Voltage pulse unveils a far away perturbation

Auteurs

CEA, Néel

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Abstract

Here we present the unveilling of dynamical interferences in a Fermi liquid by a voltage pulse propagating through a quantum wire. Time-dependent simulations were performed with the open-source Python package Kwant and its extension Tkwant.

INTRODUCTION

MODEL

We consider a N sites one dimensional wire connected to two semi-infinite leads. We model the system by the following tight-binding Hamiltonian:

$$\mathbf{H} = -\gamma \sum_{i} c_{i}^{\dagger} c_{i+1} + \sum_{i} w(t)\theta(-i)c_{i}^{\dagger} c_{i} + \sum_{i} V_{QPC}(i)c_{i}^{\dagger} c_{i}$$

$$\tag{1}$$

The scattering region will be indexed by $i \in [1, N]$, whereas left (right) lead sites are indexed with $i \leq 0$ ($i \geq N + 1$). The voltage pulse is modeled by the second term, with $\theta(x)$ being the Heaviside function and w(t):

$$w(t) = V_P e^{-2((t-t_0)/\tau)^2}$$
(2)

 V_P being the amplitude, τ the width and t_0 the starting time of the perturbation. We introduced a Quantum Point Contact (QPC) which consists in a filtering potential barrier:

$$V_{QPC}(i) = V_0 e^{-((x-x_0)/\xi)^2}$$
(3)

 V_0 being the amplitude, ξ the width and x_0 the position of the QPC. The time-dependent peturbation can be absorbed by the gauge transforming leading to a redefinition of the hopping parameter the left lead where the pulse is applied and the scattering region:

$$\gamma \to \gamma e^{-i\phi(t)} \tag{4}$$

with:

$$\phi(t) = \frac{e}{\hbar} \int_{-\infty}^{t} w(u) du \tag{5}$$

SIMULATION

The density is obtained by integrating over the energy of all onebody states:

$$n(i,t) = \int \frac{dE}{2\pi} f(E) |\psi(i,E)|^2 \tag{6}$$

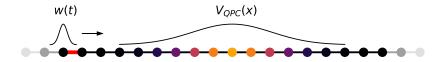


FIG. 1: A schematic of the system. The wire is connected to two semi-infinte leads. The site color reflects the QPC's voltage amplitude. The time-dependent hopping parameter is displayed in red (here $L=18,\,V_0=1,\,x_0=L/2,\,\xi=3$).

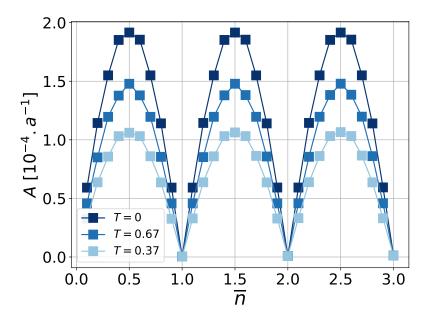


FIG. 2: The oscillation amplitude A vs the injected charge overlinen for different values for transmission T. Here the parameters were: $W=1, L=3000, \xi=50, \tau=500, t_0=1200.$ For $T=0, V_0=3, \xi=50$, for T=0.67 (T=0.37), $V_0=0.4$ ($V_0=0.24$) and $\xi=1$.

We define the injected charge \overline{n} as:

$$\overline{n} = \int I(t)dt \tag{7}$$

By using the Landauer-Bttiker formula I = eV/h, we obtain the following definition of \overline{n} :

$$\overline{n} = \frac{eV_P \tau}{h} \sqrt{\frac{\pi}{2}} \tag{8}$$

INTERFERENCES

For one-dimensional chain, the scattering states have the simple form:

$$\psi(x,t) = \frac{1}{\sqrt{v(k)}} e^{ikx - iEt} \tag{9}$$

In our case, the QPC reflects the incoming states, and we obtain the following wavefunction:

$$\Psi(x,t) = \frac{1}{\sqrt{v(k)}} e^{ikx - iEt} + r \frac{1}{\sqrt{v(k)}} e^{-ikx - iEt}$$
(10)

