n}$ .

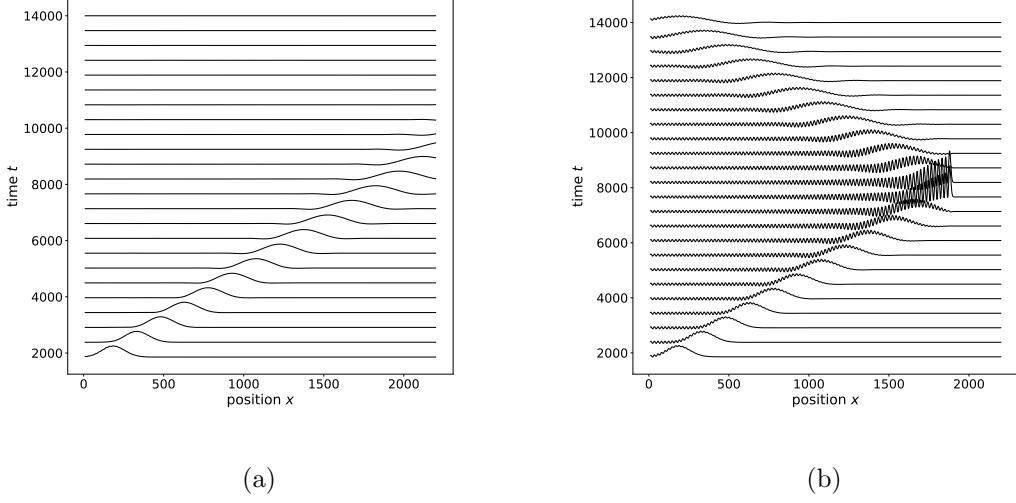


FIG. 4: The charge density as a function of time  $t$  and position  $x$  (a) without (b) QPC. The parameters were  $\bar{n} = 0.01$ ,  $\tau = 500$ ,  $t_0 = 1200$ , (a)  $V_0 = 0$  (b)  $V_0 = 3$ ,  $\xi = 50$ ,  $E_F = 2.02$ .