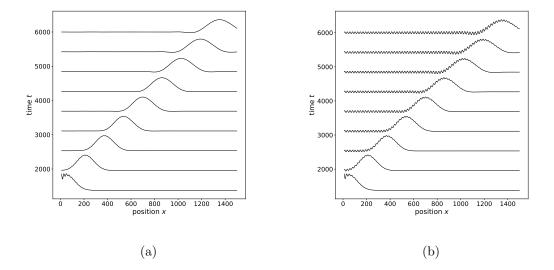


FIG. 3: The charge density as a function of time t and position x (a) without (b) QPC.

The parameters were
$$\overline{n}=0.01,\, \tau=500,\, t_0=1200,$$
 (a) $V_0=0$ (b) $V_0=3,\, \xi=50,$
$$E_F=2.02.$$

This leads to destructive interferences. However, as the voltage pulse progagates, the incoming state phase shift dynamically:

$$\psi(x,t) = \frac{1}{\sqrt{v(k)}} e^{ikx - iEt + i\phi(t)} \tag{11}$$

which unveils the presence of interferences by changing the phase differences between the incoming and reflected states by $\phi(t \gg \tau) = 2\pi \overline{n}$.

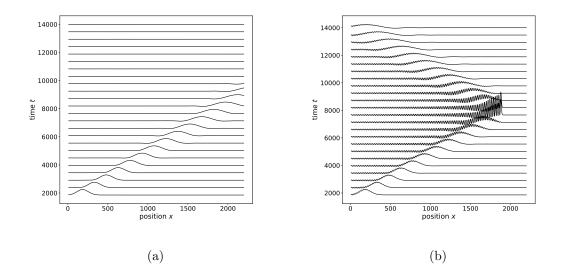


FIG. 4: The charge density as a function of time t and position x (a) without (b) QPC. The parameters were $\overline{n}=0.01,\,\tau=500,\,t_0=1200,$ (a) $V_0=0$ (b) $V_0=3,\,\xi=50,$ $E_F=2.02.$

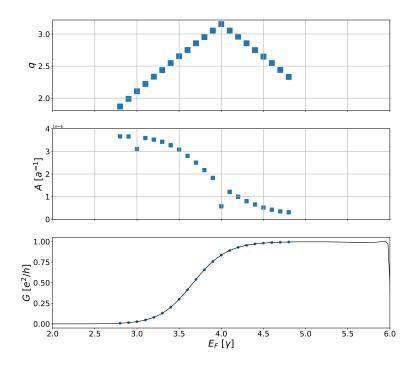


FIG. 5: $L = 4000, V_0 = 1.7, \xi = 2.1$

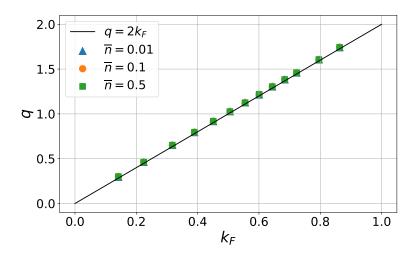


FIG. 6: The oscillation wavevector q for different k_F and \overline{n} values. Here the parameters were: $L=3000,\,V_0=3,\,\xi=50,\,\tau=100,\,t_0=1200.$

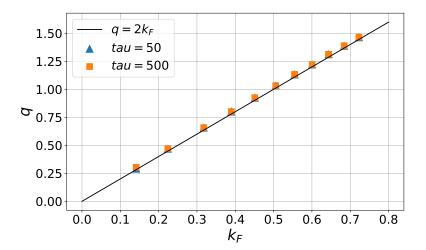


FIG. 7: The oscillation wavevector q for different k_F and τ values. Here the parameters were: $L=3000,\ V_0=3,\ \xi=50,\ \overline{n}=0.01,\ t_0=1200.$