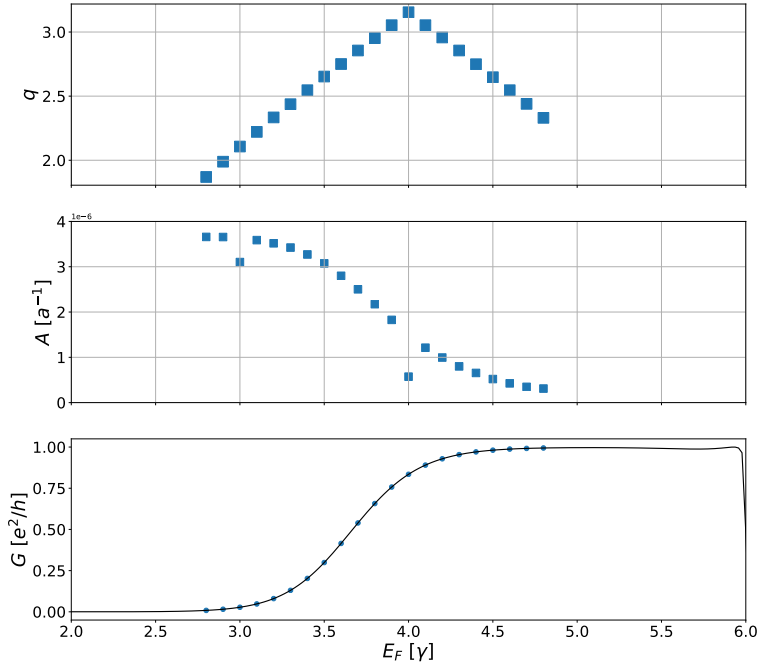


FIG. 5:  $L = 4000$ ,  $V_0 = 1.7$ ,  $\xi = 2.1$

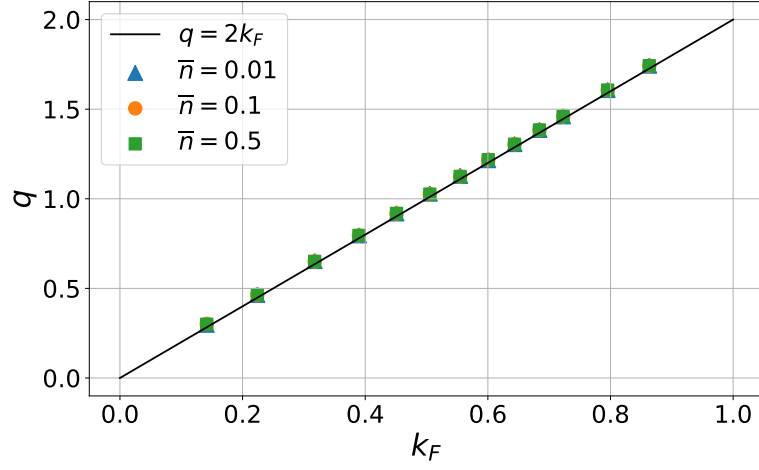


FIG. 6: The oscillation wavevector  $q$  for different  $k_F$  and  $\bar{n}$  values. Here the parameters were:  $L = 3000$ ,  $V_0 = 3$ ,  $\xi = 50$ ,  $\tau = 100$ ,  $t_0 = 1200$ .

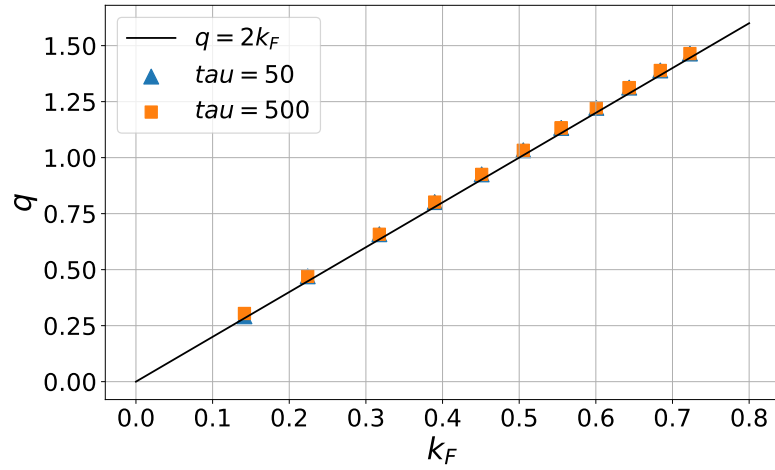


FIG. 7: The oscillation wavevector  $q$  for different  $k_F$  and  $\tau$  values. Here the parameters were:  $L = 3000$ ,  $V_0 = 3$ ,  $\xi = 50$ ,  $\bar{n} = 0.01$ ,  $t_0 = 1200$ .

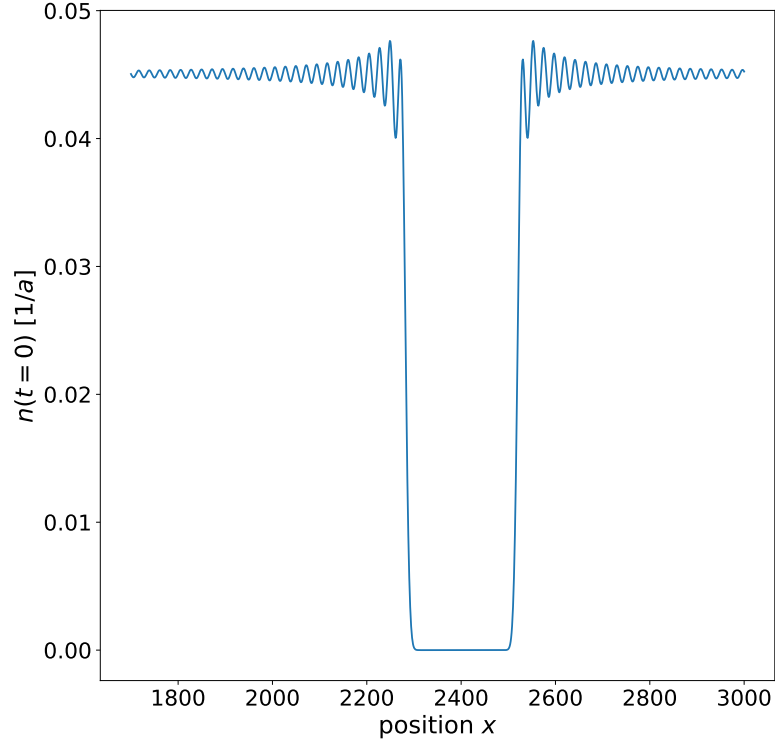


FIG. 8: Density  $n$  at  $t=0$  for  $\bar{n} = 0.01$ ,  $V_0 = 3$ ,  $E_F = 2.02$ . We observe oscillation on both side of the QPC.

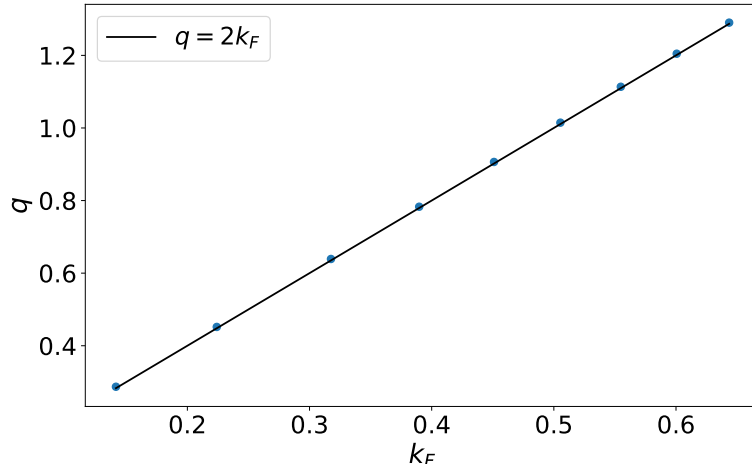


FIG. 9: Oscillation wavevector  $q$  of the oscillation at  $t=0$  at different  $k_F$  values. There is a perfect match with the line  $q = 2k_F$ .

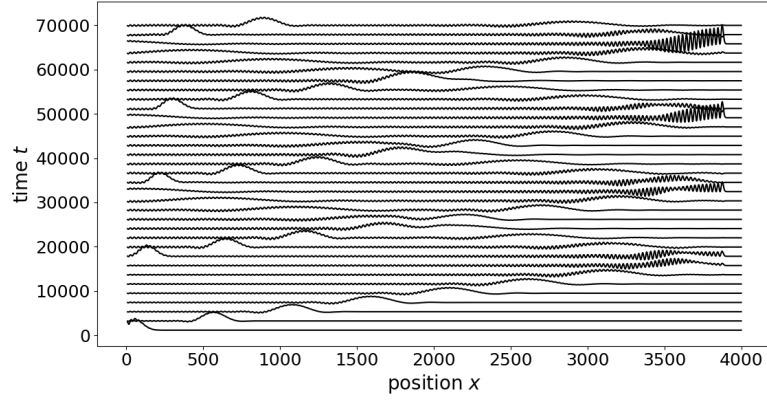


FIG. 10: For  $V_0 = 3$ ,  $\bar{n} = 0.01$ ,  $E_F = 2.015$ , we send a new pulse everytime a pulse reaches the QPC. In some region, the oscillation seems to be quasi-stationnary.